Halo EFT: an effective bridge between ab initio nuclear-structure calculations and nuclear-reaction modelling

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Halo nuclei

Halo nuclei are found far from stability Exhibit peculiar quantal structure :

- Light, n-rich nuclei
- \bullet Low S_n or S_{2n}

With large matter radius

due to strongly clusterised structure :

neutrons tunnel far from the core and form a diffuse halo

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One-neutron halo
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^{11}Be \equiv ^{10}Be + n
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{}^{15}C \equiv {}^{14}C + n
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Two-neutron halo 6 He $\equiv {}^4$ He + n + n

¹¹Li ≡ ⁹Li + n + n

This exotic structure challenges nuclear-structure models

Reactions with halo nuclei

Halo nuclei are fascinating objects Some have been calculated *ab initio* [Calci et al. PRL 117, 242501 (2016)] However difficult to study experimentally $[\tau_{1/2}({}^{11}Be)=13$ s]

How can one probe their structure? test the ab initio predictions?

⇒require indirect techniques, like reactions :

- breakup : 11 Be + Pb/C \rightarrow 10 Be + n + Pb/C
- transfer : 10 Be(d,p)¹¹Be
- knockout : 11 Be + Be \rightarrow 10 Be + X

Need good understanding of the reaction mechanism

(i.e. a good reaction model)

to know what nuclear-structure information is probed Here, we couple precise reaction models with Halo EFT (For a short review, see [P.C. Few Body Syst 63, 14 (2022)]) We consider $¹¹$ Be, the archetypical one-neutron halo nucleus</sup>

[Introduction : halo nuclei](#page-1-0)

 (2) [Description of](#page-4-0) $¹¹Be$ </sup> \bullet [Ab initio calculation of](#page-4-0) 11 Be **• [EFT description](#page-5-0)**

[Reactions with](#page-7-0) 11 Be

- **•** [Breakup](#page-7-0)
- [Role of core excitation](#page-10-0)
- **•** [Transfer](#page-12-0)
- \bullet [KO](#page-13-0)

Ab initio description of 11 Be $NCSMC$ calculation of 11 Be [Calci *et al. PRL 117, 242501 (2016)*]

FIG. 2. NCSMC spectrum of ¹¹Be with respect to the $n + {}^{10}$ Be threshold. Dashed black lines indicate the energies of the ¹⁰Be states. Light boxes indicate resonance widths. Experimental energies are taken from Refs. [1,51].

 \bullet $\frac{1}{2}^+$ ground state : $\epsilon_{1+}^2 = -0.500 \text{ MeV}$
 $\epsilon_2^2 = 0.796 \text{ fm}^{-1/2}$ \mathcal{C}_{\perp^+} = 0.786 fm^{-1/2} $S_{1,1}^2=0.90$ $\frac{1}{2}$ $\frac{1}{2}$ bound excited state : ϵ_{1}^{2} = −0.184 MeV $\mathcal{C}_{\frac{1}{2}}^2 = 0.129$ fm^{-1/2} $S_{0p\frac{1}{2}}^2 = 0.85$ $\frac{1}{2}$ 2 $⁺$ ground state :</sup> $\mathcal{C}_{\frac{1}{2}^+}^2 = 0.786$ fm^{-1/2} $S_{1s\frac{1}{2}}^{2}=0.90$ 2 − bound excited state :

Calci *et al.* also predict the 10 Be-n phaseshift $\overline{}$

¹⁰Be-n Halo-EFT potential

Replace $10B$ Be-n interaction by effective potential in each partial wave Use Halo EFT : clear separation of scales (in energy or in distance) \Rightarrow provides an expansion parameter (small scale / large scale) along which the low-energy behaviour is expanded

[C. Bertulani, H.-W. Hammer, U. Van Kolck, NPA 712, 37 (2002)]

[H.-W. Hammer, C. Ji, D. R. Phillips JPG 44, 103002 (2017)]

Use narrow Gaussian potentials @ NLO

$$
V_{lj}(r) = V_0^{lj} e^{-\frac{r^2}{2\sigma^2}} + V_2^{lj} r^2 e^{-\frac{r^2}{2\sigma^2}}
$$

- In $s\frac{1}{2}$ $\frac{1}{2}$ and $p\frac{1}{2}$ $\frac{1}{2}$: fit V_0^{lj} V_0^{lj} and V_2^{lj} n_2^{u} to reproduce
	- ϵ_{nlj} (known experimentally)
 ϵ_{nlj} (predicted *ab initio*)
	- \rightarrow C_{nl} ; (predicted ab initio) [Calci et al. PRL 117, 242501 (2016)]
- $V_{p3/2} = 0$ to reproduce ab initio $\delta_{3/2}$ ∼ 0
- For $l > 1$: $V_{li} = 0$ @ NLO

 σ = 1.2, 1.5 or 2 fm evaluates the sensitivity to short-range physics

$s_{\frac{1}{2}}^1$ $\frac{1}{2}$: @ NLO potentials fitted to $\epsilon_{\frac{1}{2}^+}$ and $C_{\frac{1}{2}^-}$ +

Potentials fitted to $\epsilon_{1s\frac{1}{2}} = -0.503$ MeV and $C_{1s\frac{1}{2}} = 0.786$ fm^{-1/2}

• Wave functions : same asymptotics but different interior

- δ*s* 1 2 : all effective potentials are in good agreement with ab initio up to 1.5 MeV (same effective-range expansion)
- Similar results obtained for $p\frac{1}{2}$ $\frac{1}{2}$ (excited bound state)

Exp : [Fukuda et al. PRC 70, 054606 (2004)] Th. : [P.C., Phillips & Hammer, PRC 98, 034610]

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- All calculations provide very similar results $\forall \sigma$ despite the difference in the internal part of the wave function \Rightarrow reaction is peripheral [P.C. & Nunes PRC 75, 054609 (2007)]
- Excellent agreement with data on Pb (no fitting parameter) \Rightarrow confirms ab initio ANC and phaseshift
- On C, breakup strength missing at the $5/2^+$ and $3/2^+$

$d_{\frac{5}{2}}^{5}$ $\frac{5}{2}$: potentials fitted to $\epsilon_{\frac{5}{2}^+}^{\text{res}}$ $\frac{1}{2}^{\rm res}$ and $\Gamma_{\frac{5}{2}}$ +

- ldentical *δ_{d*၌} up to 1.5 MeV
up to 5 MeV for the narrow up to 5 MeV for the narrow potentials (σ = 1.2 or 1.5 fm)
Excellent agreement with ab initio results up to 2 MeV
- Excellent agreement with ab initio results up to 2 MeV

¹¹Be+C→¹⁰Be+n+C @ 67*A*MeV (beyond NLO)

• In nuclear breakup, resonances play significant role [P.C., Goldstein & Baye PRC 70, 064605 (2004)]

• Still, resonant breakup not correctly described degrees of freedom [$^{10}\text{Be}(2^*)$] missing in the effective model [Moro & Lay PRL 109, 232502 (2012)]

- \bullet 3-b force can efficiently simulate 10 Be excitation [P.C., Phillips & Hammer PLB 825, 136847 (2022)]
- Range in the *c*-*T* distance should equal that of $V_{cT} R_0 = 3.5$ fm
	- \rightarrow too small ($R_0 = 2$ fm) : no effect
	- \rightarrow too large ($R_0 = 6$ fm) : erroneous angular distribution

Including core excitation in Halo-EFT (PhD Kubushishi)

To account for core excitation within Halo-EFT :

- \bullet ¹⁰Be seen as deformed rotor [Nunes *et al.* NPA 596, 171 (1996)]
- deformation $β$ treated perturbatively to couple 0^+ and 2^+_1
equations solved with B-matrix using Lagrange radial m $_1^+$ states
- equations solved with R-matrix using Lagrange radial mesh

[L.-P. Kubushishi & P.C. (in preparation)]

- $\theta \approx 0.5$ improves agreement with ab initio radial wave function • improves $\delta_{1/2^+}$ up to 4 MeV (similar results $\forall \sigma$)
- Stay tuned for reaction calculations...

 $10Be(d,p)^{11}Be$

- This idea can be extended to transfer [Yang & P.C. PRC 98, 054602 (2018)]
- Various descriptions of 11 Be (@ LO) with σ = 0.4 – 2.0 fm
show that ¹⁰Be(d,p)¹¹Be is <mark>peripheral</mark> at fwd angle and low *E^d*
- \bullet This enables to reliably infer 11 Be ANC Provides a value identical to ab initio
- Excellent agreement with data Schmidt et al. PRL 108, 192701 (2012)]

 11 Be+⁹Be→¹⁰Be+X @ 60AMeV

Using Halo-EFT within eikonal model of KO gives also good results [Hebborn & P.C. PRC 100, 054607 (2019), ibid 104, 024616 (2021)]

- Excellent agreement with experiment [Aumann PRL 84, 35 (2000)]
- Wave functions with same ANC give same $\sigma_{K_O} \Rightarrow$ peripheral
- Insensitive to description of continuum \Rightarrow good probe of ANC
- For deeply bound projectile $\sigma_{\rm KO} \propto r_{\rm rms}^2 \Rightarrow$ not SF. . .
1. Hebborn & P.C. PLB 848 [Hebborn & P.C. PLB 848, 138413 (2024)]

Summary and prospect

- Halo nuclei studied mostly through reactions
- Mechanism of reactions with halo nuclei understood How to relate ab initio calculations to reaction observables ? Halo EFT : [P.C., Phillips, Hammer, PRC 98, 034610 (2018)] Efficient way to include the significant degrees of freedom
- \bullet Using one Halo-EFT description of 11 Be, we reproduce
	- \triangleright Breakup : [P.C., Phillips, Hammer, PRC 98, 034610 (2018)]
		- \star On Pb : only ANC and δ_{lj} matter
		- \star On C : effect of core excitation [Kubushishi, P.C. in preparation]
	- \blacktriangleright 10 Be(d,p) : [Yang & P.C., PRC 98, 054602 (2018)]
	- \triangleright KO : [Hebborn, P.C., PRC 104, 024616 (2021)]
- Validate the *ab initio* predictions
- Same results on ${}^{15}C$: [Moschini, Yang & P.C., PRC 100, 044615 (2019)]
- Future :
	- Include Halo EFT with core excitation in reaction models
	- Extend to other nuclei (e.g., 31 Ne)

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