

# TOWARDS QUANTUM SIMULATIONS OF NUCLEAR REACTIONS

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# OUTLINE

- Real time propagation with a quantum computer
- Simulation of nuclear dynamics
- Approximate time evolution of two interacting neutrons: simulations and actual quantum calculations on the LLNL and AQT testbeds
- Control-centric quantum computing

# UNITARY TRANSFORMATIONS

The standard operations performed by a quantum computer on qubits is a **UNITARY TRANSFORMATION**. An example are rotations from one point to another of the Bloch sphere.

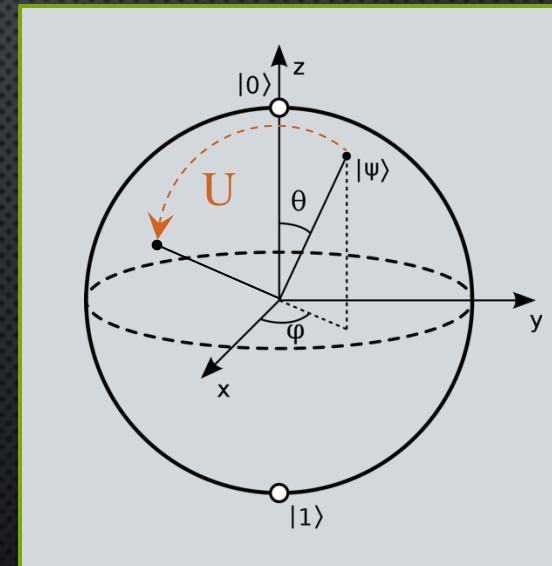
$$R_x(\theta) \equiv e^{-i\frac{\theta}{2}\sigma_x} = \cos\frac{\theta}{2}I - i\sin\frac{\theta}{2}\sigma_x = \begin{bmatrix} \cos(\theta/2) & -i\sin(\theta/2) \\ -i\sin(\theta/2) & \cos(\theta/2) \end{bmatrix}$$

However, the most relevant example for our problem is a time propagator:

$$U(t) = \exp(i\frac{H}{\hbar}t)$$

These operations can be performed on  $N$  qubits spanning a Hilbert space of needed dimensionality. And in principle the cost might just be polynomial in the number of degrees of freedom....

**Bloch Sphere**





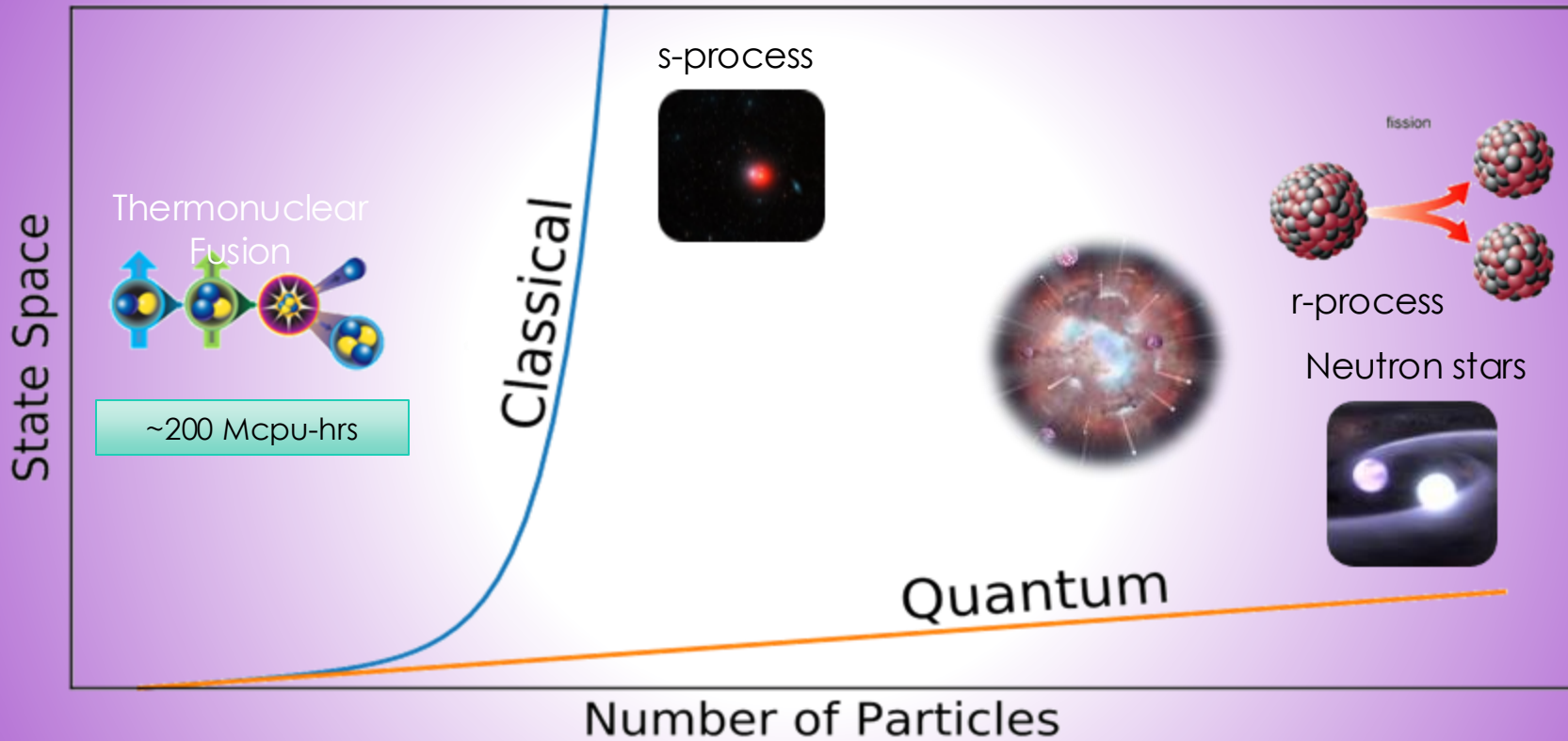
# SIMULATING REACTIONS

The concept is somewhat simple. A reaction is just a time dependent process. A quantum computer is a tool that implements the time evolution of a quantum system so it should be relatively “easy” to describe a process just as it occurs, without approximations, provided that:

- 1) We actually know the Hamiltonian of the system (and for nuclear physics this is already not trivial, unless you do full fledged QCD...)
- 2) We can implement computations:
  - a) In a reasonably large model Hilbert space (we don't really need it to be infinite, right?)
  - b) Such that noise doesn't kill the signal over the time scale relevant for the process
- 3) We know how to tame dissipation (since nature gave us quasi elastic and inelastic processes...)

And yes, at this point a reasonable person should just turn around and give up, but if it is not difficult, it is not fun!

# THE DREAM.....



# REALITY: A SIMPLE, YET NON-TRIVIAL, NUCLEAR PHYSICS PROBLEM...

Describe how **two neutrons** evolve in time under the effect of their mutual interaction

- Interaction at leading-order (LO) of chiral effective field theory (**spin tensor!**)
- Implementation of real-time evolution of the system of 2 neutrons
- Realistic device-level simulations gauged on LLNL's QPU without/with measured noise
- Measurements/calculations on the LLNL testbed
- A simple extension on a digital machine (AQT testbed @ LBL) including some kind of evolution of the coordinates



# NUCLEON NUCLEON INTERACTION (LO)

One-pion exchange (OPE)



Regularized contact  
(cutoff in momentum)

$$H_{\text{int}}^{\text{LO}} = V_{\text{OPE}} [1 - \delta_{R_0}(\vec{r})] - [C_0 + C_1 \vec{\sigma}^1 \cdot \vec{\sigma}^2] \delta_{R_0}(\vec{r})$$

regulator function

$$V_{\text{OPE}} = \frac{f_\pi^2 m_\pi}{12\pi} \left[ T_\pi(r) S_{12} - \left( Y_\pi(r) - \frac{4\pi}{m_\pi^3} \delta(\vec{r}) \right) \vec{\sigma}^1 \cdot \vec{\sigma}^2 \right] \vec{\tau}^1 \cdot \vec{\tau}^2$$

Spin independent  
Spin dependent

# TIME PROPAGATION

Formal solution of the time-dependent Schroedinger equation:

$$\exp\left(-\frac{i}{\hbar}\hat{H}t\right) = \exp\left[-\frac{i}{\hbar}\left(\hat{T} + \hat{V}_{SI} + \hat{V}_{SD}\right)t\right]$$

- $V_{SI}$ : **SPIN-INDEPENDENT** part of the interaction
- $V_{SD}$ : **SPIN-DEPENDENT** part of the interaction

In the short-time limit, we can separate the terms depending on  $V_{SI}$  and  $V_{SD}$ :

$$\exp\left[-\frac{i}{\hbar}\left(\hat{T} + \hat{V}_{SI}\right)\delta t\right] \exp\left[-\frac{i}{\hbar}\hat{V}_{SD}\delta t\right] + o(\delta t^2)$$



# TIME PROPAGATION

First application: “frozen” nucleons



Spin/isospin Hamiltonian only

$$\exp \left[ -\frac{i}{\hbar} \hat{V}_{SD} \delta t \right] = \exp \left[ -\frac{i}{\hbar} \left( \sum_{i,j=1}^A \sum_{\alpha,\beta=x,y,z} \sigma_{i\alpha} A(r_{ij})_{ij;\alpha\beta} \sigma_{j\beta} \right) \delta t \right]$$

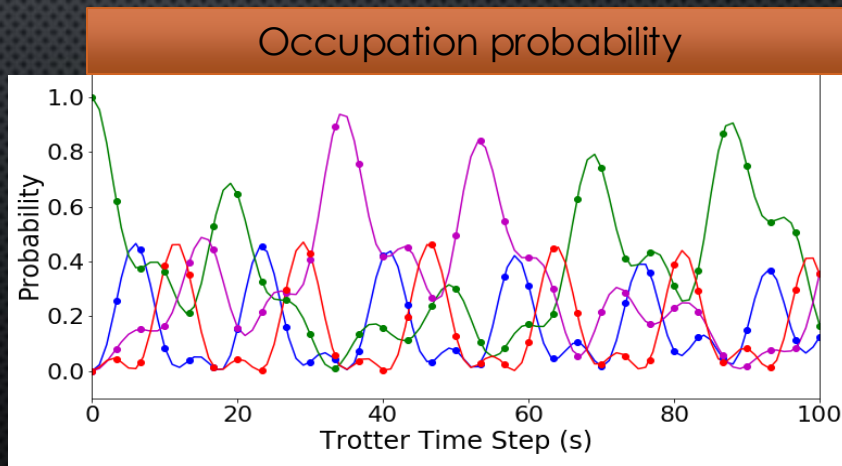
*coordinates appear as “parameters”*

Computer simulation of the actual device

**Lines:** analytic results

**Circles:** synthetic data: probability of measuring state  $|\alpha\rangle$  after a time  $t$  (obtained by repeating the process over and over)

The time evolution with a tensor Hamiltonian introduces spin-parallel components



$$|0\rangle = |\uparrow\uparrow\rangle \quad |1\rangle = |\uparrow\downarrow\rangle \quad |2\rangle = |\downarrow\uparrow\rangle \quad |3\rangle = |\downarrow\downarrow\rangle$$

# “COPROCESSING” SCHEME FOR FULL TIME EVOLUTION

$$\exp\left(-\frac{i}{\hbar}\hat{H}t\right) = \exp\left[-\frac{i}{\hbar}\left(\hat{T} + \hat{V}_{SI} + \hat{V}_{SD}\right)t\right]$$

As a first step in the direction of simulating the full dynamical evolution of a scattering process, we employed a **mixed scheme** in which the coordinates are evolved classically.

Remind that:

$$\exp\left[-\frac{i}{\hbar}\left(\hat{T} + \hat{V}_{SI}\right)\delta t\right] \exp\left[-\frac{i}{\hbar}\hat{V}_{SD}\delta t\right] + o(\delta t^2)$$

We can make the (very crude) approximation of evolving the coordinates of the nucleons **classically** (using  $V_{SI}$ ), and evolving the spin of the nucleons with the second factor in the propagator. Since the coordinates are evolved on a classical computer, we call this “**coprocessing**” scheme.



With T. Chistolini, A. Hashim, Y. Kim,  
W. Livingston, D. Santiago, I. Siddiq  
@LBLi, Phys. Rev. A 108, 032417  
(2023)



# “COPROCESSING” SCHEME FOR FULL TIME EVOLUTION

This approach has severe limits. In particular:

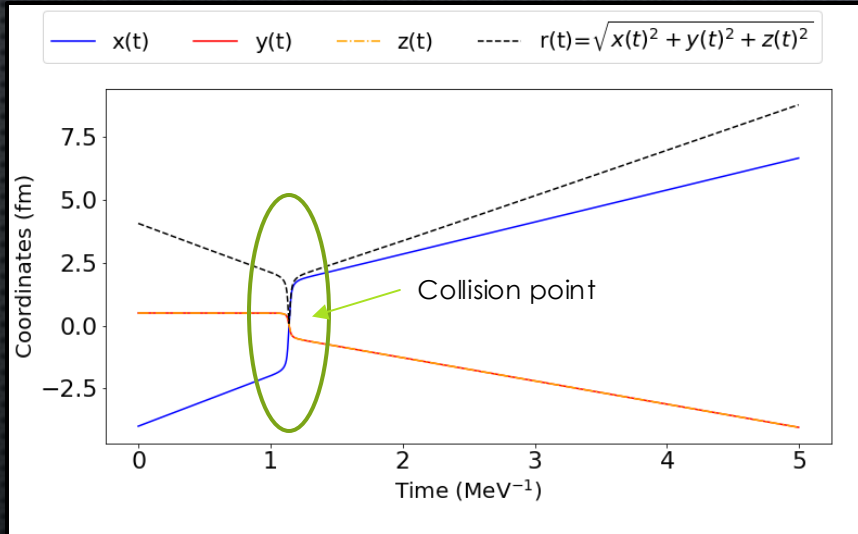
- The true space evolution is obviously **not** classical
- while there is feedback of the special evolution on the spin evolution, **the opposite does not happen**. This could be fixed in several ways.

For the moment the aim is essentially to study the **stability of a quantum computation in which the Hamiltonian has a heavy parametric dependence on time dependent quantities**.

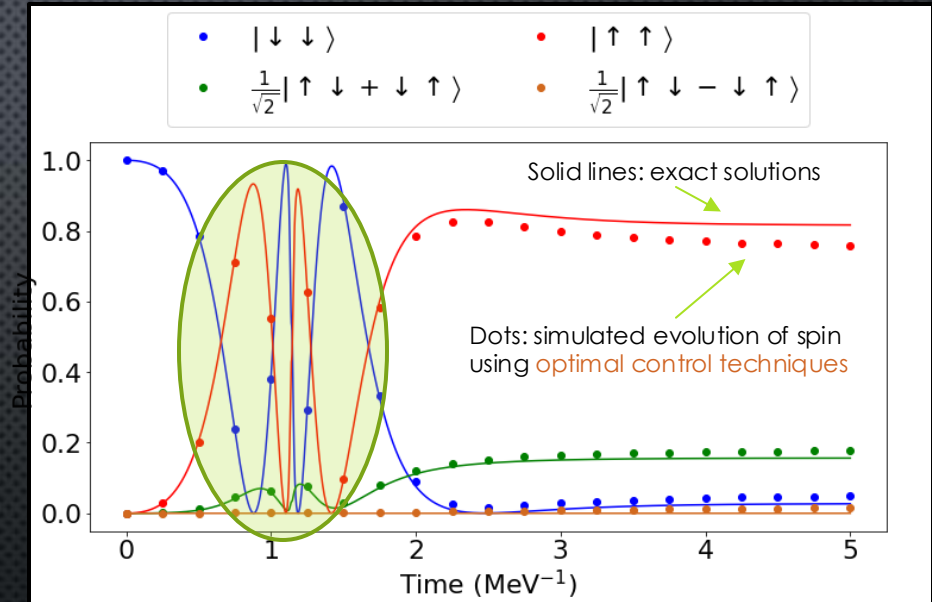
Essentially, one is studying the evolution of a state of a **time-dependent Hamiltonian** (which is one of the important features of the propagator when we use the Trotter decomposition)



# TIME EVOLUTION OF SPIN (IDEAL CASE)

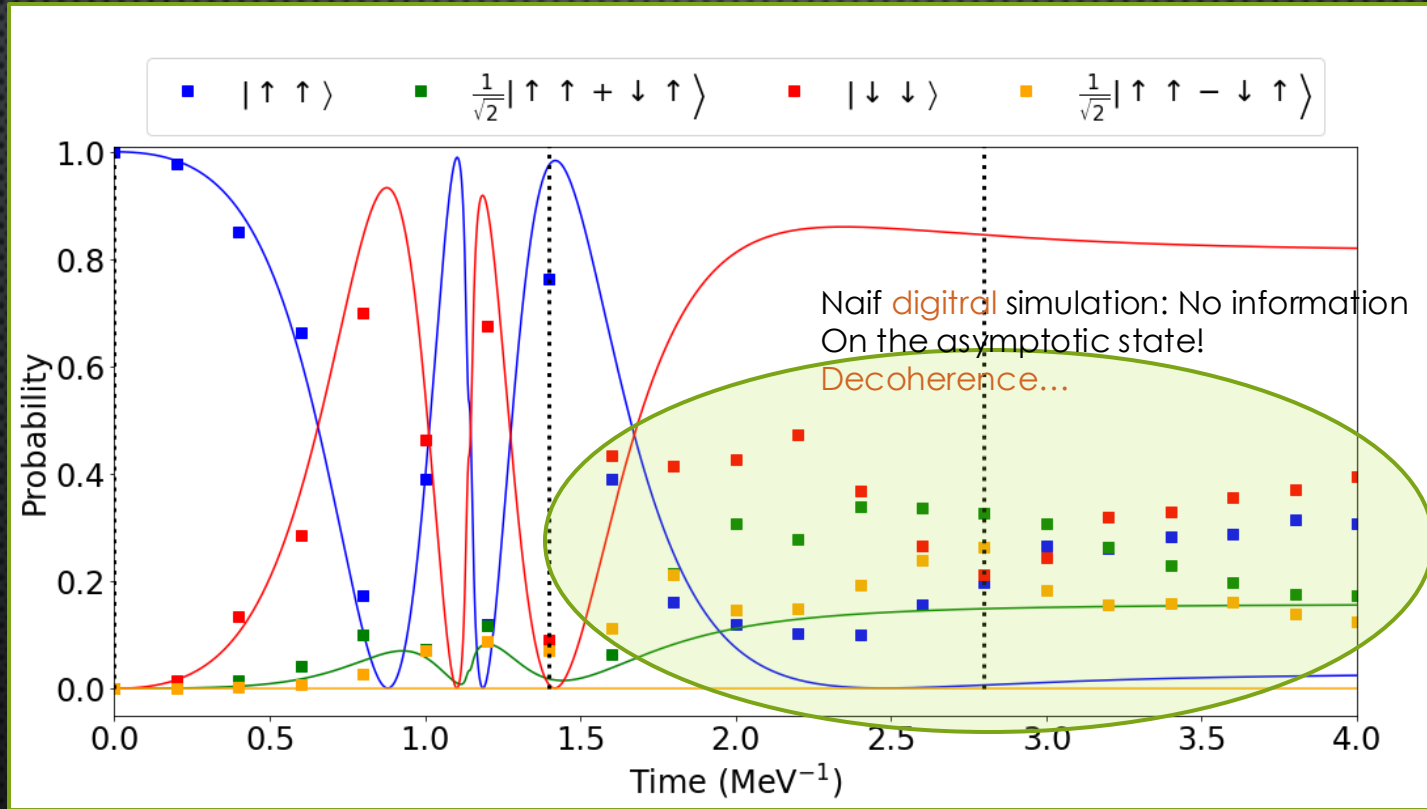


Evolution of the relative coordinate

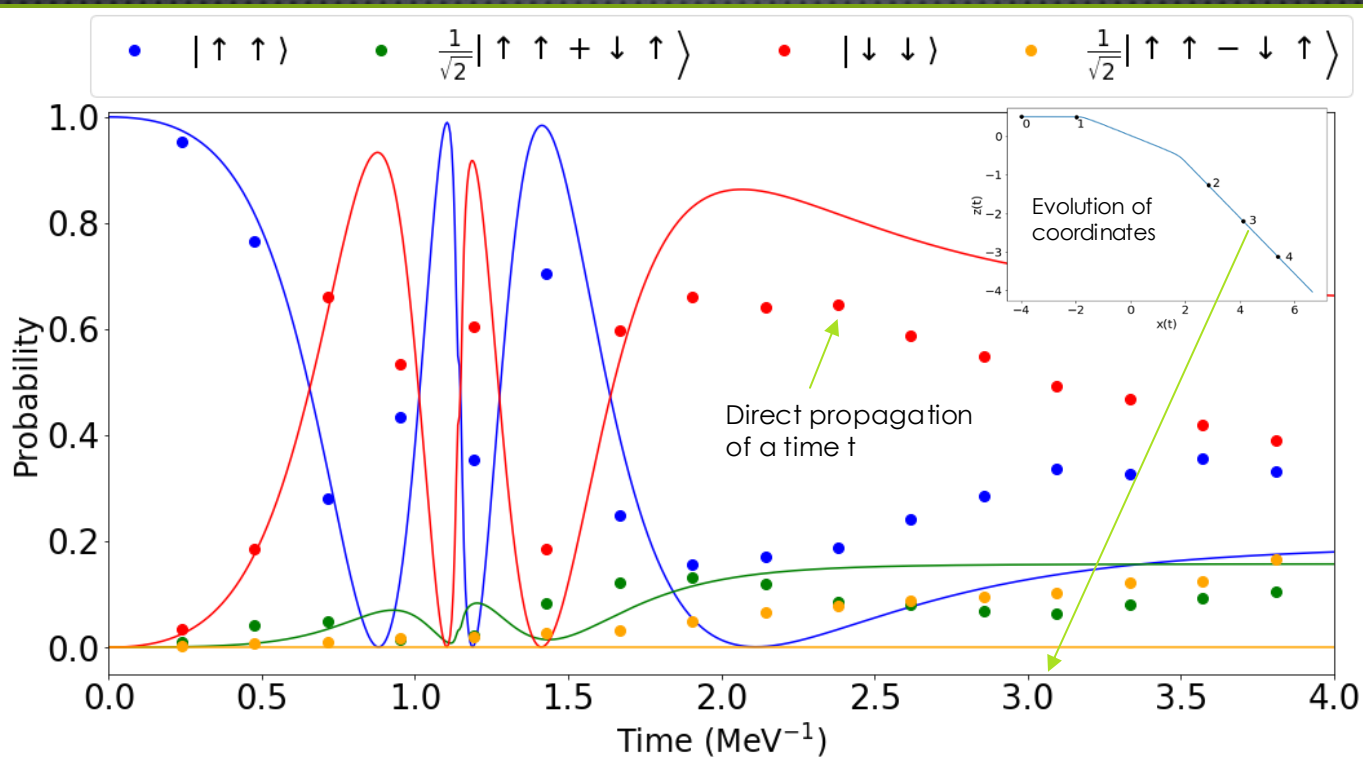


Evolution of the spin state (simulated, Optimal control)

# TIME EVOLUTION OF SPIN (AQT@LBL)



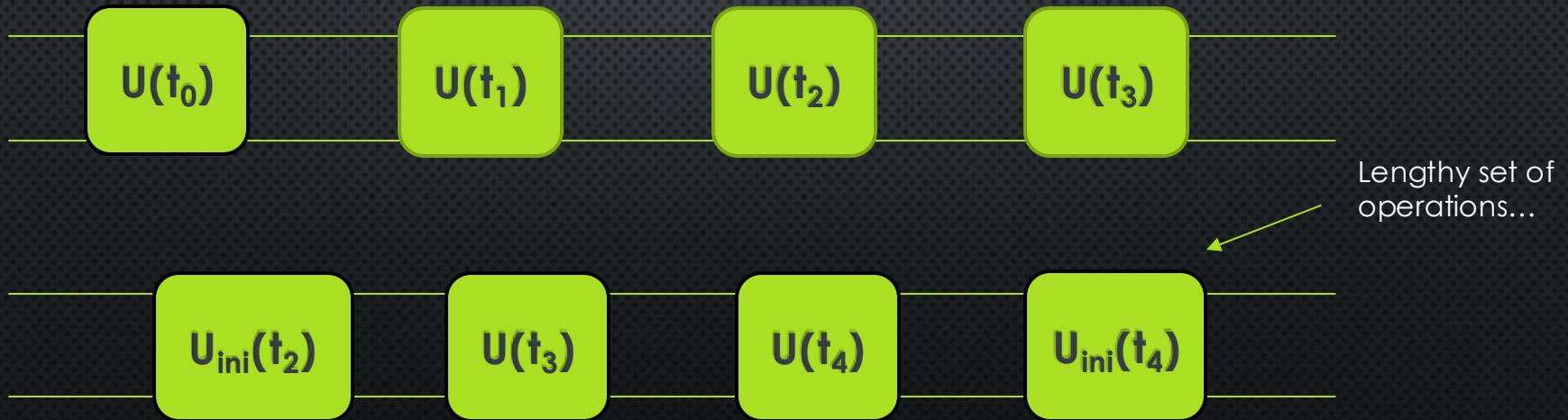
# "LONG TIME" PROPAGATOR





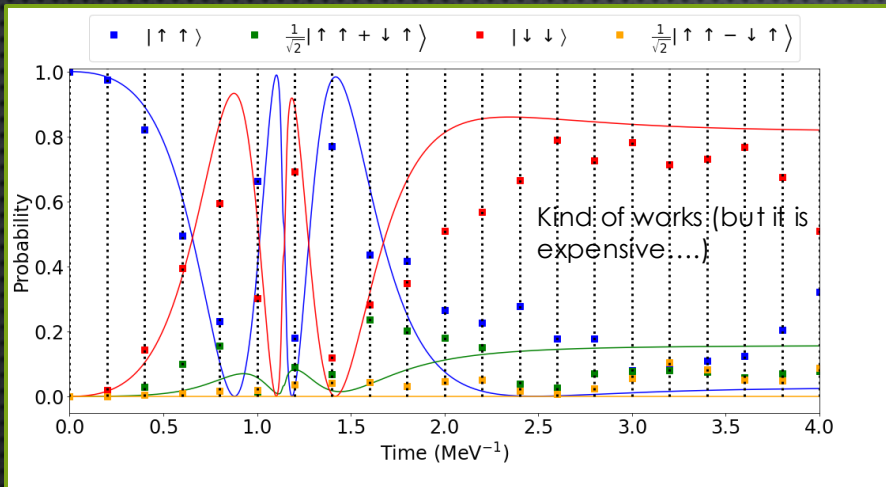
# REINIZIALIZATION

A possible, non scalable, approach consists of re-initializing the state at each time-step. Of course this requires **STATE TOMOGRAPHY**, followed by a circuit (set of (controlled) rotations) giving back the required state.



# REINIZIALIZATION

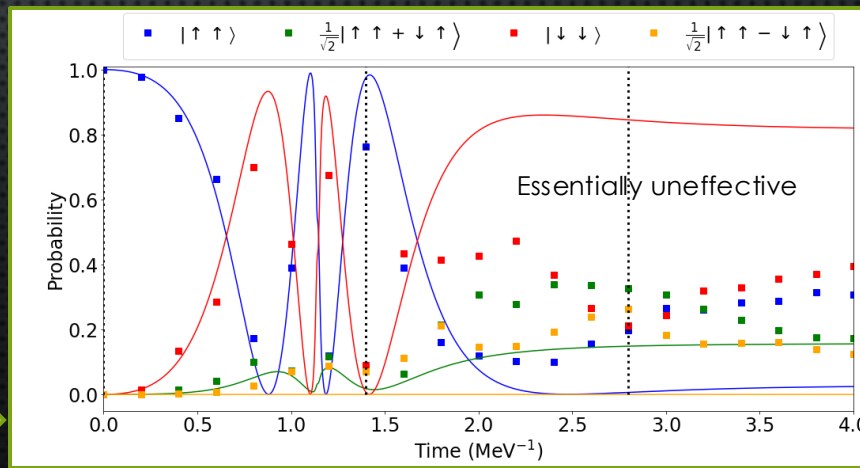
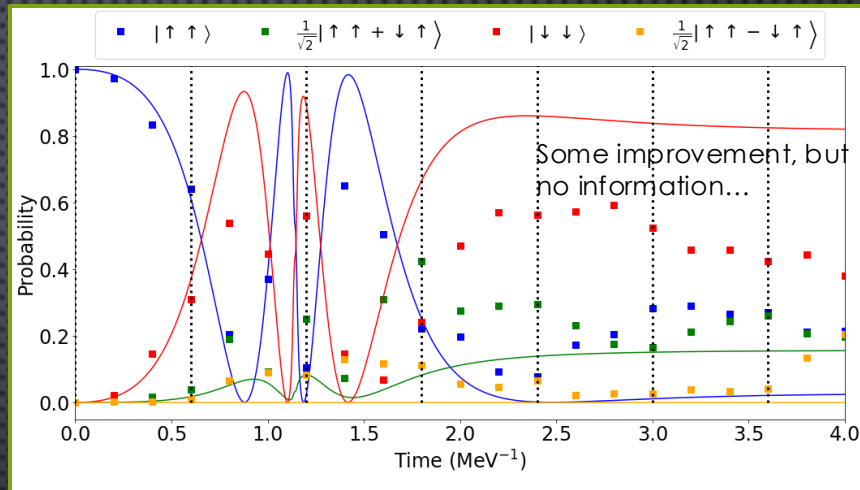
Every 3 time steps



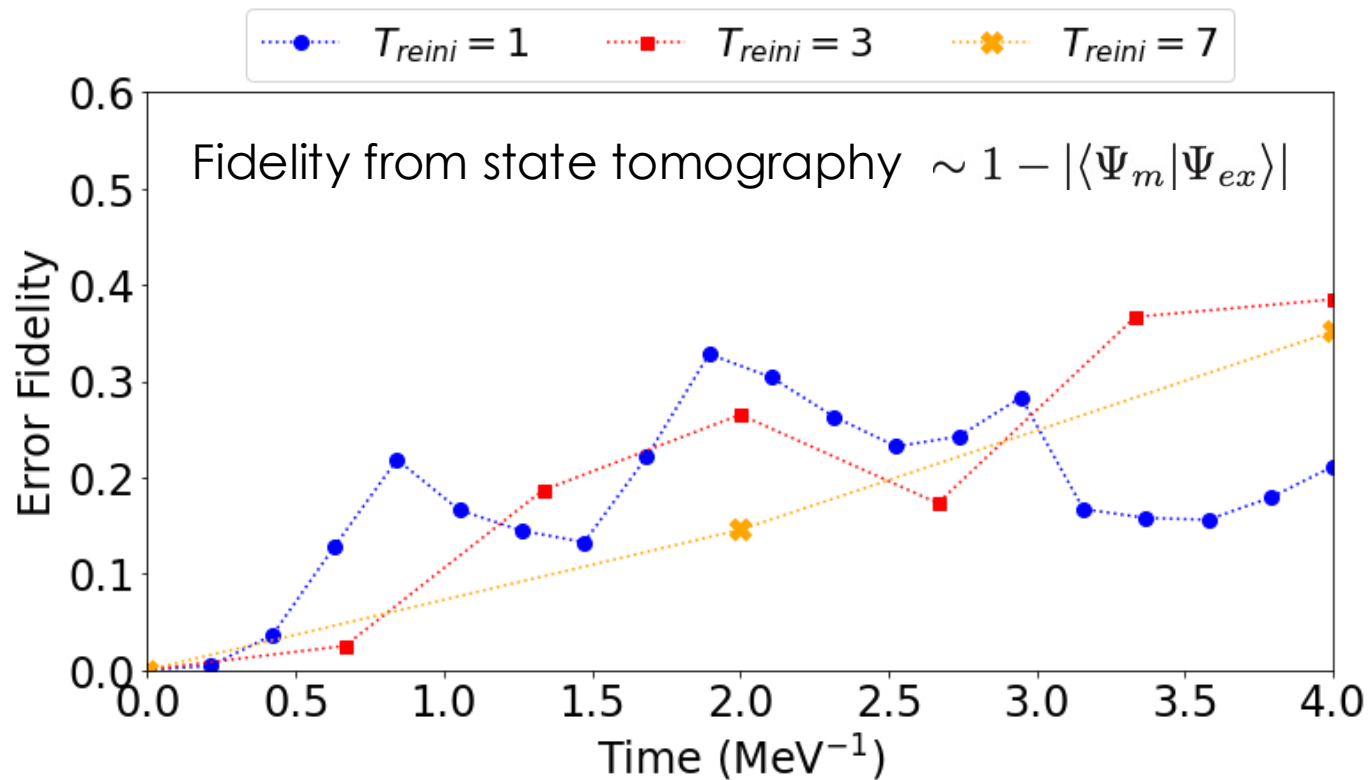
Every time step



Every 7 time steps

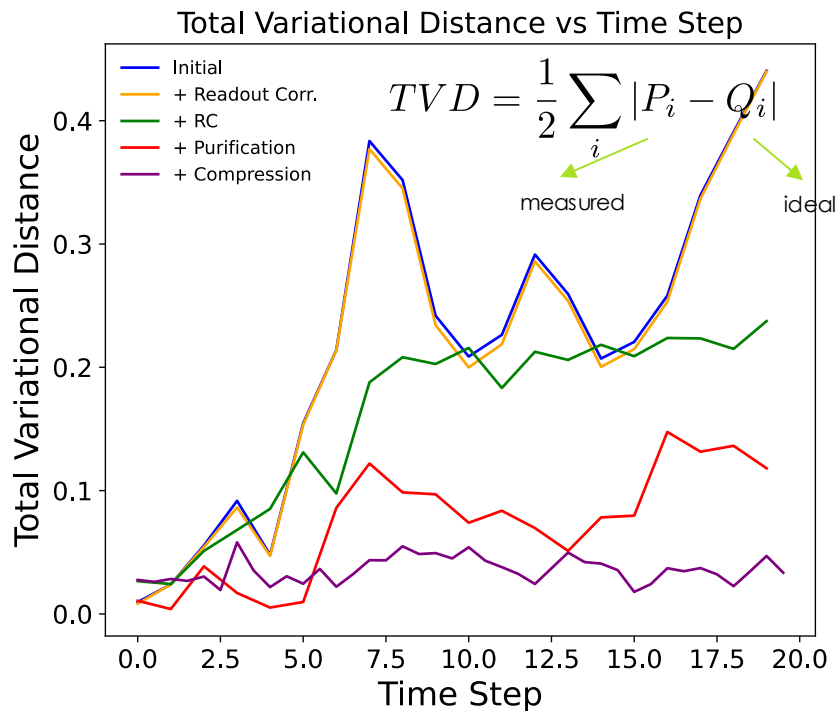


# FIDELITY IN THE REINIZIALIZATION PROCEDURE





# FURTHER IMPROVEMENTS



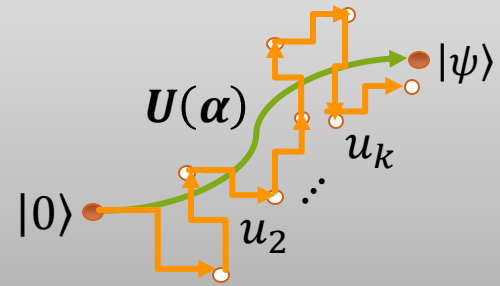
The results shown for the reinitialization strategy were obtained without implementing any error correction procedures. These might help to increase the time interval between successive reinitializations.

# "DIGITAL" QUANTUM COMPUTING

## Useful theorem (Solovay-Kitaev)

Given an Hilbert space  $\mathcal{H}$ , any unitary transformation  $U \in U(\mathcal{H})$  can be approximated with arbitrary precision  $\delta$  by a circuit of size  $\text{poly}(1/\delta)$  over the standard gates basis, possibly using ancillary qubits (i.e there exists a

This guarantees that we can approximate for instance an arbitrary time evolution under a Hamiltonian  $H$  with a circuit depth that will increase at most a polinomially with the required accuracy.



$$U(\alpha, t) = u_1(\alpha, \Delta t)u_2(\alpha, \Delta t)\dots$$

**THIS THEOREM IS AT THE BASIS OF UNIVERSAL QC SCHEME**

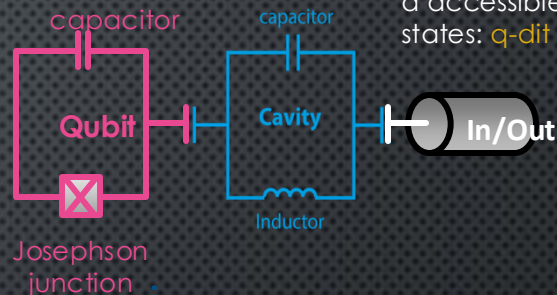
**WARNING:** other problems like decoherence (thermal effects) and dissipation

# “CONTROL-CENTRIC” APPROACH TO QUANTUM COMPUTATION

Given some control parameterization  $f(t, \alpha) = \sum_k \alpha_k \phi_k(t)$

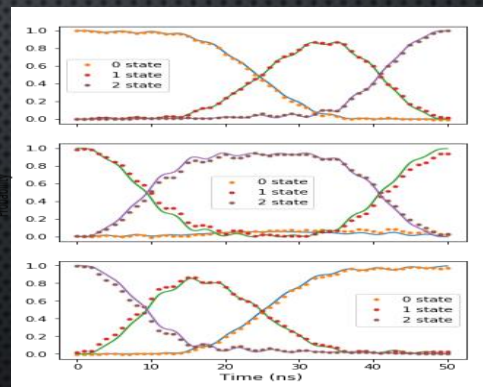
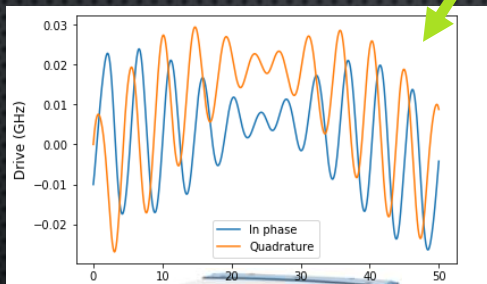
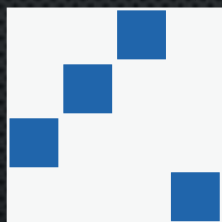
all unitary transformations in time T can be expressed as

$$U(t, \alpha) = \mathcal{T} e^{-i \int_0^t dt' [\hat{H}_0 + f(t', \alpha) \hat{H}_c]}$$



Device with d accessible states: q-dit

$|01\rangle$   $|00\rangle$   $|10\rangle$   $|11\rangle$

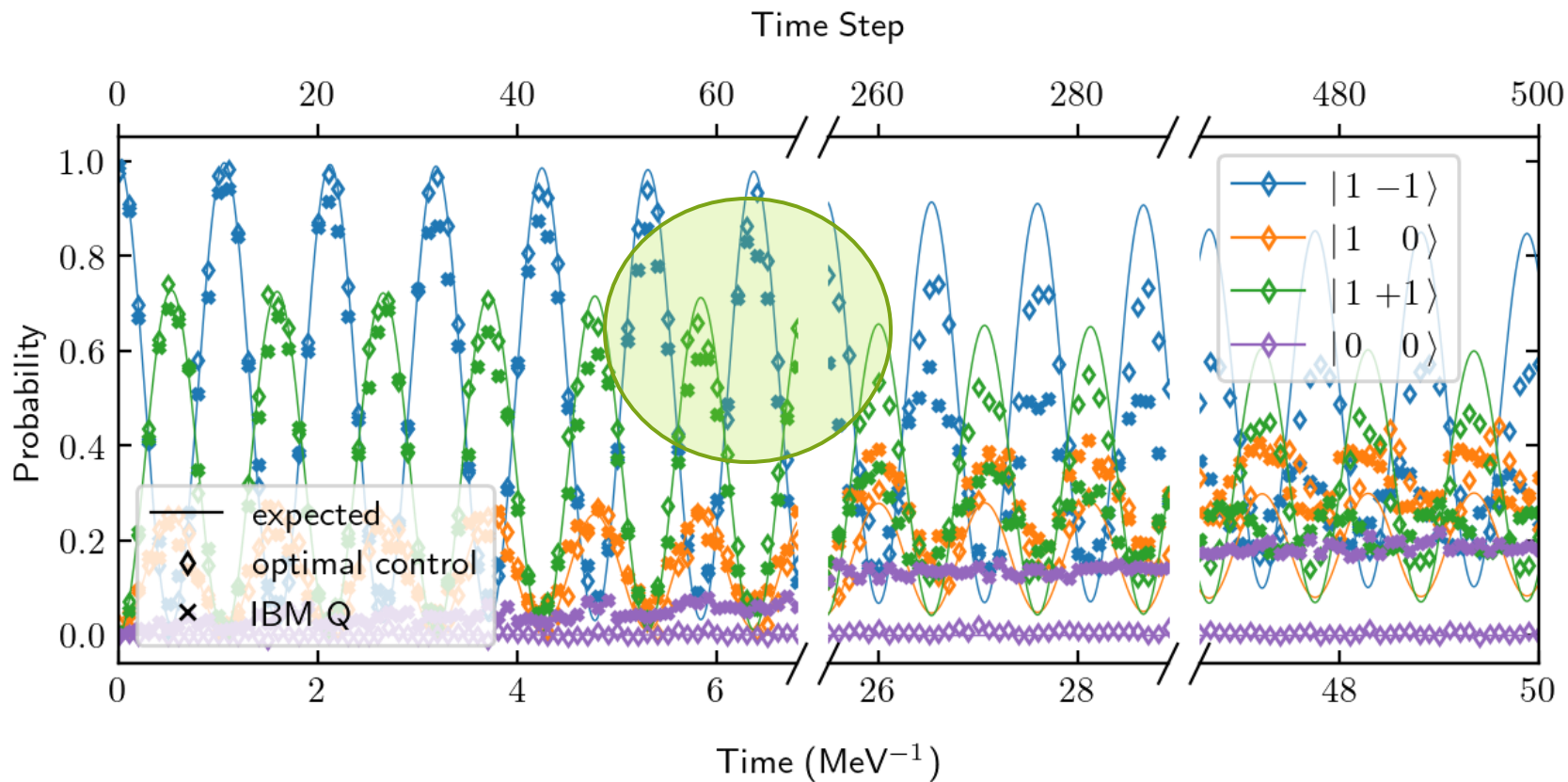


See e.g. S. Schirmer, “Hamiltonian Engineering for Quantum Systems.”, Proceedings of the 7th International Conference on Cooperative Control and Optimization (2007)

Optimization methods, machine learning (Q@Tn ML-QFORGE project (Jonathan Dubois, Kyle Wendt (LLNL), Simone Taioli, Paolo Trevisanutto (ECT\*-LISC), FP, Piero Luchi, (UNITN)) – PRA 101, 062307 (2020) – INFN Quart&T Project (UNIMIB, FBK, TIFPA, INFN-LNF and others, 2024)



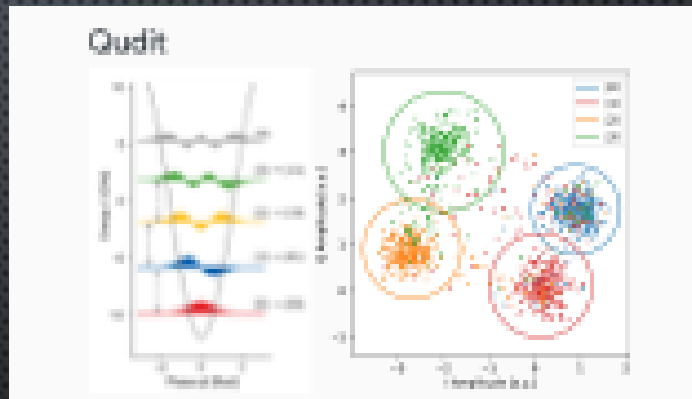
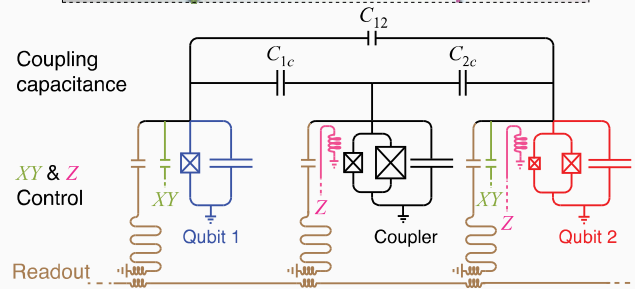
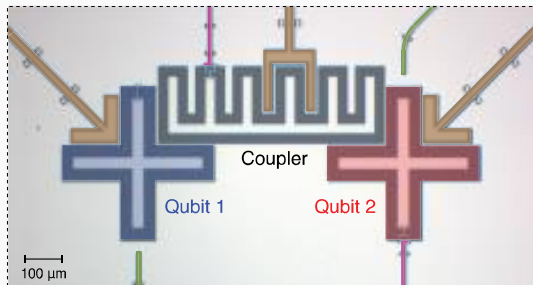
# RESULTS ON LLNL TEST BED SIMULATION





# QUART&T @INFN

Tunable coupler



Implementation of qudits with all-to-all connectivity.

Tunable coupling: possibility of controlling  
The coupling between qubits  
(change the machine Hamiltonian)



CO-DESIGN



# CONCLUSIONS

- Optimal control techniques can definitely help to realize proofs of concept for quantum algorithms that need to be robust over time (as for the case of scattering studies)
- A path for the study of nuclear reactions with realistic interactions is visible. However, there are still many open problems (like a more accurate implementation of space propagators with a reasonable scaling) which still need to be addressed.
- Digital QC requires a more elaborated circuit optimization, but some interesting demonstration might not be too far.
- In both cases, it is very important to have access to low level optimization -> **work in strict cooperation with the experimental groups!**

THANK YOU!