## Next-to-Leading Order Unitarity Fits in the (Extended) Georgi-Machacek Model

based on

[arXív: 2111.14195] with A. Kundu and P. B. Pal [arXív: 2404.18996] with D. Chowdhury and S. Samanta

> Poulami Mondal IIT Kanpur, India

Higgs Hunting Workshop, Orsay/Paris September 25, 2024







#### Beyond the Standard Model

The Standard Model provides a framework for explaining much of the observed results @ expts.

#### **Major Issues**

- \*Dark Matter
- \* Massive Neutrinos
- \*Baryon asymmetry of the Universe (BAU)
- \*Origin of electroweak (EW) symmetry breaking and the type of EWPT

# Search for BSM @ forefront of particle physics research



#### Higgs Triplet Models (HTMS)



\*Exciting phenomenological aspects in colliders: LHC, HL-LHC, ILC, Muon collider, FCC etc.

\*Can be probed in cosmological observatories: LISA, DECIGO, PRIME, Roman Telescope etc.

# $\rho = \frac{m_W^2}{m_Z^2 Cos^2 \theta_W} = 1, \text{ at tree level} \qquad \begin{array}{l} \textbf{Custodial Symmetry (CS)} \\ \phi = \left( \frac{\phi_+}{\phi_0} \right), \quad \chi = \left( \begin{array}{c} \chi_{++} \\ \chi_+ \\ \chi_0 \end{array} \right), \quad \xi = \left( \begin{array}{c} \xi_+ \\ \xi_0 \\ -\xi_- \end{array} \right) \\ \textbf{Y=1/2} \qquad \textbf{Y=1} \qquad \textbf{Y=0} \end{array}$

 $V = -m_{\phi}^{2}(\phi^{\dagger}\phi) - m_{\xi}^{2}(\xi^{\dagger}\xi) - m_{\chi}^{2}(\chi^{\dagger}\chi) + \mu_{1}(\chi^{\dagger}t_{a}\chi)\xi_{a} + \mu_{2}(\phi^{\dagger}\tau_{a}\phi)\xi_{a}$  $+ \mu_{3}\left[(\phi^{T}\epsilon\tau_{a}\phi)\tilde{\chi}_{a} + \text{h.c.}\right] + \lambda_{\phi}(\phi^{\dagger}\phi)^{2} + \lambda_{\xi}(\xi^{\dagger}\xi)^{2} + \lambda_{\chi}(\chi^{\dagger}\chi)^{2}$  $+ \tilde{\lambda}_{\chi}|\tilde{\chi}^{\dagger}\chi|^{2} + \lambda_{\phi\xi}(\phi^{\dagger}\phi)(\xi^{\dagger}\xi) + \lambda_{\phi\chi}(\phi^{\dagger}\phi)(\chi^{\dagger}\chi) + \lambda_{\chi\xi}(\chi^{\dagger}\chi)(\xi^{\dagger}\xi)$  $+ \kappa_{1}|\xi^{\dagger}\chi|^{2} + \kappa_{2}(\phi^{\dagger}\tau_{a}\phi)(\chi^{\dagger}t_{a}\chi) + \kappa_{3}\left[(\phi^{T}\epsilon\tau_{a}\phi)(\chi^{\dagger}t_{a}\xi) + \text{h.c.}\right],$ 

 $\langle \phi \rangle = v_{\phi}, \langle \xi \rangle = v_{\xi}, \langle \chi \rangle = v_{\chi}$ EW symmetry  $\rho = \frac{v_{\phi}^2 + 4(v_{\xi}^2 + v_{\chi}^2)}{v_{\phi}^2 + 8v_{\chi}^2}, \qquad \rho = 1 \to v_{\chi} = v_{\xi}$ breaking

## Georgi Machacek Model

+ In 1985, GM model was first proposed by Georgi and Machacek as a minimal HTM with  $\rho = 1$ 

$$V = \frac{1}{2}m_2^2 \operatorname{Tr}(\Phi^{\dagger}\Phi) + \frac{1}{2}m_3^2 \operatorname{Tr}(X^{\dagger}X) - M_1 \operatorname{Tr}(\Phi^{\dagger}\tau_a^{\dagger}\Phi\tau_b)X_{ab} - M_2 \operatorname{Tr}(X^{\dagger}t_a^{\dagger}Xt_b)X_{ab} + \lambda_1 (\operatorname{Tr}\Phi^{\dagger}\Phi)^2 + \lambda_2 (\operatorname{Tr}X^{\dagger}X)^2 + \lambda_3 \operatorname{Tr}(X^{\dagger}XX^{\dagger}X) + \lambda_4 (\operatorname{Tr}\Phi^{\dagger}\Phi) \operatorname{Tr}(X^{\dagger}X) - \lambda_5 \operatorname{Tr}(\Phi^{\dagger}\tau_a^{\dagger}\Phi\tau_b) \operatorname{Tr}(X^{\dagger}t_a^{\dagger}Xt_b),$$

On the centre stage of BSM searches @collider and cosmological expts.

#### EUROPEAN ORGANISATION FOR NUCLEAR RESEARCH (CERN)





Combination of searches for singly and doubly charged Higgs bosons produced via vector-boson fusion in proton–proton collisions at  $\sqrt{s} = 13$  TeV with the ATLAS detector

The ATLAS Collaboration

PHYSICAL REVIEW D 106, 055019 (2022)

#### Updated constraints on the Georgi-Machacek model and its electroweak phase transition and associated gravitational waves

Ting-Kuo Chen<sup>®</sup>,<sup>1,\*</sup> Cheng-Wei Chiang<sup>®</sup>,<sup>1,2,†</sup> Cheng-Tse Huang<sup>®</sup>,<sup>1,‡</sup> and Bo-Qiang Lu<sup>3,§</sup> <sup>1</sup>Department of Physics, National Taiwan University, Taipei, Taiwan 10617, Republic of China <sup>2</sup>Physics Division, National Center for Theoretical Sciences, Taipei, Taiwan 10617, Republic of China <sup>3</sup>School of Science, Huzhou University, Huzhou, Zhejiang 313000, People's Republic of China</sup>

(Received 15 May 2022; accepted 22 August 2022; published 19 September 2022)

#### PHYSICAL REVIEW D 90, 015007 (2014)

#### The decoupling limit in the Georgi-Machacek model

Katy Hartling,<sup>\*</sup> Kunal Kumar,<sup>†</sup> and Heather E. Logan<sup>‡</sup> Ottawa-Carleton Institute for Physics, Carleton University, 1125 Colonel By Drive, Ottawa, Ontario K1S 5B6, Canada (Received 25 April 2014; published 9 July 2014)

#### PHYSICAL REVIEW D 91, 015013 (2015)

#### Indirect constraints on the Georgi-Machacek model and implications for Higgs boson couplings

Katy Hartling,<sup>\*</sup> Kunal Kumar,<sup>†</sup> and Heather E. Logan<sup>‡</sup> Ottawa-Carleton Institute for Physics, Carleton University, 1125 Colonel By Drive, Ottawa, Ontario K1S 5B6, Canada (Received 4 November 2014; published 15 January 2015)



Published for SISSA by 🖉 Springer

RECEIVED: December 19, 2018 REVISED: January 16, 2019 ACCEPTED: January 24, 2019 PUBLISHED: January 29, 2019

#### Electroweak phase transition and Higgs phenomenology in the Georgi-Machacek model

Ruiyu Zhou,  $^a$  Wei Cheng,  $^a$  Xin Deng,  $^a$  Ligong  $\mathrm{Bian}^{a,b,1}$  and Yongcheng  $\mathrm{Wu}^c$ 

Published for SISSA by 2 Springer

RECEIVED: November 15, 2012 ACCEPTED: December 8, 2012 PUBLISHED: January 3, 2013

## Testing the custodial symmetry in the Higgs sector of the Georgi-Machacek model

Cheng-Wei Chiang<sup>*a,b,c*</sup> and Kei Yagyu<sup>*a*</sup>





vs.



 Scalar multiplets mass degenerate

• Divergent contribution to  $\rho$  parameter @ one-loop

• Only  $H^+$  couples to fermions.

Scalar multiplets non mass degenerate

new decay modes open up

search for new physics @ colliders

All the counter-terms are present in the Lagrangian

**Both** H<sup>+</sup> and F<sup>+</sup>couple to fermions

Much richer flavour physics phenomenology Theory Constraints on the parameter space

#### **\* Why theory bounds are important?**

Ans: More statistically robust than expt data with errors

**Theory Constraints** 

- Higgs potential must be bounded from below
- Yukawa and quartic couplings need to be in perturbative regime
- Eigenvalues of the S matrix of  $2\times 2$  scattering should satisfy NLO unitarity bounds
- NLO corrections to the LO eigenvalues should be smaller in magnitude.

#### Bounded From Below (BFB)

Make sure that the scalar potential must be bounded from in any direction of the field scape









3 and all 13 nonvanishing field directions generating large overlapping parameter space unitarity constraints

Prior to Higgs discovery : unitarity bound @ tree level@ 1-loopLee, Quigg, Thacker '77 $\lambda \leq \frac{8\pi}{3}$  $\lambda \leq 2-2.5$ Dawson, Eillenbrock '89; Durand, Johnson,

2 loop calculation shows no revised limit

Lopez'92

Durand, Maher, Ríesselmann, 92

#### Weakly interacting Higgs scenario

$\mathbf{Q}^{\mathbf{Y}}$	0		1/2		1		3/2	2	Goldstone boson	
	$\phi^{0*}\phi^0$	$\frac{\xi^0\xi^0}{\sqrt{2}}$	$\chi^{0*}\chi^0$	$\phi^0 \xi^0$	$\phi^{0*}\chi^0$	$\frac{\phi^0\phi^0}{\sqrt{2}}$	$\chi^{0}\xi^{0}$	$\phi^0\chi^0$	$\frac{\chi^0\chi^0}{\sqrt{2}}$	equivalence theorem
0	$\phi^+\phi^-$	$\xi^+\xi^-$	$\chi^+\chi^-$	$\phi^+\xi^-$	$\chi^+ \phi^-$		$\chi^+\xi^-$			$\mathcal{M}(W_L^{\pm}, Z_L, h, \ldots) = (iC)^k \mathcal{M}(w^{\pm}, z, h, \ldots) .$
			<i>x</i> <sup>++</sup> <i>x</i> <sup></sup>							
-	$\phi^{0*}\phi^+$	$\xi^0 \xi^+$	$\chi^{0*}\chi^+$	$\phi^0 \xi^+$	$\phi^{0*}\chi^+$	$\phi^0 \phi^+$	$\xi^0 \chi^+$	$\phi^0 \chi^+$	$\chi^+\chi^0$	
1			$\chi^{-}\chi^{++}$	$\phi^+\xi^0$	$\chi^{++}\phi^-$	$\chi^{0}\xi^{+}$	$\chi^{++}\xi^{-}$	$\phi^+\chi^0$		
0		$\frac{\xi^+\xi^+}{\sqrt{2}}$	$\chi^{0*}\chi^{++}$	$\phi^+\xi^+$	$\phi^{0*}\chi^{++}$	$\frac{\phi^+\phi^+}{\sqrt{2}}$	$\xi^+\chi^+$	$\phi^+\chi^+$	$\chi^0 \chi^{++}$	
2							$\xi^0 \chi^{++}$	$\phi^0 \chi^{++}$	$\frac{\chi^+\chi^+}{\sqrt{2}}$	Unitarity in eGM/GM
3		X			X		$\xi^+\chi^{++}$	$\phi^+\chi^{++}$	$\chi^+\chi^{++}$	
4		X			X		X	X	$\frac{\chi^{++}\chi^{++}}{\sqrt{2}}$	

# Tree-level Unitarity

16,15,11,3,1 unique tree level eigenvalues for the Block Q = 0,1,2,3,4

9 eigenvalues are independent

NLO unitarity



[Grinstein, Murphy, Uttayarat '15; Cacchio, Chowdhury, Murphy, Eberhardt'16]

GM

$$R_{1} = \frac{\left|a_{0}^{\rm NLO}\right|}{\left|a_{0}^{\rm LO} + a_{0}^{\rm NLO}\right|}, \qquad R_{1}' = \frac{\left|a_{0}^{\rm NLO}\right|}{\left|a_{0}^{\rm LO}\right|},$$

Perturbative expansion is not valid at NLO when  $R_1 = 1$  or  $R'_1 = 1$ 

# Higgs Signal Strength (Run 2)











#### Status of GIM model



The maximum mass splitting for heavy Higgs boson masses > 700 GeV is 400 GeV for GM model

which is reduced  $\sim 100 \,\text{GeV}$  from the literature



• Maximum Mass difference within same multiplet in eGM model is < 210 GeV

Allowed mass differences and quartic couplings



## Summary

- Minimal two triplet extension of SM with  $\rho = 1$  gives eGM model
- Quartic couplings in GM and eGM model gets strongly constrained by NLO unitarity
- Mixing angles and vevs get constrained from the latest LHC Higgs signal strength data
- Updated theory constraints (NLO unitarity, BFB) alone exclude a large part of the parameter space



# **Backup Slides**



#### Mass plane of eGM model



1000



## GIM model









# Higgs Signal Strength

**ATLAS** 

CMS

Signal strength	Value	Correlation matrix [	$\mathcal{L}$ $[\mathbf{fb}^{-1}]$	Source	Signal strength	Value	Correlation matrix	$\frac{\mathcal{L}}{[\mathbf{f}\mathbf{b}^{-1}]}$	Source
$\mu^{\gamma\gamma}_{ m ggF,bbh} \ \mu^{\gamma\gamma}_{ m VBF} \ \mu^{\gamma\gamma}_{ m Wh} \ \mu^{\gamma\gamma}_{ m Zh} \ \mu^{\gamma\gamma}_{ m th} \ \mu^{\gamma\gamma}_{ m Zh} \ \mu^{\gamma\gamma}_{ m th} \ \mu^{\gamma\gamma}_{ m th}$	$\begin{array}{c} 1.04 \pm 0.10 \\ 1.20 \pm 0.26 \\ 1.5 \pm 0.55 \\ -0.2 \pm 0.55 \\ 0.89 \pm 0.31 \\ 3 \pm 3.5 \end{array}$	$ \begin{vmatrix} 1 & -0.13 & 0 & 0 & 0 & 0 \\ -0.13 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -0.37 & 0 & -0.11 \\ 0 & 0 & -0.37 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & -0.44 \\ 0 & 0 & -0.11 & 0 & -0.44 & 1 \end{vmatrix} $	139	[18]	$\mu_{ m ggh,bbh}^{\gamma\gamma} \ \mu_{ m VBF}^{\gamma\gamma} \ \mu_{ m Vh}^{\gamma\gamma} \ \mu_{ m Vh}^{\gamma\gamma} \ \mu_{ m th,th}^{\gamma\gamma}$	$1.07 \pm 0.11$ $1.04 \pm 0.32$ $1.34 \pm 0.34$ $1.35 \pm 0.31$		137	[11]
$\mu^{ZZ}_{ m ggF} \ \mu^{ZZ}_{ m VBF} \ \mu^{ZZ}_{ m VBF}$	$\begin{array}{c} 0.95 \pm 0.1 \\ 1.19 \pm 0.45 \\ 1.43 \pm 1.0 \end{array}$	$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	139	[4]	$\mu^{ZZ}_{ m ggh,bbh,tth,th}$ $\mu^{ZZ}_{ m VBF,Vh}$	$0.95 \pm 0.13$ $0.82 \pm 0.34$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	137	[10]
$\mu^{ZZ}_{ m tth}  onumber\ \mu^{ZZ}_{ m incl.}$	$1.69 \pm 1.45$ $1.0 \pm 0.1$	0 0 -0.18 1	139	[4]	$\mu^{WW}_{ m ggh}  onumber \ \mu^{WW}_{ m VBF}  onumber \ ww$	$0.92 \pm 0.11$ $0.71 \pm 0.26$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	120	[16]
$\mu^{WW}_{ m ggF,bbh}$ $\mu^{WW}_{ m VBF}$ $\mu^{WW}_{ m grF}$ $\mu^{WW}_{ m grF}$ $\mu^{WW}_{ m grF}$	$1.15 \pm 0.135$ $0.93 \pm 0.21$ $1.09 \pm 0.11$		139	[17]	$\mu^{WW}_{ m Zh} \ \mu^{WW}_{ m Wh}$	$2.0 \pm 0.7 \\ 2.2 \pm 0.6$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	130	
$\mu_{\rm VBF}^{\tau\tau}$ $\mu_{\rm ggF,bbh}^{\tau\tau}$ $\mu_{\rm Vh}^{\tau\tau}$ $\mu_{\rm Vh}^{\tau\tau}$ $\mu_{\rm tth,th}^{\tau\tau}$	$0.90 \pm 0.18$ $0.96 \pm 0.31$ $0.98 \pm 0.60$ $1.06 \pm 1.18$	$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$	139	[13]	$egin{aligned} \mu^{ au au}_{ ext{incl.}} \ \mu^{ au au}_{ ext{ggh}} \ \mu^{ au au}_{ ext{qqh}} \ \mu^{ au au}_{ ext{qqh}} \ \mu^{ au au}_{ ext{Vh}} \end{aligned}$	$0.93 \pm 0.12$ $0.97 \pm 0.19$ $0.68 \pm 0.23$ $1.80 \pm 0.44$		138	[15]
$\mu^{bb}_{ m VBF} \ \mu^{bb}_{ m Wh} \ \mu^{bb}_{ m Zh} \ \mu^{bb}_{ m bb}$	$0.95 \pm 0.37$ $0.95 \pm 0.26$ $1.08 \pm 0.24$		126 139 139	[9] [6] [6]	$\mu^{bb}_{ m qqh} \ \mu^{bb}_{ m ggh}$	$1.59 \pm 0.60$ $-2.7 \pm 3.89$	$egin{array}{cccc} 1 & -0.75 \ -0.75 & 1 \end{array}$	90.8	[19]
$\mu_{ m Vh} \ \mu_{ m tth,th}^{bb} \ \mu_{ m pp}^{\mu\mu}$	$ \begin{array}{c} 1.02 \pm 0.17 \\ 0.35 \pm 0.35 \\ 1.2 \pm 0.6 \end{array} $		139 139 139	[0] [12] [7]	$\mu^{\mu\mu}_{ m ggh,tth} \ \mu^{\mu\mu}_{ m VBF,Vh}$	$0.66 \pm 0.67$ $1.85 \pm 0.86$	$egin{array}{cccc} 1 & -0.24 \ -0.24 & 1 \end{array}$	137	[8]
$\mu^{Z\gamma}_{ m pp}$	$2.0\pm0.95$		139	[5]	$\mu^{Z\gamma}_{ m pp}$	$2.4\pm0.9$		138	[14]

Signal Value		Correlation matrix							
strength									
$\mu_{ m ggF}^{\gamma\gamma}$	$1.10\pm0.23$	1	-0.25	0	-0.14	0			
$\mu_{ m VBF}^{\gamma\gamma}$	$1.3\pm0.5$	-0.25	1	0	0	0			
$\mu_{ m Wh}^{\gamma\gamma}$	$0.5 \pm 1.3$	0	0	1	-0.64	0			
$\mu_{ m Zh}^{\gamma\gamma}$	$0.5 \pm 2.8$	-0.14	0	-0.64	1	-0.11			
$\mu_{ m tth}^{\gamma\gamma}$	$2.2 \pm 1.5$	0	0	0	-0.11	1			
$\mu^{ZZ}_{ m ggF}$	$1.13\pm0.33$	1	-0.26						
$\mu_{ m VBF}^{ZZ}$	$0.1 \pm 0.9$	-0.26	1						
$\mu^{WW}_{ m ggF}$	$0.84 \pm 0.17$	1	-0.16		enter 1				
$\mu^{WW}_{ ext{VBF}}$	$1.2 \pm 0.4$	-0.16	1						
$\mu_{ m ggF}^{ au au}$	$1.0 \pm 0.6$	1	-0.37	0	-0.25	0	-0.21		
$\mu_{ m VBF}^{ au au}$	$1.3 \pm 0.4$	-0.37	1	0	0	0	0		
$\mu^{WW}_{ m Wh}$	$1.6 \pm 1.1$	0	0	1	-0.12	-0.12	0		
$\mu_{ m Wh}^{ au au}$	$-1.4 \pm 1.4$	-0.25	0	-0.12	1	0	0		
$\mu^{WW}_{ m Zh}$	$5.9 \pm 2.4$	0	0	-0.12	0	1	0		
$\mu_{ m Zh}^{ au au}$	$2.2 \pm 2.0$	-0.21	0	0	0	0	1		
$\mu^{WW}_{ m tth}$	$5.0 \pm 1.8$	1	-0.47						
$\mu_{ ext{tth}}^{ au au}$	$-1.9 \pm 3.5$	-0.47	1						
$\mu^{bb}_{ m Wh}$	$1.0 \pm 0.5$								
$\mu^{bb}_{ m Zh}$	$0.4 \pm 0.4$								
$\mu^{bb}_{ m tth}$	$1.1 \pm 1.0$								
$\mu^{\mu\mu}_{ m pp}$	$0.1 \pm 2.5$				State 1				

h signal strength from official ATLAS and CMS combination for run 1