Vector boson scattering in the ATLAS detector



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The standard model of particle physics



- The theory describing elementary particles and their interactions
- Well tested experimentally

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Some very rare processes

Vector boson scattering (VBS)

• Electroweak production of vector bosons associated with jets: scattering («collision») of vector bosons

 $V_1 V_2 jj \rightarrow V_3 V_4 jj$

- V_i are electroweak gauge bosons (Z, W[±], photon)
- Quarks hadronize and form jets (j) in detectors



• High energy needed for such process: need to collide gauge bosons

VBS is very rich

Lots of different gauge couplings can be involved



Gauge couplings

• Not all couplings are allowed in the Standard Model (SM)



- Of course charge should be conserved at the vertex
- Furthermore there is no neutral couplings in the SM
- Search for deviation from SM in gauge couplings !

Gauge couplings

- VBS (and triboson) processes are the only ones that are sensitive to quartic gauge couplings
- Unique way of probing physics beyond the SM (BSM) affecting these couplings
 - New QGC (e.g. neutral ones)

Anomalous quartic gauge couplings (aQGC)

- Alteration of couplings existing in the SM
- Can be studied in the Effective Field Theory (EFT) framework

We haven't seen BSM physics for now

- It does not exist ? No, we know that SM is incomplete
- It is too weakly coupled ?
- It is hidden in SM backgrounds ?

Need to increase statistics and improve theory predictions

New accelerators

• It is present at too high energy for the current accelerators ?

EFT studies !

EFT approach





- Expand the SM Lagrangian (mass dimension 4) to higher dimensions
- Effective Lagrangian $\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \frac{1}{\Lambda}\mathcal{L}_5 + \frac{1}{\Lambda^2}\mathcal{L}_6 + \frac{1}{\Lambda^3}\mathcal{L}_7 + \frac{1}{\Lambda^4}\mathcal{L}_8 + \dots$

Dimension-n Lagrangians

• At a given dimension-n: $\mathcal{L}_n = \sum_i C_i^n Q_i^n$

Wilson coefficients

Dimension-n operators

- In the SM there is no high dimension term, Wilson coefficients are 0
- Operators are uniquely associated to Wilson coefficients and form a complete basis
- Odd-dimension operators violate lepton or baryon number conservation and are usually ignored

Dimension-n operators and aQGC

- Wilson coefficients associated to dimension-6 operators are constrained since dimension-6 can be probed with lots of analysis (different final states)
- Dimension-8 is not well known and can induce aQGC: VBS opportunity !
- Link with experiment ? Need observables from EFT
- For instance dimension-8 gives some amplitude (hence cross-section prediction):

$$\begin{split} \mathcal{A}^2 &= |\mathcal{A}_{SM}|^2 + 2\sum_i \frac{C_i}{\Lambda^4} Re(\mathcal{A}_i^*\mathcal{A}_{SM}) + 2\sum_i \frac{C_i^2}{\Lambda^8} |\mathcal{A}_i|^2 + 2\sum_{i\neq j} \frac{C_i C_j}{\Lambda^8} Re(\mathcal{A}_i^*\mathcal{A}_j) \\ \text{Pure SM} \quad \text{EFT-SM interference (linear)} \quad \text{Pure EFT (quadratic) Interference between EFT operators} \end{split}$$

Eboli model

• Complete classification of dimension-8 operators respecting symmetries

$$\begin{split} \mathcal{O}_{S,0} &= \left[(D_{\mu}\Phi)^{\dagger} D_{\nu}\Phi \right] \times \left[(D^{\mu}\Phi)^{\dagger} D^{\nu}\Phi \right] \\ \mathcal{O}_{S,1} &= \left[(D_{\mu}\Phi)^{\dagger} D^{\mu}\Phi \right] \times \left[(D_{\nu}\Phi)^{\dagger} D^{\nu}\Phi \right] \\ \mathcal{O}_{S,1} &= \left[(D_{\mu}\Phi)^{\dagger} D^{\mu}\Phi \right] \times \left[(D_{\nu}\Phi)^{\dagger} D^{\nu}\Phi \right] \\ \mathcal{O}_{M,2} &= \left[B_{\mu\nu}B^{\mu\nu} \right] \times \left[(D_{\beta}\Phi)^{\dagger} D^{\beta}\Phi \right] \\ \mathcal{O}_{M,4} &= \left[(D_{\mu}\Phi)^{\dagger} \widehat{W}_{\beta\nu}D^{\mu}\Phi \right] \times B^{\beta\nu} \\ \mathcal{O}_{M,5} &= \left[(D_{\mu}\Phi)^{\dagger} \widehat{W}_{\beta\nu}D^{\nu}\Phi \right] \times B^{\beta\mu} + h.c. \\ \mathcal{O}_{M,7} &= \left[(D_{\mu}\Phi)^{\dagger} \widehat{W}_{\beta\nu}\widehat{W}^{\beta\mu} \right] \times \mathrm{Tr} \left[\widehat{W}_{\alpha\beta}\widehat{W}^{\alpha\beta} \right] \\ \mathcal{O}_{T,2} &= \mathrm{Tr} \left[\widehat{W}_{\alpha\mu}\widehat{W}^{\mu\beta} \right] \times \mathrm{Tr} \left[\widehat{W}_{\alpha\mu}\widehat{W}^{\mu\beta} \right] \times B_{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,6} &= \mathrm{Tr} \left[\widehat{W}_{\alpha\nu}\widehat{W}^{\mu\beta} \right] \times B_{\mu\beta}B^{\alpha\nu} \\ \mathcal{O}_{T,6} &= \mathrm{Tr} \left[\widehat{W}_{\alpha\nu}\widehat{W}^{\mu\beta} \right] \times B_{\mu\beta}B^{\alpha\mu} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,9} &= B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\nu\alpha} \\ \mathcal{O}_{T,9} &= B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\nu\alpha} \\ \mathcal{O}_{T,9} &= B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\nu\alpha} \\ \mathcal{O}_{T,9} &= B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\nu\alpha} \\ \mathcal{O}_{T,9} &= B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\mu\alpha} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,9} &= B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\nu\alpha} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{\mu\nu}B^{\alpha\beta}B^{\alpha\beta} \\ \mathcal{O}_{T,8} &= B_{\mu\nu}B^{$$

• Different vertices impacted

	SM				Not SM				
Operators	WWWW	WWZZ	$WW\gamma\gamma$	$WW\gamma Z$	ZZZZ	$\mathrm{ZZZ}\gamma$	$ZZ\gamma\gamma$	$Z\gamma\gamma\gamma$	$\gamma\gamma\gamma\gamma\gamma$
FS0, FS1	1	1			1				
FM0, FM1, FM7	1	1	1	1	1	1	1		
FM2, $FM3$, $FM4$, $FM5$		1	1	1	1	1	1		
FT0, FT1, FT2	1	1	1	1	1	1	1	\checkmark	1
FT5, FT6, FT7		1	1	1	1	1	1	1	1
FT8, FT9					1	1	1	\checkmark	1

The ATLAS experiment



We look at the collision and decays products

Semileptonic final states

- Inclusive VVjj (V=W,Z) production in LHC Run-2 data
- Semileptonic final states: one gauge boson decays hadronically (quarks pair) and the other one decays leptonically (leptons pair)



What we want to measure



Analysis results

- Semileptonic VBS observed with a significance higher than 5σ
- Measurement compatible with SM signal expectation
- Observed fiducial cross section 33.0 ± 5.5 fb
- aQGC interpretation
- Publication during the next months



aQGC limits

- Inclusive analysis: all operators of the Eboli model can be constrained
- Dedicated EFT samples added to the Monte-Carlo (SM background and VBS signal)
- SM VBS is now a background for aQGC
- Constraints (95% CI) on different operators are very competitive

Operator	Expected [TeV ⁻⁴]	Observed [TeV ⁻⁴]	- 2
FS0/A ⁴	[-3.22, 3.22]	[-3.72, 3.73]	ATLAS Work In Progress
FS1/A ⁴	[-6.84, 6.85]	[-7.62, 7.63]	$\nabla = 13 \text{ TeV}, 139 \text{ fb}^{-1}$
FM0/Λ ⁴	[-1.12, 1.12]	[-1.20, 1.19]	1.5
FM1/Λ ⁴	[-3.24, 3.24]	[-3.77, 3.77]	
FM2/Λ ⁴	[-1.66, 1.66]	[-1.76, 1.76]	0.5 68% CL
FT0/Λ ⁴	[-0.20, 0.18]	[-0.24, 0.21]	
FT6/Λ ⁴	[-0.76, 0.72]	[-0.71, 0.68]	-0.0 -0.4 -0.2 0 0.2 0.4 0.0 f_{T8}/Λ^4 [TeV ⁻⁴]

QGC combination

- Run-2 analyses are going to be finalized (Run-3 already started in 2022)
- Lots of different and complementary ATLAS VBS analyses
 - Access to aQGC
 - Opportunity to constrain dimension-8 operators
 - VBS in a nice candidate for a combination !
- Started the effort with semileptonic, WW, WZ, Wy, Zy and ZZ analyses
- Complementary final states, hence operators, involved
- Expecting great results, improved 1D and 2D limits



• Challenges: harmonization and correlations between analyses...

Conclusion

- VBS processes probe the most fundamental structure of electroweak interactions
- They are very rare and provide high sensitivity to BSM physics affecting gauge and Higgs couplings
- This can be studied in the framework of EFT through anomalous gauge couplings
- The semileptonic final states analysis allows to study a lot of couplings
- It is part of the ATLAS Full Run-2 aQGC combination: great results expected !

THANK YOU !