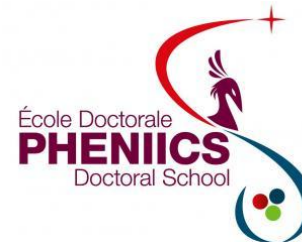


# Nuclear TMD analysis with SIDIS in RG-D Experiment at Jefferson Lab

Partonic Structure of Nuclei at the “Jlab/EIC” group in **IJCLab**  
**PHENIICS 2024**

DANIEL MATAMOROS

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# Motivations: Nuclear TMDs with CLAS12

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Understand the structure of nuclei in terms of quarks and gluons through 3D momentum space distribution of partons -> Transverse Momentum Distribution Functions (TMDs).

Decipher the effect of in-medium modifications TMDs in nuclei.

**This work:** Nuclear TMDs

- Study of the modification of final state interactions
- Cross section components can be linked to parton level effect

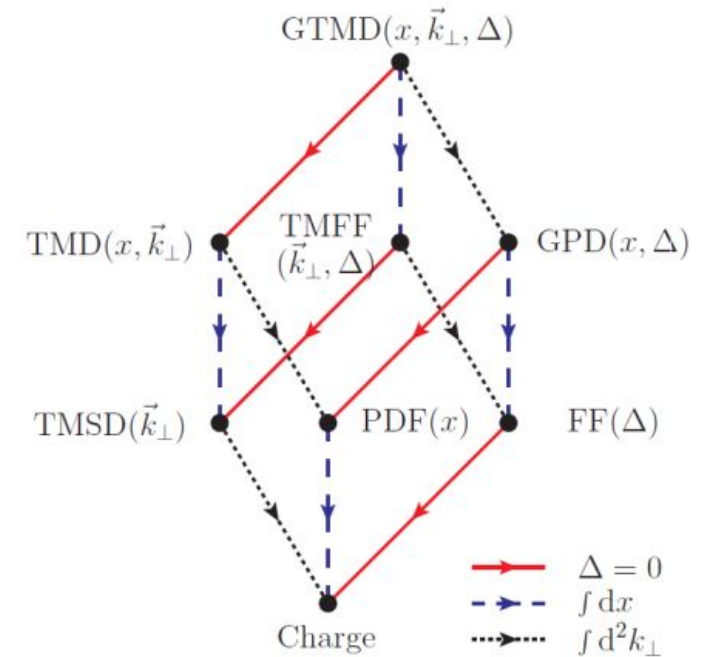
# Transverse Momentum Dependent (TMD) functions. **From PDFs to TMDs**

**PDF:** Describes the probability of finding a parton carrying a fraction of the nucleon's **longitudinal** momentum (**1D**)

**TMD:** Incorporate **transverse** momentum of partons allowing to study the partons' motion and interactions in greater detail (**3D**)

**This work:** Nuclear TMDs

- Study of the modification of final state interactions
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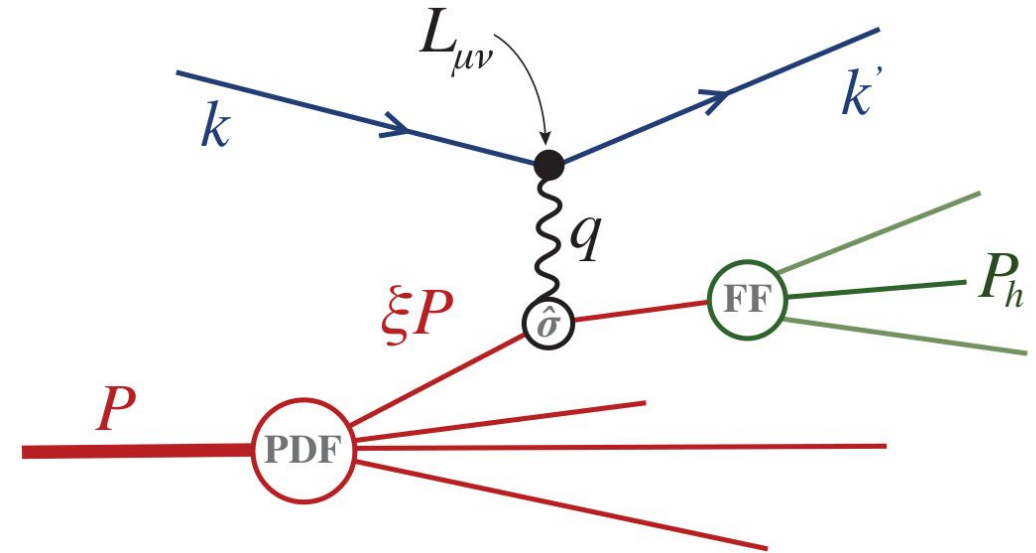
N/q	U	L	T
U	$f_1$		$h_1^{\perp}$
L		$g_1$	$h_{1T}^L$
T	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	$h_1 \quad h_{1T}^{\perp}$

# Semi Inclusive Deep Inelastic Scattering (SIDIS)

- Hadron production through  $\gamma^*$

$$e(k) + N(p) \rightarrow e(k') + X(p') + h(P_h)$$

- detection of one of the produced hadrons
- cross section  $\rightarrow$  convolution of TMD parton distributions and TMD fragmentation functions

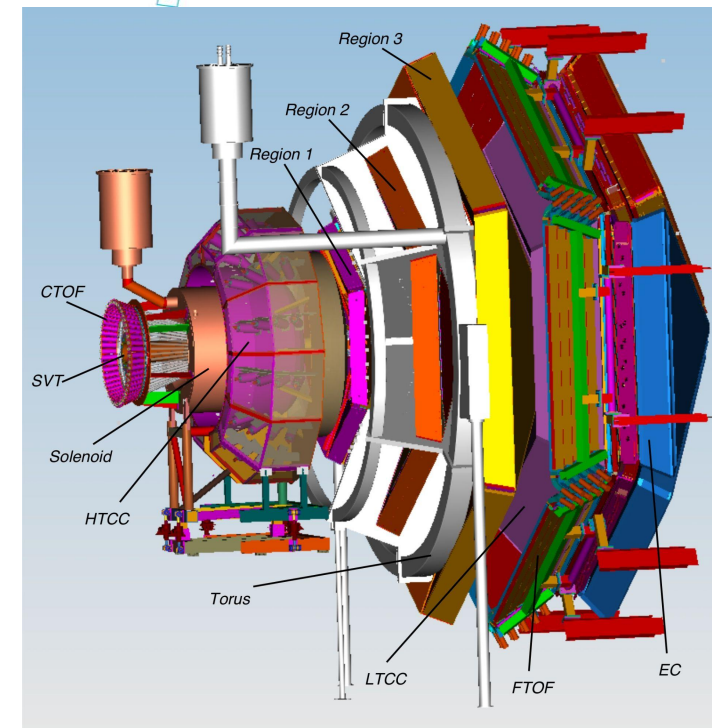


# TMD study with JLab's CLAS12

- Located at Jefferson Lab in Newport News (VA), USA
- Use of the 12 GeV electron beam at JeffersonLab
- Use of the CLAS12 spectrometer in Hall B.

Using the CLAS12 spectrometer @ JeffersonLab provides several advantages for SIDIS study

- Use of multiple targets, and possibility of beam polarization.
- Large Acceptance allowing the detection of scattered leptons and produced hadrons over a wide range of angles
- Precise measurements of various kinematic variables, including momentum, angle, and transverse momentum



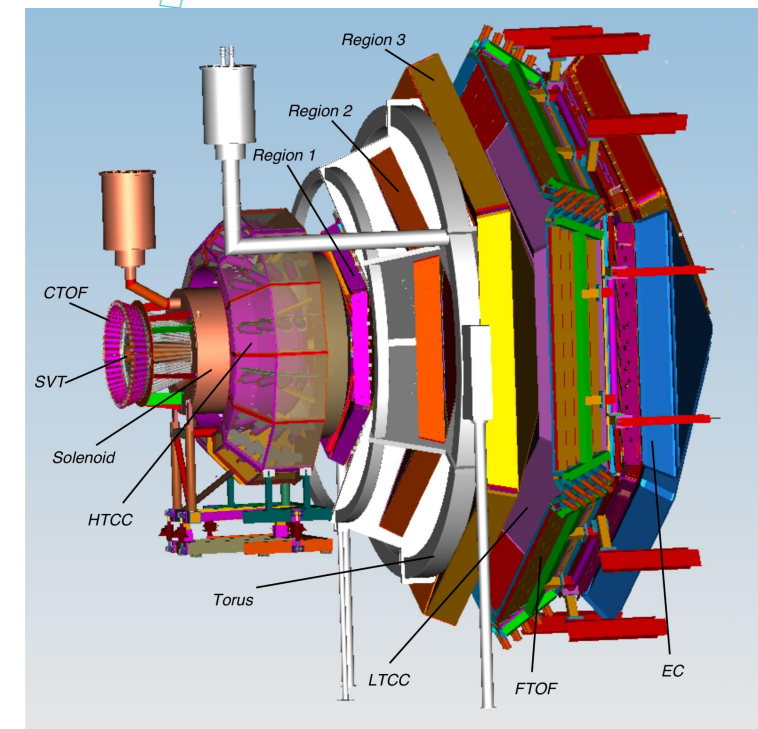
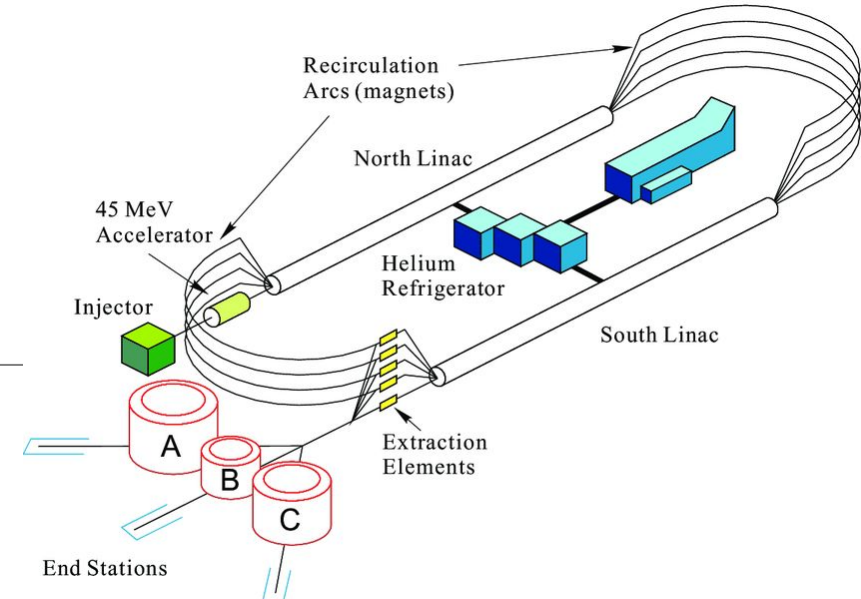


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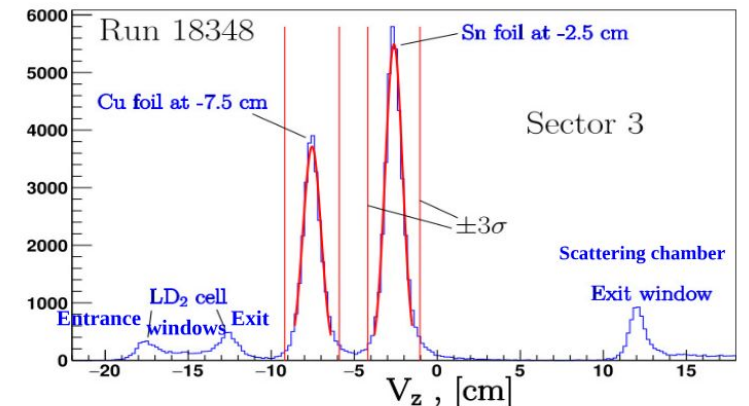
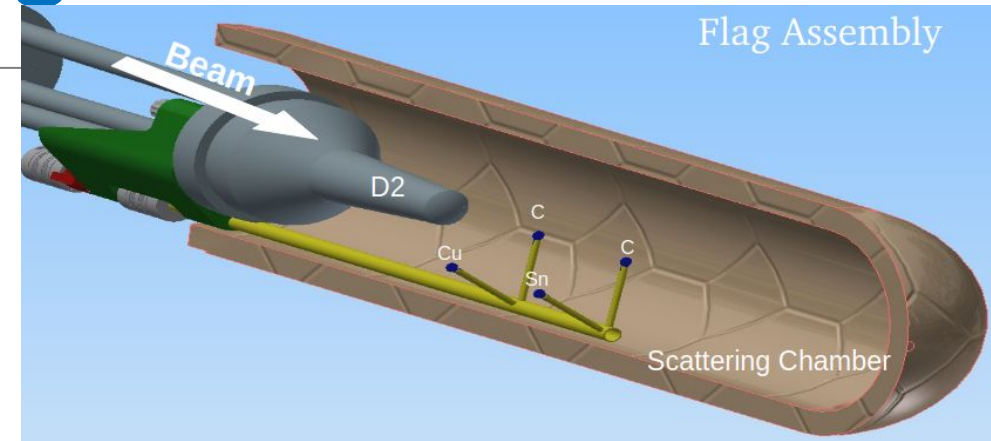


# Experimental configuration

10-11 GeV electron beam on nuclear targets

- Targets: nuclear solid foils; **Deuterium**, Copper; Carbon and Tin.
- we will develop analysis tools to monitor the reconstructed particles yields such as  $e^-$ ,  $\pi$ ,  $k$ , etc

**This work:** Successful completion of data taking (Oct. 04, 2023 - Dec. 15, 2023). Data Calibration ongoing  
Analysis Implemented:  $\pi^+$  selection



calibration cooking  $V_z$  vertex reconstruction

# Experimental Observables

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Comparison study of data with a nucleon target N and a nuclear target A

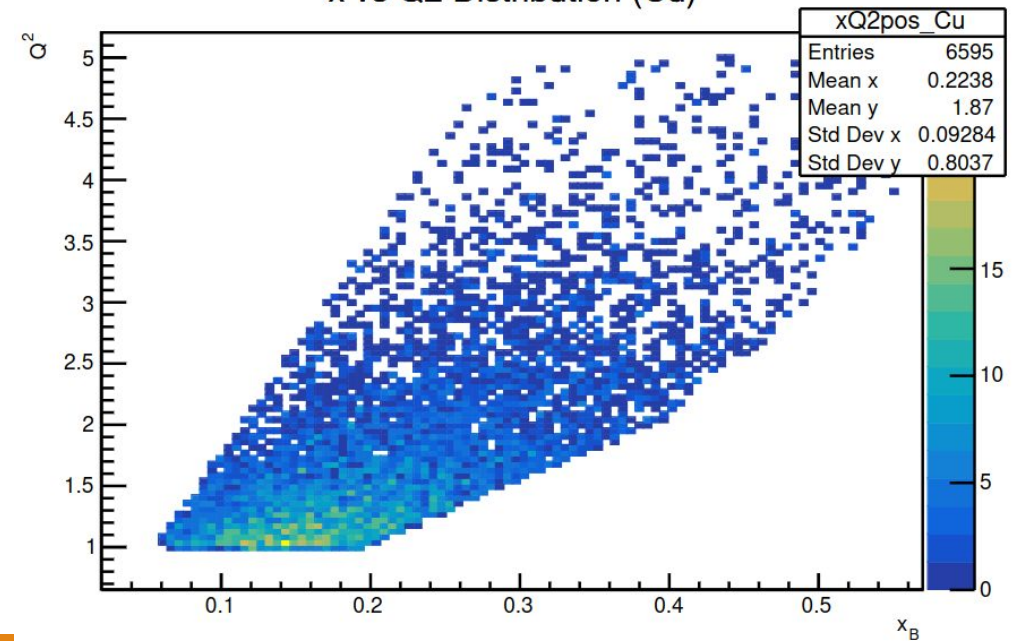
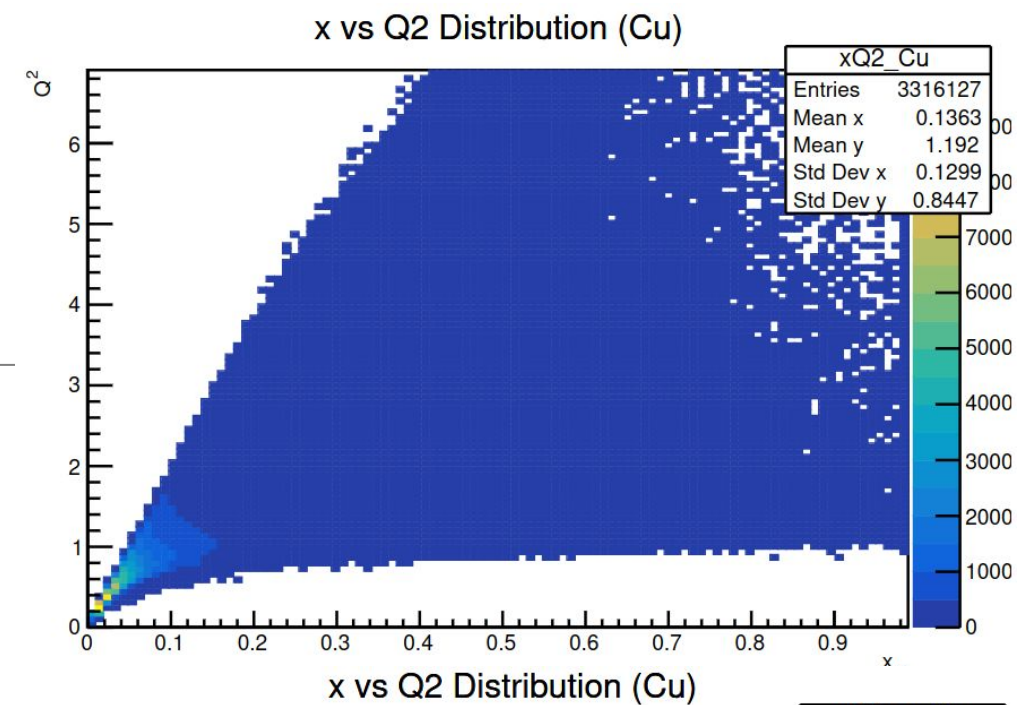
- Multiplicity Ratio R: 
$$R_A^h = \frac{(N_h/N_e)_A}{(N_h/N_e)_D}$$
- $p_t^2$  difference: 
$$\Delta \langle p_t^2 \rangle = \langle p_t^2 \rangle_A - \langle p_t^2 \rangle_N$$

Both Observables can be studied as a function of different kinematic variables.



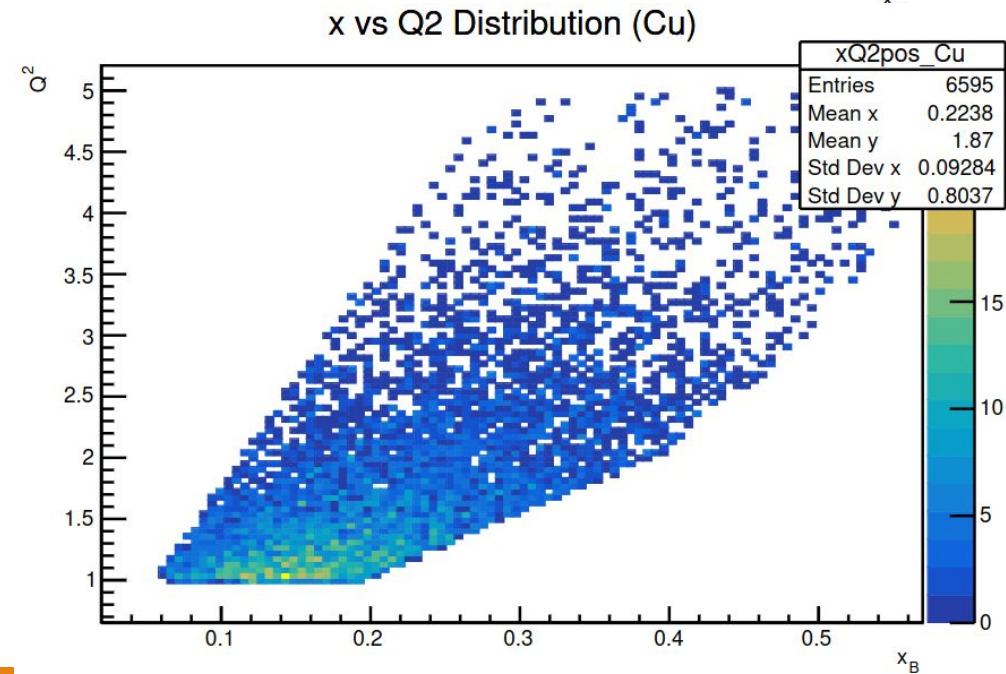
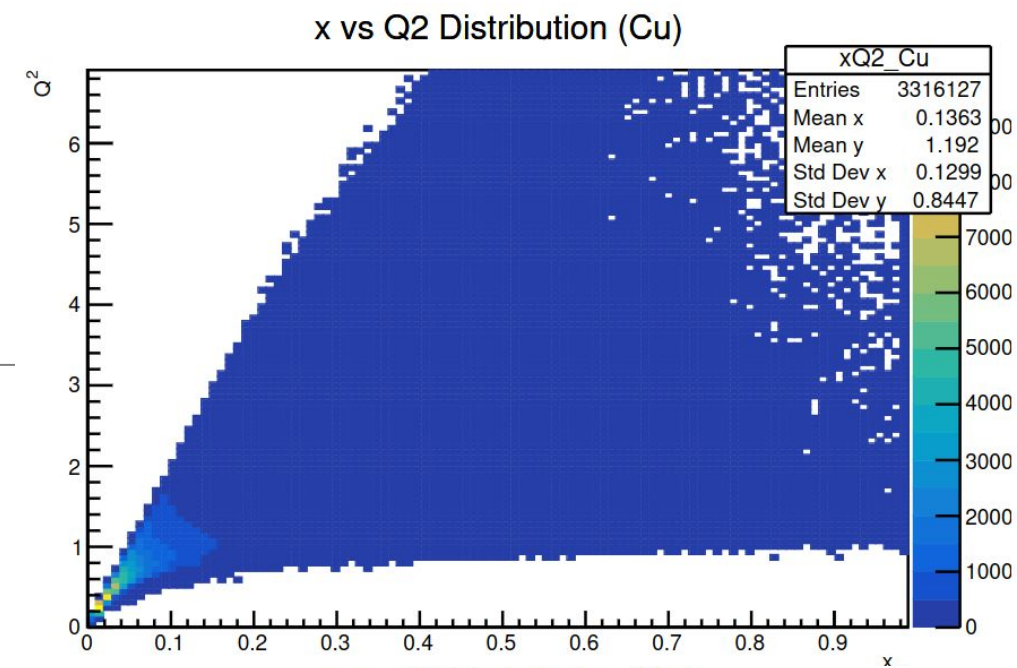
# Kinematic Variables

- Considering first coincidences on  $e^-$  and  $\pi^+$
- different kinematic variables can be considered we will focus here on the following variables from the **scattered electron**.
  - $Q^2 = -\gamma^2$
  - $v = E_{e^-} - E_{e'}$
  - $x_b = Q^2 / (2.M.v)$



# Data Selection

- Considering first coincidences on  $e^-$  and  $\pi^+$
- Kinematic variables considered and cuts for example:
  - $Q^2 > 1 \text{ GeV}^2$
  - Setting a missing mass threshold
  - Specific detector cuts

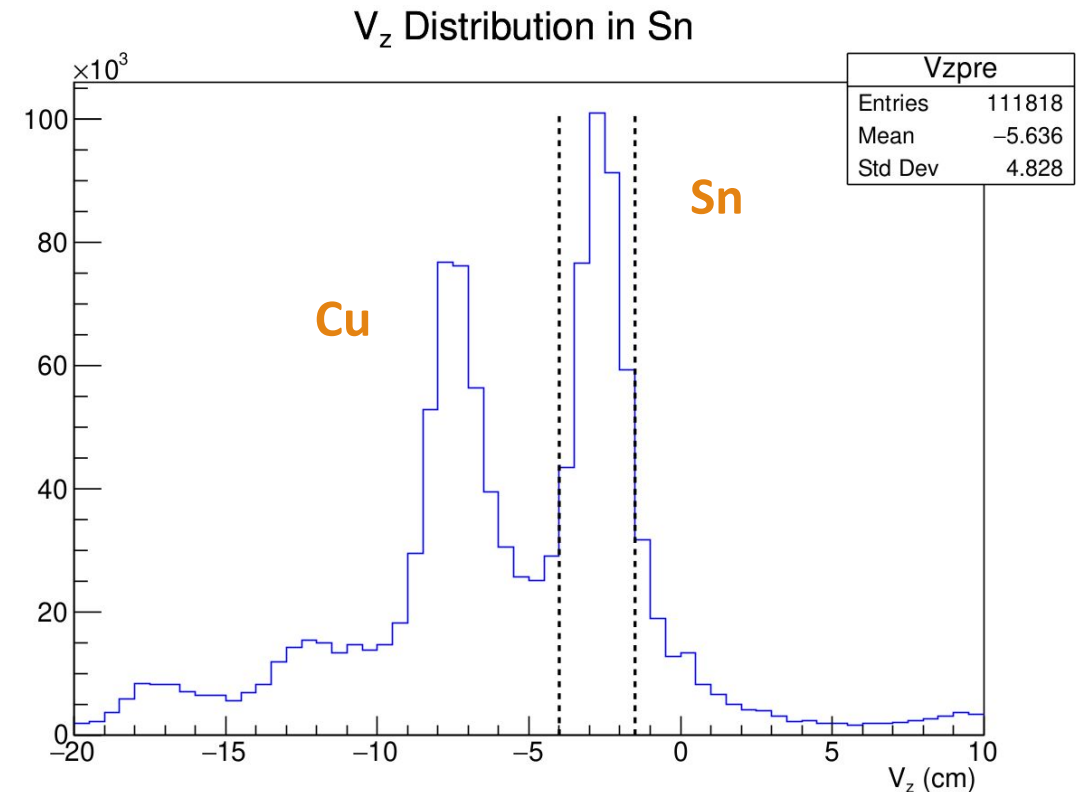


# specific preliminary analysis

- Monitoring data from available data
- Events considered with  $\pi^+$  production
- Kinematic Variables specific to the **hadron** for TMDs:
  - $z$  = Fraction of the virtual photon energy carried by the hadron.
  - $p_t^2$  = transverse momentum of hadrons
- Vertex z cuts need to be considered according to target positions
  - Preliminary Arbitrary Cuts on  $V_z$  for all targets

Using determined variables to determine multiplicity Ratio as follows:

$$R_A^\pi(Q^2, \nu, z, p_t^2) = \frac{N_\pi^{Sn}(Q^2, \nu, z, p_t^2) / N_e^{Sn}(Q^2, \nu)}{N_\pi^D(Q^2, \nu, z, p_t^2) / N_e^D(Q^2, \nu)}$$



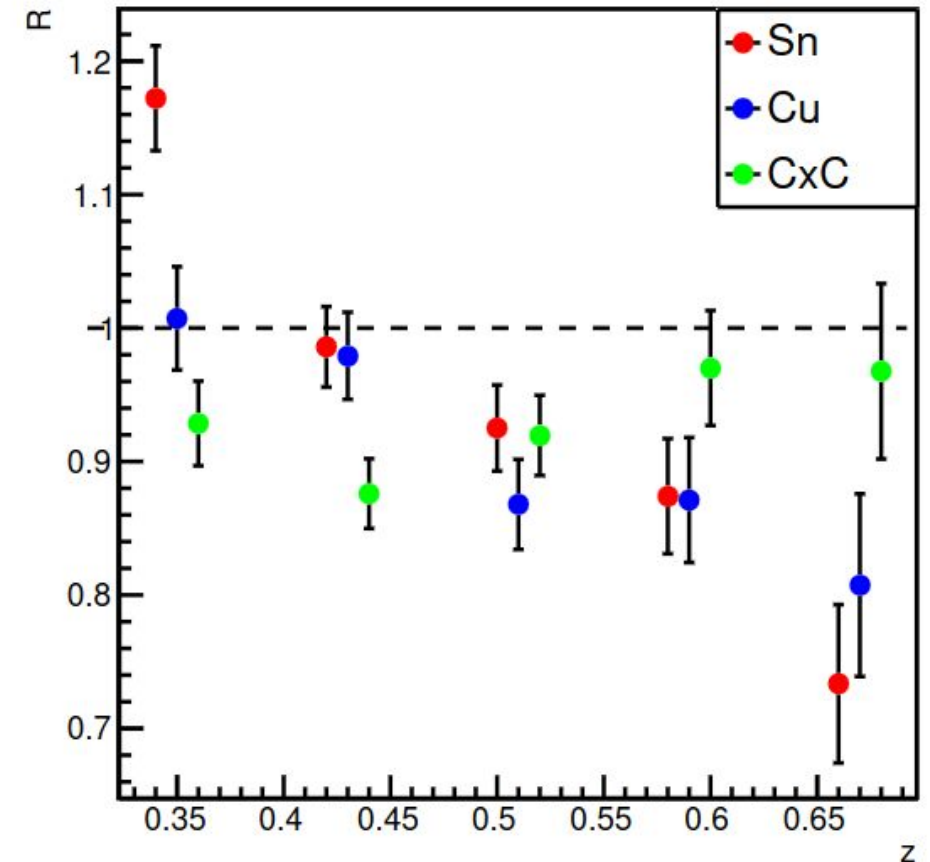
# Multiplicity Ratio

$$R_A^\pi(\nu, z, p_t^2) = \frac{N_\pi^{Sn}(\nu, z, p_t^2)/N_e^{Sn}(\nu)}{N_\pi^D(\nu, z, p_t^2)/N_e^D(\nu)}$$

R vs z, pt2=0.75

Multi dimensional analysis can be implemented in order to consider cross-variable correlations.  
Currently considering a three-Dimensional analysis with  $\nu$ ,  $z$  and  $pt^2$   
Analysis with no corrections implemented (on going)

Here, plot of multiplicity Ratio for different nuclei targets compared to reference Deuterium for  $\nu=4.5$ ;  $pt^2 = 0.75$  as a function of  $z$ .



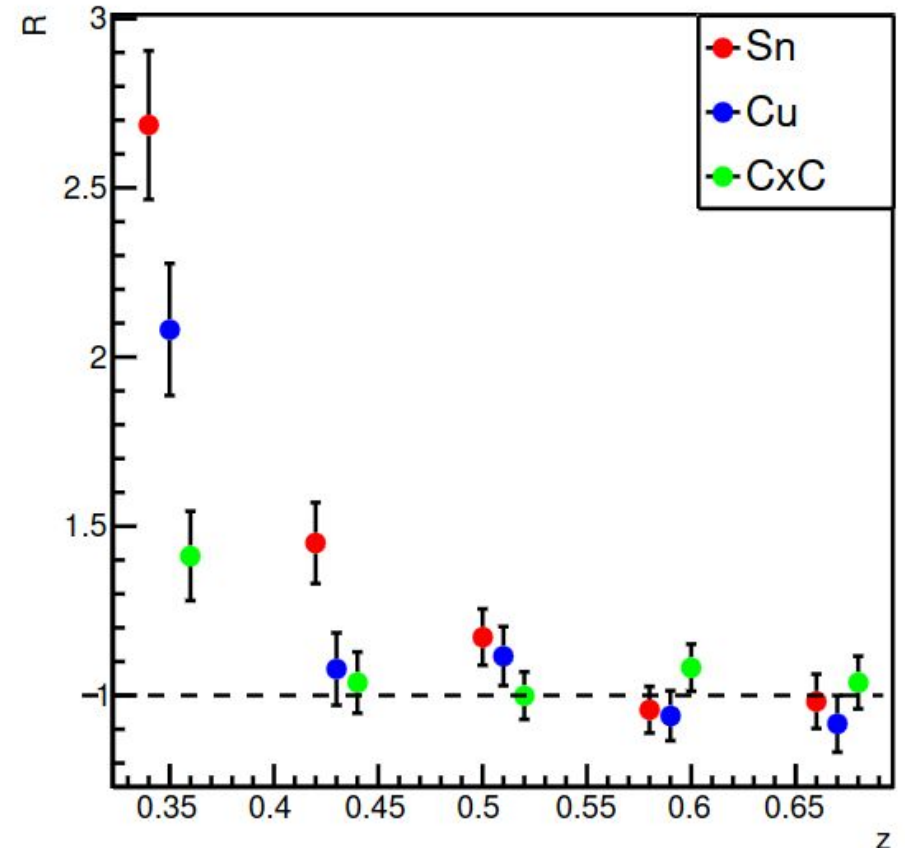
# Multiplicity Ratio $R_A^\pi(Q^2, \nu, z, p_t^2) = \frac{N_\pi^{Sn}(Q^2, \nu, z, p_t^2) / N_e^{Sn}(Q^2, \nu)}{N_\pi^D(Q^2, \nu, z, p_t^2) / N_e^D(Q^2, \nu)}$

Multi dimensional analysis can be implemented in order to consider cross-variable correlations.  
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 Analysis with no corrections implemented (on going)

Here, plot of multiplicity Ratio for different nuclei targets compared to reference Deuterium for  $\nu=4.5$ ;  $p_t^2 = 1.35$  as a function of  $z$ .

We can observe a deviation for high values of transverse momentum. But non conclusive results yet, calibrations and corrections remain to be implemented

R vs z, pt2=1.35

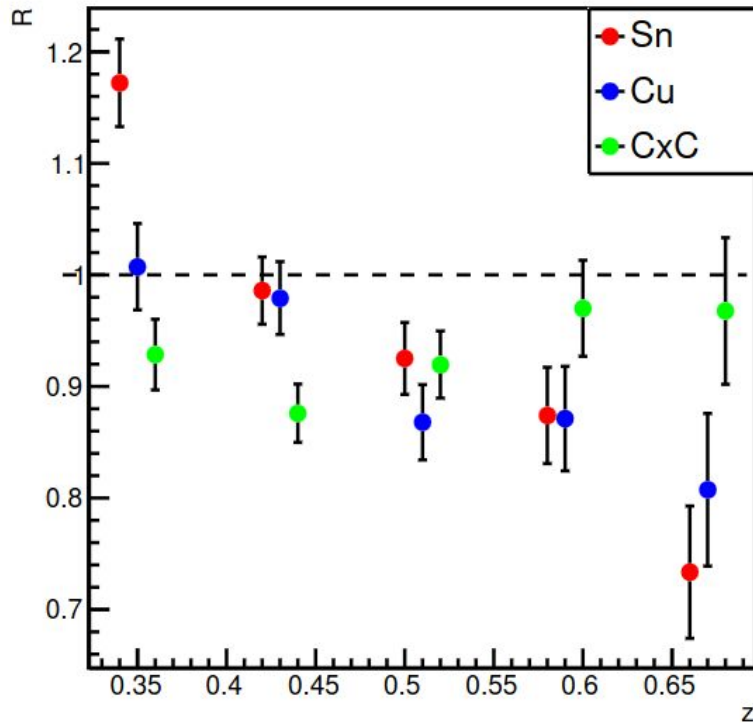




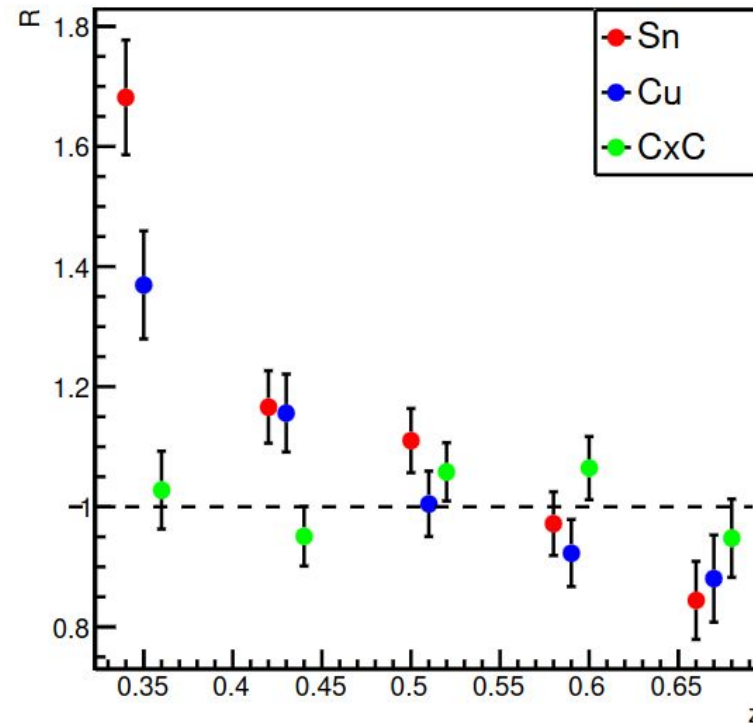
# Multiplicity Ratio

$$R_A^\pi(Q^2, \nu, z, p_t^2) = \frac{N_\pi^{Sn}(Q^2, \nu, z, p_t^2) / N_e^{Sn}(Q^2, \nu)}{N_\pi^D(Q^2, \nu, z, p_t^2) / N_e^D(Q^2, \nu)}$$

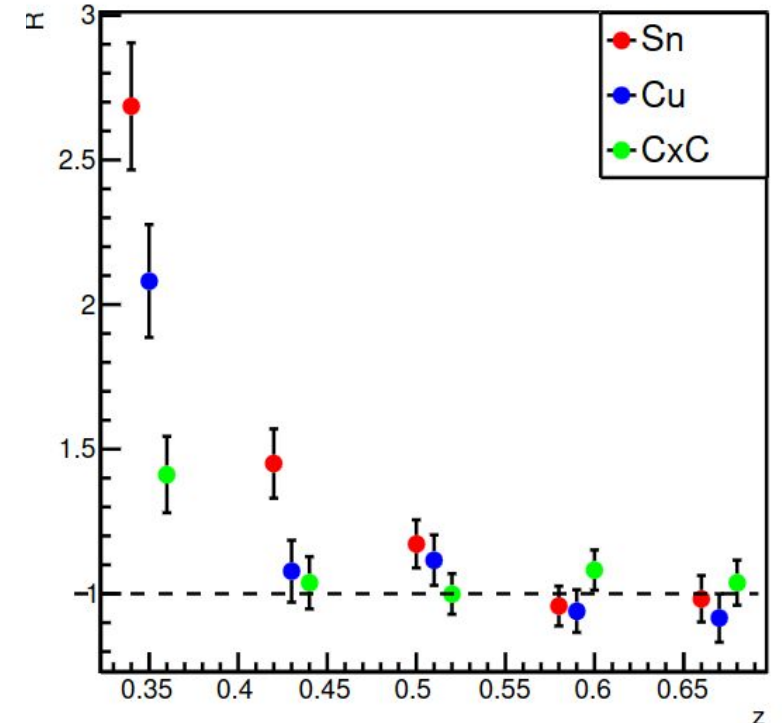
R vs z, pt2=0.75



R vs z, pt2=1.05



R vs z, pt2=1.35



Preliminary Studies show coherent behaviour between targets.  
Potentially large amount of results with small uncertainties.



# Summary

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## Implemented work

- Preliminary Data analysis for a study of TMDs on SIDIS with most recent available data.
- Coherent/expected results on available data with initial analysis.

## Study of the Simulation

- [In preparation] Very preliminary tests of a simulation non specific for TMD study nuclear effects
- Simulation used to properly compare with available data replicating experimental conditions as possible

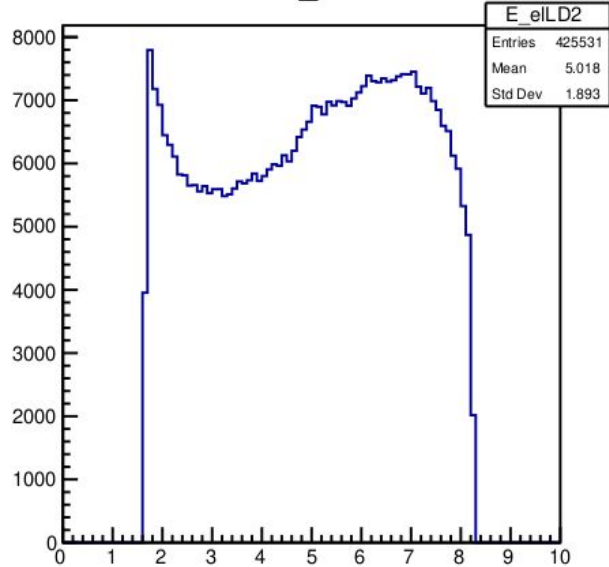
## On going work

- Calibration on available data remains to be done/perfected.
- Acceptance and Radiative effects need to be considered for corrections.
- Consider other experimental observables

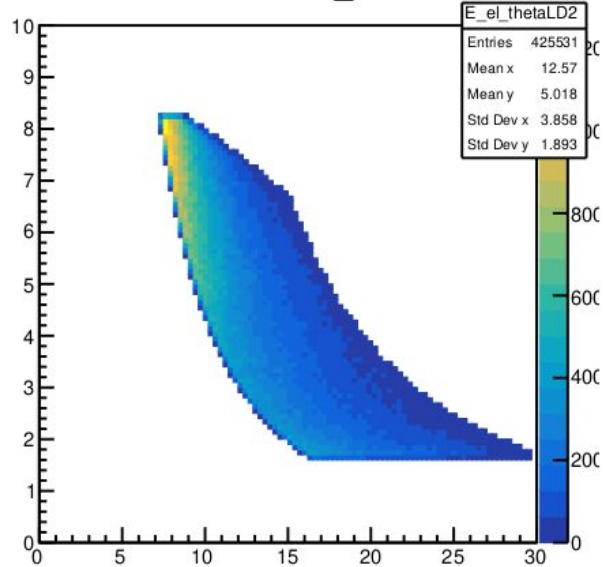
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# Back Ups

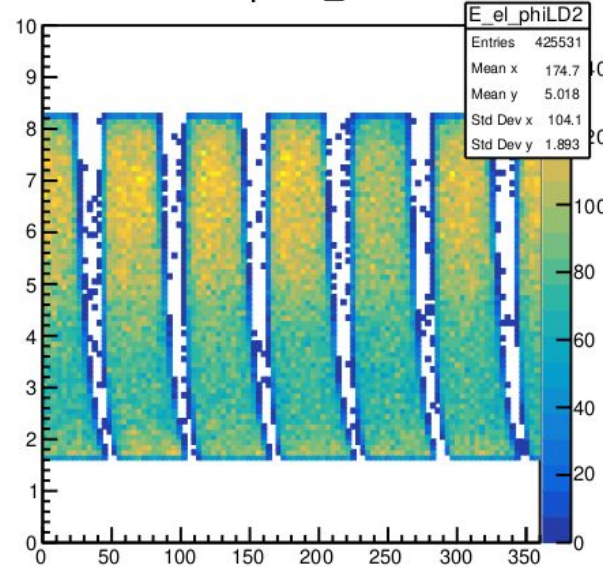
E el\_LD2



E vs theta el\_LD2



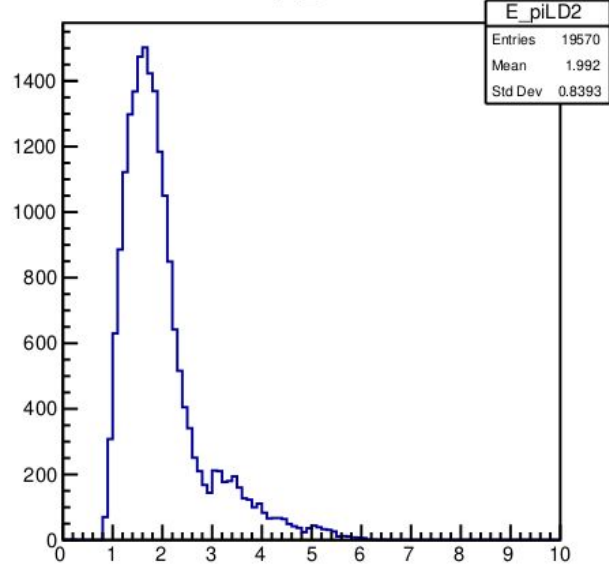
E vs phi el\_LD2



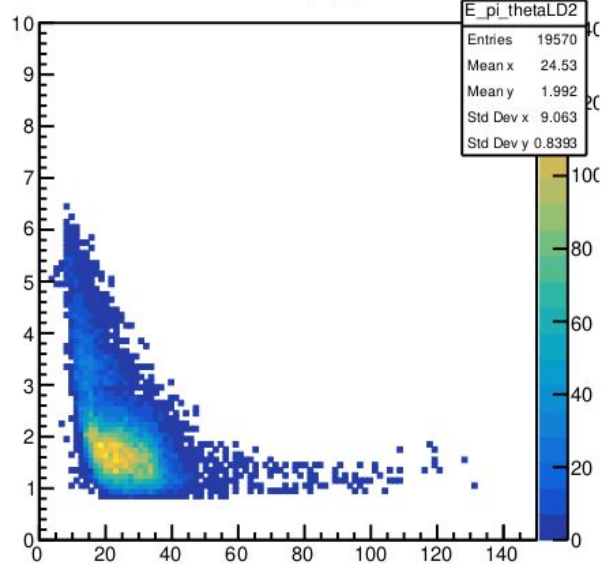
# Particle Energies

Electron energy vs azimuthal and polar angle

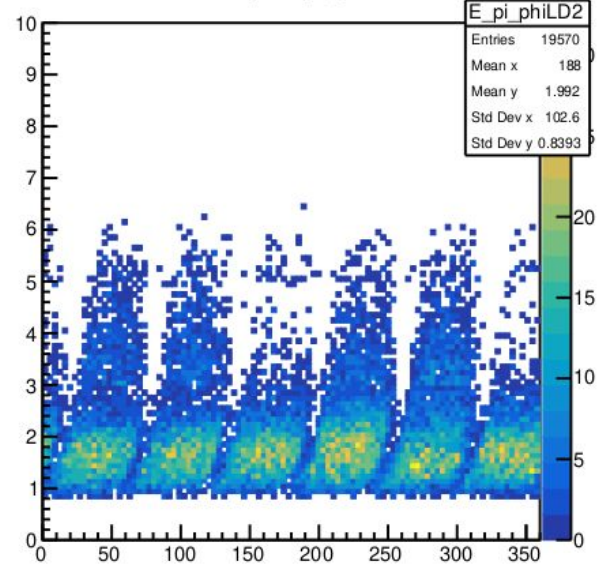
E pi\_LD2



E vs theta pi\_LD2



E vs phi pi\_LD2

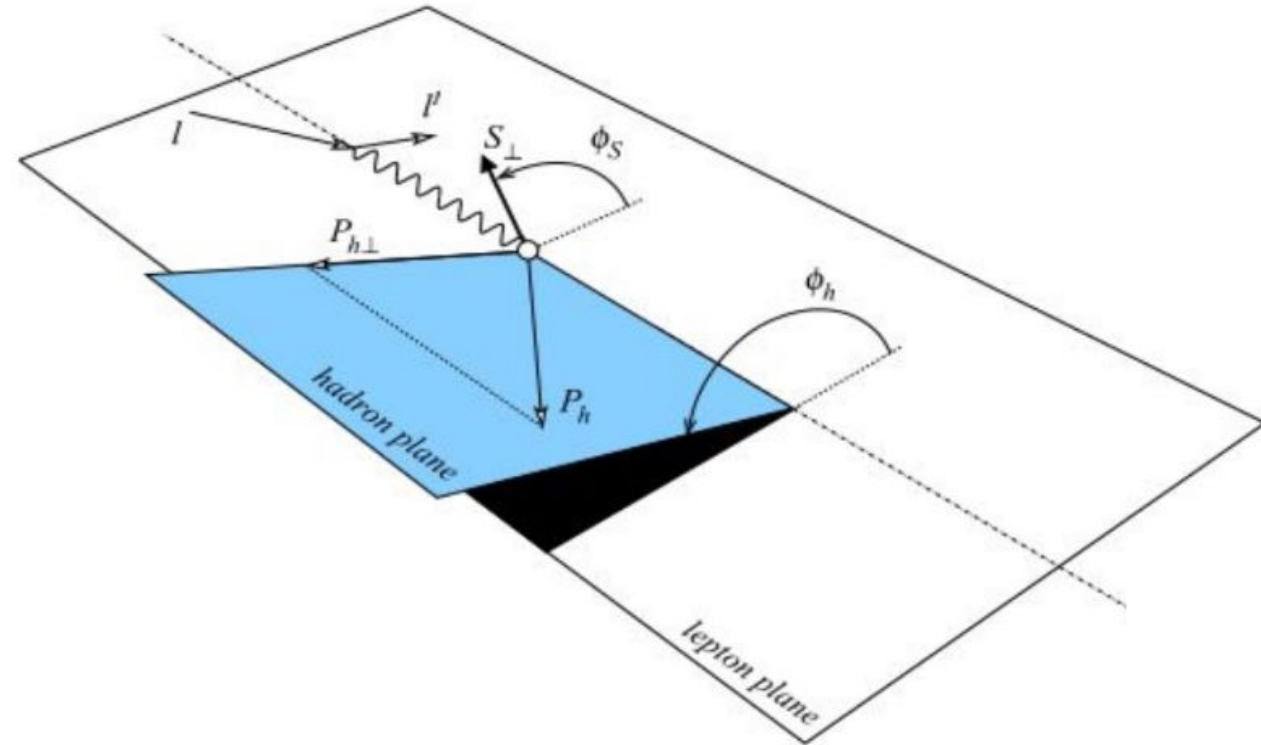


Pion energy vs azimuthal and polar angle

# Observables

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- Experimental Observables: Cross section, Beam Spin Asymmetry

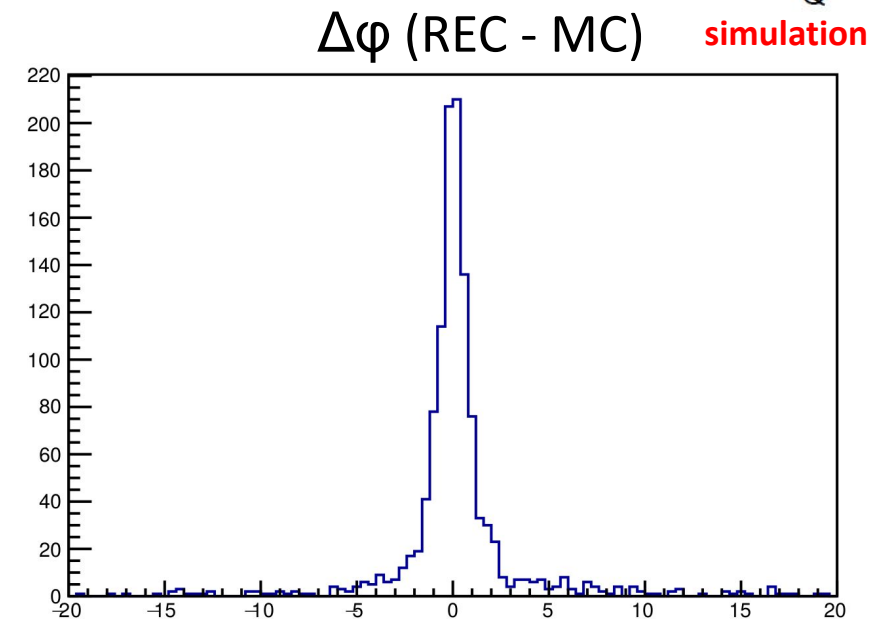
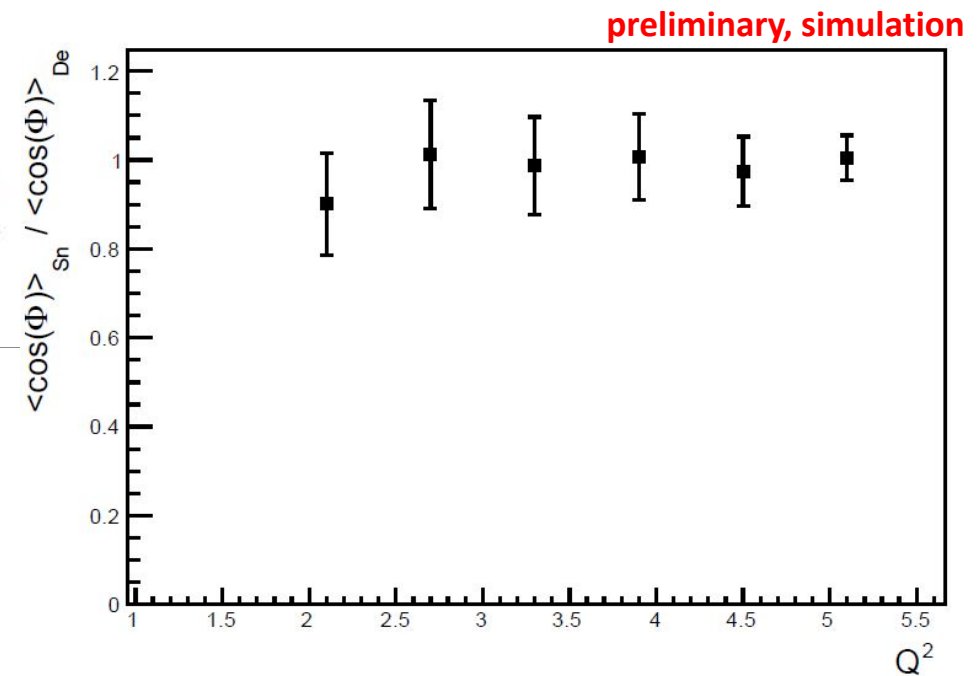


**This Work:** We use unpolarized nuclear targets where only the  $\cos \varphi$ ,  $\cos 2\varphi$  and  $\sin \varphi$  components will contribute in this cross section.

# cos $\Phi$ ratio

$$\frac{\langle \cos\phi_{\pi} \rangle_{Sn}}{\langle \cos\phi_{\pi} \rangle_D}$$

- $\phi$  resolution determined by comparison of Simulated and reconstructed  $\phi$  values
- cos  $\phi$  and sin  $\phi$  ratios remain to be done
- Preliminary result
- Importance of Radiative effects remain to be determined



$$\begin{aligned}
\frac{d\sigma}{dx dy dz d\phi d\phi_S dP_{h\perp}^2} = & \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi F_{UU}^{\cos\phi} \right. \\
& + \varepsilon \cos(2\phi) F_{UU}^{\cos 2\phi} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi F_{LU}^{\sin\phi} \\
& + S_{\parallel} \left[ \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi F_{UL}^{\sin\phi} + \varepsilon \sin(2\phi) F_{UL}^{\sin 2\phi} \right] \\
& + S_{\parallel} \lambda_e \left[ \sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi F_{LL}^{\cos\phi} \right] \\
& + |\mathbf{S}_{\perp}| \left[ \sin(\phi - \phi_S) \left( F_{UT,T}^{\sin(\phi-\phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi-\phi_S)} \right) \right. \\
& + \varepsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi+\phi_S)} + \varepsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi-\phi_S)} \\
& \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi-\phi_S)} \right] \\
& + |\mathbf{S}_{\perp}| \lambda_e \left[ \sqrt{1-\varepsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi-\phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
& \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi-\phi_S)} \right] \left. \right\},
\end{aligned}$$