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UNIVERSITÄT
MÜNCHEN

New Constraints from the Abundance of South Pole Telescope-selected Clusters and the Large-Scale Structure

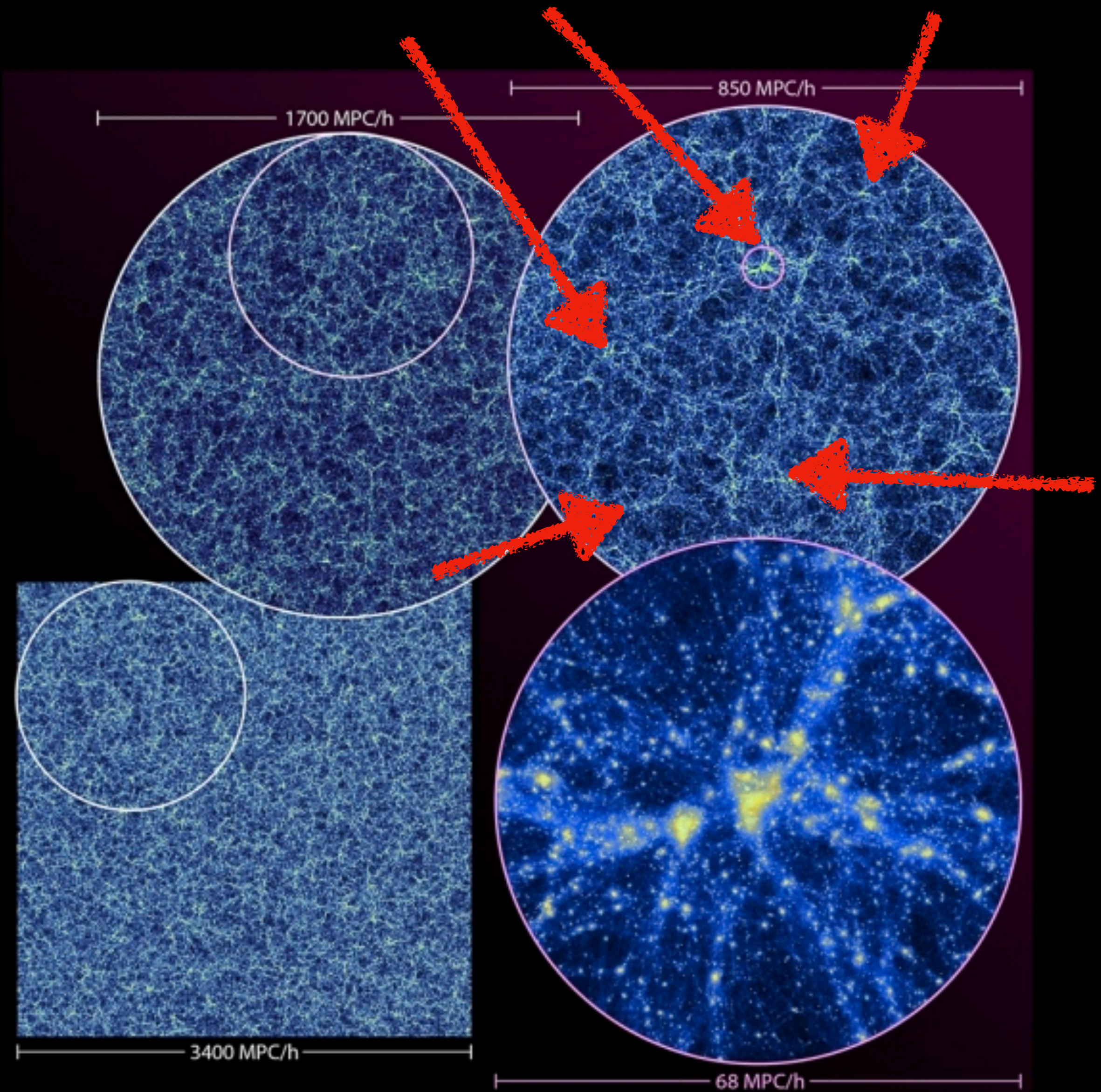
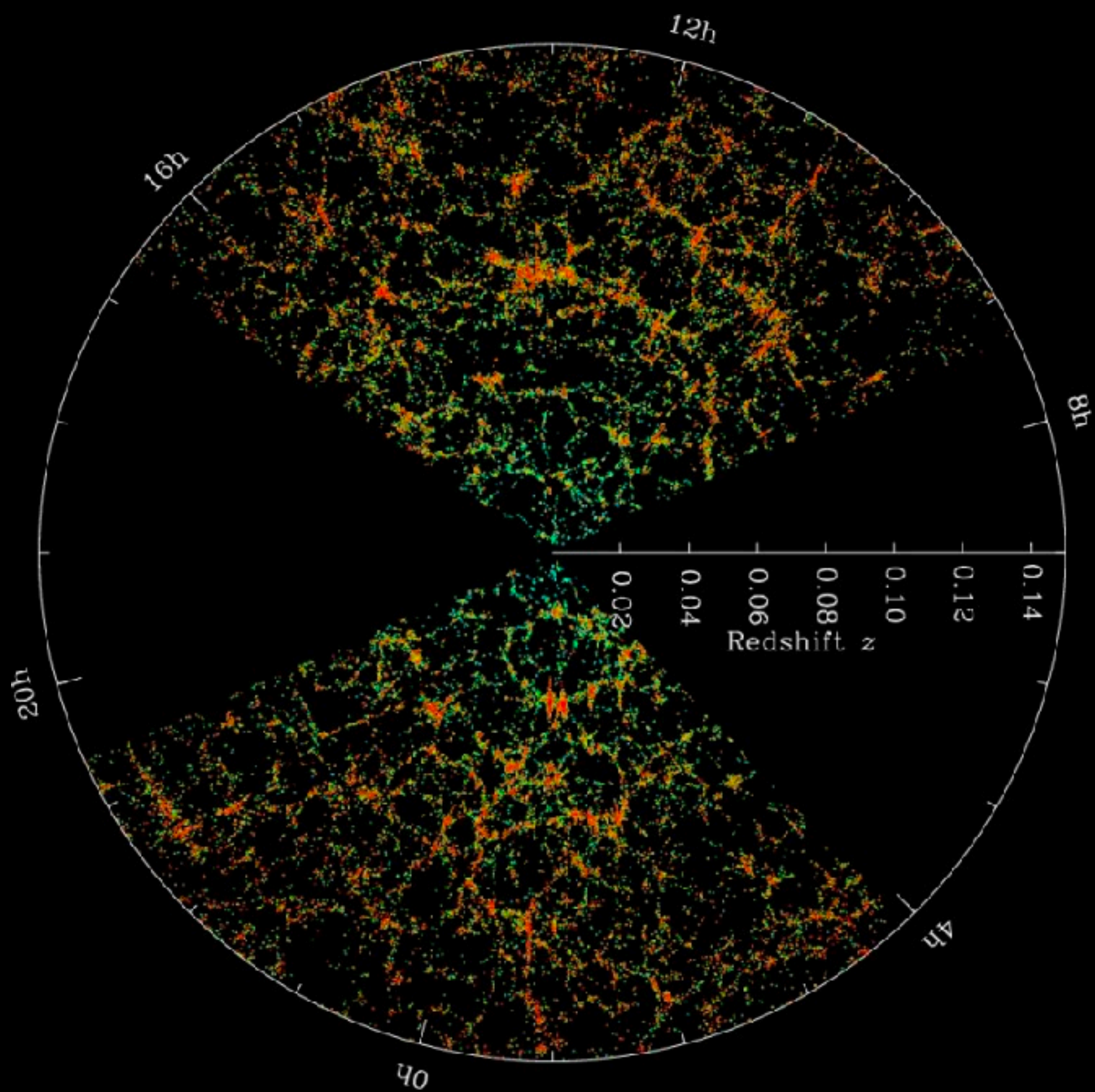
Sebastian Bocquet, LMU Munich

with Sebastian Grandis, Lindsey Bleem, Matthias Klein, Joe Mohr, Tim Schrabback,
Elisabeth Krause, Chun-Hao To,
and the *South Pole Telescope (SPT)* and *Dark Energy Survey (DES)* collaborations



Image credit: SPT 2024 winter-overs Josh + Kevin

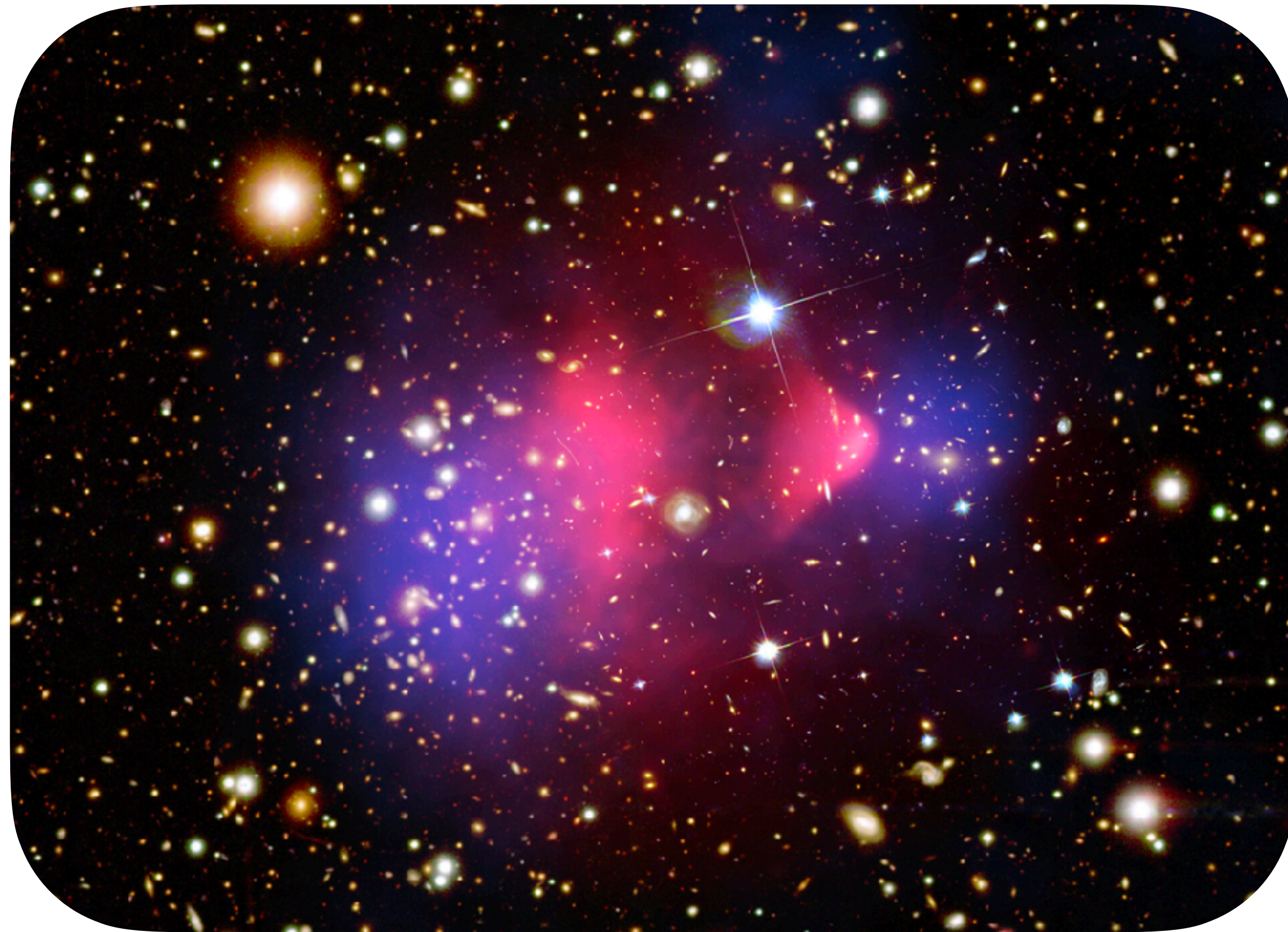
Massive Halos $\approx 10^{14} M_{\text{sun}}$... trace the large-scale structure



Last Journey (on Mira supercomputer) (Heitmann+)

Cluster Cosmology

The most massive collapsed objects $\gtrsim 10^{14} M_{\odot}$

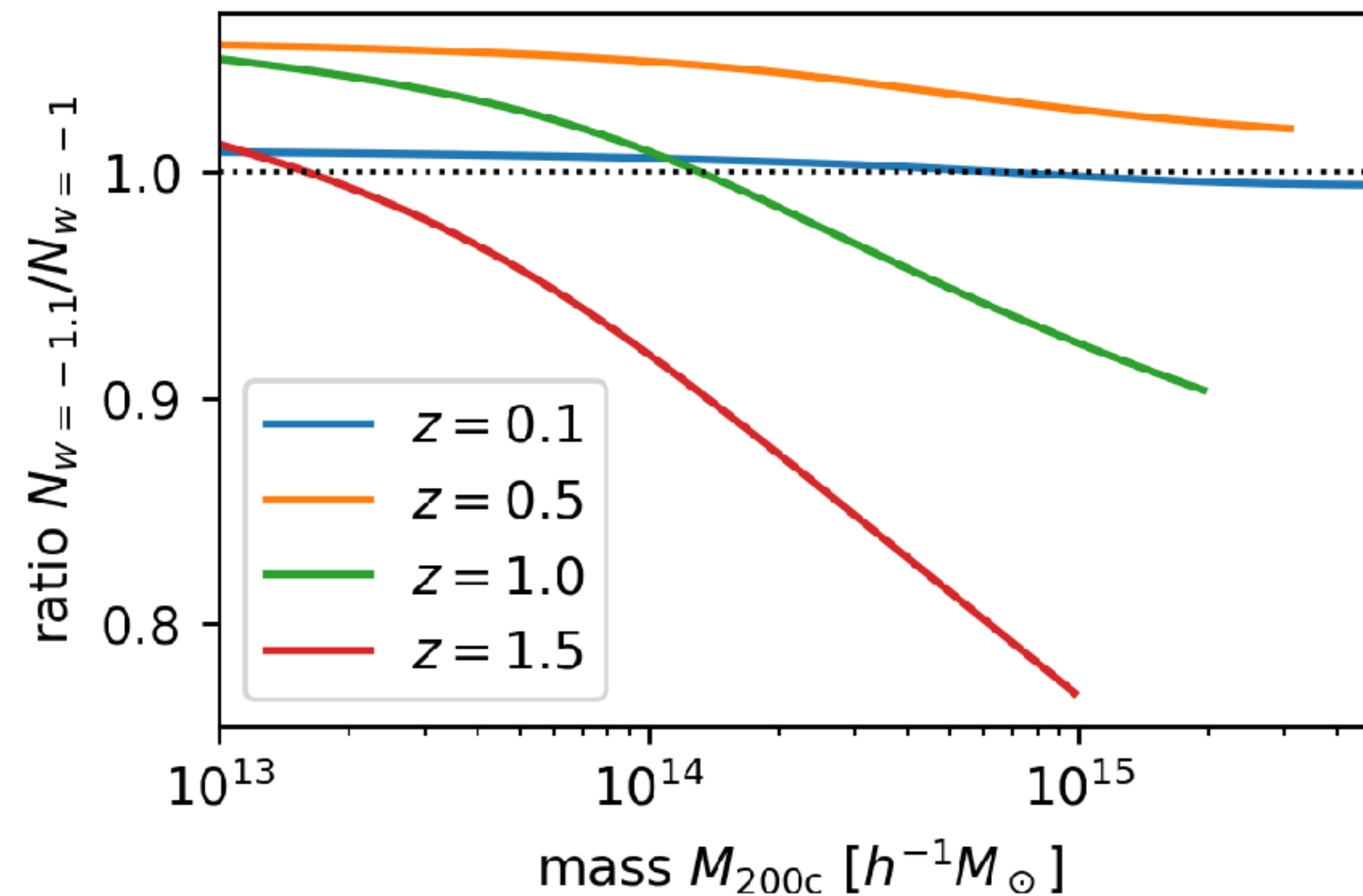


Bullet Cluster. X-ray: NASA/CXC/CfA/M.Markevitch, Optical and lensing map: NASA/STScI, Magellan/U.Arizona/D.Clowe, Lensing map: ESO WFI
AstroParticle Symposium 2024

- Composition
 - 85–90% dark matter
 - 10–15% ordinary matter, of which
 - $\sim 75\%$ (gravitationally heated) gas
 - $\sim 25\%$ galaxies/stars
- Somewhat arbitrary (but useful) definition
 - Halo \equiv *entire* thing
 - Cluster \equiv galaxies & gas (what we see)

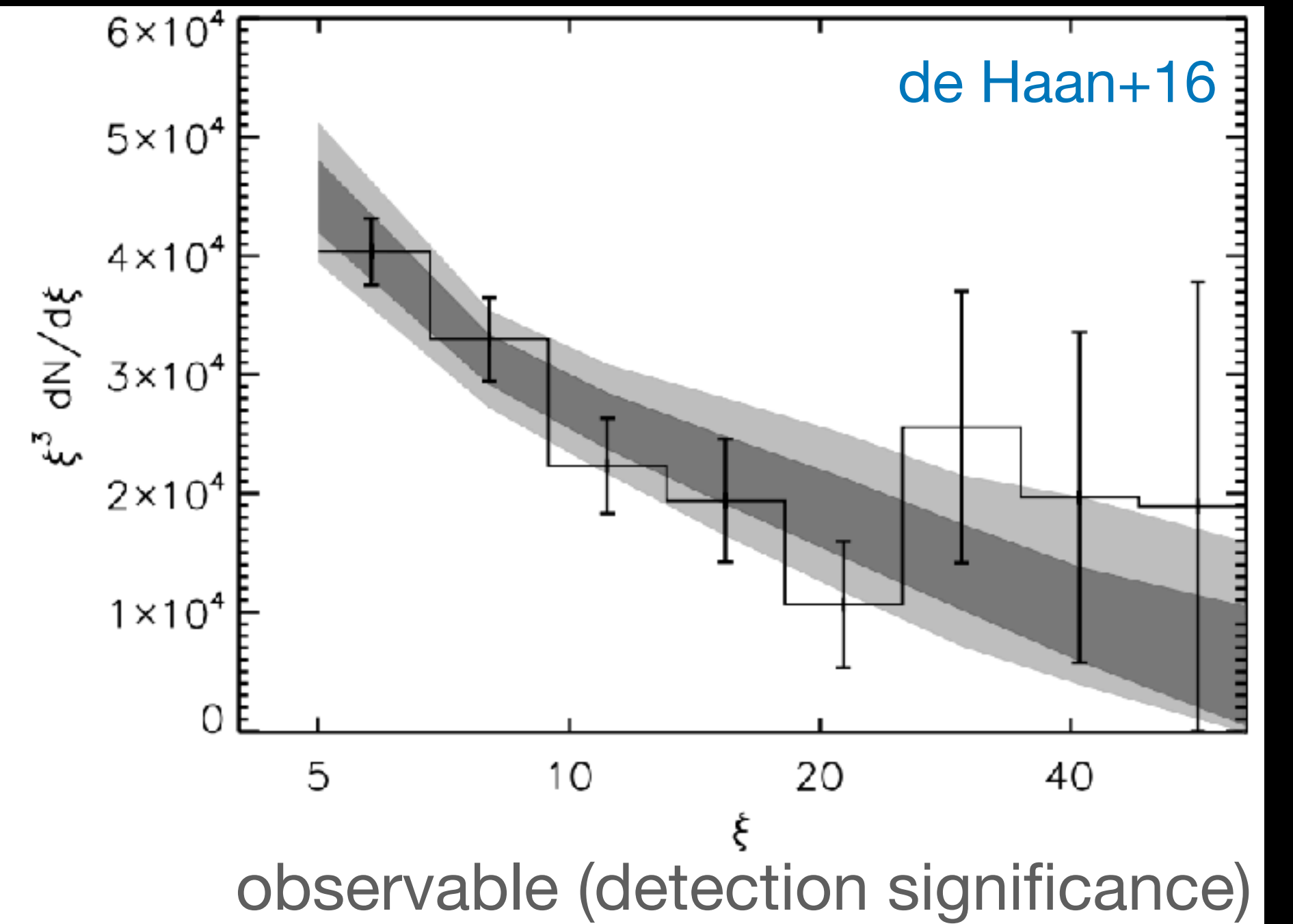
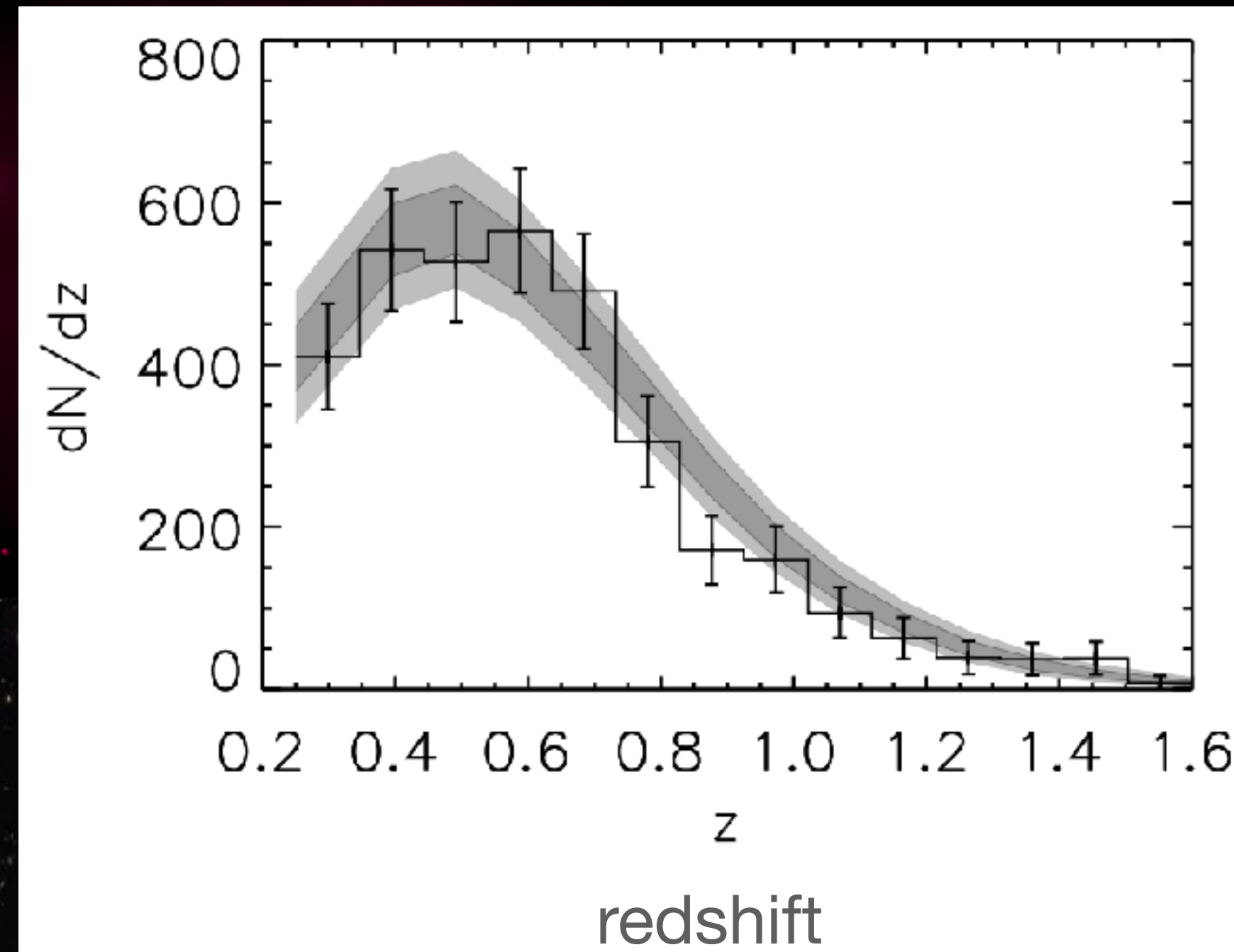
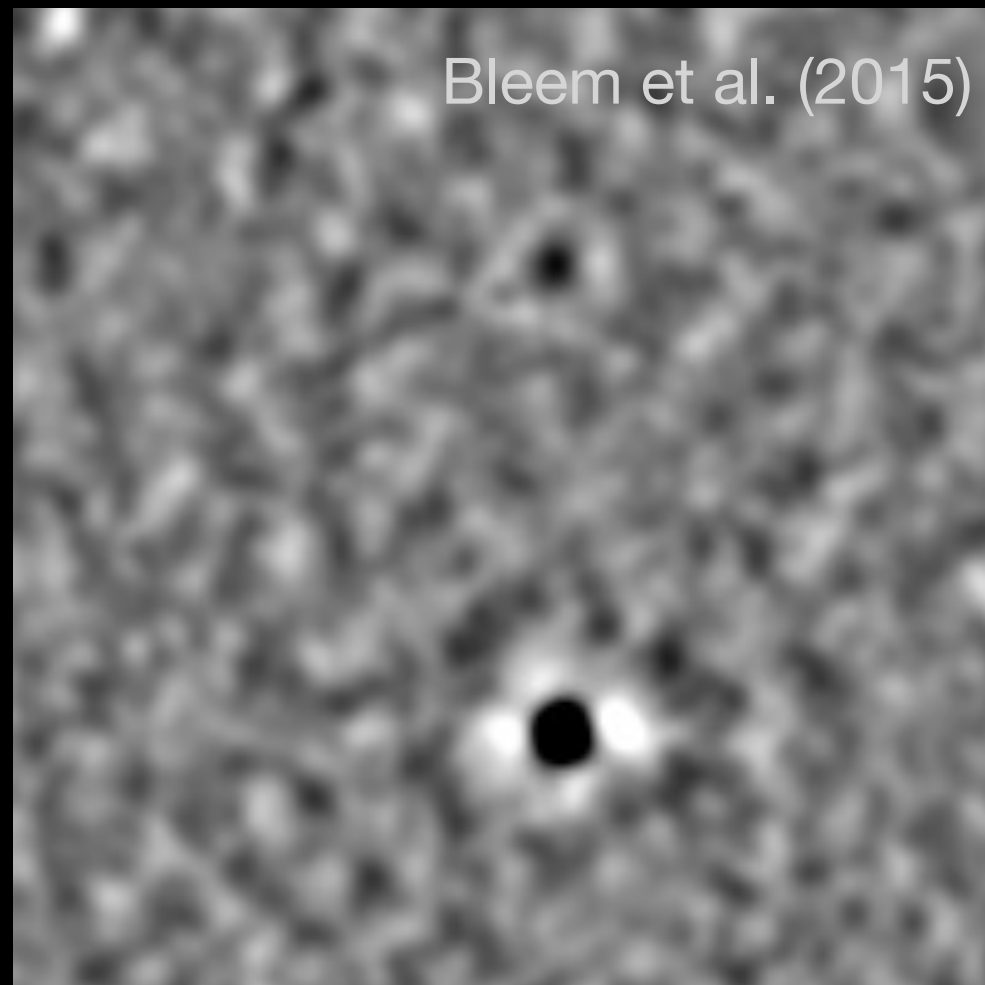
Halo Mass Function

Impact of changing dark energy equation of state parameter by 0.1



Back to reality

“Halo Observable Function”



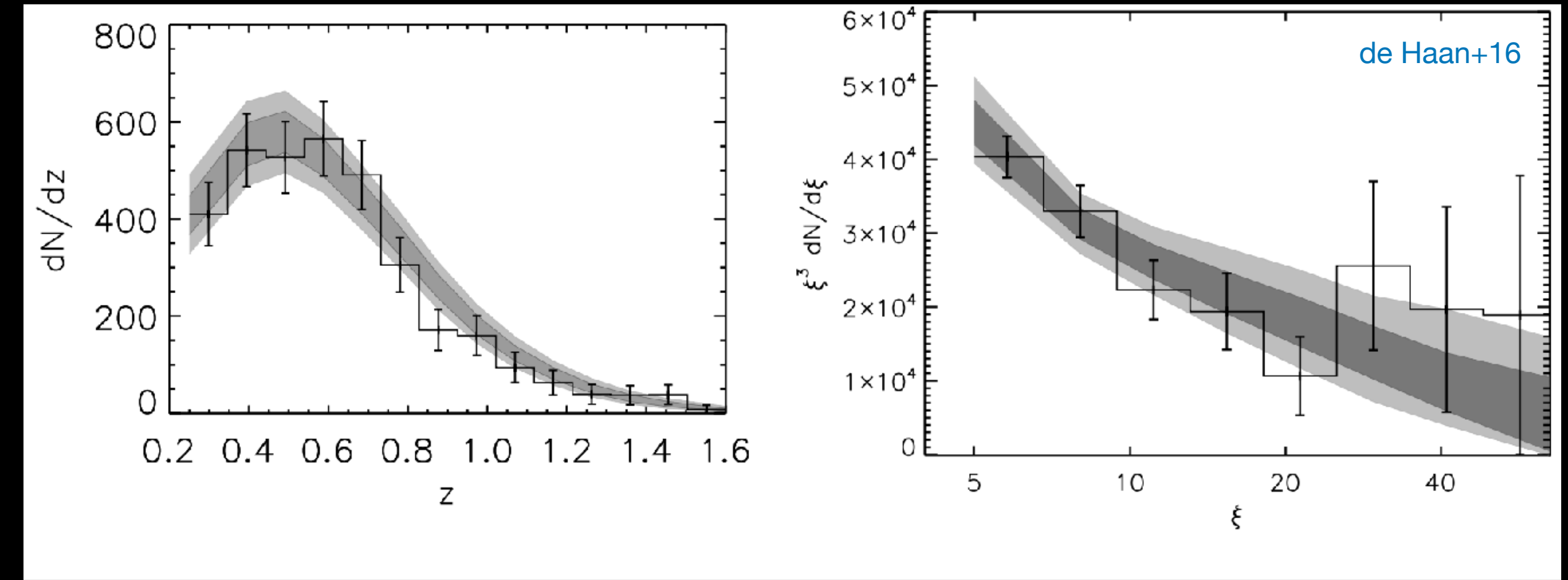
Credit: NASA, ESA, the Hubble Heritage Team (STScI/AURA), J. Blakeslee (NRC Herzberg Astrophysics Program, Dominion Astrophysical Observatory), and H. Ford (JHU) <http://www.spacetelescope.org/images/heic1317a>

Credit: NASA, ESA, and J. Lotz, M. Mountain, A. Koekemoer, and the HFF Team (STScI) <http://www.spacetelescope.org/images/heic1401a/>

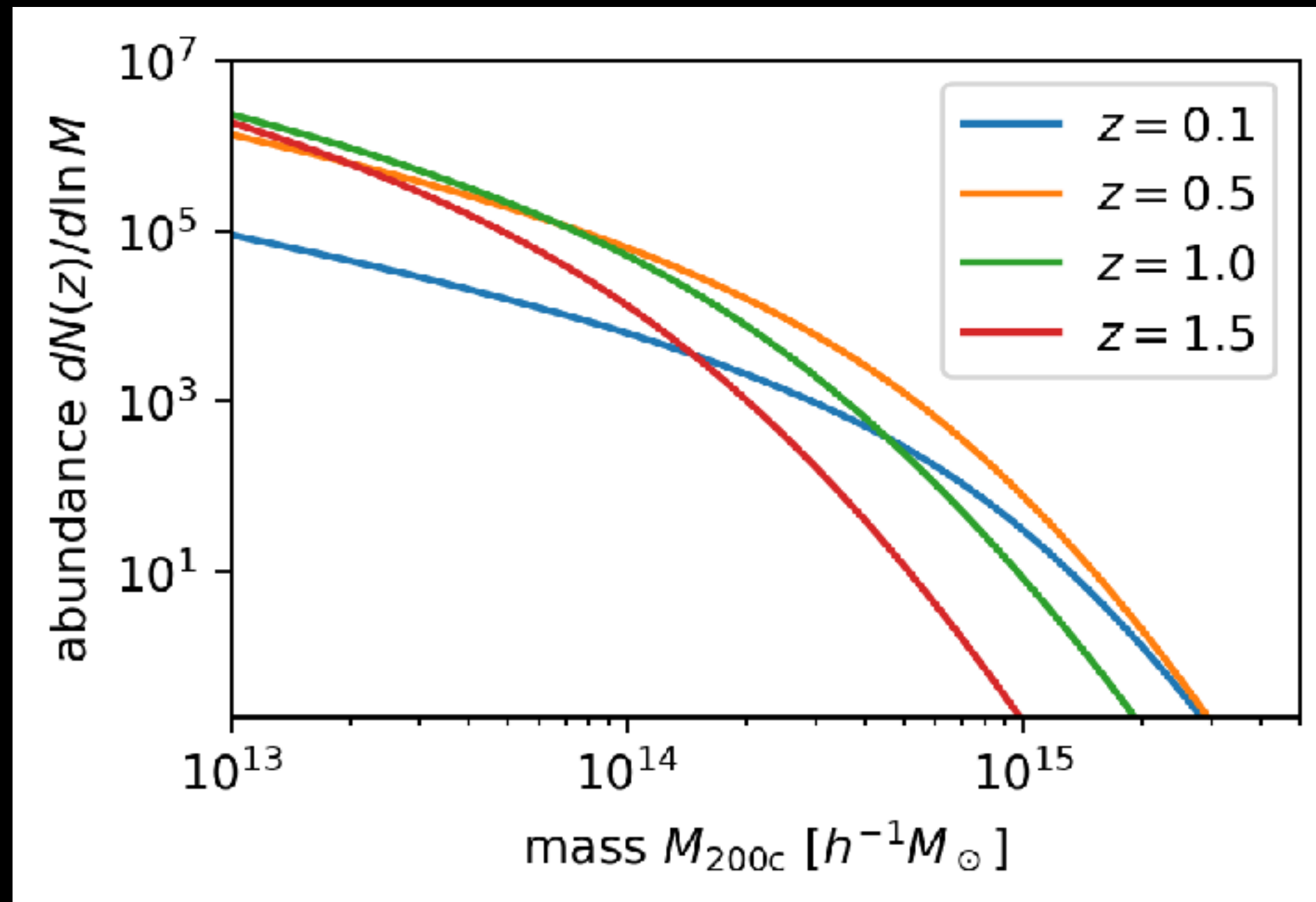
Observable vs. Mass

clusters

halos



halo mass function



observable—mass relation

$$\frac{dN}{d_{\text{obs}}} = \int dM P(\text{obs} | M) \frac{dN}{dM}$$

halo mass function

cluster cosmology = cluster selection + mass calibration

SPT Clusters with DES and HST Weak Lensing. I. Cluster Lensing and Bayesian Population Modeling of Multi-Wavelength Cluster Datasets

Bocquet, Grandis, Bleem, Klein, Mohr, DES, SPT
(arXiv:2310:12213 — Phys. Rev. D 2024, 110, 083509)

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Bocquet, Grandis, Bleem, Klein, Mohr, Schrabback, SPT, DES
(arXiv:2310:12213 — Phys. Rev. D 2024, 110, 083510)



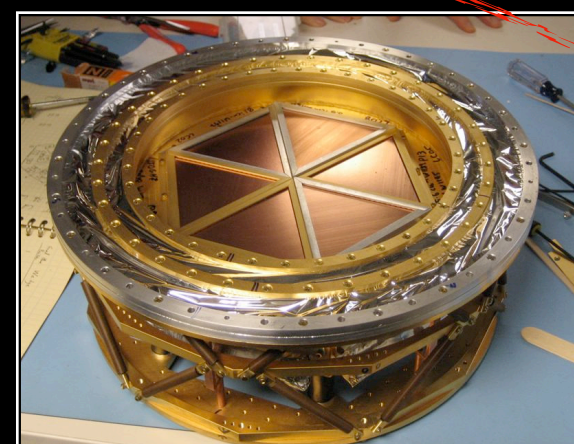
The South Pole Telescope (SPT)

10-meter sub-mm quality wavelength telescope

90, 150, 220 GHz and
1.6, 1.2, 1.0 arcmin resolution

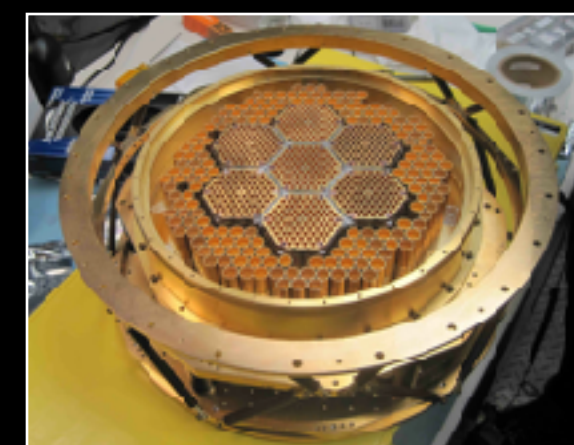
2007: SPT-SZ

960 detectors
90, 150, 220 GHz



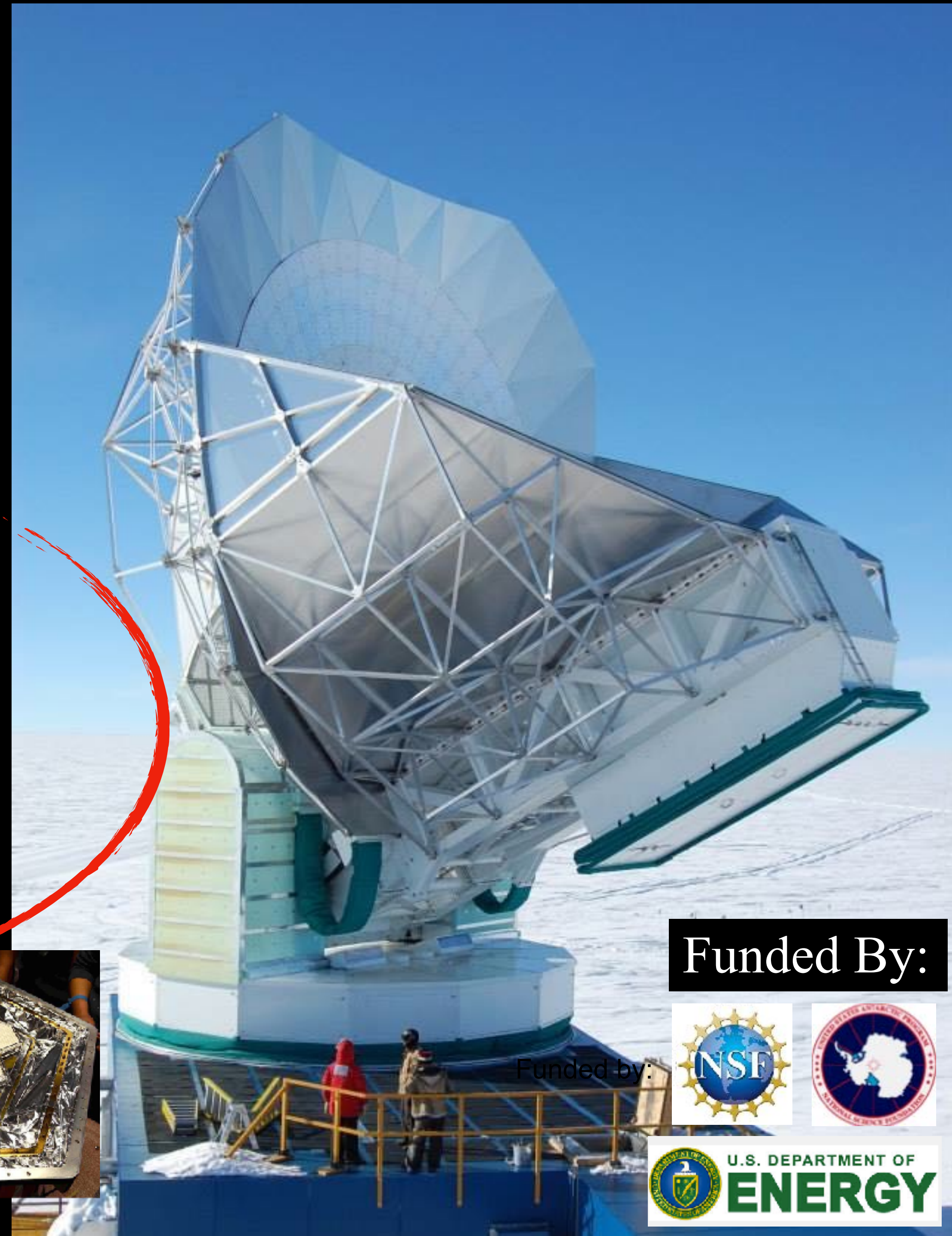
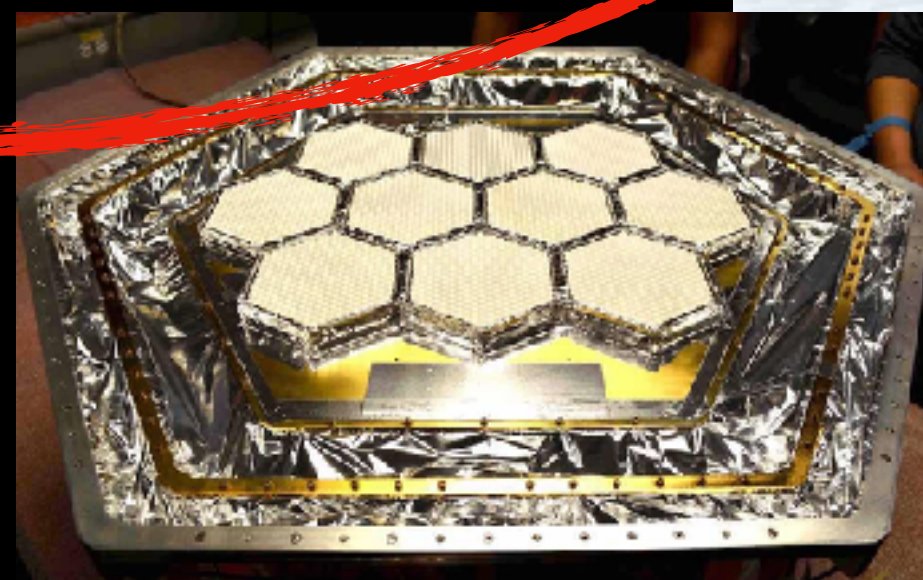
2012: SPTpol

1600 detectors
90, 150 GHz
+Polarization



2017: SPT-3G

~15,200 detectors
90, 150, 220 GHz
+Polarization

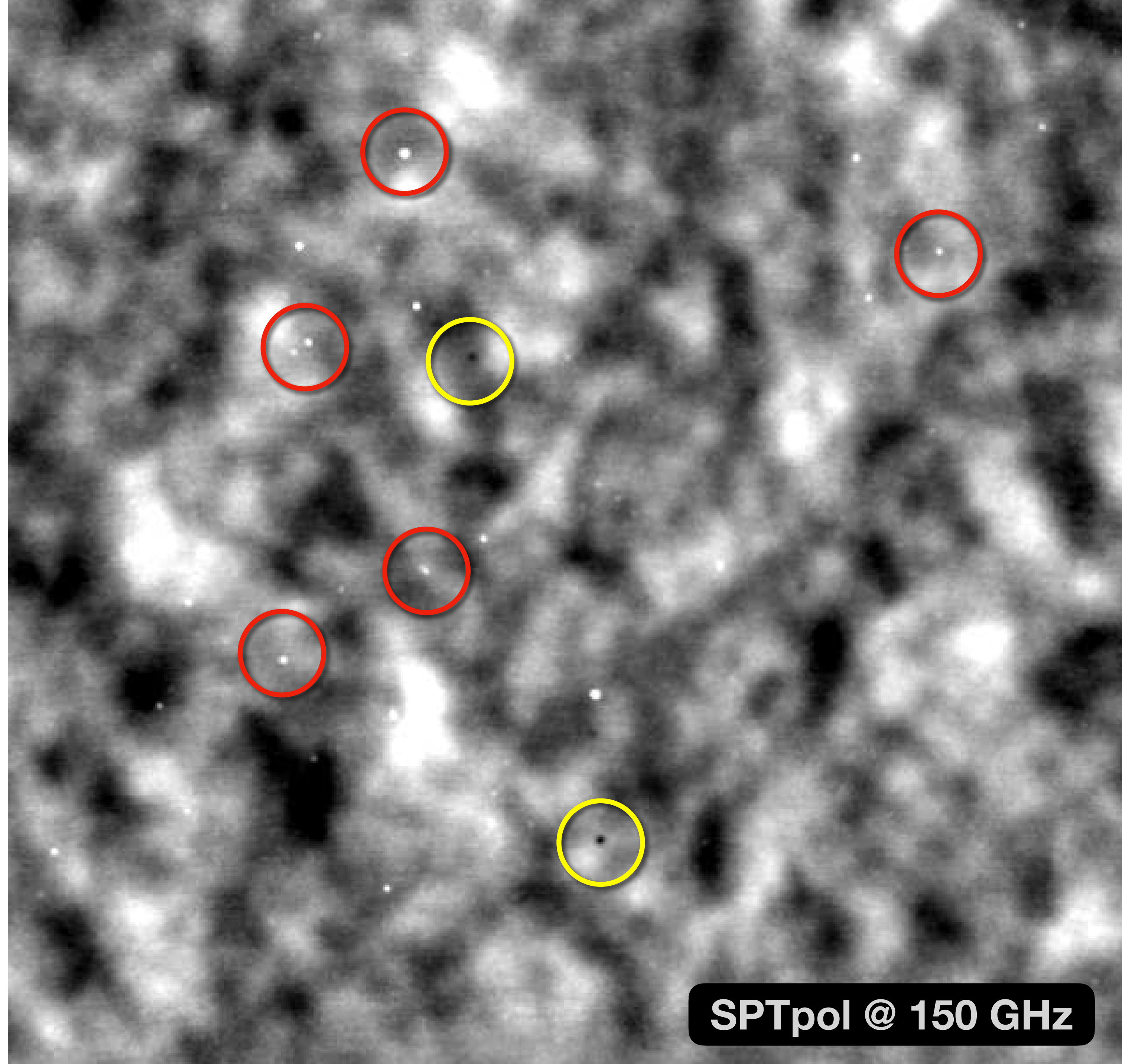
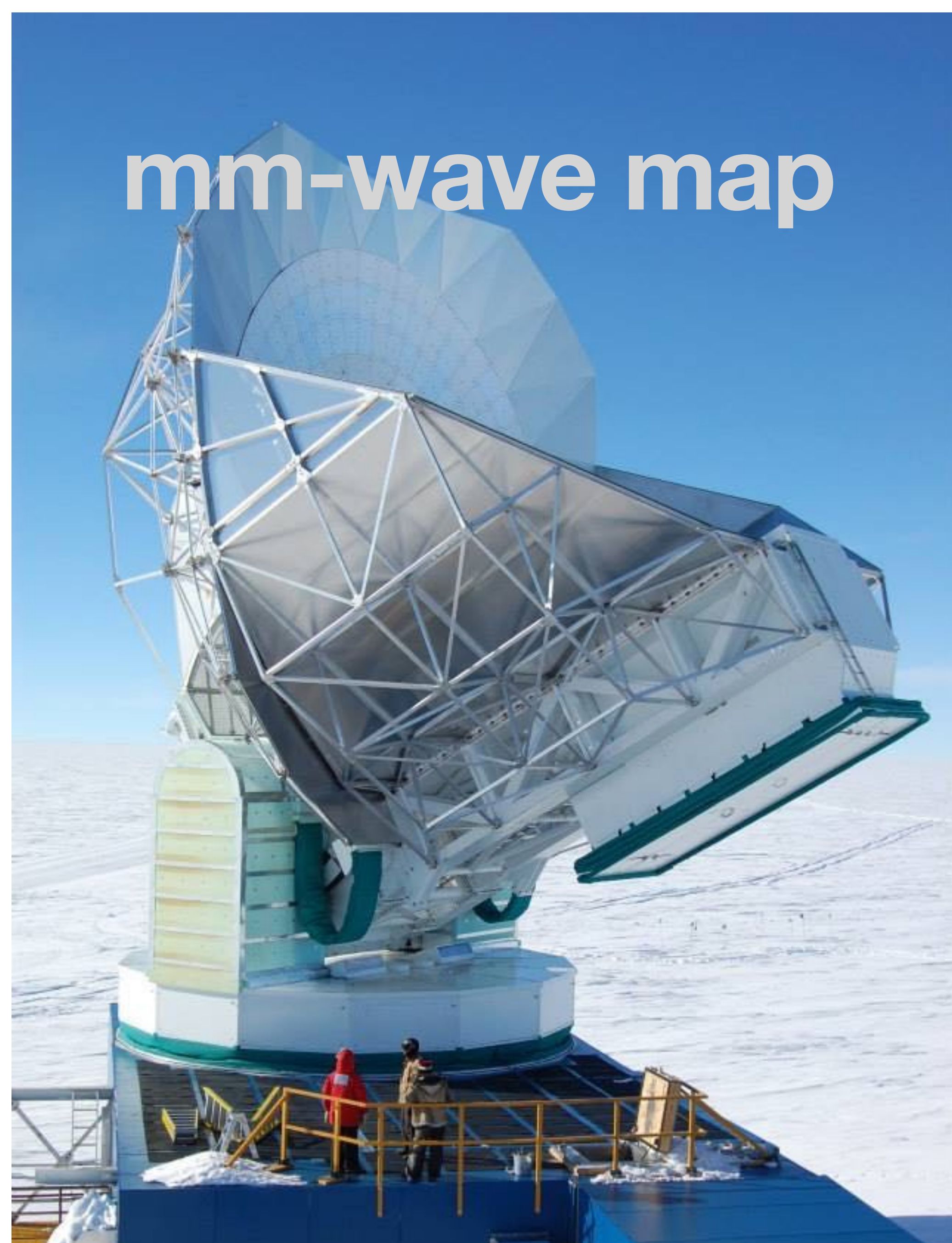


Funded By:



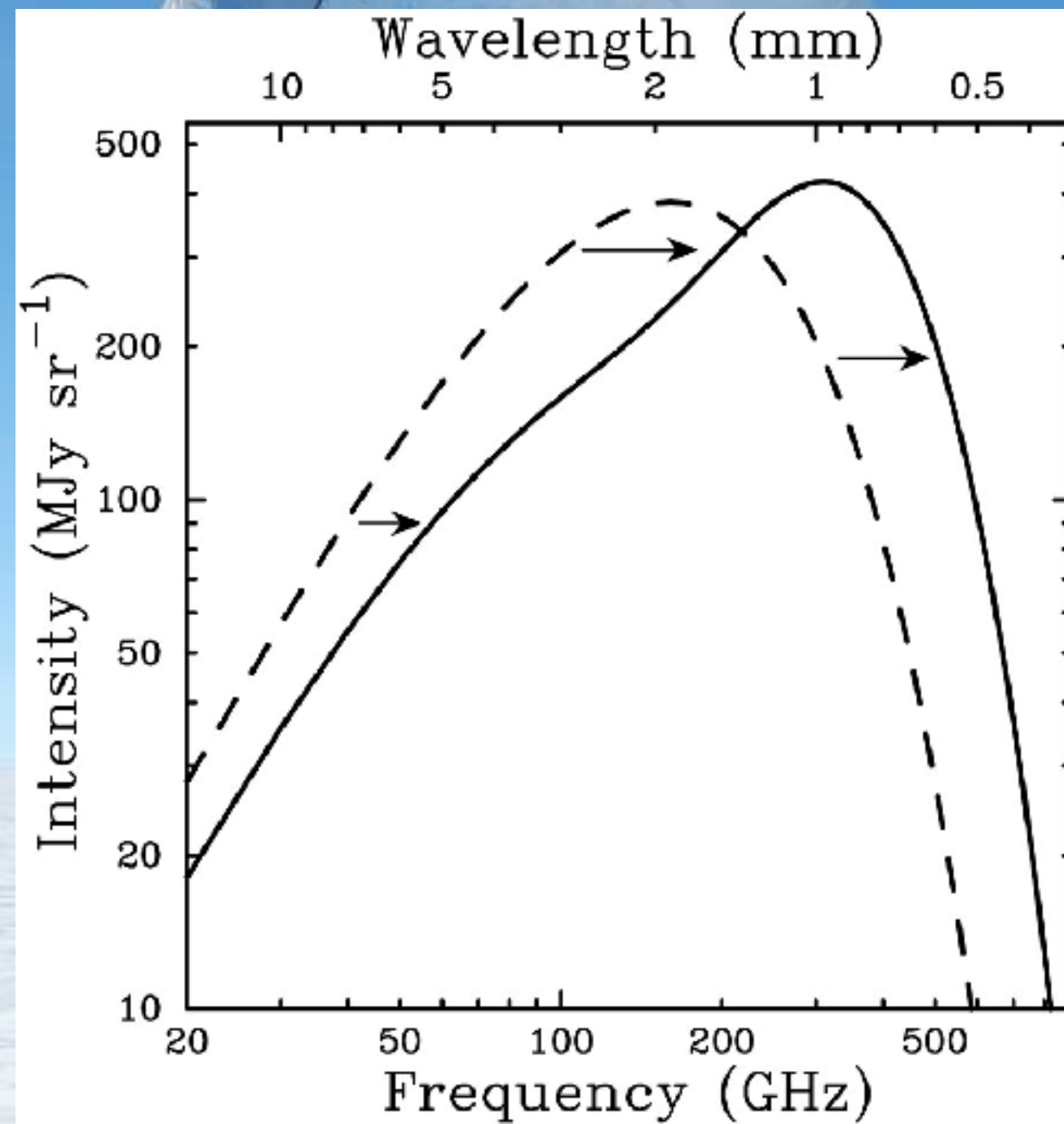
U.S. DEPARTMENT OF
ENERGY

mm-wave map

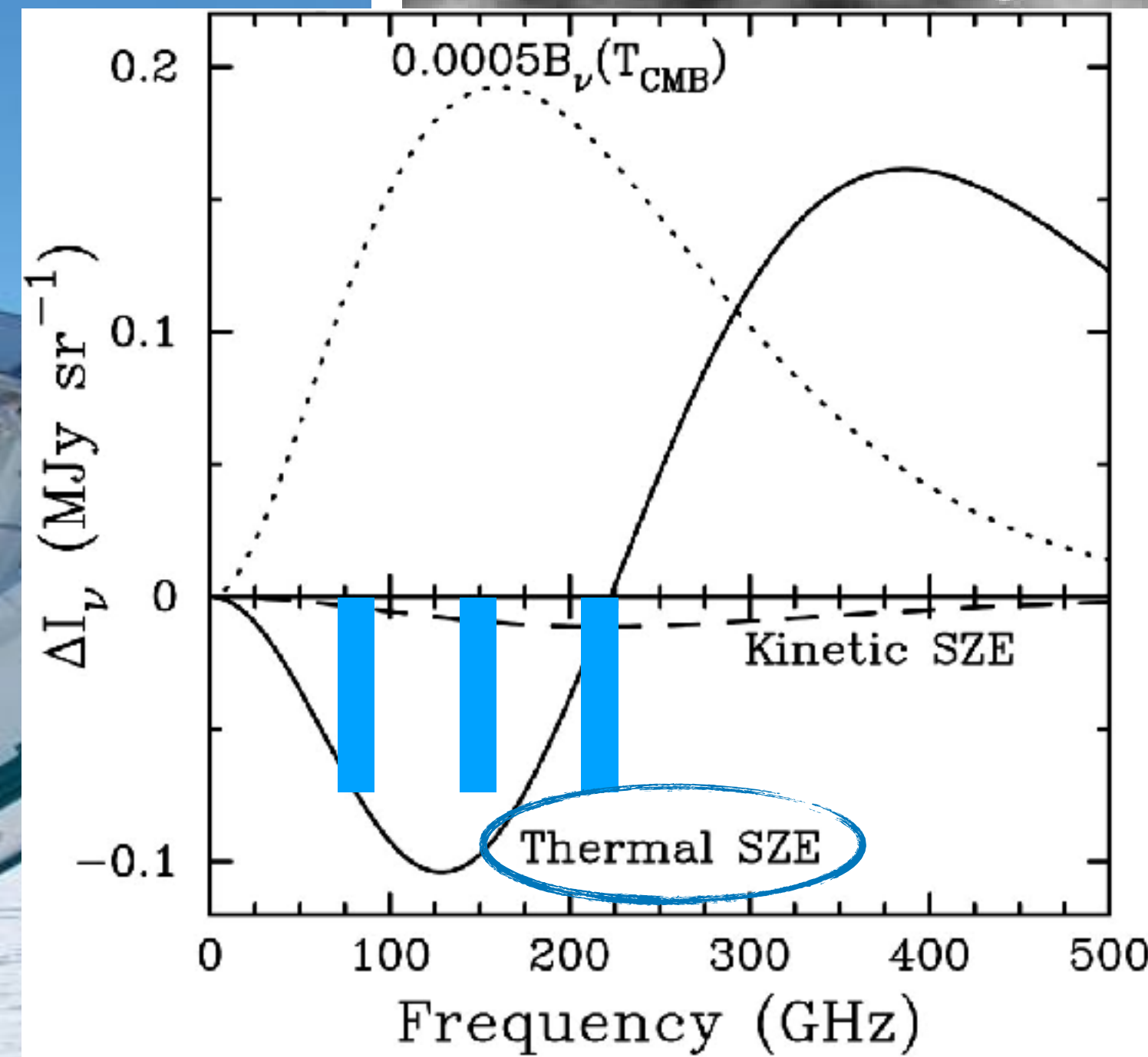


Find clusters

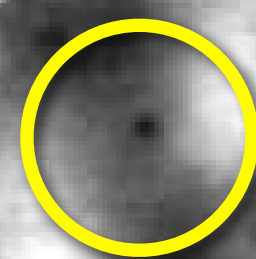
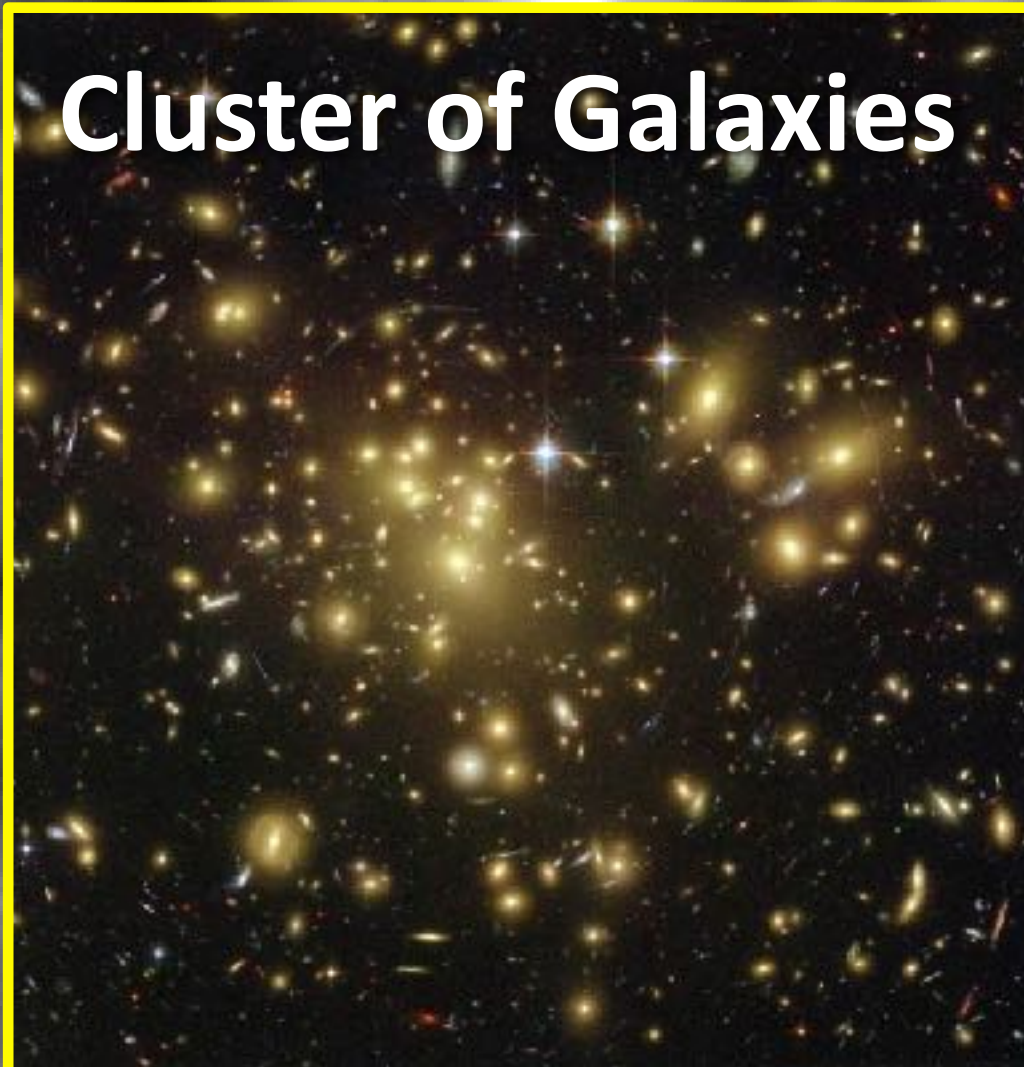
Sunyaev-Zel'dovich (SZ) Effect



CMB spectrum

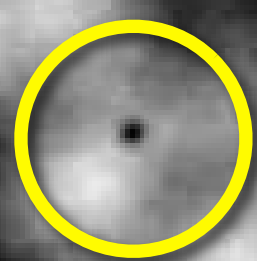


SZ spectrum



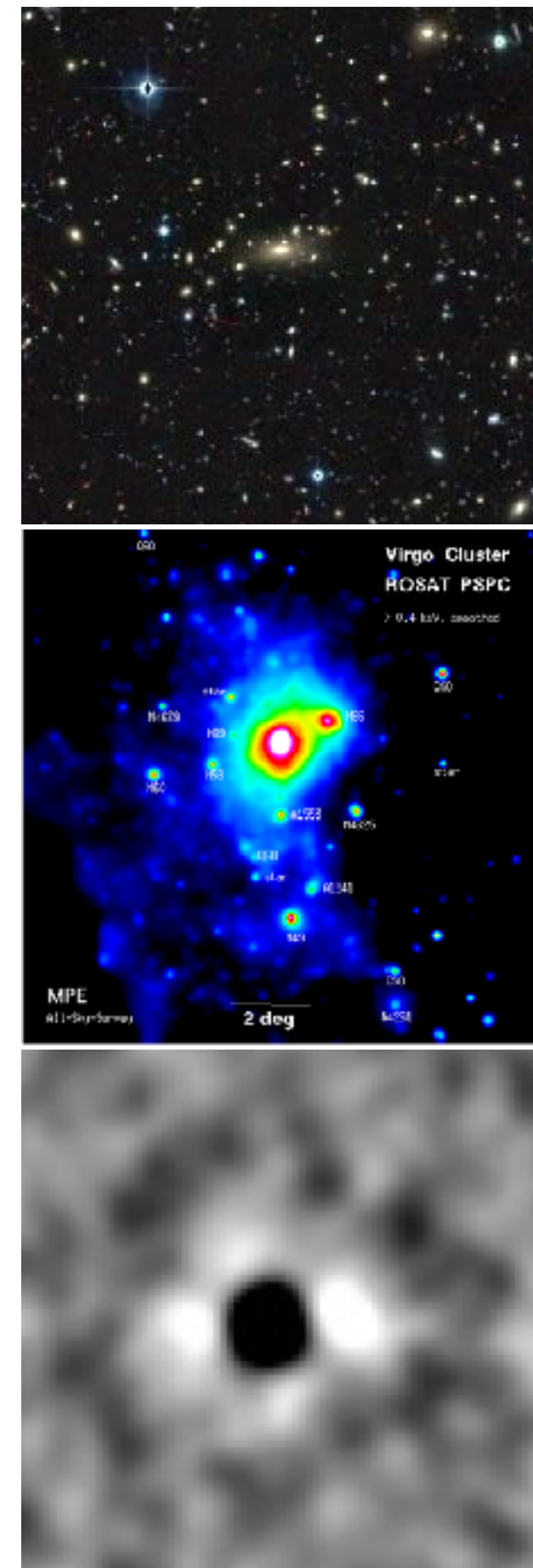
Clean and well-understood selection of cluster candidates

Out to highest redshifts where clusters exist!

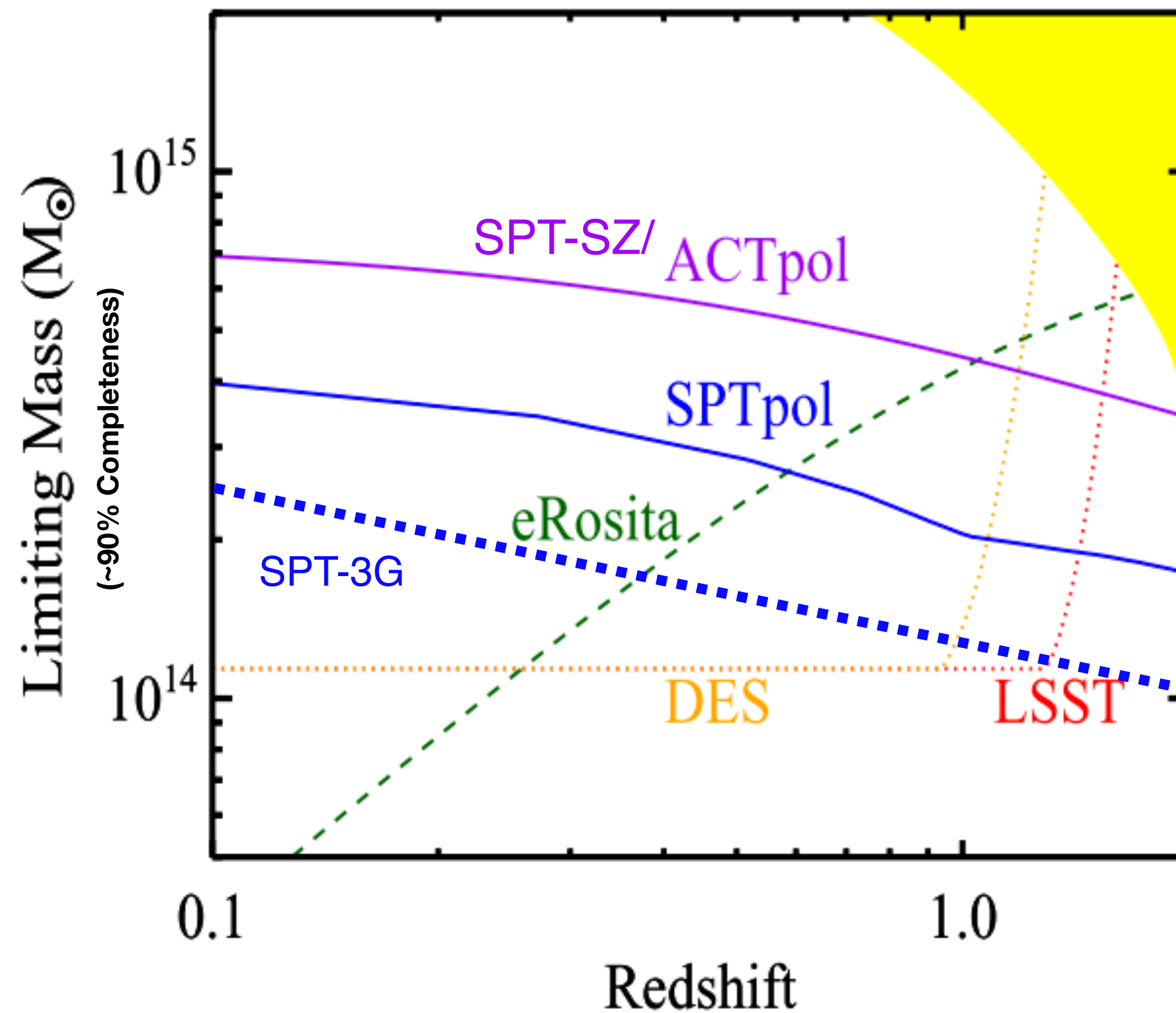


Why use SZ-selected clusters?

Three approaches: **X-ray**, **Optical**, **SZ**

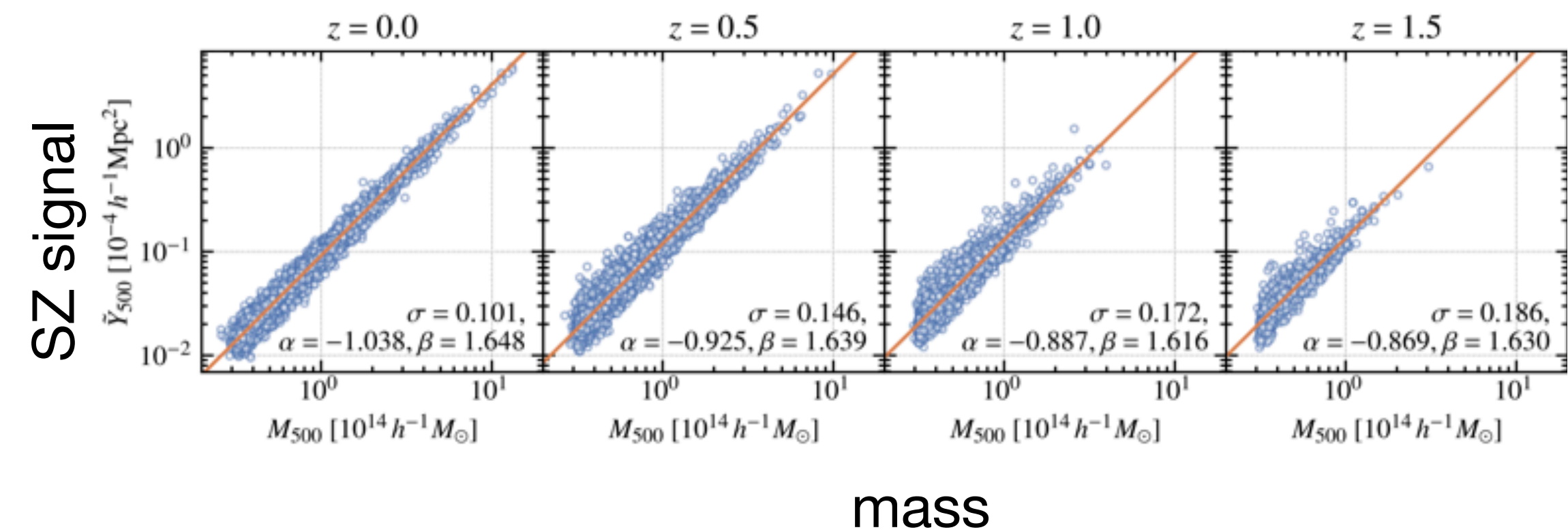


Weinberg et al. 2013

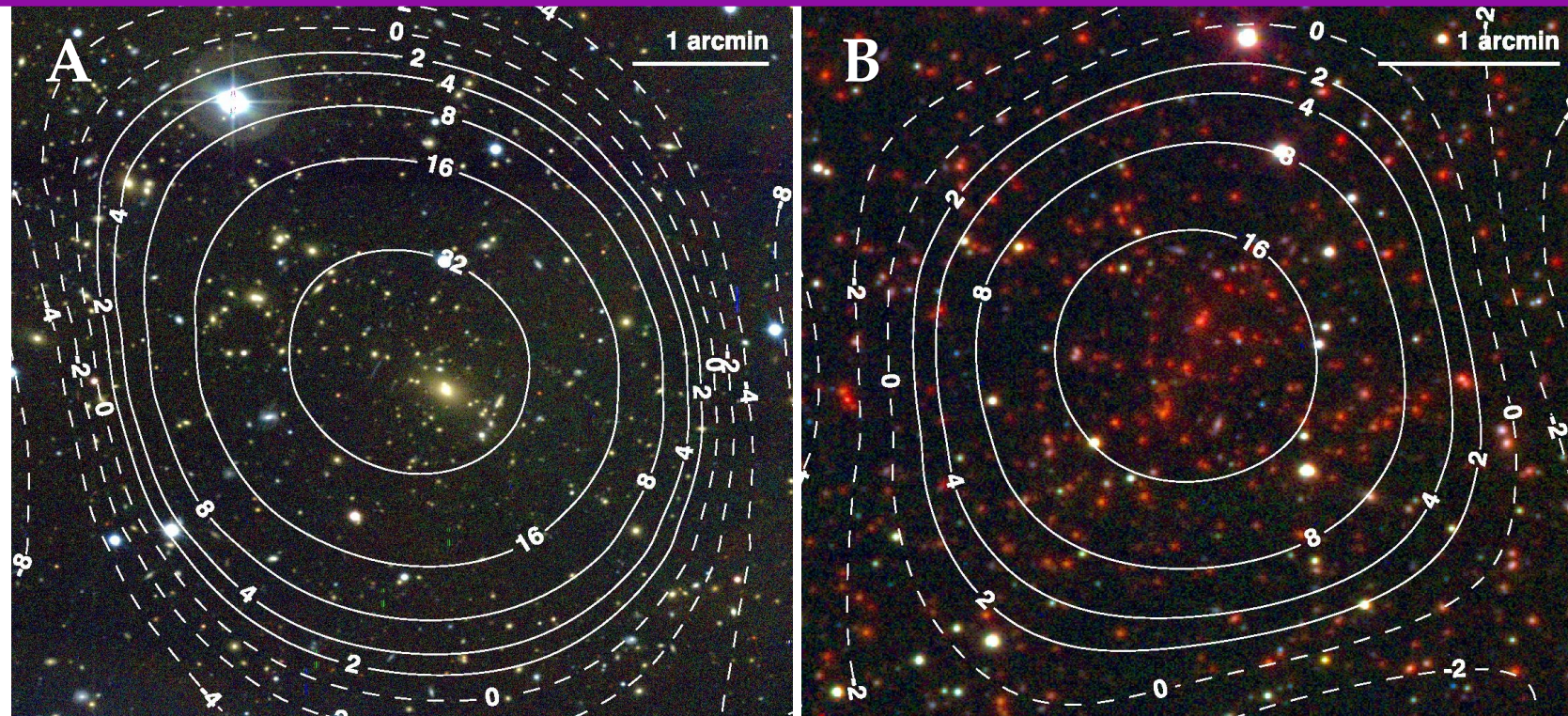


SZ: clean, well-understood selection
Complementarity with other methods

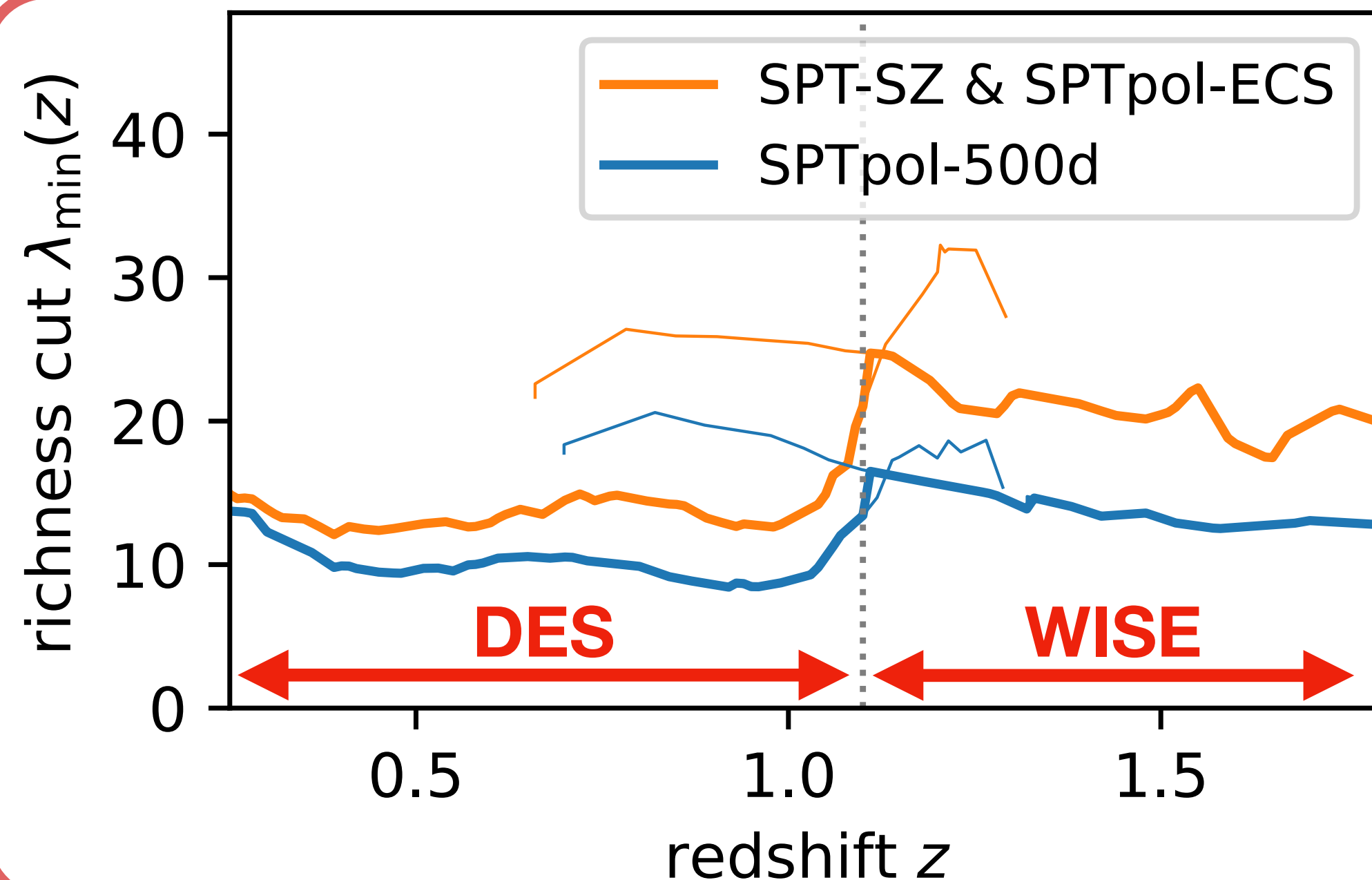
Kéruzoré et al. 2023



How to confirm SZ candidates?



Measure richness (\cong number of cluster member galaxies) and redshift



Get rid of chance associations (with SPT noise fluctuation)

Calibrate probability of chance association by measuring (λ, z) at random locations

Establish $\lambda_{\min}(z)$ to achieve target purity ($> 98\%$)

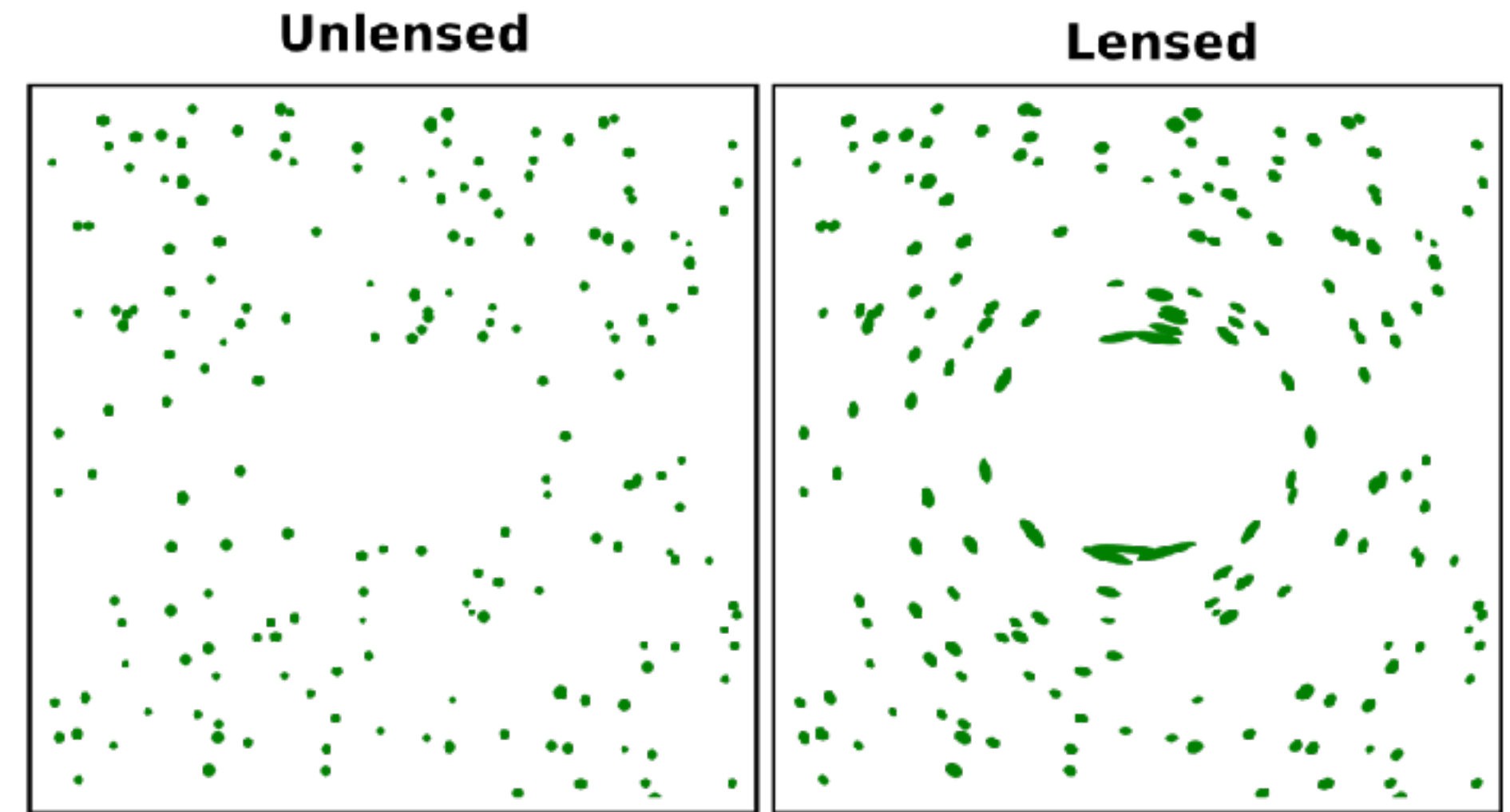
(Klein+18,24; Bleem+24)

Mass Calibration I

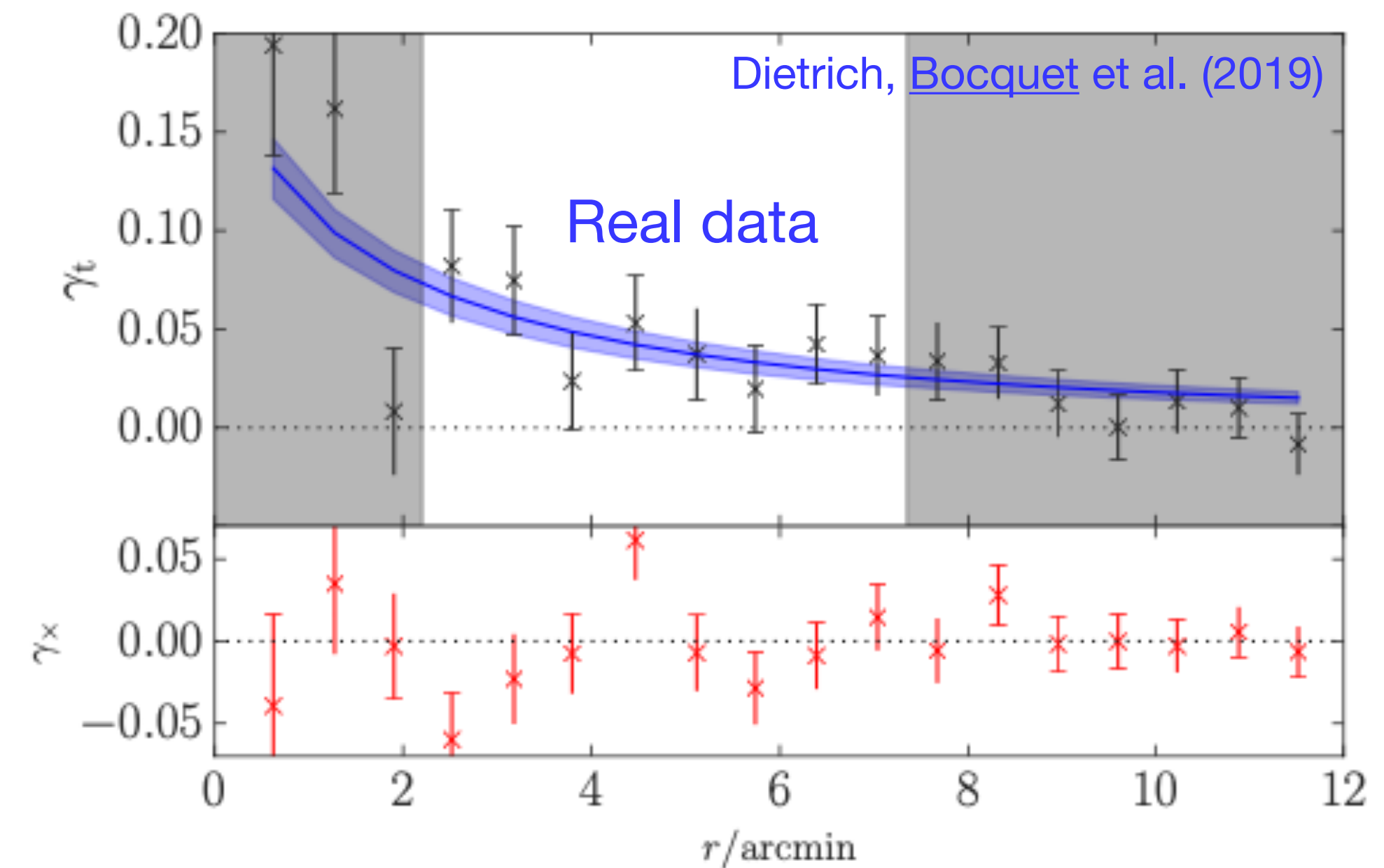
How do the observables relate to halo mass?

- We *could* use predictions from first principles (e.g., hydrostatic equilibrium) or numerical simulations
 - Systematically limited by uncertain astrophysics
- Weak-lensing-to-mass relation is known within few percents

Idealized (exaggerated) situation



By TallJimbo - Own work, CC BY-SA 3.0, <https://commons.wikimedia.org/w/index.php?curid=4150002>



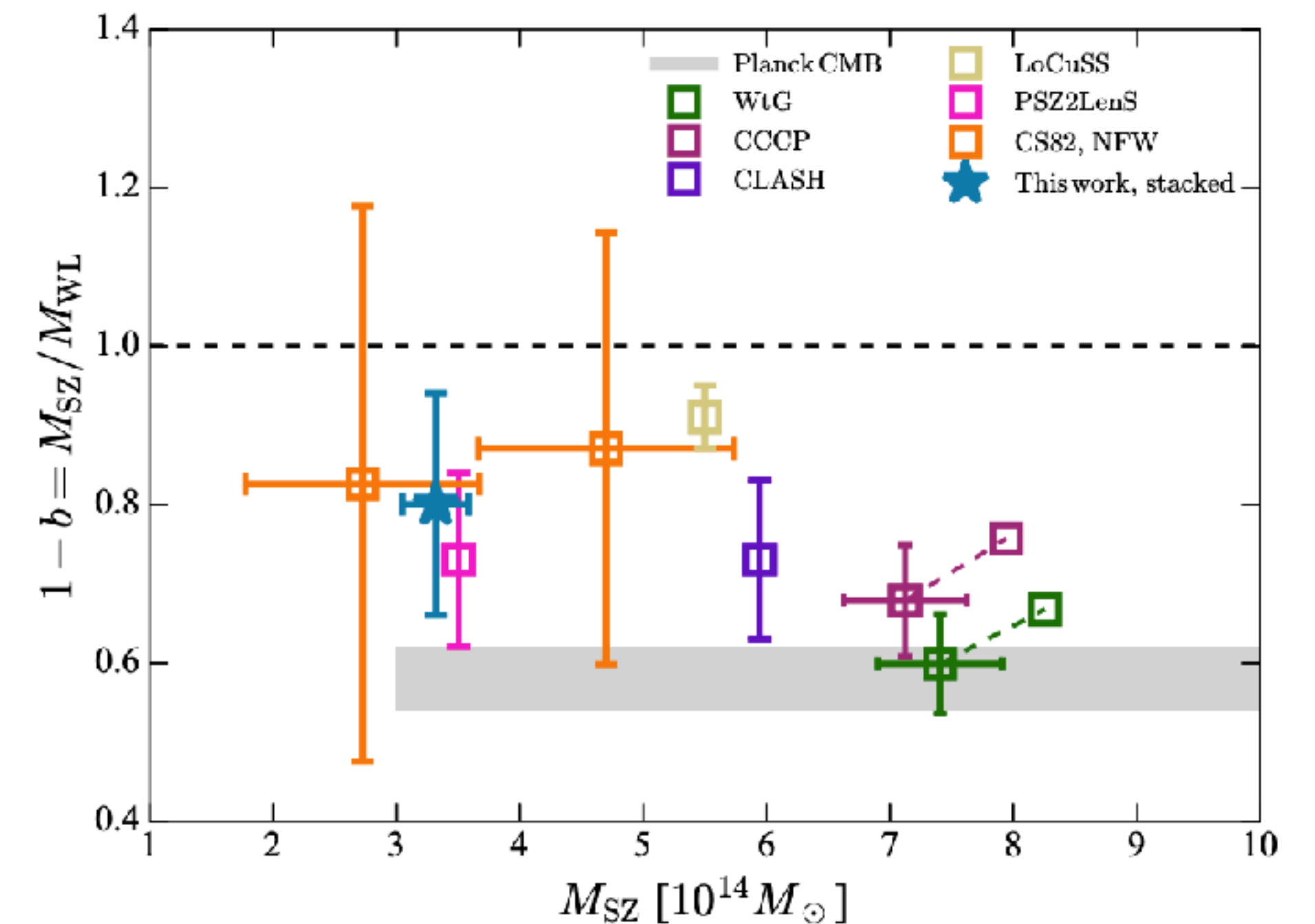
(b) Tangential shear profile of SPT-CL J0254–5857.

Mass Calibration II. Weak Lensing

Robust observable–mass relations

- We *could* use predictions from first principles (e.g., hydrostatic equilibrium) or numerical simulations
 - Systematically limited by uncertain astrophysics
- Weak-lensing-to-mass relation is known within few percents
 - Used to demonstrate that **hydrostatic mass \neq halo mass**
 - ▶ **With lensing** measurements of sample clusters, **we empirically calibrate the observable–mass relations**

Medezinski+ 18



The Dark Energy Survey

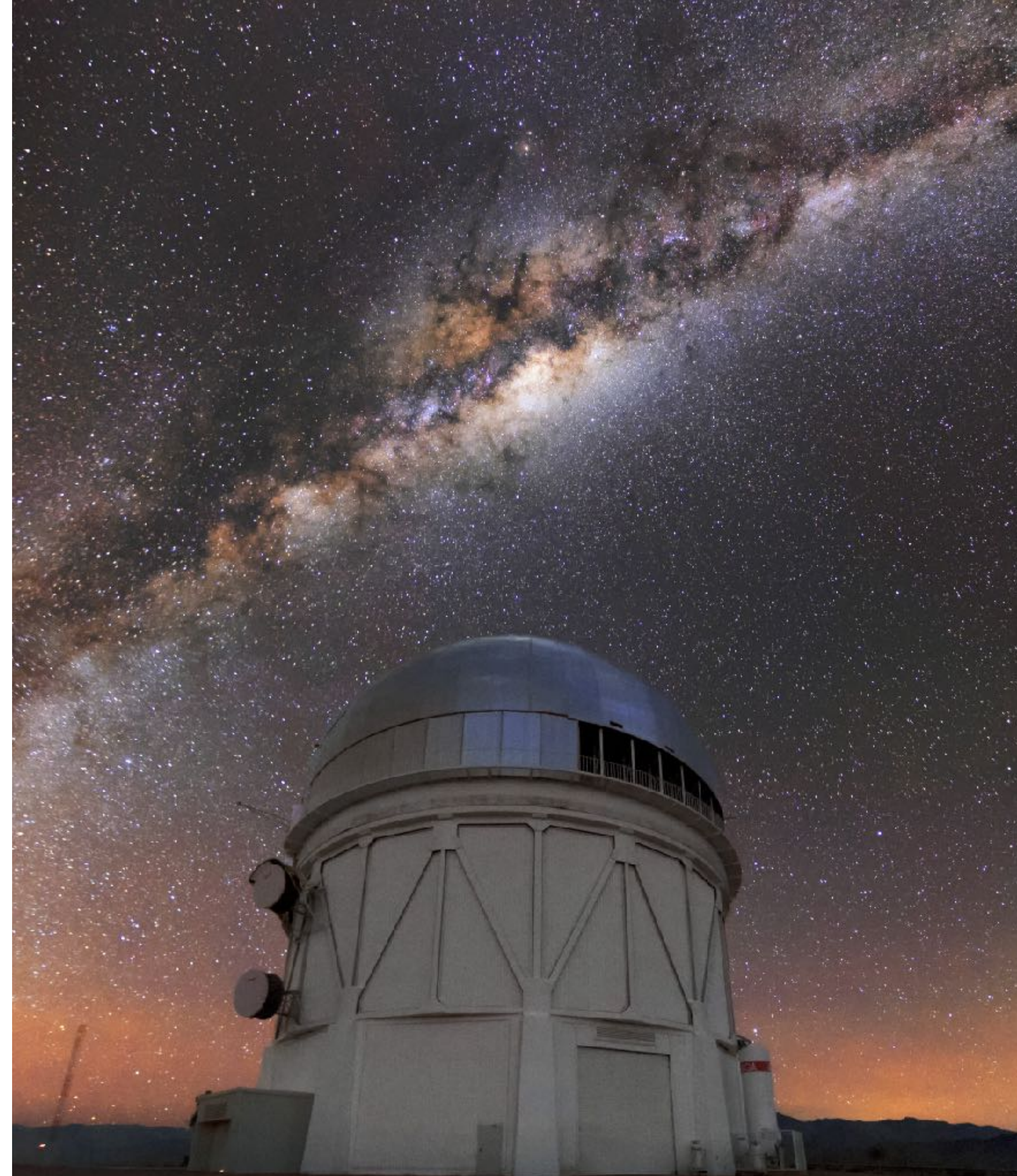
5000 deg² galaxies & weak lensing

Catalog of SPT-selected cluster candidates needs

- Confirmation
- Cluster redshifts
- Weak-lensing (mass) measurement

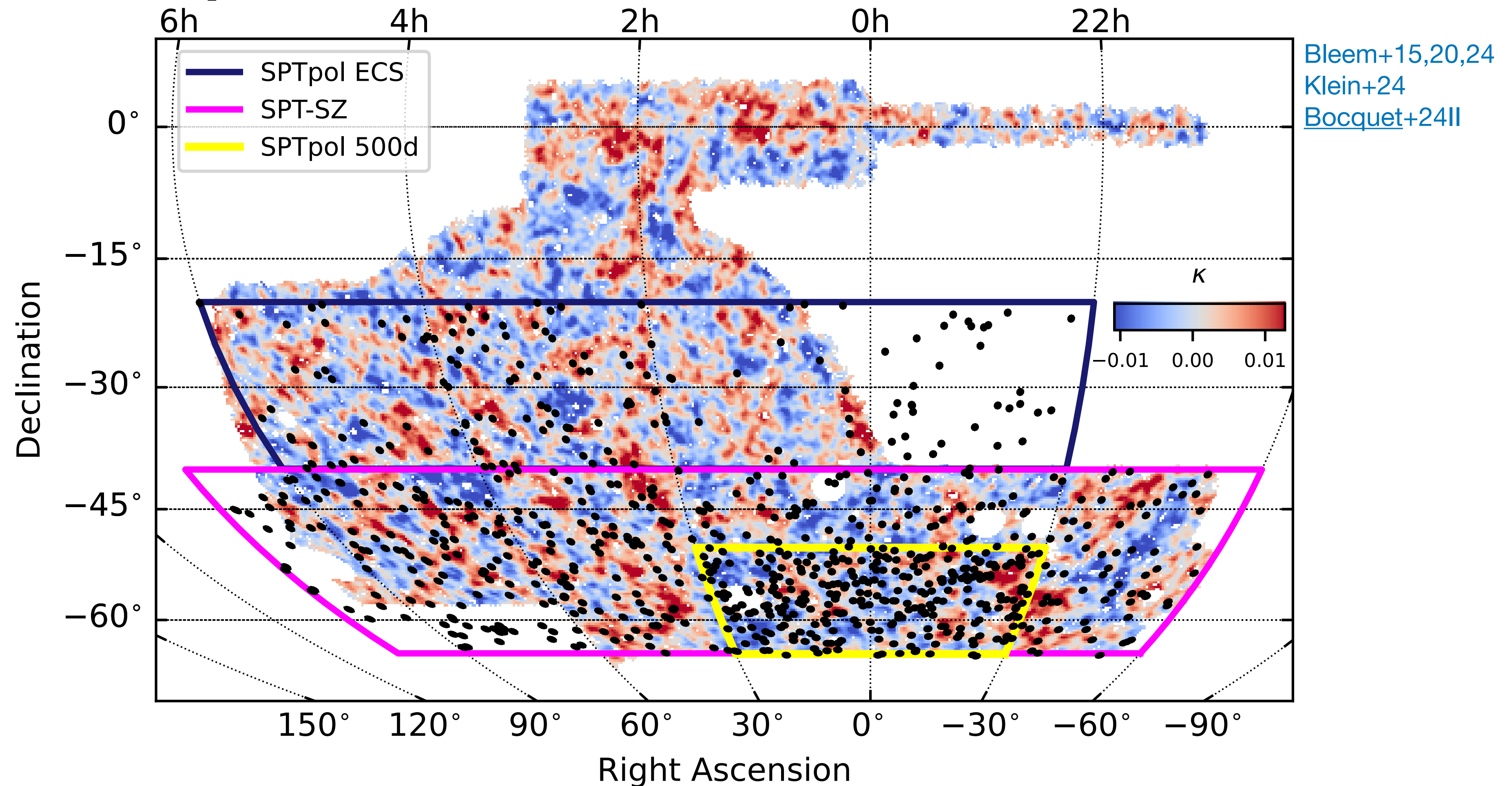
all of which DES was designed for

(here we use DES Year 3 data = Y3)



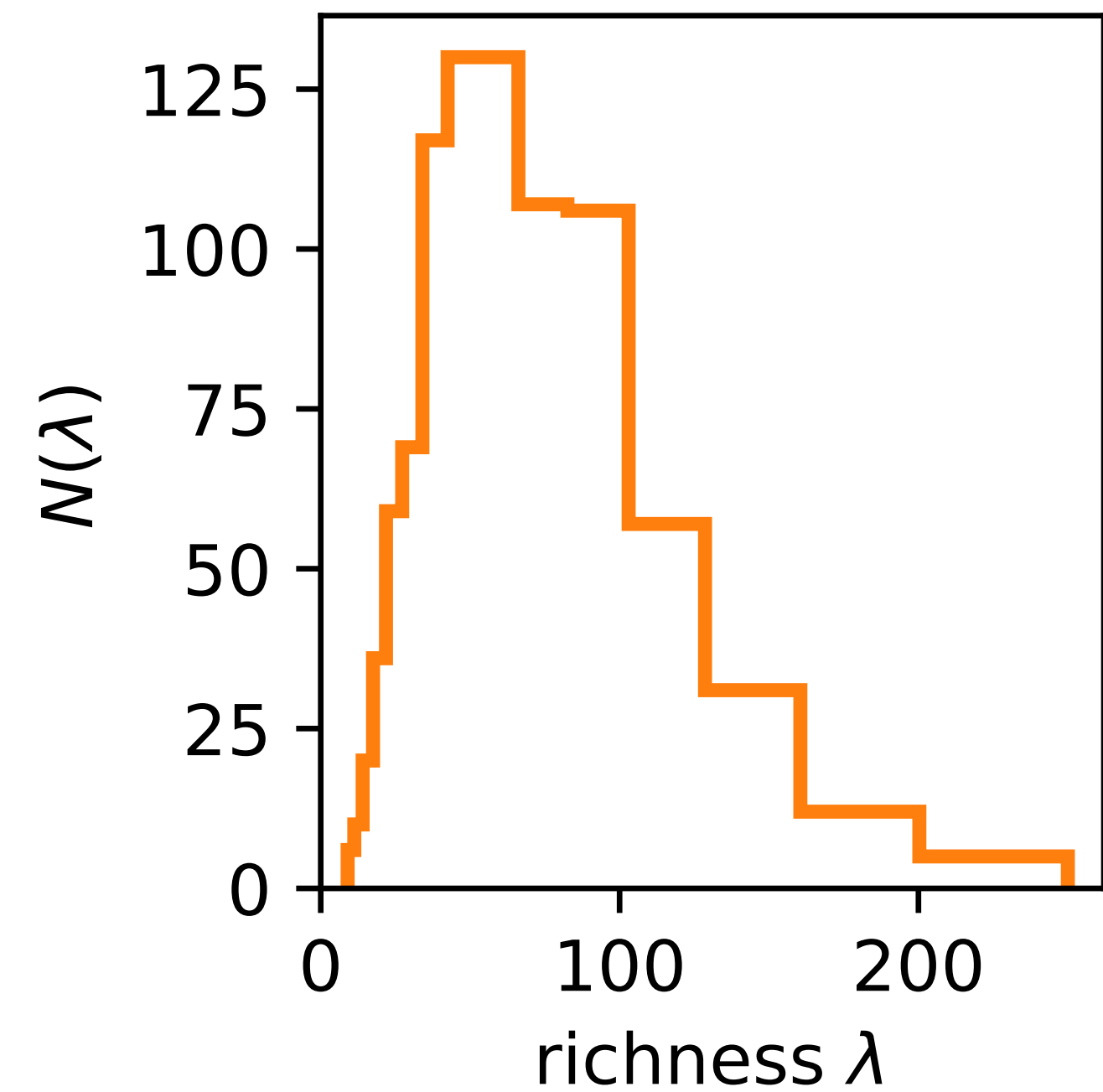
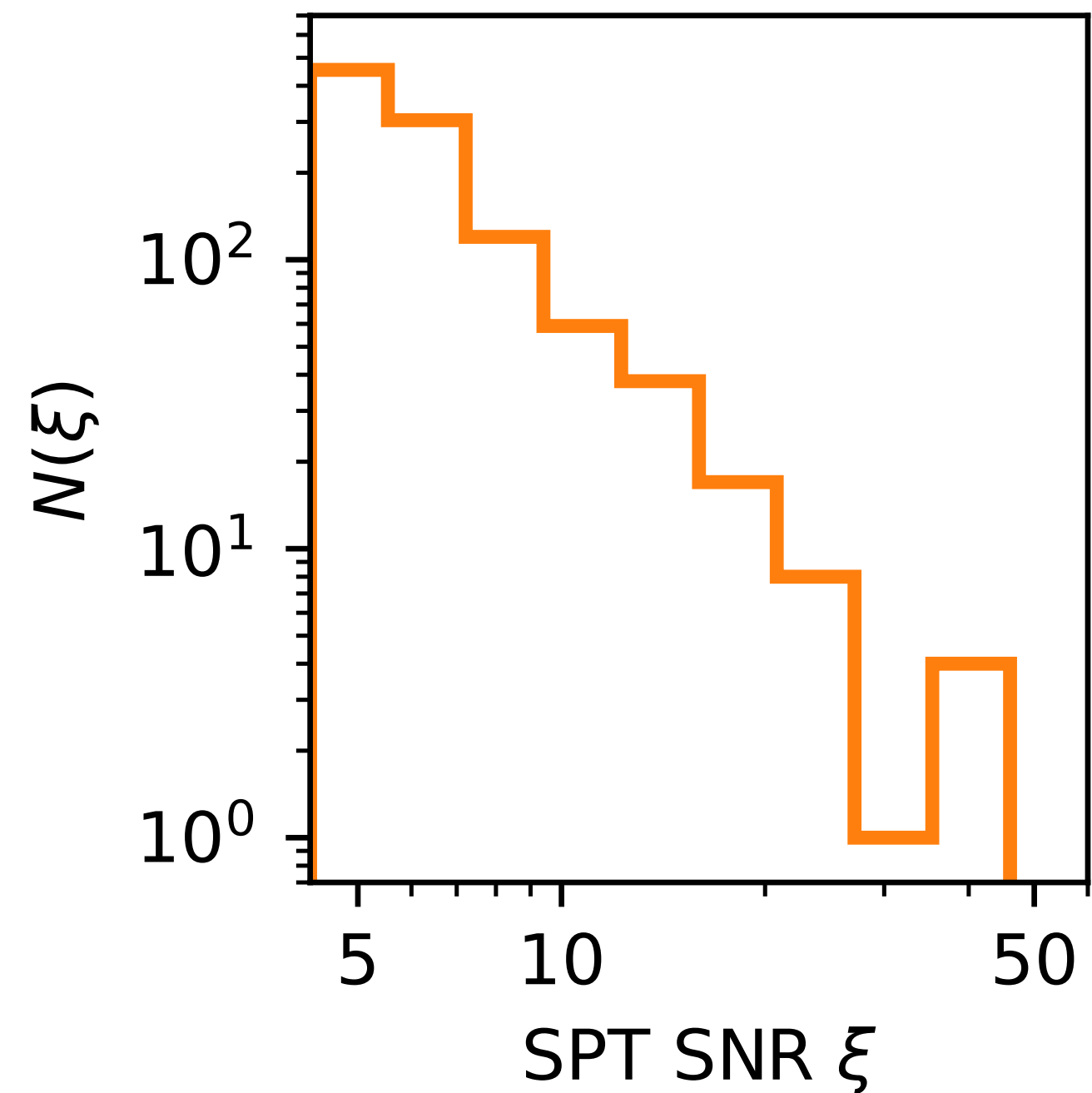
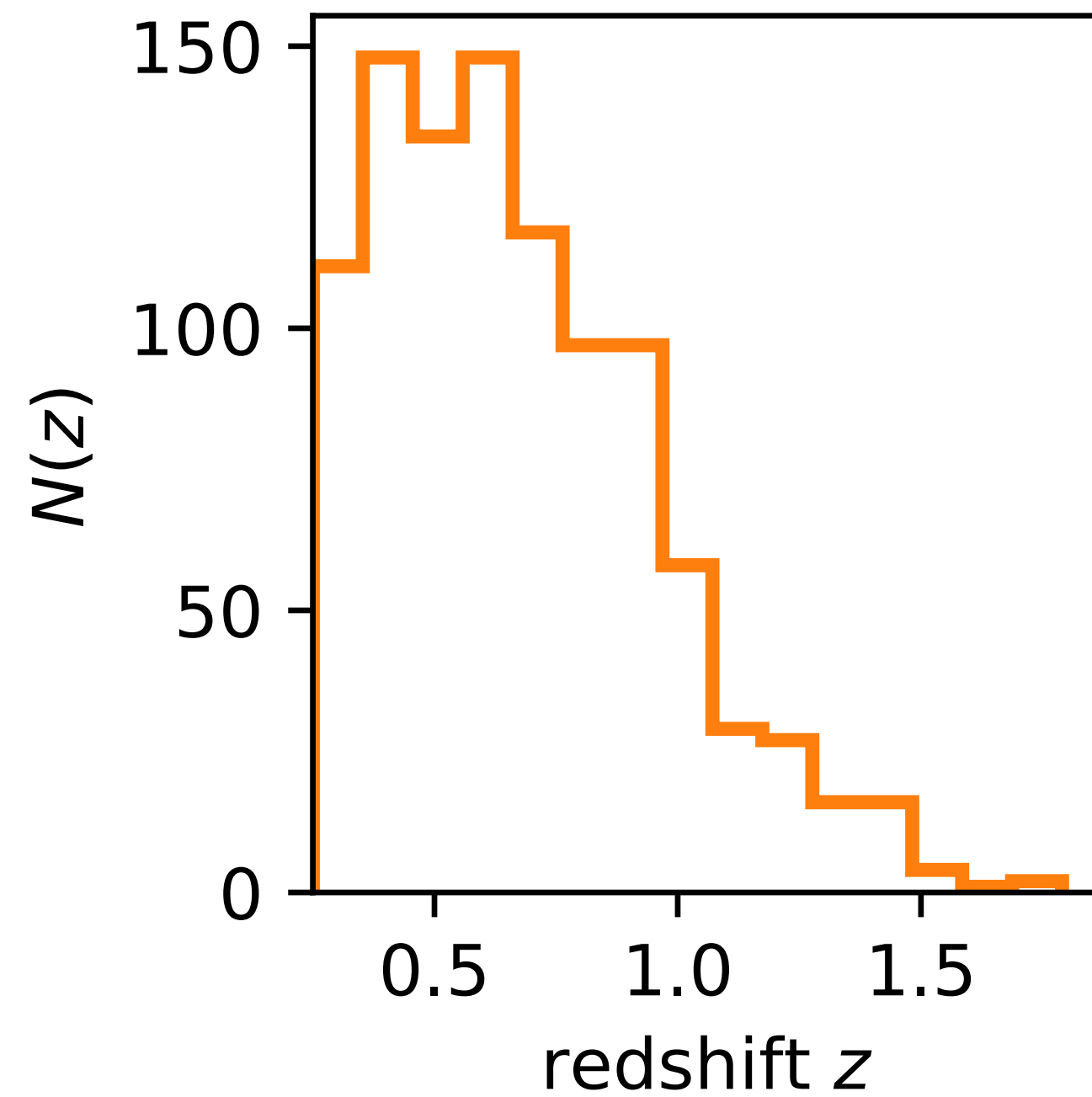
SPT Clusters and the Dark Energy Survey

3,600 deg² overlap



SPT(SZ+pol) Cluster Sample

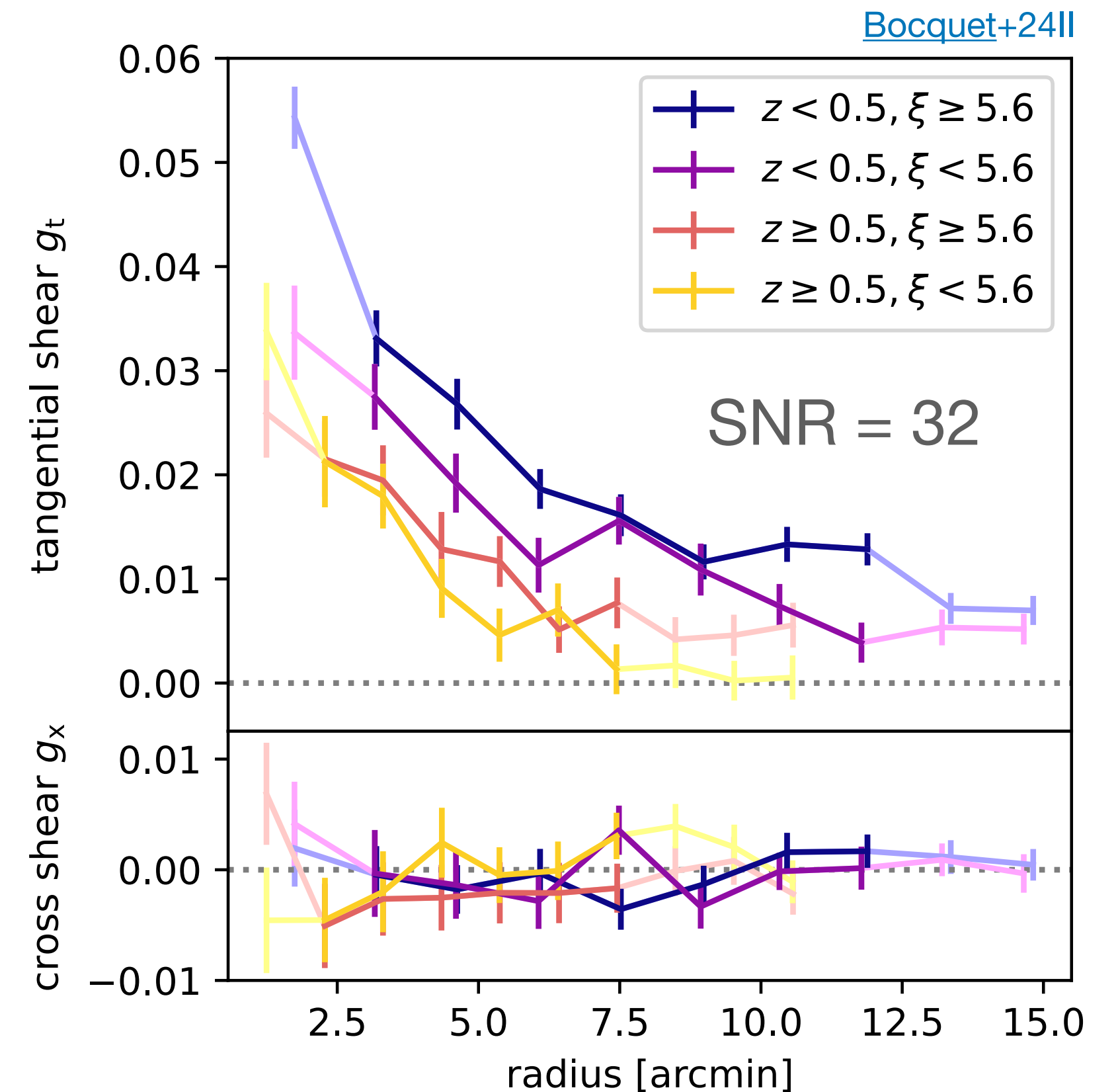
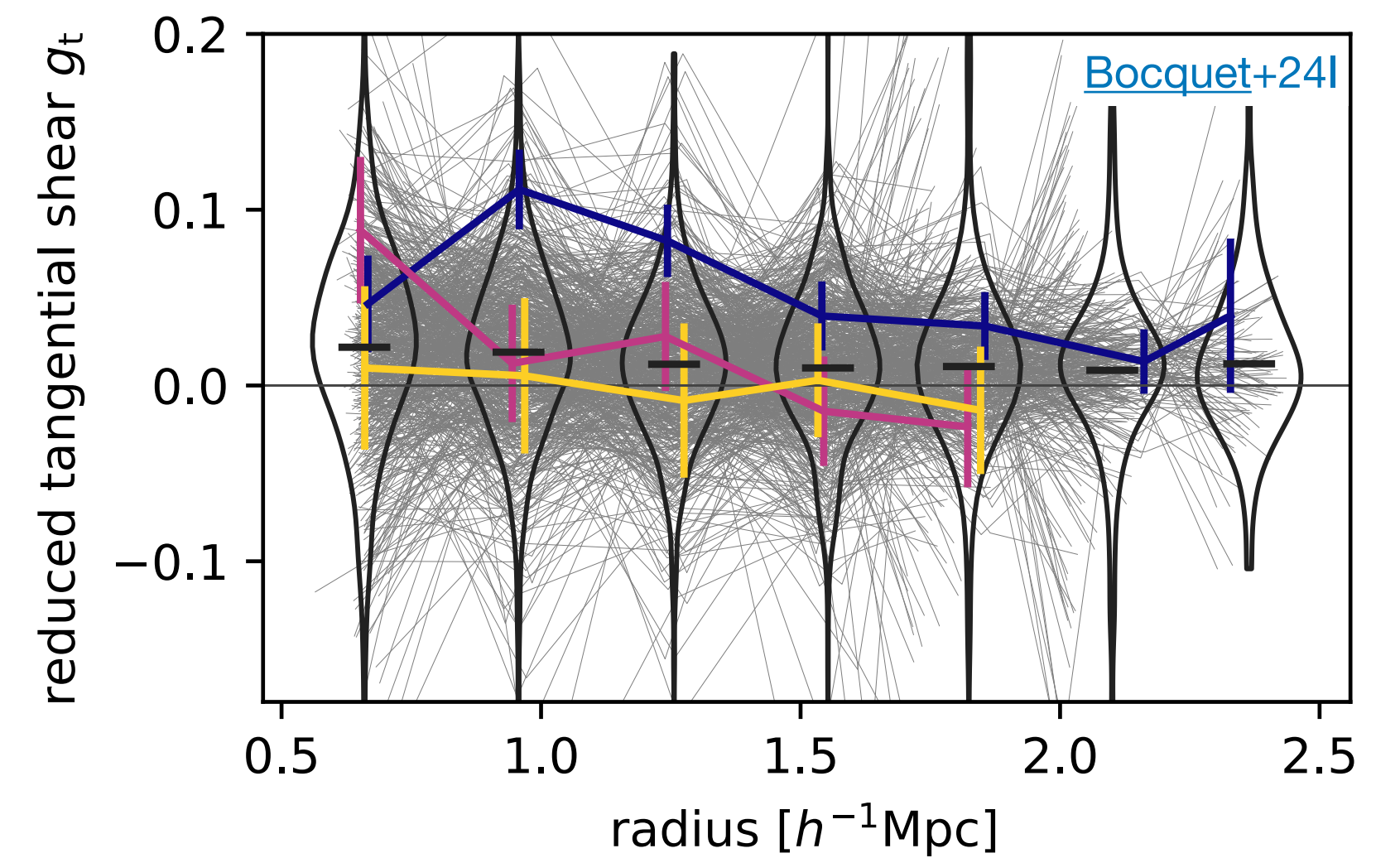
1,005 confirmed clusters above $z > 0.25$ over 5,200 deg²



Cluster lensing analysis

Shear profiles

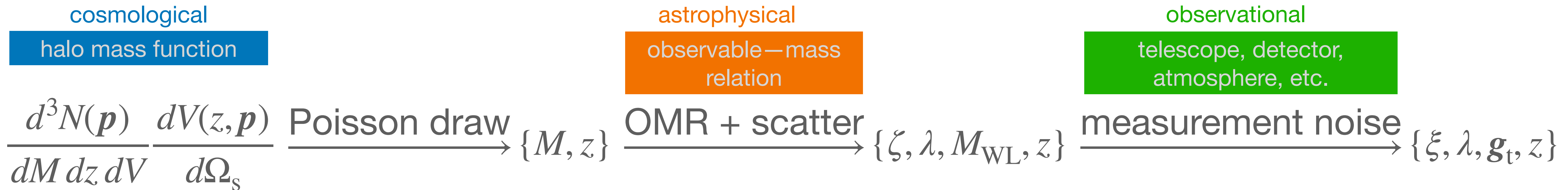
- Almost 700 SPT clusters (redshift 0.25—0.95) with DES Y3 shear
 - For the experts:
 - Analysis uses individual cluster shear profiles (Stacks are shown for visualization purposes)
 - Same source selection as in DES Y3 3x2pt
 - Same photo-z and shear calibrations
 - Radial range: $0.5 < r [h^{-1}\text{Mpc}] < 3.2 / (1 + z)$
(avoid cluster centers, stay in 1-halo term regime)
 - 39 high-redshift clusters (redshift 0.6—1.7) with the Hubble Space Telescope
[Schraback+18](#), [Schraback](#), [Bocquet+21](#), [Zohren](#), [Schraback](#), [Bocquet+22](#)



Likelihood Function I

Bayesian Population Modeling

Let us generate a cluster dataset!



Differential multi-observable cluster abundance

$$\frac{d^4 N(\mathbf{p})}{d\xi d\lambda d\mathbf{g}_t dz} = \int \dots \int dM d\zeta d\tilde{\lambda} dM_{\text{WL}} d\Omega_s P(\xi | \zeta) P(\lambda | \tilde{\lambda}) P(\mathbf{g}_t | M_{\text{WL}}) P(\zeta, \lambda, M_{\text{WL}} | M, z, \mathbf{p}) \frac{d^2 N(\mathbf{p})}{dM dV} \frac{d^2 V(z, \mathbf{p})}{dz d\Omega_s}$$

marginalize over
latent variables

Likelihood Function II

Poisson likelihood function: $\mathcal{L}(k \text{ events} \mid \text{rate } \mu) \propto \mu^k e^{-\mu} \Rightarrow \ln \mathcal{L} = k \ln(\mu) - \mu$

$$\ln \mathcal{L}(\mathbf{p}) = \sum_i \ln \left. \frac{d^4 N(\mathbf{p})}{d\xi d\lambda d\mathbf{g}_t dz} \right|_{\xi_i, \lambda_i, \mathbf{g}_{t,i}, z_i} - \int \dots \int d\xi d\lambda d\mathbf{g}_t dz \frac{d^4 N(\mathbf{p})}{d\xi d\lambda d\mathbf{g}_t dz} \Theta_s(\xi, \lambda, z) + \text{const.}$$

Likelihood Function II

Poisson likelihood function: $\mathcal{L}(k \text{ events} \mid \text{rate } \mu) \propto \mu^k e^{-\mu} \Rightarrow \ln \mathcal{L} = k \ln(\mu) - \mu$

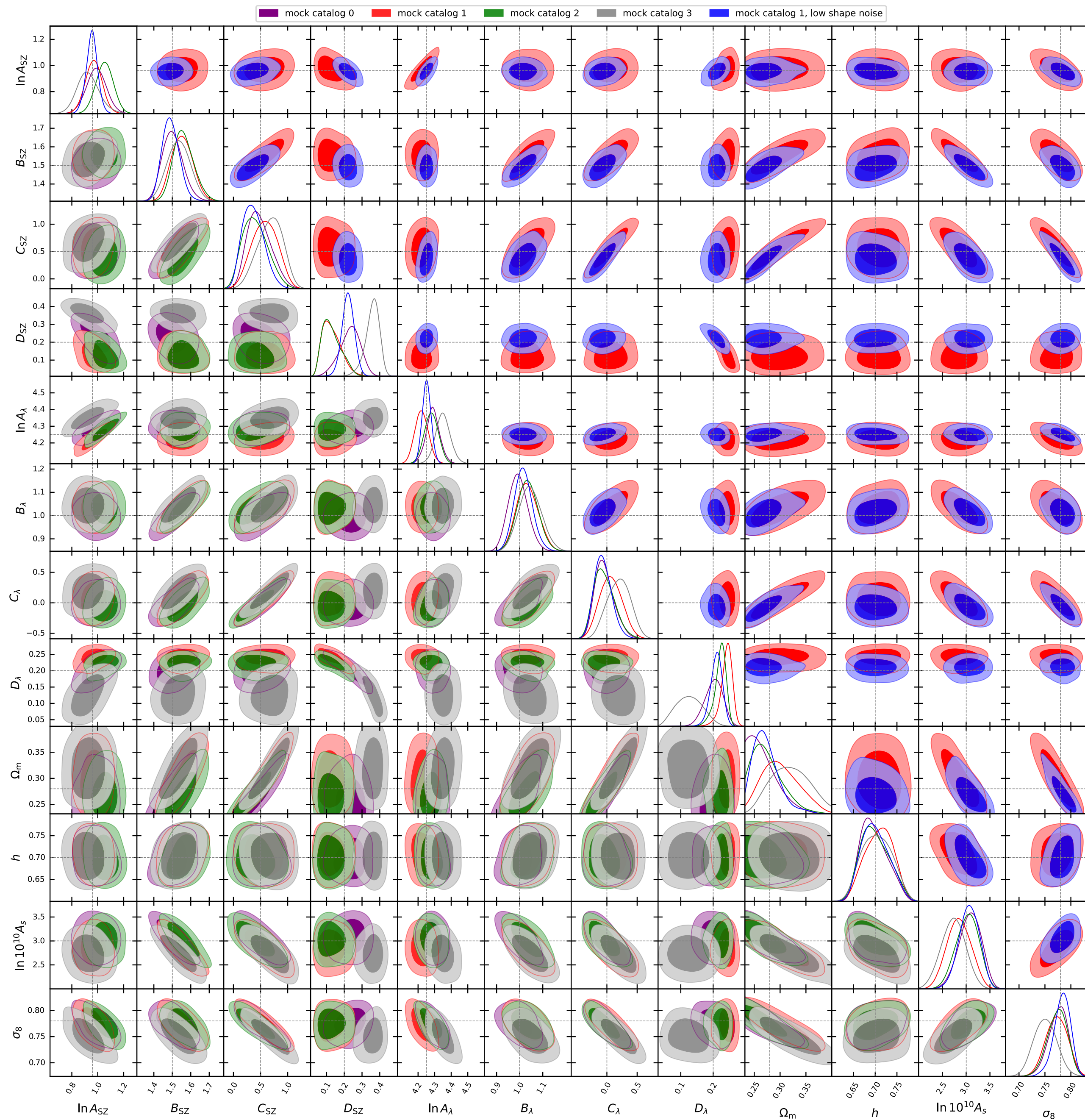
$$\ln \mathcal{L}(\mathbf{p}) = \sum_i \ln \left. \frac{d^4 N(\mathbf{p})}{d\xi d\lambda d\mathbf{g}_t dz} \right|_{\xi_i, \lambda_i, \mathbf{g}_{t,i}, z_i} - \int \dots \int d\xi d\lambda d\mathbf{g}_t dz \frac{d^4 N(\mathbf{p})}{d\xi d\lambda d\mathbf{g}_t dz} \Theta_s(\xi, \lambda, z) + \text{const.}$$

can be re-written as

$$\ln \mathcal{L}(\mathbf{p}) = \underbrace{\sum_i \ln \int_{\lambda_{\text{cut}}}^{\infty} d\lambda \left. \frac{d^3 N(\mathbf{p})}{d\xi d\lambda dz} \right|_{\xi_i, z_i} - \int_{z_{\text{cut}}}^{\infty} dz \int_{\xi_{\text{cut}}}^{\infty} d\xi \int_{\lambda_{\text{cut}}}^{\infty} d\lambda \frac{d^3 N(\mathbf{p})}{d\xi d\lambda dz}}_{\text{cluster abundance likelihood}} + \sum_i \ln \frac{\left. \frac{d^4 N(\mathbf{p})}{d\xi d\lambda d\mathbf{g}_t dz} \right|_{\xi_i, \lambda_i, \mathbf{g}_{t,i}, z_i}}{\int_{\lambda_{\text{cut}}}^{\infty} d\lambda \left. \frac{d^3 N(\mathbf{p})}{d\xi d\lambda dz} \right|_{\xi_i, z_i}} + \text{const.}$$

$$\frac{\left. \frac{d^4 N(\mathbf{p})}{d\xi d\lambda d\mathbf{g}_t dz} \right|_{\xi_i, \lambda_i, \mathbf{g}_{t,i}, z_i}}{\int_{\lambda_{\text{cut}}}^{\infty} d\lambda \left. \frac{d^3 N(\mathbf{p})}{d\xi d\lambda dz} \right|_{\xi_i, z_i}} = \frac{P(\lambda, \mathbf{g}_t, \xi, z \mid \mathbf{p})}{P(\lambda > \lambda_{\text{cut}}, \xi, z \mid \mathbf{p})} \equiv P(\lambda, \mathbf{g}_t \mid \lambda > \lambda_{\text{cut}}, \xi, z, \mathbf{p})$$

conditional “mass calibration likelihood”



Pipeline Verification

using mock datasets created from the model

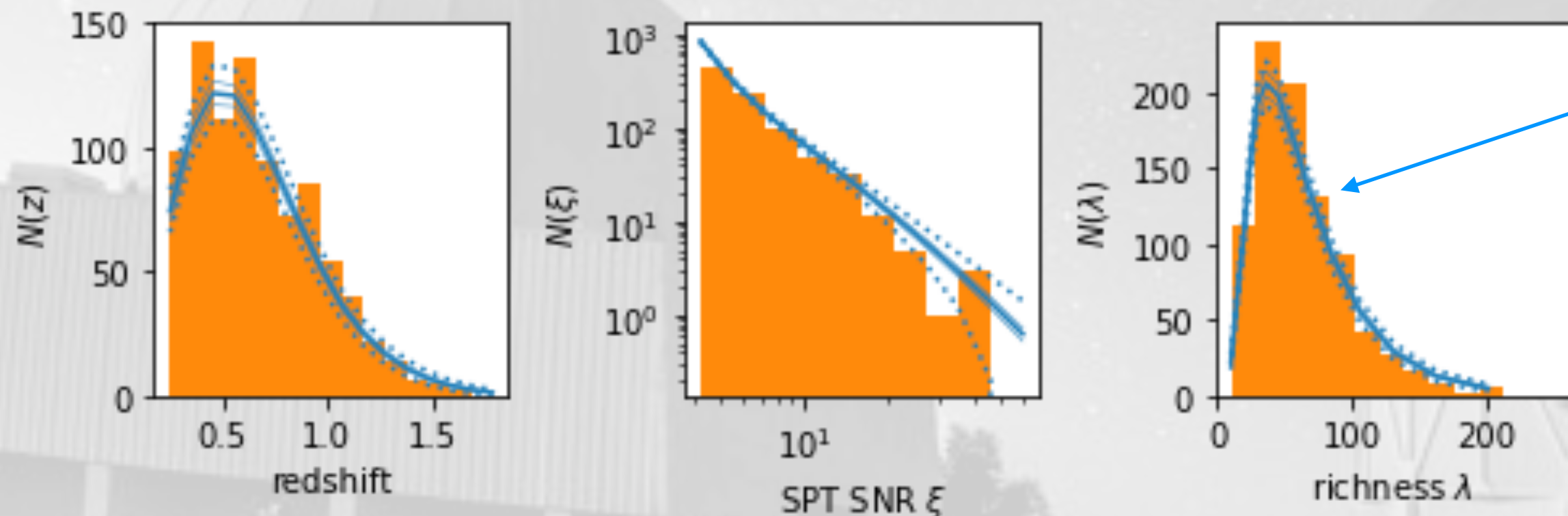
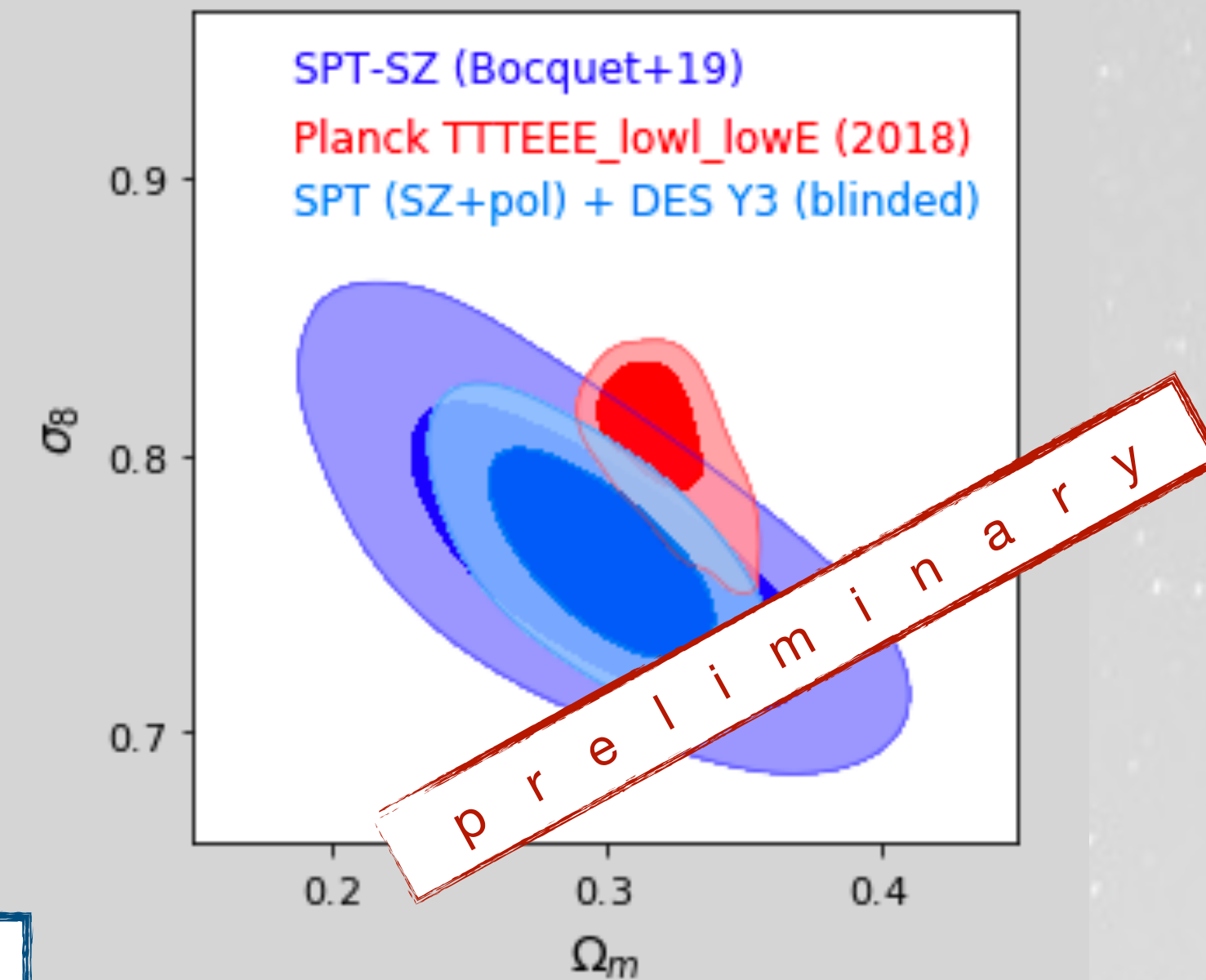
- Create synthetic clusters from the halo mass function using observable—mass relations
- Analyze several statistically independent mock realizations
- Pipeline recovers input values
- We correctly implemented the analysis framework!

SPT-SZ + SPTpol + DES Year 3 Weak Lensing

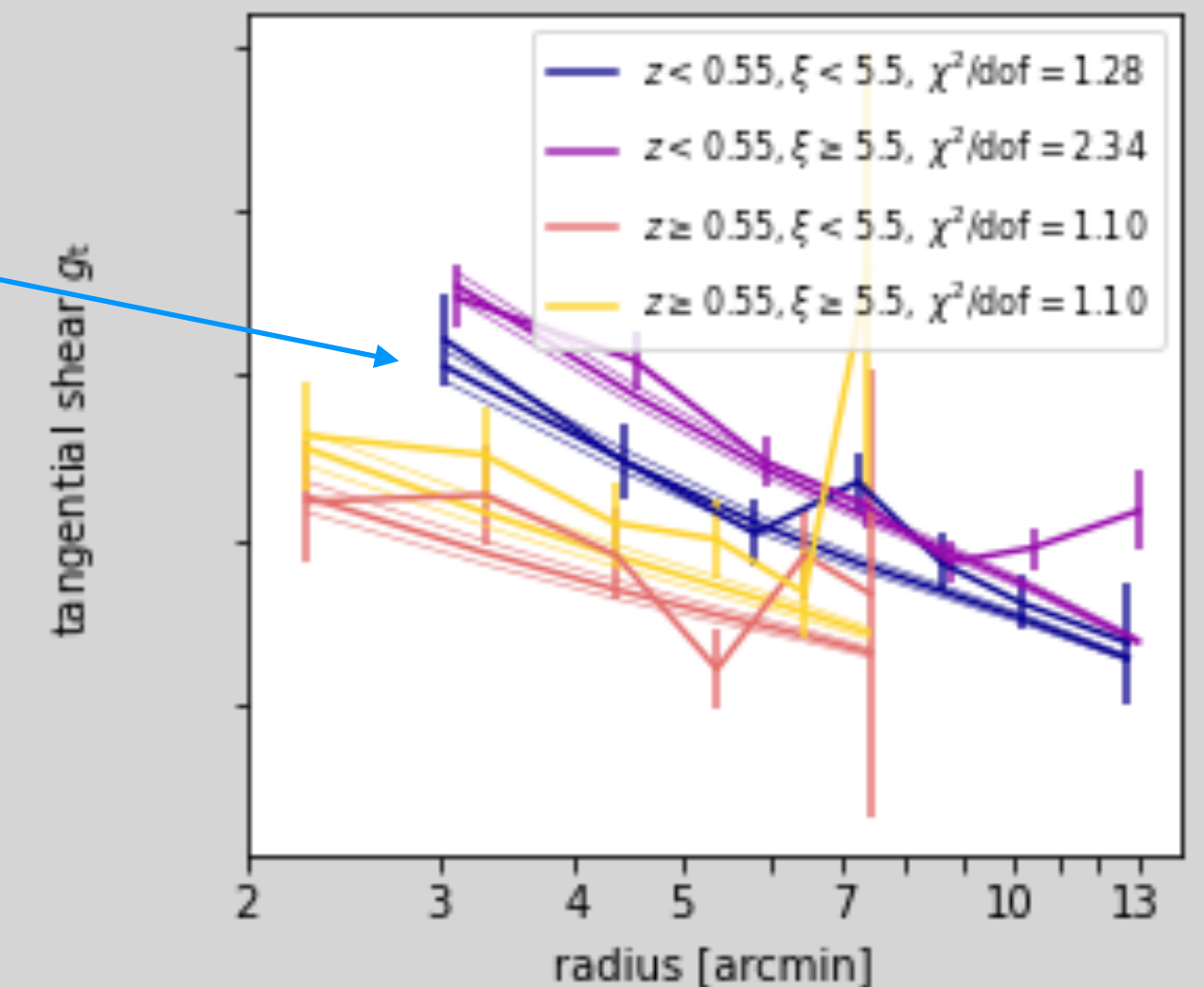
Update on blinded results

- Confirm robustness of lensing modeling blindly (centering, boost factors, radial cuts)
- Model is a good description of the data
- No roadblock toward unblinding!

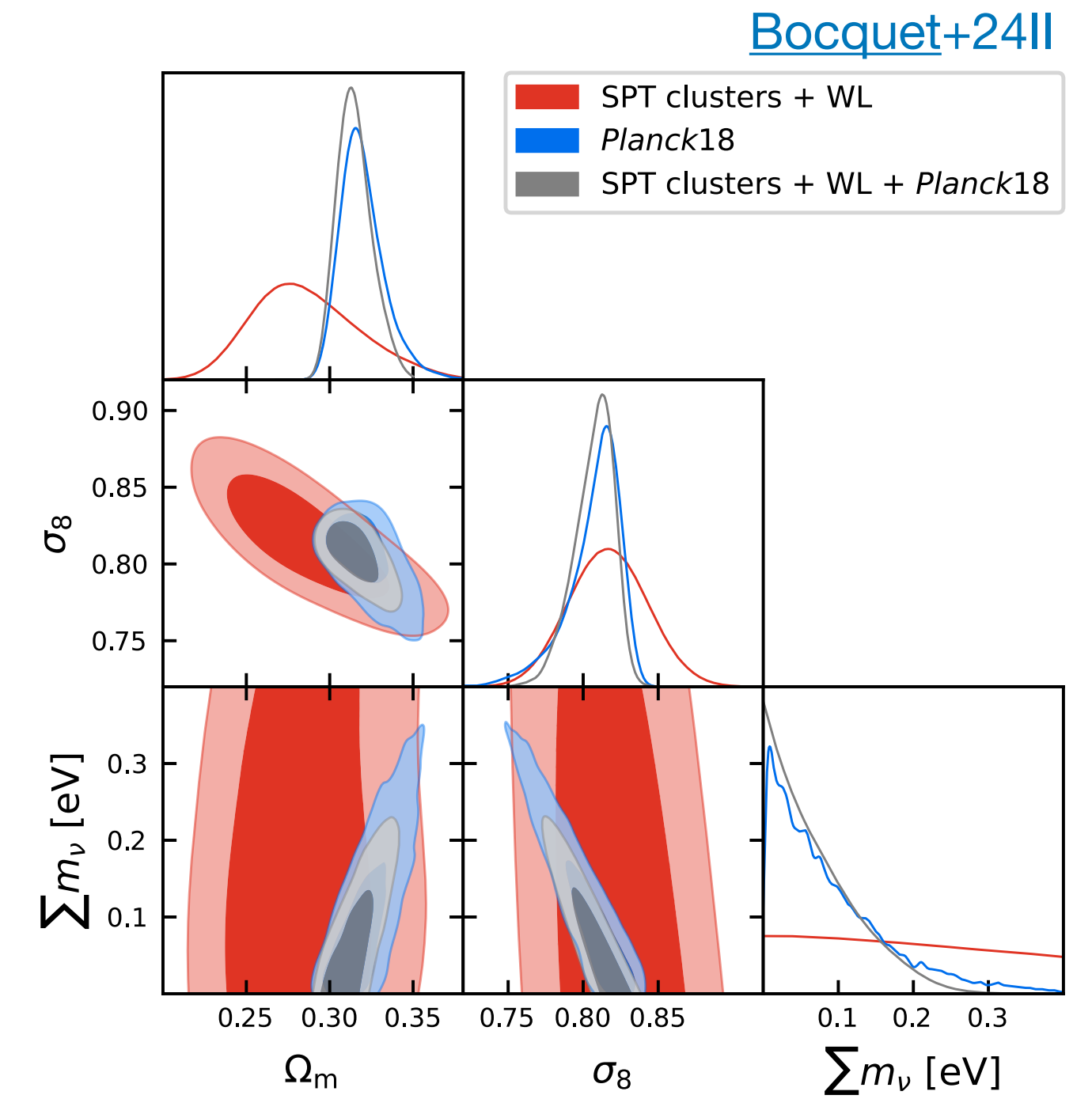
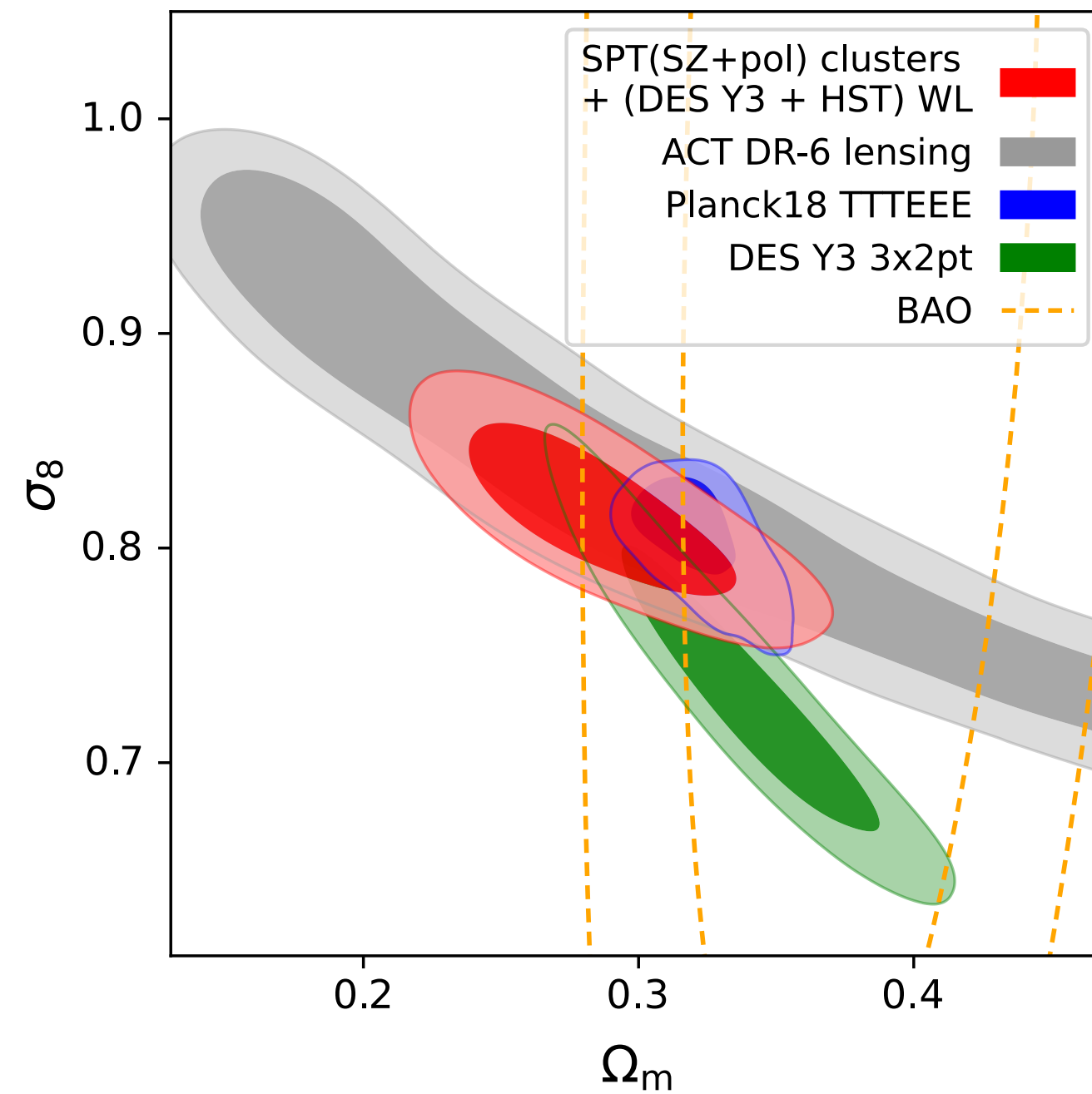
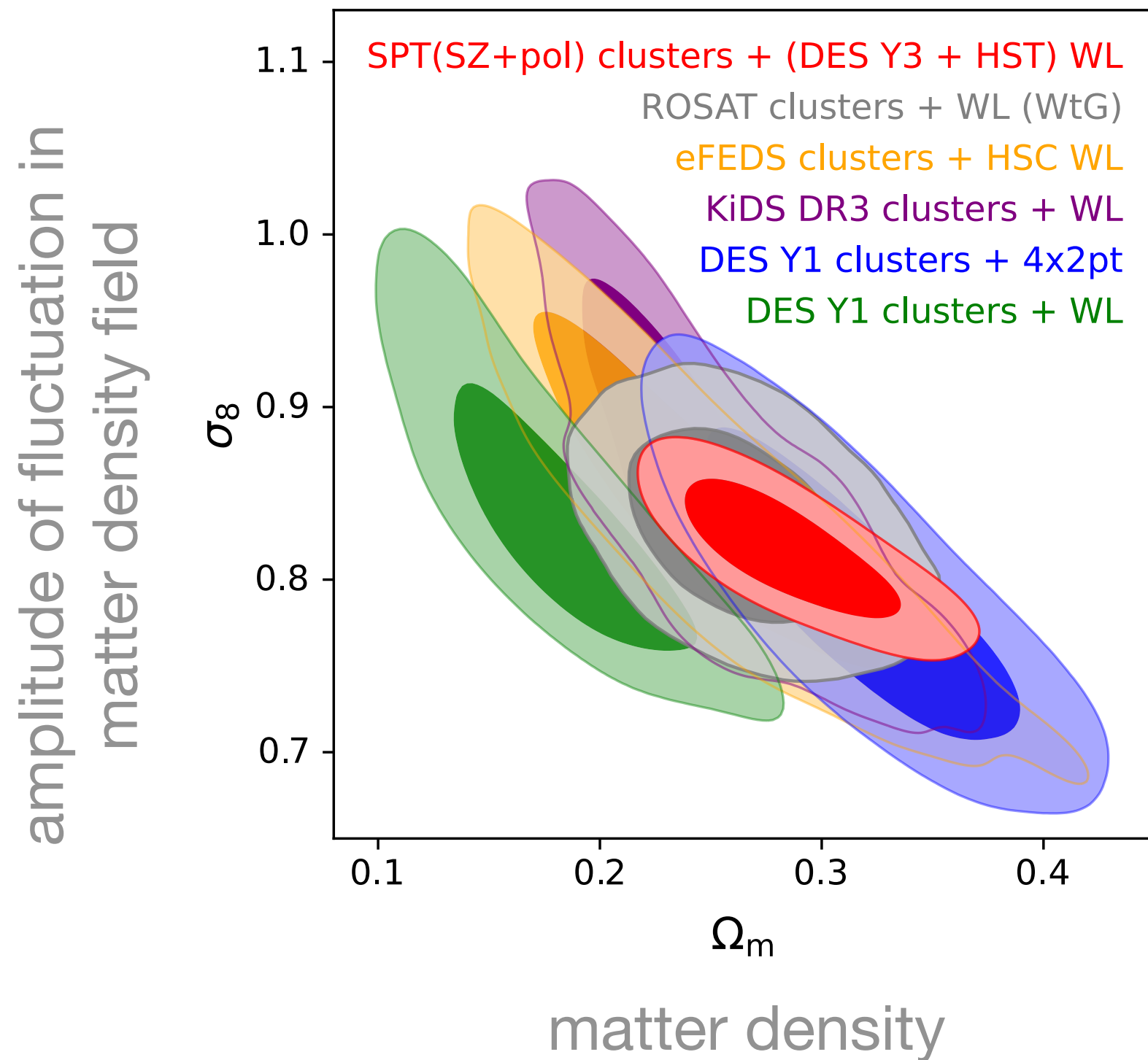
Slide from
AstroParticle Symposium 2022!



Model lines from
blinded run



Λ CDM with massive neutrinos

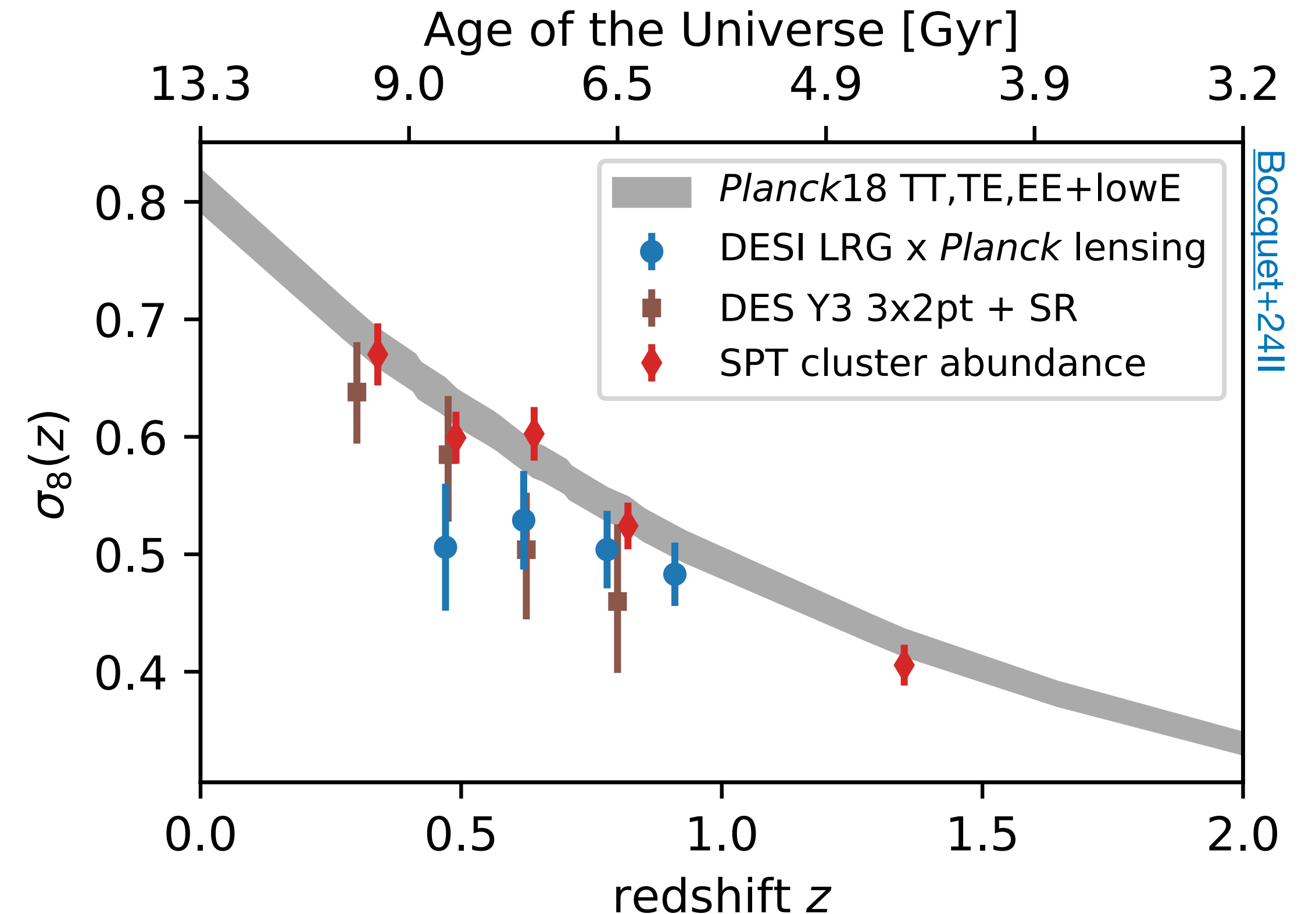


- Competitive constraints, especially on $S_8^{\text{opt}} \equiv \sigma_8 (\Omega_m/0.3)^{0.25}$
- No evidence for S_8 tension (difference with Planck 1.1σ)
- In combination with Planck $\sum m_\nu < 0.18 \text{ eV}$ (95 % C.L.)

Tracing the Growth of Structure

Phenomenological test

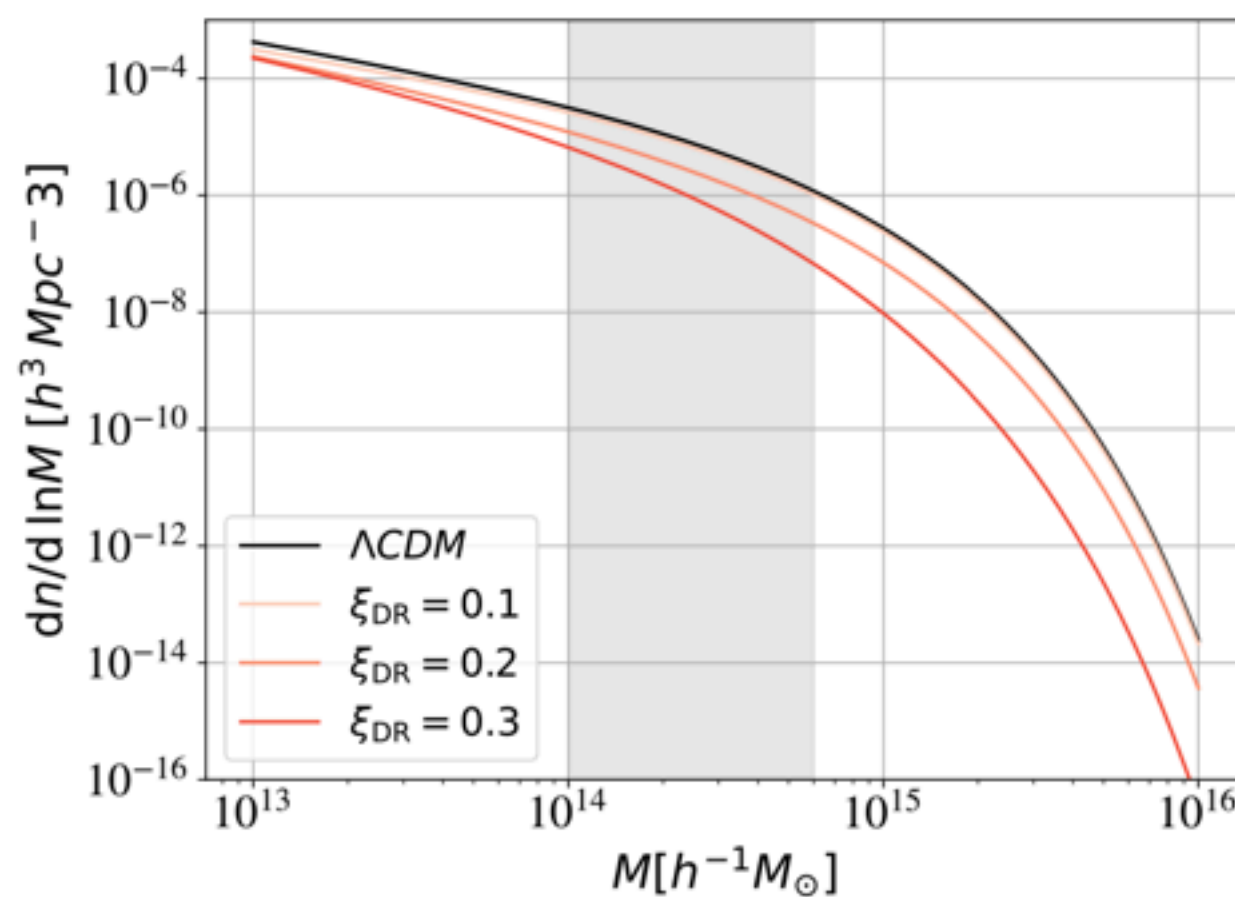
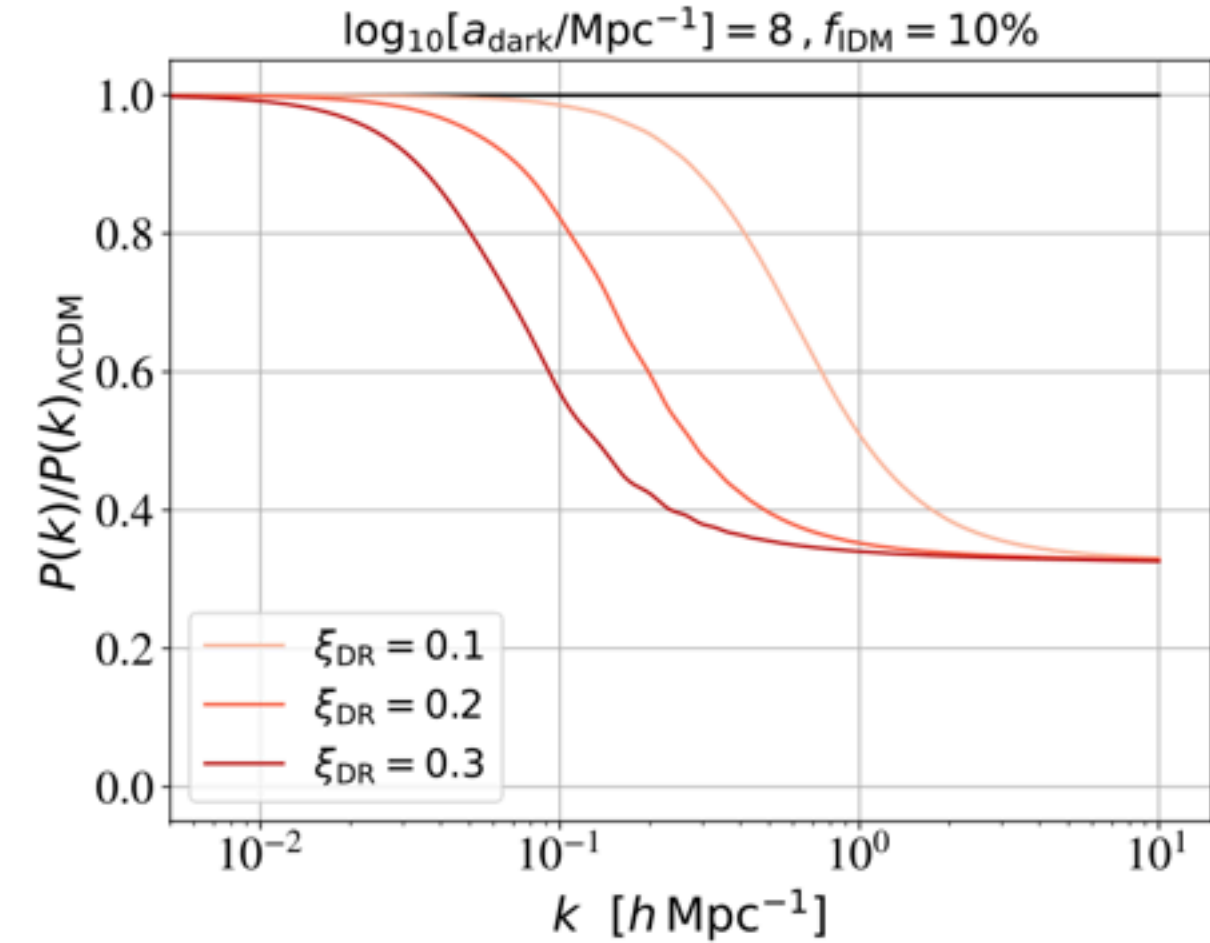
- Five bins in redshift with equal number of clusters
- Fit for independent amplitudes $\sigma_8(z)$
- With loose prior on Ω_m from the sound horizon at recombination θ_*
- Good agreement with Λ CDM model and *Planck* parameters from $z = 0.25$ to $z = 1.8$



Outlook

Select Work by PhD Students

Mazoun, Bocquet, Garny, Mohr, Rubira, Vogt 24
(arXiv:2312.17622)



Asmaa Mazoun

Interacting dark sector models

Analysis of SPT+DES dataset ongoing
(Mazoun+ in prep.)



Sophie Vogt

f(R) and nDGP models

Analysis of SPT+DES dataset done
(Vogt+ arXiv:2409.13556)

Vogt, Bocquet, Davies, Mohr, Schmidt 24
(arXiv:2401.09959)

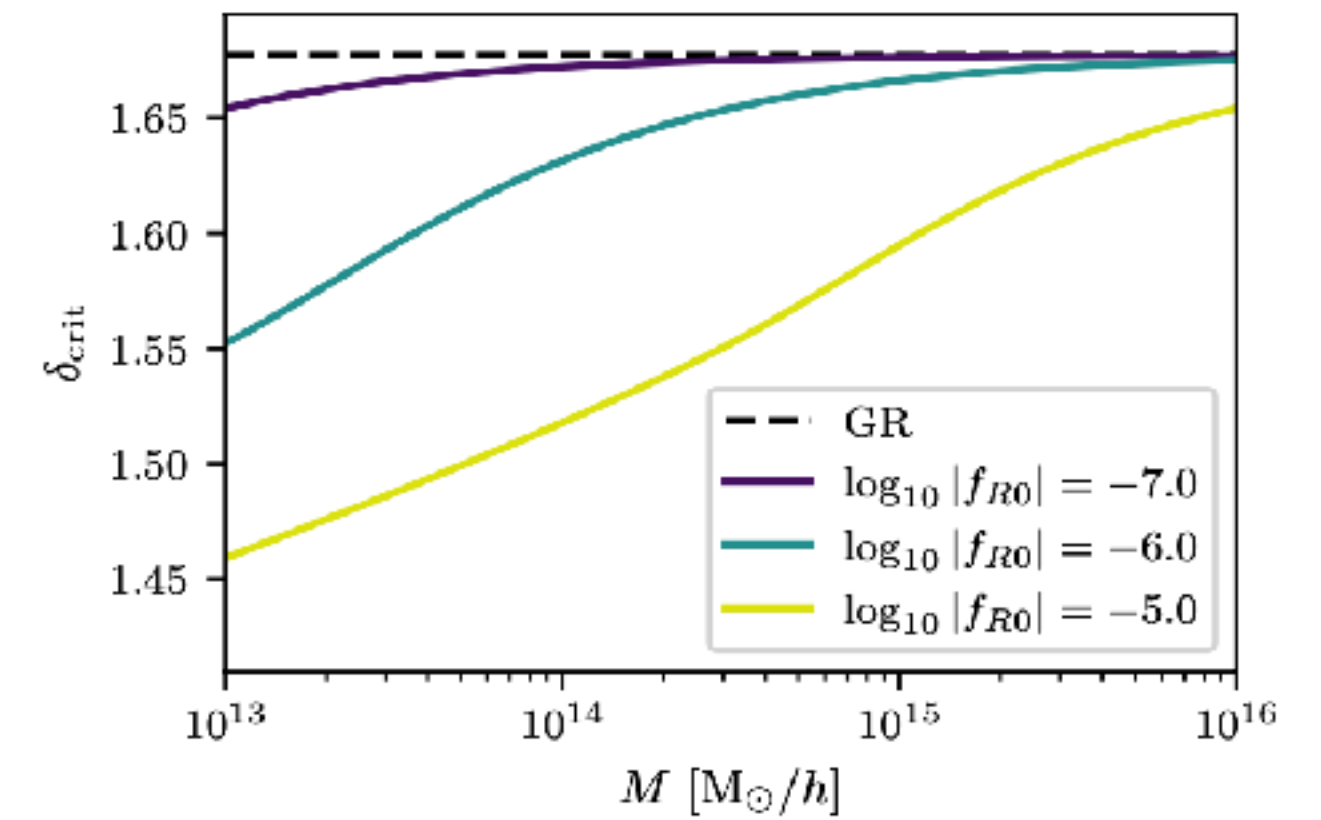


FIG. 1. The critical overdensity δ_{crit} for spherical collapse in $f(R)$ gravity (Eq. (12)) for different values of $\log_{10}|f_{R0}|$ at collapse redshift $z_c = 0$ in colored solid lines. The dashed black line represents δ_{crit} in a corresponding GR cosmology (Eq. (13)).

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(arXiv:2310.12213 — Phys. Rev. D 2024, 110, 083509)

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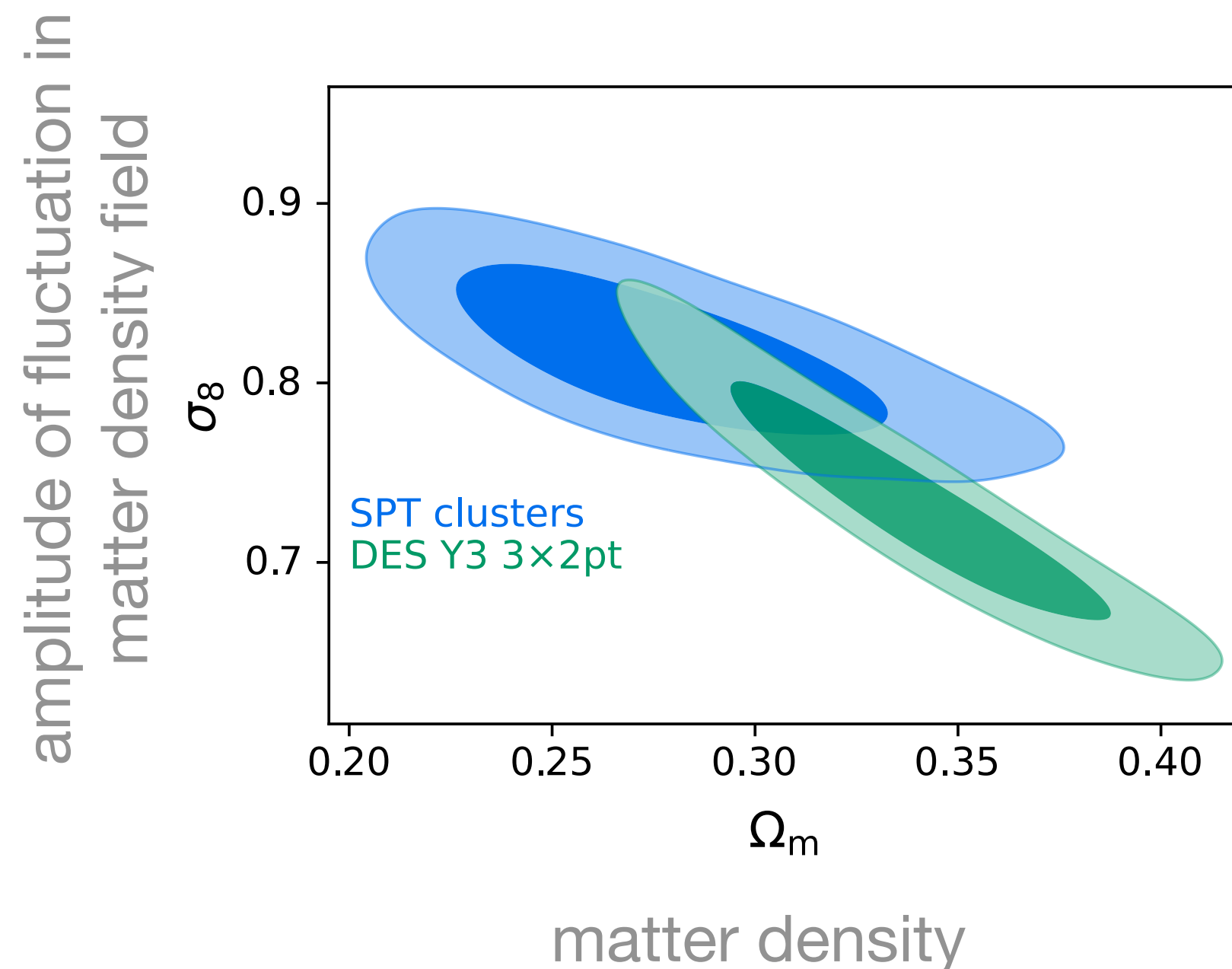
Bocquet, Grandis, Bleem, Klein, Mohr, Schrabback, SPT, DES
(arXiv:2401.02075 — Phys. Rev. D 2024, 110, 083510)

Multiprobe Cosmology from the Abundance of SPT Clusters and DES Galaxy Clustering and Weak Lensing

Bocquet, Grandis, Krause, To, SPT, DES
(to be submitted)

Outlook: Joint Constraints

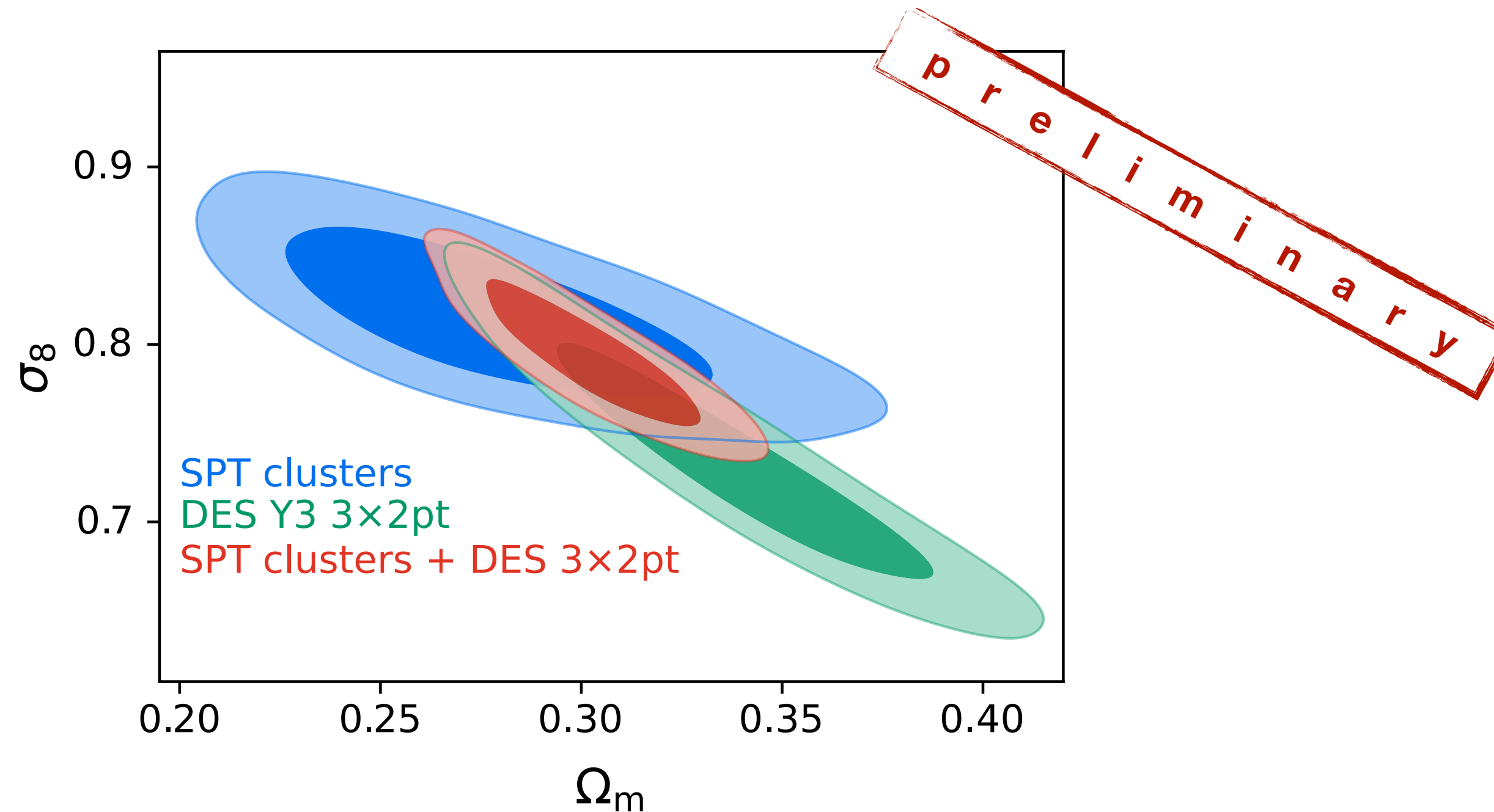
SPT Cluster Abundance + DES 3x2 pt = Multiprobe Cosmology



- Joint analysis (w/ Chun-Hao To, Elisabeth Krause, Sebastian Grandis)
- Cosmological covariance (negligible)
 - SPT cluster abundance is dominated by shot noise
 - SPT cluster mass calibration limited by lensing shape noise
- Shared systematics (same DES Y3 lensing data)
- Expect powerful constraints on $z < 2$ large-scale structure
- Ideal complement to high-redshift CMB measurements by *Planck*

Λ CDM with massive neutrinos

SPT clusters + DES 3x2pt



Contours are only 15% wider than *Planck* 2018 TT,TE,EE

Independent constraint on Hubble parameter

No strong suggestion for S8 tension (1.7 σ difference with *Planck*)

Summary

Cluster abundance as a cosmological probe

SZ-selection + weak-lensing mass calibration
= excellent control over systematics

Latest analysis of SPT (SZ+pol) clusters with
DES Y3 + HST lensing is compatible with and
complementary to other probes

Joint SPT clusters + DES 3x2pt analysis yields
tight constraints