



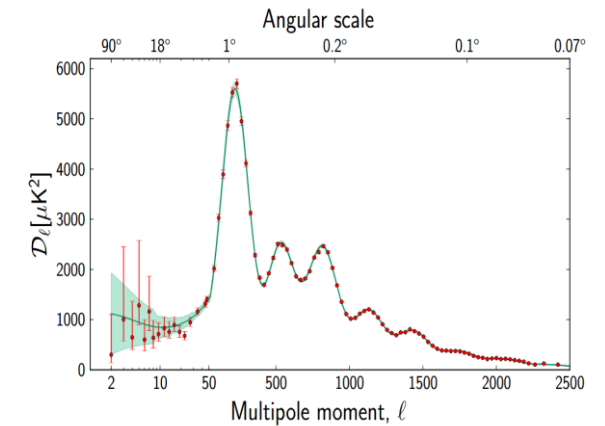
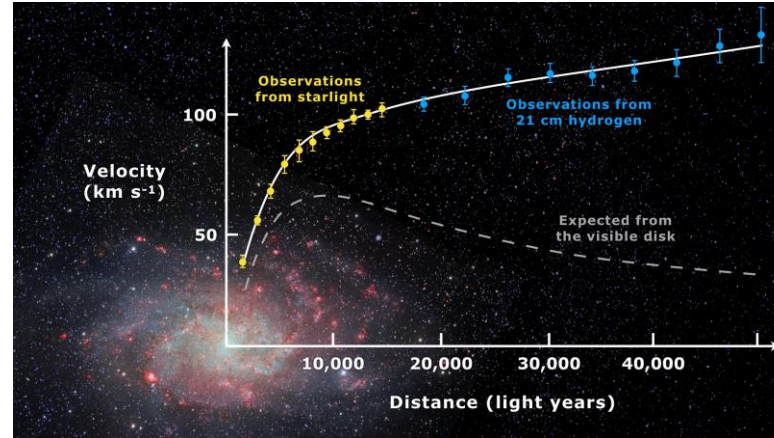
université
PARIS-SACLAY

Direct Dark Matter Research : The Damic-M* experiment

MELLOUKI ISSAM-EDDINE

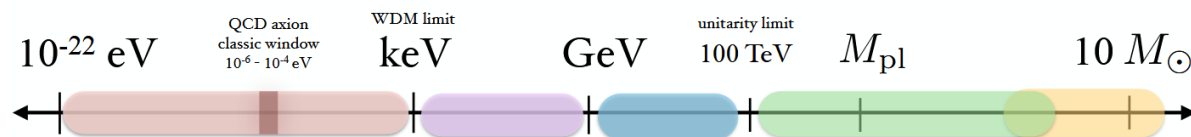
Evidence for Dark Matter in the Universe

- ❖ Galaxy Rotation Curves
- ❖ Lensing
- ❖ Cluster Virialization, Mergers
- ❖ CMB Anisotropies
- ❖ Etc.



Mass scale of dark matter

(not to scale)



“Ultralight” DM

non-thermal
bosonic fields

“Light” DM

dark sectors
sterile ν
can be thermal

WIMP

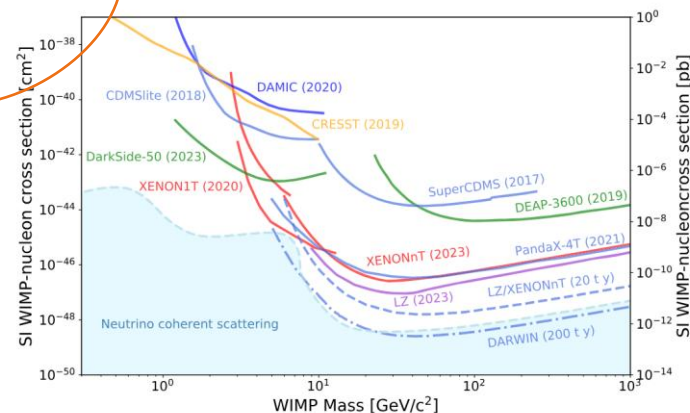
Composite DM

(Q-balls, nuggets, etc)

Primordial
black holes

T. Lin, arxiv:1904.07915

WIMP Waning,
Nearly Excluded

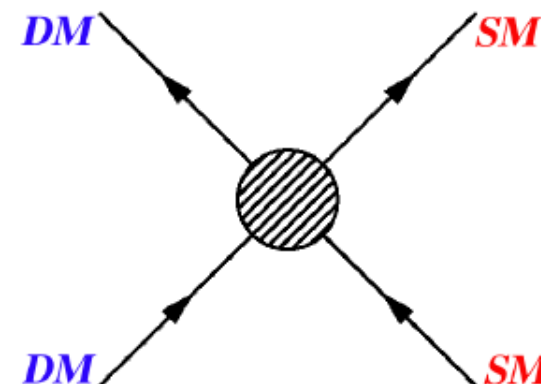


Detection of electronic/nuclear recoils

Detection of particles produced by the annihilation of SM particles

thermal freeze-out (early Univ.)
indirect detection (now)

direct detection



production at colliders

Search for missing
transverse energy

An alternative to the Wimp paradigm : The Hidden Sector

- ❖ The SM is not complete (neutrino, dark matter, ...)
- ❖ New Interactions may exist !
- ❖ Simple natural extension of the Standard Model : new U(1) symmetry
 - ❖ Dark QED
 - ❖ Electroweak extension
- ❖ Leads to MeV-scale Dark Matter

New bosons expected to mediate new interactions

New spin-1 bosons \leftrightarrow

new gauge symmetries beyond $SU(3) \times SU(2) \times U(1)$

Simplest possibility

$$SU(3) \times SU(2) \times U(1) \times \text{extra } U(1)$$

new gauge coupling (g'') \leftrightarrow intensity of new interaction ($\propto g''^2$)

Pierre FAYET, "The U BOSON as a generalized DARK PHOTON"

$$\mathcal{D}_\mu = \partial_\mu - igT_3W_\mu^3 - \frac{i}{2}g'YB_\mu - \frac{i}{2}g''XC_\mu, \quad \text{Symmetry breaking}$$

Modification of the Z boson's coupling constant and a new gauge boson in the electroweak theory

$$\begin{aligned} A_\mu &= \sin \theta_W W_\mu^3 + \cos \theta_W B_\mu, \\ \tilde{Z}_\mu &= \cos \xi \cos \theta_W W_\mu^3 - \cos \xi \sin \theta_W B_\mu - \sin \xi C_\mu, \\ A'_\mu &= \sin \xi \cos \theta_W W_\mu^3 - \sin \xi \sin \theta_W B_\mu + \cos \xi C_\mu. \end{aligned}$$

$$g_A(T_3, Y, X) = eQ,$$

$$g_{\tilde{Z}}(T_3, Y, X) = g \cos \theta_W \cos \xi T_3 + g' \sin \theta_W \cos \xi \frac{Y}{2} + g'' \sin \xi \frac{X}{2},$$

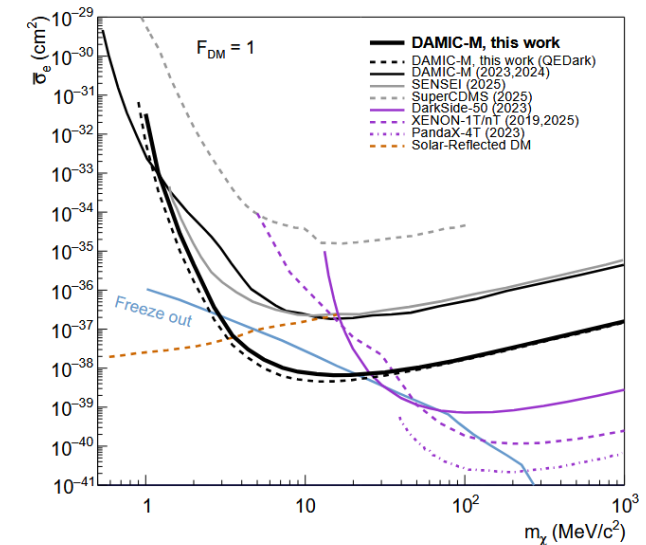
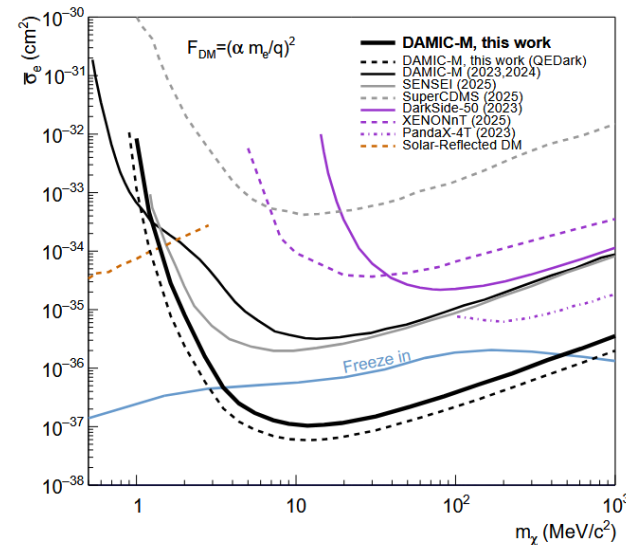
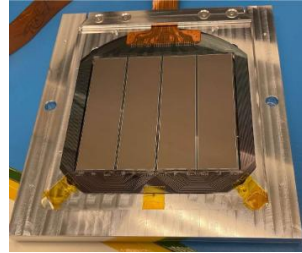
$$g_{A'}(T_3, Y, X) = g \cos \theta_W \sin \xi T_3 + g' \sin \theta_W \sin \xi \frac{Y}{2} + g'' \cos \xi \frac{X}{2}.$$

Olivier Deligny, DAMIC-M internal note

Damic-M in Dark Matter research

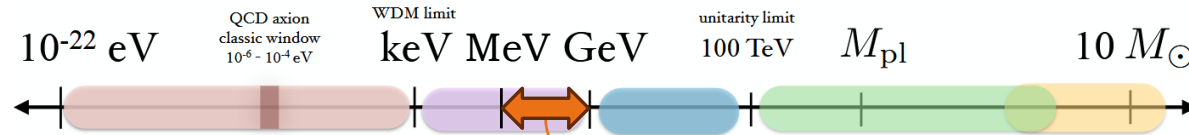
- ❖ Direct Dark Matter experiment
- ❖ Skipper-CCD detectors
- ❖ Single electron detection capability
 - ❖ For low-mass dark matter :

$$E_{Target} < 4 \frac{M_{DM}}{M_{Target}} E_{DM}$$
- ❖ Sub-electron readout noise
- ❖ World-leading limits on hidden sector Dark Matter



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(not to scale)



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“Light” DM

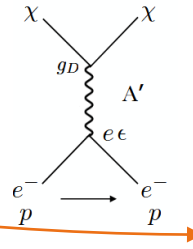
dark sectors
sterile ν
can be thermal

WIMP

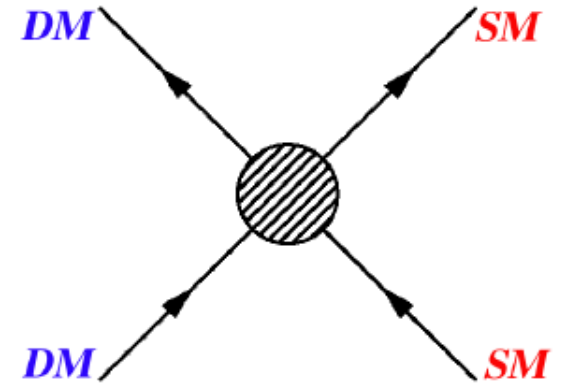
Composite DM
(Q-balls, nuggets, etc)

Primordial
black holes

DAMIC-M

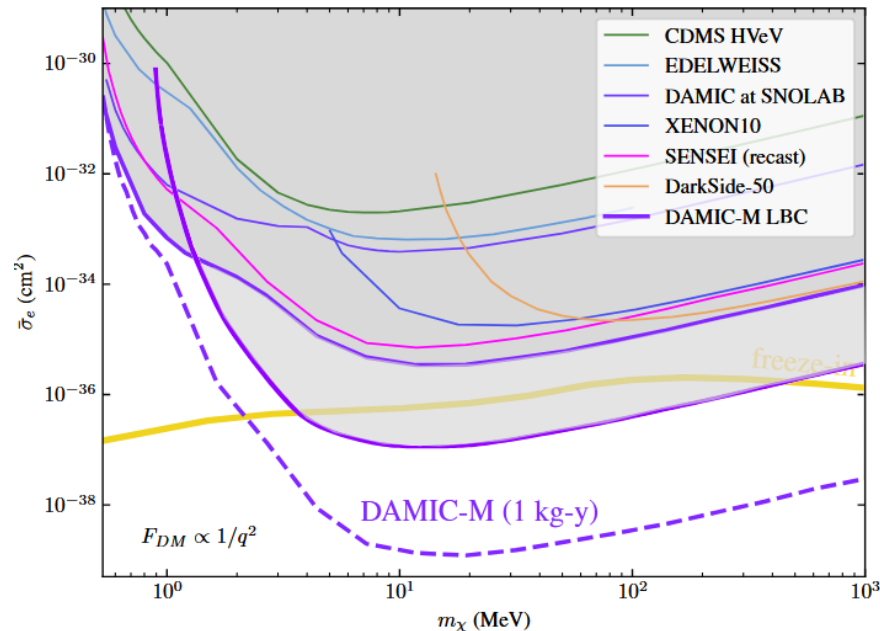


thermal freeze-out (early Univ.)
indirect detection (now)



production at colliders

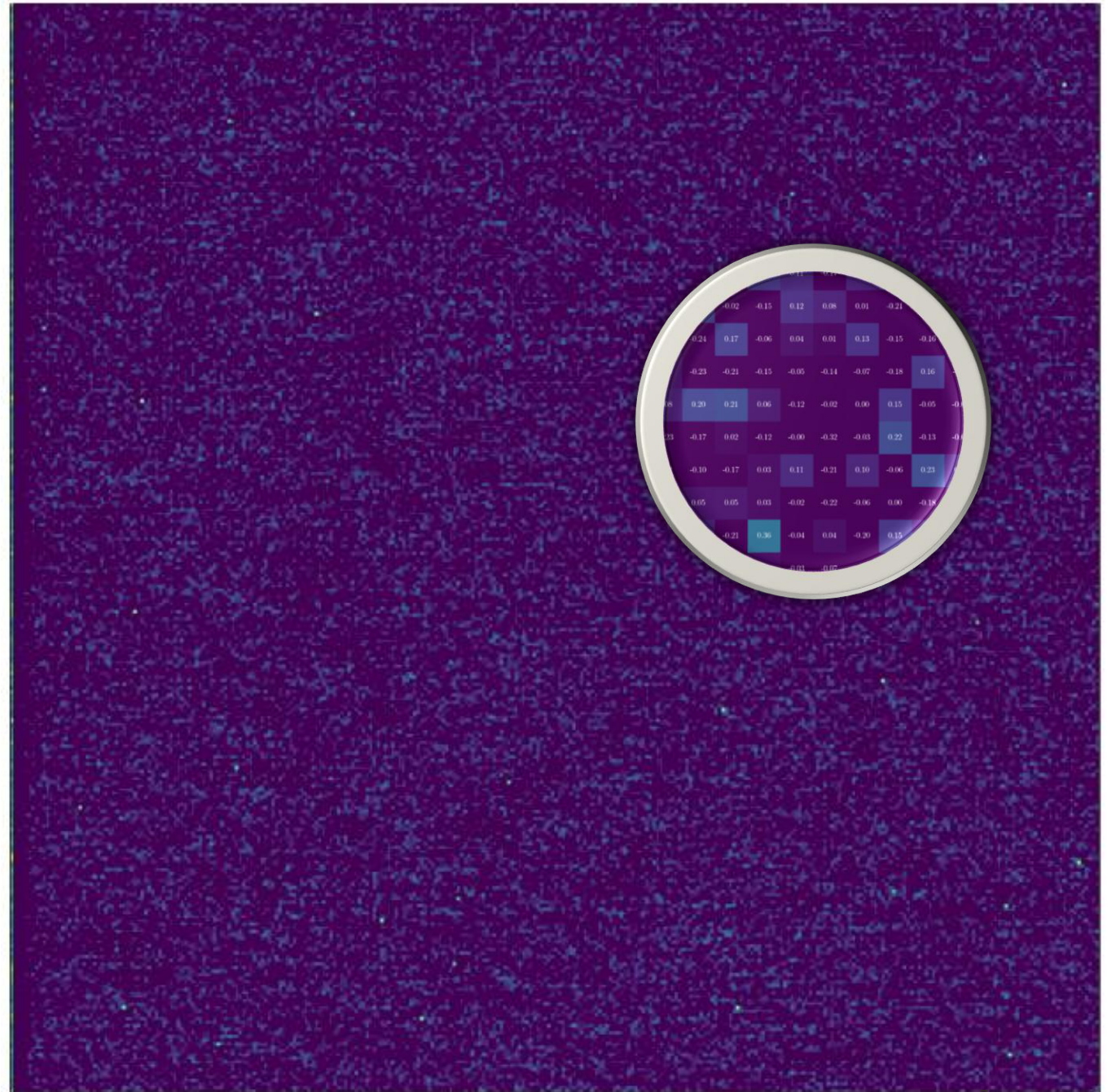
Max Planck Institute for Nuclear Physics



T. Lin, arxiv:1904.07915

The data

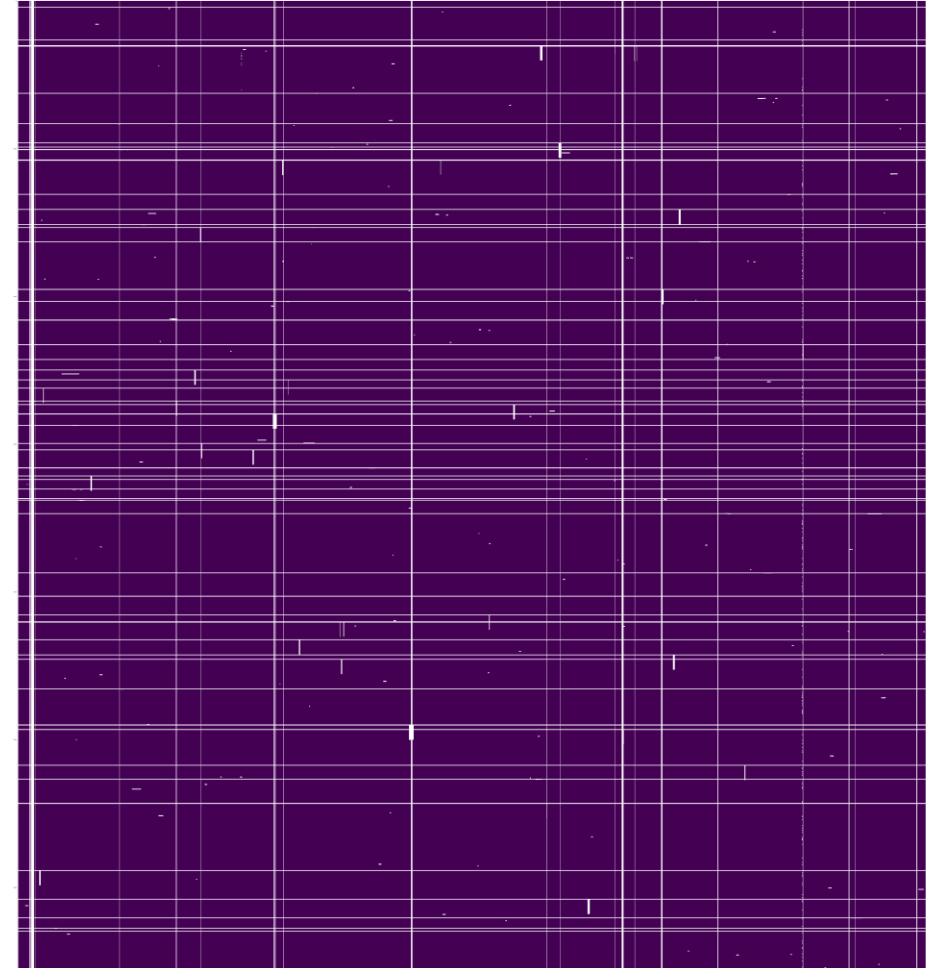
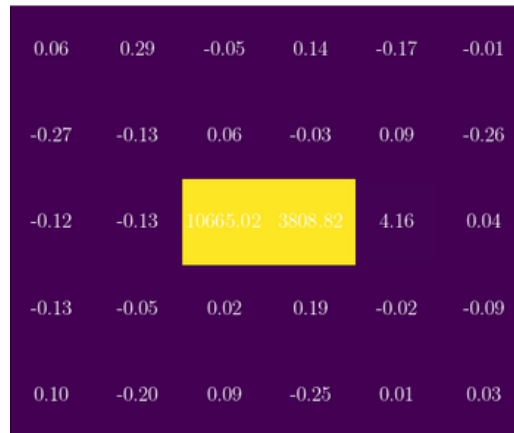
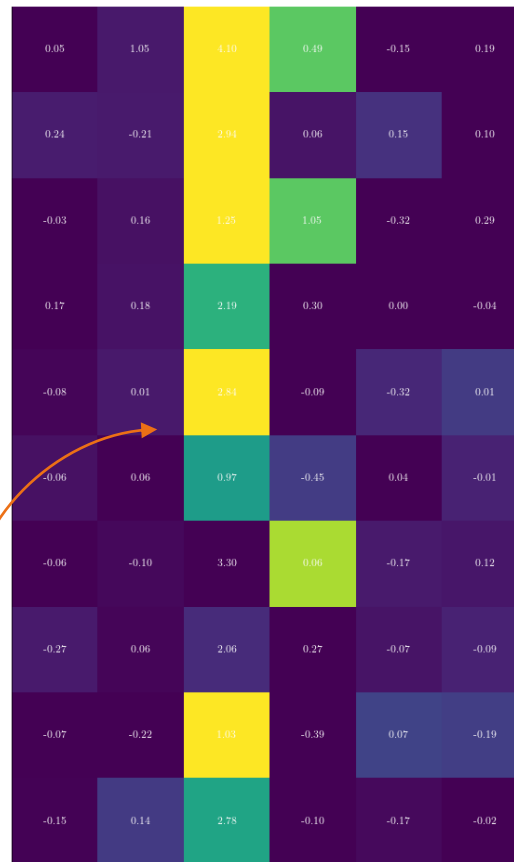
- ❖ ~8 days of DAMIC-M LBC data
- ❖ 4 CCDs → 4 images of 6560×6148 pixels
- ❖ $\sim 2.9 \cdot 10^{-4}$ kg.year exposure
- ❖ 500 skips for $\sigma \sim 0.17 e^-$



Portion of the image from CCD 1

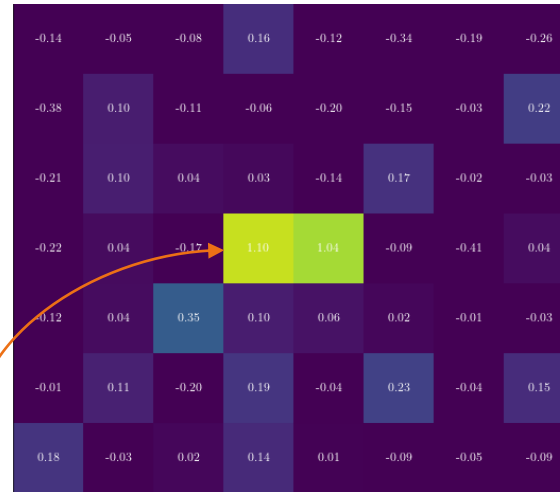
Image masking

- ❖ Hot region (large $1e^-$ rate)
- ❖ Defects generating charges continuously
- ❖ Clusters of high-charge pixels ($>5 e^-$)
- ❖ Cross-talk between CCDs
- ❖ 79.75% of data are kept



Pattern analysis and exclusion limits

- ❖ Searching for {11}, {111} or {12}/{21} events
- ❖ $\sim 2.9 \cdot 10^{-4}$ kg.year exposure
- ❖ ~ 13.5 expected events / 14 observed
- ❖ Comparison with observed events
 - ❖ Background is compatible with observations
 - ❖ No evidence for Dark Matter
- ❖ Dark matter could be hidden in our events, even if they are compatible with the background
- ❖ Feldman-Cousins table provides confidence interval limits for Poisson distributions

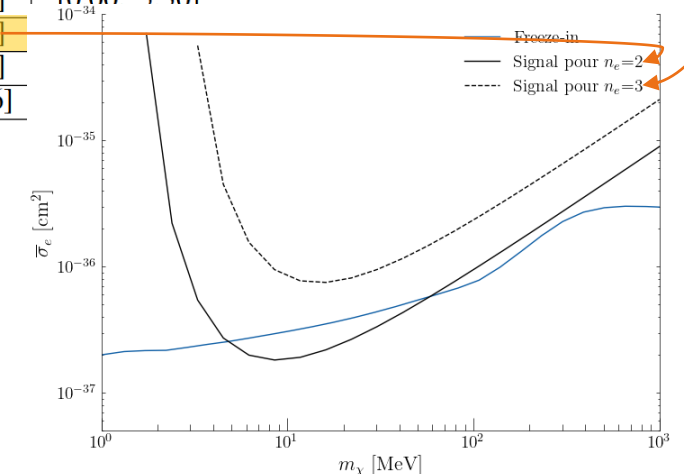


Attendu	CCD 1	CCD 2	CCD 3	CCD 4
{11}	3.2	2.6	4.9	2.8
{111}	$1 \cdot 10^{-3}$	$7 \cdot 10^{-4}$	$2 \cdot 10^{-3}$	$8 \cdot 10^{-4}$
{21}, {12}	$5 \cdot 10^{-4}$	$4 \cdot 10^{-4}$	$9 \cdot 10^{-4}$	$4 \cdot 10^{-4}$

Observations are compatible with the prediction based on the DC

Observé	CCD 1	CCD 2	CCD 3	CCD 4
{11}	3	3	8	0
{111}	0	0	0	0
{21}, {12}	0	0	0	0

$n_0 \setminus b$	0	13	14	15
0	[0.00, 2.44]	[0.00, 0.92]	[0.00, 0.92]	[0.00, 0.92]
2	[0.53, 5.91]	[0.00, 1.05]	[0.00, 1.06]	[0.00, 1.07]
4	[1.47, 8.60]	[0.00, 1.36]	[0.00, 1.39]	[0.00, 1.42]
6	[2.21, 11.47]	[0.00, 1.52]	[0.00, 1.55]	[0.00, 1.59]
8	[3.96, 13.99]	[0.00, 2.06]	[0.00, 2.11]	[0.00, 2.16]
10	[5.50, 16.50]	[0.00, 3.20]	[0.00, 3.27]	[0.00, 3.35]
12	[7.01, 19.00]	[0.00, 5.09]	[0.00, 5.19]	[0.00, 5.30]
14	[8.50, 21.50]	[0.00, 7.41]	[0.00, 7.50]	[0.00, 7.61]
16	[9.99, 23.99]	[0.00, 9.90]	[0.00, 9.99]	[0.00, 10.09]
18	[11.47, 26.16]	[0.00, 12.07]	[0.00, 12.16]	[0.00, 12.26]



Annexe

Dark Matter production scenarios

❖ Multiple production scenarios :

- ❖ Freeze-in/out
- ❖ Gravitational
- ❖ Etc.

❖ Freeze-out assumption : thermal equilibrium in early universe

❖ Freeze-in assumption : production from thermal bath species over time

- ❖ Y^2 is neglected

Boltzman equation : $L[f] = \mathcal{C}[f]$ $ds^2 = dt^2 - a^2(t)(dx^2 + dy^2 + dz^2)$

$$L[f] = E \frac{\partial f}{\partial t} - \|\mathbf{p}^2\| \frac{\dot{a}}{a} \frac{\partial f}{\partial E}$$

$$\frac{g}{(2\pi^3)} \int d^3p \frac{L[f]}{E} = \frac{dn}{dt} + 3Hn \quad n = \frac{g}{(2\pi^3)} \int d^3p f(p, T)$$

||

$$\frac{g_1}{(2\pi)^3} \int d^3p_1 \frac{\mathcal{C}[f_1]}{E_1} = \int d\pi_1 d\pi_2 d\pi_3 d\pi_4 (2\pi)^4 \delta_D^{(4)}\left(\sum_{i=1}^4 \mathbf{p}_i\right) \times$$

$$[f_1 f_2 (1 \pm f_3)(1 \pm f_4) |M_{12 \rightarrow 34}|^2 - f_3 f_4 (1 \pm f_1)(1 \pm f_2) |M_{34 \rightarrow 12}|^2]$$

$$\frac{dY}{dx} = \frac{\langle \sigma v \rangle}{Hx} s(Y^2 - Y_{Eq}^2) \left(1 + \frac{1}{3} \frac{T}{h} \frac{dh}{dT}\right)$$

Freeze-in setup

$$\frac{dY}{dx} = -\frac{\langle \sigma v \rangle}{Hx} s Y_{Eq}^2 \left(1 + \frac{1}{3} \frac{T}{h} \frac{dh}{dT}\right) \simeq -\frac{\langle \sigma v \rangle}{Hx} s Y_{Eq}^2$$

$$\frac{dY}{dT} = -\frac{M_P}{(2\pi)^2 T^6} \left(\frac{45}{\pi}\right)^{\frac{3}{2}} \frac{1}{g_{*s} \sqrt{g_{\text{eff}}}} \times$$

$$\left(\sum_i (\bar{n}_i^2 \langle \sigma v \rangle_{ii \rightarrow \chi \bar{\chi}}) + \sum_j (\bar{n}_j^2 \Gamma_{j \rightarrow \chi \bar{\chi}}) \right)$$

$$\langle \sigma v \rangle = \frac{\int \sigma v dn_{\chi}^{Eq} dn_{\bar{\chi}}^{Eq}}{n_{\chi}^{Eq} n_{\bar{\chi}}^{Eq}}$$

$$Y = \frac{n}{s}, \quad x = \frac{m_{\chi}}{T}$$

$$s = \frac{2\pi^2}{45} g_{*s}(T) T^3, \quad h = g_{\text{eff}}(T) \frac{\pi^2}{30} T^4$$

Freeze-In Production in a Dark QED Sector

$$\mathcal{D}_\mu = \partial_\mu - igT_3 W_\mu^3 - \frac{i}{2} \mathcal{Y} g' B_\mu - \frac{i}{2} Q_X g'' C_\mu \quad \mathcal{L}_{\text{mix}} = -\frac{\epsilon}{2} B_{\mu\nu} C^{\mu\nu}$$

$$\begin{aligned} \mathcal{L} &= -\frac{1}{4} W_{\mu\nu}^3 W^{3\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} C_{\mu\nu} C^{\mu\nu} - \frac{\epsilon}{2} B_{\mu\nu} C^{\mu\nu} \\ &= -\frac{1}{4} \hat{W}_{\mu\nu}^3 \hat{W}^{3\mu\nu} - \frac{1}{4} \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} - \frac{1}{4} \hat{C}_{\mu\nu} \hat{C}^{\mu\nu} \end{aligned}$$

Modification of the Z boson's coupling constant and a new gauge boson in the electroweak theory

$$\begin{pmatrix} \hat{W}_\mu^3 \\ \hat{B}_\mu \\ \hat{C}_\mu \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & \epsilon \\ 0 & 0 & \frac{1}{\sqrt{1-\epsilon^2}} \end{pmatrix} \begin{pmatrix} W_\mu^3 \\ B_\mu \\ C_\mu \end{pmatrix}$$

$$g_A = eQ$$

$$g_Z \simeq g \cos \theta_W T_3 - g' \sin \theta_W \frac{\mathcal{Y}}{2} - g'' \epsilon \sin \theta_W \frac{Q_X}{2}$$

$$g_{A'} \simeq \epsilon e \cos \theta_W T_3 - g' \epsilon (1 + \cos^2 \theta_W) \frac{\mathcal{Y}}{2} + g'' \frac{Q_X}{2}$$

Computation of the matrix element

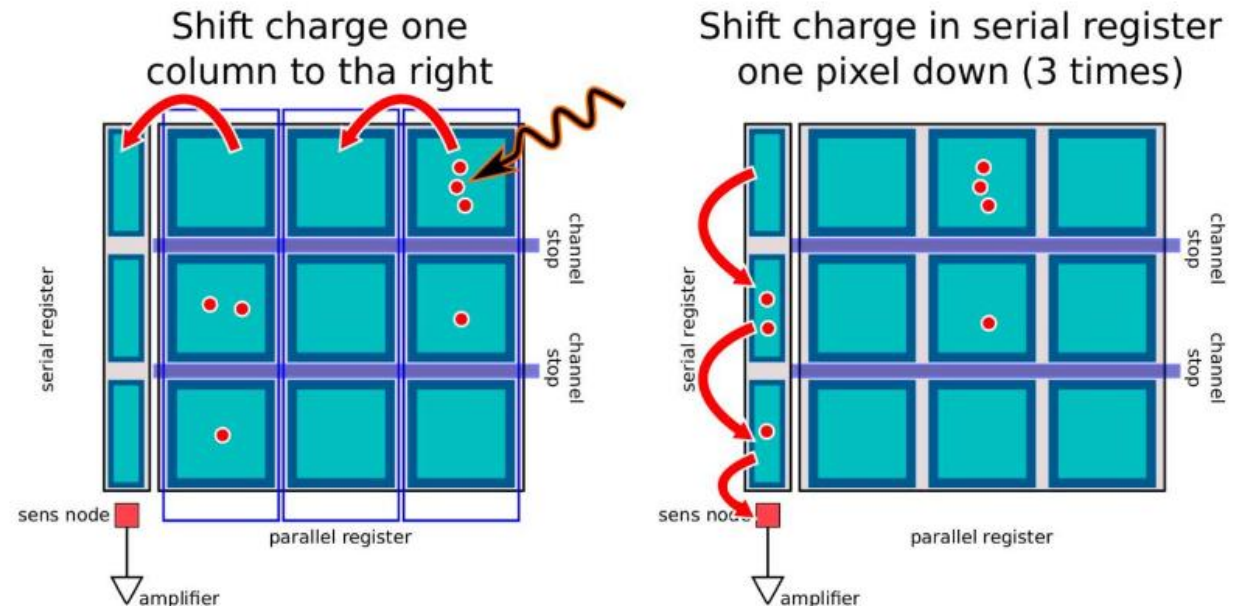
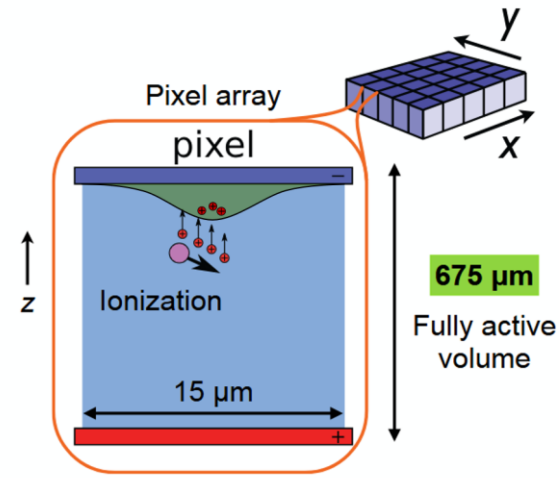
$$\int d\Omega^* |\overline{\mathcal{M}}|^2 = \frac{16\pi N_c}{3} \frac{(Qe)^2 g_{A'}^2 s^2}{(s - m_{A'}^2)^2} \left(1 + \frac{2m_i^2}{s}\right) \left(1 + \frac{2m_\chi^2}{s}\right)$$

$$\Omega_\chi h^2 = \frac{m_\chi s_{\text{today}} Y_{\text{today}}}{\rho_c / h^2} + \frac{dY}{dT} = -\frac{M_P}{(2\pi)^2 T^6} \left(\frac{45}{\pi}\right)^{\frac{3}{2}} \frac{1}{g_{*s} \sqrt{g_{\text{eff}}}} \times \left(\sum_i (\tilde{n}_i^2 \langle \sigma v \rangle_{ii \rightarrow \chi \bar{\chi}}) + \sum_j (\tilde{n}_j^2 \Gamma_{j \rightarrow \chi \bar{\chi}}) \right)$$

$$g_{A'}^2 = \frac{\Omega_\chi \rho_c}{s_0 m_\chi} \left[\int \frac{2M_P \left(\frac{45}{\pi}\right)^{3/2}}{(2\pi)^2 T^5 g_{*s}(T) \sqrt{g_{\text{eff}}(T)}} \int ds \frac{N_c \alpha_{\text{QED}} Q^2}{24\pi^4} \times \left(\frac{(s - 4m_\chi^2)(s - 4m_i^2)}{s} \right)^{1/2} K_1 \left(\frac{\sqrt{s}}{T} \right) \left(1 + \frac{2m_i^2}{s}\right) \left(1 + \frac{2m_\chi^2}{s}\right) \right]$$

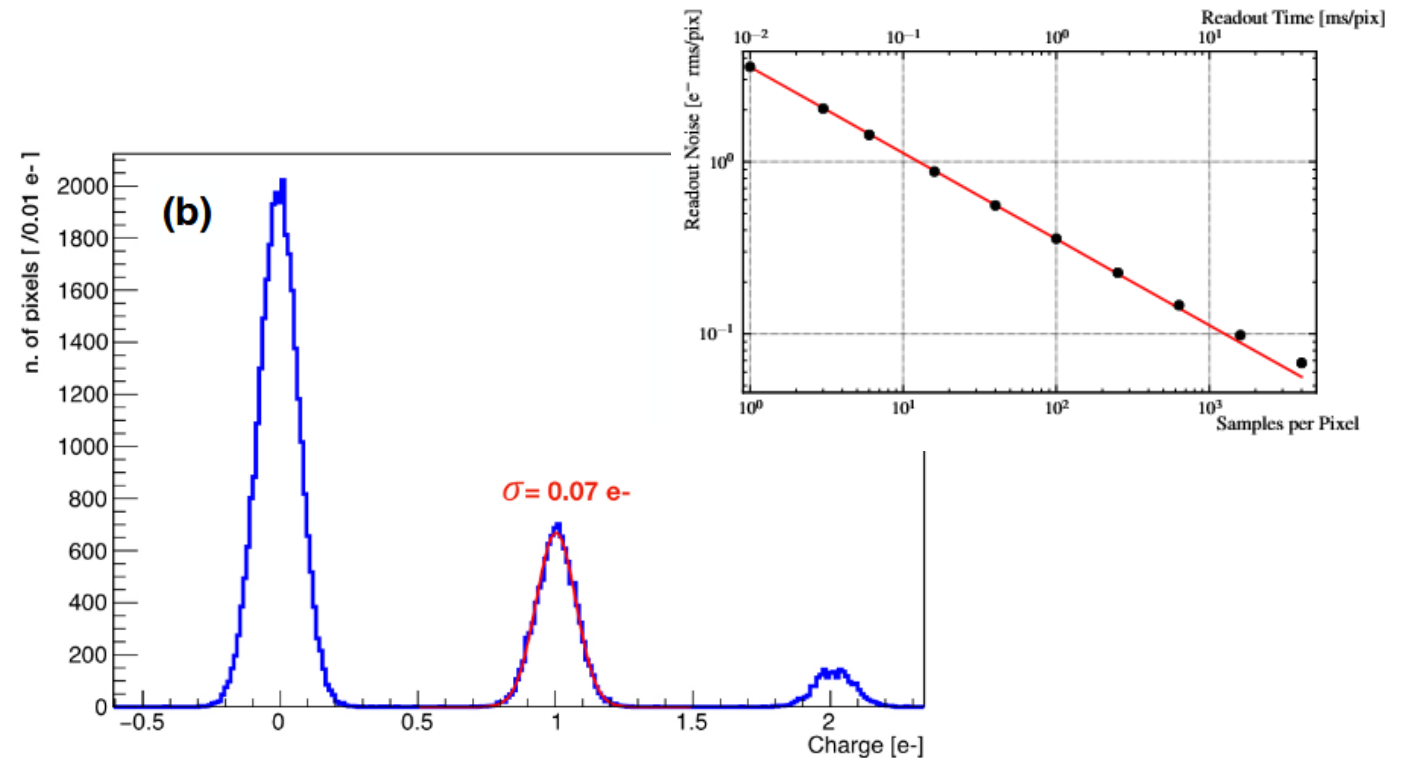
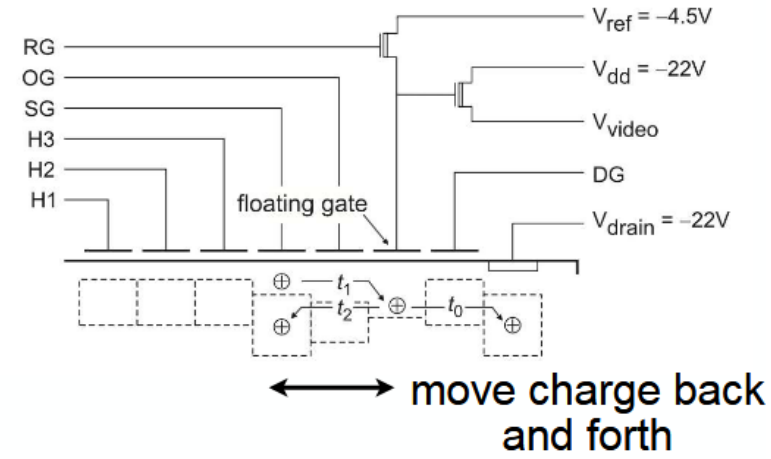
Classical CCDs

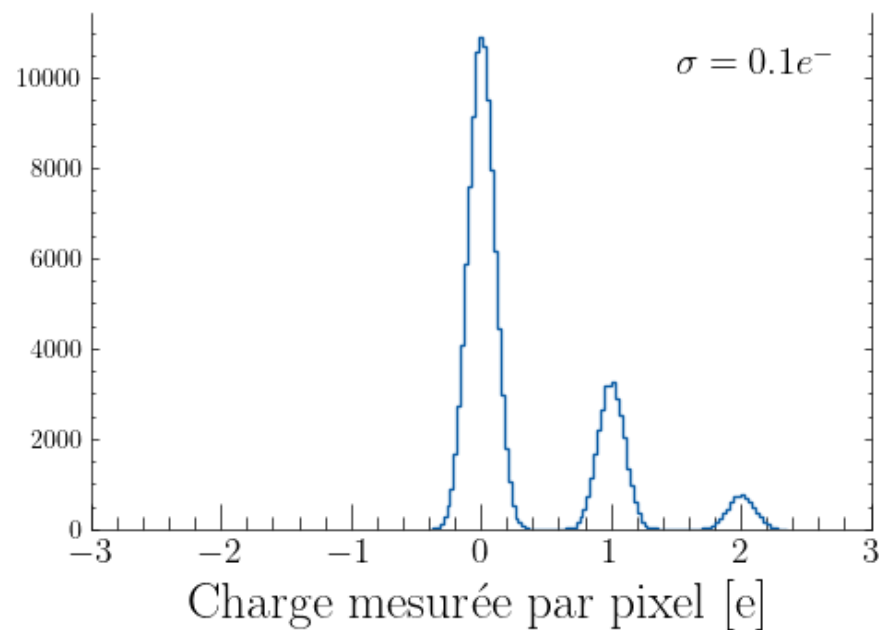
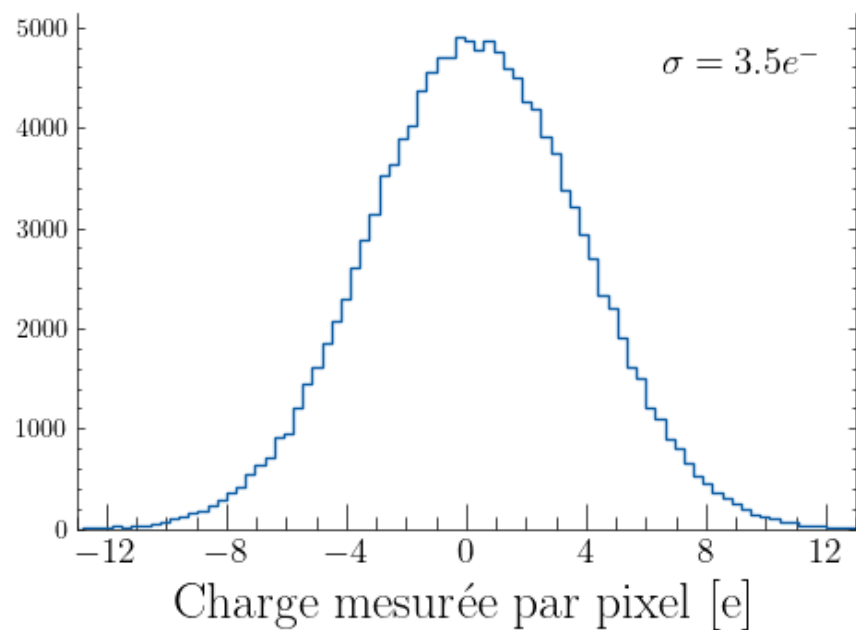
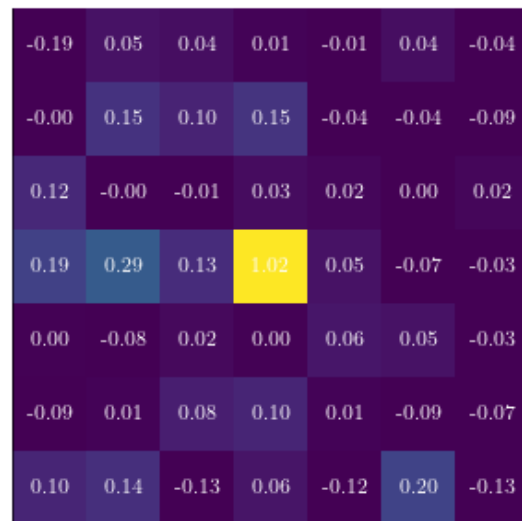
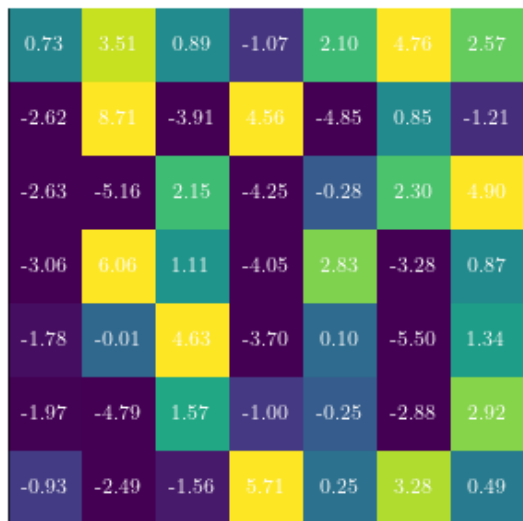
- ❖ interaction ionizes the medium and generates charges
- ❖ Charges are trapped in potential wells
- ❖ Then, the charges are moved line by line, pixel by pixel, toward the serial register
- ❖ Measurement of the charge on each pixel in the sens node



Skipper-CCD

- ❖ Multiple non-destructive charge measurements
 - ❖ Sub-electron reading noise
- ❖ Readout noise divided by $\sqrt{N_S}$
- ❖ Allows us to explore small energy deposits





Clustering and high charge pixels

❖ Clustering

- ❖ We mask two pixels before and after the cluster, and one pixel above and below

❖ High charge pixels

- ❖ We mask two pixels before and after the cluster, and one pixel above and below it

-0.17	-0.11	-0.15	0.09	0.44	0.06	-0.06	0.03	-0.09	0.01	-0.15	0.05
0.16	-0.05	0.11	0.03	0.29	0.10	-0.10	-0.01	-0.01	0.05	-0.14	0.00
0.42	0.01	0.15	0.16	0.22	0.05	0.05	0.34	0.10	0.27	-0.35	0.26
0.03	-0.03	0.02	6.92	48.65	105.17	73.53	16.17	0.71	-0.14	-0.04	-0.19
0.13	0.32	0.00	0.14	-0.09	0.03	0.22	-0.09	0.09	0.06	0.04	-0.16
-0.05	-0.11	0.02	-0.10	0.08	-0.11	-0.12	0.05	-0.06	-0.07	-0.20	0.03
0.28	0.11	-0.17	0.46	-0.12	0.20	0.05	-0.15	-0.28	-0.18	0.12	0.07

-0.17	-0.11	-0.15	0.09	0.44	0.06	-0.06	0.03	-0.09	0.01	-0.15	0.05
0.16	-0.05	0.11	0.03	0.29	0.10	-0.10	-0.01	-0.01	0.05	-0.14	0.00
0.42	0.01	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-0.35	0.26
0.03	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-0.19
0.13	0.32	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	-98.00	0.04	-0.16
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0.16	-0.05	0.11	0.03	0.29	0.10	-0.10	-0.01	-0.01	0.05	-0.14	0.00
0.42	0.01	0.15	0.16	0.22	0.05	0.05	0.34	0.10	0.27	-0.35	0.26
-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00
0.13	0.32	0.00	0.14	-0.09	-95.00	0.22	-0.09	0.09	0.06	0.04	-0.16
-0.05	-0.11	0.02	-0.10	0.08	-95.00	-0.12	0.05	-0.06	-0.07	-0.20	0.03
0.28	0.11	-0.17	0.46	-0.12	-95.00	0.05	-0.15	-0.28	-0.18	0.12	0.07

-0.17	-0.11	-0.15	0.09	0.44	0.06	-0.06	0.03	-0.09	0.01	-0.15	0.05
0.16	-0.05	0.11	0.03	0.29	0.10	-0.10	-0.01	-0.01	0.05	-0.14	0.00
0.42	0.01	-96.00	-96.00	-96.00	-96.00	-96.00	-96.00	-96.00	-96.00	-0.35	0.26
-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00	-95.00
0.13	0.32	-96.00	-96.00	-96.00	-95.00	-96.00	-96.00	-96.00	-96.00	0.04	-0.16
-0.05	-0.11	0.02	-0.10	0.08	-95.00	-0.12	0.05	-0.06	-0.07	-0.20	0.03
0.28	0.11	-0.17	0.46	-0.12	-95.00	0.05	-0.15	-0.28	-0.18	0.12	0.07