Universität Zürich[UZH]

Flavour Deconstructing the Composite Higgs*

16.07.2025, Orsay & Paris

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*[2407.10950] with S. Covone , J. Davighi, G. Isidori & M.P

Motivation

- Higgs is « new » physics ! -> First (seemingly) *elementary* scalar in the SM
- The Higgs is at the *heart* of many SM mysteries: flavour, vacuum stability, naturalness,…

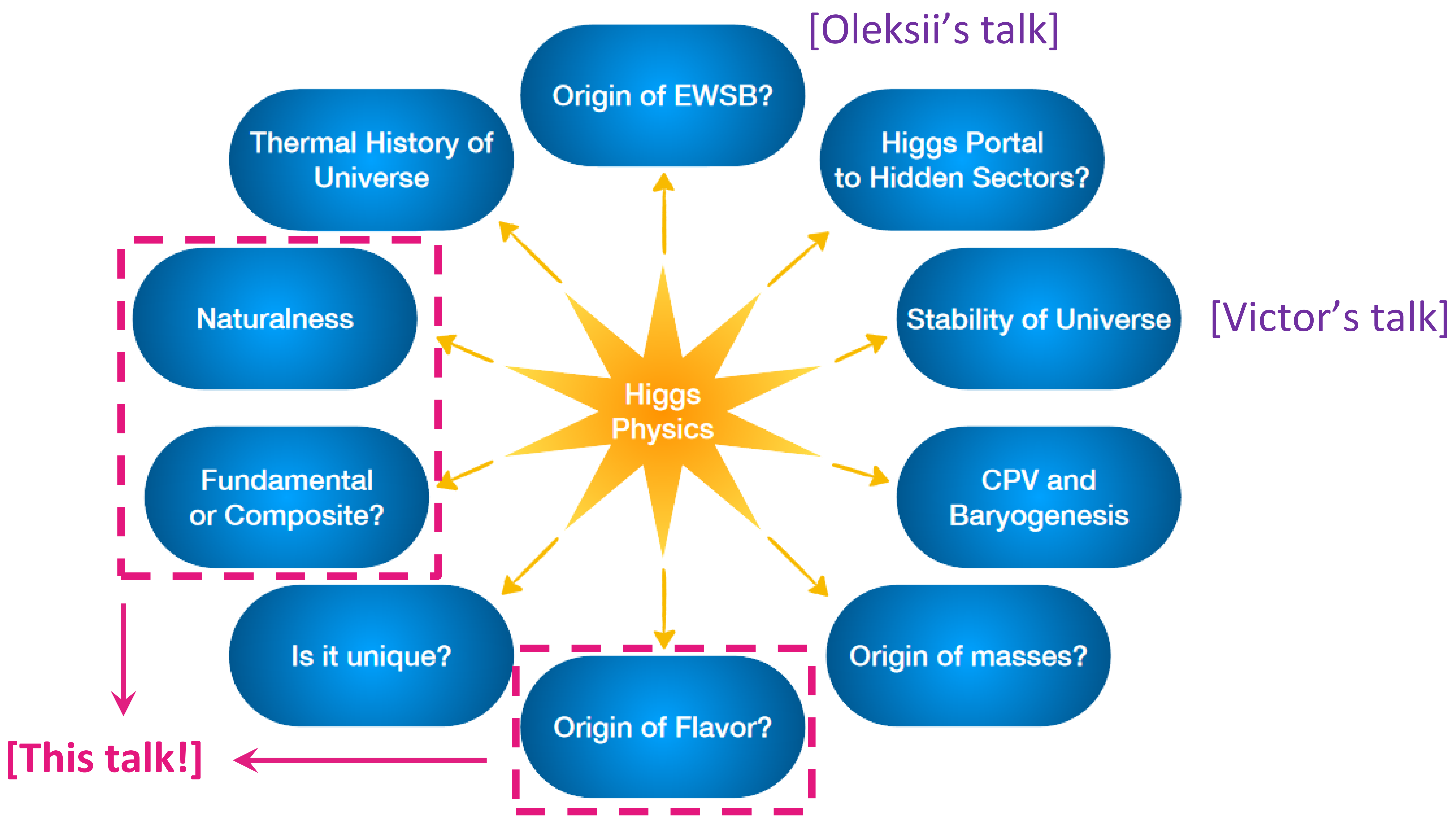
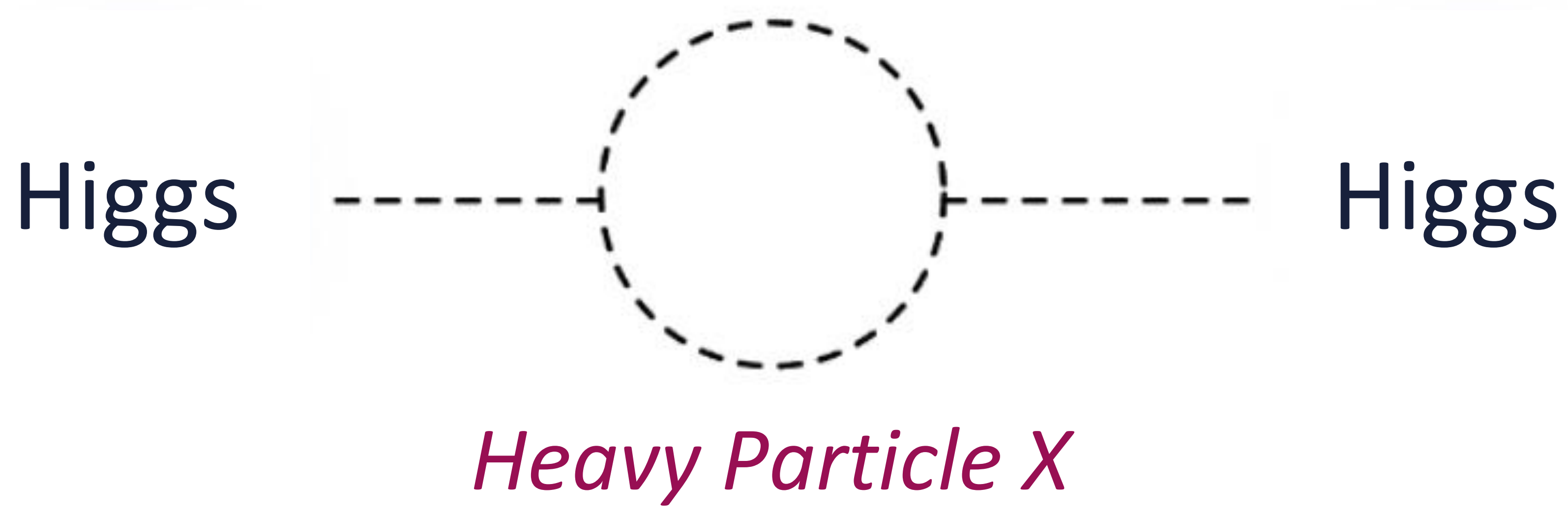


Fig. from [2209.07510]

The Hierarchy Problem

Any heavy NP will *destabilize* the Higgs mass



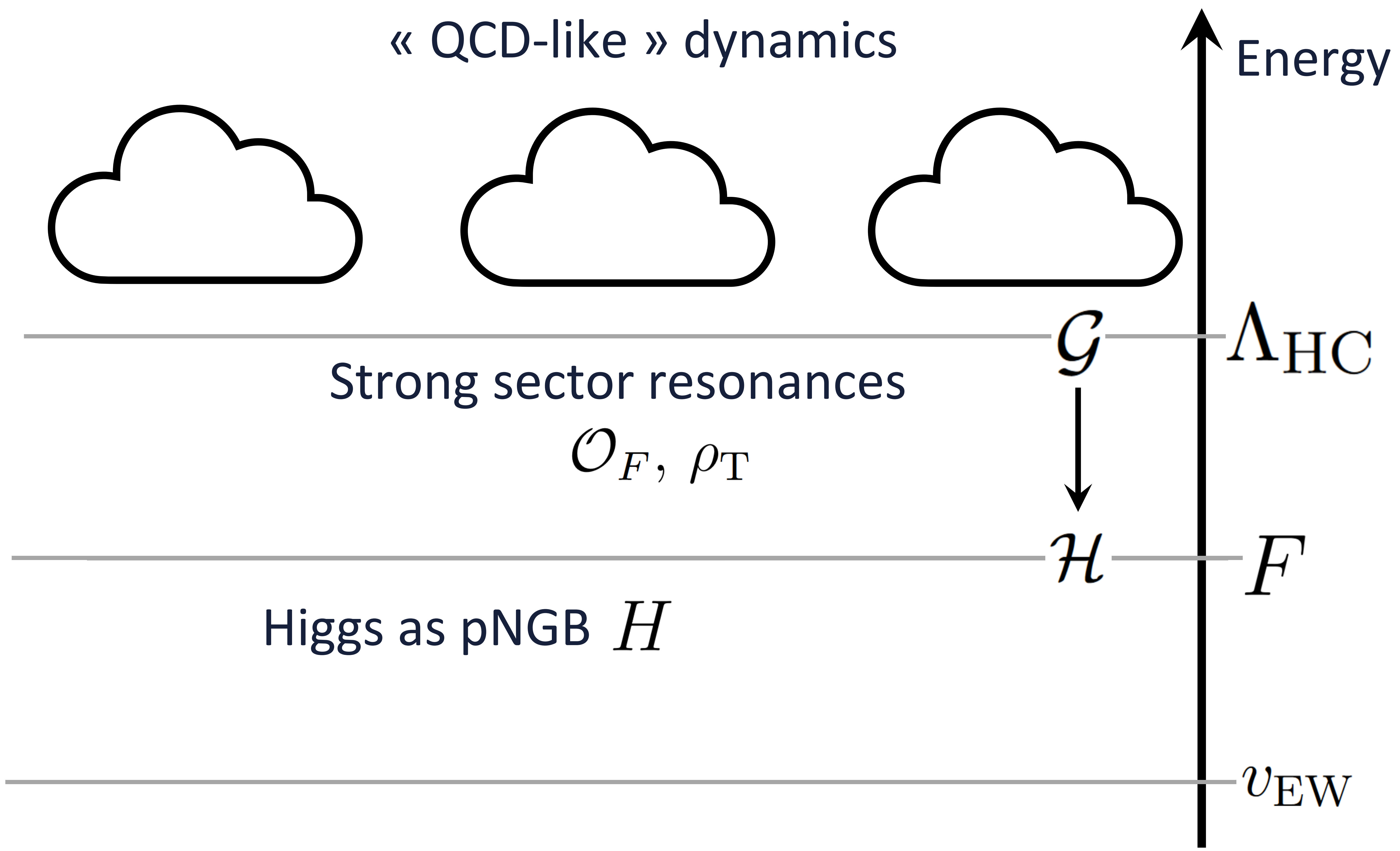
$$\Rightarrow \quad \delta m_H^2 \sim \frac{1}{16\pi^2} g^2 {m_X}^2$$

$$m_H^2 \sim \Lambda_{\rm NP}^2 \quad \textit{vs} \quad m_H = 125\,\mathrm{GeV}$$

➡ Protection mechanism (e.g. SUSY, Composite Higgs) as low scale as possible

Higgs Compositeness

Wulzer & Panico 2015
Agashe et al. 2005,
…



Compositeness scale cuts off quantum corrections to the Higgs potential *(like pions in QCD !)*

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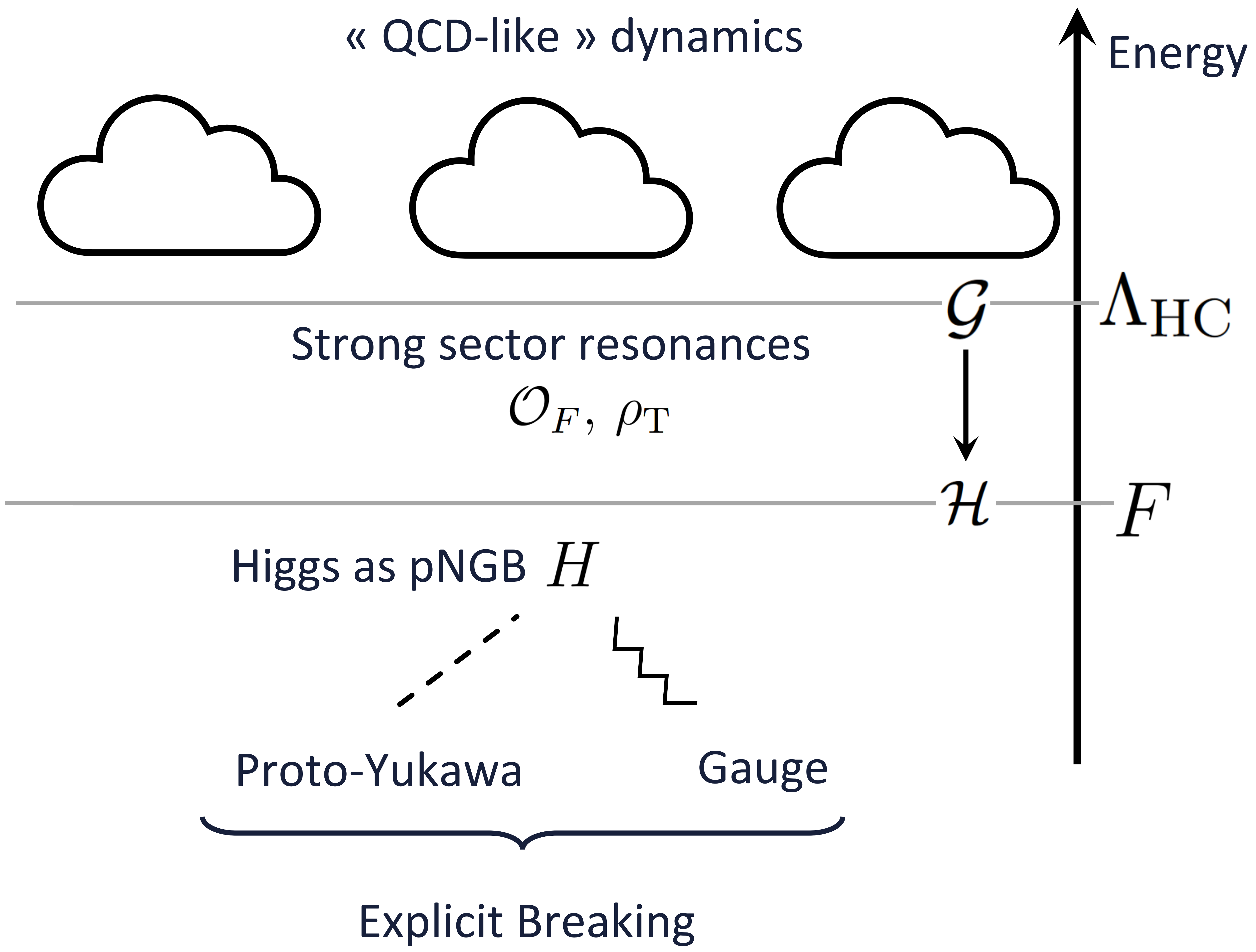

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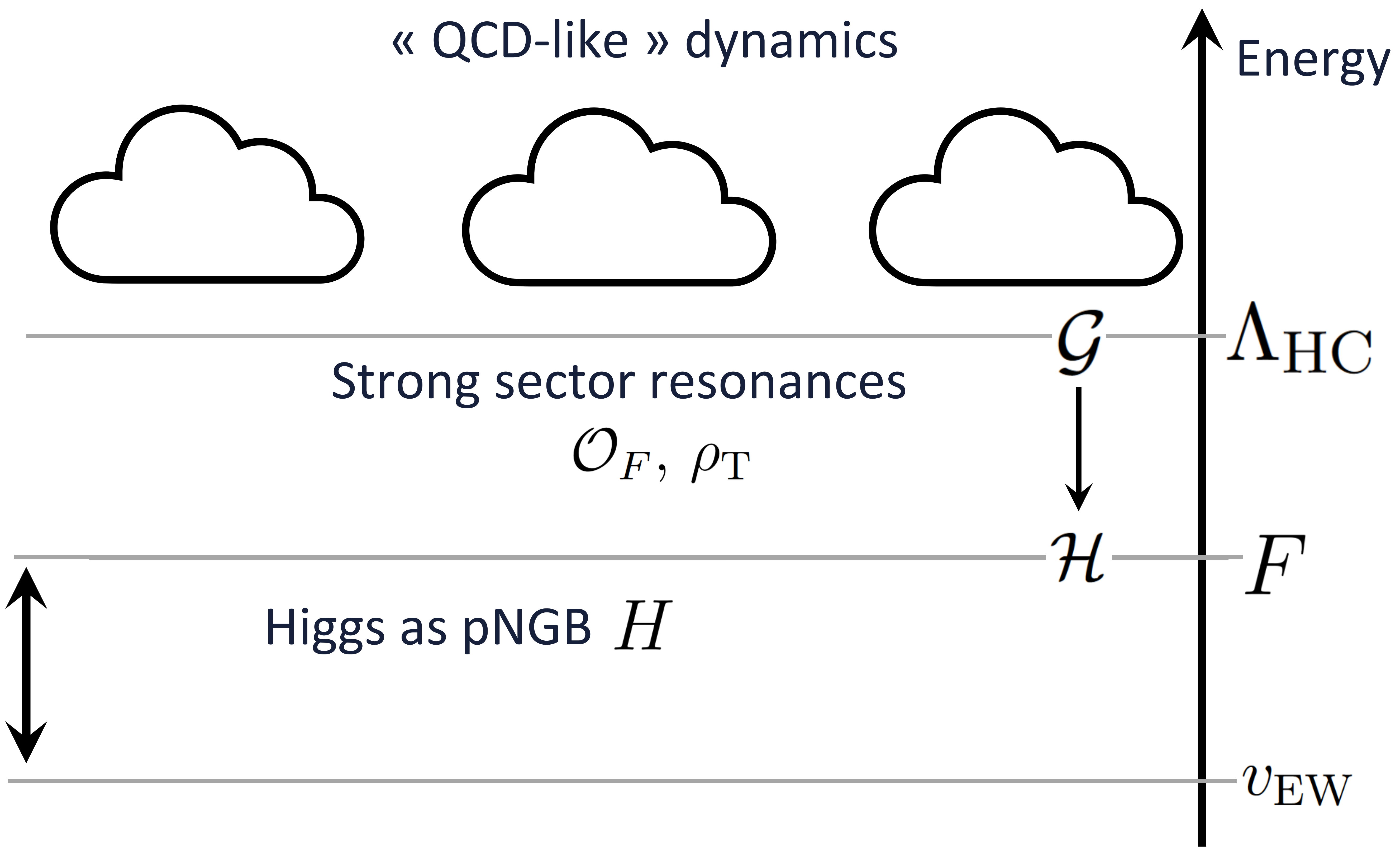

Higgs is SM-like to good approx !

$$\xi \equiv \frac{v_{\rm EW}^2}{F^2} \lesssim 0.06$$

Compositeness corrects SM predictions

$$g_{VVh} = g_{VVh}^{\rm SM}\sqrt{1-\xi}$$

$$g_{VVhh} = g_{VVhh}^{\rm SM}\,(1-2\xi)$$

Measuring Higgs couplings is important to probe CHM!
[many exp. talks @HH2025]

And flavour ?

Natural Higgs -> low scale NP … what about flavour ?

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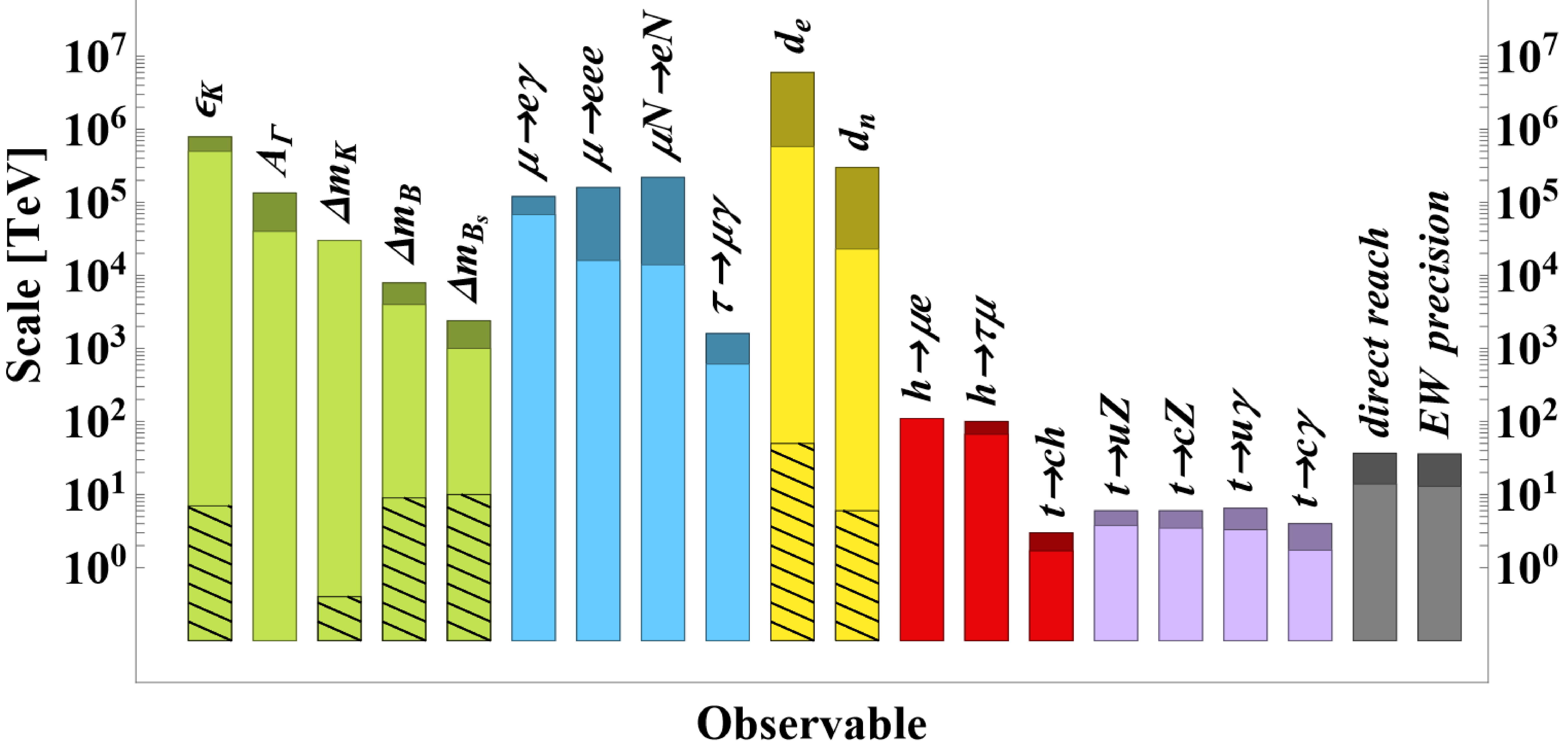
➡ While the hierarchy problem points to scale $M \sim$ TeV, flavour points to much *higher scales*!

Physics Briefing Book 2020, 1910.11775



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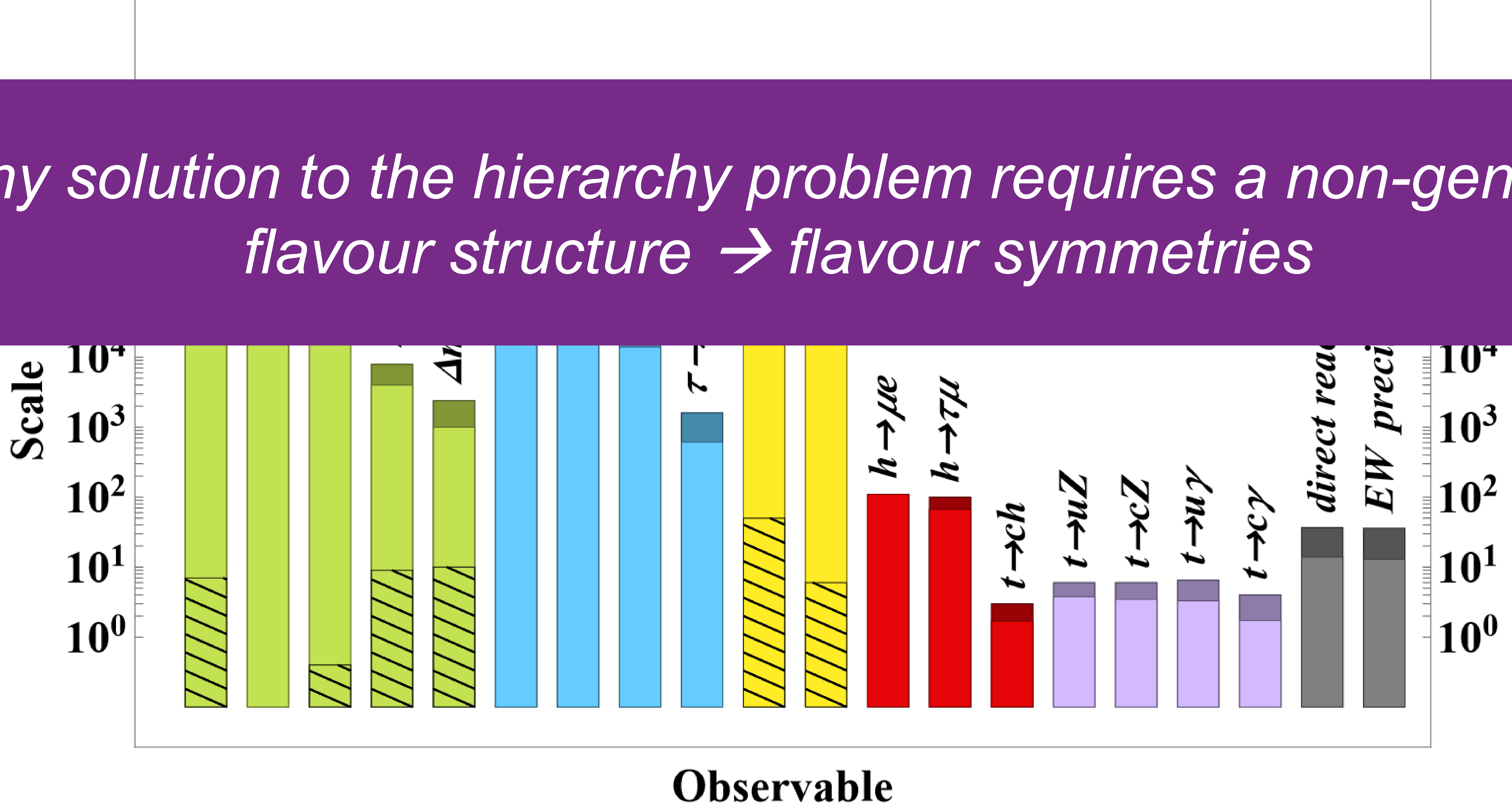

Any solution to the hierarchy problem requires a non-generic flavour structure → flavour symmetries

Flavour Non-Universality

Lessons from SM & EXP:

- Exact $U(3)^5$ flavour symmetry in the gauge and fermion sectors of the SM
- Peculiar breaking $U(3)^5 \rightarrow U(2)^n$ with only $y_t \sim \mathcal{O}(1)$ --> $Y_u \sim \begin{pmatrix} & < 0.01 & 0.04 \\ & & 1 \end{pmatrix}$
- No large breaking of $U(2)$ @TeV & stringent flavour bounds on light families

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$G_{12} \times G_{3+H}$ * -> G_{SM}

Flavour symmetries *encoded* in the gauge!

*Different options to Flavour deconstruct: Davighi & Isidori [2303.01520]

Exploring the Flavor Symmetry Landscape [2402.09503]
-> Flavour symmetries to reduce scale of CHM

Flavour Deconstructing the Composite Higgs [2407.10950]

$$SU(3)_c \times Sp(4) \times U(1)^{[3]}_{B-L} \times U(1)^{[12]}_Y$$

Non-perturbative Dynamics

Λ_{HC} $\downarrow$ $H \sim (\mathbf{2},\mathbf{2})$

$$SU(3)_c \times SU(2)_L \times SU(2)^{[3]}_R \times U(1)^{[3]}_{B-L} \times U(1)^{[12]}_Y$$

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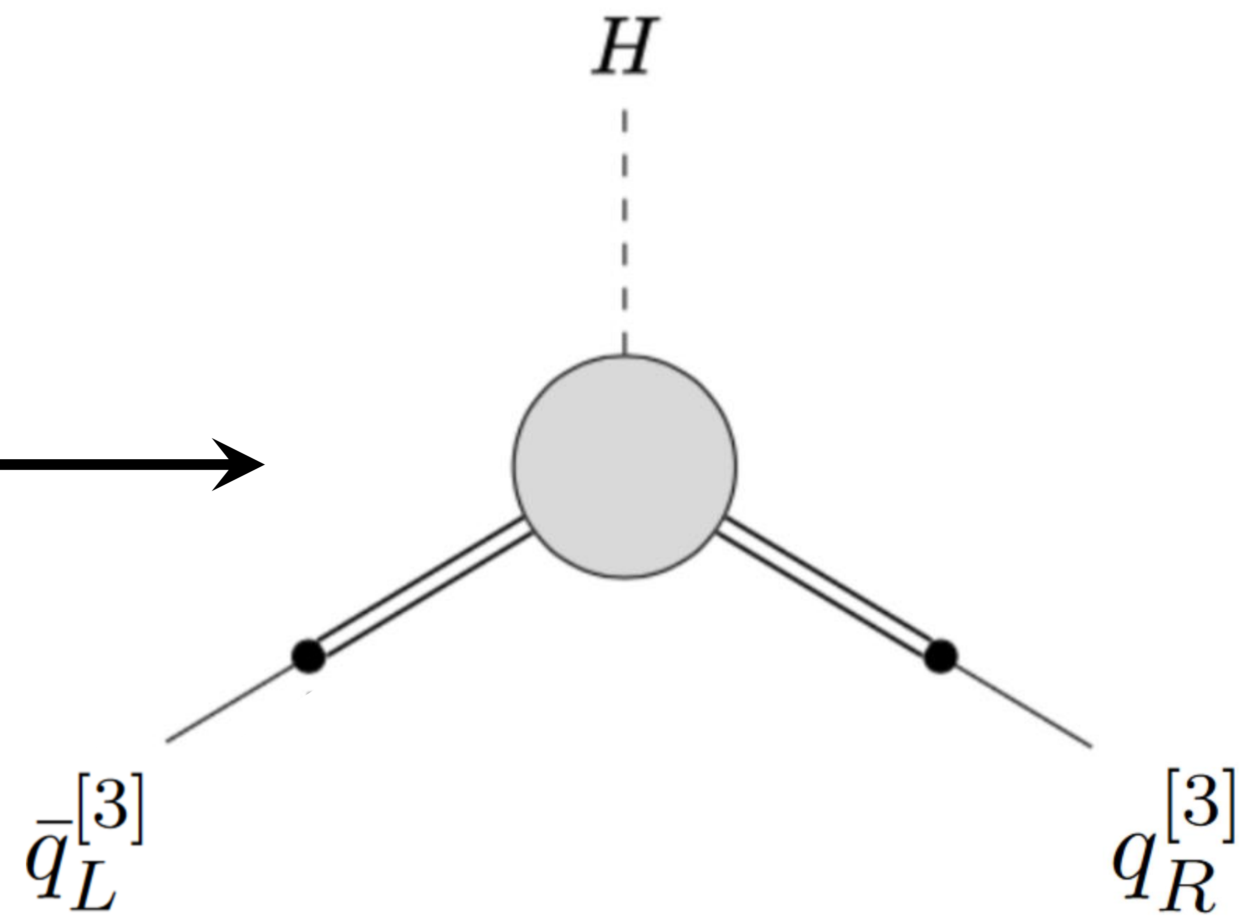
Third-family Partial Compositeness:

$$\mathcal{L} \supset \lambda_L \bar{q}^{[3]}_L \mathcal{O}_R \longrightarrow$$



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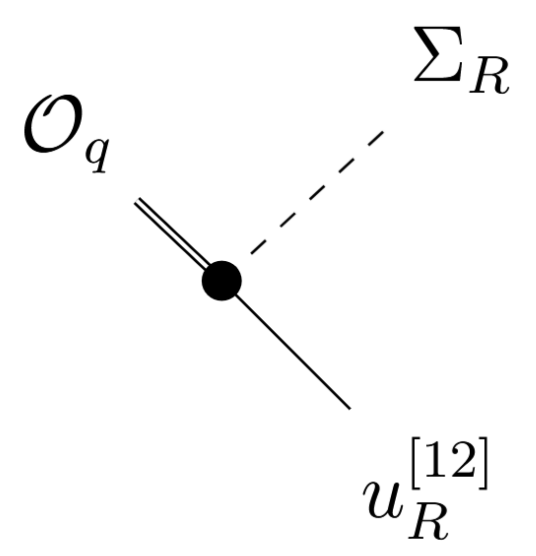

Flavour for light fam.

Horizontal Breaking

$\langle \Sigma_R \rangle$ $\quad \langle \Omega \rangle$ *

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

Two Z′ and a pair of $W_R^\pm$

(Most of the pheno!)

- $B \to X_s \gamma$
- Bound on Z-pole obs.
- Bounds on Z′ masses from FCNC
- LHC bound from Drell-Yan data

** elementary scalars here, w.i.p. to embed them as composite states*

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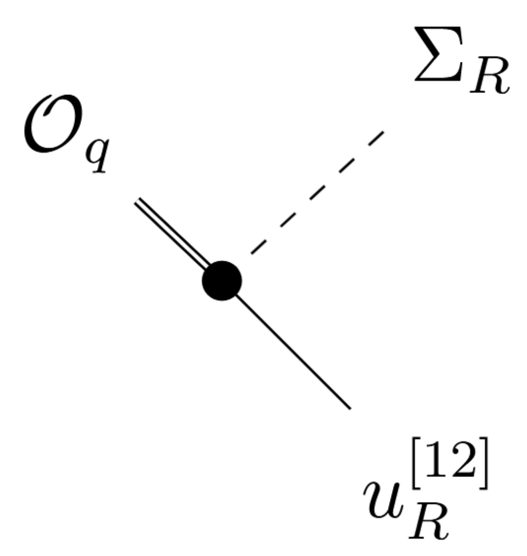

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(Most of the pheno!)

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

Parameter space motivated by naturalness:

- $g_{R,3} \approx O(1) >> g_{Y,12}$
- $M^2_{W_R} \sim g^2_{R,3} v^2_\Sigma \lesssim M^2_\rho$

** elementary scalars here, w.i.p. to embed them as composite states*

Flavour Deconstructing the Composite Higgs [2407.10950]

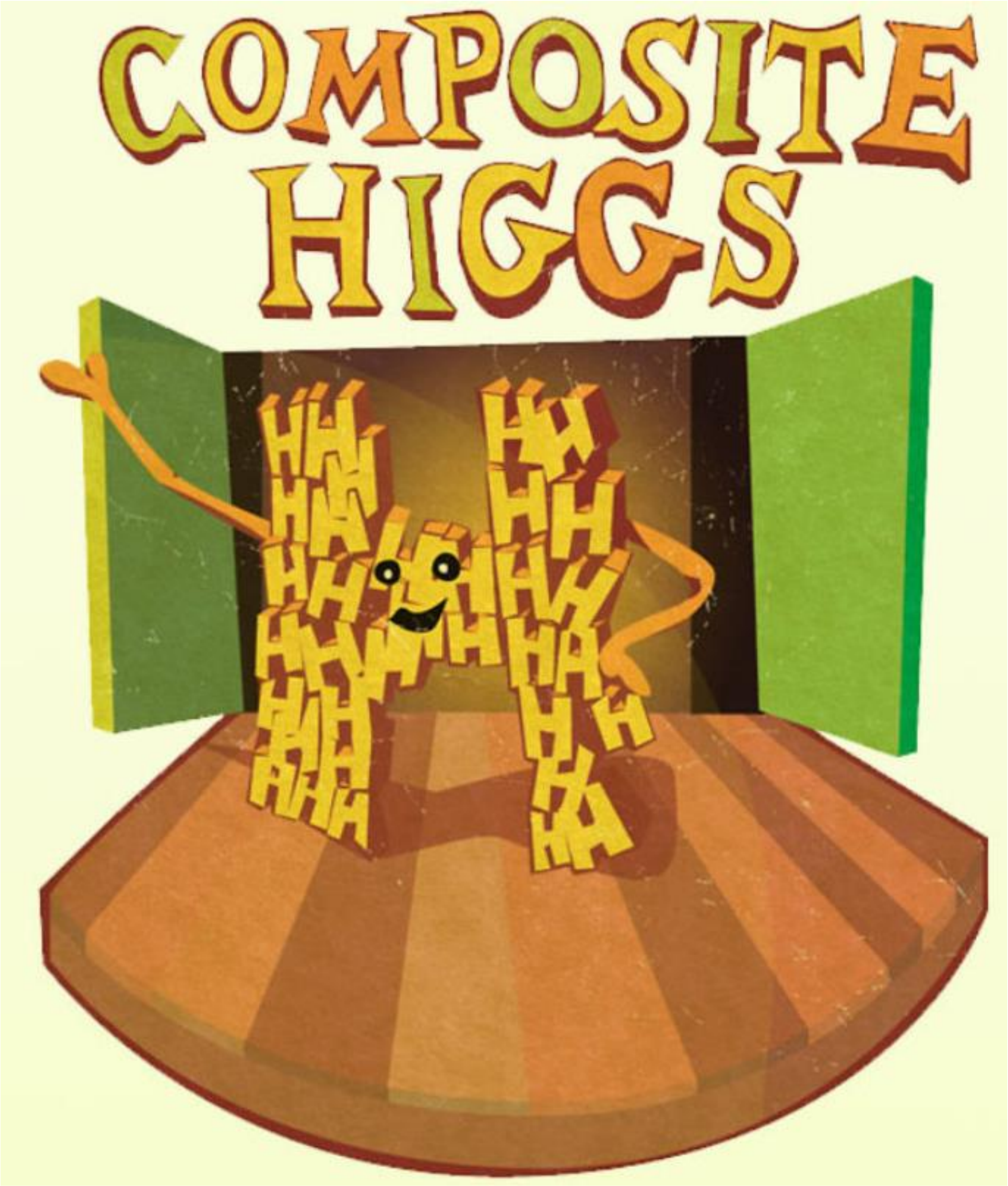
Typical scenario

- Large 3rd gen. RH gauge coupling: $g_{R,3} = O(1)$
- Light Top partner $M_T \approx 2$ TeV and $M_\rho \approx 10$ TeV
- Scale of flavour deconstruction $v_\Sigma \approx 3$ TeV

} → Minimize the tuning in the potential
→ $O(1\%)$ corrections to Higgs couplings

➡ *Pheno. viable* TeV-scale model to stabilize the Higgs and address the flavour puzzle + provide testable signatures at current and near future colliders !

Thank You !



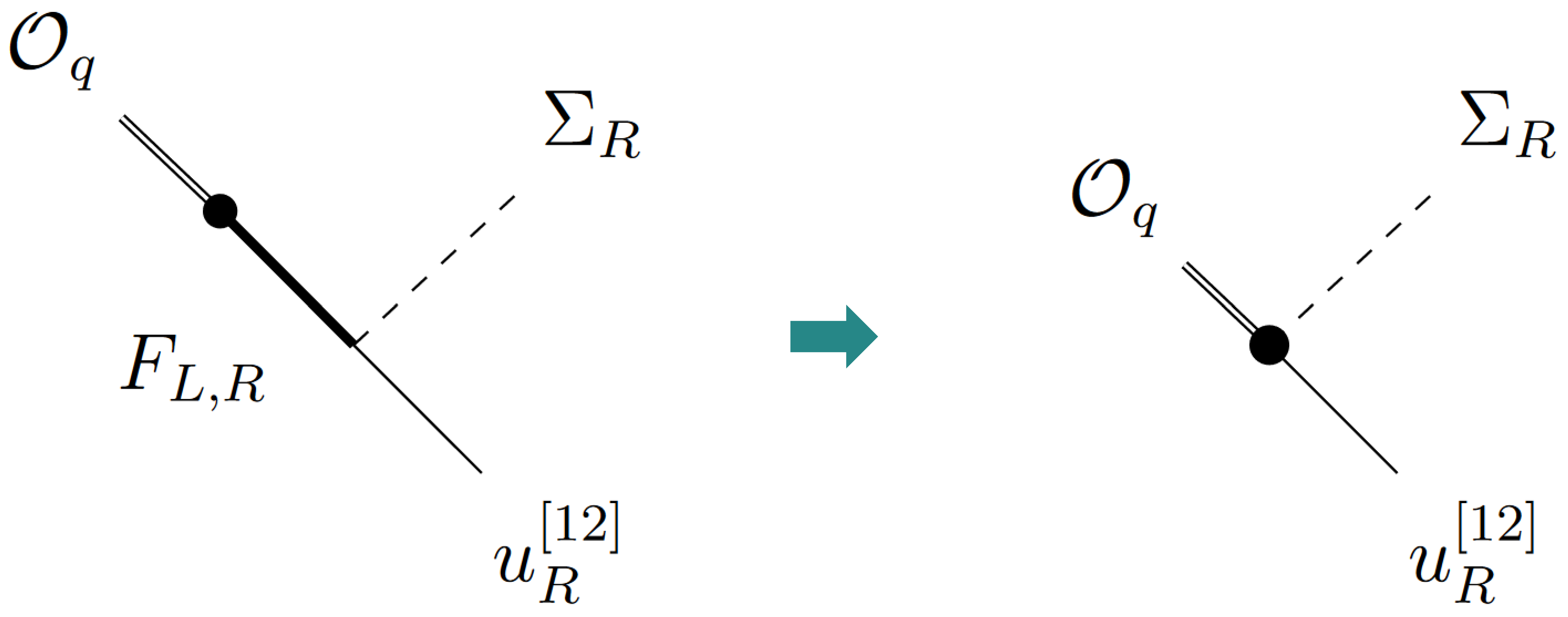
Backup Slides

Flavour Deconstructing the Composite Higgs

$$SU(3)_c \times SU(2)_L \times SU(2)_R^{[3]} \times U(1)_{B-L}^{[3]} \times U(1)_Y^{[12]}$$

Horizontal Breaking $\langle \Sigma_R \rangle$ $\langle \Omega \rangle$

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$



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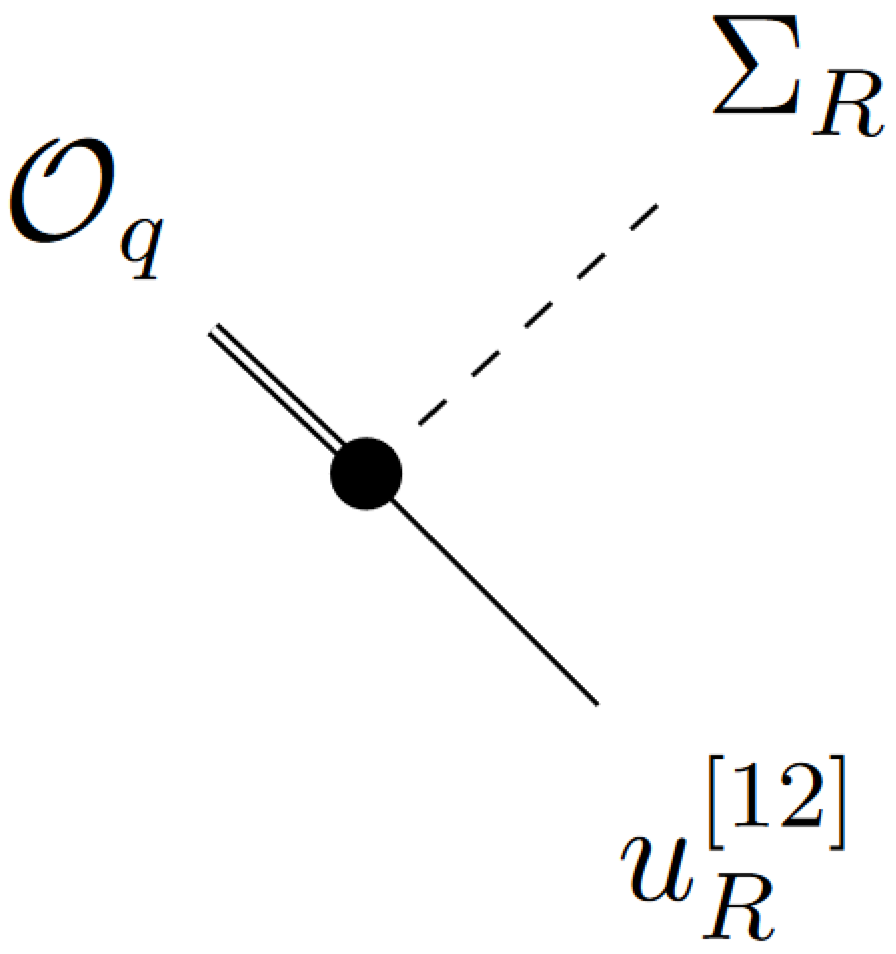
Horizontal Breaking: $\langle \Sigma_R \rangle$ $\langle \Omega \rangle$

$$\downarrow$$

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$



$$\epsilon_R = \frac{\langle \Sigma \rangle}{M_F}$$

$$\epsilon_\Omega = \frac{\langle \Omega \rangle}{M_F}$$

$$\longrightarrow \quad Y_{u,d,e} \sim \begin{pmatrix} \epsilon_R & \epsilon_\Omega \\ \epsilon_R \epsilon_\Omega & 1 \end{pmatrix}$$

Flavour Deconstructing the Composite Higgs

$$Y_{u,d,e} \sim \begin{pmatrix} \epsilon_R & \epsilon_\Omega \\ \epsilon_R \epsilon_\Omega & 1 \end{pmatrix}$$

$$\epsilon_\Omega = O\left(|V_{cb}|\right) = O\left(10^{-1}\right)$$

$$\epsilon_R = O\left(m_c/m_t\right) = O\left(10^{-2}\right)$$

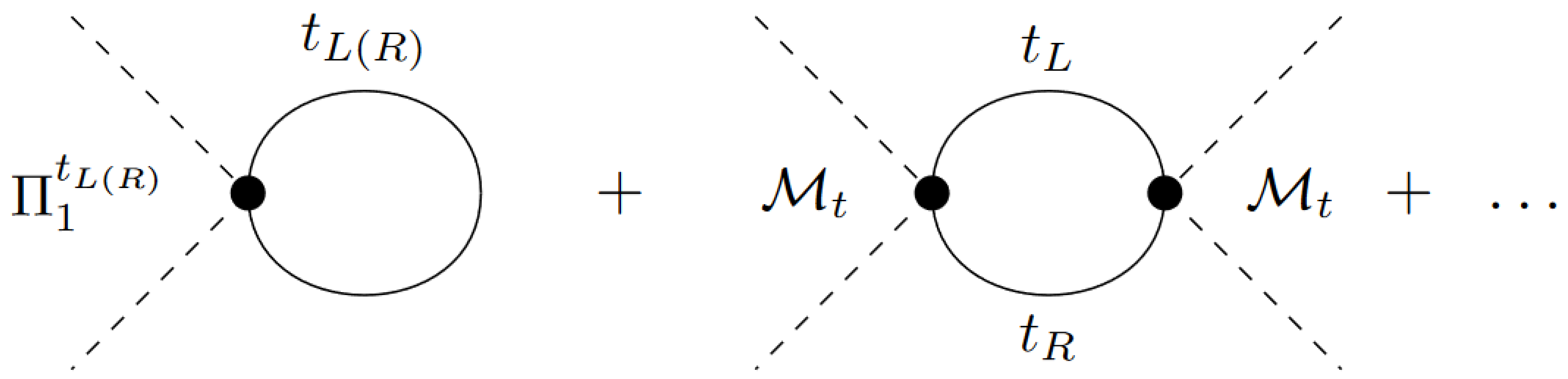
⟶ The deconstruction scale is ***anchored*** by its impact on the Higgs potential

Higgs Potential

➡ Higgs potential induced at 1-loop

Fermion contribution

$\Delta V(h)_f$



$$\mathcal{L}_{\text{eff}} \supset \bar{q}_L \not{p} \left[\Pi_0^{q_L}(p^2)\mathbb{1} + \Pi_1^{t_L}(p^2) u_L^\dagger \Delta_+ u_L \right] q_L + \left\{ \bar{q}_L \left[\mathcal{M}_t(p^2) u_L^\dagger \Delta_+ u_R \right] q_R + \text{h.c.} \right\}$$

$$\Pi_1^{t_L}(0) = \frac{F^2}{M_T^2} \left(\lambda_L^t\right)^2 \kappa_L^t$$

$$|\mathcal{M}_t(0)| = y_t \sqrt{2} F$$

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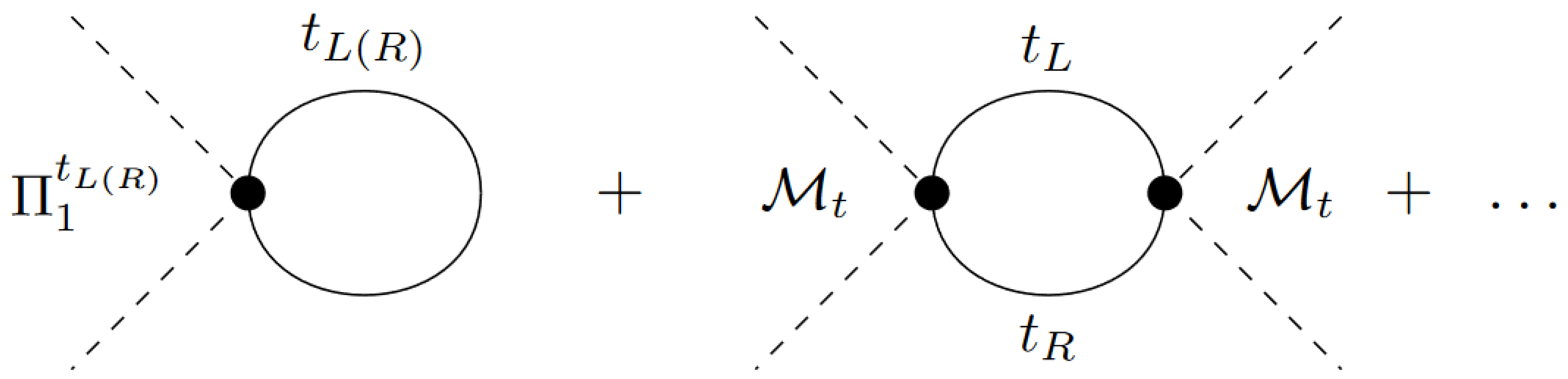


Coleman-Weinberg potential ⟶

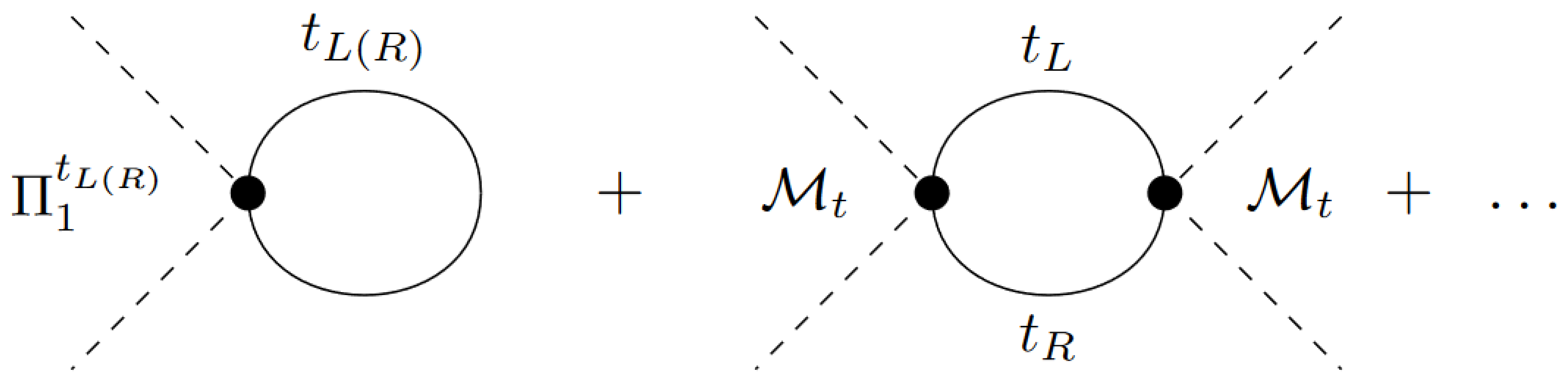
$$\Delta V(h)_f = -2N_c \int \frac{d^4p_E}{(2\pi)^4} \left\{ \log\left[1 + \frac{\Pi_1^{t_L}}{\Pi_0^{q_L}} \sin^2\left(\frac{h}{2F}\right)\right] + \log\left[1 + \frac{|\mathcal{M}_t|^2 \sin^2\left(\frac{h}{2F}\right)\cos^2\left(\frac{h}{2F}\right)}{p_E^2\left(\Pi_0^{q_L} + \Pi_1^{q_L}\sin^2\left(\frac{h}{2F}\right)\right)\left(\Pi_0^{q_R} - \Pi_1^{q_R}\sin^2\left(\frac{h}{2F}\right)\right)}\right]\right\}$$

Higgs Potential

Higgs potential induced at 1-loop

Fermion contribution

$\Delta V(h)_f$



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$$\mathcal{M}_t(q^2) = \mathcal{M}_t(0) \times \frac{M_T^2}{M_T^2 - q^2}$$

$$\frac{\Pi_1^{t_L}(q^2)}{\Pi_1^{t_L}(0)} \frac{\Pi_0^{q_L}(0)}{\Pi_0^{q_L}(q^2)} = \frac{M_T^2}{M_T^2 - q^2} \frac{M_f^2}{M_f^2 - q^2}$$

Explicit expression in terms of model parameters

Flavour Deconstructing the Composite Higgs [2407.10950]

$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

$$\frac{c_2}{F^4} = \frac{N_c y_t^2}{4\pi^2}\frac{M_T^2}{F^2} +$$ Gauge contributions (suppressed)

↑ Top partner → $M_T \approx 2.5F$

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Fermionic resonances — Top partner — Gauge contributions

$$\frac{c_1}{F^4} = \frac{N_c}{8\pi^2}\left[\left(\lambda_R^t\right)^2 \kappa_R^t - \left(\lambda_L^t\right)^2 \kappa_L^t\right]\frac{M_f^2}{F^2} + \frac{N_c y_t^2}{4\pi^2}\frac{M_T^2}{F^2} - \frac{9 g_R^2}{32\pi^2}\left(1 - \frac{g_R^2 v_\Sigma^2}{2M_\rho^2}\right)\frac{M_\rho^2}{F^2} + \mathcal{O}\left(g_L g_R, g_L^2\right)$$

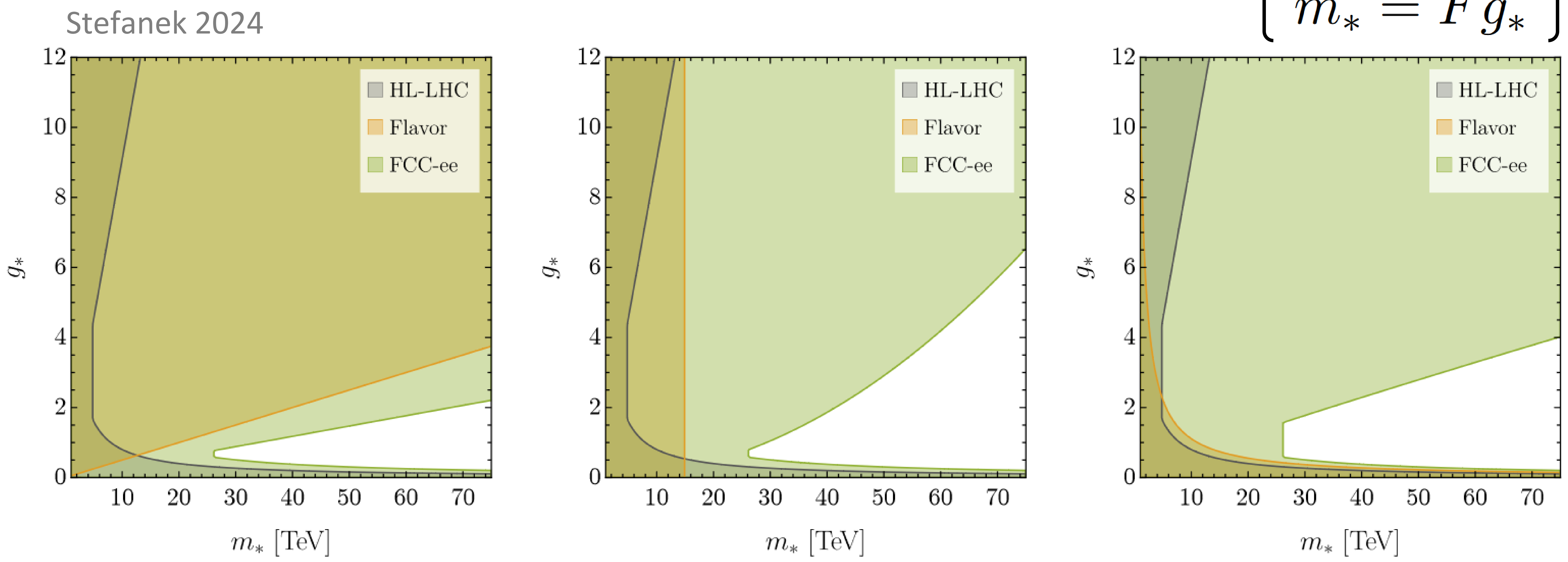
- Increase size of gauge contribution ⟶ $g_{R,3} = O(1) \gg g_{R,12} \approx g_Y^{\mathrm{SM}}$

  (Natural in flavour non-universal scenario !)

- Avoid suppression / sign flip ⟶ $M_{W_R}^2 = \frac{1}{4} g_R^2 v_\Sigma^2 < \frac{1}{2} M_\rho^2$

Composite Higgs @ HL-LHC and FCC-ee

$[\, m_* = F g_* \,]$



(a) Left compositeness (b) Mixed compositeness (c) Right compositeness

➡ With improved precision: RG-running into EWPO become **crucial**

➡ Composite Higgs will be put under a microscope @ FCC-ee!

$$m_* \gtrsim 25 \text{ TeV}$$

Flavour Non-Universal Composite Higgs

Ingredients:

- Spontaneously broken strong sector: $\mathcal{G} \equiv Sp(4) \xrightarrow{\Lambda_{\rm HC}} SU(2)_L \times SU(2)_R^{[3]} \equiv \mathcal{H}$

- Field Content:

Elementary fields		$U(1)_{B-L}^{[3]}$	$U(1)_Y^{[12]}$	$SU(2)_L$	$SU(2)_R^{[3]}$
chiral light quarks	$q_L^{[12]}$	0	1/6	$\mathbf{2}$	$\mathbf{1}$
	$u_R^{[12]}$	0	2/3	$\mathbf{1}$	$\mathbf{1}$
	$d_R^{[12]}$	0	−1/3	$\mathbf{1}$	$\mathbf{1}$
chiral $3^{\rm rd}$ gen. quarks	$q_L^{[3]}$	1/6	0	$\mathbf{2}$	$\mathbf{1}$
	$q_R^{[3]}$	1/6	0	$\mathbf{1}$	$\mathbf{2}$
vector-like quarks	F_L^q	1/6	0	$\mathbf{2}$	$\mathbf{1}$
	F_R^q	0	1/6	$\mathbf{1}$	$\mathbf{2}$
scalar link fields	Σ_R	0	1/2	$\mathbf{1}$	$\mathbf{2}$
	Ω_q	−1/6	1/6	$\mathbf{1}$	$\mathbf{1}$
	Ω_ℓ	1/2	−1/2	$\mathbf{1}$	$\mathbf{1}$

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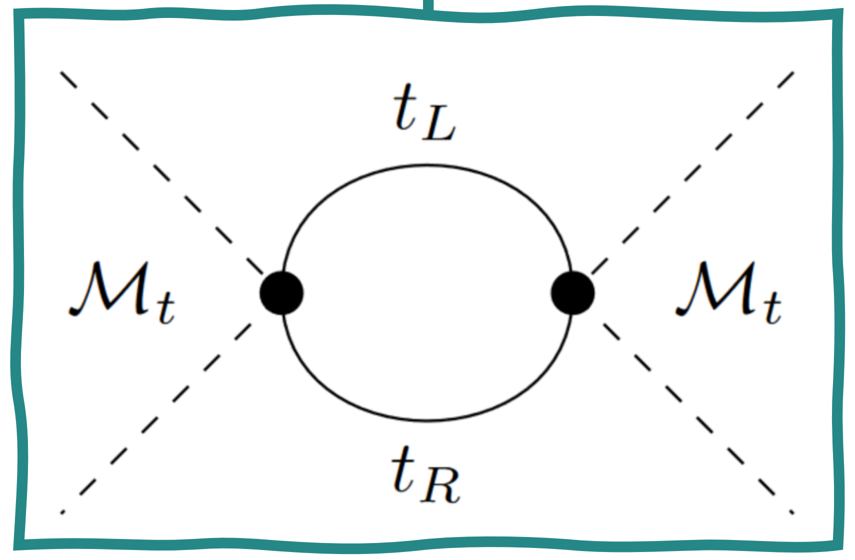
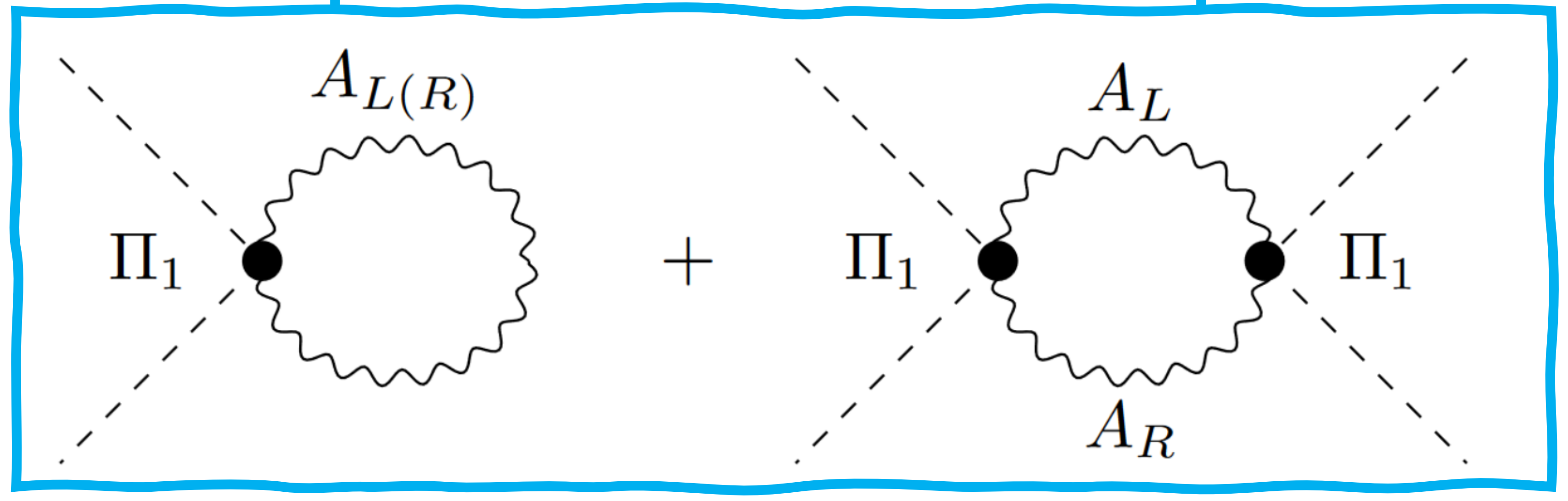
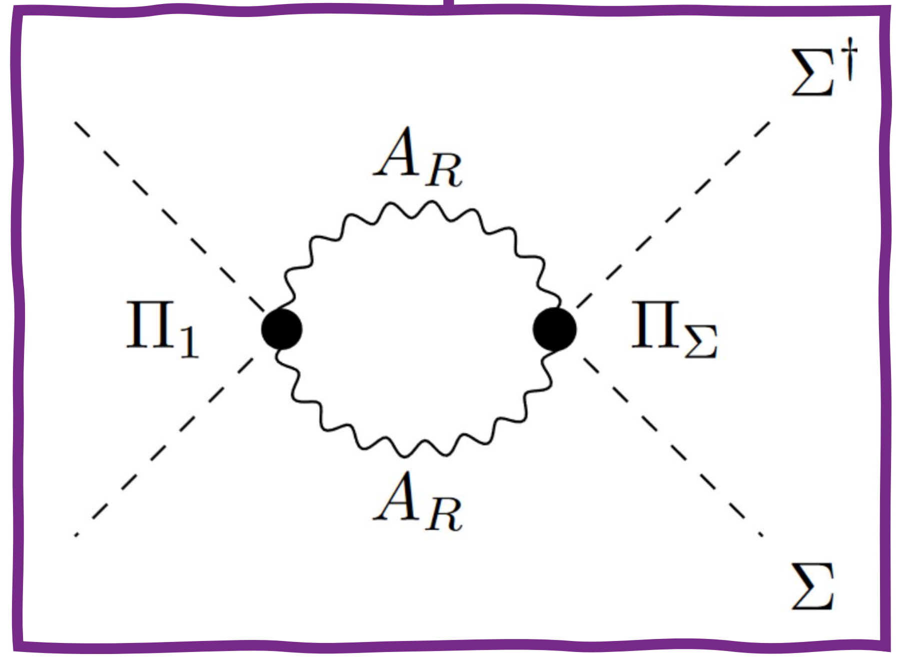
$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

Gauge contributions

$$\frac{c_2}{F^4} = \frac{N_c y_t^2}{4\pi^2}\frac{M_T^2}{F^2} + \frac{9g_R^2}{32\pi^2}\delta_\pi\left(1 - \frac{g_R^2 v_\Sigma^2}{2M_\rho^2}\right)\frac{M_\rho^2}{F^2} - \frac{9g_R^4}{64\pi^2}\log\left(\frac{M_\rho^2}{M_{W_R}^2}\right) + \mathcal{O}(g_L g_R, g_L^2)$$







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Light Top partner

Negligible for $g_R \lesssim 2$

$\delta_\pi \ll 1$ from spurion analysis $\longrightarrow$ $M_T \approx 2.5F$

Flavour and Higgs Compositeness

How do we couple fermions to the Higgs ?

OG approach in strongly-coupled EWSB models:

Strong resonances
-> Naturalness

$$\mathcal{L} \supset \frac{\lambda_b}{\Lambda_{\rm UV}^{d-1}} \bar{q}_L \mathcal{O}_S b_R \longrightarrow y_{t,b} \simeq \lambda_{t,b} \left(\frac{F}{\Lambda_{\rm UV}} \right)^{d-1}$$

Scalar Op.
from the
strong sector

Strong interactions are
resolved
-> Flavour

- Difficult to have $y_t \sim \mathcal{O}(1)$ and $\Lambda_{\rm UV}$ high enough to avoid extra flavour-violation
- Reintroduces the Hierarchy problem for $\mathcal{O}_S^2$
- Enforce hierarchy of $\lambda_{t,b}$ in the UV because only one scalar op $\mathcal{O}_S$

Wulzer & Panico 2015
Agashe et al. 2005,
Rattazzi et al. 2008

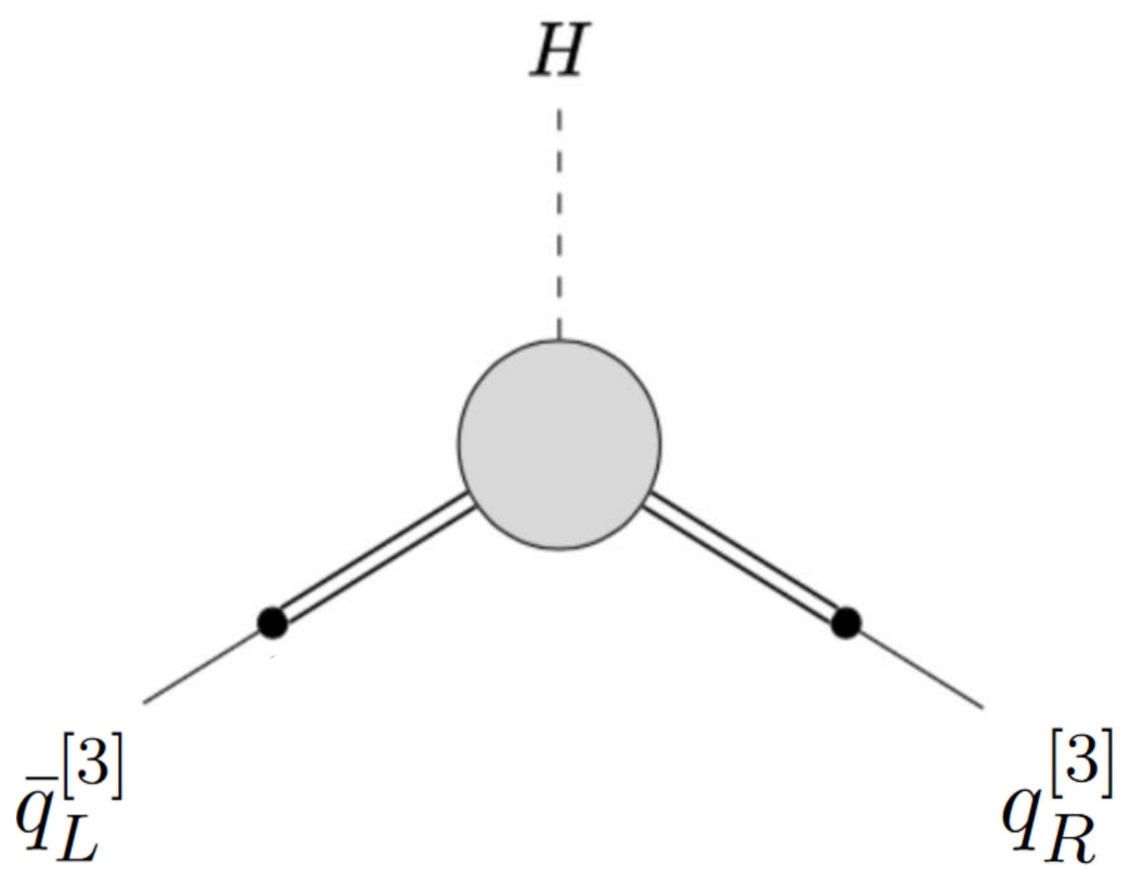
Partial Compositeness

➡ How do we couple fermions to the Higgs ?

Partial Compositeness:

$$\mathcal{L} \supset \lambda_L \bar{q}_L \mathcal{O}_R \longrightarrow |y_q| = \lambda_L^q \lambda_R^{q*} \kappa_{LR}^q \frac{F}{\sqrt{2} M_q}$$

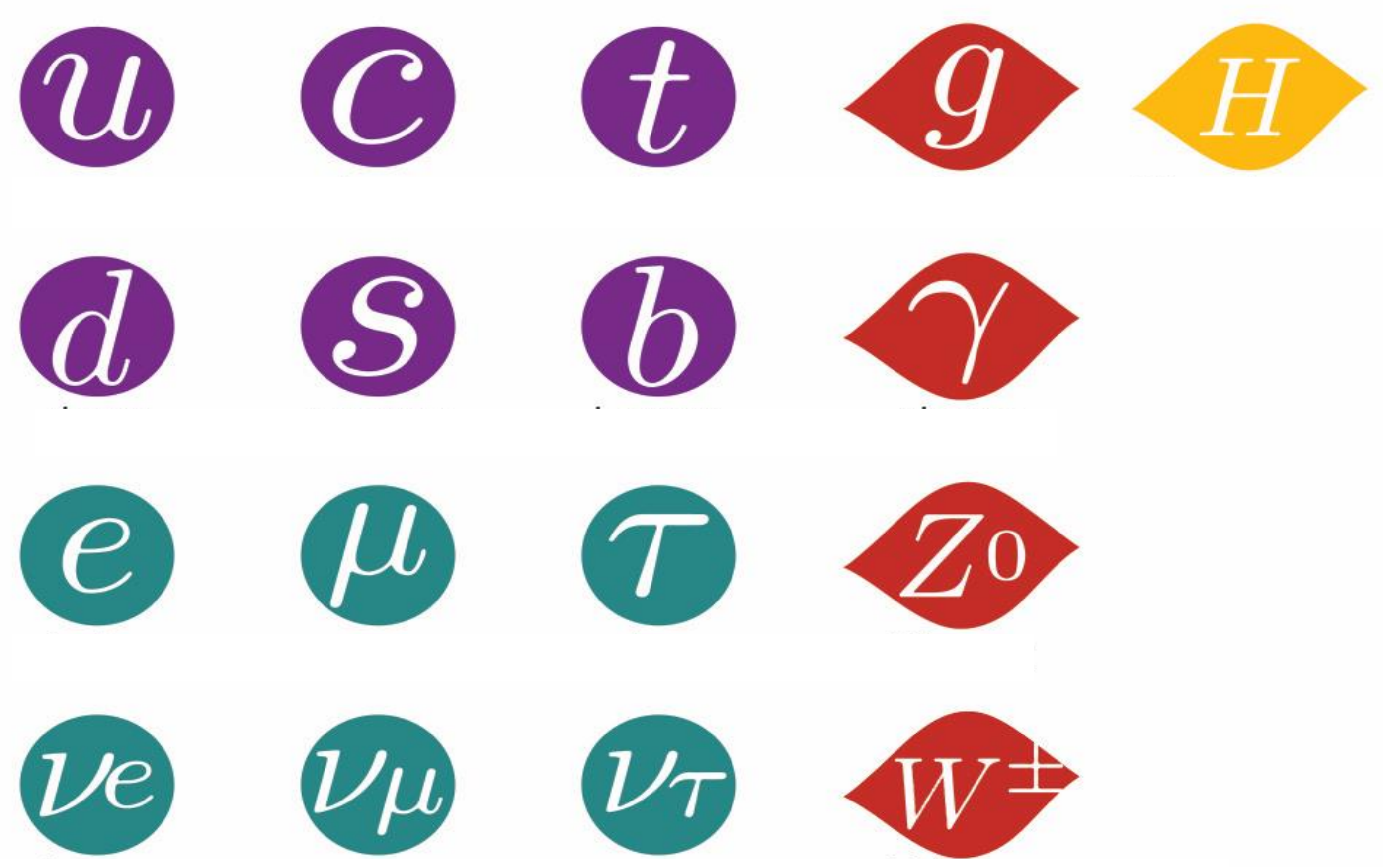
Composite partner

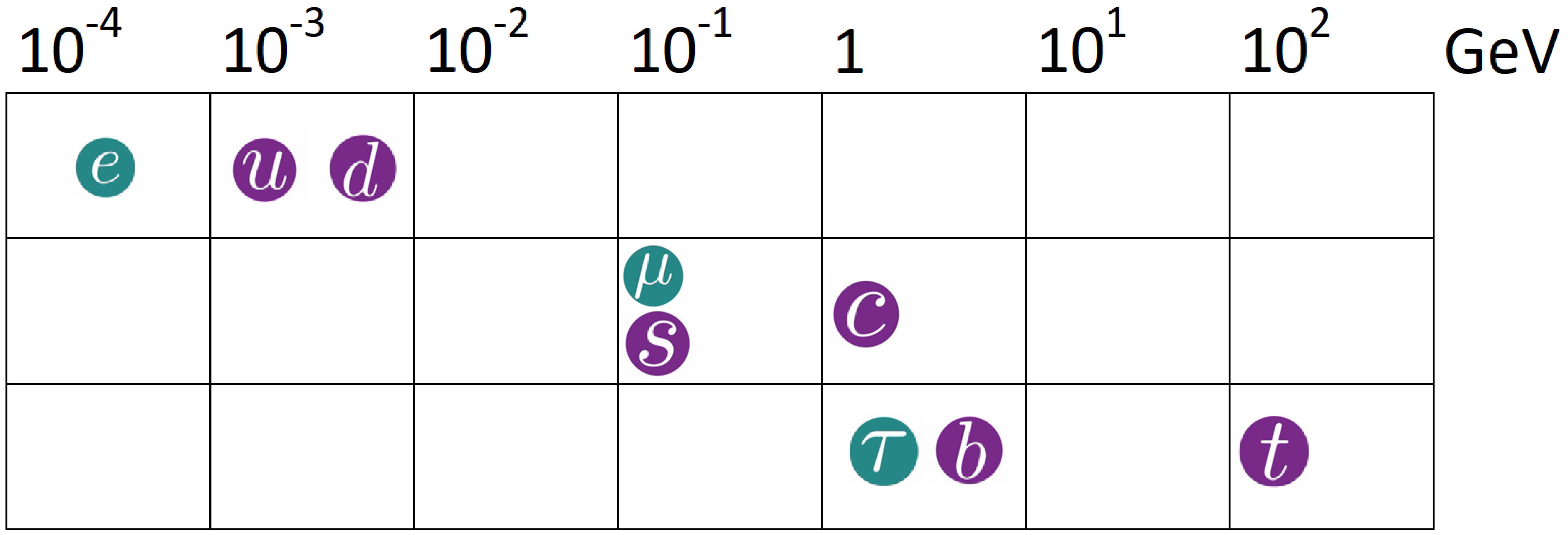


- Fermionic Ops -> No risk of reintroducing a hierarchy problem for $\mathcal{O}_F^2$
- Partners for each fermions -> can reproduce Yukawa pattern

Kaplan 1991
Wulzer & Panico 2015

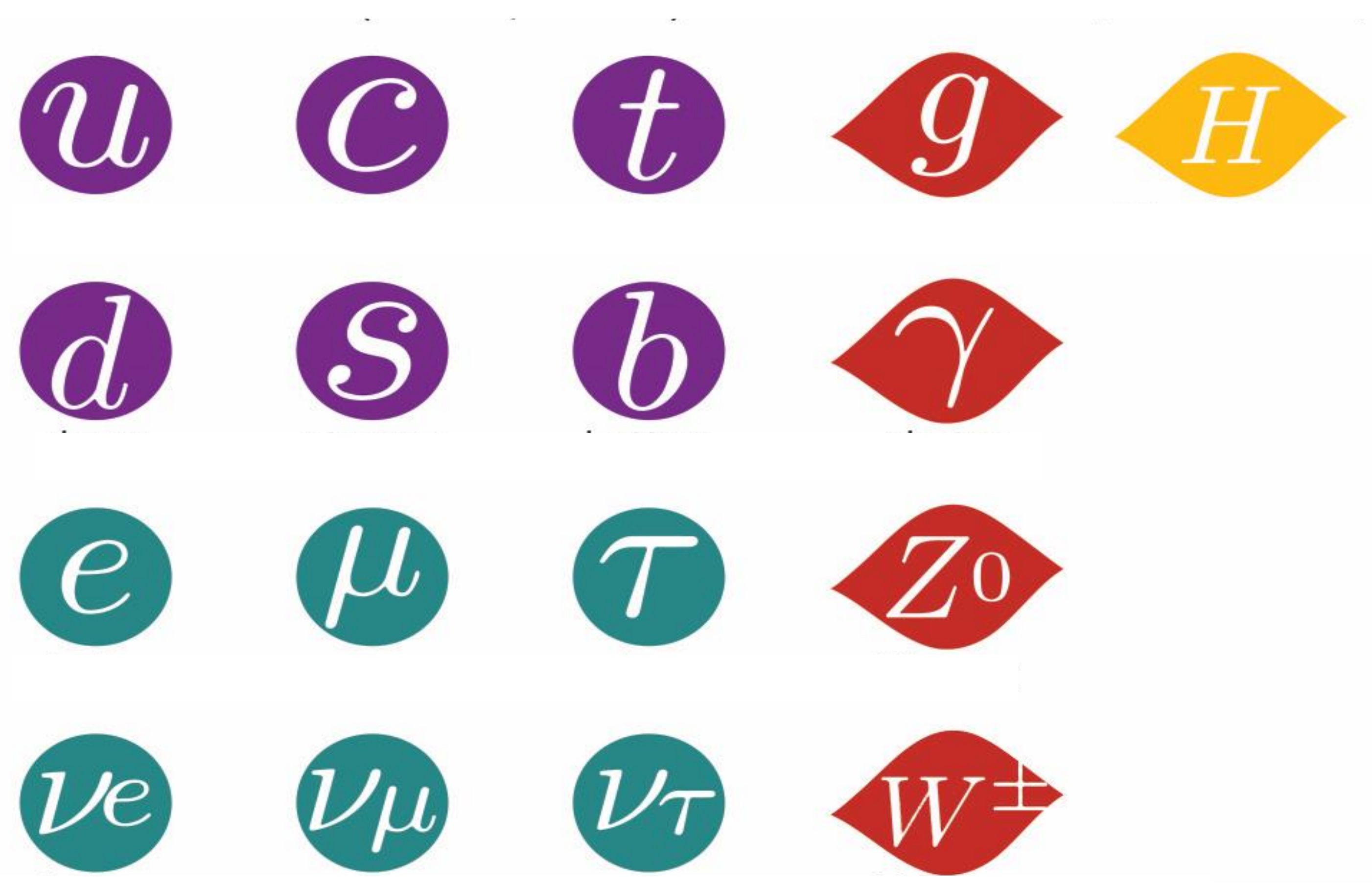
Flavour Puzzle

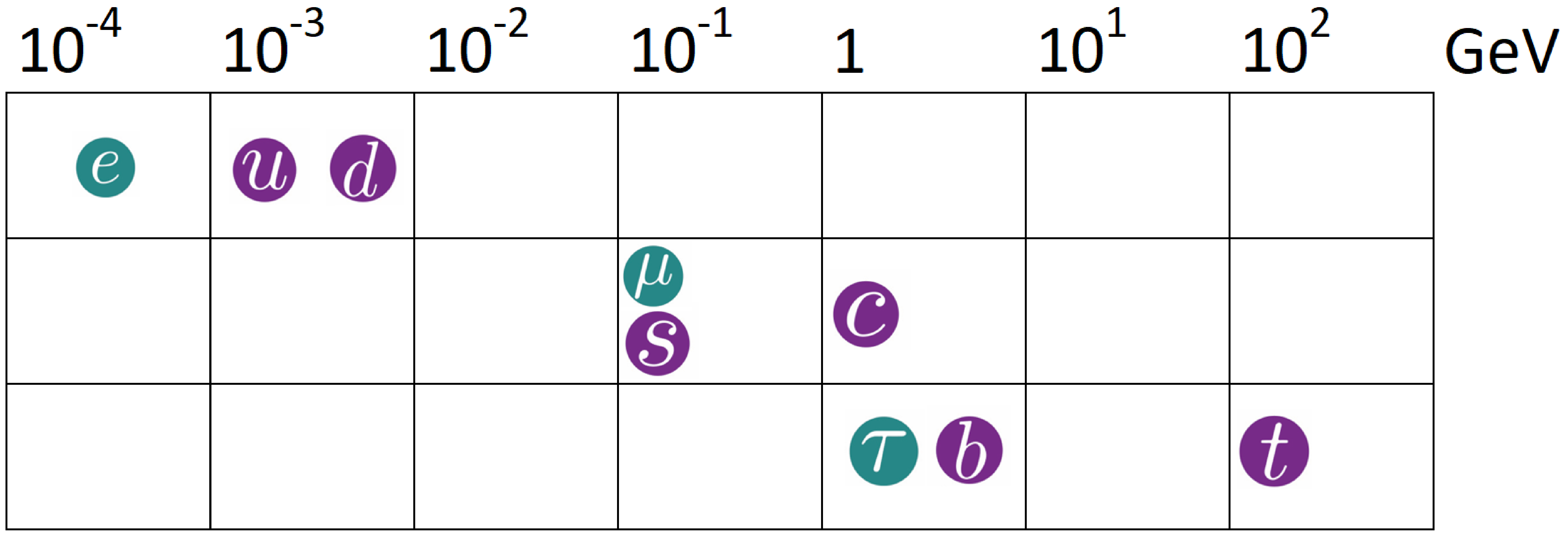




[See Gino's & Barbieri's talk]

Flavour Puzzle





[See Gino's & Barbieri's talk]

$$\mathcal{L}^{d\leq 4} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}}$$

Flavour-universal — Yukawa hierarchies

$$Y_u \sim \begin{pmatrix} & < 0.01 & 0.04 \\ & & 1 \end{pmatrix}$$

In the basis where Y_d is diagonal

$\Rightarrow U(2)^n$ approx. flavour symmetries

Barbieri et al. 2011, Isidori & Straub 2012,
Kagan et al. 2009, Blankenburg et al. 2012

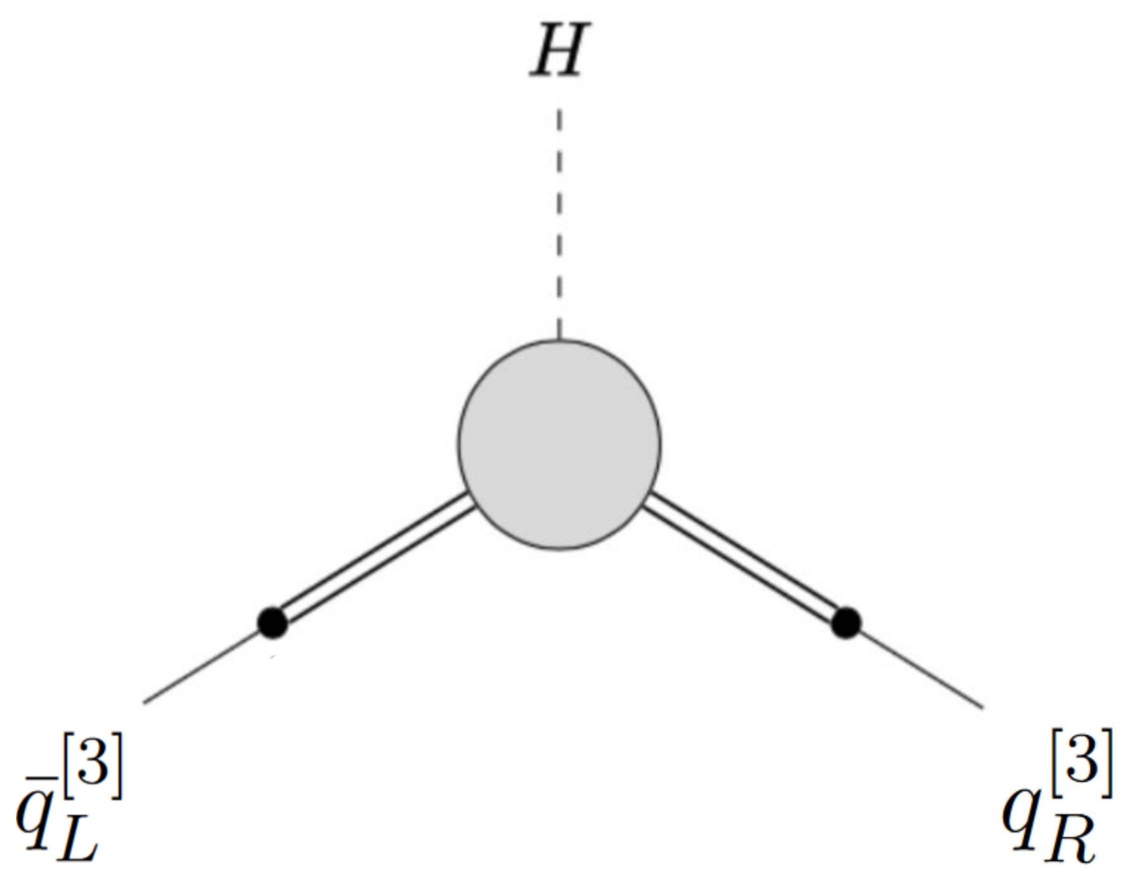
Partial Compositeness

➡ How do we couple fermions to the Higgs ?

Partial Compositeness:

$$\mathcal{L} \supset \lambda_L \bar{q}_L \mathcal{O}_R \longrightarrow |y_q| = \lambda_L^q \lambda_R^{q*} \kappa_{LR}^q \frac{F}{\sqrt{2} M_q}$$

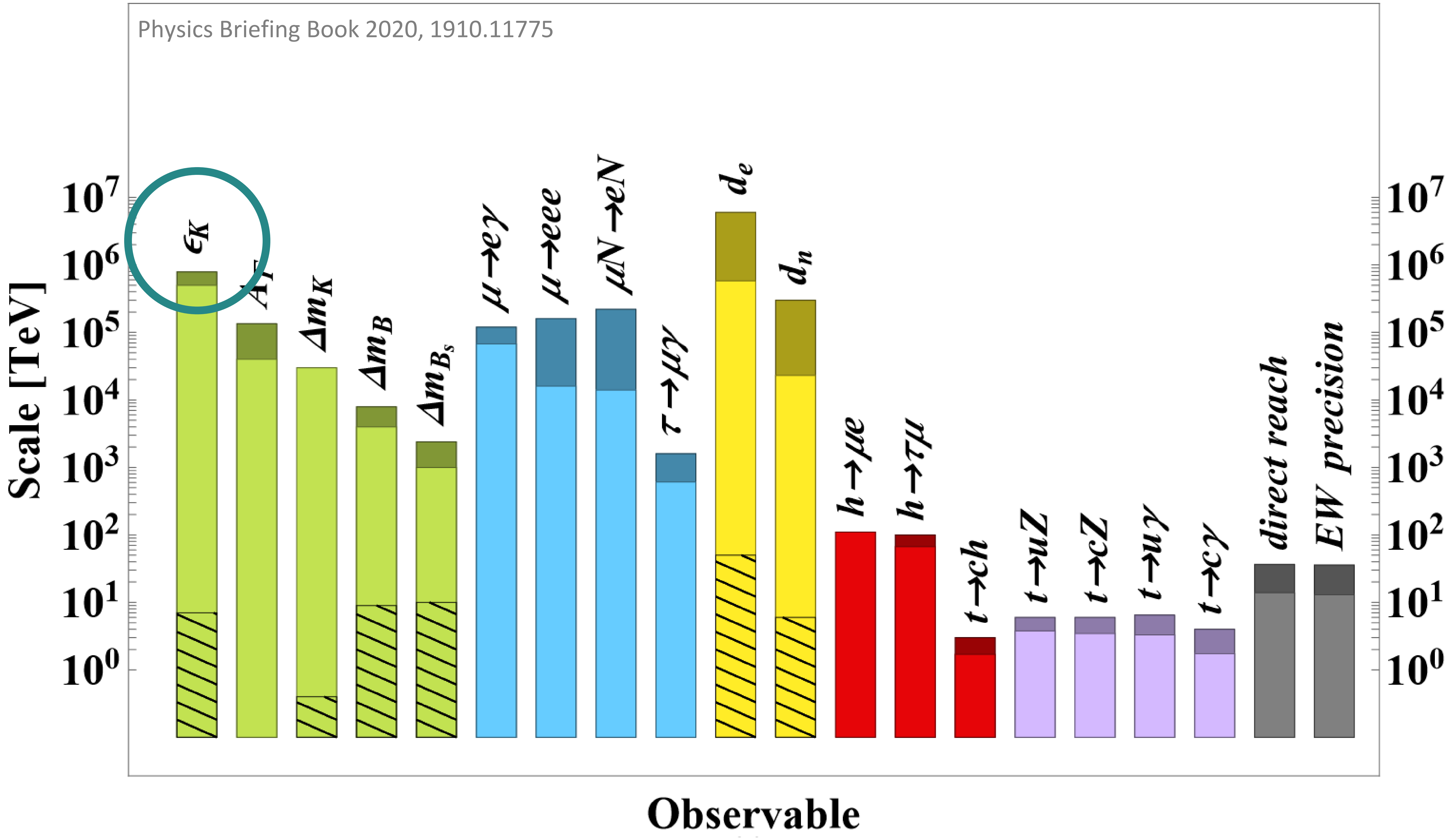
↑ Composite partner



- Fermionic Ops -> No risk of reintroducing a hierarchy problem for $\mathcal{O}_F^2$
- Partners for each fermions -> can reproduce Yukawa pattern

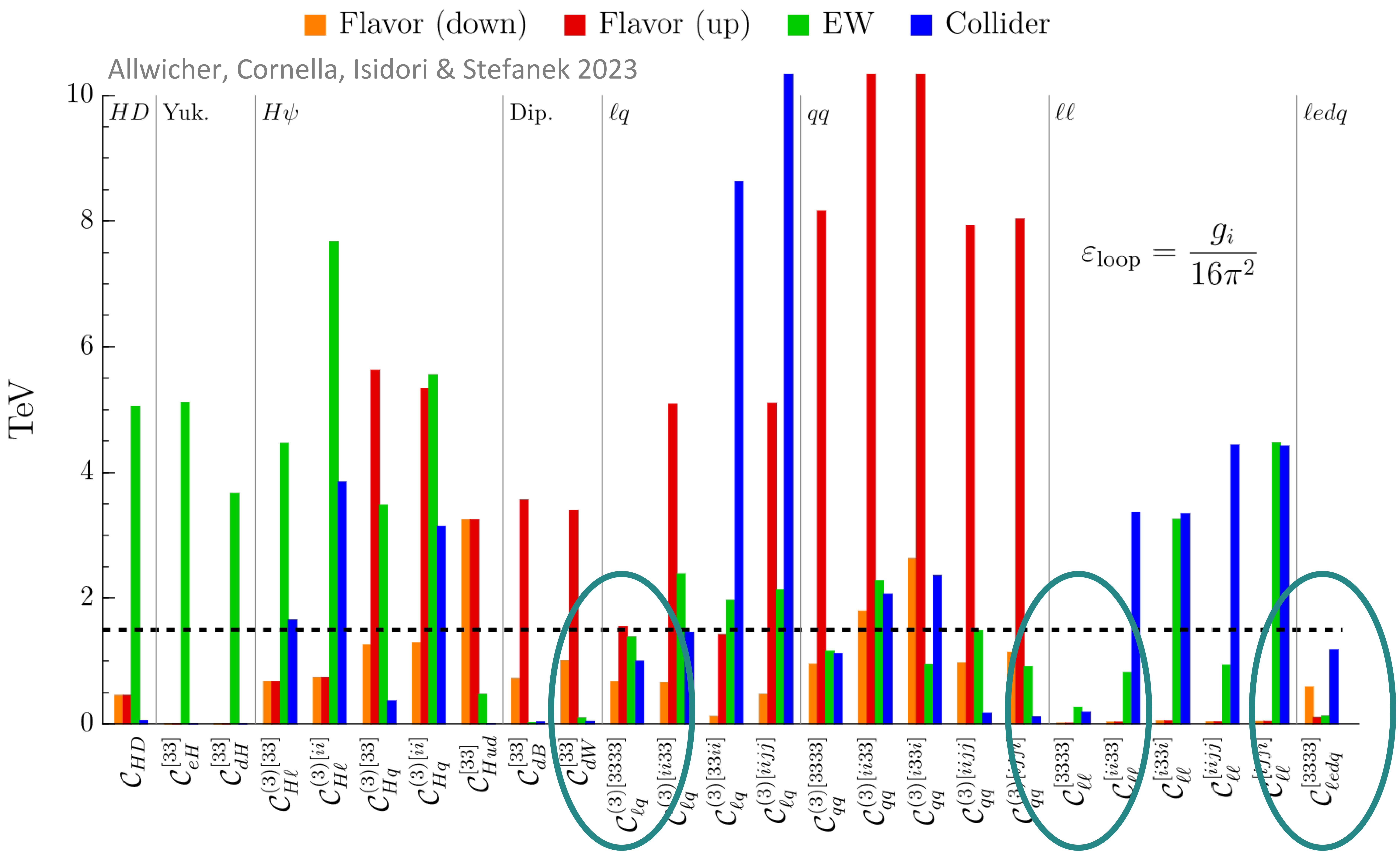
Experimental Constraints on NP Scale

- ➢ No clear direct signals of NP -> Mass gap is a « fact » of life
- ➢ Proton decay, neutrino masses, EDMs, … -> NP scale could be very high
- ➢ **Flavour** probes very high scale too!



Back to Flavour

3rd Gen. is the least constrained!

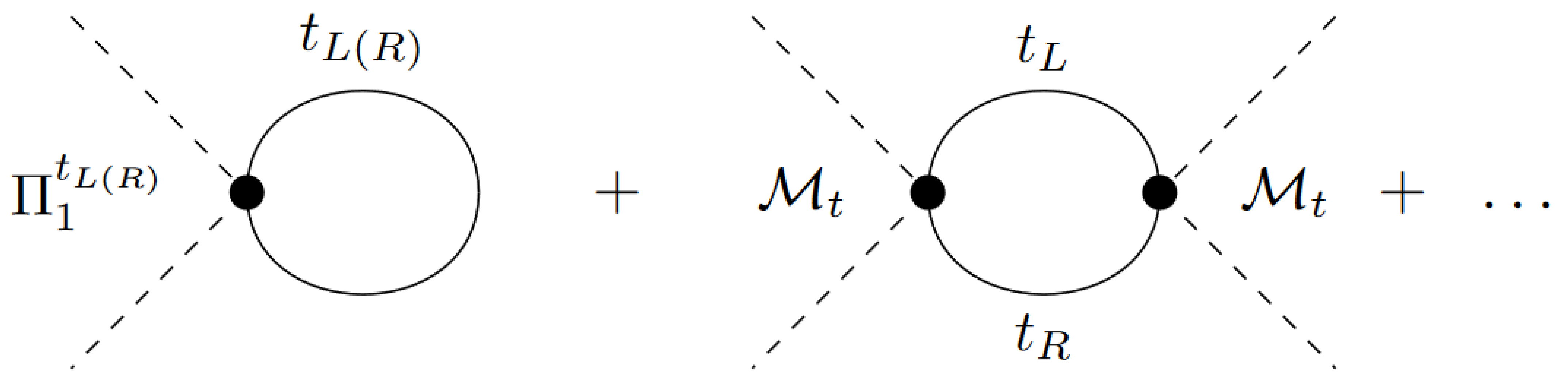


Higgs Potential

$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

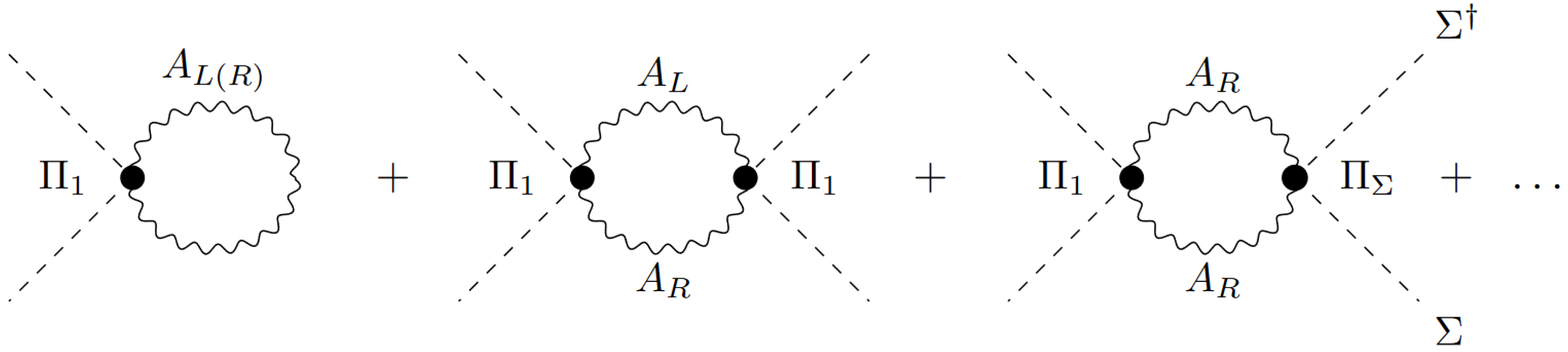
Fermion contribution

$\Delta V(h)_f$



Gauge contribution

$\Delta V(h)_A$



Higgs Potential

$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

$$\left.\frac{c_1}{F^4}\right|_{\mathrm{phys.}} = \frac{m_h^2}{F^2} \lesssim 0.06$$ (from Exp.)

$$\left.\frac{c_2}{F^4}\right|_{\mathrm{phys.}} = \frac{2m_h^2}{v^2} \approx \frac{1}{2}$$ Natural

Higgs Potential

$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

$$\frac{c_2}{F^4} = \frac{N_c y_t^2}{4\pi^2}\frac{M_T^2}{F^2} + \text{Gauge contributions (suppressed)}$$

Gauge contributions
(suppressed)

↑

Top partner → $M_T \approx 2.5F$

Higgs Potential

$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

Fermionic resonances — Top partner — Gauge contributions

$$\frac{c_1}{F^4} = \frac{N_c}{8\pi^2}\left[\left(\lambda_R^t\right)^2 \kappa_R^t - \left(\lambda_L^t\right)^2 \kappa_L^t\right]\frac{M_f^2}{F^2} + \frac{N_c y_t^2}{4\pi^2}\frac{M_T^2}{F^2} - \frac{9 g_R^2}{32\pi^2}\left(1 - \frac{g_R^2 v_\Sigma^2}{2M_\rho^2}\right)\frac{M_\rho^2}{F^2} + \mathcal{O}\left(g_L g_R, g_L^2\right)$$

- Increase size of gauge contribution ⟶ $g_{R,3} = O(1) \gg g_{R,12} \approx g_Y^{\text{SM}}$

  (Natural in flavour non-universal scenario !)

- Avoid suppression / sign flip ⟶ $M_{W_R}^2 = \frac{1}{4} g_R^2 v_\Sigma^2 < \frac{1}{2} M_\rho^2$

Phenomenological Constraints

- Constraints related to strong dynamics
  - Modification of VVh- and VVhh-couplings

$$F \gtrsim 500\,\mathrm{GeV}$$

  - Top partners and heavy resonances searches

$$M_T \gtrsim 1.5\ \mathrm{TeV} \longrightarrow F \gtrsim 600\,\mathrm{GeV}$$

$$M_\rho \gtrsim 5\ \mathrm{TeV}$$

  - EWPO (S and T parameters)

$$g^2_{L,R}\frac{v^2}{M^2_\rho} \lesssim 10^{-3}$$

- Constraints related to flavoured gauge bosons
  - $B \to X_s\gamma$
  - Bound on Z-pole obs.

$$v_\Sigma \gtrsim 3\ \mathrm{TeV}$$

  - Bounds on Z′ masses from FCNC (B_s-mixing)

$$v_\Omega \gtrsim 2.7\ \mathrm{TeV}$$ (up- vs down-alignment)

  - LHC bound from Drell-Yan data

$$v_\Sigma \gtrsim 2.0\ \mathrm{TeV}$$

Phenomenological Constraints

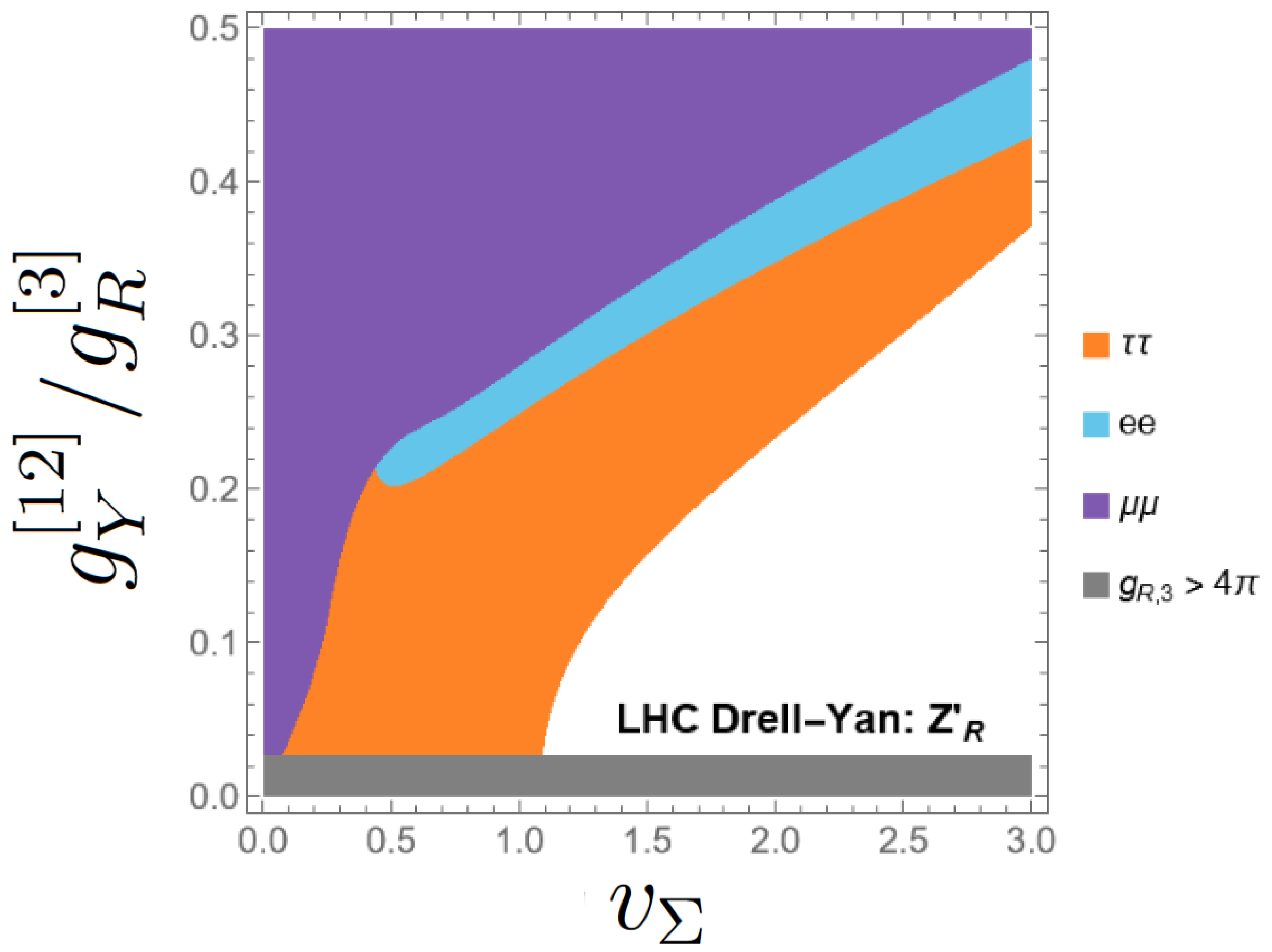
Typical scenario

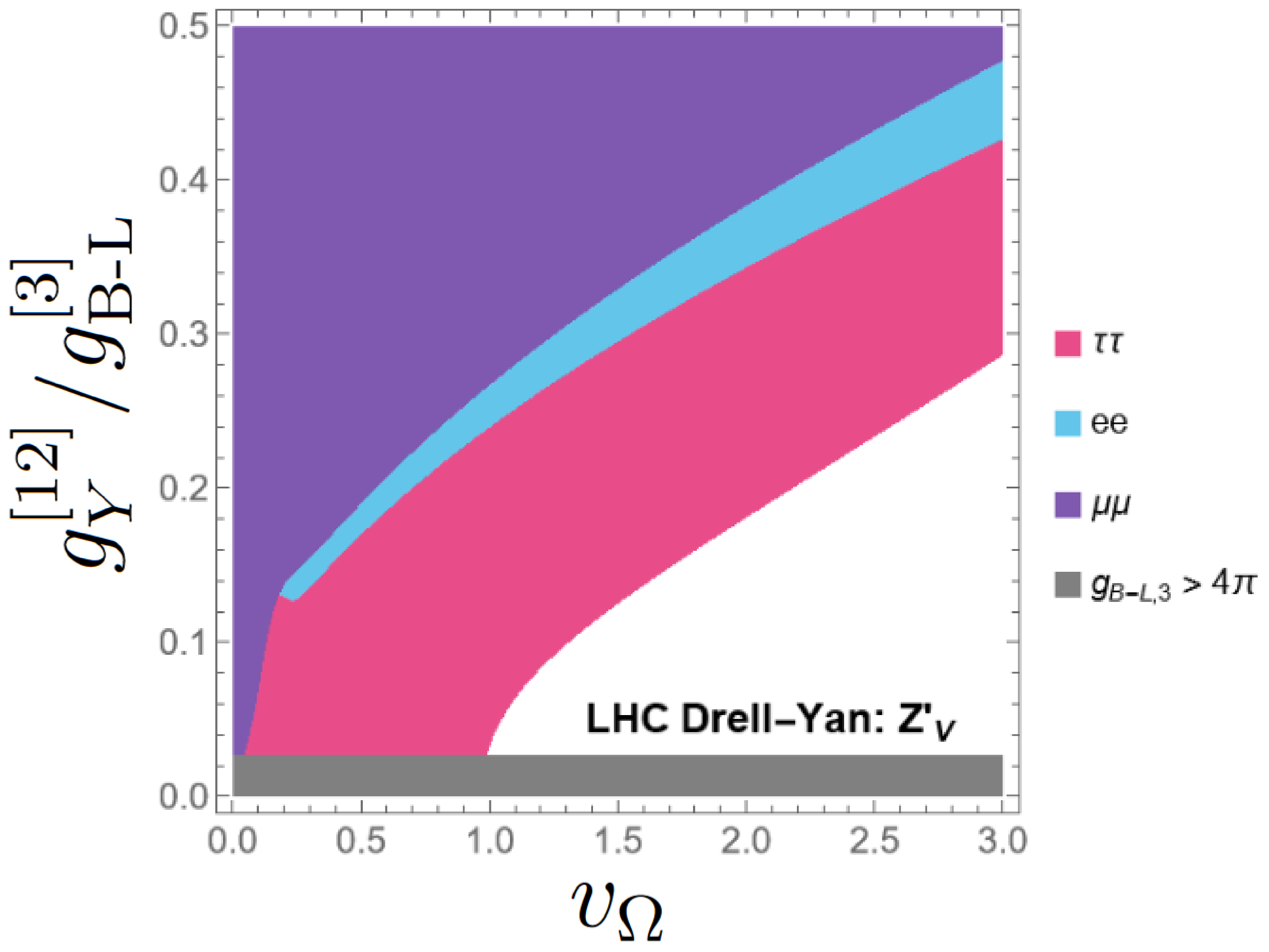
- Large 3rd gen. RH gauge coupling: $g_{R,3} = O(1)$
- Light Top partner $M_T \approx 2$ TeV and $M_\rho \approx 10$ TeV
- Flavour deconstruction breaking $v_\Sigma \approx 3$ TeV

All constraints are satisfied and $\delta_{\rm EW} \lesssim 10^{-3}$

→ 3% tuning in the potential

→ $O(1\%)$ corrections to Higgs couplings





Allwicher et al. 2022

Conclusion

➡ Null results @LHC put pressure on *natural* solutions to the hierarchy problem…

Naturalness has played a crucial role in NP searches in the past …

➡ Bounds on flavour violation suggest either a high NP scale or non-generic flavour of BSM

Approx. U(2)-preserving + 3rd gen. NP compatible with TeV scale

- Flavour non-universal NP @TeV compatible with exp. bounds & accessible at current and near future exp.

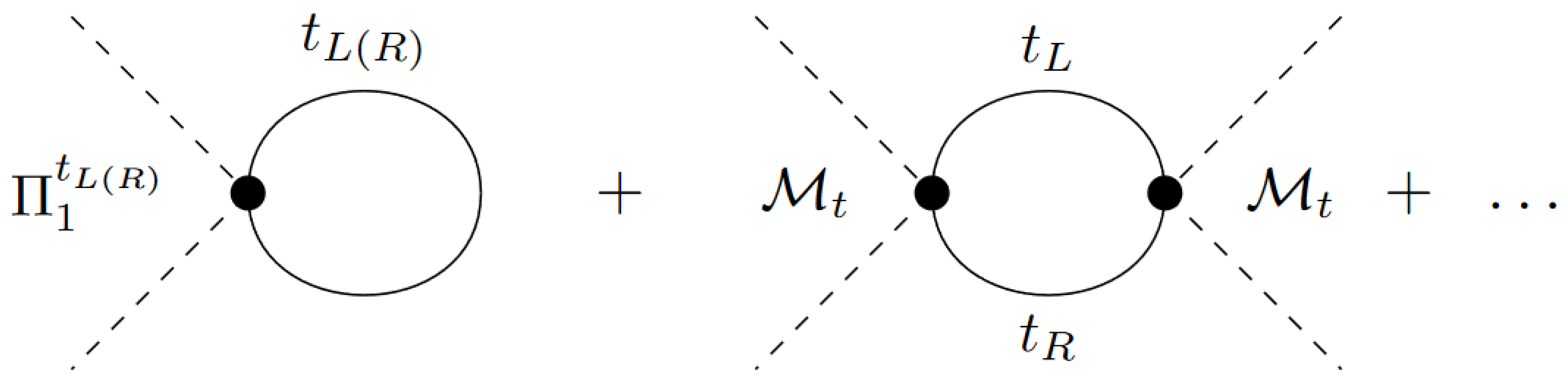
- Well-motivated *model* for addressing *simultaneously* Higgs & flavour

Higgs Potential

➡ Higgs potential induced at 1-loop

Fermion contribution

$\Delta V(h)_f$



$$\mathcal{L}_{\rm eff} \supset \bar{q}_L \not{p} \left[\Pi_0^{q_L}(p^2) \mathbb{1} + \Pi_1^{t_L}(p^2) u_L^\dagger \Delta_+ u_L \right] q_L$$
$$+ \left\{ \bar{q}_L \left[\mathcal{M}_t(p^2) u_L^\dagger \Delta_+ u_R \right] q_R + \text{h.c.} \right\}$$

$$\Pi_1^{t_L}(0) = \frac{F^2}{M_T^2} \left(\lambda_L^t\right)^2 \kappa_L^t$$

$$|\mathcal{M}_t(0)| = y_t \sqrt{2} F$$

Higgs Potential

➡ Higgs potential induced at 1-loop

Fermion contribution

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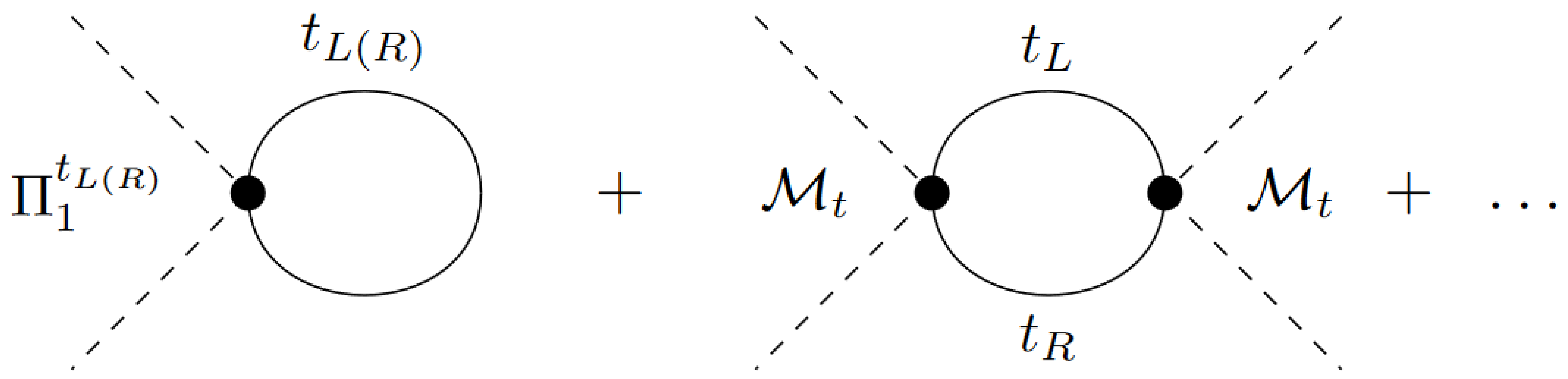

Coleman-Weinberg potential ⟶

$$\Delta V(h)_f = -2N_c \int \frac{d^4 p_E}{(2\pi)^4} \left\{ \log\left[1 + \frac{\Pi_1^{t_L}}{\Pi_0^{q_L}} \sin^2\left(\frac{h}{2F}\right)\right] + \log\left[1 + \frac{|\mathcal{M}_t|^2 \sin^2\left(\frac{h}{2F}\right)\cos^2\left(\frac{h}{2F}\right)}{p_E^2\left(\Pi_0^{q_L} + \Pi_1^{q_L}\sin^2\left(\frac{h}{2F}\right)\right)\left(\Pi_0^{q_R} - \Pi_1^{q_R}\sin^2\left(\frac{h}{2F}\right)\right)}\right]\right\}$$

Higgs Potential

➡ Higgs potential induced at 1-loop

Fermion contribution

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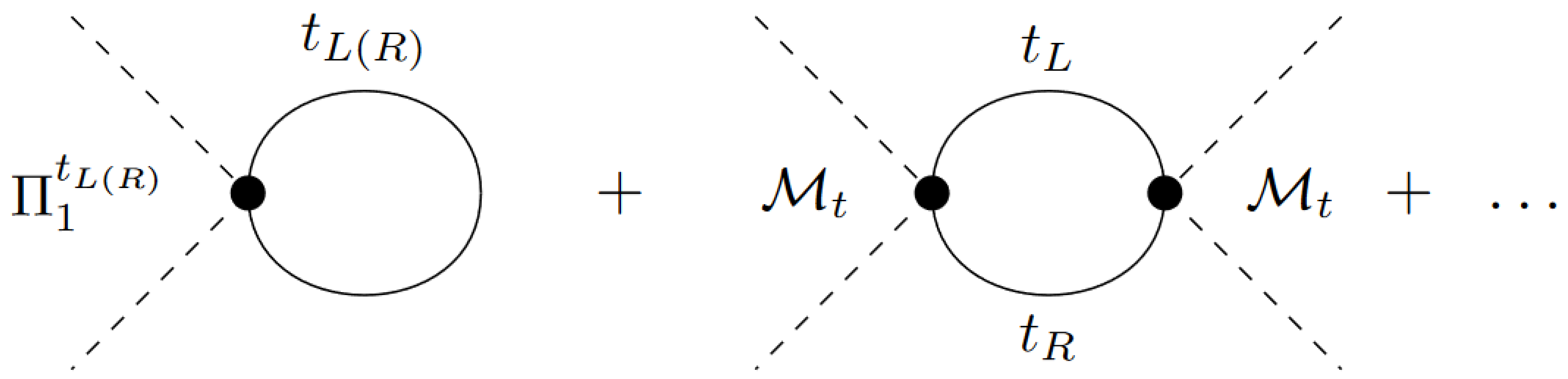

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$$\mathcal{M}_t(q^2) = \mathcal{M}_t(0) \times \frac{M_T^2}{M_T^2 - q^2}$$

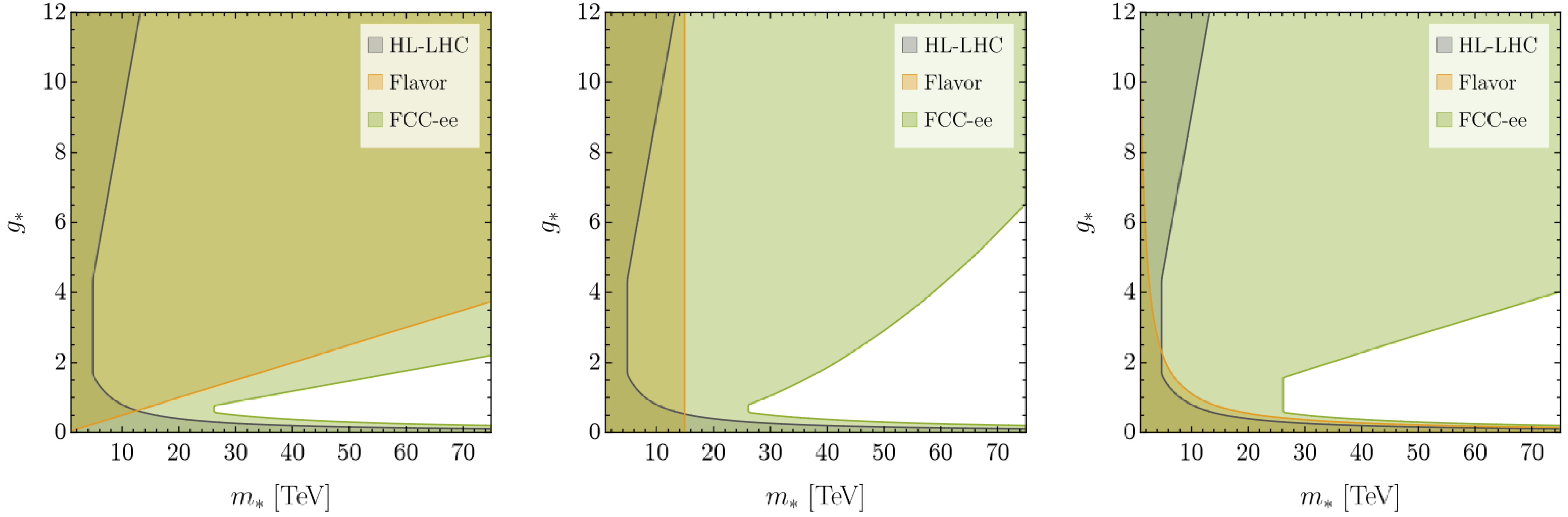
$$\frac{\Pi_1^{t_L}(q^2)}{\Pi_1^{t_L}(0)} \frac{\Pi_0^{q_L}(0)}{\Pi_0^{q_L}(q^2)} = \frac{M_T^2}{M_T^2 - q^2} \frac{M_f^2}{M_f^2 - q^2}$$

⟶ Explicit expression in terms of model parameters

Composite Higgs @ HL-LHC and FCC-ee

$\left[\ m_* = F g_*\ \right]$

Stefanek 2024



(a) Left compositeness (b) Mixed compositeness (c) Right compositeness

➡ With improved precision: RG-running into EWPO become **crucial**

➡ Composite Higgs will be put under a microscope @ FCC-ee!

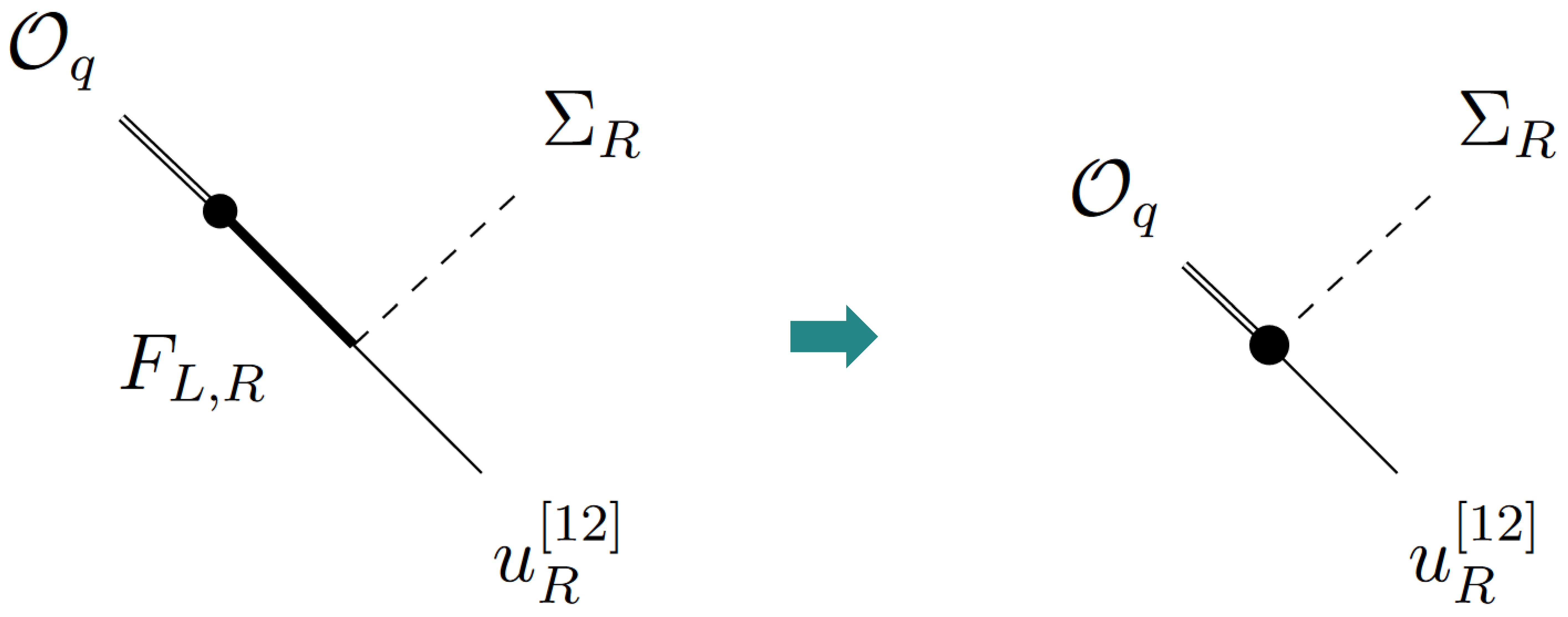
$$m_* \gtrsim 25\ \text{TeV}$$

Flavour Deconstructing the Composite Higgs

$$SU(3)_c \times SU(2)_L \times SU(2)_R^{[3]} \times U(1)_{B-L}^{[3]} \times U(1)_Y^{[12]}$$

Horizontal Breaking $\langle \Sigma_R \rangle$ $\langle \Omega \rangle$

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$



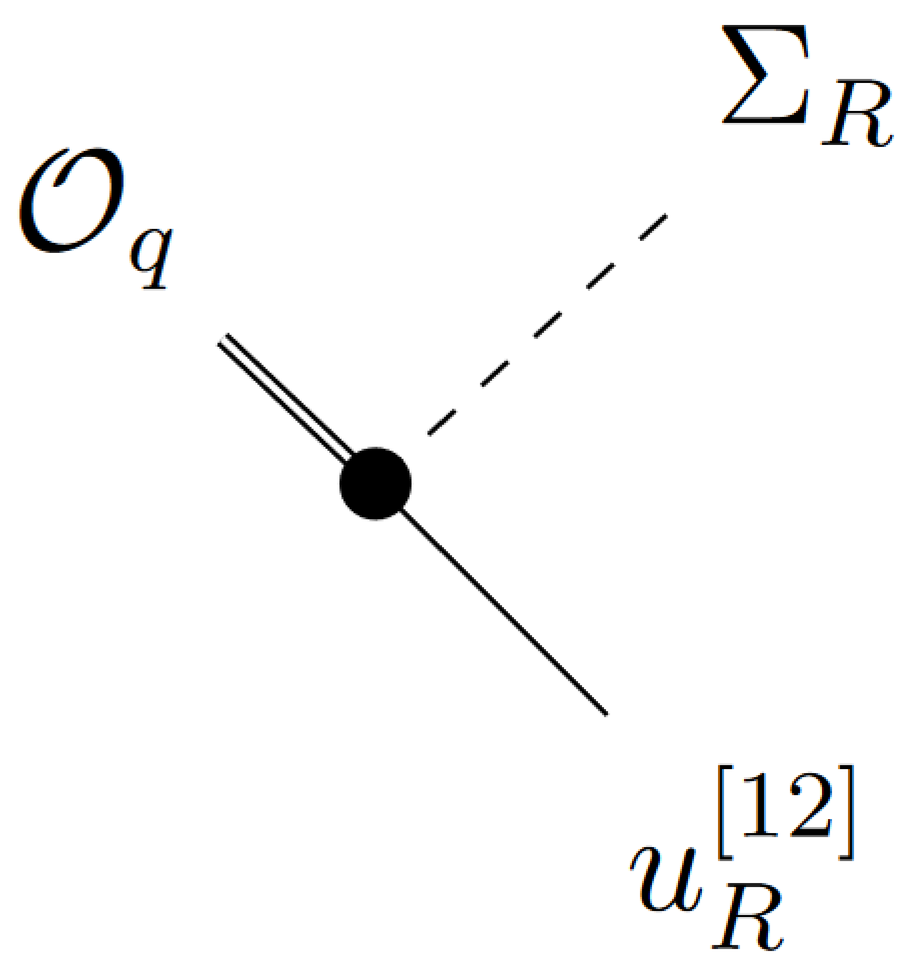
Flavour Deconstructing the Composite Higgs

$$SU(3)_c \times SU(2)_L \times SU(2)_R^{[3]} \times U(1)_{B-L}^{[3]} \times U(1)_Y^{[12]}$$

Horizontal Breaking: $\langle \Sigma_R \rangle$ $\langle \Omega \rangle$

$$\downarrow$$

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$



$$\epsilon_R = \frac{\langle \Sigma \rangle}{M_F}$$

$$\longrightarrow$$

$$\epsilon_\Omega = \frac{\langle \Omega \rangle}{M_F}$$

$$Y_{u,d,e} \sim \begin{pmatrix} \epsilon_R & \epsilon_\Omega \\ \epsilon_R \epsilon_\Omega & 1 \end{pmatrix}$$

Flavour Deconstructing the Composite Higgs

$$Y_{u,d,e} \sim \begin{pmatrix} \epsilon_R & \epsilon_\Omega \\ \epsilon_R \epsilon_\Omega & 1 \end{pmatrix}$$

$$\epsilon_\Omega = O\left(|V_{cb}|\right) = O\left(10^{-1}\right)$$

$$\epsilon_R = O\left(m_c/m_t\right) = O\left(10^{-2}\right)$$

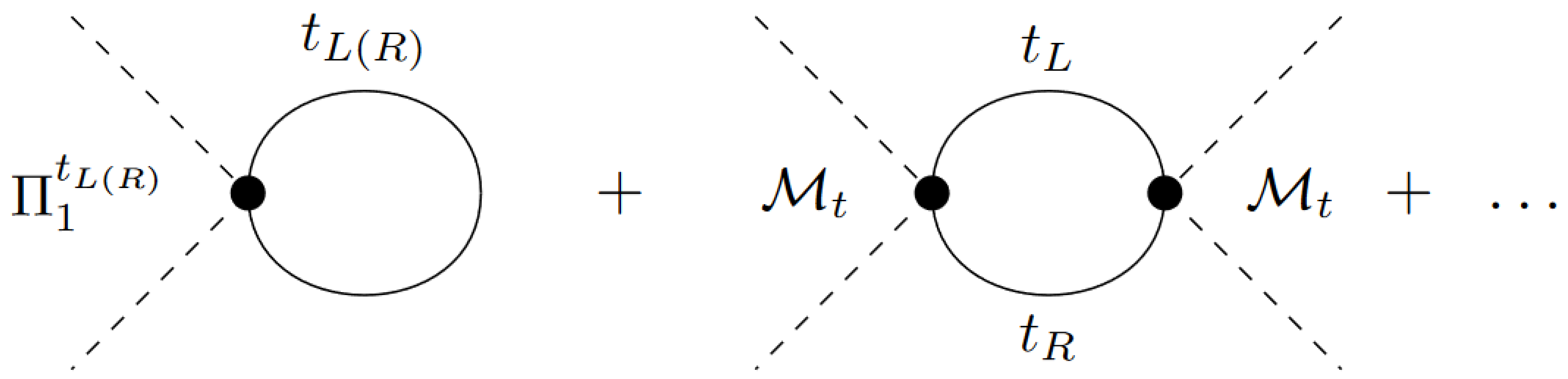
⟶ The deconstruction scale is ***anchored*** by its impact on the Higgs potential

Higgs Potential

➡ Higgs potential induced at 1-loop

Fermion contribution

$\Delta V(h)_f$



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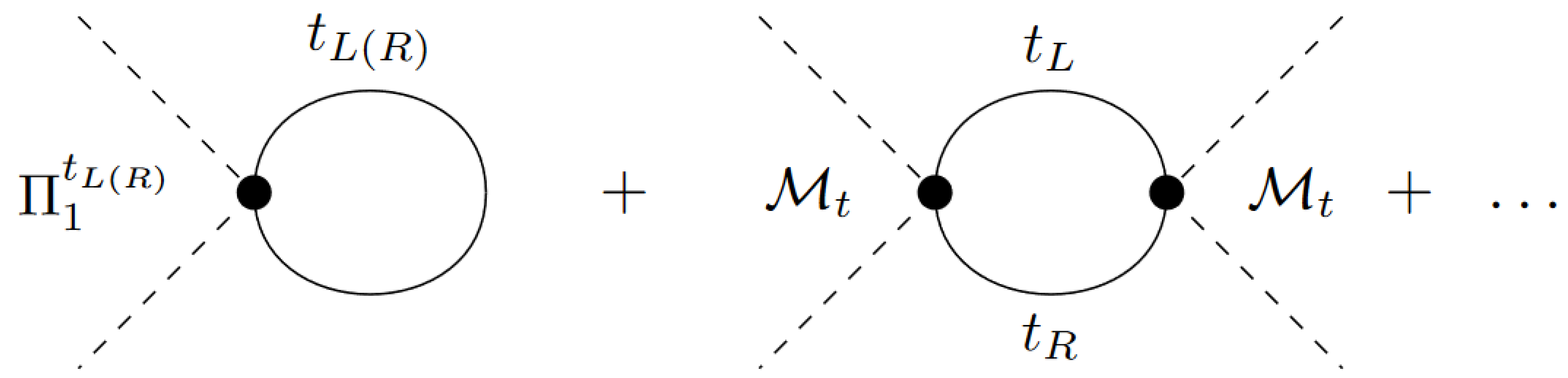

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Higgs Potential

➡ Higgs potential induced at 1-loop

Fermion contribution

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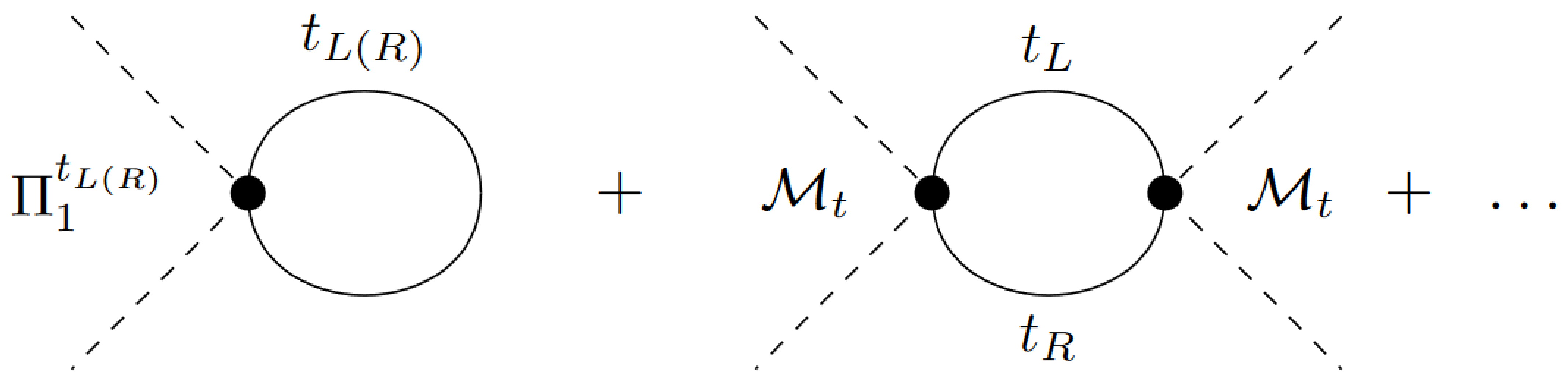

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⟶ Explicit expression in terms of model parameters

The Hierarchy Problem

What possible theoretical frameworks can address the Hierarchy problem ?

« Beyond QFT / EFT » option

↓

Solution: « Elsewhere »

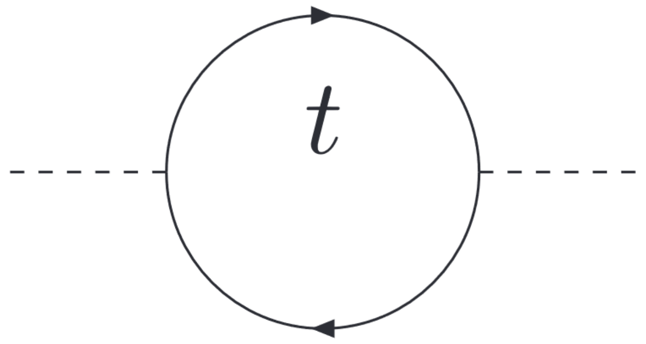
Multiverse / Anthropics
Cosmological evolution
(Failure of EFT)

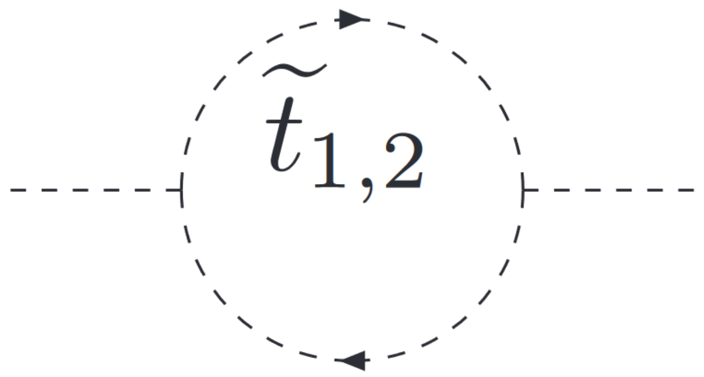
Agrawal et al., 1997
Kawai & Okada, 2011
Giudice et al. 2021
Kephart & Päs 2024, …

SUSY

↓

Elementary scalars are protected by symmetry





Higgs Compositeness

↓

No elementary scalars

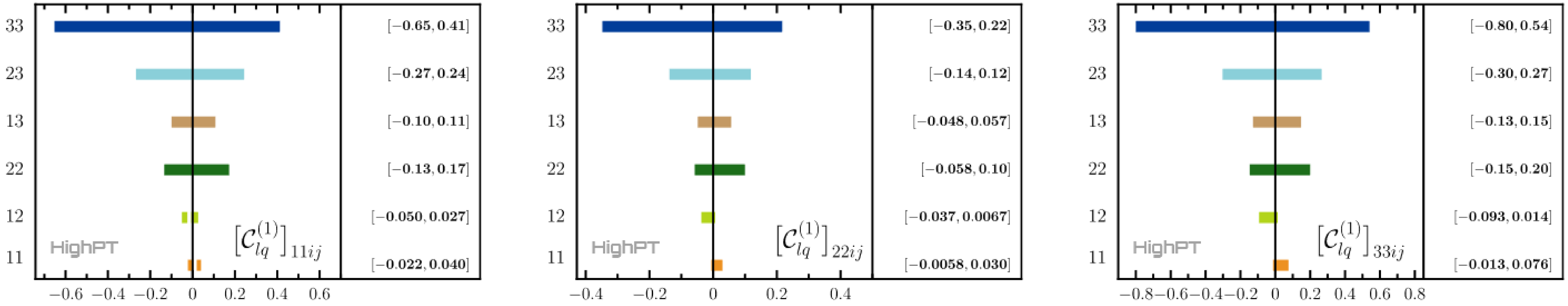
Higgs emerges as a composite pseudo-Goldstone boson of S.S.B

Dugan et al. 1985,
Agashe, Contino &
Pomarol 2005,…

Back to Flavour

3rd Gen. is the least constrained

High-pt Drell-Yan tails



Allwicher et al. 2022

Allwicher et al. 2022
D'Ambrosio et al. 2002
Li & Ma 1981

Back to Flavour

3rd Gen. is the least constrained

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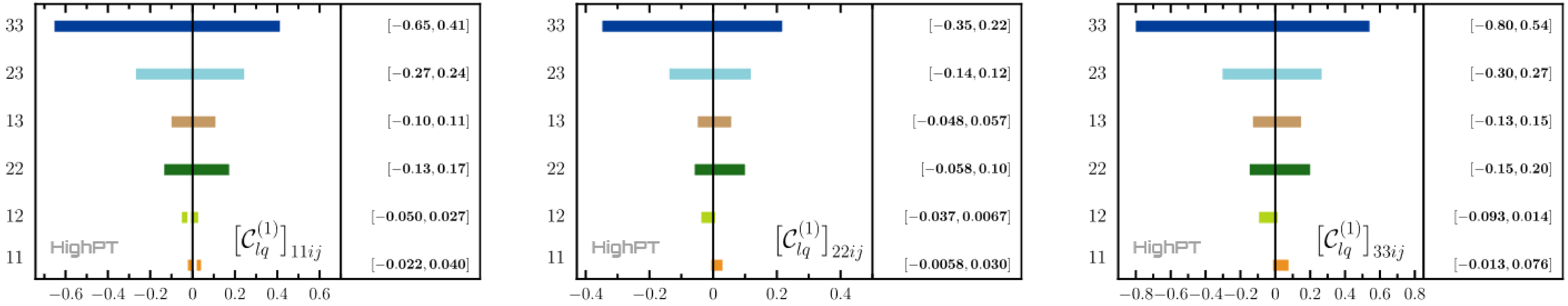


Allwicher et al. 2022

Experiments have imposed strong bounds on flavour **universal** NP

(Naturalness 😑)

➡ Flavour **non-universal** NP @TeV-scale, mainly coupled to 3rd gen.

Allwicher et al. 2022
D'Ambrosio et al. 2002
Li & Ma 1981

Flavour Non-Universal Composite Higgs

Ingredients:

- Spontaneously broken strong sector: $\mathcal{G} \equiv Sp(4) \xrightarrow{\Lambda_{\rm HC}} SU(2)_L \times SU(2)_R^{[3]} \equiv \mathcal{H}$

- Field Content:

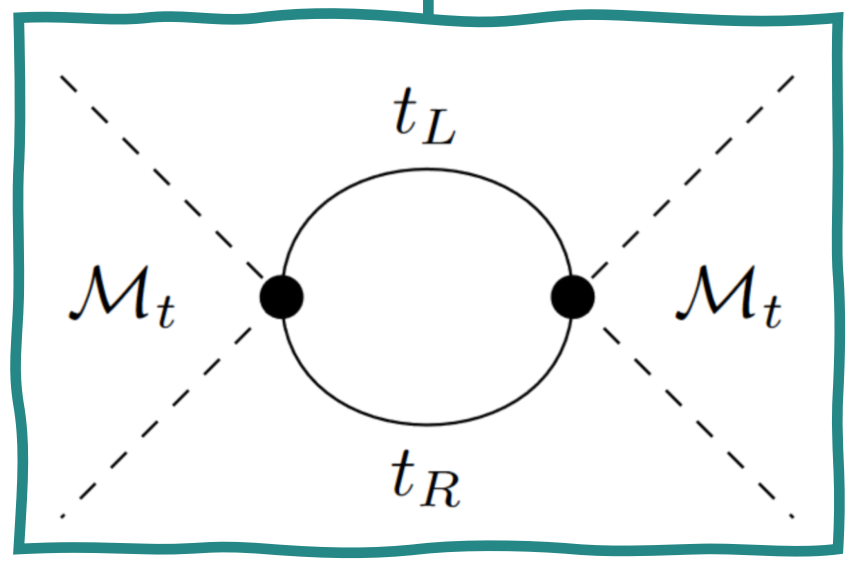
Elementary fields		$U(1)_{B-L}^{[3]}$	$U(1)_Y^{[12]}$	$SU(2)_L$	$SU(2)_R^{[3]}$
chiral light quarks	$q_L^{[12]}$	0	1/6	**2**	**1**
	$u_R^{[12]}$	0	2/3	**1**	**1**
	$d_R^{[12]}$	0	−1/3	**1**	**1**
chiral 3rd gen. quarks	$q_L^{[3]}$	1/6	0	**2**	**1**
	$q_R^{[3]}$	1/6	0	**1**	**2**
vector-like quarks	F_L^q	1/6	0	**2**	**1**
	F_R^q	0	1/6	**1**	**2**
scalar link fields	Σ_R	0	1/2	**1**	**2**
	Ω_q	−1/6	1/6	**1**	**1**
	Ω_ℓ	1/2	−1/2	**1**	**1**

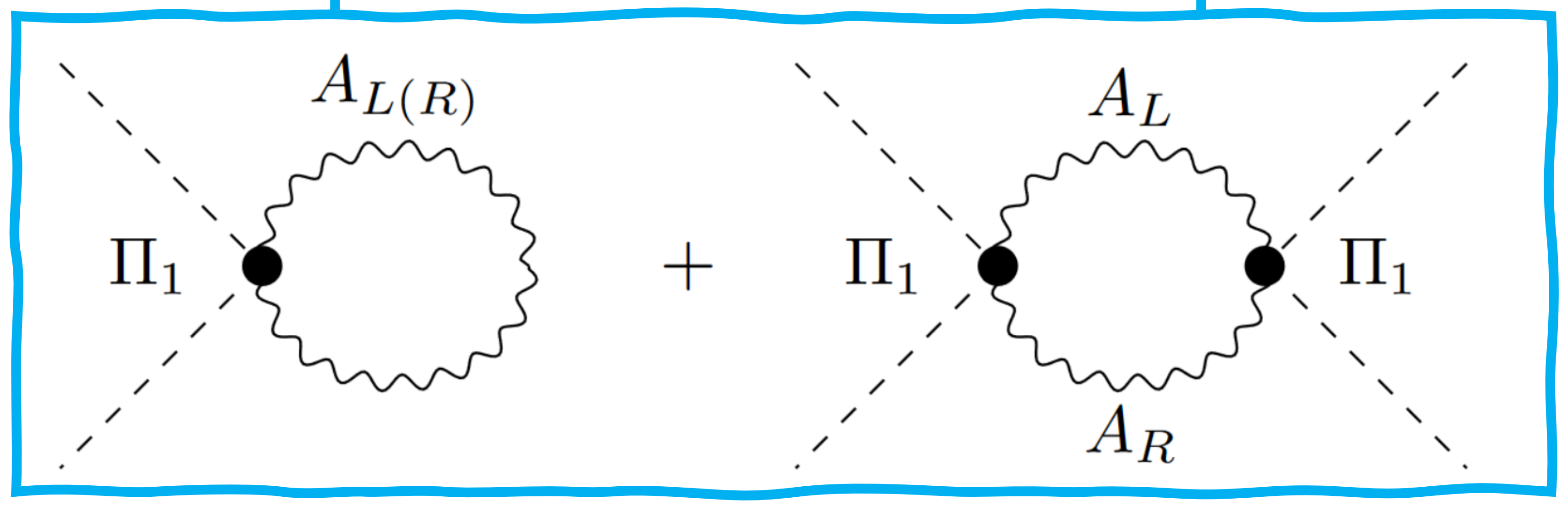
Higgs Potential

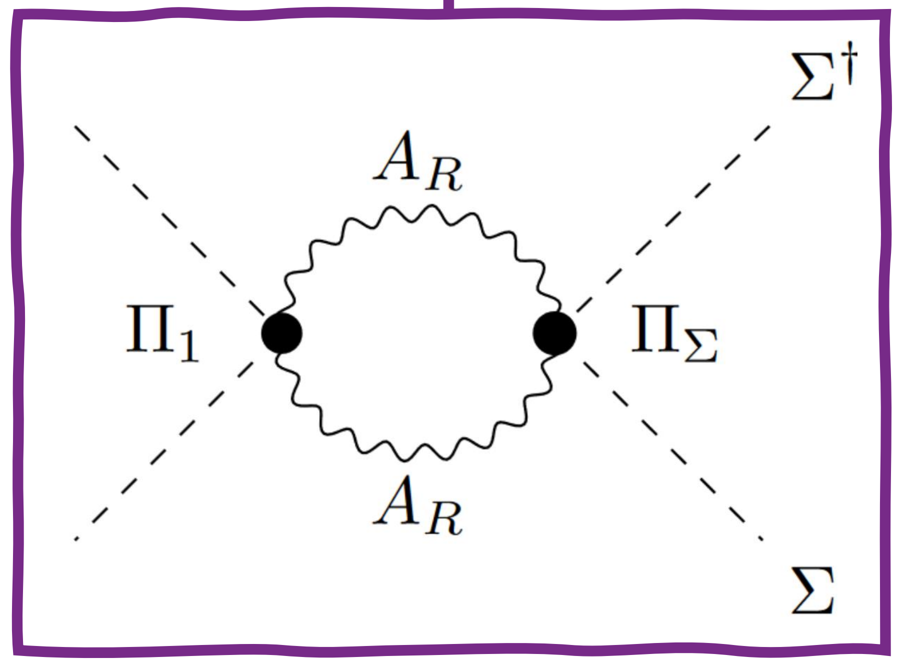
$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

Gauge contributions

$$\frac{c_2}{F^4} = \frac{N_c y_t^2}{4\pi^2}\frac{M_T^2}{F^2} + \frac{9 g_R^2}{32\pi^2}\delta_\pi\left(1 - \frac{g_R^2 v_\Sigma^2}{2M_\rho^2}\right)\frac{M_\rho^2}{F^2} - \frac{9 g_R^4}{64\pi^2}\log\left(\frac{M_\rho^2}{M_{W_R}^2}\right) + \mathcal{O}(g_L g_R, g_L^2)$$







Flavour and Higgs Compositeness

How do we couple fermions to the Higgs ?

OG approach in strongly-coupled EWSB models:

Strong resonances
-> Naturalness

$$\mathcal{L} \supset \frac{\lambda_b}{\Lambda_{\rm UV}^{d-1}} \bar{q}_L \mathcal{O}_S b_R \longrightarrow y_{t,b} \simeq \lambda_{t,b} \left(\frac{F}{\Lambda_{\rm UV}} \right)^{d-1}$$

Scalar Op. from the strong sector

Strong interactions are resolved
-> Flavour

- Difficult to have $y_t \sim \mathcal{O}(1)$ and $\Lambda_{\rm UV}$ high enough to avoid extra flavour-violation
- Reintroduces the Hierarchy problem for $\mathcal{O}_S^2$
- Enforce hierarchy of $\lambda_{t,b}$ in the UV because only one scalar op $\mathcal{O}_S$

Wulzer & Panico 2015
Agashe et al. 2005,
Rattazzi et al. 2008

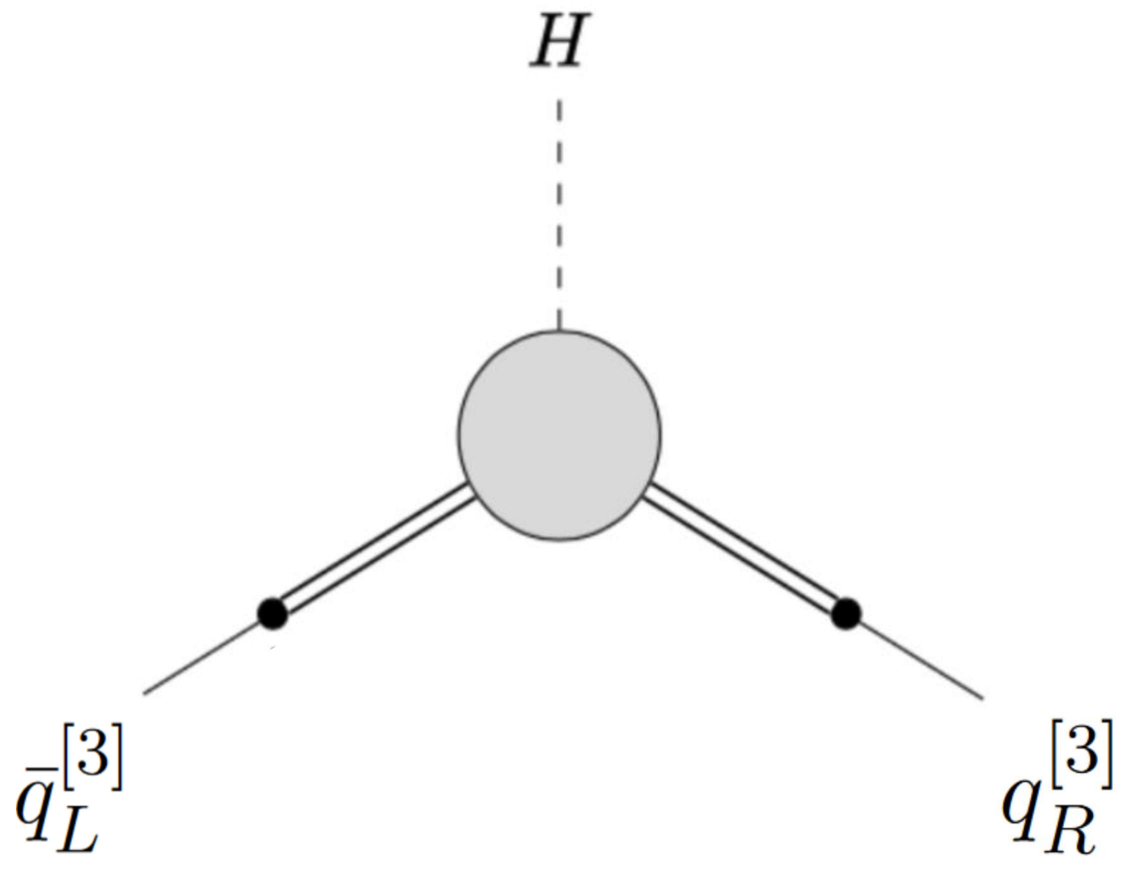
Partial Compositeness

➡ How do we couple fermions to the Higgs ?

Partial Compositeness:

$$\mathcal{L} \supset \lambda_L \bar{q}_L \mathcal{O}_R \longrightarrow |y_q| = \lambda_L^q \lambda_R^{q*} \kappa_{LR}^q \frac{F}{\sqrt{2} M_q}$$

Composite partner

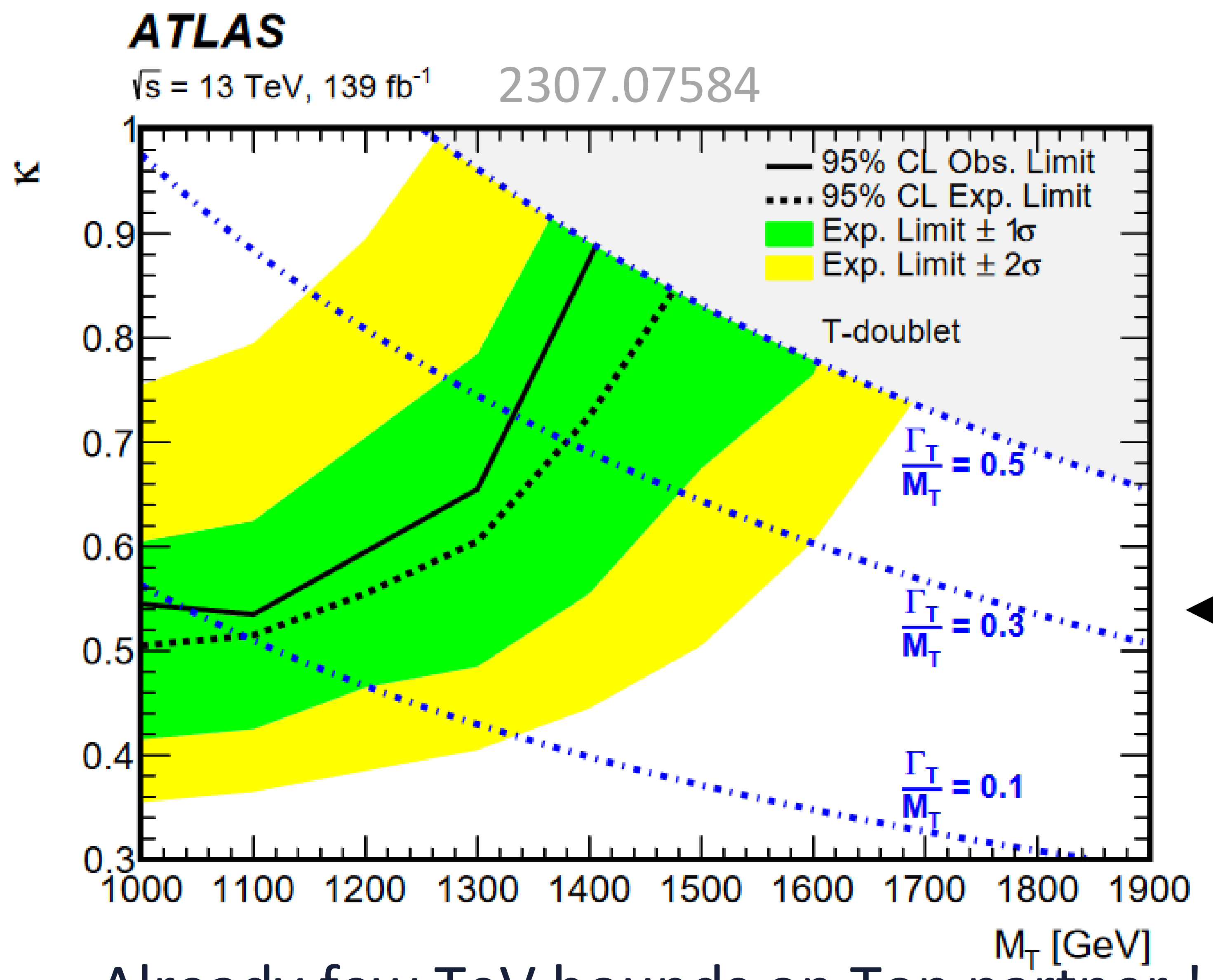


- Fermionic Ops -> No risk of reintroducing a hierarchy problem for $\mathcal{O}_F^2$
- Partners for each fermions -> can reproduce Yukawa pattern

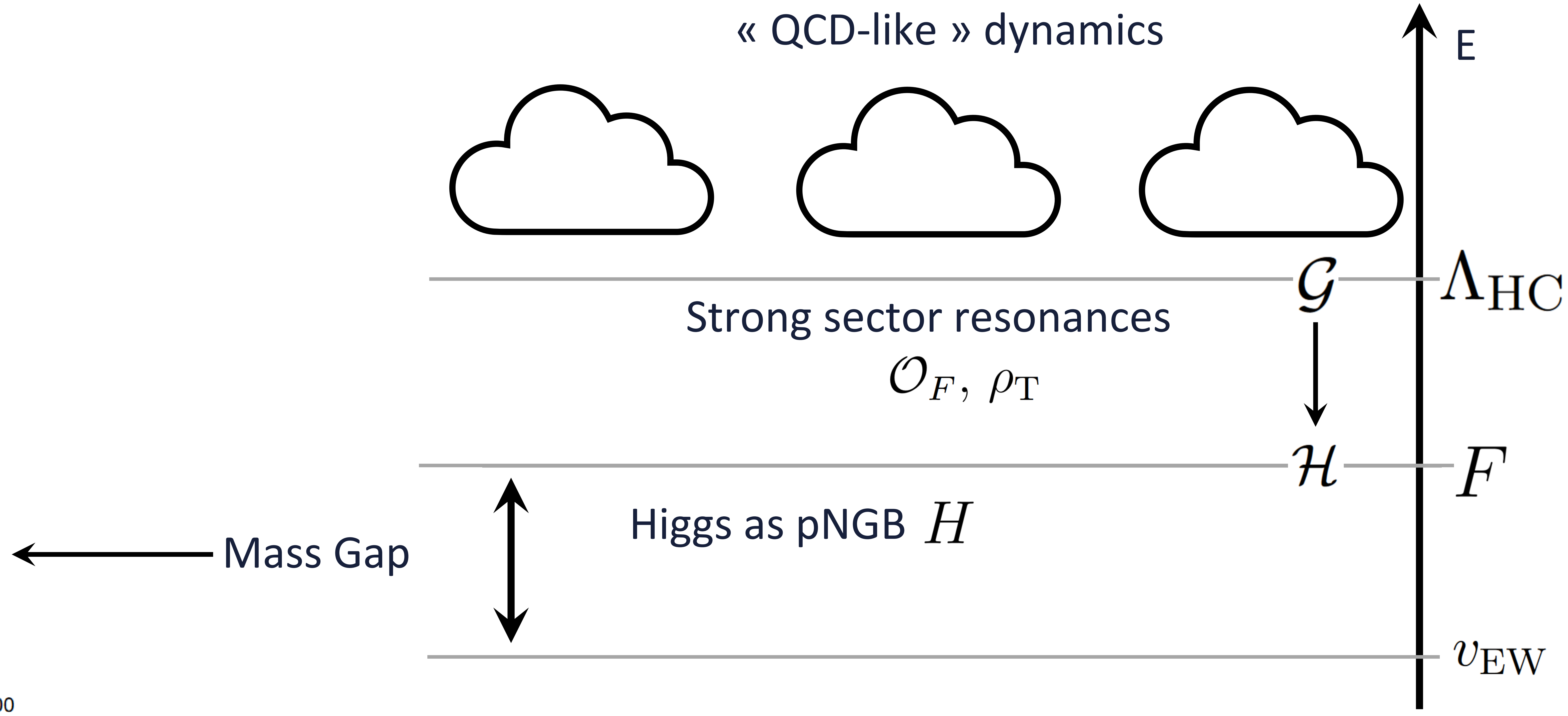
Kaplan 1991
Wulzer & Panico 2015

Higgs Compositeness

Compositeness scale cannot be too low!

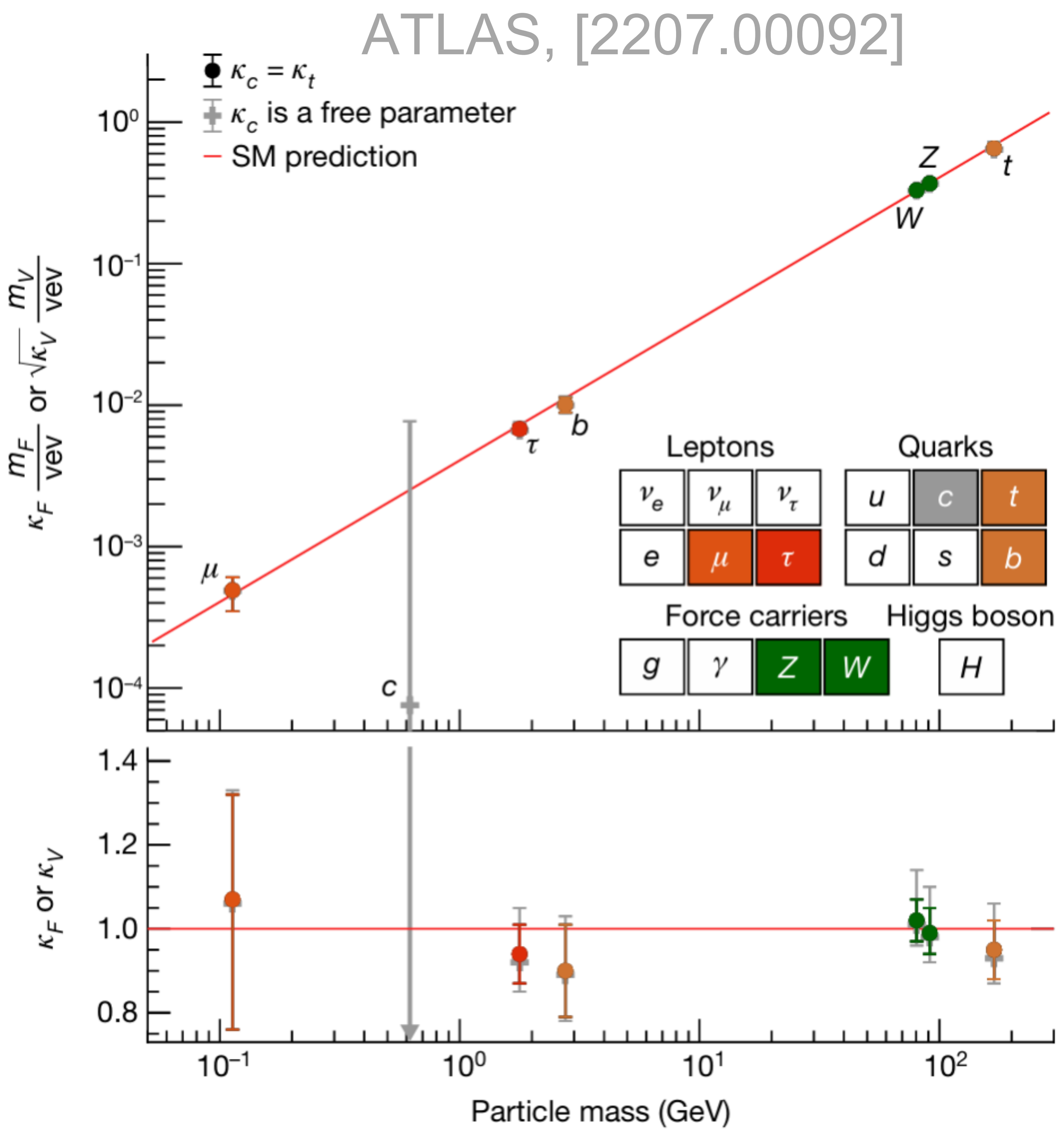


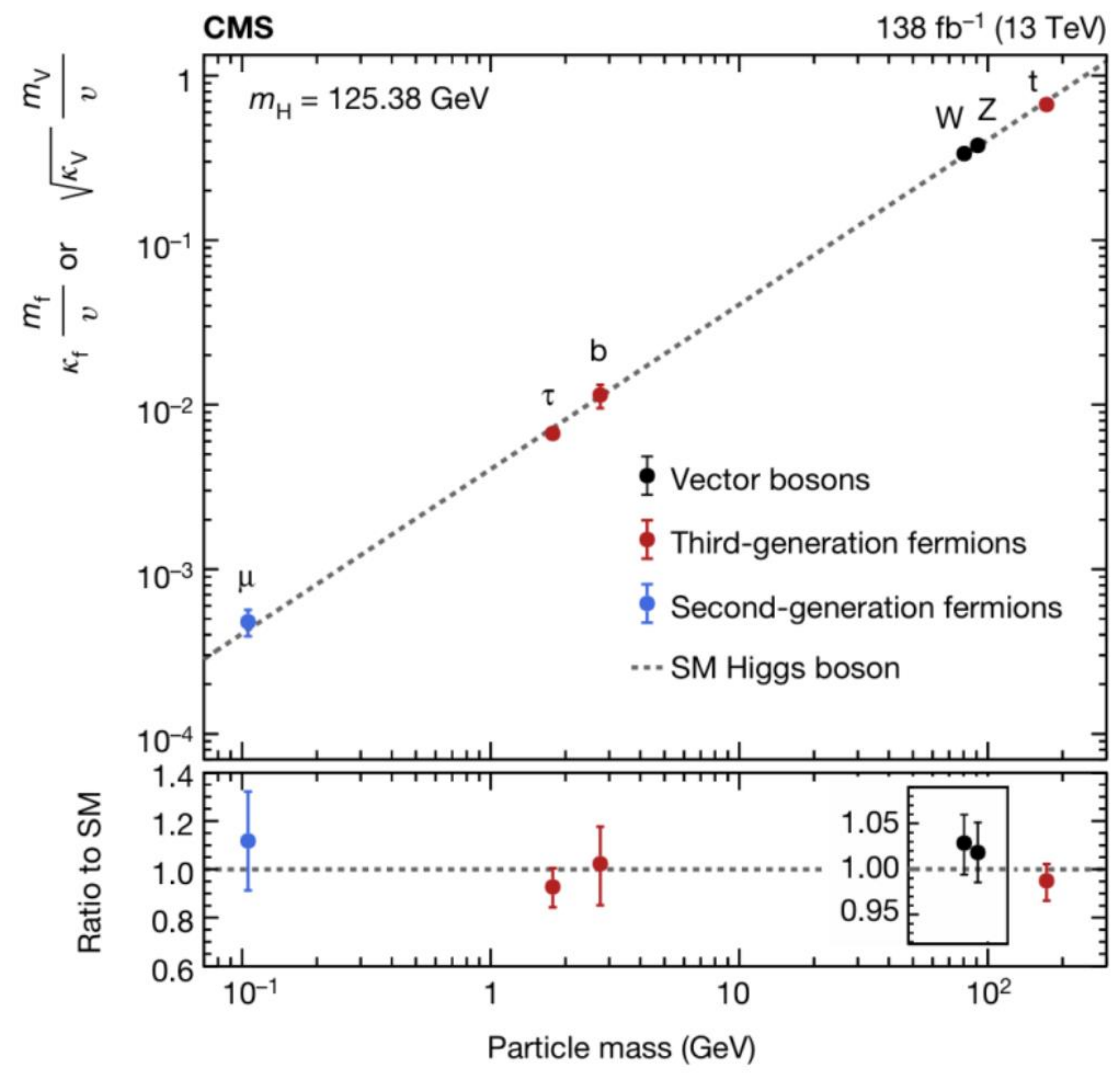
Already few TeV bounds on Top partner !



Wulzer & Panico 2015
Agashe et al. 2005,
...

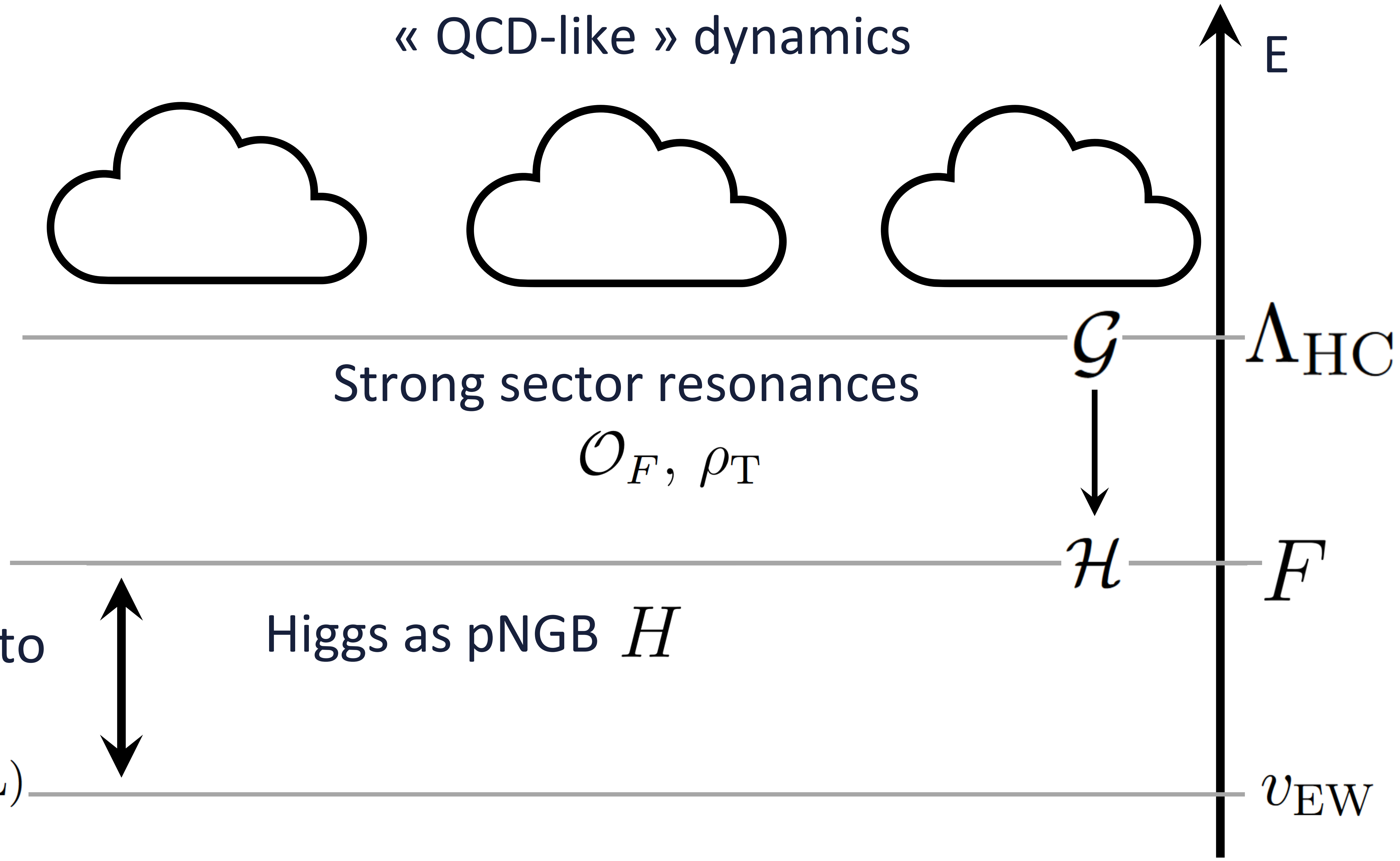
Higgs Compositeness





Higgs is SM-like to good approx.

$\xi < 0.06$ (95% CL)



Compositeness corrects SM predictions

$$g_{VVh} = g_{VVh}^{\mathrm{SM}} \sqrt{1-\xi}$$

$$g_{VVhh} = g_{VVhh}^{\mathrm{SM}} (1-2\xi)$$

$$\xi = \frac{v_{\mathrm{EW}}^2}{4F^2}$$

Wulzer & Panico 2015
Agashe et al. 2005,
...

Flavour Deconstructing the Composite Higgs

$G_{12} \times G_{3+H} \longrightarrow G_{SM}$

$$SU(2)_R^{[3]} \times U(1)_{B-L}^{[3]} \times U(1)_Y^{[12]}$$

$$\langle \Sigma_R \rangle \;\Big\downarrow\; \langle \Omega \rangle$$

$$U(1)_Y$$

Composite Higgs

$$Sp(4)^{\,*}$$

$$\Big\downarrow\; H \sim (\mathbf{2}, \mathbf{2})$$

$$SU(2)_L \times SU(2)_R^{[3]}$$

Davighi & Isidori 2023
Glioti et al. 2024

$^{*}\,\mathfrak{sp}(4) \cong \mathfrak{so}(5)$

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- Flavour non-universal NP @TeV mainly coupled to 3rd generation

  - $U(2)^5$ protection
  - Low compositeness scale -> naturalness
  - Explain SM flavour

Davighi & Isidori 2023
Glioti et al. 2024

$^*\mathfrak{sp}(4) \cong \mathfrak{so}(5)$

Flavour Deconstructing the Composite Higgs

$$Sp(6)_{\text{global}} \longrightarrow SU(2)_L \times SU(2)_R^{[3]} \times SU(2)_R^{[12]}$$

- Composite scalars needed for flavour deconstruction breaking
- Suppression in light Yukawas from heavy pNGBs -> no VLFs needed