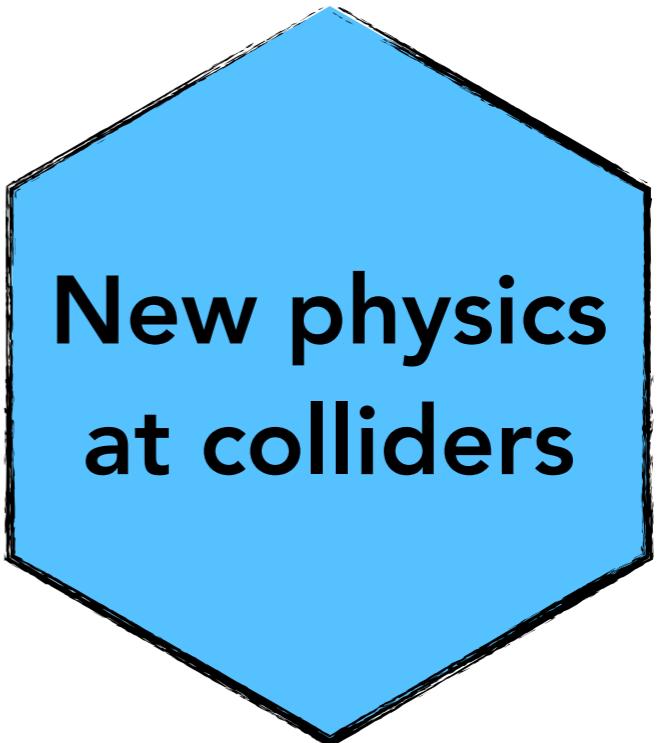


Footprints of EW Phase Transition at Colliders

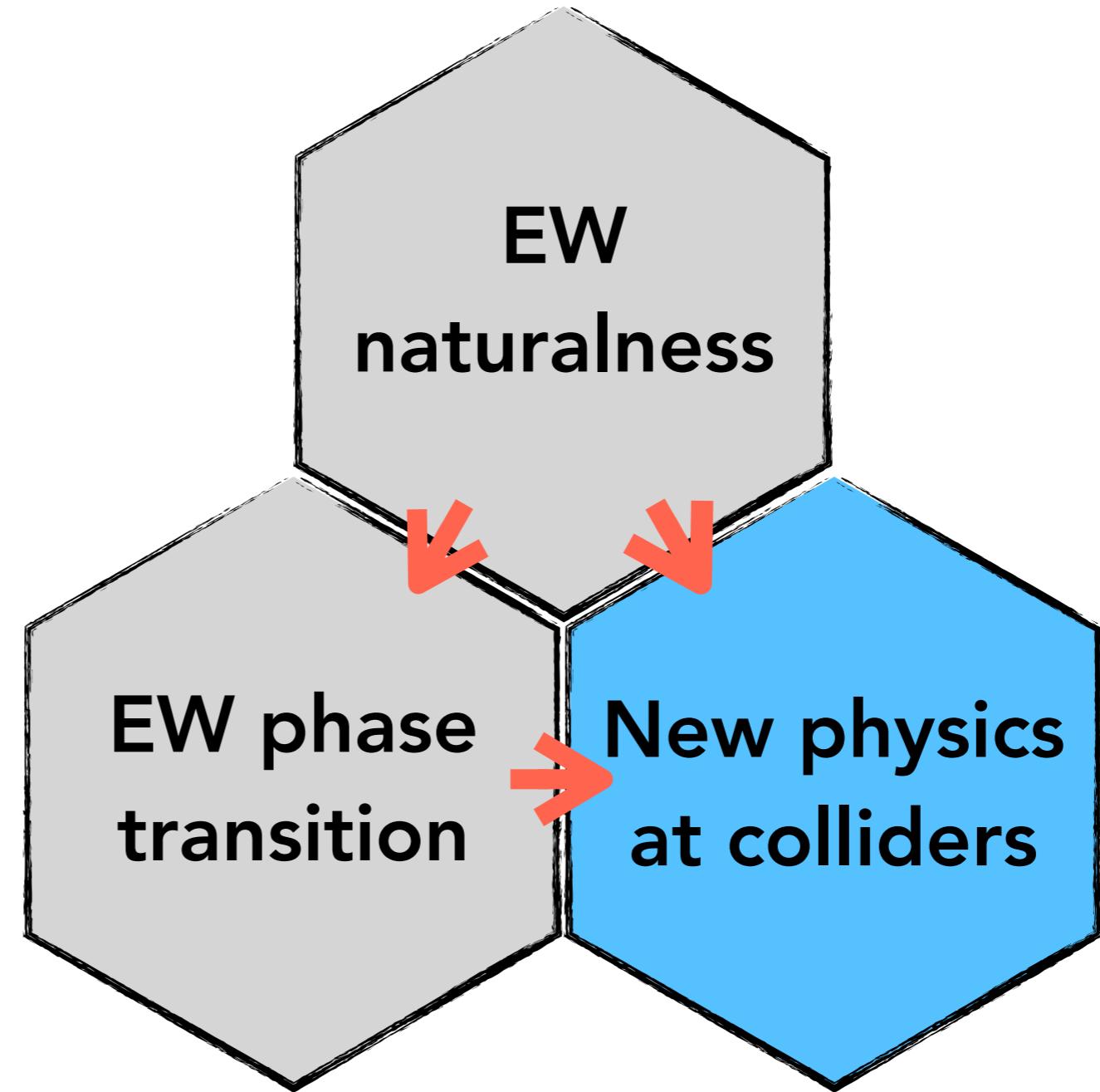
Oleksii Matsedonskyi

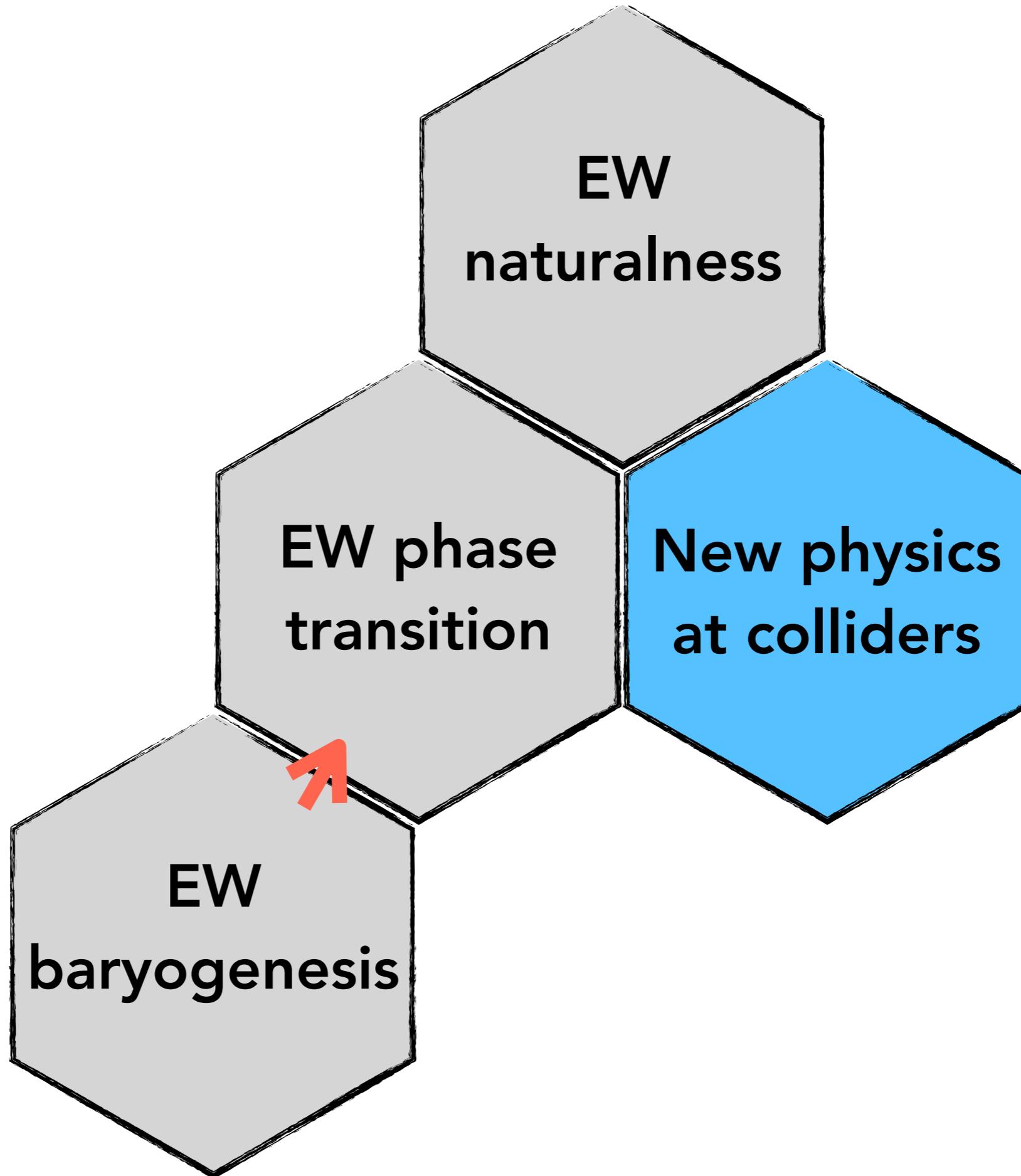
Technical University of Munich

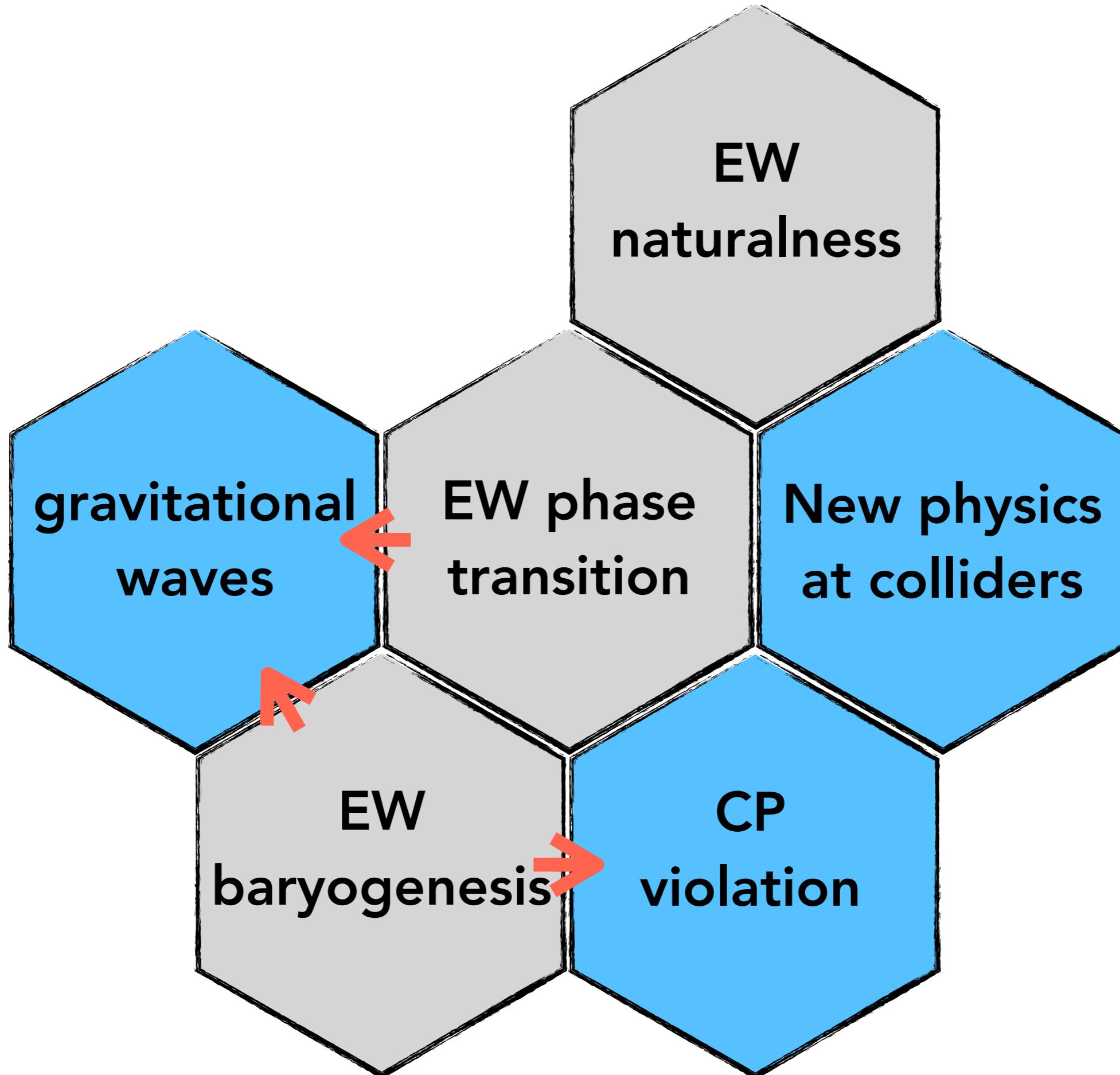
Higgs Hunting, 2025



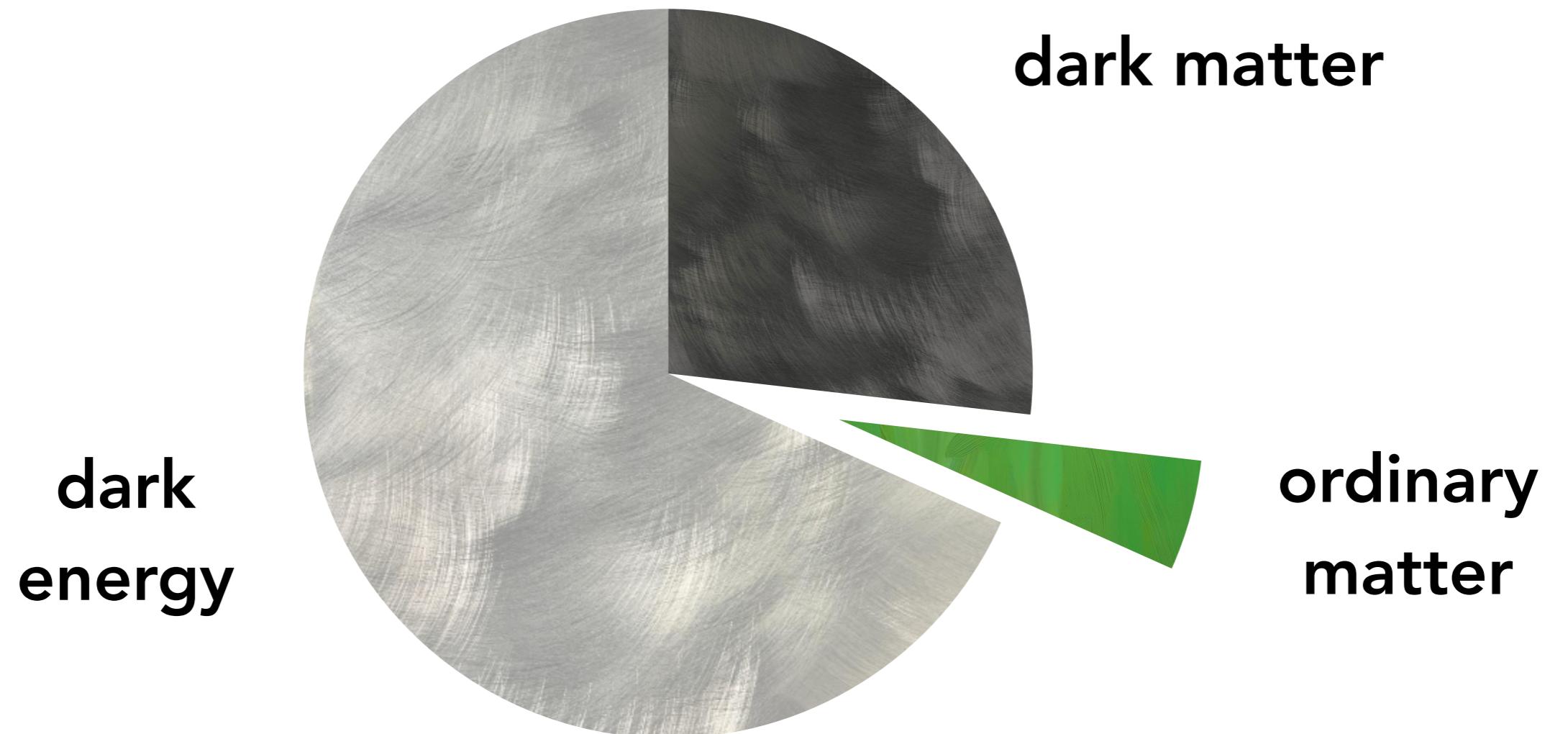
**New physics
at colliders**



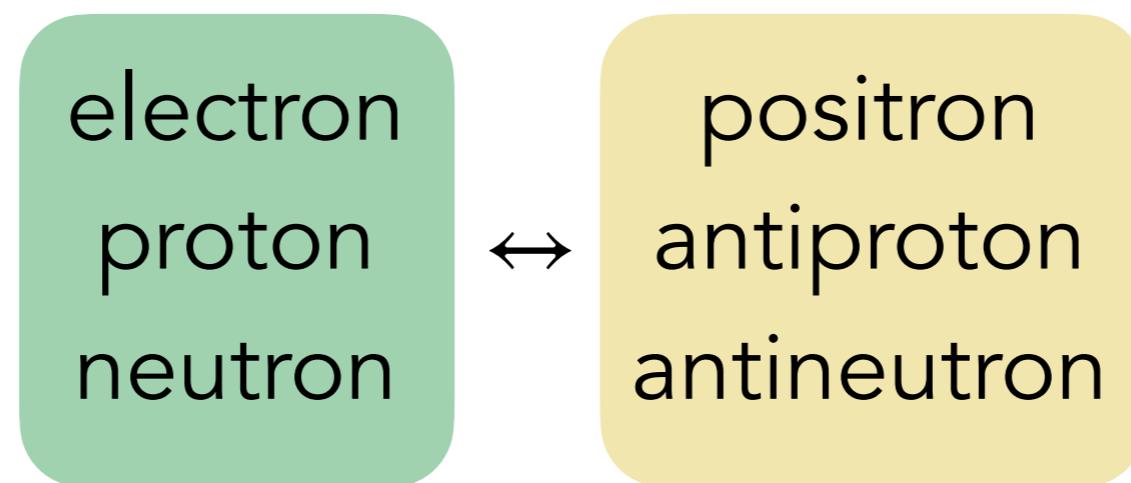




Matter-Antimatter Asymmetry



- QFT predicts particles to have partners - antiparticles



- But no such symmetry around us:

- we are made of matter
- in cosmic rays: $\bar{p}/p \sim 10^{-4}$

- cosmology: $\frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-9}$

baryons

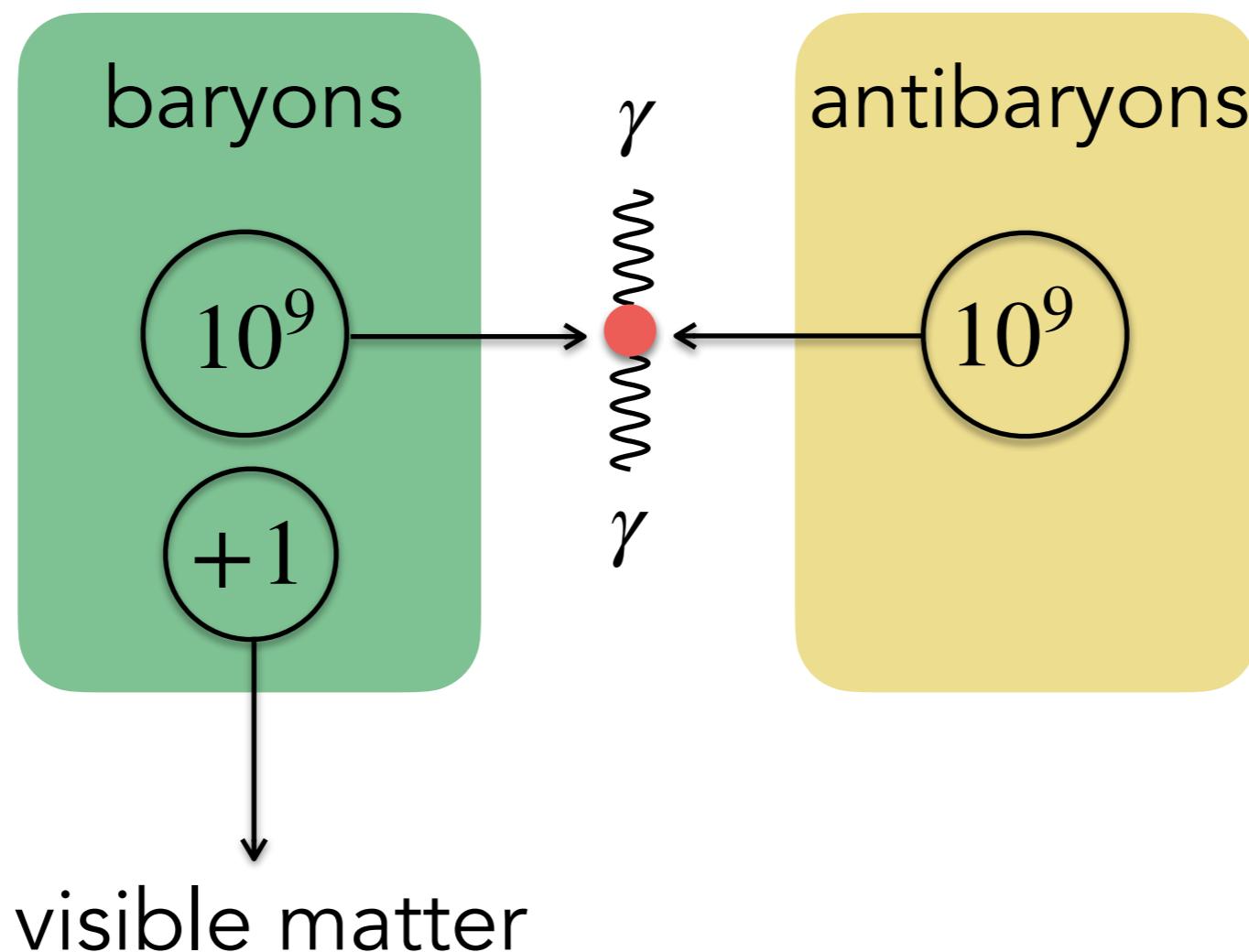
10^9

+1

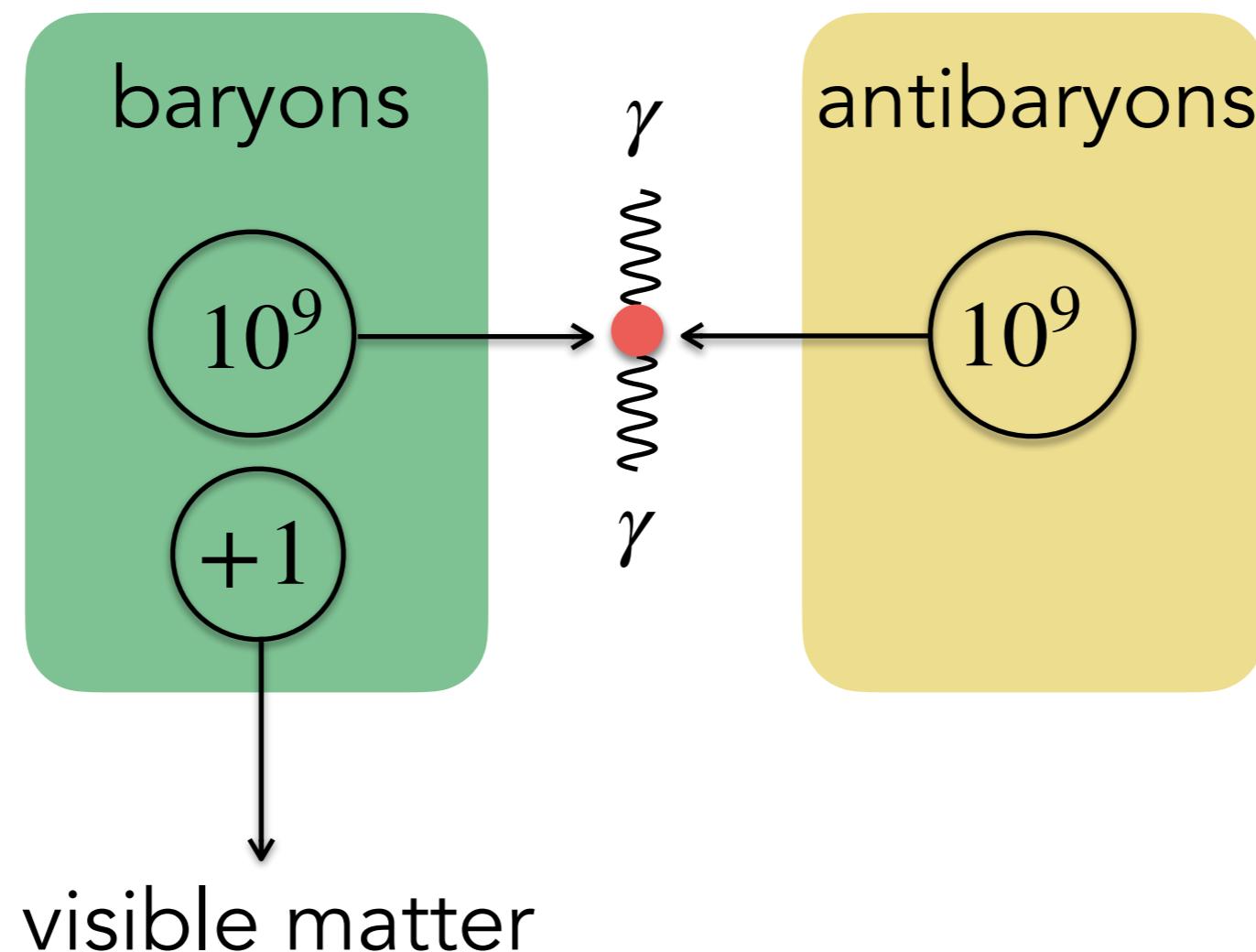
antibaryons

10^9

- cosmology: $\frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-9}$



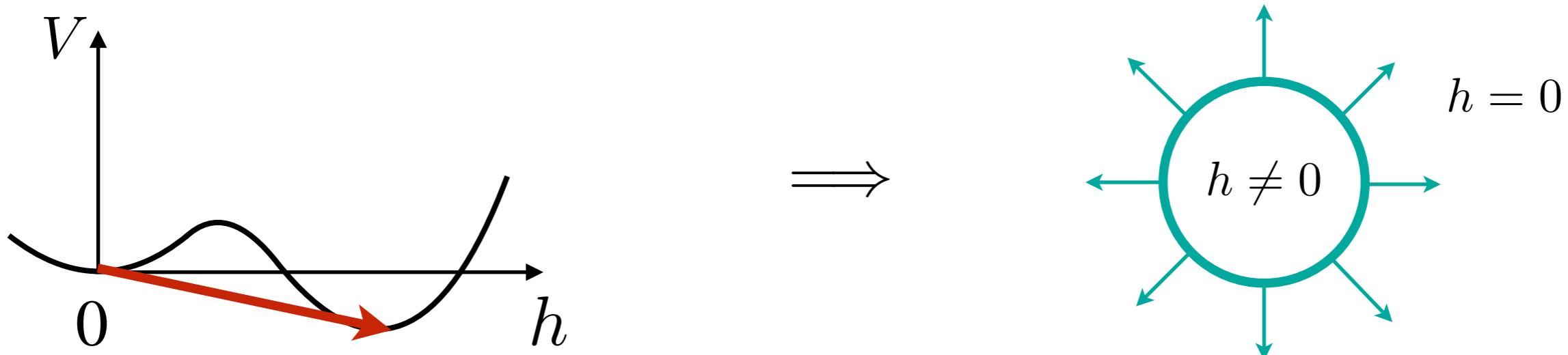
- cosmology: $\frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-9}$
- Which is ... **too large.**
No explanation within SM.



Electroweak Baryogenesis

Electroweak Baryogenesis

First order EW phase transition proceeds through bubble nucleation:

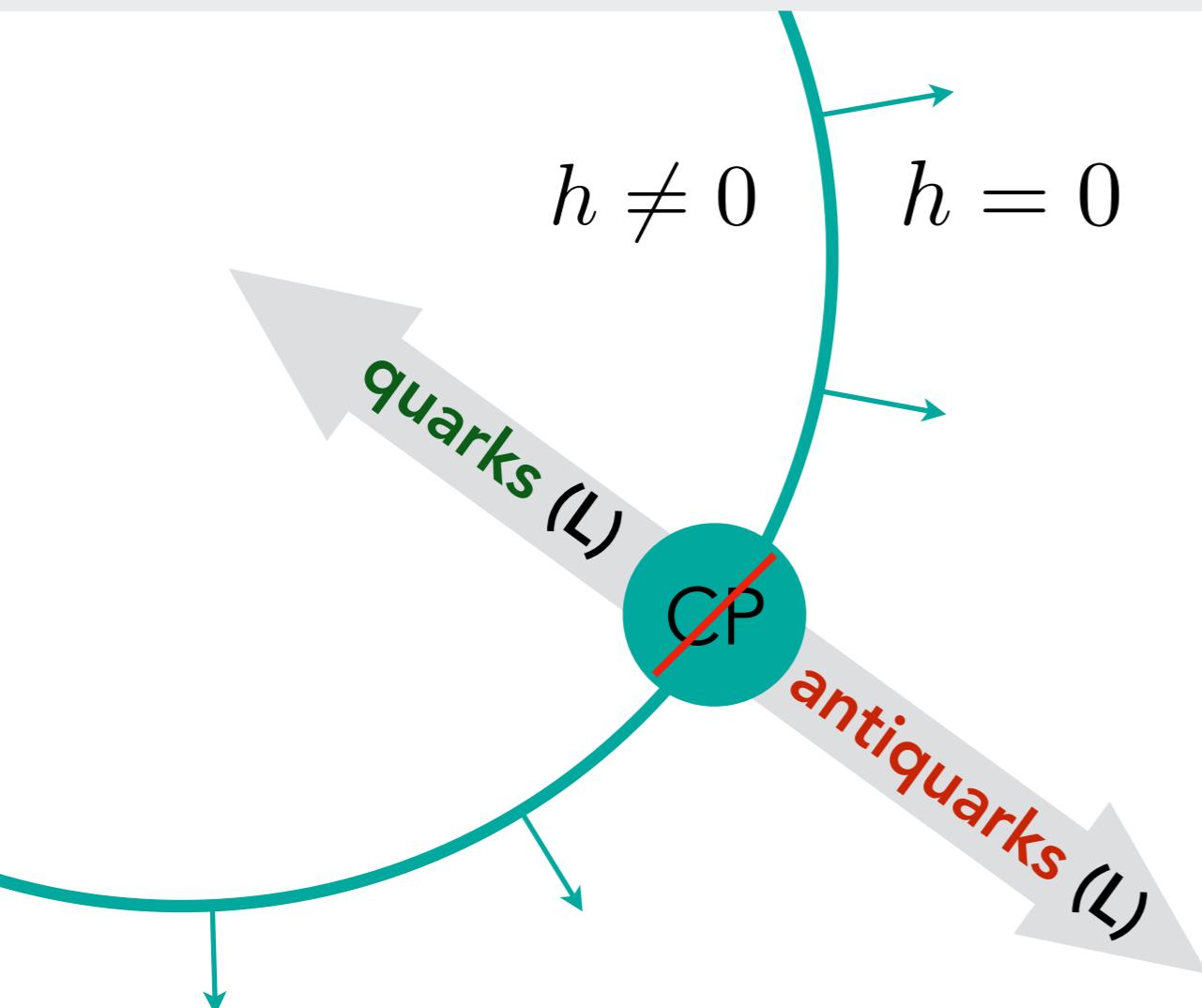


Shaposhnikov '87
Cohen,Kaplan,Nelson '91

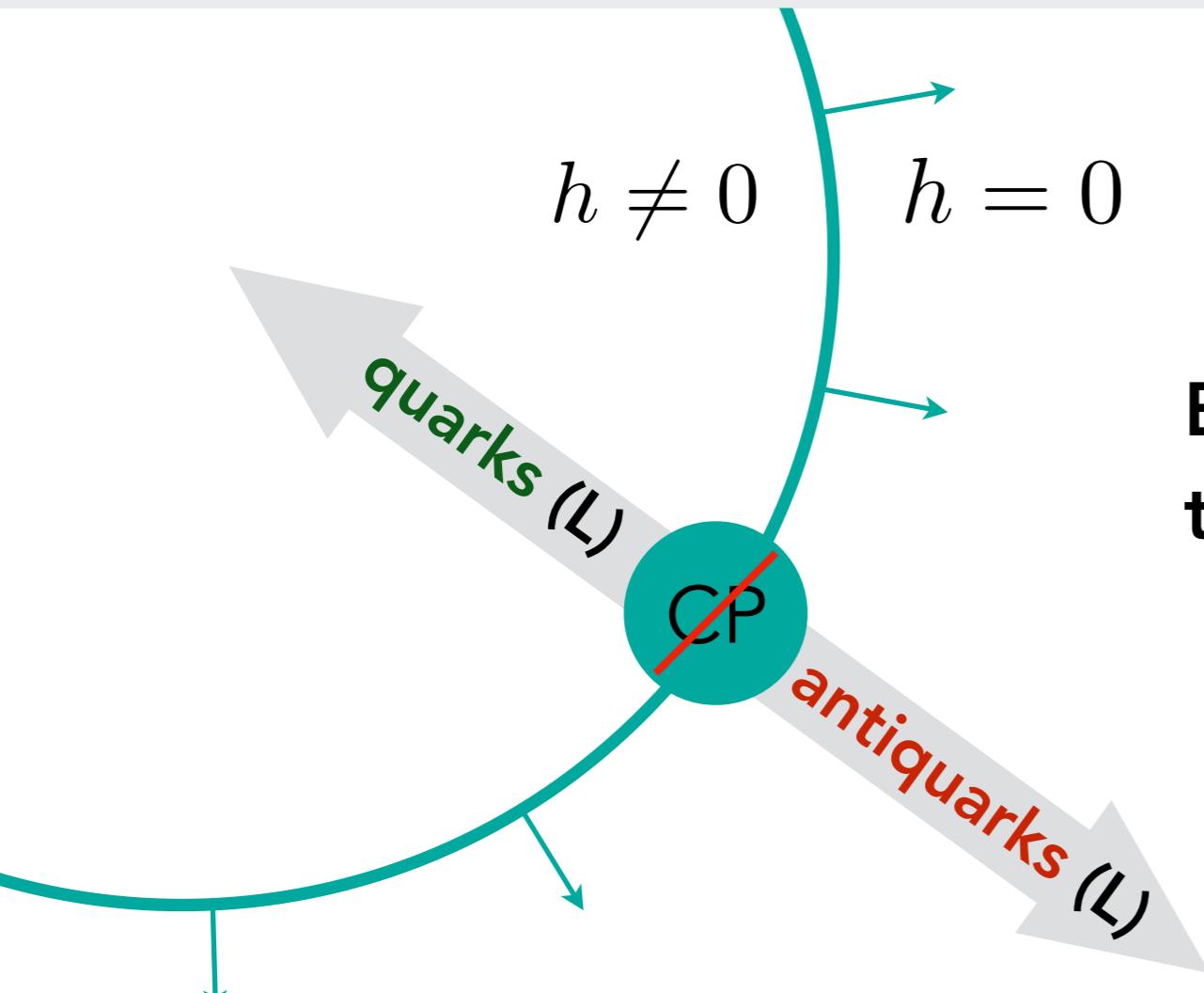
~



Electroweak Baryogenesis



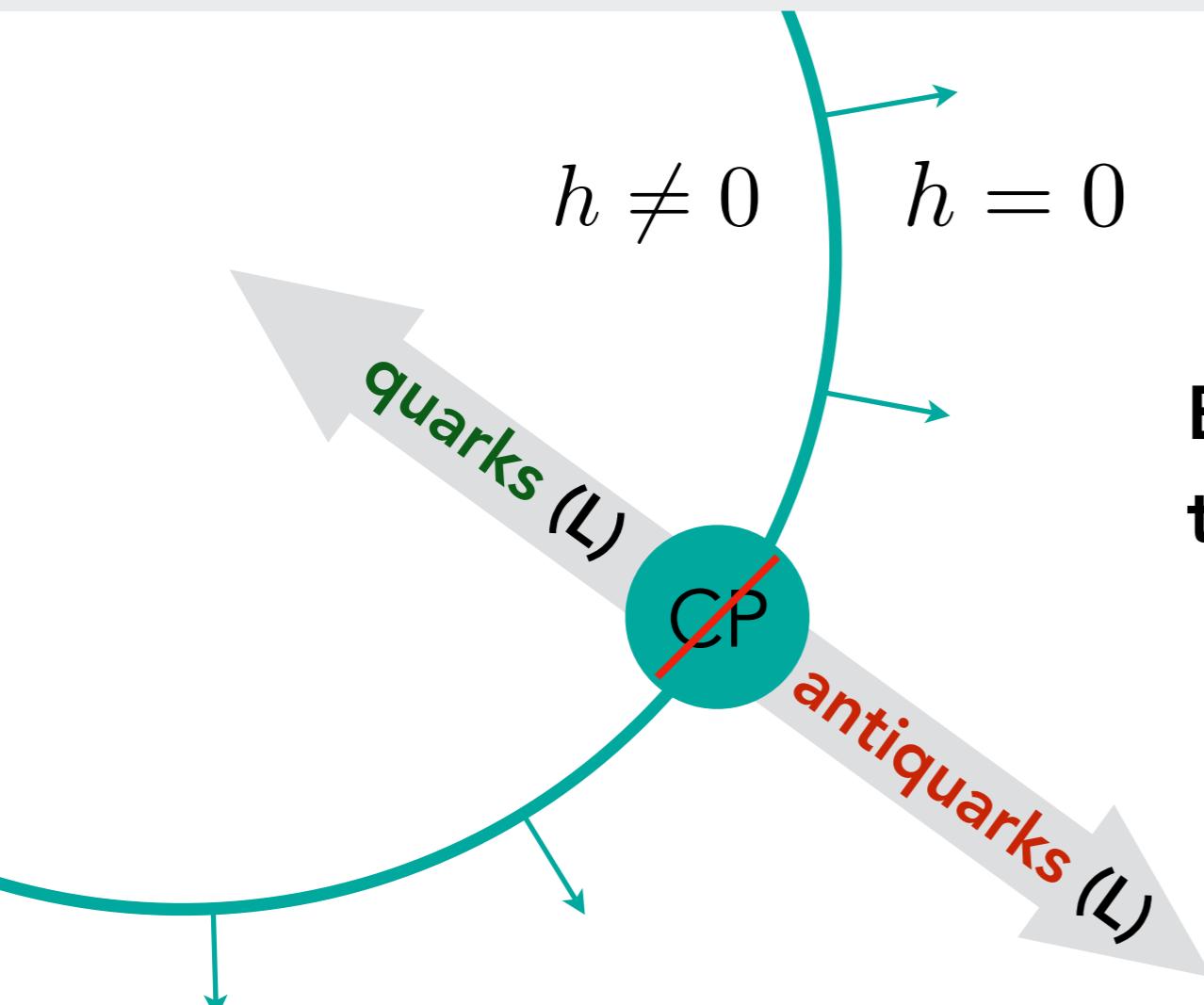
Electroweak Baryogenesis



Baryon asymmetry is created close to bubble walls:

- B+L violation from EW sphalerons
- $$\partial_\mu j_B^\mu = \partial_\mu j_l^\mu \\ = n_f \left(\frac{g^2}{32\pi^2} W_{\mu\nu}^a \tilde{W}^{a\mu\nu} - \frac{g'^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \right)$$

Electroweak Baryogenesis



Baryon asymmetry is created close to bubble walls:

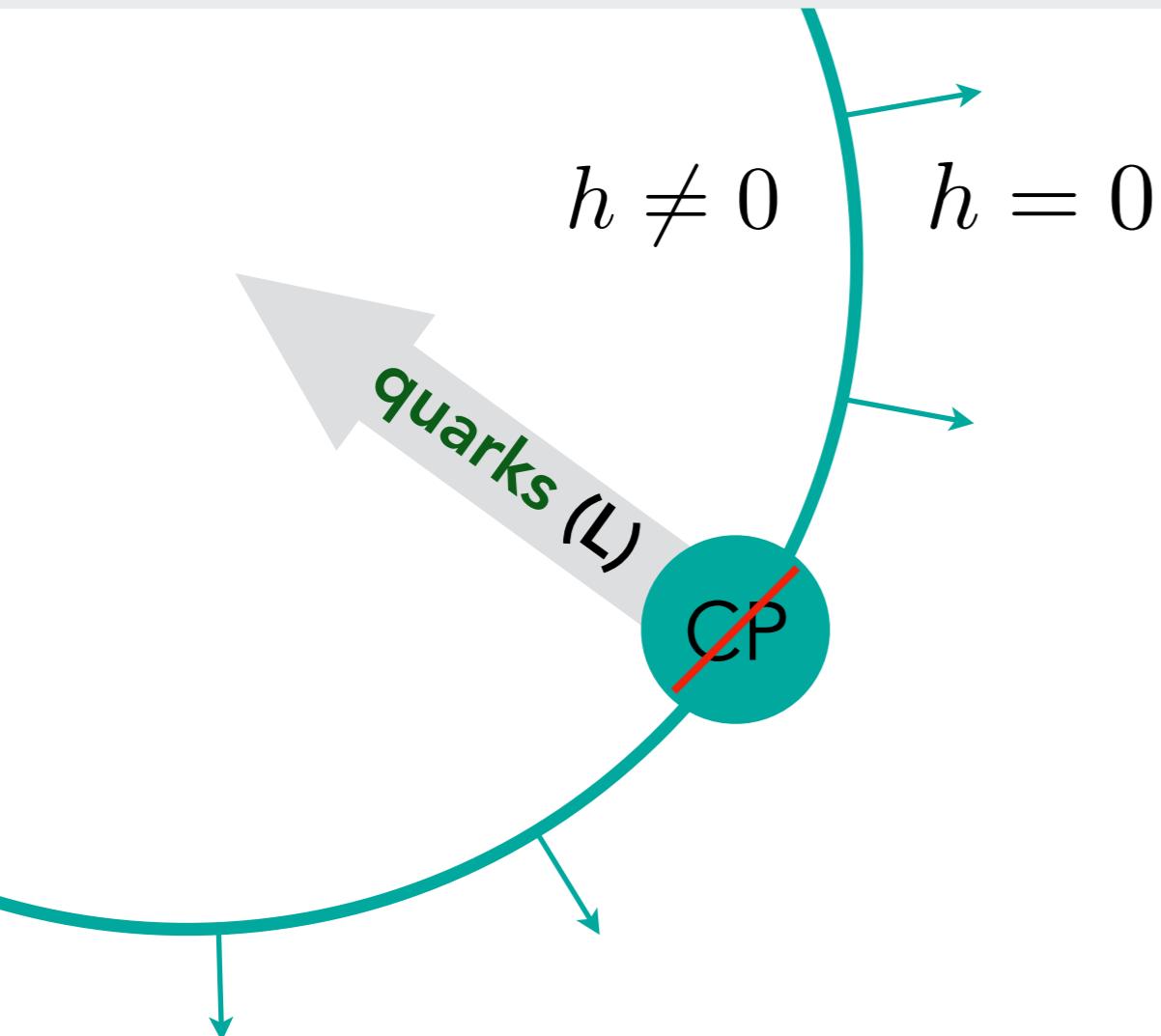
- B+L violation from EW sphalerons
- rate is Boltzmann-suppressed:

$$\sim \exp[-E_{\text{sph}}/T]$$

$$\underbrace{}_{\propto h}$$

⇒ convert -B into L at $h/T < 1$

Electroweak Baryogenesis



Baryon asymmetry is created close to bubble walls:

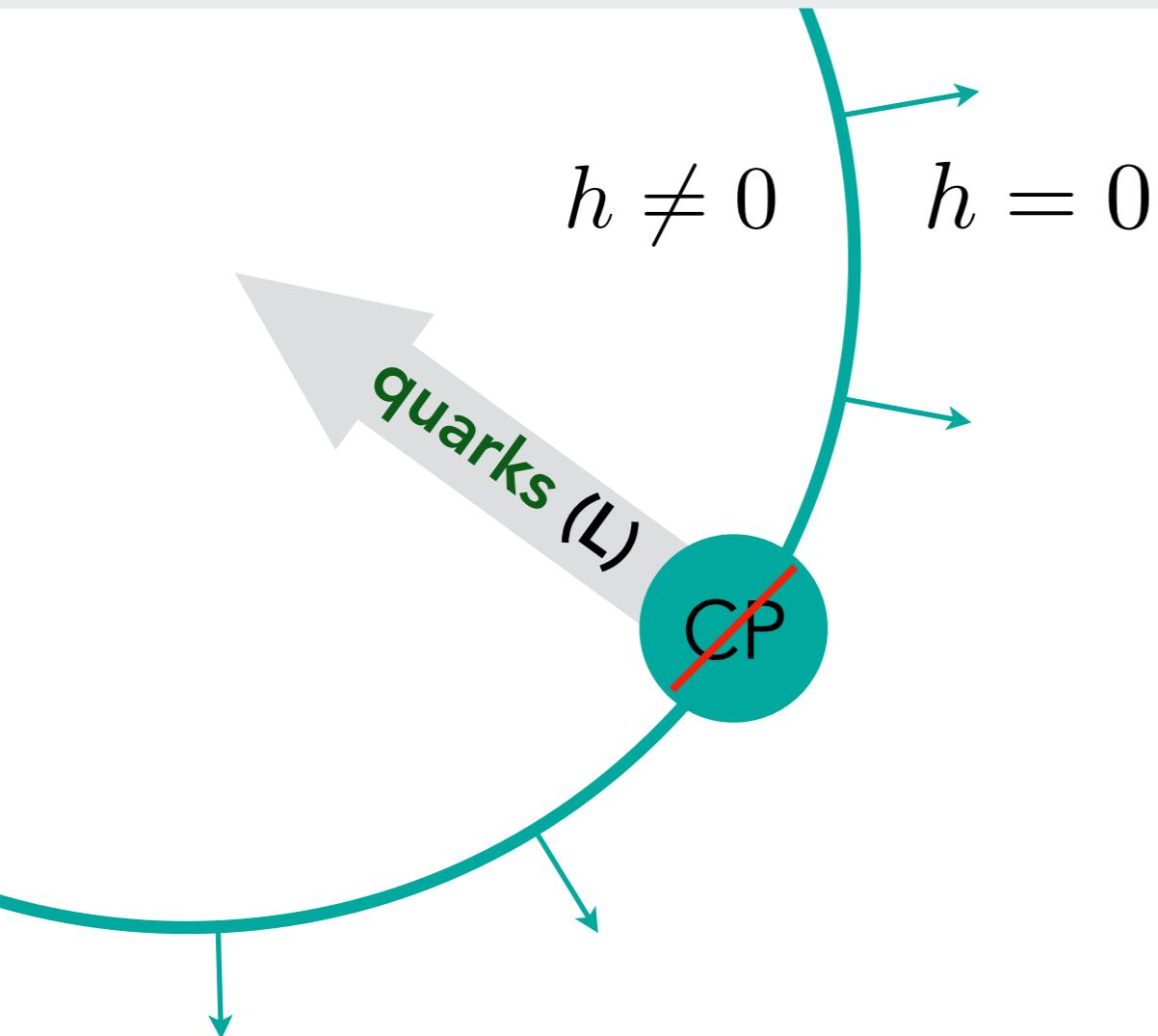
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$$\underbrace{}_{\propto h}$$

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Electroweak Baryogenesis



New Physics needed for:

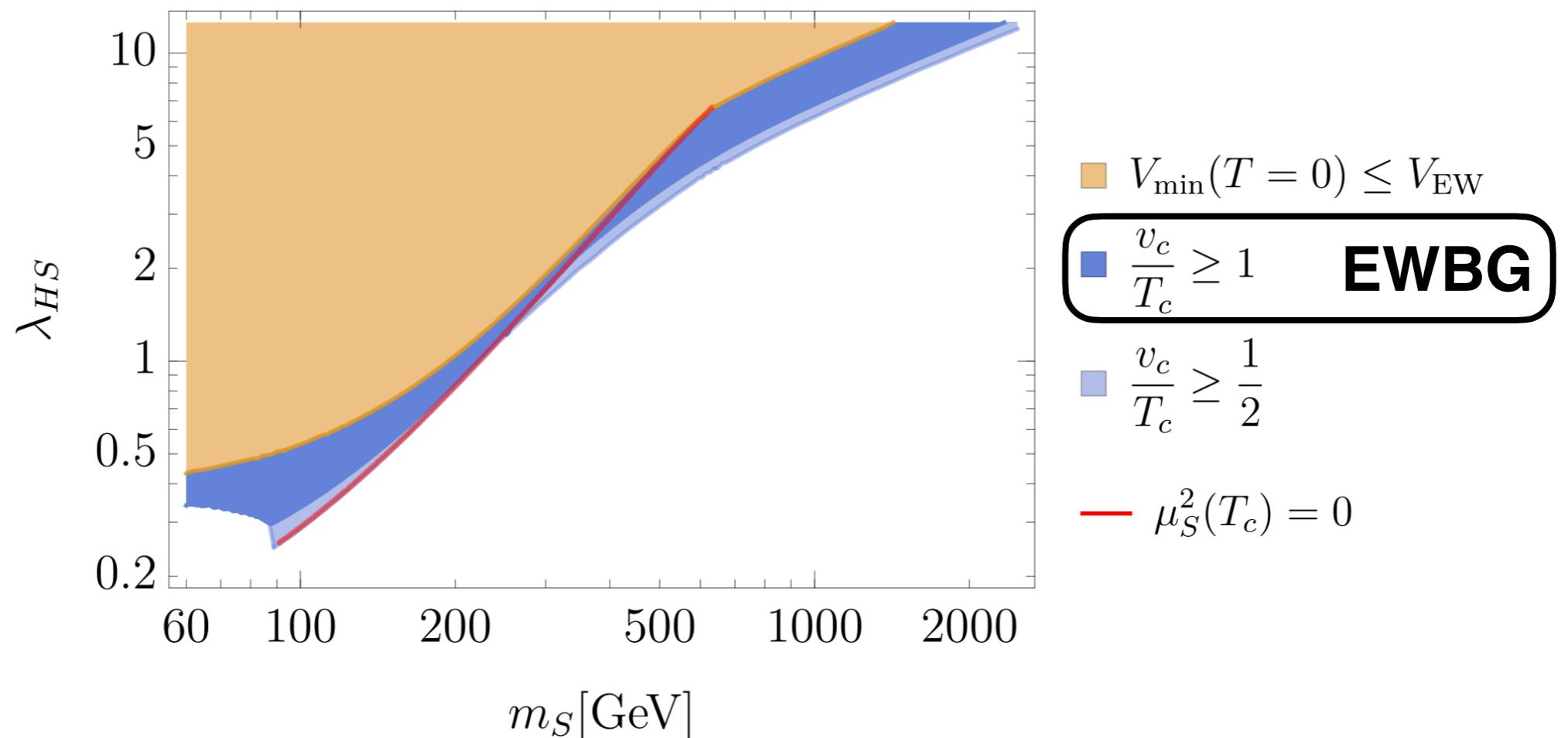
- strongly first-order EW phase transition
as compared to cross-over in SM
- CP-violation

Minimal EWBG Model

SM + Singlet

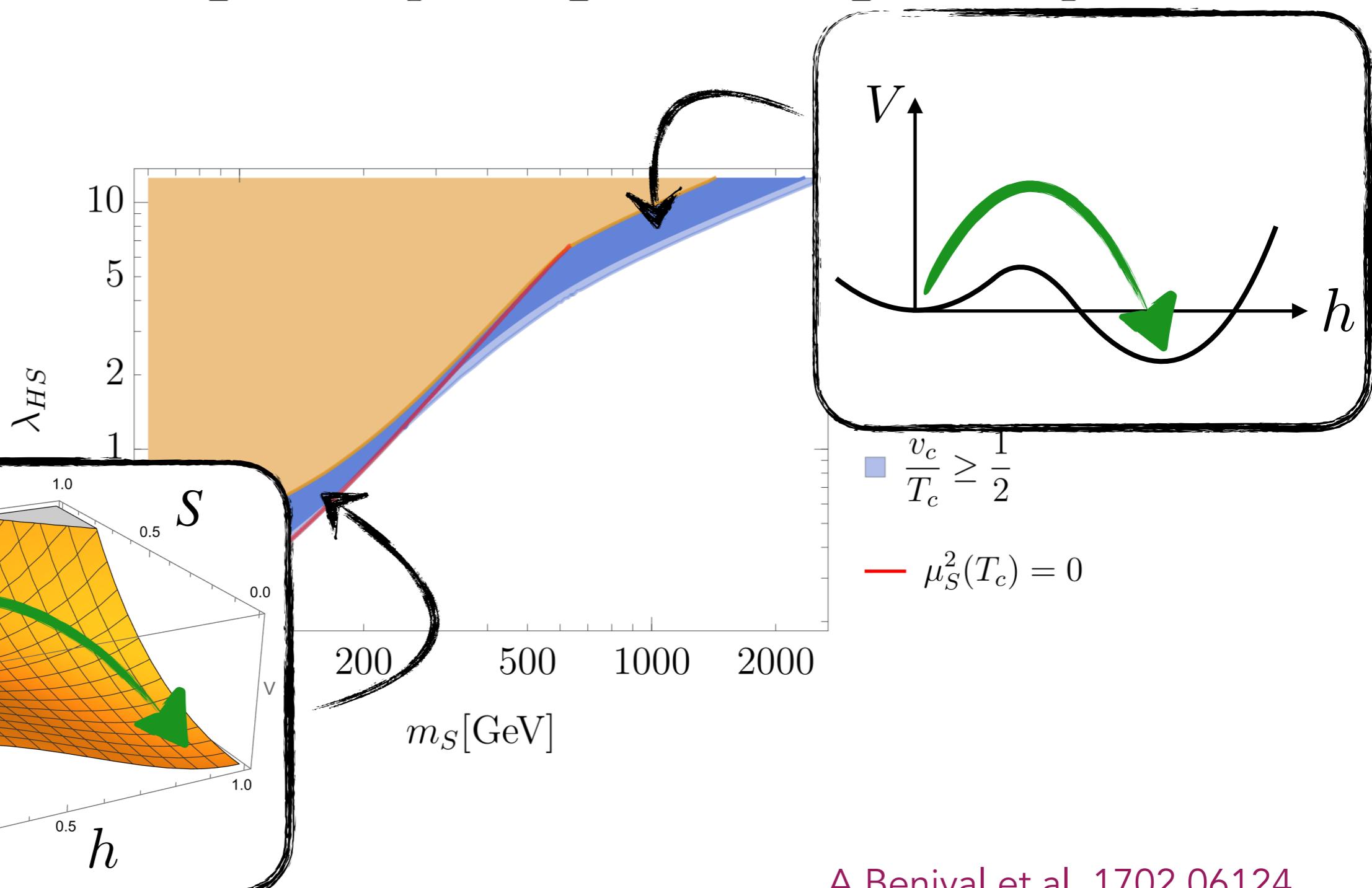
$$V_{\text{tree}}(h, S) = -\frac{1}{2}\mu^2 h^2 + \frac{1}{4}\lambda h^4 + \frac{1}{2}\lambda_{HS} h^2 S^2 + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{4}\lambda_S S^4$$

new



SM + Singlet

$$V_{\text{tree}}(h, S) = -\frac{1}{2}\mu^2 h^2 + \frac{1}{4}\lambda h^4 + \frac{1}{2}\lambda_{HS} h^2 S^2 + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{4}\lambda_S S^4$$



SM + Singlet

Pheno: S-h mixing

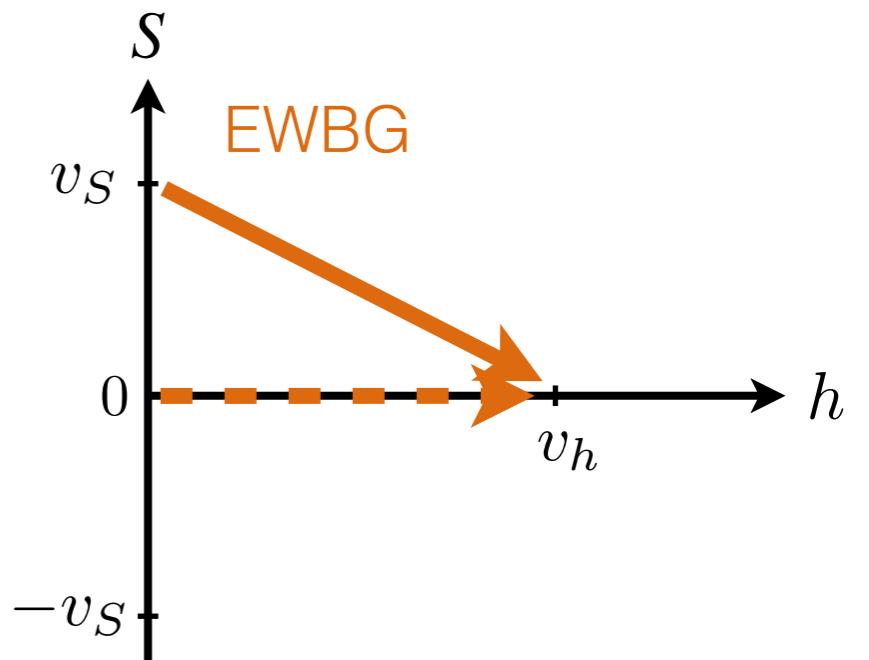
$$V_{\text{tree}}(h, S) = -\frac{1}{2}\mu^2 h^2 + \frac{1}{4}\lambda h^4 + \frac{1}{2}\lambda_{HS} h^2 S^2 + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{4}\lambda_S S^4$$

- $S \rightarrow -S$ symmetry:

⇒ no sizeable Higgs-S mixing

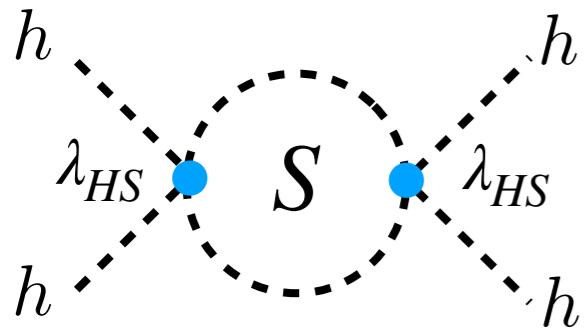
$$\sin \theta \propto \lambda_{HS} \langle h \rangle \langle S \rangle$$

⇒ search for loop-induced effects of λ_{HS}



SM + Singlet

Pheno: c_H



$$\mathcal{L} \ni \frac{c_H}{\Lambda^2} \mathcal{O}_H \quad \text{where} \quad \mathcal{O}_H = \frac{1}{2} (\partial_\mu |H|^2)^2$$

$$\frac{c_H}{\Lambda^2} = \frac{\lambda_{HS}^2}{48\pi^2} \frac{1}{m_S^2}$$

M.Carena et al, 2104.00638

future sensitivities (1σ):

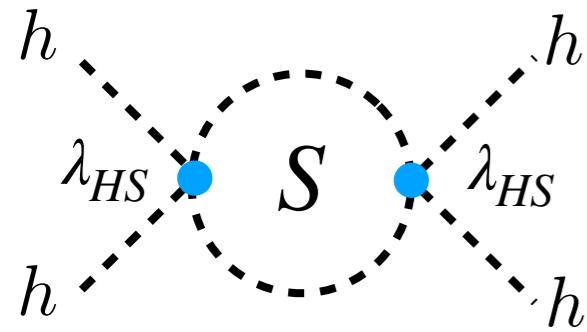
HL-LHC: $\Lambda/\sqrt{|c_H|} < 1.4(1.8) \text{ TeV}$

+FCC-ee: $\Lambda/\sqrt{|c_H|} < 3.2(5) \text{ TeV}$

J de Blas, Eur. Phys. J. Plus (2021) 136:897

SM + Singlet

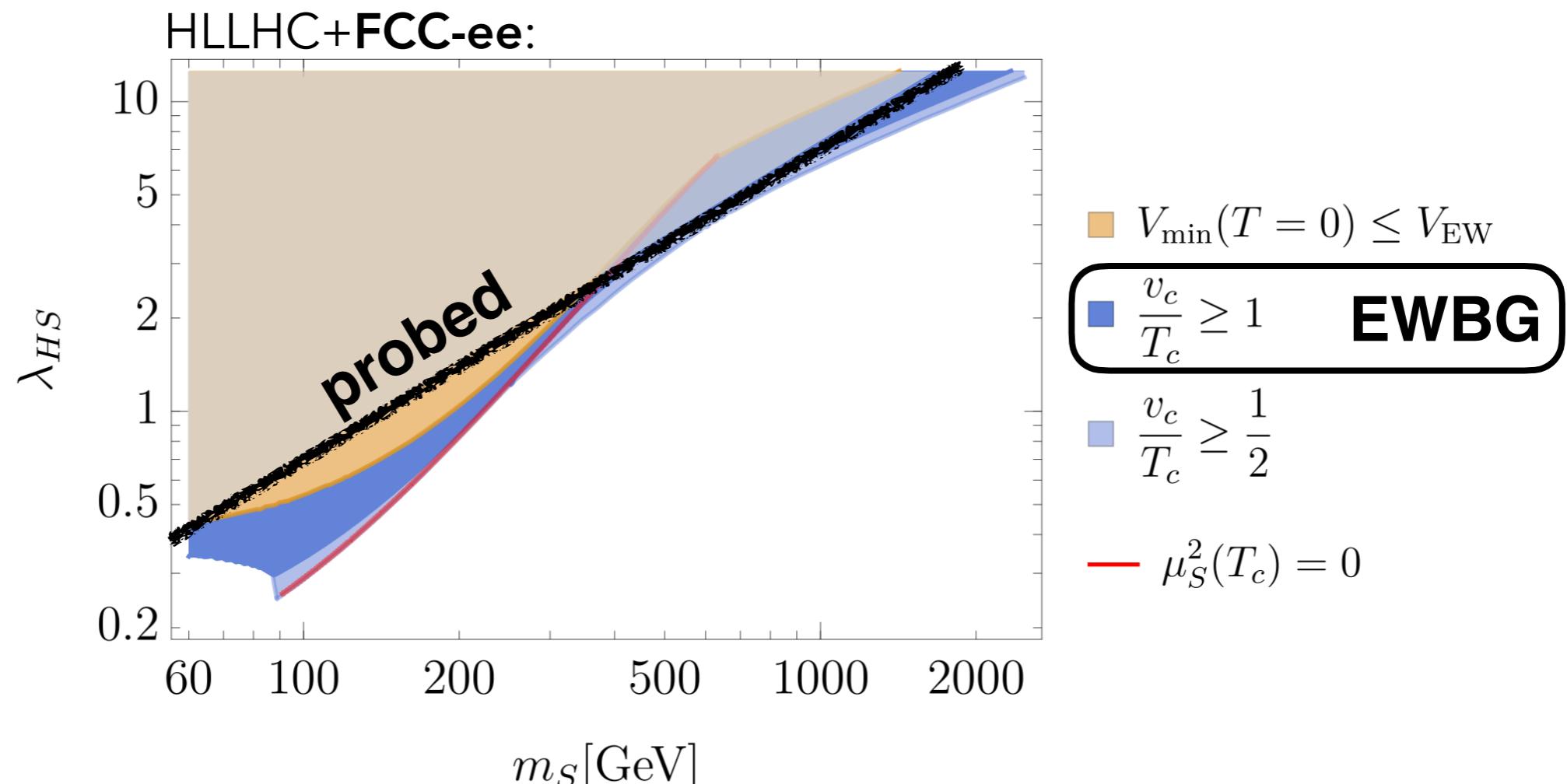
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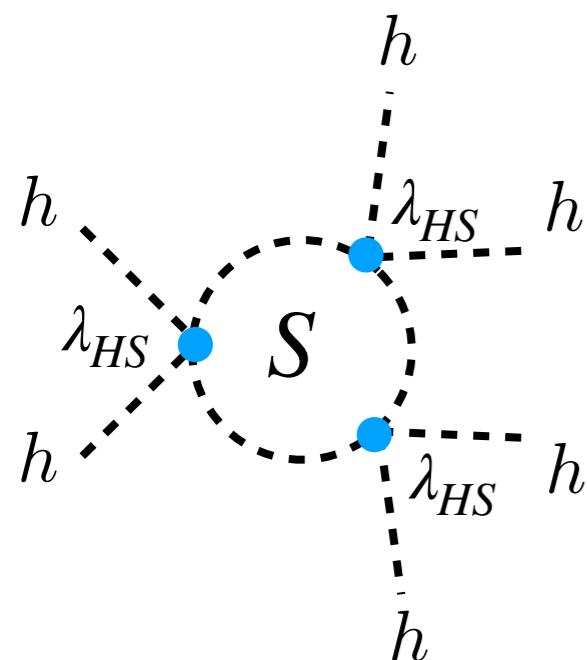
$$\frac{c_H}{\Lambda^2} = \frac{\lambda_{HS}^2}{48\pi^2} \frac{1}{m_S^2}$$

M.Carena et al, 2104.00638



SM + Singlet

Pheno: h^3

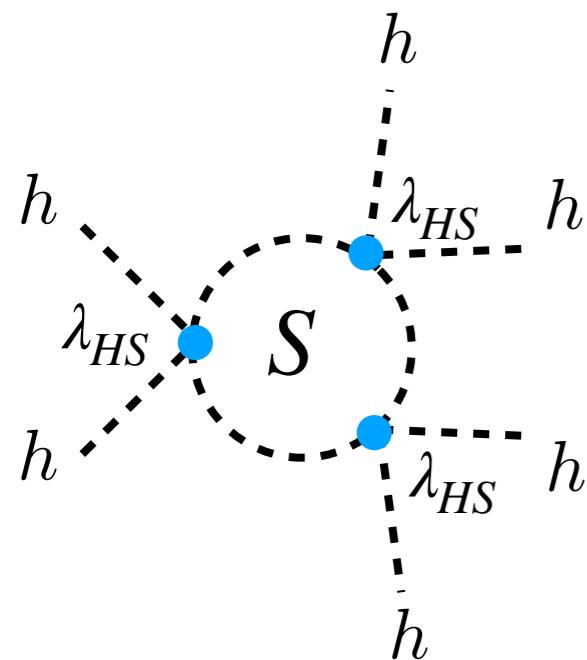


$$\lambda_3 = \frac{1}{6} \frac{\partial^3 V(h, S = 0, T = 0)}{\partial h^3} \Big|_{h=v_0} \approx \frac{m_h^2}{2v_0} + \frac{\lambda_{HS}^3 v_0^3}{24\pi^2 m_S^2}$$

A.Benival et al, 1702.06124

SM + Singlet

Pheno: h^3



$$\lambda_3 = \frac{1}{6} \frac{\partial^3 V(h, S = 0, T = 0)}{\partial h^3} \Big|_{h=v_0} \approx \frac{m_h^2}{2v_0} + \frac{\lambda_{HS}^3 v_0^3}{24\pi^2 m_S^2}$$

A.Benival et al, 1702.06124

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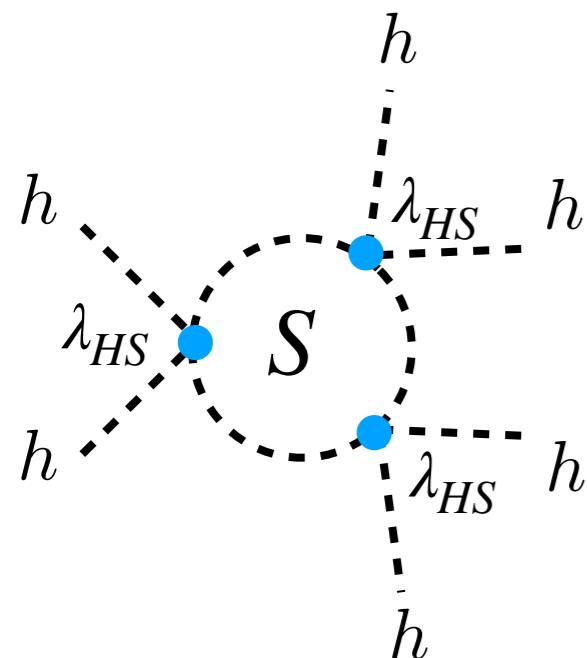
HL-LHC: $\delta\kappa \sim 0.5$

+FCC-ee: $\delta\kappa \sim 0.2 - 0.3$

J de Blas et al, 1905.03764

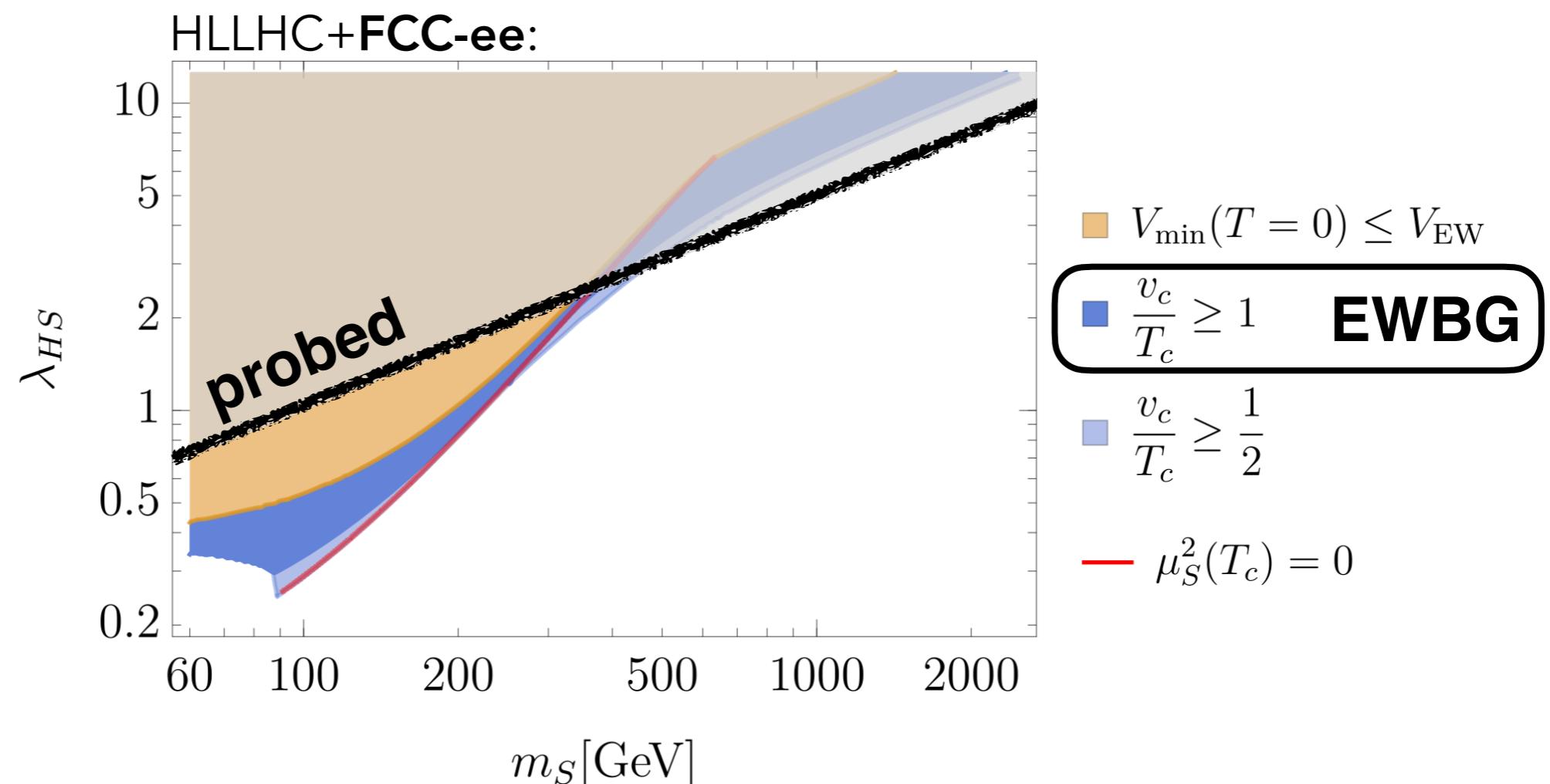
SM + Singlet

Pheno: h^3



$$\lambda_3 = \frac{1}{6} \frac{\partial^3 V(h, S = 0, T = 0)}{\partial h^3} \Big|_{h=v_0} \approx \frac{m_h^2}{2v_0} + \frac{\lambda_{HS}^3 v_0^3}{24\pi^2 m_S^2}$$

A.Benival et al, 1702.06124

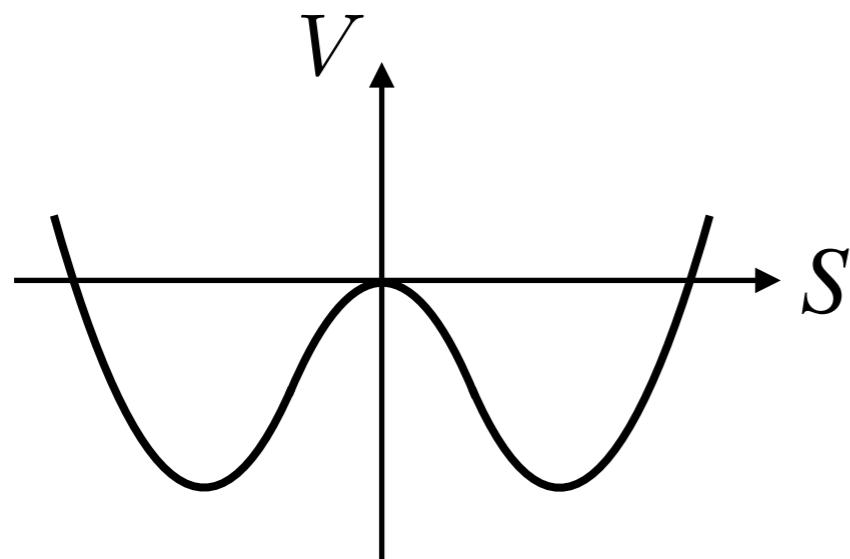


Minimal *Defect-Mediated EWBG Model*

SM + Singlet: EWBG via Domain Walls

- Brandenberger, Davis, Prokopec, Trodden
Phys. Rev. D 53 (1996) 4257–4266
- J.Azzola, OM, A.Weiler JHEP 04 (2025)

discrete symmetry of
vacuum manifold

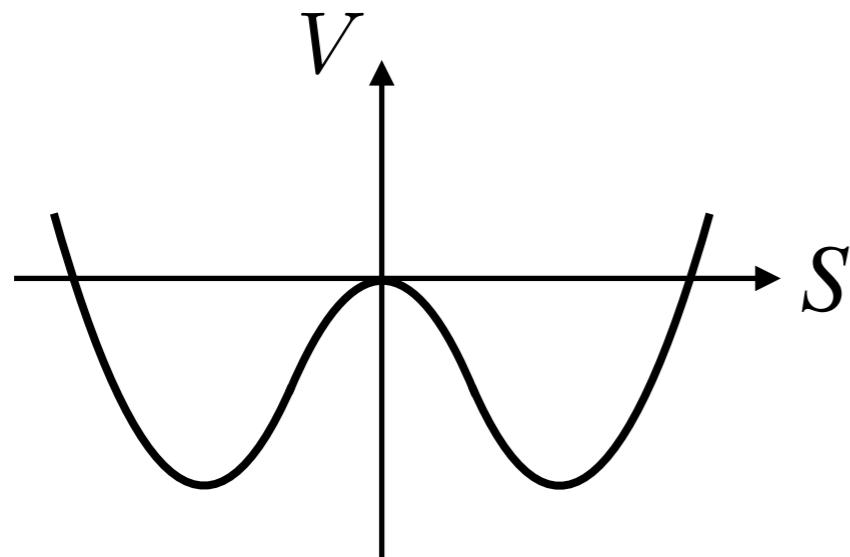


$$S = v_S \tanh \frac{m_S z}{2}$$

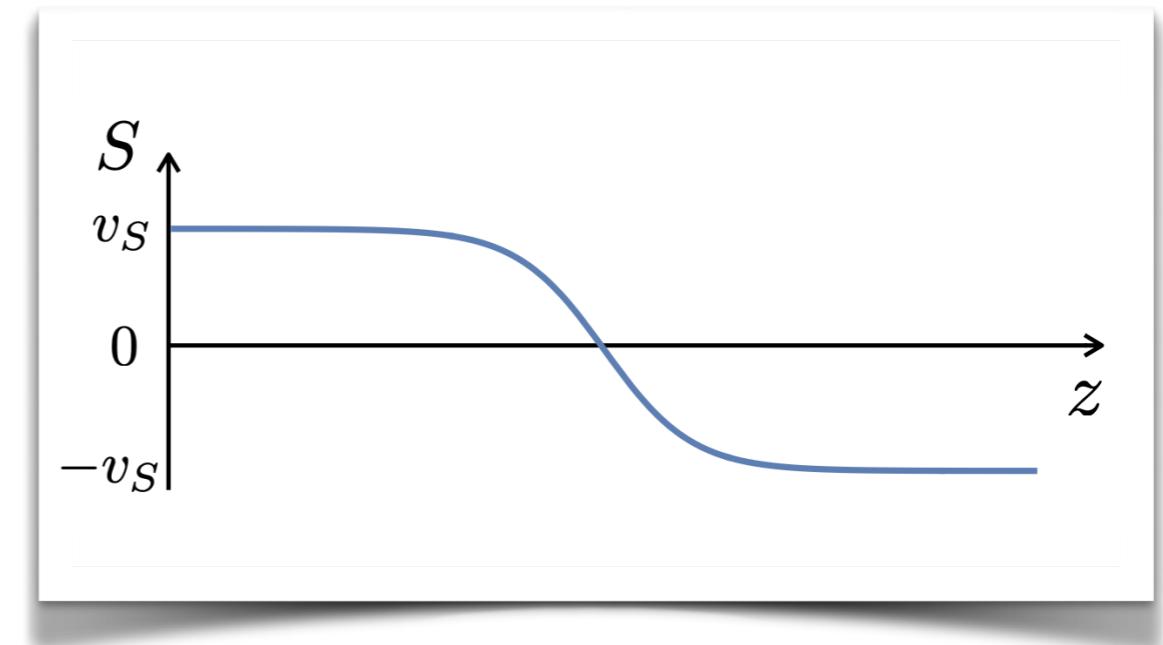
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discrete symmetry of
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domain walls can be
formed in Z2 breaking
phase transition

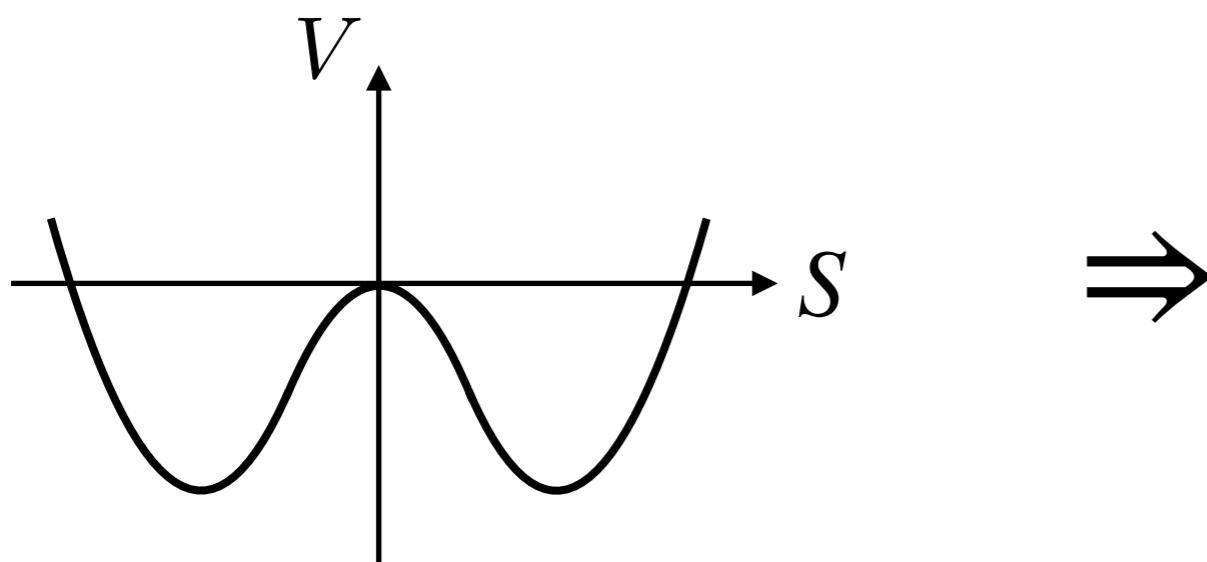


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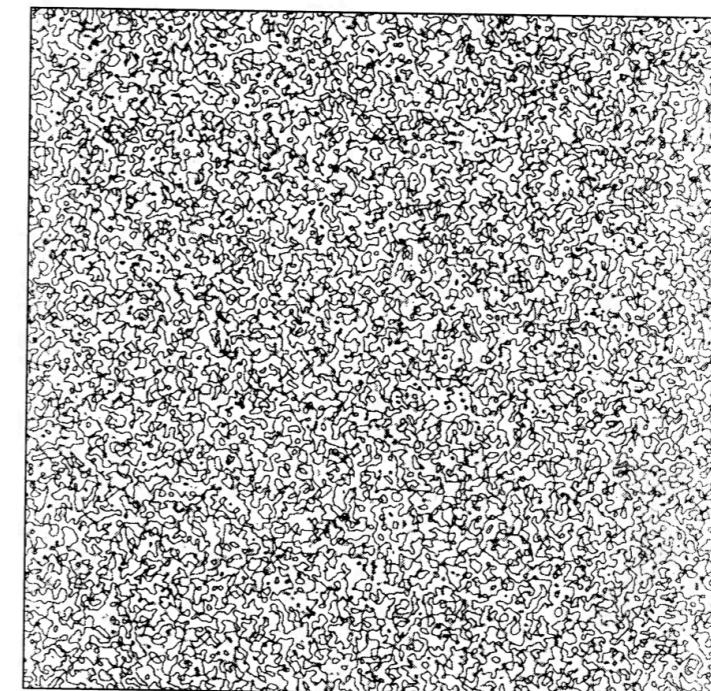
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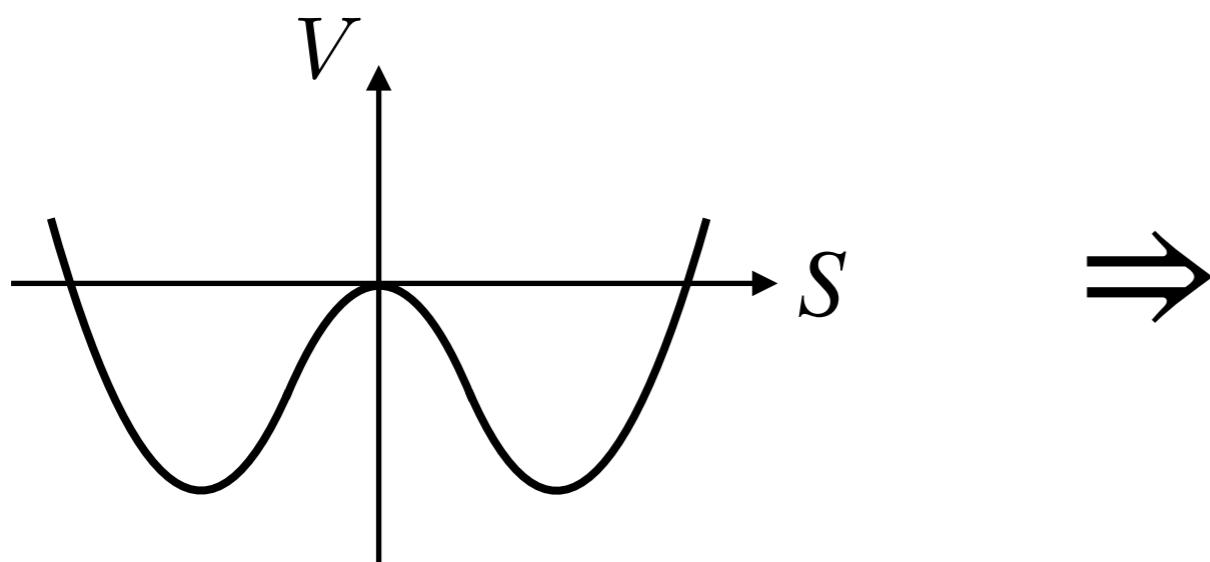
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Press, Ryden, Spergel,
Astrophys.J. 347 (1989)

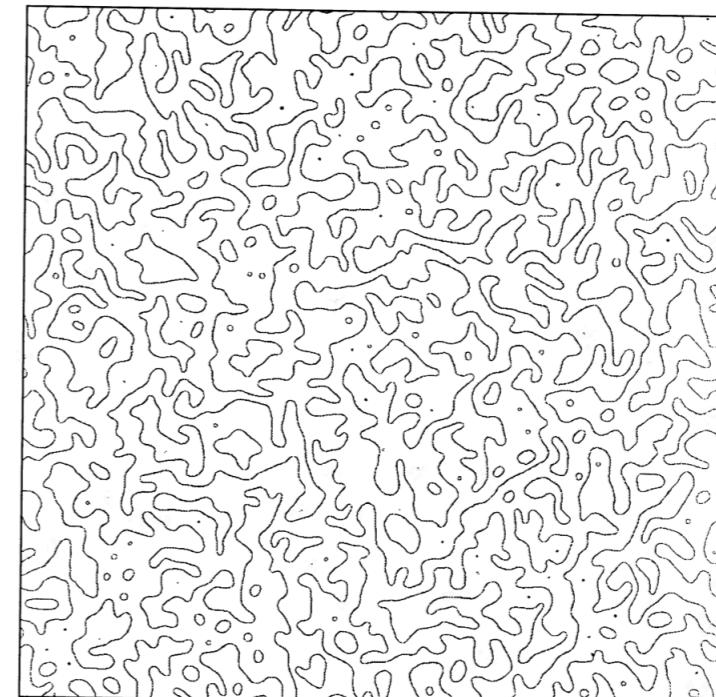
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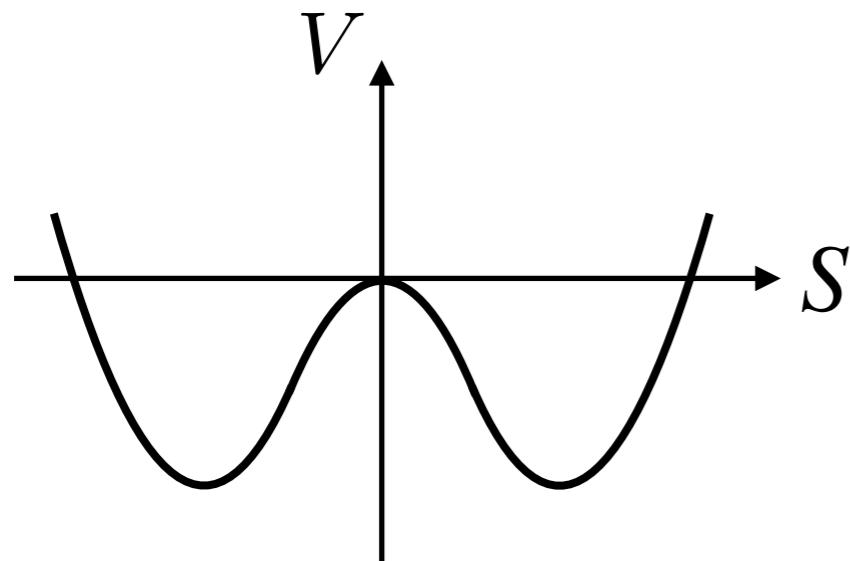
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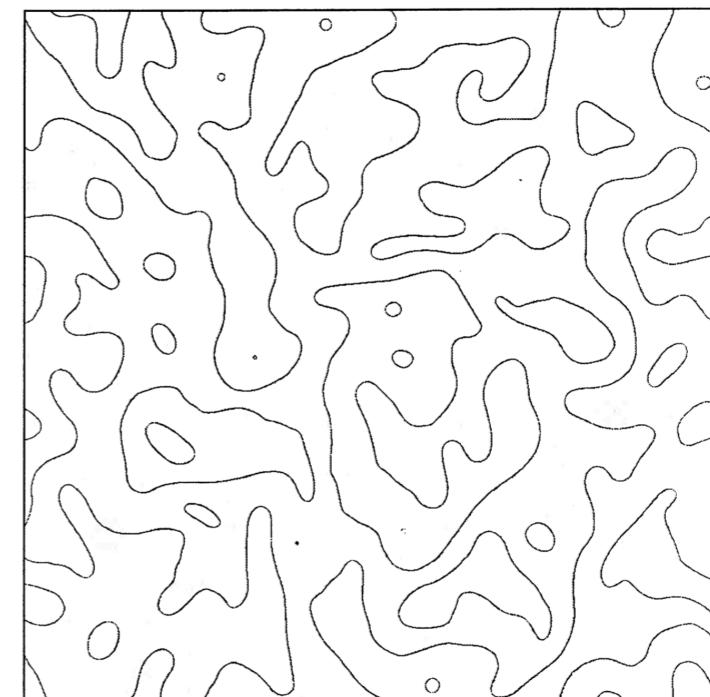
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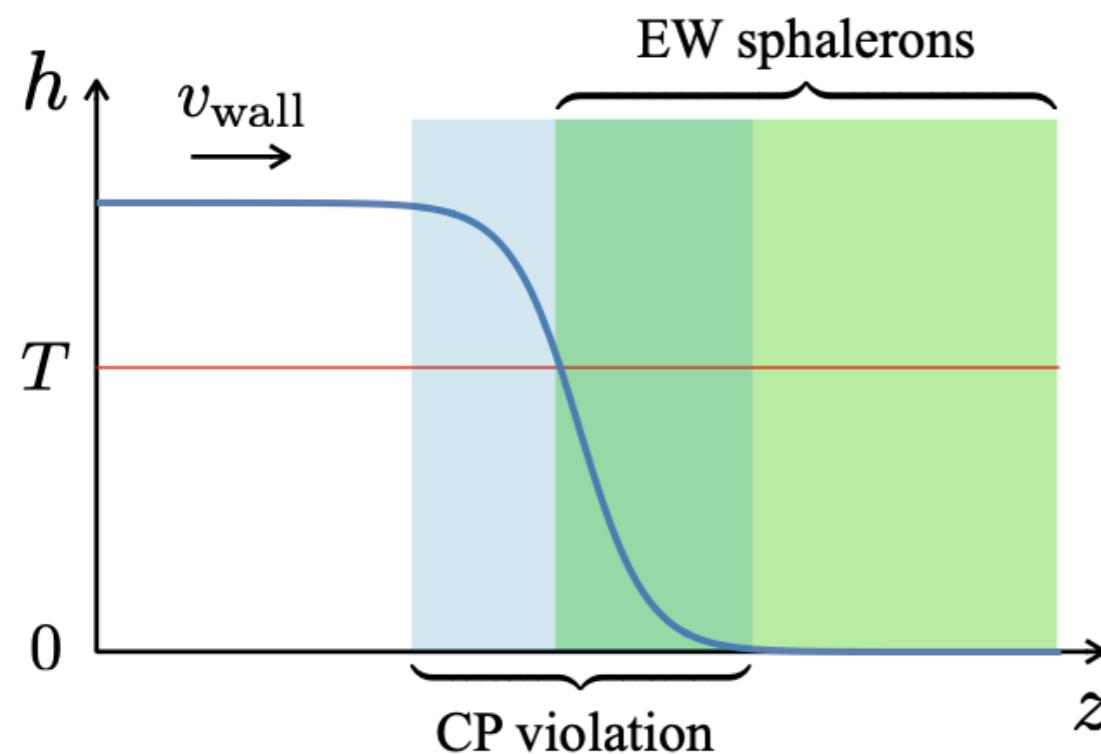


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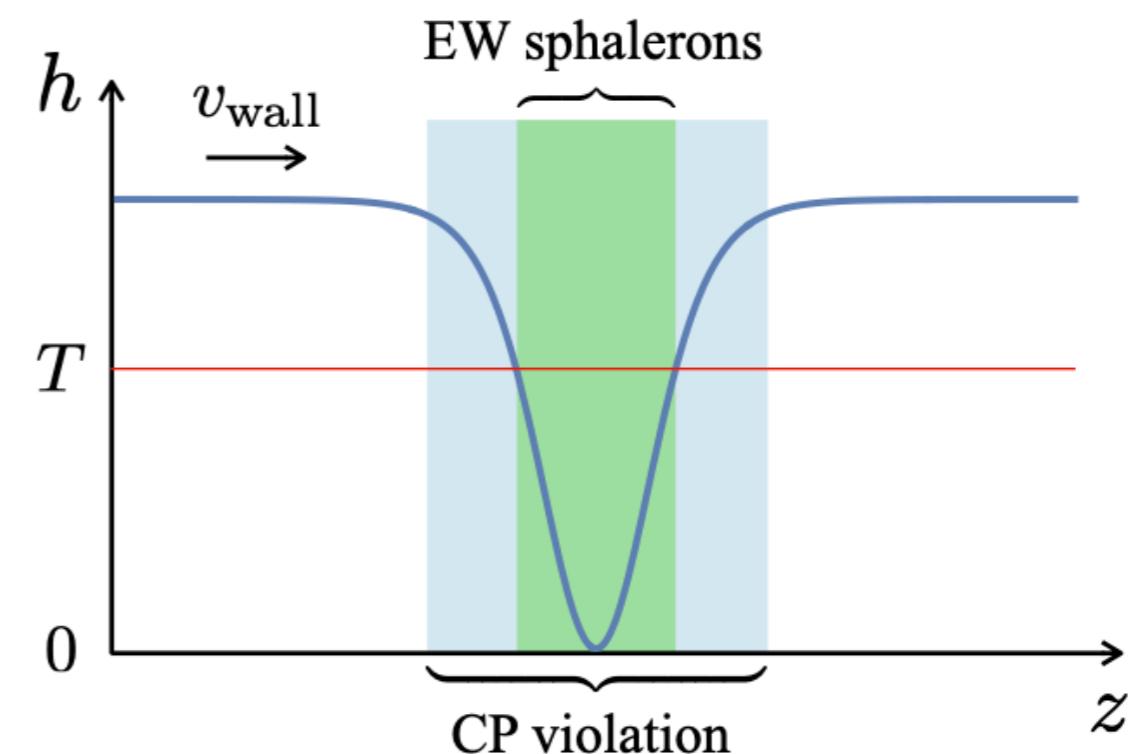
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Astrophys.J. 347 (1989)

SM + Singlet: EWBG via Domain Walls

1st order EWPT



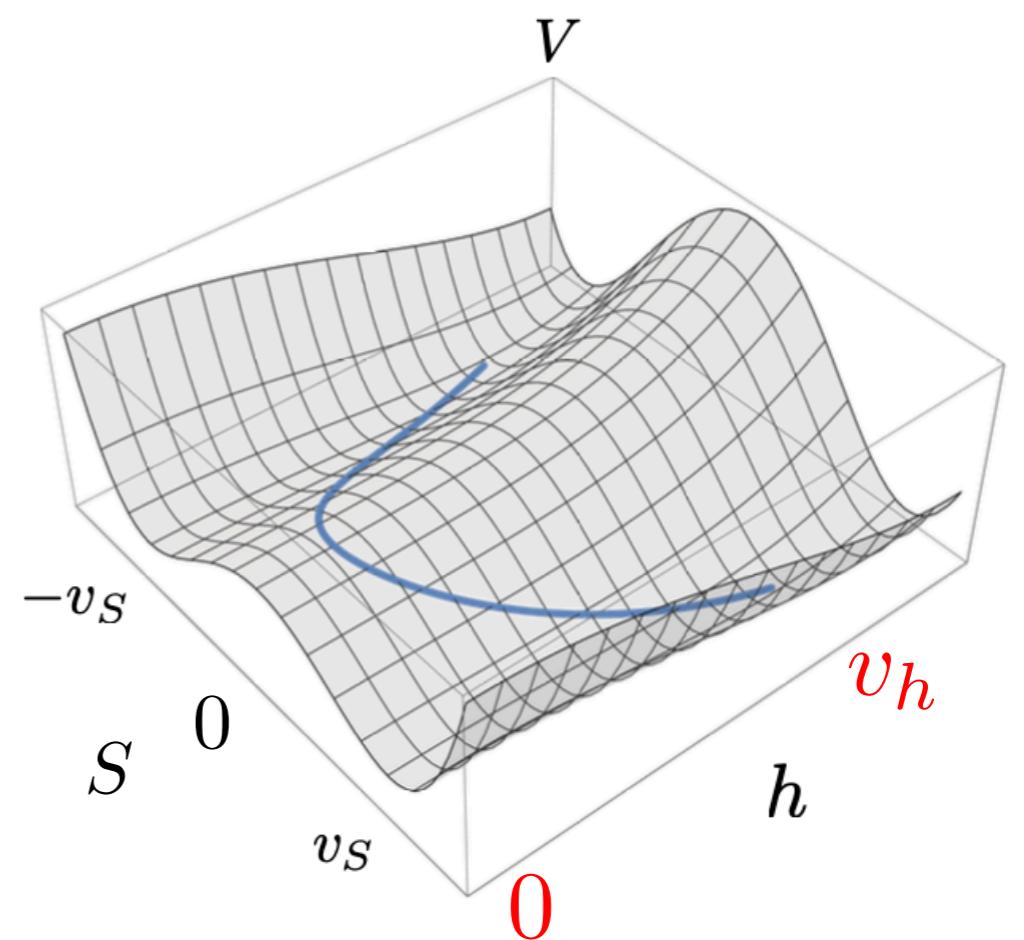
EW restoration in walls



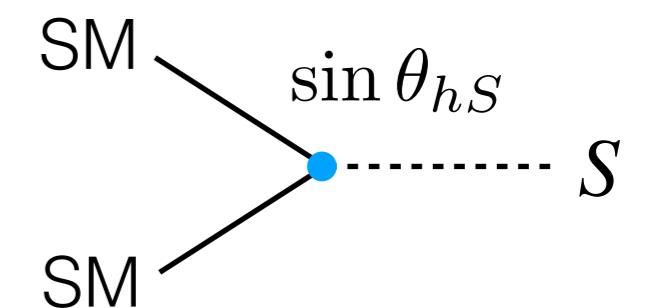
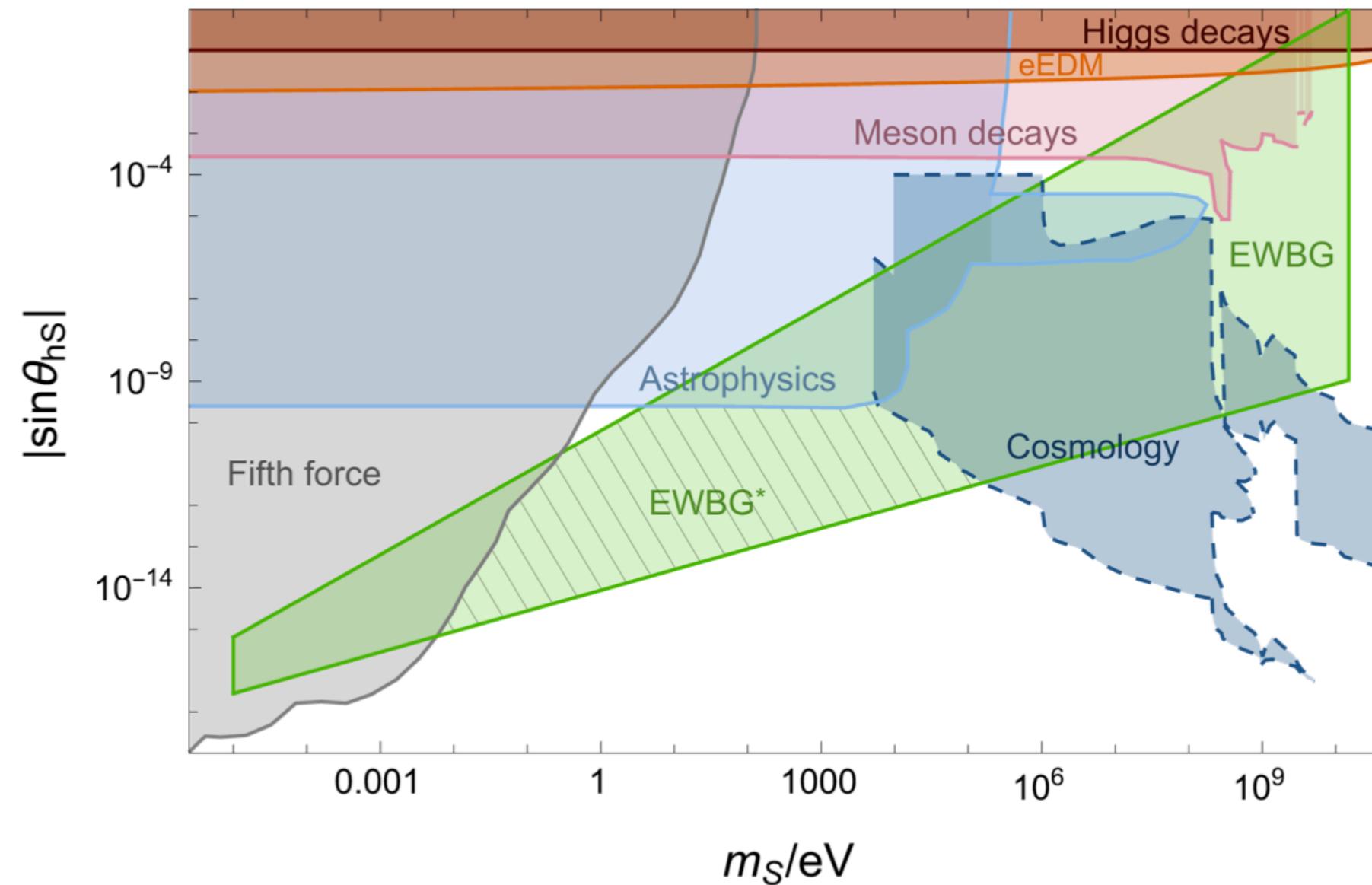
SM + Singlet: EWBG via Domain Walls

- h-S trajectory curved towards h=0

$$\mathcal{L} \supset -\frac{1}{2} |\lambda_{HS}| |H|^2 S^2$$

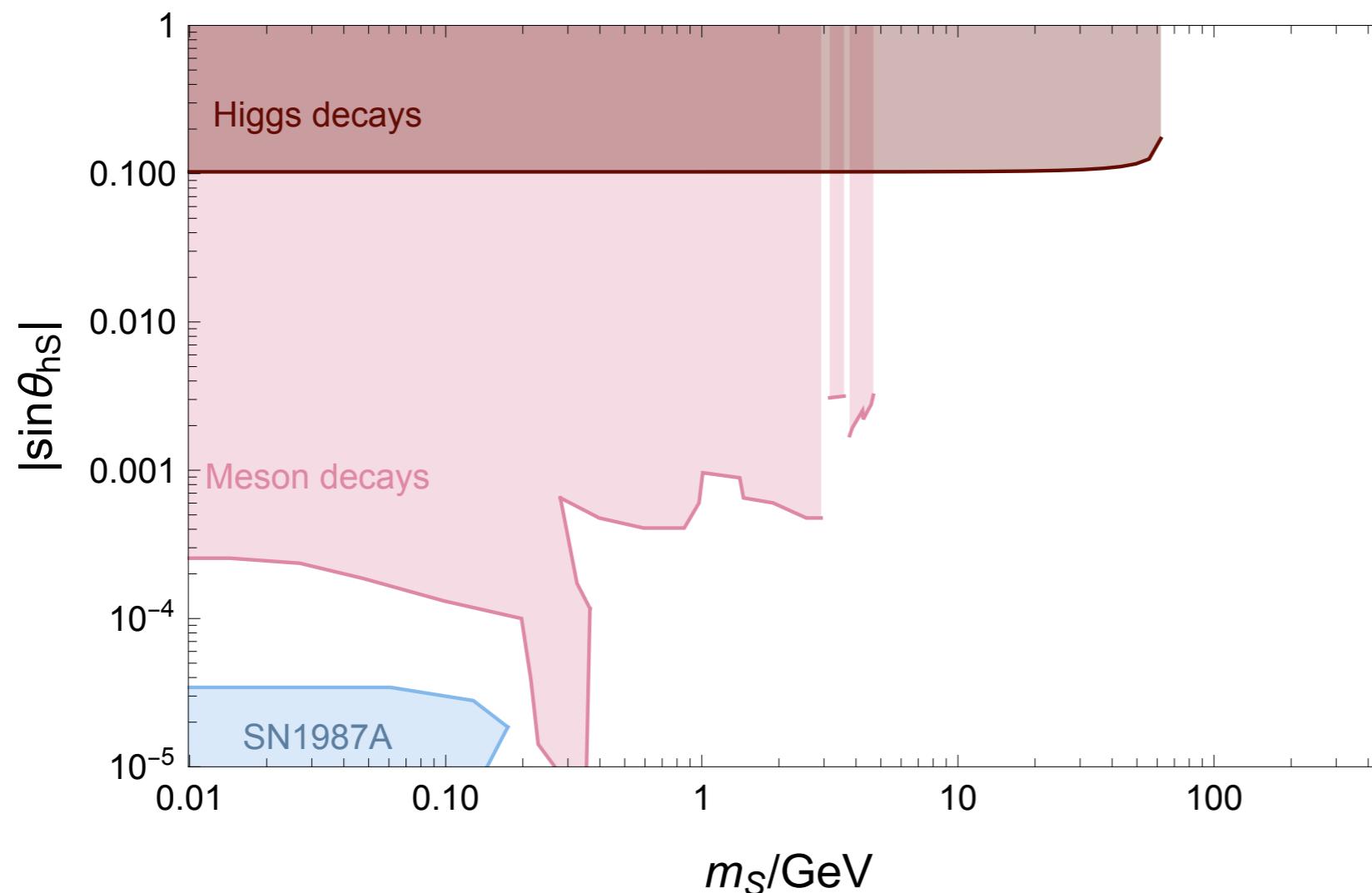


Parameter Space



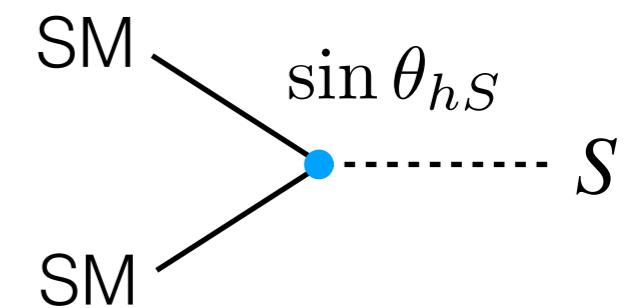
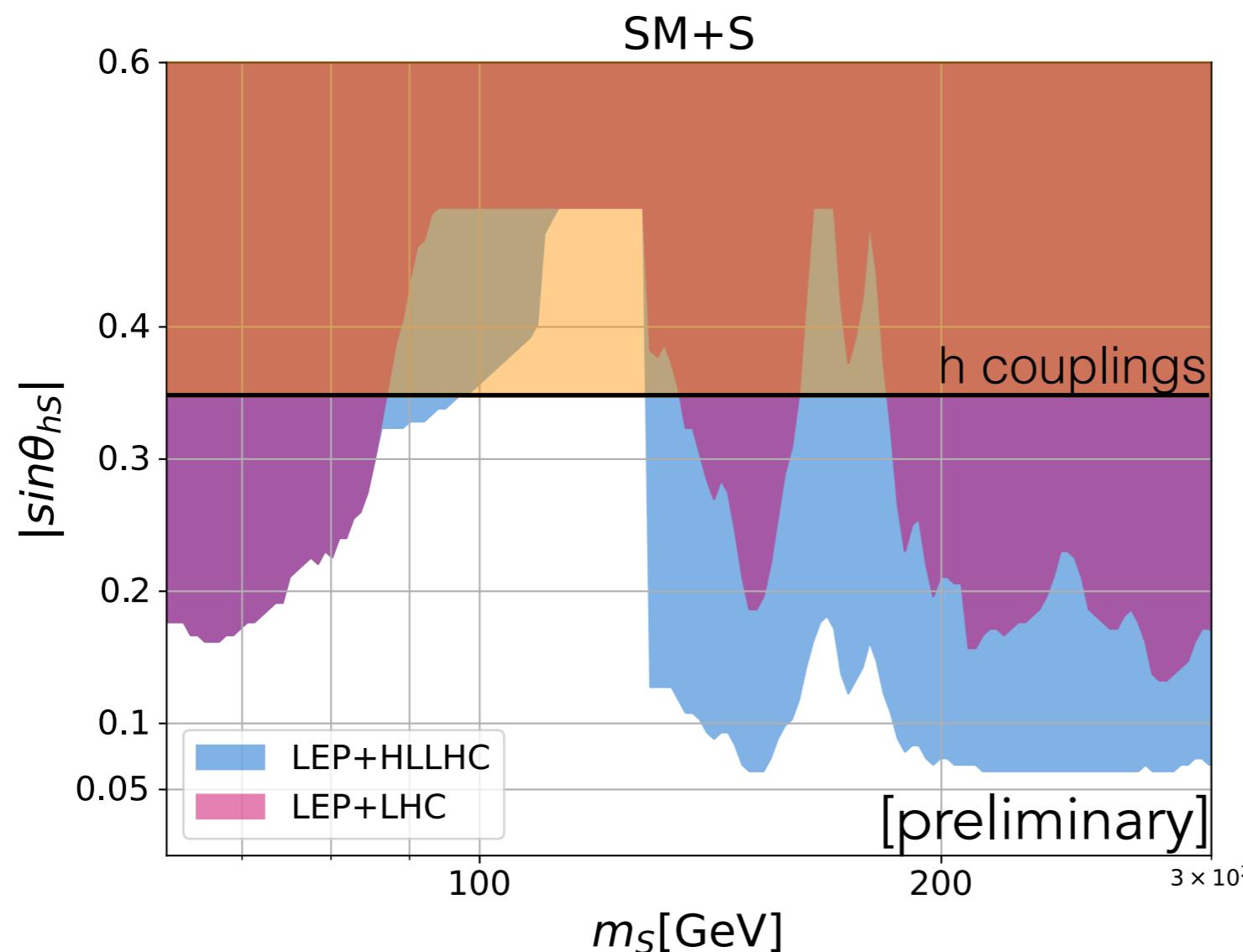
Parameter Space

- High mass region:



Parameter Space

- High mass region:



Naturalness Perspective

Naturalness Perspective

Gauge hierarchy problem:

$$\delta m_h^2 \propto \Lambda^2 \leftarrow \text{any physics that Higgs interacts with}$$

e.g. $\frac{m_P^2}{m_h^2} \sim 10^{34}$

Naturalness Perspective

Traditional approaches:

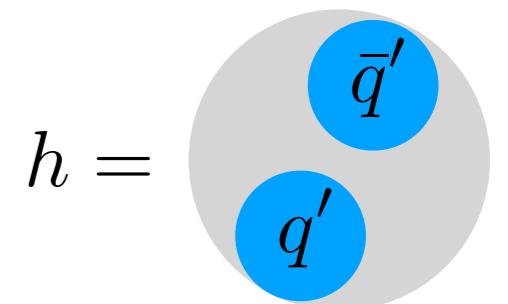
$$\delta m_h^2 = \cancel{0} \Lambda^2 + \mathcal{O}(100\text{GeV})$$



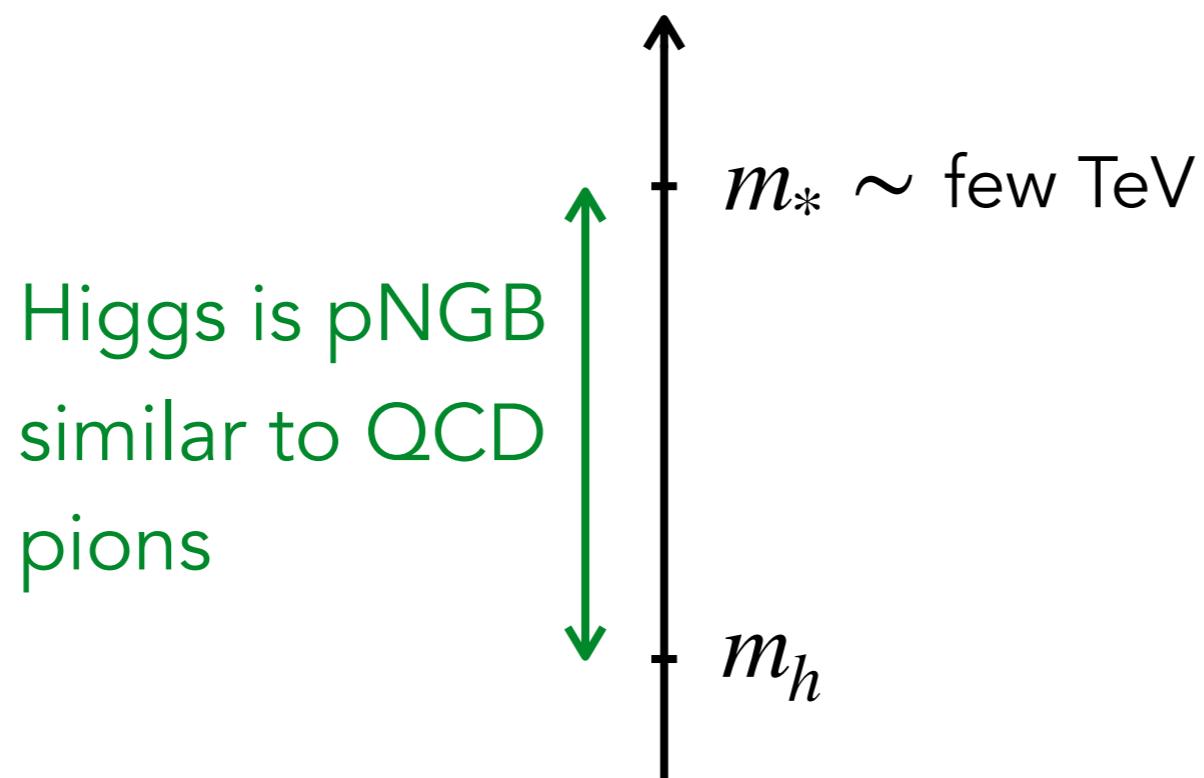
supersymmetry
compositeness
extra dimensions

Composite Higgs

Higgs is a bound state of new strong interactions
confining at $f \sim 1\text{TeV}$



spectrum:



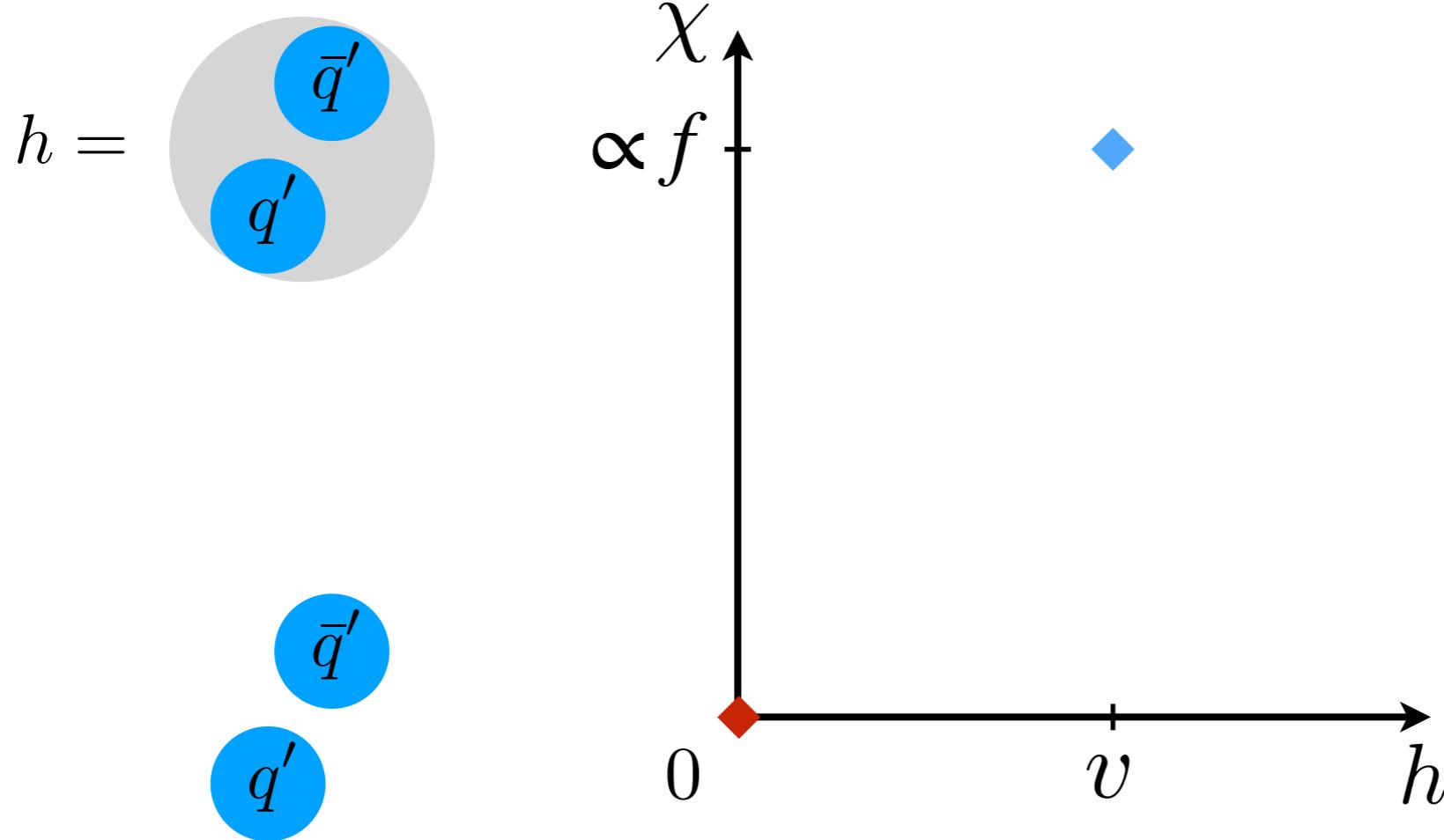
Kaplan,Georgi '84

Agashe,Contino,Pomarol '04

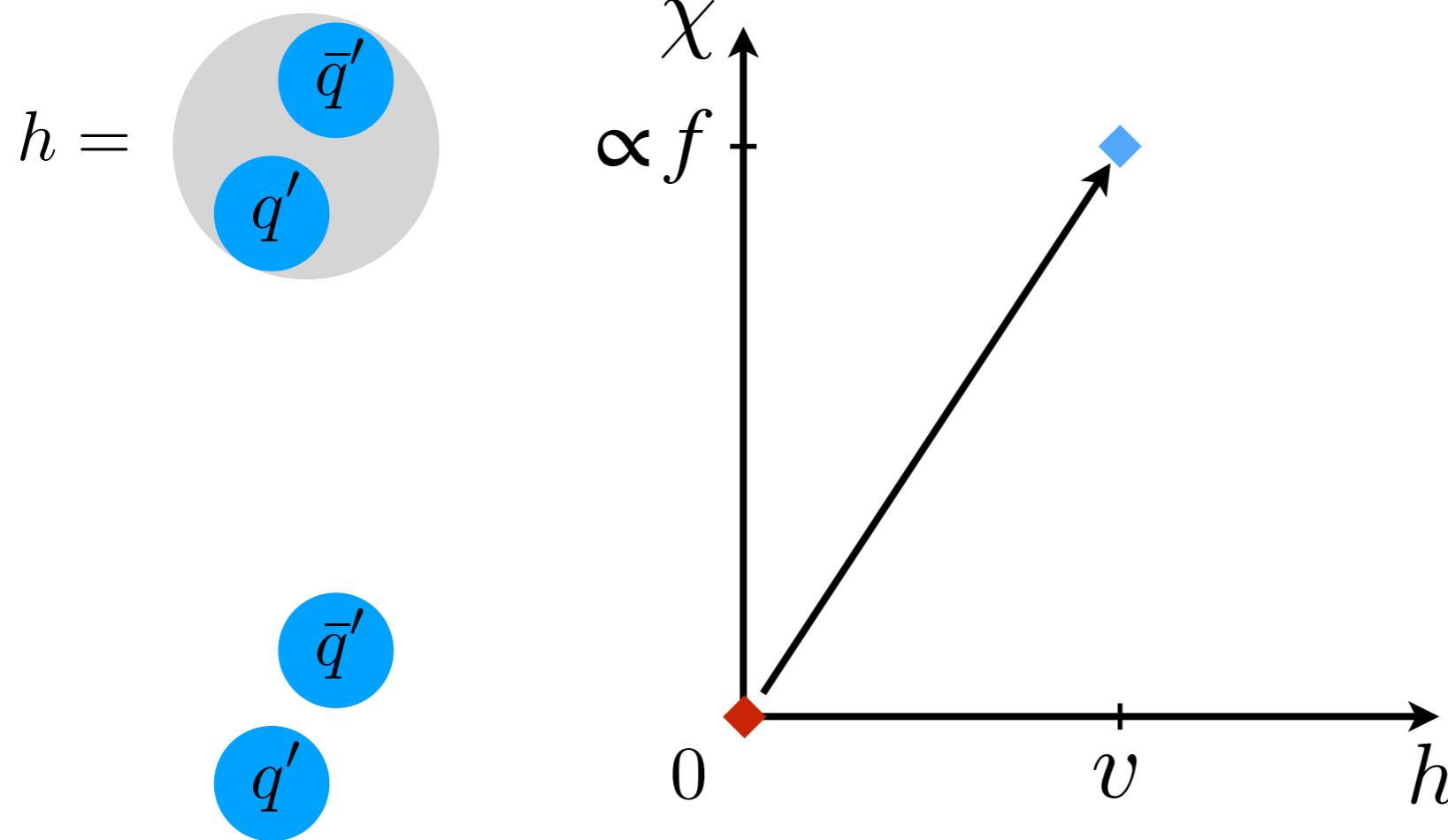
Phase Transitions in CH models



Phase Transitions in CH models



Phase Transitions in CH models

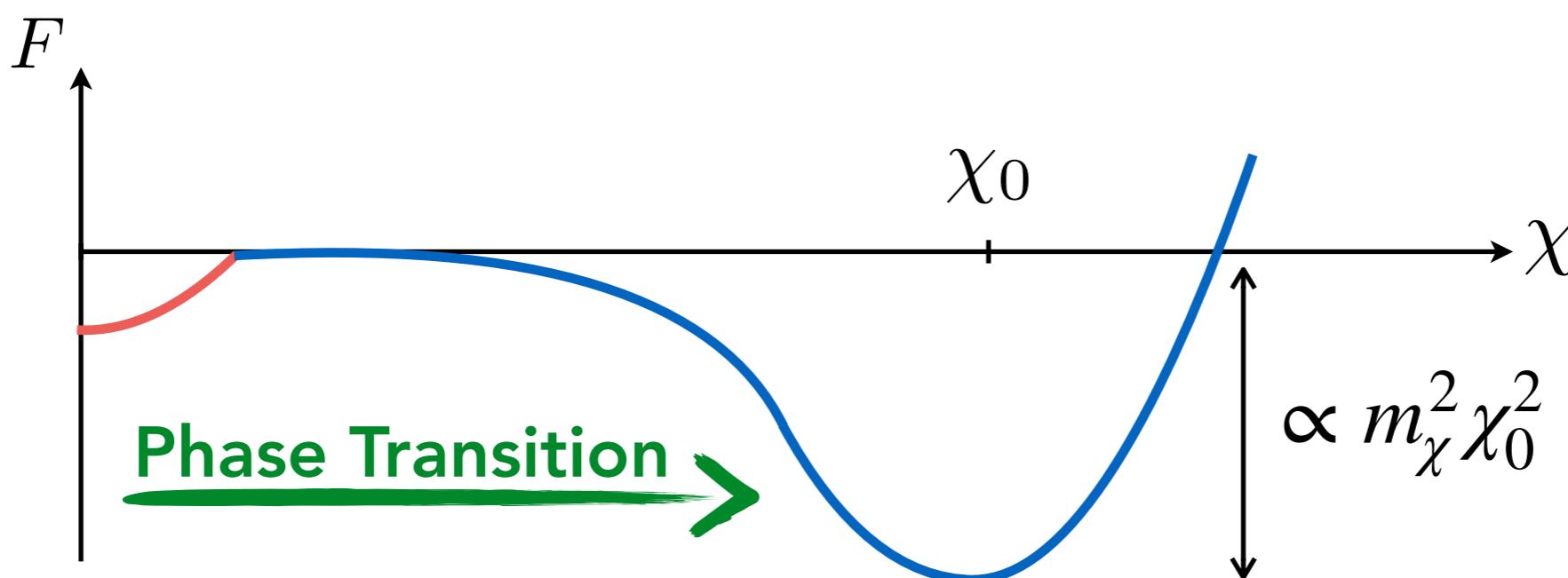


1-step: if $T(\text{confinement}) < T(\text{EWSB})$

$h \propto \chi$ and EWPT is 1st order if confinement PT is

Dilaton & Conformal Invariance

- Conformal invariance is often assumed for new strong interactions
- When broken spontaneously, a light PNGB - dilaton - can appear
- Dilaton mass and VEV are defining for the PT properties



Dilaton at the LHC

- Dilaton value sets all SM mass terms, and also renormalization scales for the running induced by the new strong dynamics.

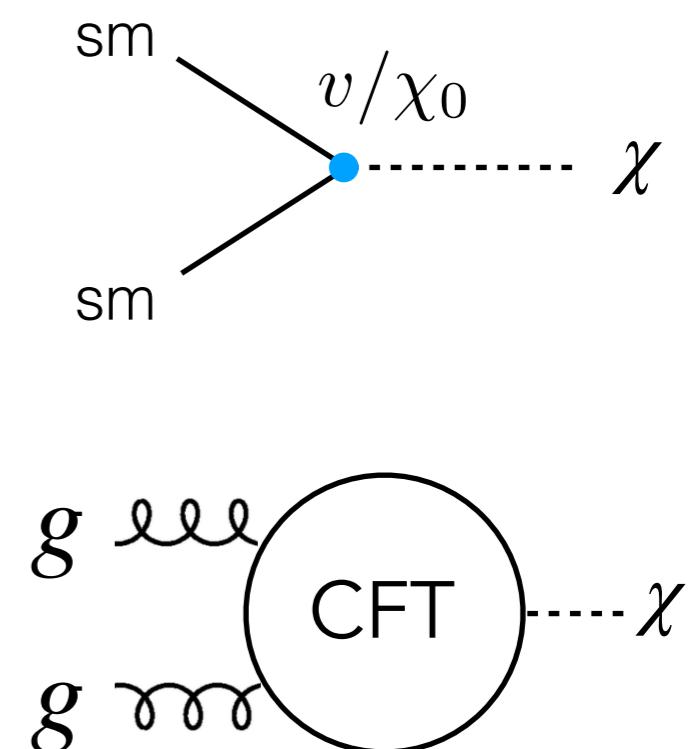
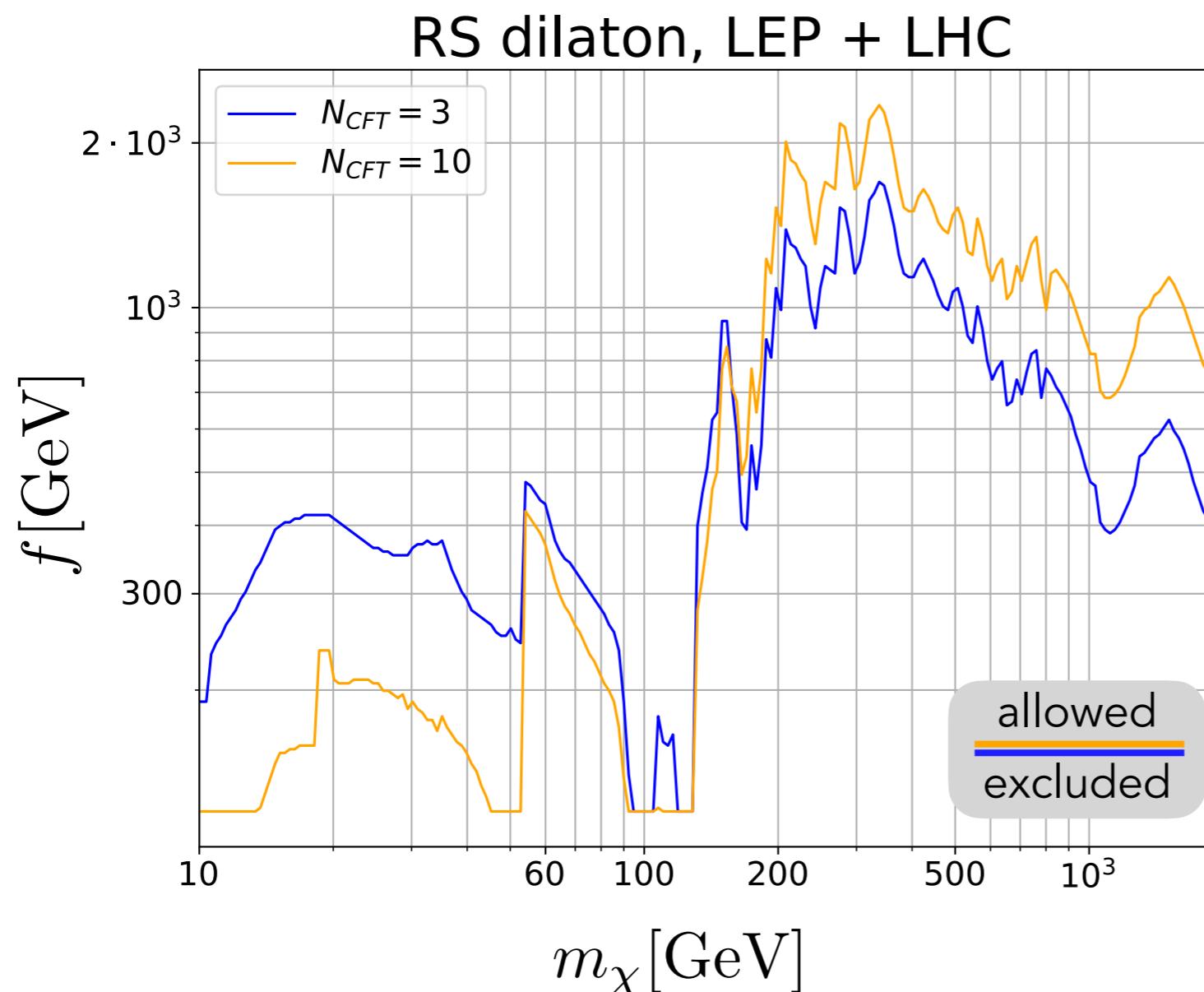
e.g. quark mass: $\underbrace{m_q \bar{q}q}$

$$m_q \frac{\chi}{\chi_0} = \frac{y_q}{\sqrt{2}} \frac{v}{\chi_0} \chi \quad \text{with} \quad \chi_0 \propto f$$

- Hence similar to the Higgs couplings, up to a factor of v/χ_0

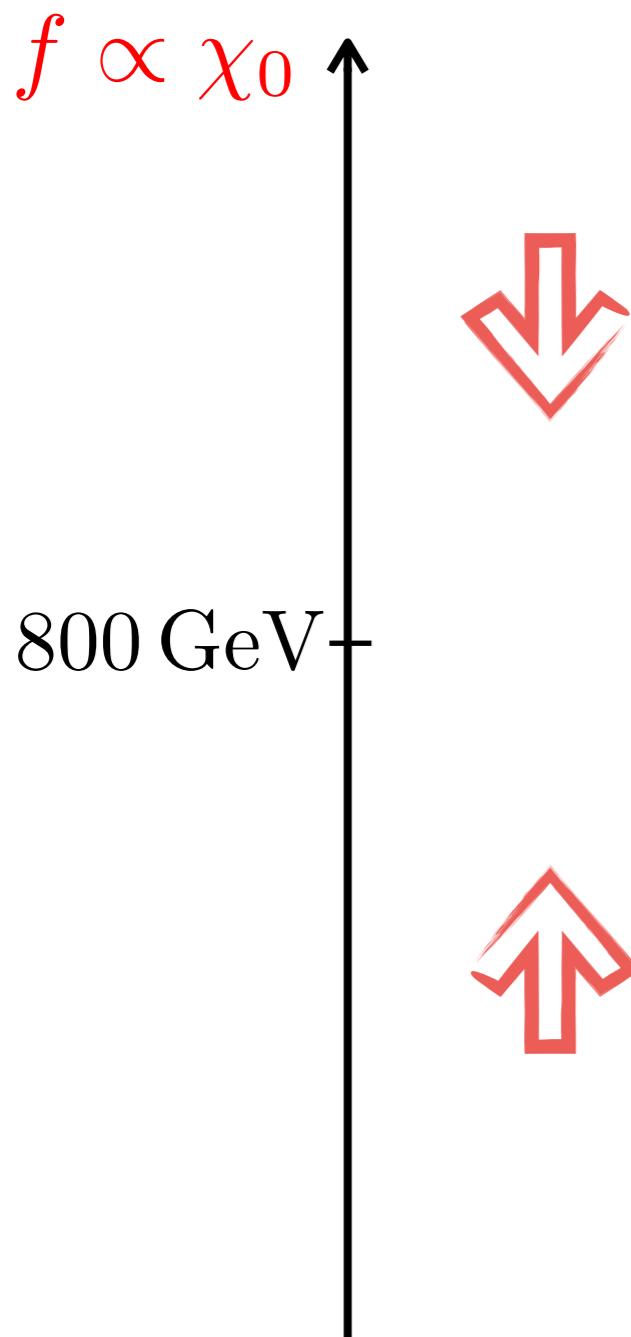
Dilaton at the LHC

- LHC bounds:



Dilaton at the LHC

- Probing compositeness as a solution to the naturalness problem:



Naturalness:

EW scale $\sim f$ times **tuning**

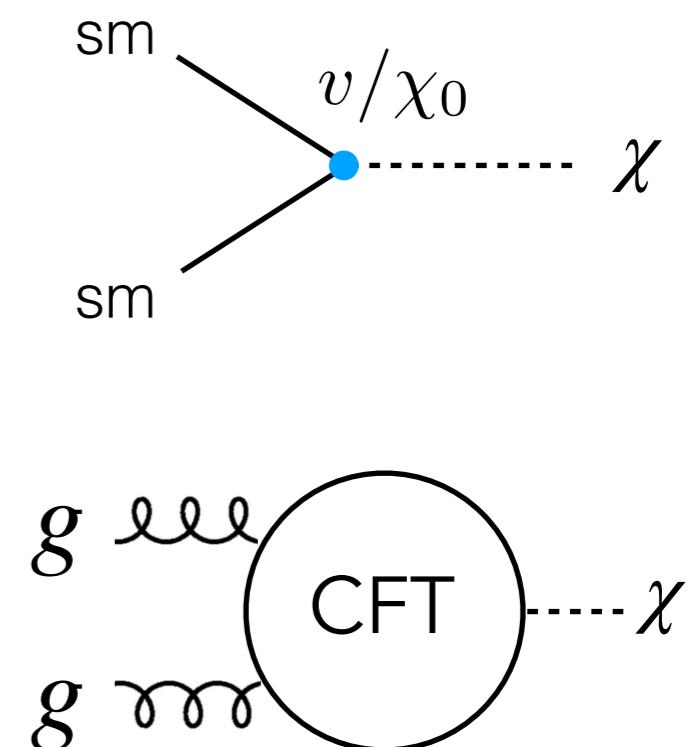
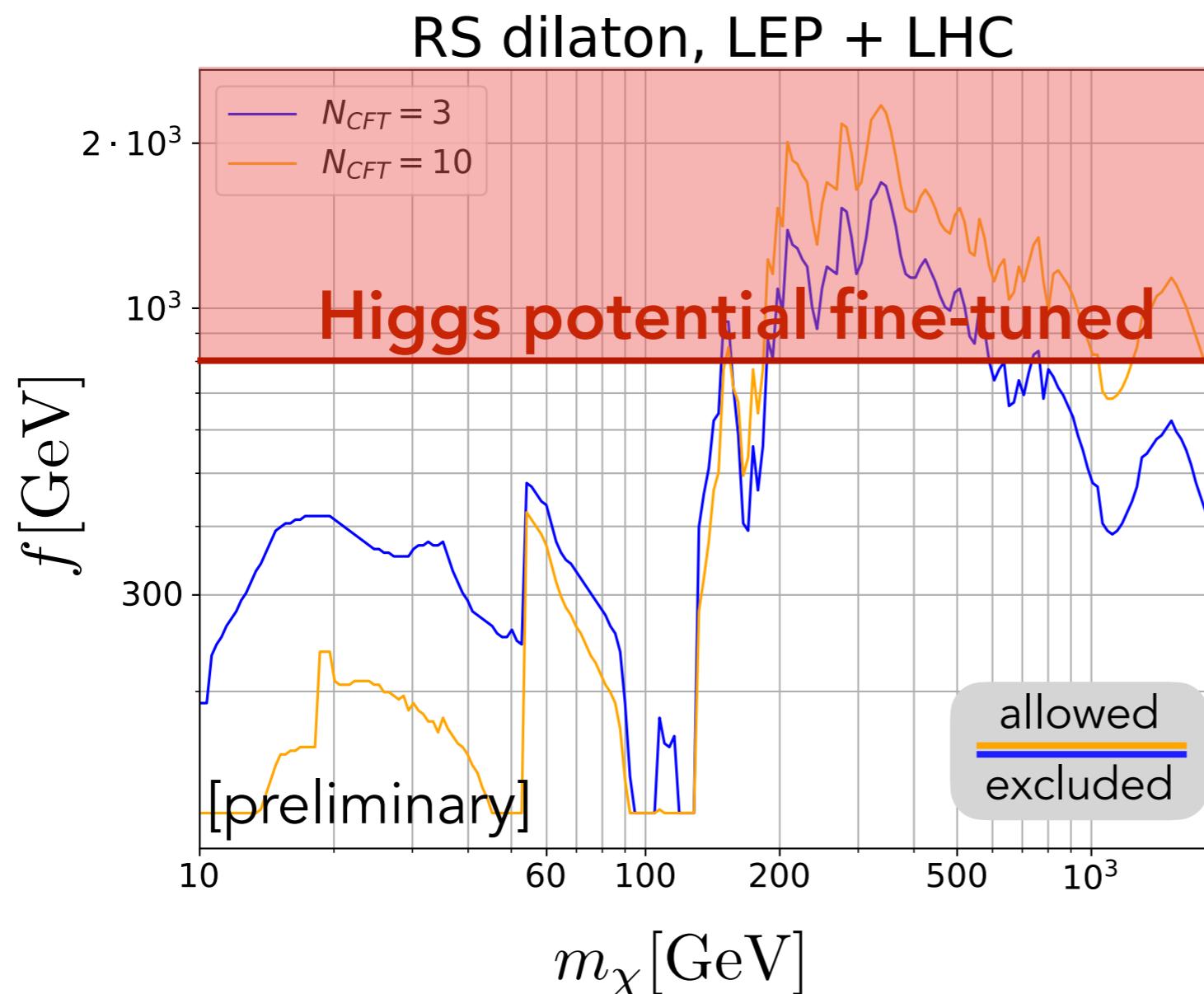
Experimental bounds:

$$\text{e.g. } m_W^2 \propto \sin^2 h/f$$

$$g_{hWW} \propto (m^2)'_h \propto \cos v/f \simeq 1 - v^2/2f^2$$

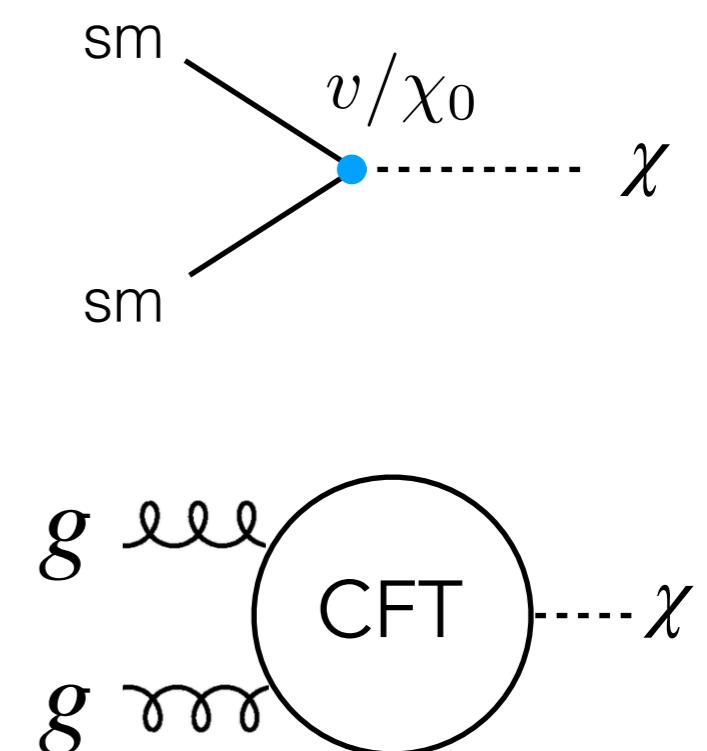
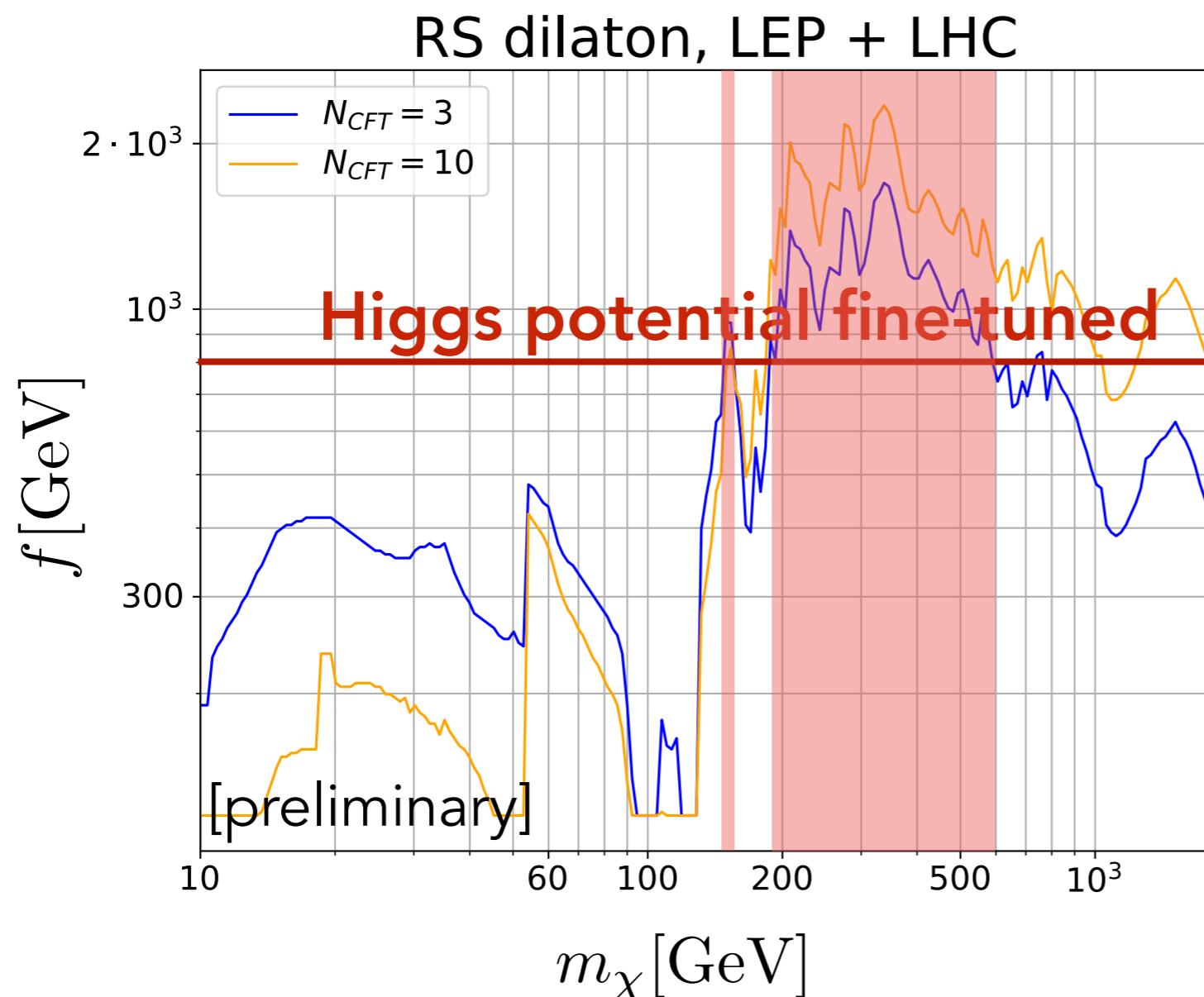
Dilaton at the LHC

- LHC bounds:



Dilaton at the LHC

- LHC bounds:



Summary

Summary

- Electroweak phase transition may be related to such fundamental particle physics problems as matter-antimatter asymmetry and gauge hierarchy problem
- EWPT can be tested at colliders
 - such indirect tests require coherent combined treatment of all the available probes (colliders, EDM, GW, ...)
 - and hence models that could describe the correlations
- Minimal models for EWBG are not fully probed yet and provide a range of potential collider signals, from Higgs couplings modifications to new resonances

Thank you!