

Long-distance reconstruction of QED corrections to the hadronic vacuum polarization for the muon g-2

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The long-distance contribution of QED corrections to the hadronic vacuum polarization is particularly challenging to compute in lattice QCD+QED. Currently, it is one of the limiting factors towards matching the precision of the recent result by the Fermilab E989 experiment for the muon g-2. In this work, we present a method for obtaining high-precision results for this contribution by reconstructing exclusive finite-volume state contributions. We find relations between the pion-photon contributions of individual diagrams and demonstrate the reconstruction method with lattice QCD+QED data at a single lattice spacing of $a^{-1} \approx 1.73$ GeV and $m_\pi \approx 275$ MeV.

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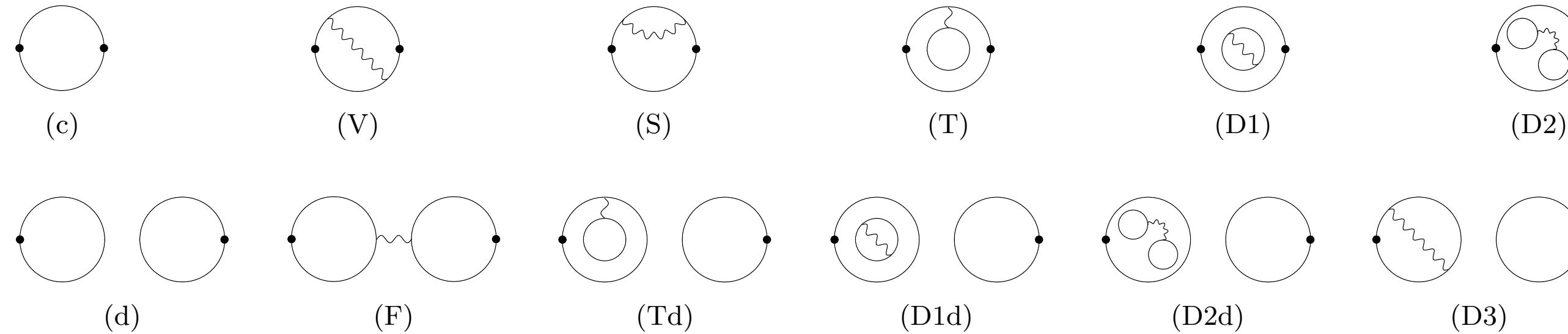
arXiv:2508.21685

The vector-vector correlator including QED corrections at order e^2 is without SIB is

$$G(t) \equiv \frac{1}{3} \sum_i \langle V_i(t, \vec{p}=0) V_i^\dagger(t=0, \vec{p}=0) \rangle = G^{(0)}(t) + e^2 G^{(2)}(t)$$

$$V_j(t, \vec{p}) = \frac{1}{\sqrt{L^3}} \sum_{\vec{x}} e^{i\vec{x} \cdot \vec{p}} \left(i \frac{2}{3} \bar{u}(t, \vec{x}) \gamma_j u(t, \vec{x}) - i \frac{1}{3} \bar{d}(t, \vec{x}) \gamma_j d(t, \vec{x}) \right)$$

which gives 12 diagrams



making the prefactors explicit

$$G = \frac{5}{9}(c) - \frac{1}{9}(d) + e^2 \left(-\frac{17}{81}((V) + 2(S)) + \frac{25}{81}(F) + \frac{14}{81}((T) + (D3)) - \frac{10}{81}(Td) \right. \\ \left. + \frac{25}{162}(D1) - \frac{5}{162}((D1d) + (D2)) + \frac{1}{162}(D2d) \right)$$

Using a QED regulator with a finite-volume transfer matrix, a spectral representation is admitted

$$G(t) = \sum_n c_n e^{-E_n t} = \sum_n c_n^{(0)} e^{-E_n^{(0)} t} + e^2 \sum_n (c_n^{(2)} - t E_n^{(2)} c_n^{(0)}) e^{-E_n^{(0)} t}$$

$$E_n = E_n^{(0)} + e^2 E_n^{(2)}$$

$$c_n = c_n^{(0)} + e^2 c_n^{(2)} = |\langle 0 | V_i | n \rangle|^2$$

Idea: study larger operator basis that includes operators that project to individual states

Minimal example: vector current and a separate operator for lowest-lying pion-photon state

$$C_{ij} = \langle O_i O_j^\dagger \rangle$$

$$O_1 = V_i$$

$$O_2 = O_{\pi\gamma}$$

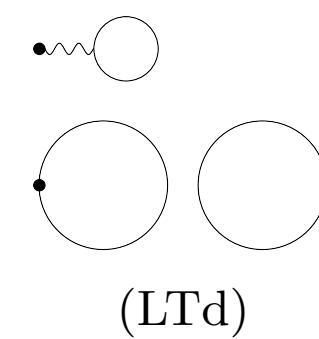
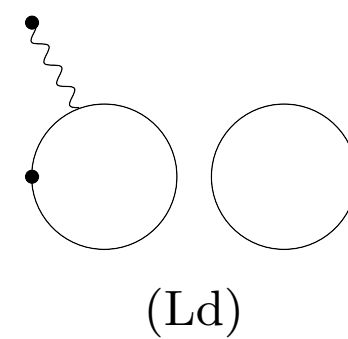
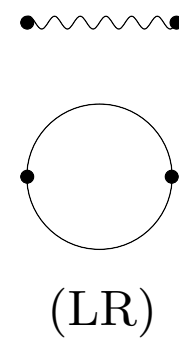
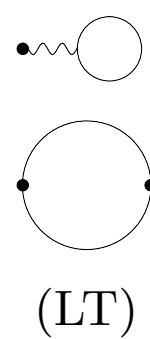
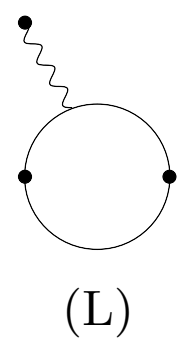
The correlation matrix

$$C(t) = \begin{pmatrix} |V_{1,\rho}^{(0)}|^2 e^{-E_\rho^{(0)}t} & 0 \\ 0 & 0 \end{pmatrix} + e^2 \begin{pmatrix} (2\text{Re}(V_{1,\rho}^{(0)} V_{1,\rho}^{(2),*}) - t|V_{1,\rho}^{(0)}|^2 E_\rho^{(2)})e^{-E_\rho^{(0)}t} + |V_{1,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)}t} & V_{1,\rho}^{(0)} V_{2,\rho}^{(2),*} e^{-E_\rho^{(0)}t} + V_{1,\pi\gamma}^{(1)} V_{2,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)}t} \\ V_{2,\rho}^{(2)} V_{1,\rho}^{(0),*} e^{-E_\rho^{(0)}t} + V_{2,\pi\gamma}^{(1)} V_{1,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)}t} & |V_{2,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)}t} \end{pmatrix}$$

$$V = V^{(0)} + eV^{(1)} + e^2V^{(2)}$$

$$V_{in} = \langle 0 | O_i | n \rangle$$

needs new diagrams



Full analysis allows for determination of individual state contributions to $G(t)$.

Pion-photon operator in finite volume:

- Needs to transform in T_1^u irrep. of octahedral group
- Here: use Feynman gauge but only couple to transversal photons
- Photon and pion have fixed back-to-back momentum

Unique operator satisfying these conditions:

$$\begin{aligned}
 O_{\pi\gamma,\vec{p},i}(t) &= \frac{e}{2} \sum_{\vec{q} \in H(\vec{p})} \left([\hat{q} \times \vec{A}^T(t, \vec{q})]_i O_\pi(t, -\vec{q}) - [\hat{q} \times \vec{A}^T(t, -\vec{q})]_i O_\pi(t, +\vec{q}) \right) \\
 &= \frac{e}{2} \sum_{\vec{q} \in H(\vec{p})} \left([\hat{q} \times \vec{A}(t, \vec{q})]_i O_\pi(t, -\vec{q}) - [\hat{q} \times \vec{A}(t, -\vec{q})]_i O_\pi(t, +\vec{q}) \right).
 \end{aligned}$$

$$A_i(t, \vec{p}) = \frac{1}{\sqrt{L^3}} \sum_{\vec{x}} e^{i\vec{x}\vec{p}} \tilde{A}_i(x)$$

$$A_j^T(t, \vec{p}) = \left(\delta_{jm} - \hat{p}_j \hat{p}_m \right) A_m(t, \vec{p}),$$

$$O_\pi(t, \vec{p}) = \frac{i}{\sqrt{L^3}} \sum_{\vec{x}} e^{i\vec{x}\vec{p}} \bar{u}(x) \gamma_5 d(x)$$

$$\hat{p} = \frac{\vec{p}}{|\vec{p}|}.$$

with $H(\vec{p})$ being the orbit under chiral octahedral group (24 elements) of vector \vec{p}

It is useful to consider $I=0$, $I=1$, and several QED charge assignments separately with

$$\begin{aligned} O_{\pi(I_3=-1)} &= i\bar{u}\gamma_5d\,, & O_{\pi(I_3=0)} &= \frac{i}{\sqrt{2}}\left(\bar{u}\gamma_5u-\bar{d}\gamma_5d\right)\,, & O_{\pi(I_3=1)} &= i\bar{d}\gamma_5u\,, \\ V_{I=1,I_3=-1}^j &= i\bar{u}\gamma_jd\,, & V_{I=1,I_3=0}^j &= \frac{i}{\sqrt{2}}\left(\bar{u}\gamma_ju-\bar{d}\gamma_jd\right)\,, & V_{I=1,I_3=1}^j &= i\bar{d}\gamma_ju\,, \\ & & V_{I=0}^j &= \frac{i}{\sqrt{2}}\left(\bar{u}\gamma_ju+\bar{d}\gamma_jd\right)\,.\end{aligned}$$

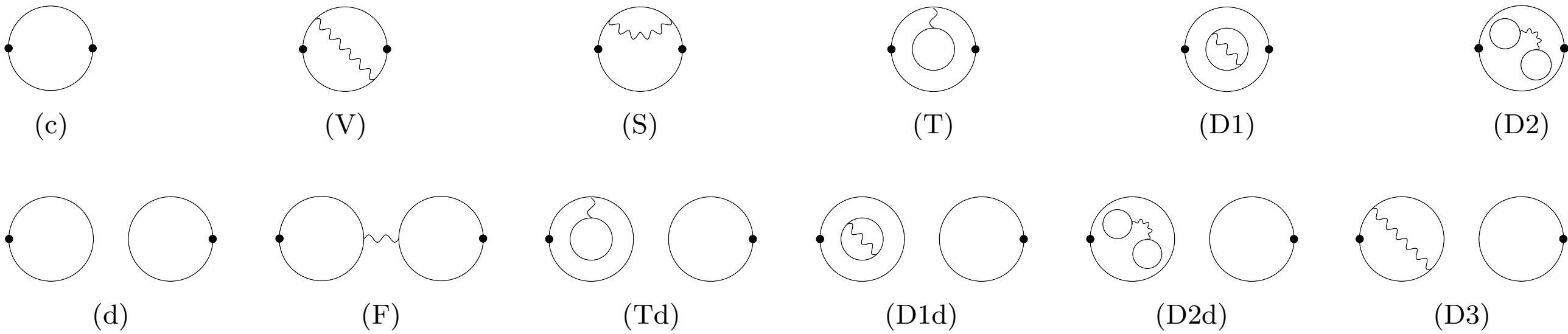
We show in the paper that various versions have the identical pion-photon contribution, which allows us to derive relations of the pion-photon contributions between individual diagrams:

$$(F)=(V)\,,\qquad 2(S)-(D1)=(V)\,,\qquad 2(T)-4(Td)-(D1d)+2(D3)+2(D2d)-(D2)=0\,.$$

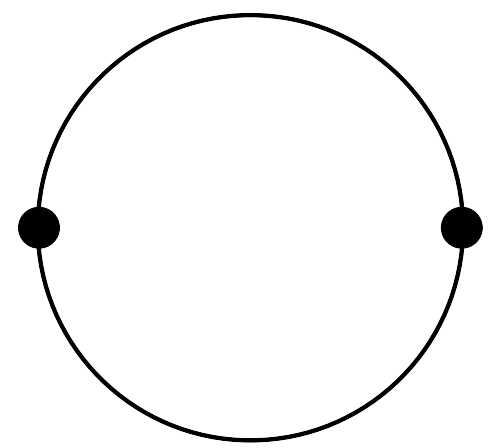
Collecting pion-photon contributions to HVP from diagrams (S), (V), (F), (D1) then gives

$$-e^2\frac{9}{162}((V)+2(S))$$

with a small QED charge pre-factor of $9/162$.



First results for $a^{-1} = 1.7312(28)$ GeV, $m_\pi = 274.8(2.5)$ MeV, $m_K = 530.1(3.1)$ MeV with $m_\pi L = 3.8$

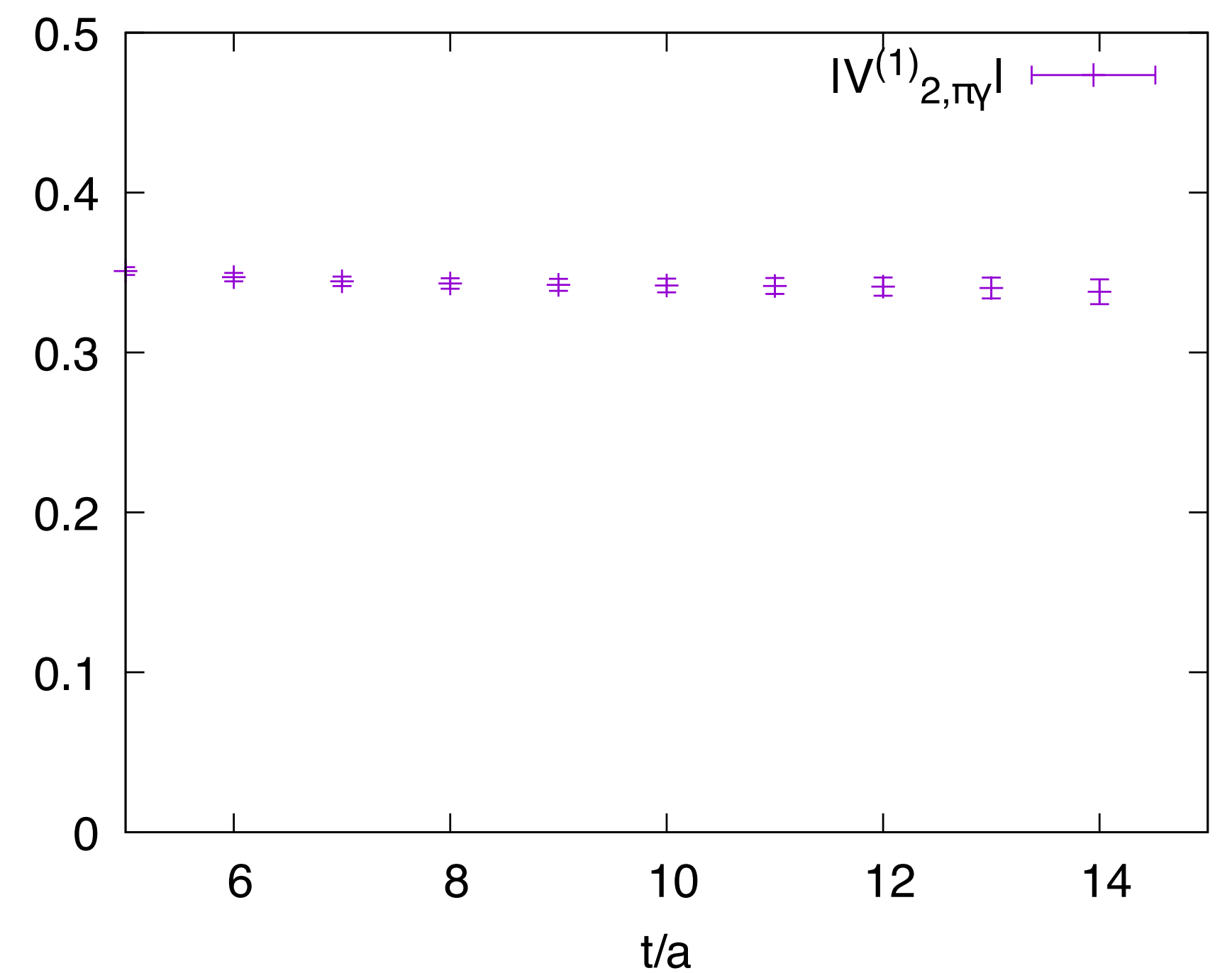
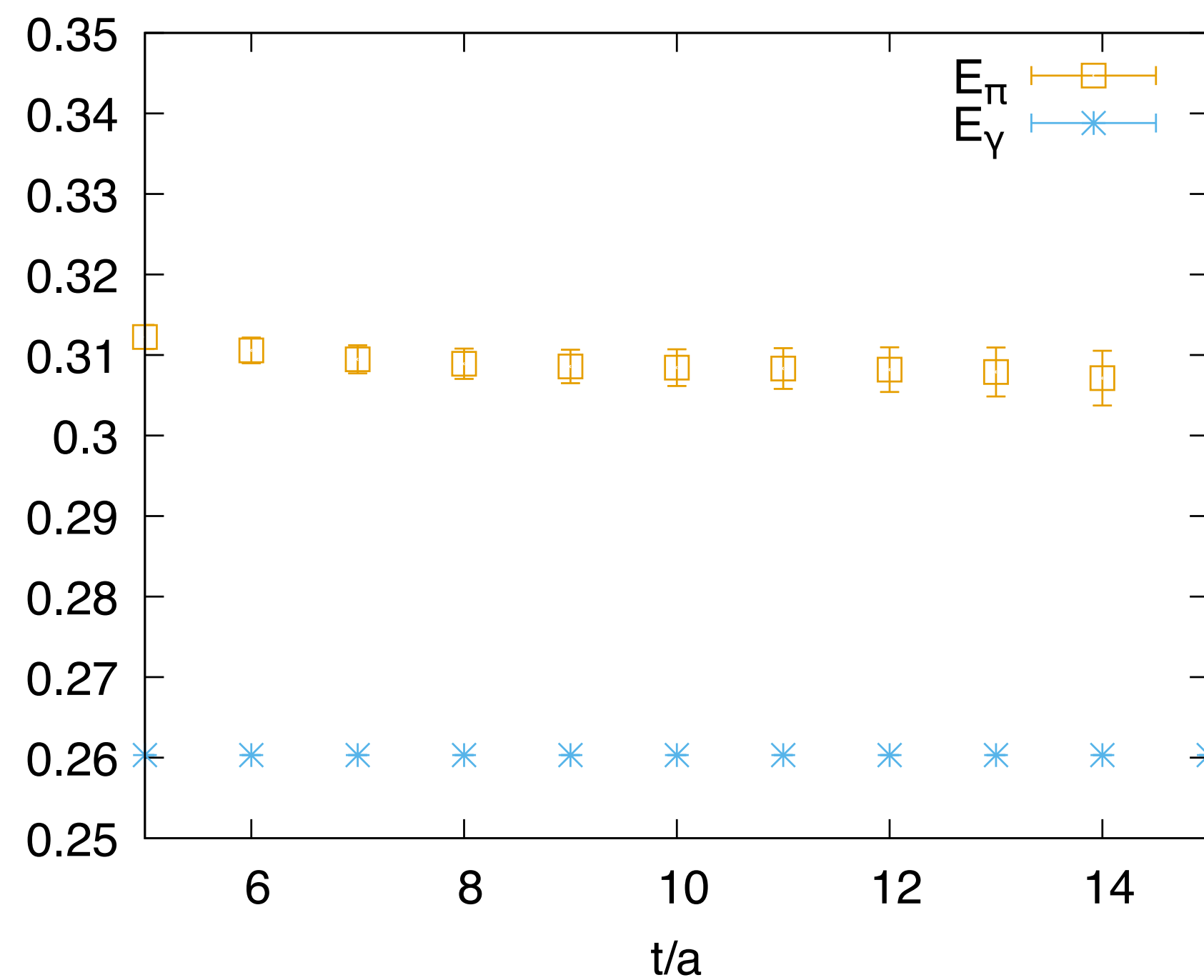


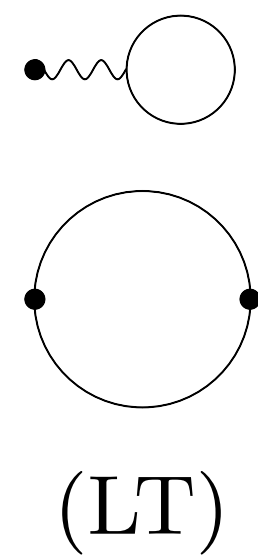
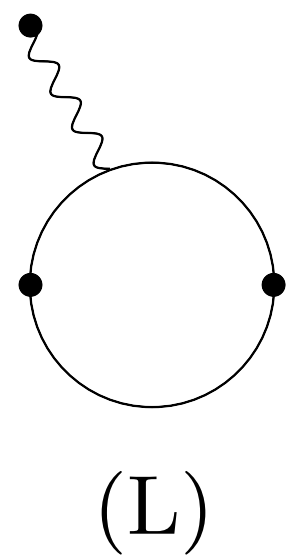
(LR)

$$V_{in} = \langle 0 | O_i | n \rangle \quad V = V^{(0)} + eV^{(1)} + e^2 V^{(2)}$$

$$O_1 = V_i \quad O_2 = O_{\pi\gamma}$$

$$C(t) = \begin{pmatrix} |V_{1,\rho}^{(0)}|^2 e^{-E_\rho^{(0)} t} & 0 \\ 0 & 0 \end{pmatrix} + e^2 \begin{pmatrix} (2\text{Re}(V_{1,\rho}^{(0)} V_{1,\rho}^{(2),*}) - t|V_{1,\rho}^{(0)}|^2 E_\rho^{(2)}) e^{-E_\rho^{(0)} t} + |V_{1,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)} t} & V_{1,\rho}^{(0)} V_{2,\rho}^{(2),*} e^{-E_\rho^{(0)} t} + V_{1,\pi\gamma}^{(1)} V_{2,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)} t} \\ V_{2,\rho}^{(2)} V_{1,\rho}^{(0),*} e^{-E_\rho^{(0)} t} + V_{2,\pi\gamma}^{(1)} V_{1,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)} t} & |V_{2,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)} t} \end{pmatrix}$$





$$C(t) = \begin{pmatrix} |V_{1,\rho}^{(0)}|^2 e^{-E_\rho^{(0)} t} & 0 \\ 0 & 0 \end{pmatrix}$$

$$+ e^2 \begin{pmatrix} (2\text{Re}(V_{1,\rho}^{(0)} V_{1,\rho}^{(2),*}) - t |V_{1,\rho}^{(0)}|^2 E_\rho^{(2)}) e^{-E_\rho^{(0)} t} + |V_{1,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)} t} & \\ V_{2,\rho}^{(2)} V_{1,\rho}^{(0),*} e^{-E_\rho^{(0)} t} + V_{2,\pi\gamma}^{(1)} V_{1,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)} t} & \end{pmatrix}$$

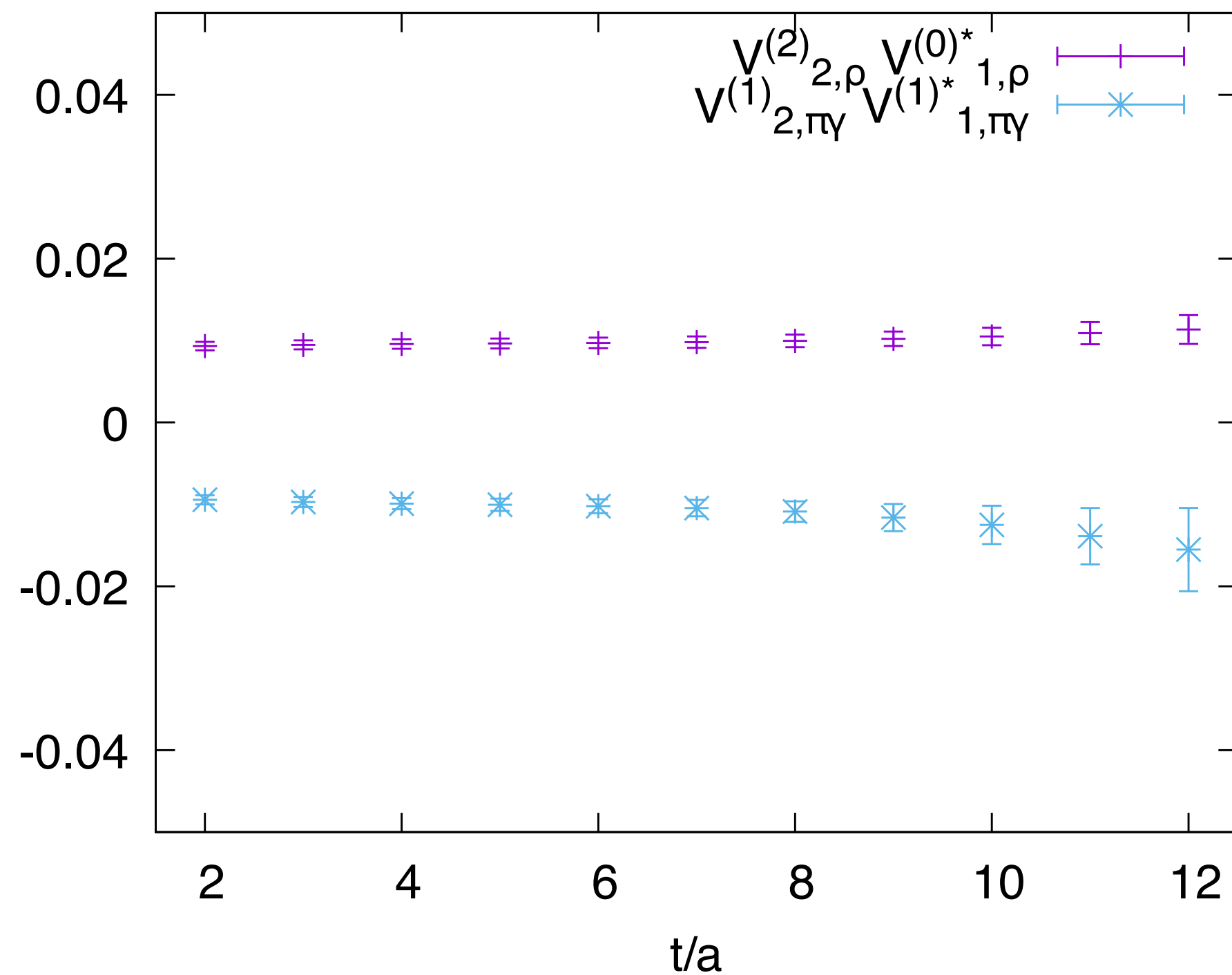
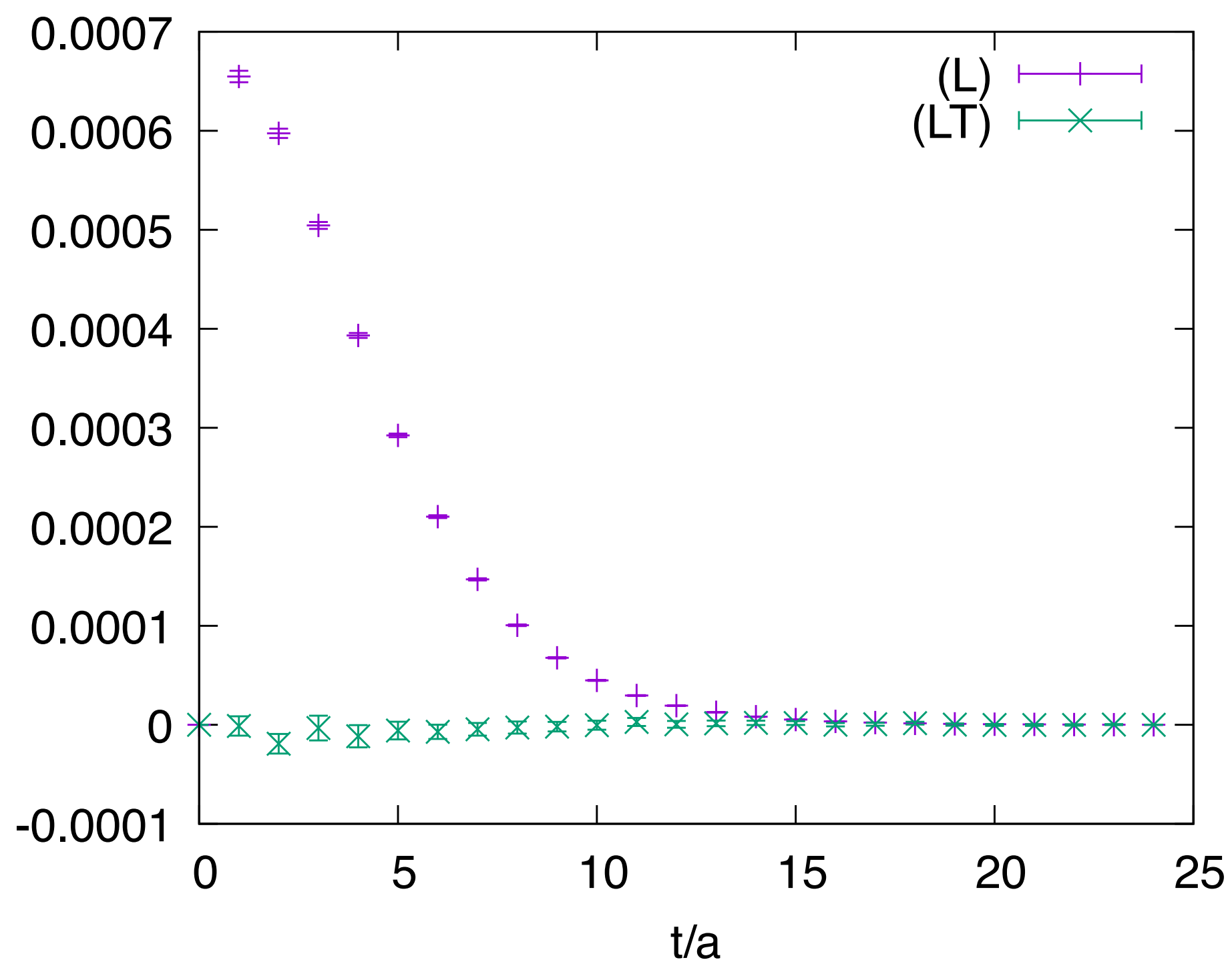
$$V_{in} = \langle 0 | O_i | n \rangle$$

$$O_1 = V_i$$

$$V = V^{(0)} + e V^{(1)} + e^2 V^{(2)}$$

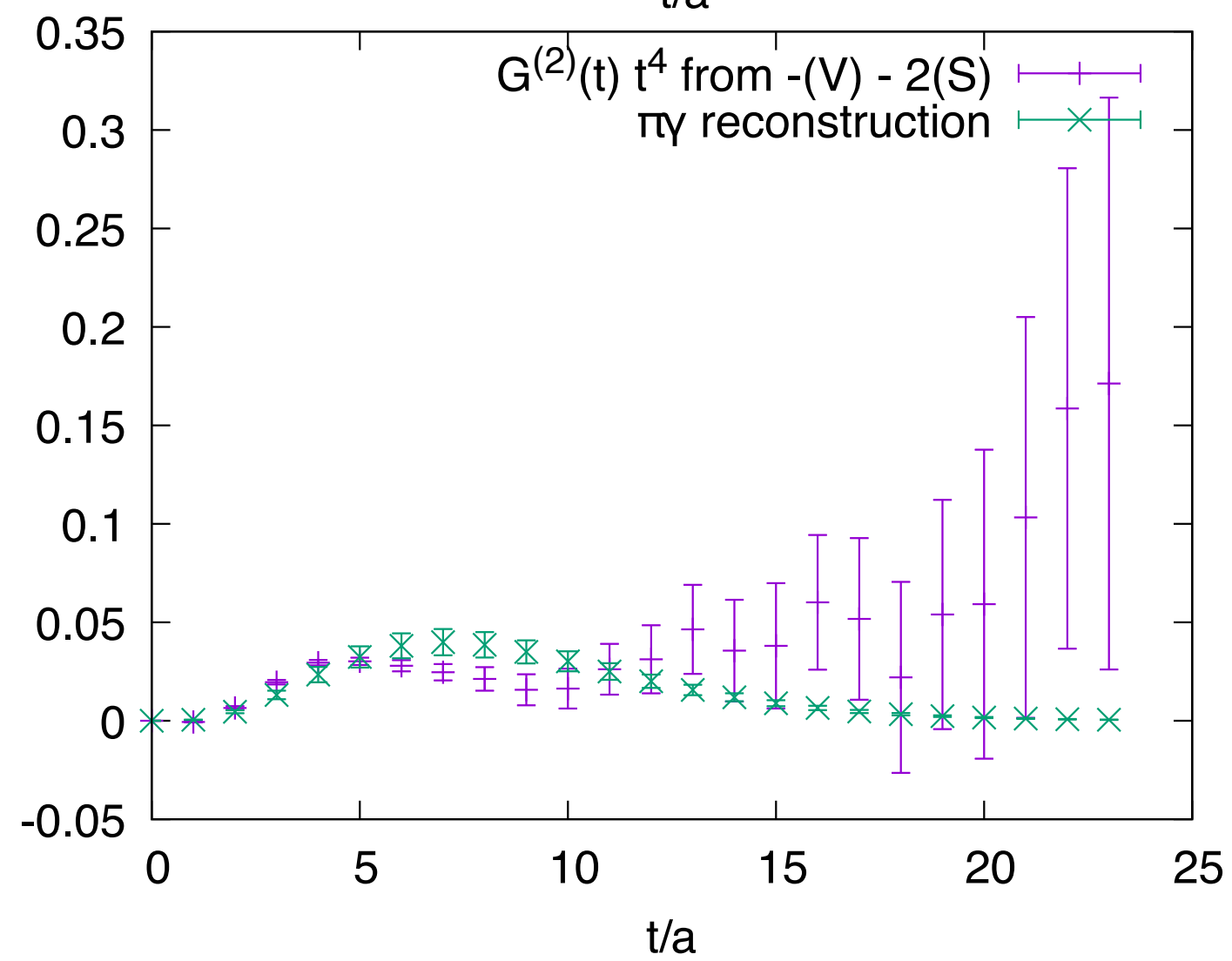
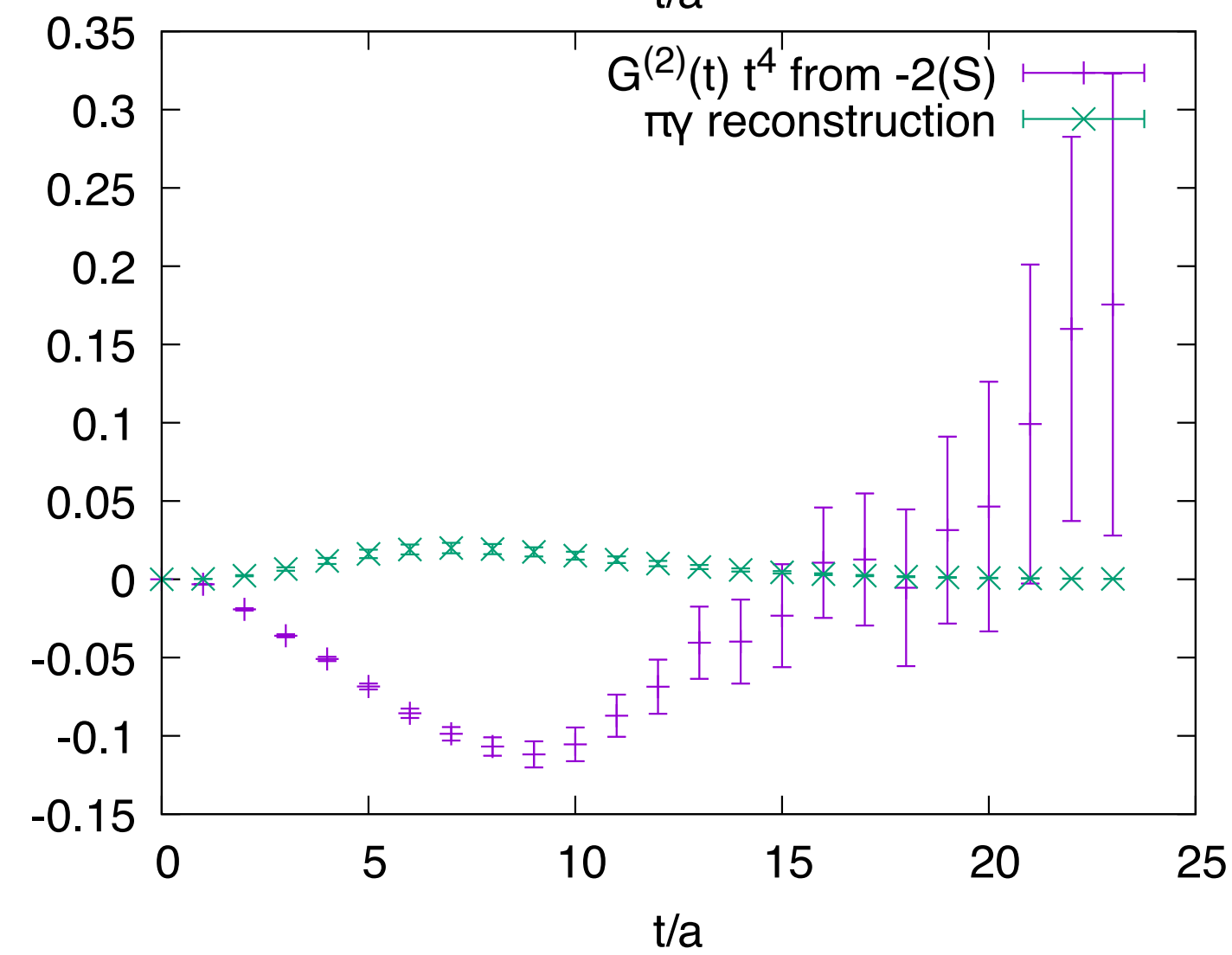
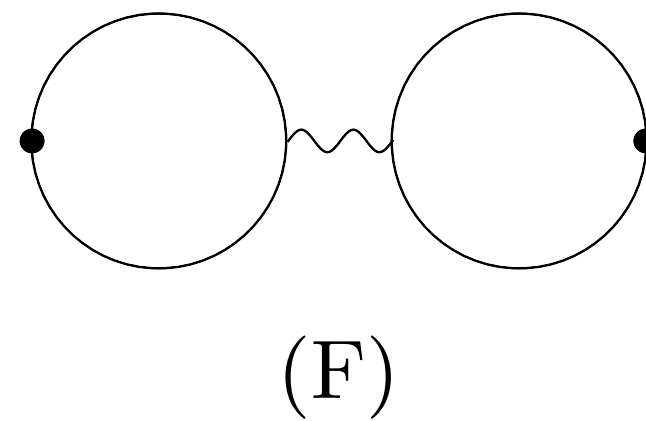
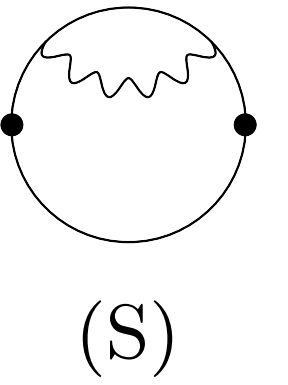
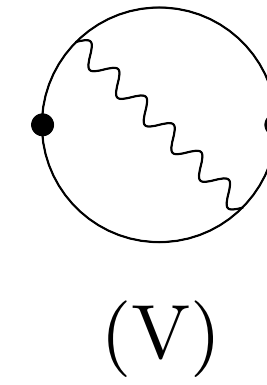
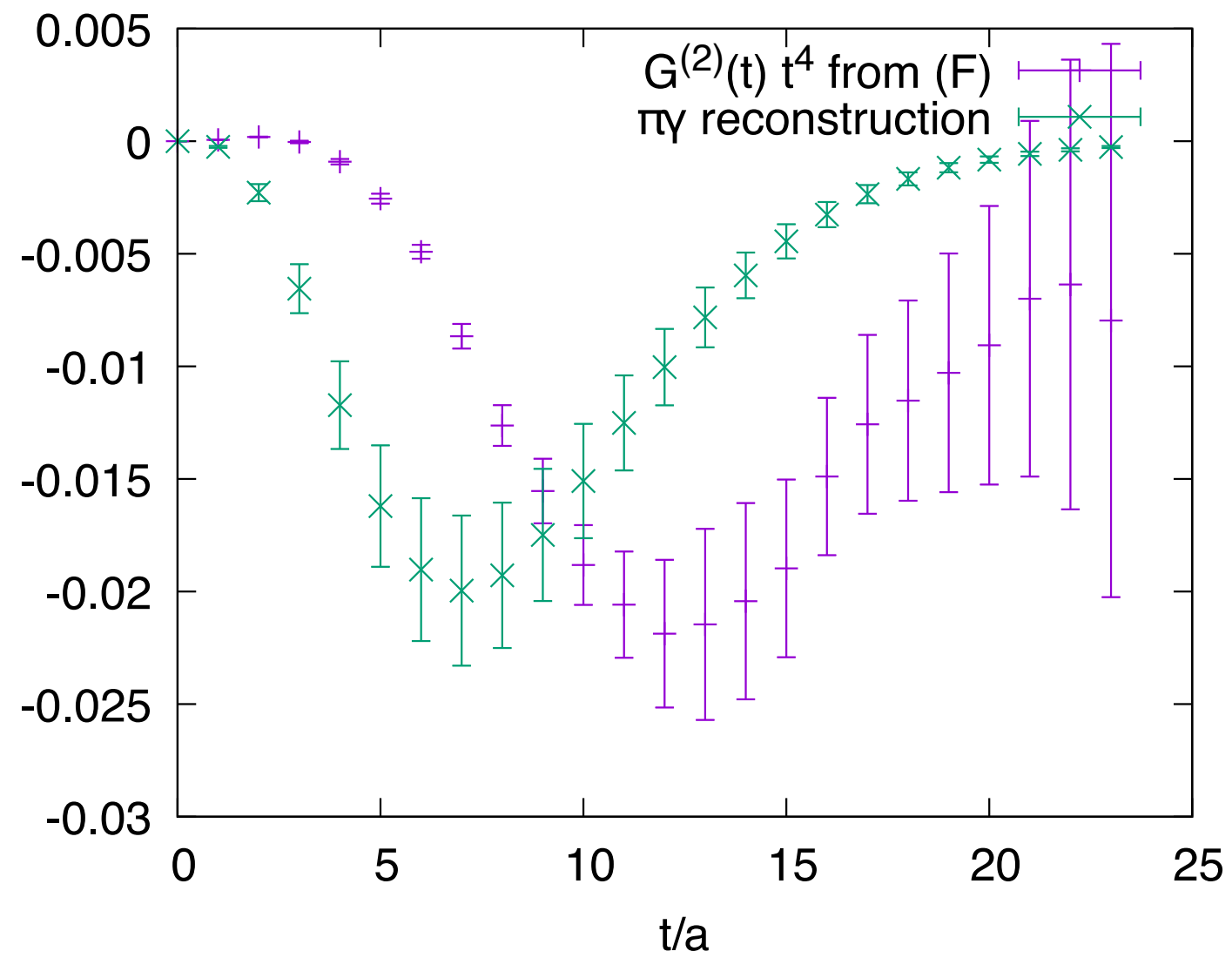
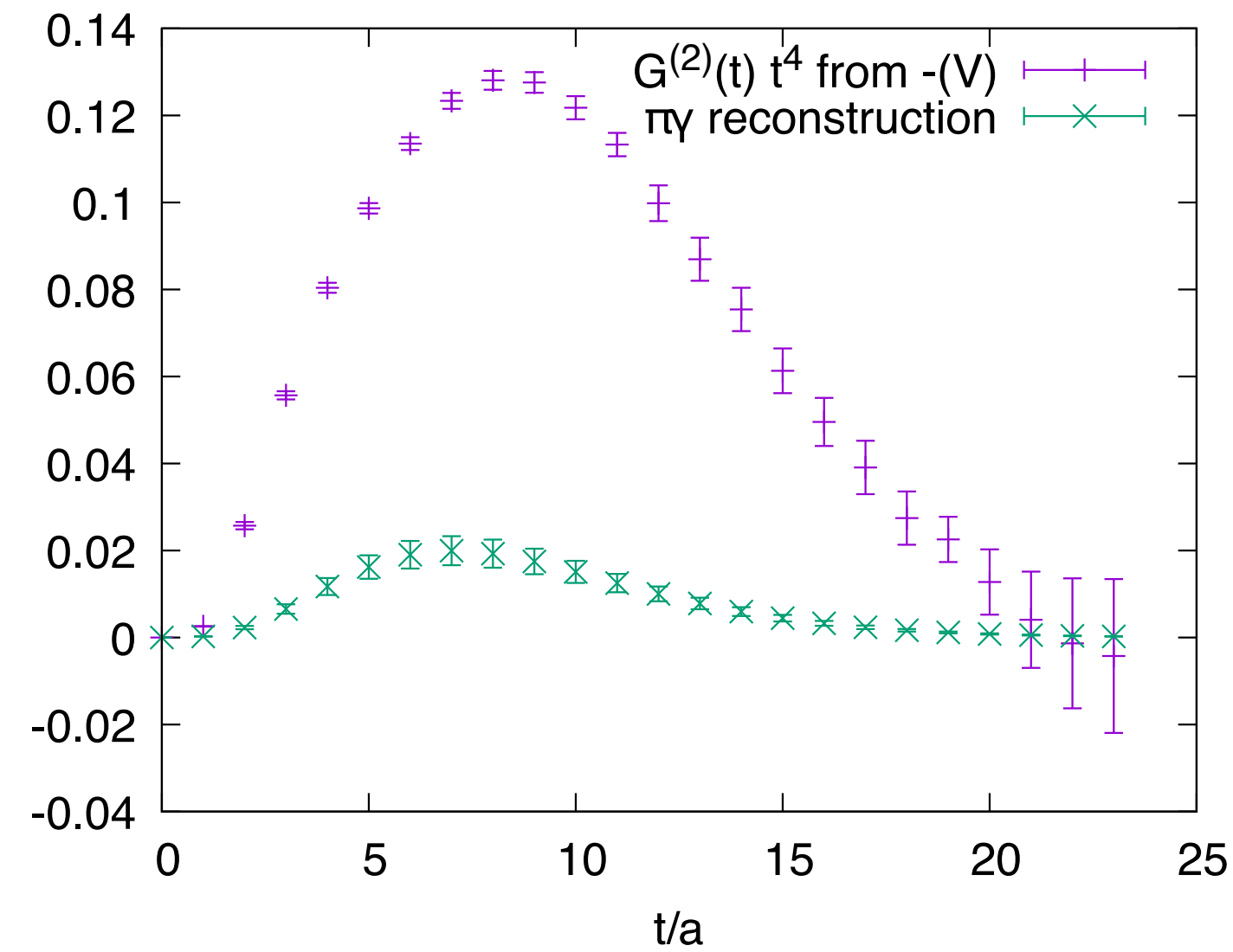
$$O_2 = O_{\pi\gamma}$$

$$\begin{pmatrix} V_{1,\rho}^{(0)} V_{2,\rho}^{(2),*} e^{-E_\rho^{(0)} t} + V_{1,\pi\gamma}^{(1)} V_{2,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)} t} \\ |V_{2,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)} t} \end{pmatrix}$$



Fit to matrix elements V (right panel) is stable versus fit range [t/a,24]

For now, only study $G(t)t^4$ which mimics behavior of HVP integrand (blinded study)



Next: subtract reconstructed pion-photon contribution
from $G(t)$ and determine QED corrections to ρ

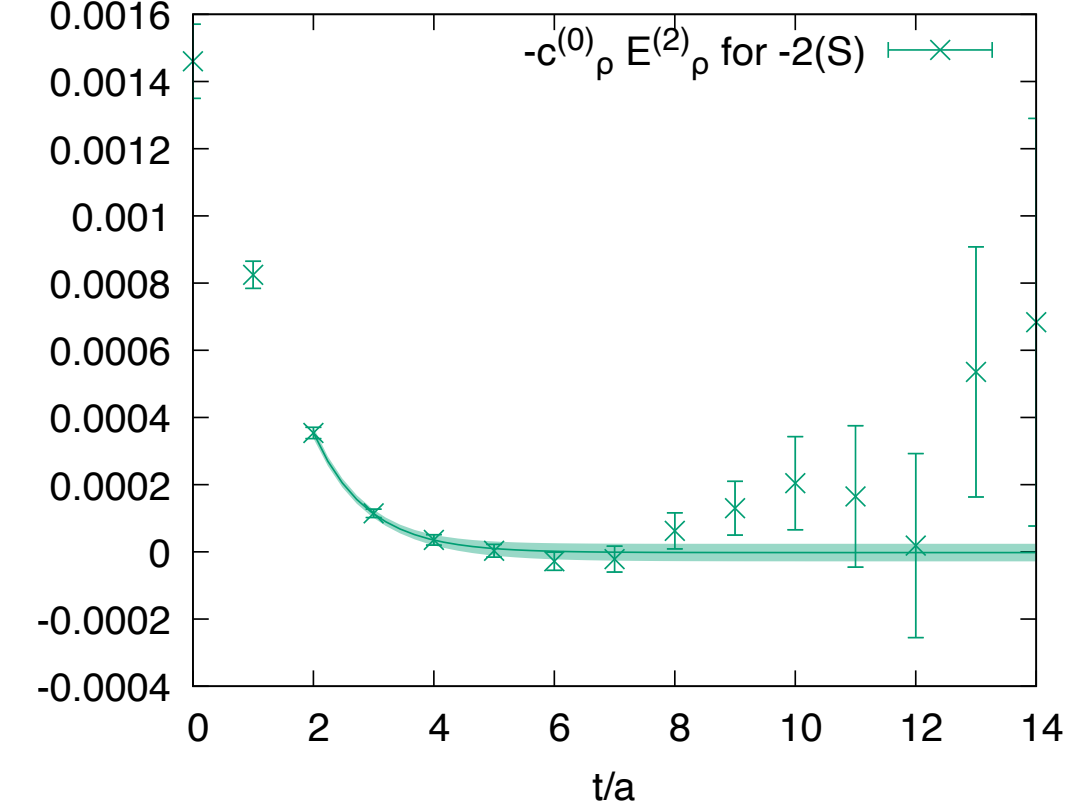
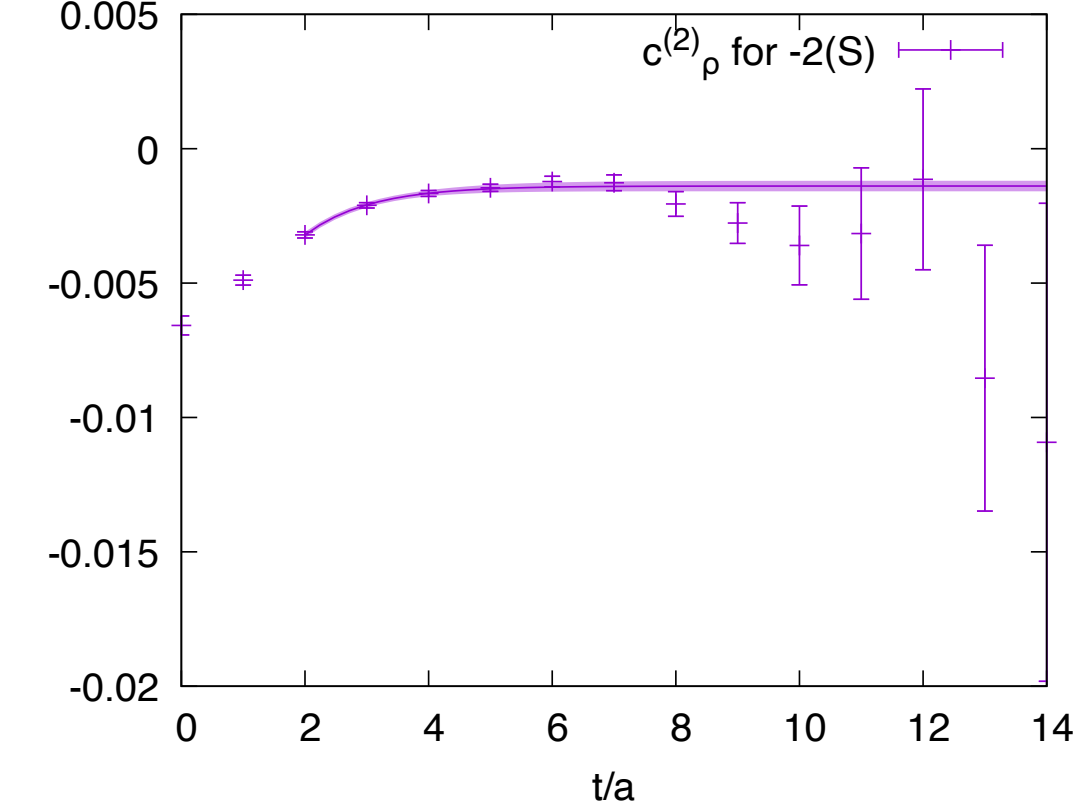
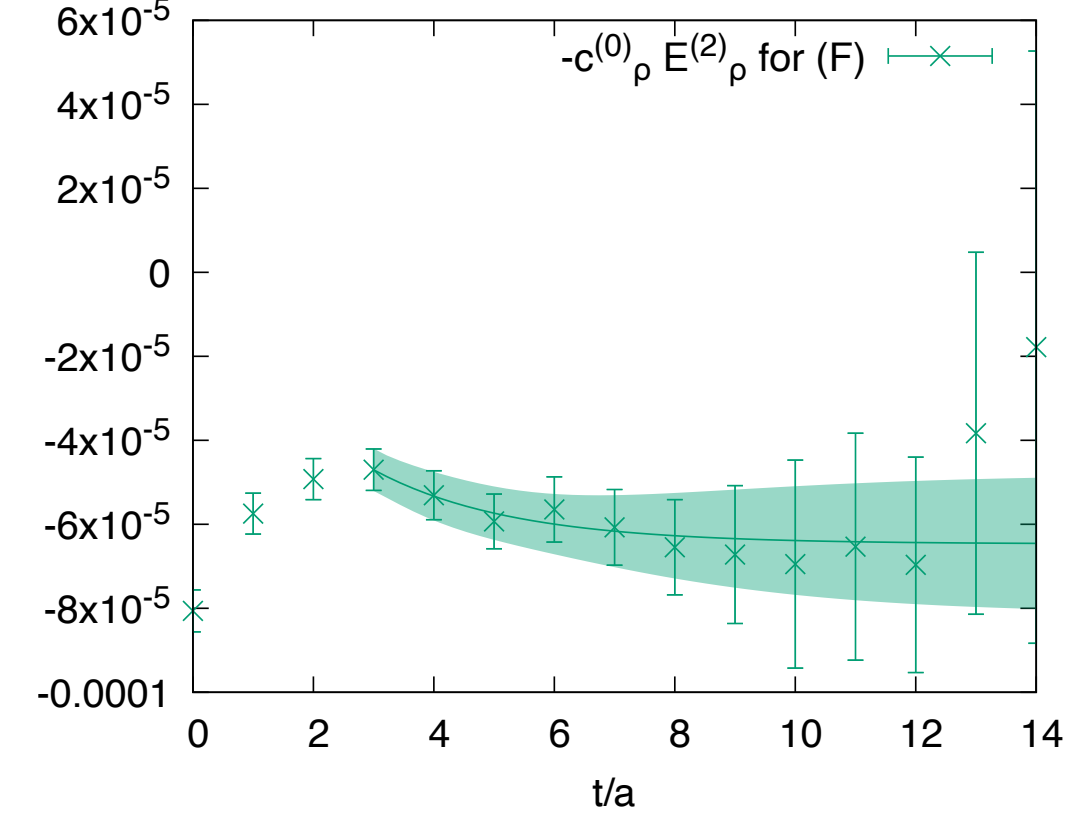
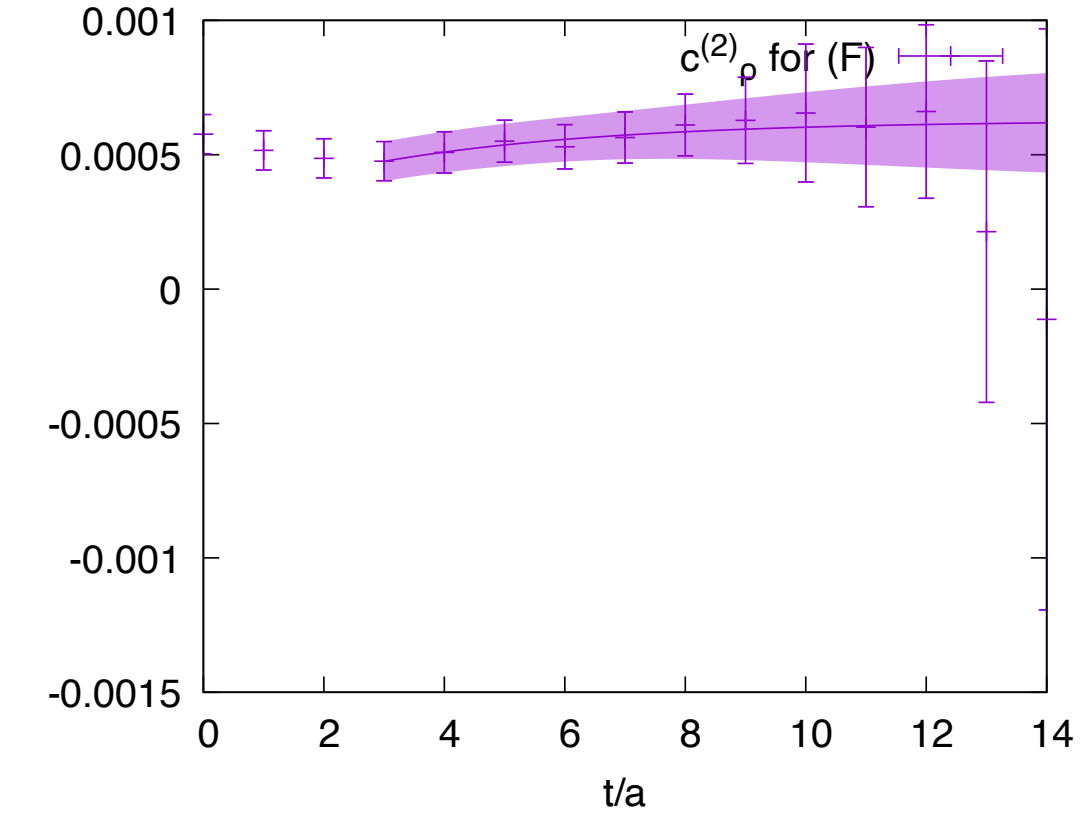
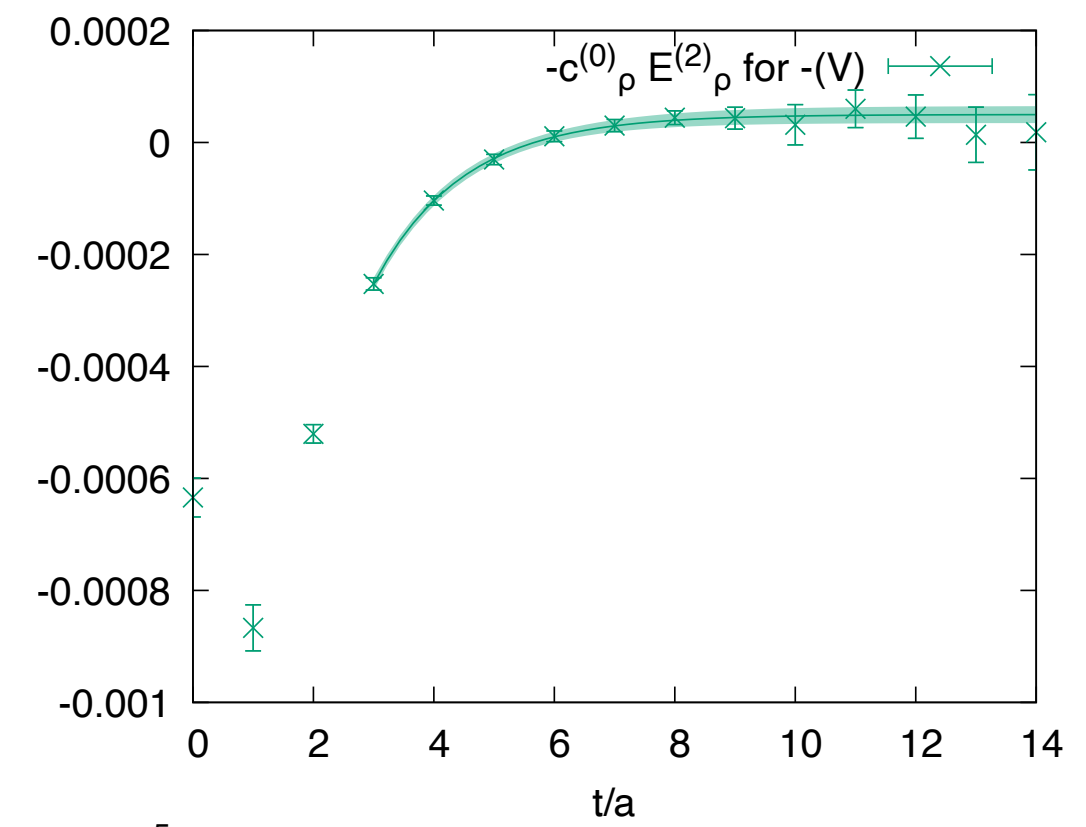
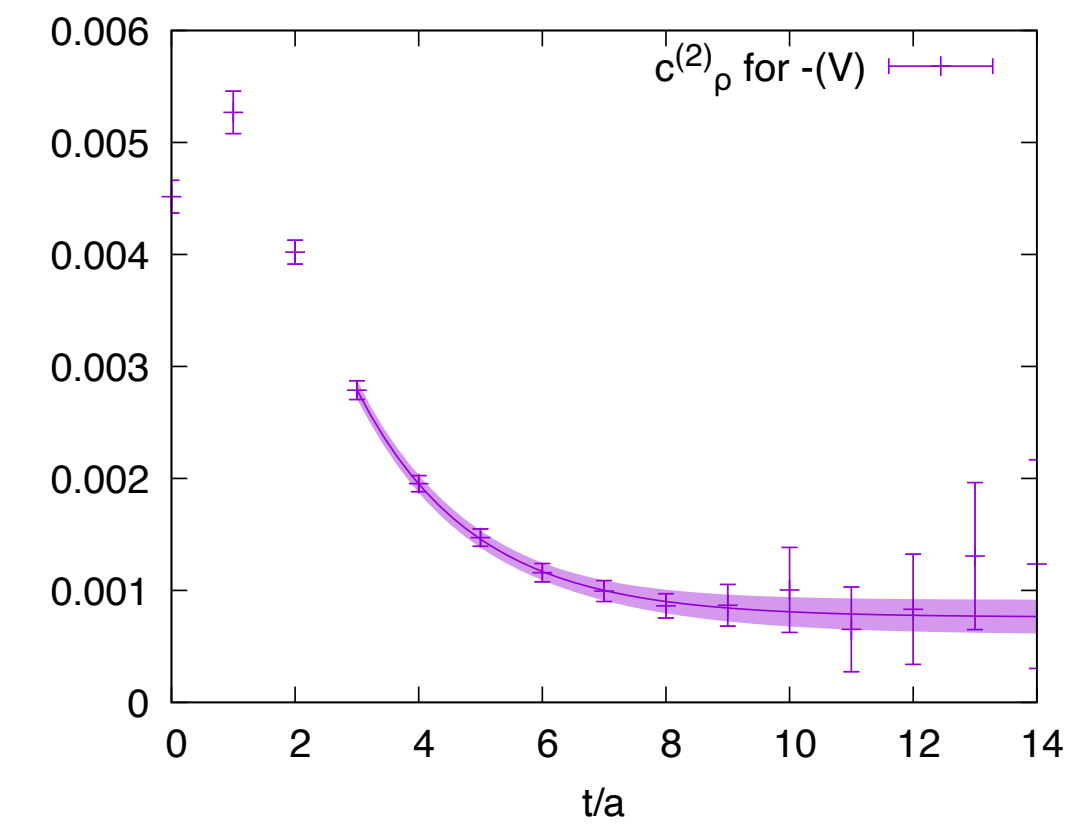
Fits are stable over fit ranges $[t/a, 24]$

$$V_{in} = \langle 0 | O_i | n \rangle \quad V = V^{(0)} + eV^{(1)} + e^2V^{(2)}$$

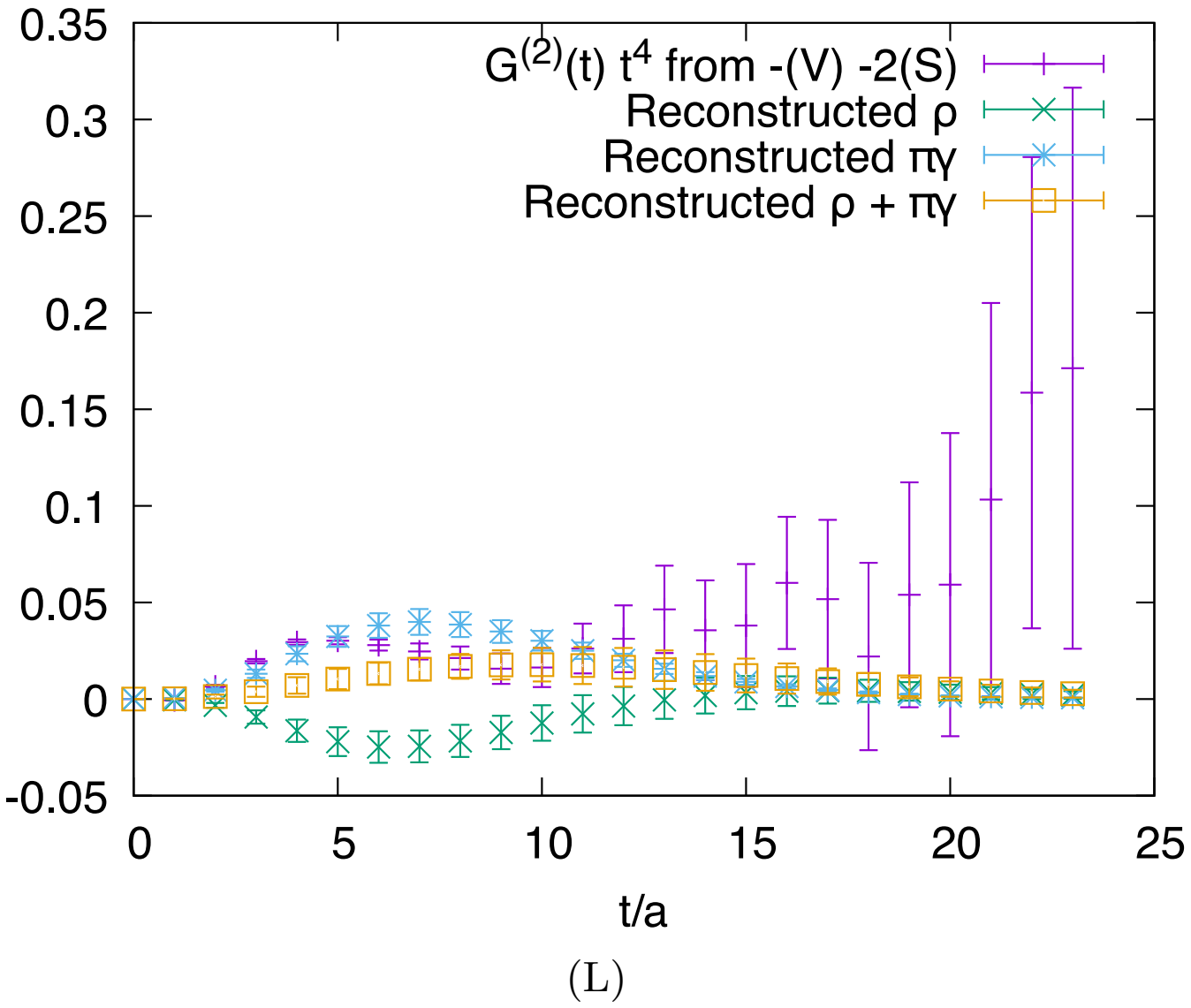
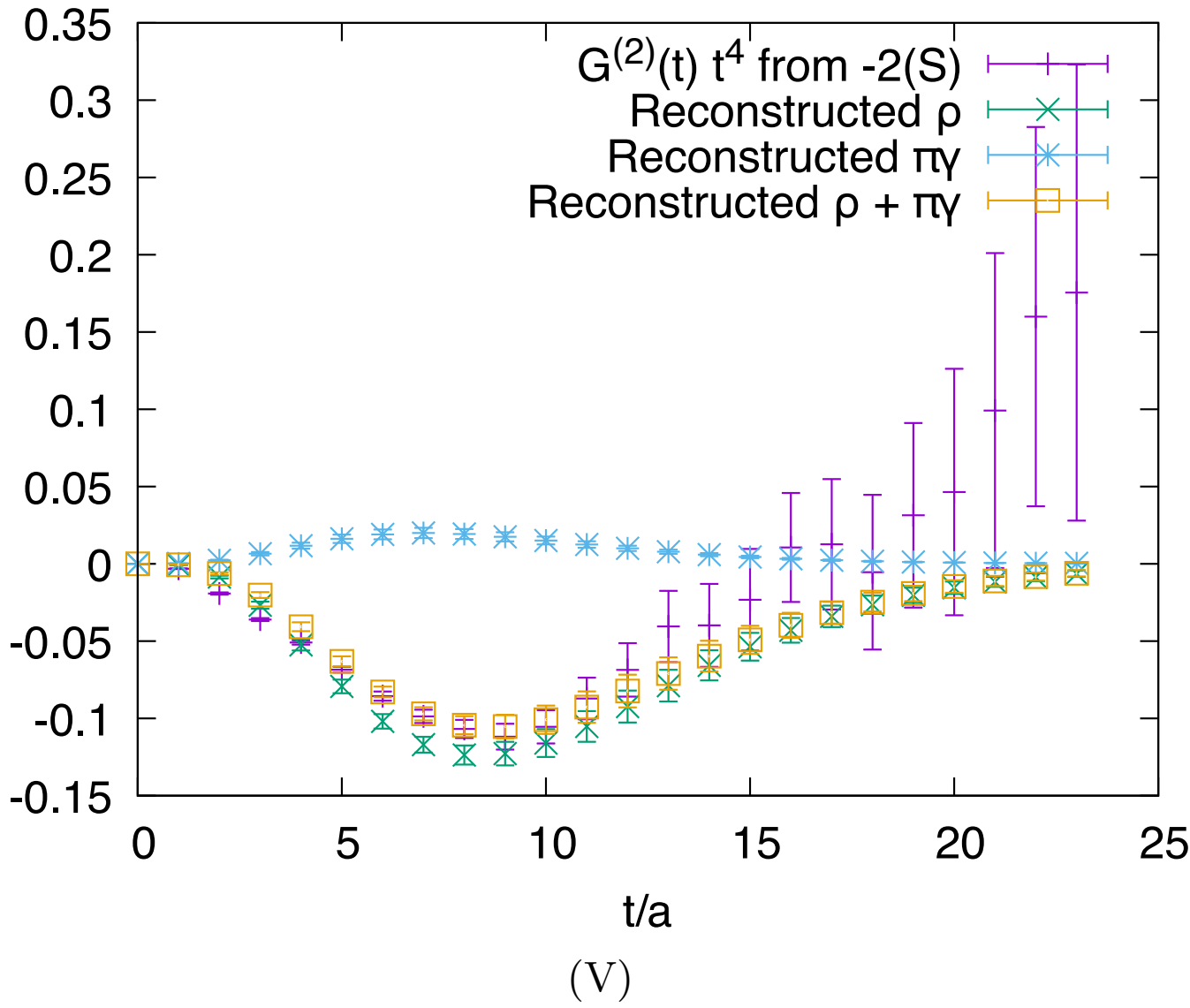
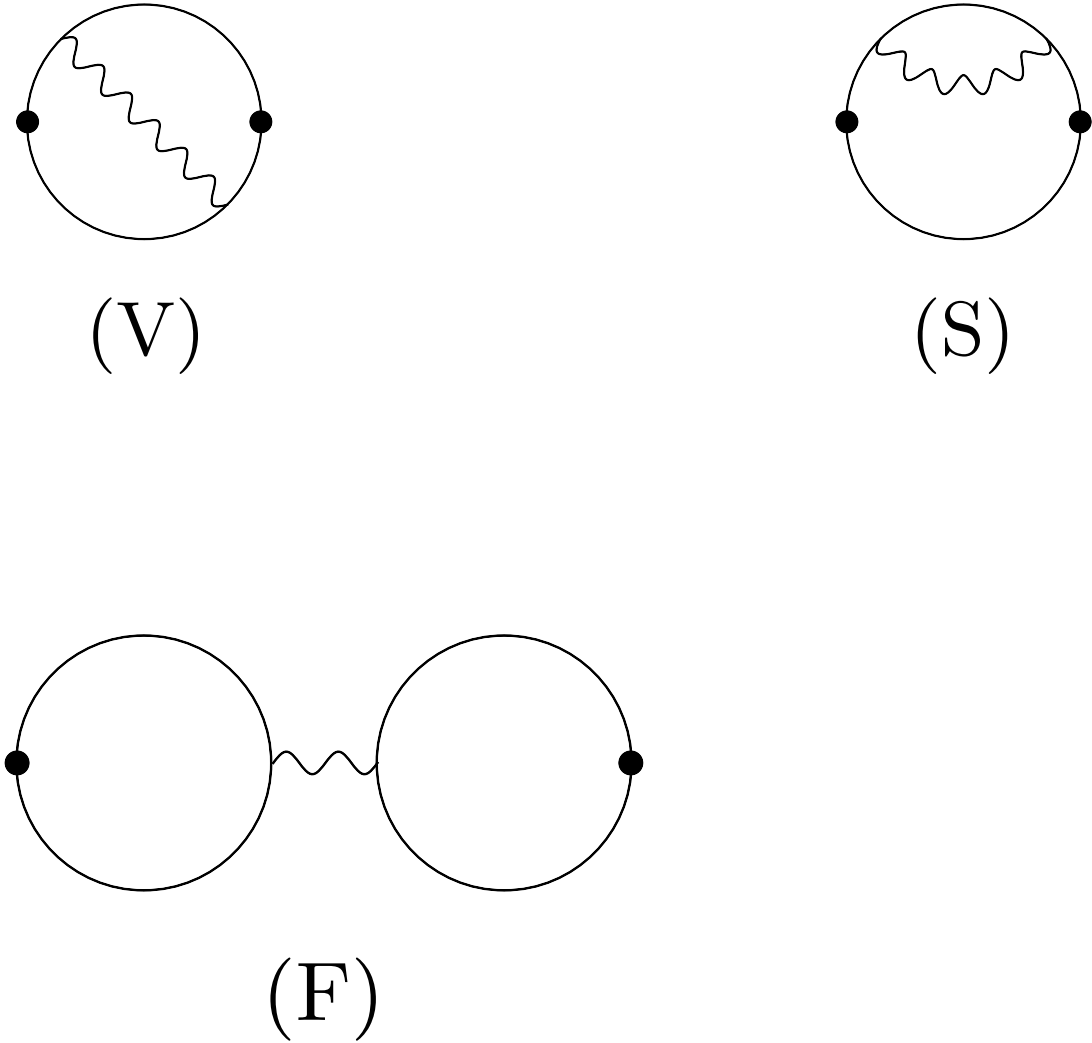
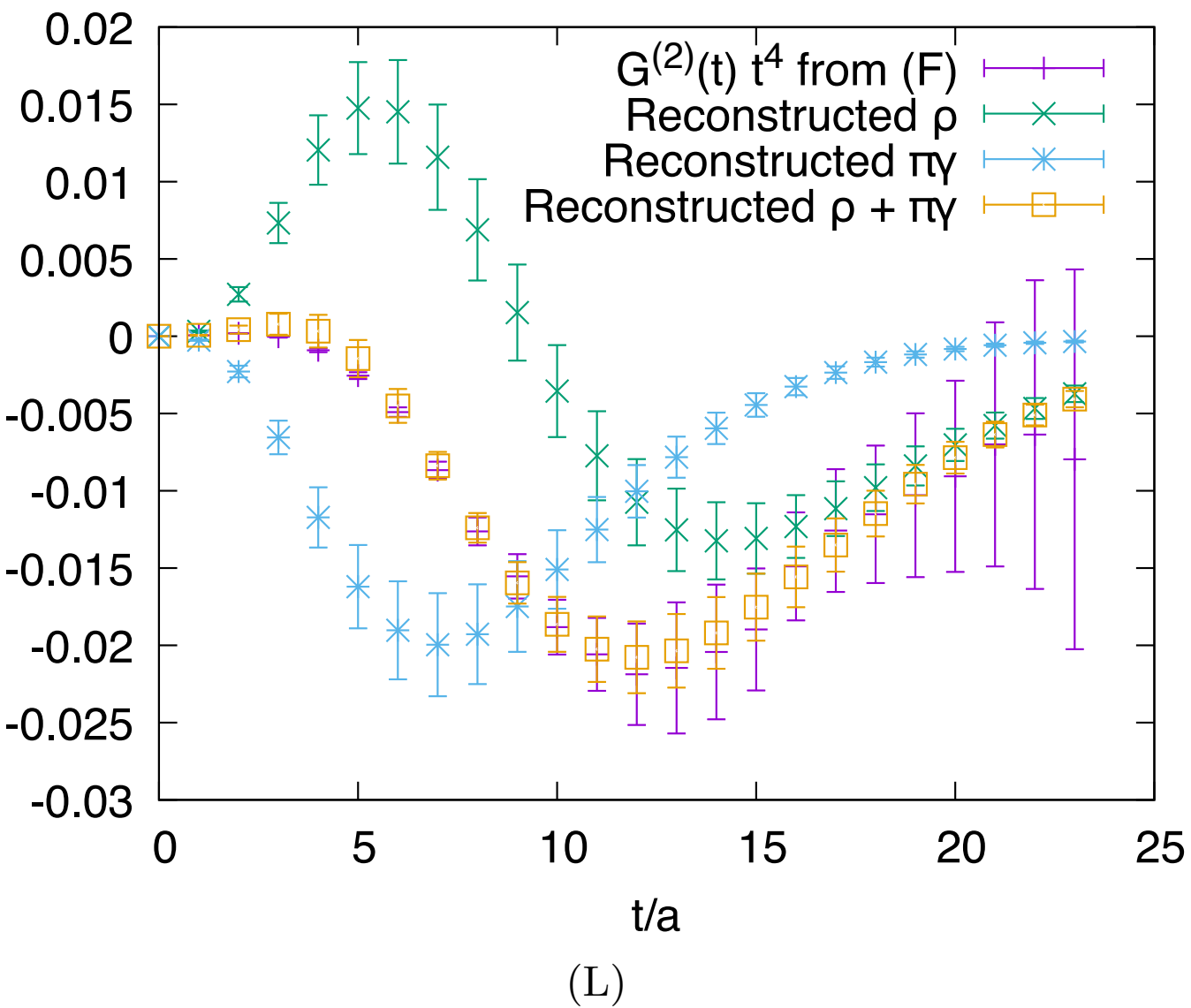
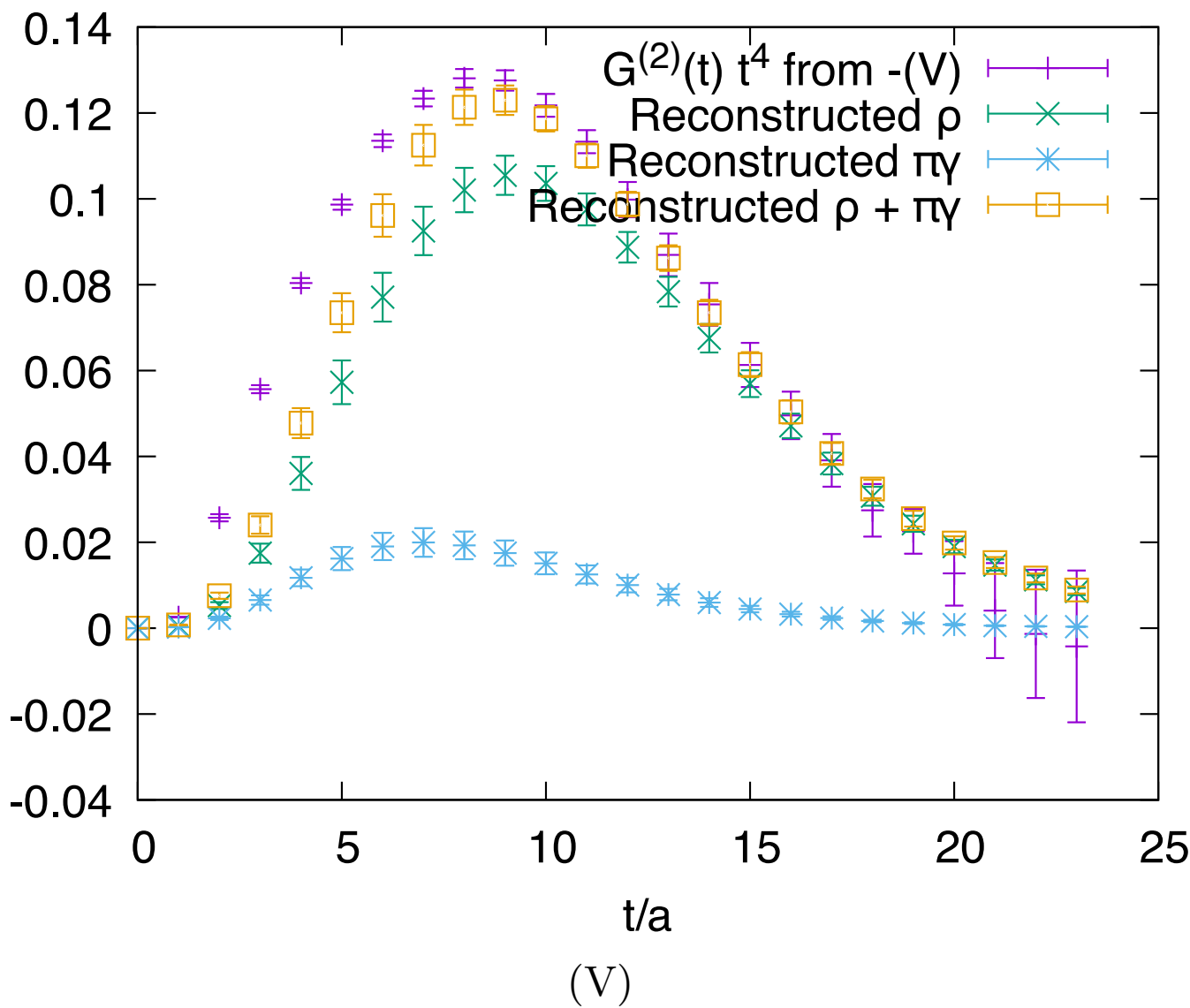
$$O_1 = V_i \quad O_2 = O_{\pi\gamma}$$

$$C(t) = \begin{pmatrix} |V_{1,\rho}^{(0)}|^2 e^{-E_\rho^{(0)}t} & 0 \\ 0 & 0 \end{pmatrix}$$

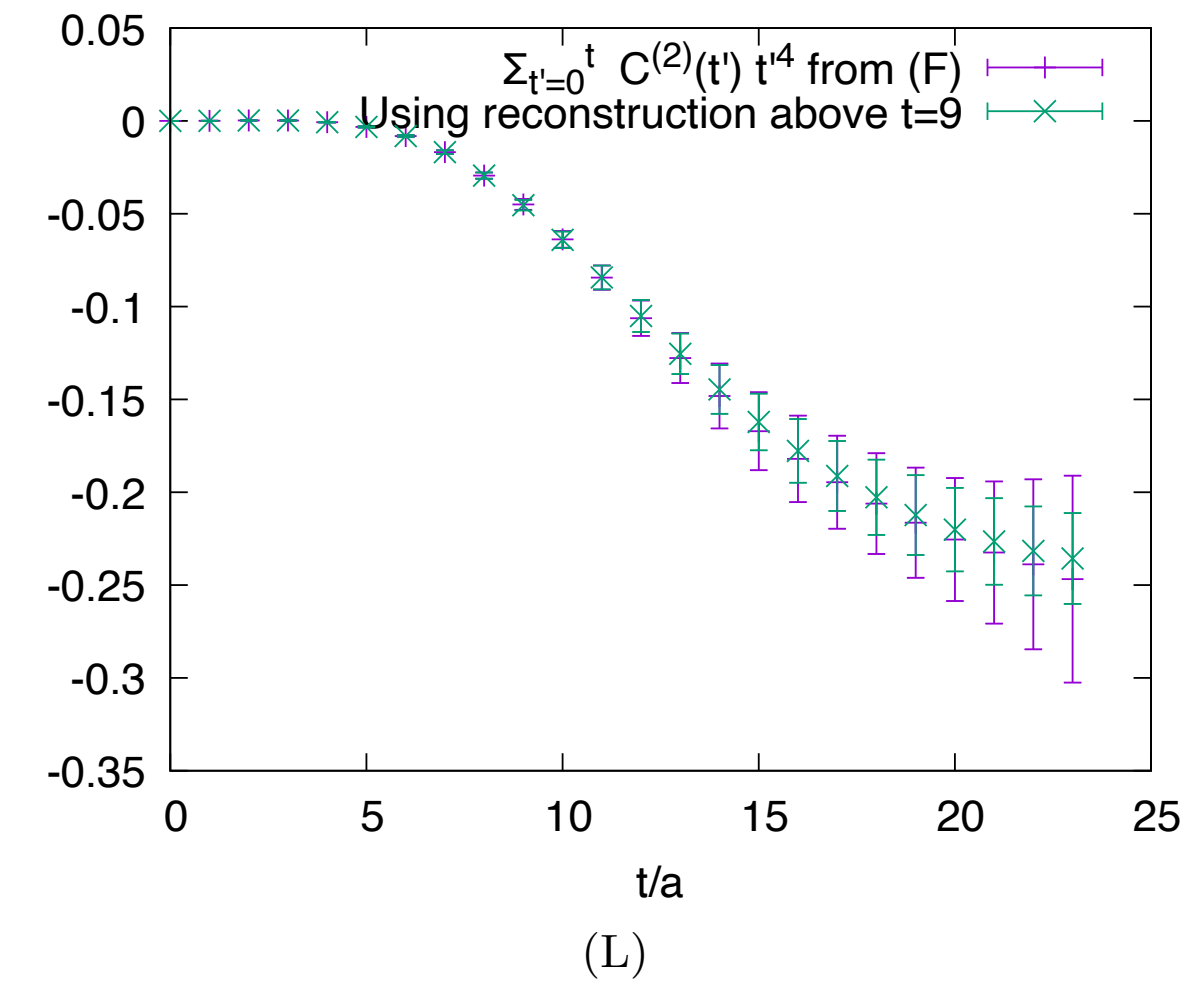
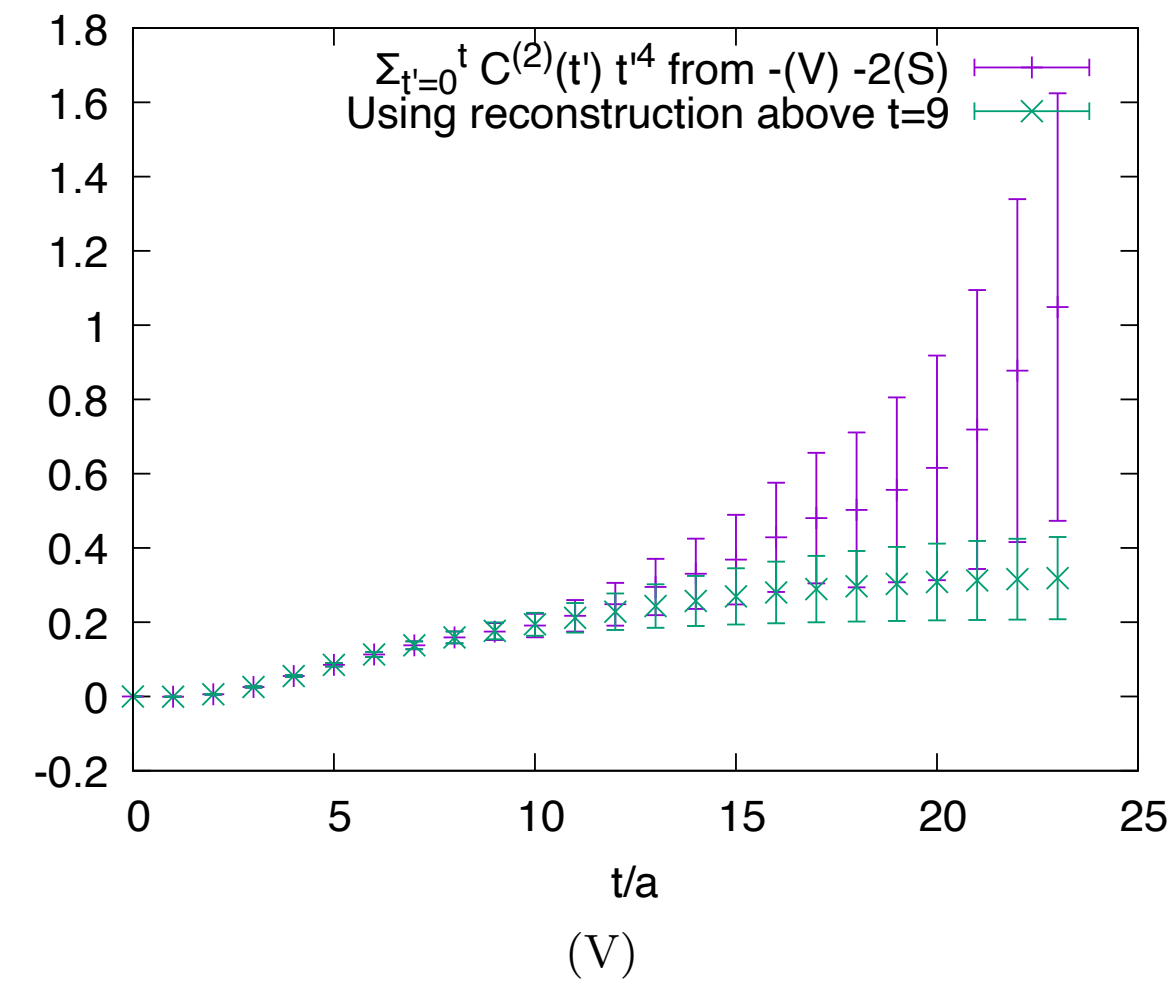
$$+ e^2 \begin{pmatrix} (2\text{Re}(V_{1,\rho}^{(0)} V_{1,\rho}^{(2),*}) - t|V_{1,\rho}^{(0)}|^2 E_\rho^{(2)})e^{-E_\rho^{(0)}t} + |V_{1,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)}t} & V_{1,\rho}^{(0)} V_{2,\rho}^{(2),*} e^{-E_\rho^{(0)}t} + V_{1,\pi\gamma}^{(1)} V_{2,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)}t} \\ V_{2,\rho}^{(2)} V_{1,\rho}^{(0),*} e^{-E_\rho^{(0)}t} + V_{2,\pi\gamma}^{(1)} V_{1,\pi\gamma}^{(1),*} e^{-E_{\pi\gamma}^{(0)}t} & |V_{2,\pi\gamma}^{(1)}|^2 e^{-E_{\pi\gamma}^{(0)}t} \end{pmatrix}$$



Compare reconstruction with pion-photon and rho compared to inclusive correlator:



Compare noise of long-distance tail with reconstruction to original:



Noise reduction by factor > 5 for noisy (S) diagram!

Next:

- extend this analysis to full list of ensembles, all diagrams, blinded analysis with multiple analysis groups
- combination of finite-volume QED with infinite-volume QED result for pion-photon contribution