

Antonio Evangelista on behalf of ETMC

Progress on the computation of the leading-order HVP contribution to the muon $g - 2$ from ETMC.



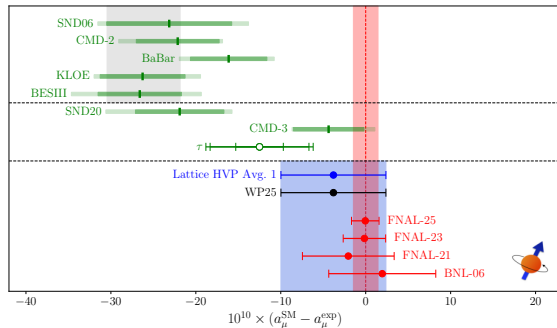
$$a_\mu = \frac{g_\mu - 2}{2} = a_\mu^{\text{QED}} + a_\mu^{\text{EW}} + a_\mu^{\text{HVP}} + a_\mu^{\text{HLbL}}$$

$$a_\mu^{\text{HVP,LO}} = 4\alpha^2 m_\mu \int_0^\infty dt C(t) t^3 K(t)$$

$$a_\mu^{\text{HVP,LO}} = \frac{\alpha^2}{3\pi^2} \int_0^\infty \frac{ds}{s} R_{\text{had}}(\sqrt{s}) \tilde{K}(s)$$

[R. Aliberti et al. – WP25, arXiv:2505.21476]

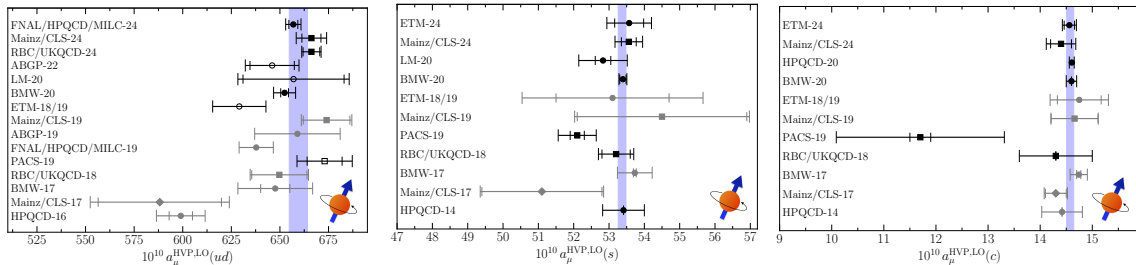
accepted on Physics Reports



- ❖ Lattice QCD $a_\mu^{\text{HVP}} \rightarrow$ no tension
- ❖ Data-driven disp. $a_\mu^{\text{HVP}} \rightarrow$ unsolved puzzle

$a_\mu^{\text{HVP,LO}}$ flavour decomposition status in the WP25 isoQCD scheme

- ❖ nice consistency among several collaborations on strange and charm quark-connected contributions
- ❖ light quark-connected contribution challenging
 - interplay of lattice artifacts, statistical errors and mass dependence
- ❖ quark-disconnected contributions: agreement among a few results, error control important too



[R. Aliberti et al. – WP25, Fig. 34]

ETMC is finalising the $a_\mu^{\text{HVP,LO}}$ in FLAG and (likely) WP25 isoQCD scheme

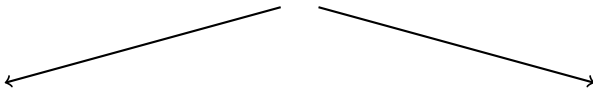
$$C(t) = -\frac{1}{3} \sum_{k=1,2,3} \int d^3 \mathbf{x} \langle 0 | \hat{J}_k(\mathbf{x}, t) \hat{J}_k(\mathbf{0}, 0) | 0 \rangle$$

$$\hat{J}_\mu(x) = \sum_{f=u,d,s,c,\dots} Z_V^f q_f \bar{\psi}_f(x) \gamma_\mu \psi_f(x)$$

mixed action setup

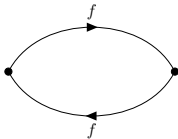


two possible definitions of vector current



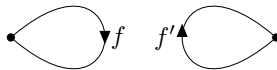
❖ connected contributions
subpercent relative error

two ways to approach the continuum



❖ disconnected contributions
larger relative error

one way to approach the continuum



ETMC strategy is using the isospin decomposition

(as proposed by Mainz/CLS)

[D. Djukanovic et al. – JHEP 04 (2025) 098]

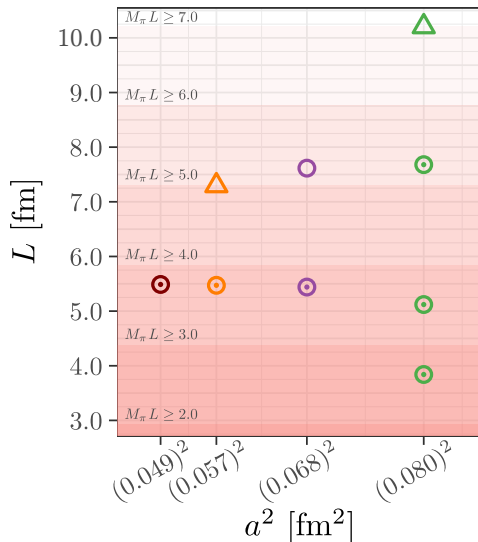
$$a_{\mu}^{\text{HVP}}(I=1) = \frac{9}{10} a_{\mu}^{\text{HVP}}(\ell)$$

$$a_{\mu}^{\text{HVP}}(I=0) = \frac{1}{10} a_{\mu}^{\text{HVP}}(\ell) + a_{\mu}^{\text{HVP}}(s) + a_{\mu}^{\text{HVP}}(c) + a_{\mu}^{\text{HVP}}(\text{disc.})$$

ensembles very close to the FLAG/Edinburgh and WP25 isoQCD definition

Simulation status: $\circ$ done $\triangle$ ongoing

Analysis status: $\odot$ done $\bigcirc$ planned



Used in this analysis:

- ❖ 4 different β 's $\rightarrow a \in 0.05 - 0.08$ fm
- ❖ Different linear size $\rightarrow L \in 3.8 - 7.6$ fm
- ❖ Careful analysis of mistuning on μ_ℓ , μ_s , μ_c and m_{cr} to match exactly the FLAG isoQCD point

❖ Bounding method

$$0 \leq C(t_c) e^{-m_{\text{eff}}(t_c)(t-t_c)} \leq C(t) \leq C(t_c) e^{-E_0(t-t_c)} \quad \forall t > t_c$$

❖ E_0 is the energy of the lowest two-pion state ($I = 1$) and three-pion state ($I = 0$)

❖ $m_{\text{eff}}(t_c)$ can be safely replaced by $m_{\text{eff}}(t^*)$ choosing a suitable $t^* < t_c$

[T. Blum et al. – Phys.Rev.Lett. 121 (2018) 2]

❖ Additive blinding of the $I = 1$ data

$$C^{\text{reg}}(an) \rightarrow C^{\text{reg}}(an) + \frac{a^3}{Z_V^{\text{reg}}} \sum_{k=1}^K A_k e^{-m_k \frac{T}{2}} \cosh \left(m_k \left(\frac{T}{2} - an \right) \right)$$

❖ large enough K (we chose $K > 12$) does not allow for unblinding data

❖ does not alter the cut-off effects $\rightarrow$ removal of free-theory lattice artifacts possible;

❖ Finite Volume Effects (FVE) not altered;

❖ Additive a - and L - independent blinding of the $I = 0$ data

$l = 1$ contribution

mistuning corrections

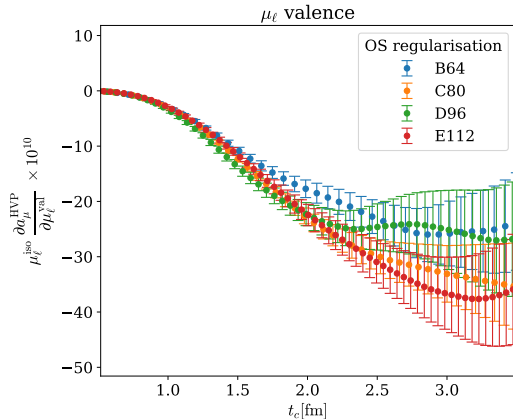
$$\mathcal{O}(\mu_f, m_{\text{cr}} | \mu_f^{\text{val}}, m_{\text{cr}}) = \mathcal{O}^{\text{sim}}(\mu_f^{\text{val}}, m_{\text{cr}}^{\text{sim}})$$

$$+ (m_{\text{cr}} - m_{\text{cr}}^{\text{sim}}) \left[\partial_{\text{cr}}^{\text{sea}} + \partial_{\text{cr}}^{\text{val}} \right] \mathcal{O}$$

$$+ \sum_{f=\ell, s, c} (\mu_f - \mu_f^{\text{sim}}) \partial_f^{\text{sea}} \mathcal{O}$$

$$+ (\mu_\ell - \mu_\ell^{\text{sim}}) \partial_\ell^{\text{val}} \mathcal{O}$$

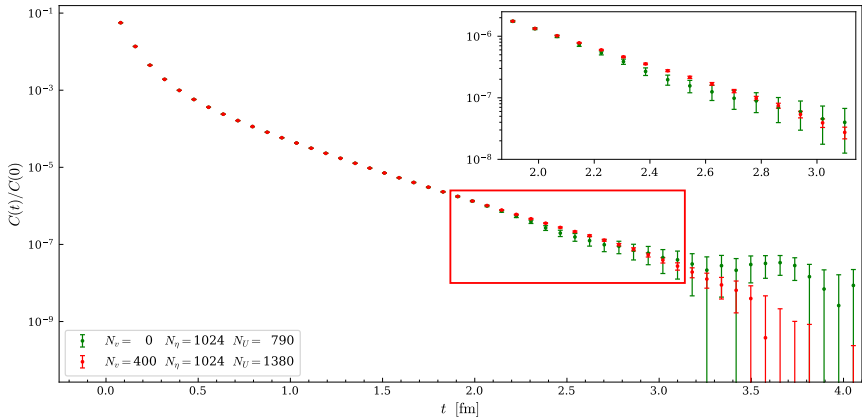
- ❖ m_{cr} both sea and valence $\longrightarrow$ small due to $\frac{m_{\text{PCAC}}}{\mu_\ell} \ll 1$ tuning, correction size is β dependent
- ❖ μ_f sea with re-weighting $\longrightarrow$ typically small or similar w.r.t. statistical uncertainty
- ❖ μ_ℓ valence $\longrightarrow$ dedicated measurements at μ_ℓ close to the isoQCD FLAG point



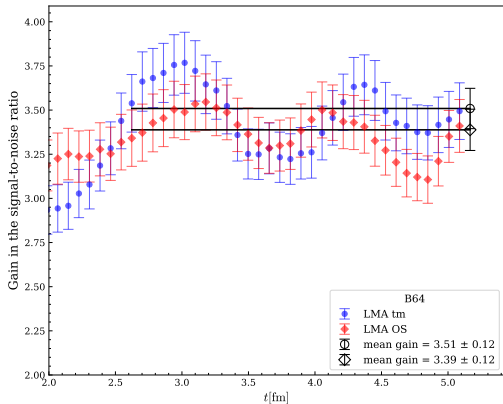
Low Mode Averaging setup

$$S_{r,\eta}^{\text{defl}}(x, y) - S_{r,\eta}(x, y) = \sum_{j=1}^{N_\nu} \frac{1}{\lambda_j + i r \mu} |v_j(x)\rangle \langle v_j(y)| \left(\mathbb{1} - |\eta(x)\rangle \langle \eta(y)| \right) \gamma_5 = S_r^{\text{IR}}(x, y) - S_{r,\eta}^{\text{IR}}(x, y).$$

$$C_{rr',\eta}^{\text{defl}}(t) = C_{rr',\eta}(t) - C_{rr',\eta}^{\text{IR}}(t) + C_{rr'}^{\text{IR}}(t)$$



S/N gain on vector-vector correlators



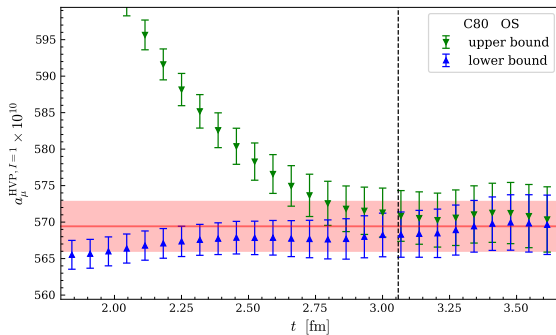
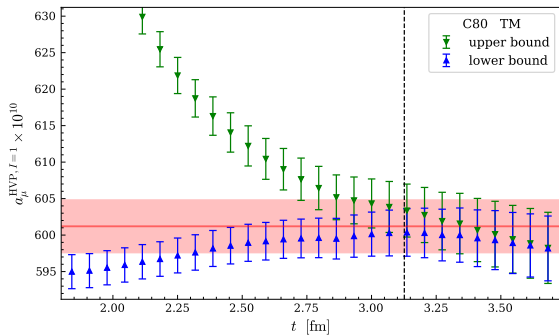
❖ N_v chosen to obtain a gain in $S/N \gtrsim 3.5$

❖ for a fixed gain the N_v values depend on physical volume

Ensemble	L [fm]	N_U	N_v	N_η	$a\mu_\ell^{\text{sea}}$
B64	5.09	1380	400	1024	0.00072
C80	5.46	870	530	1600	0.00060
D96	5.46	450	530	960	0.00054
E112	5.46	420	530	1344	0.00044

C80 ensemble ($L \simeq 5.5$ fm) bounding

- ❖ upper bound $\rightarrow$ two-pion (lattice) state with lowest relative momentum
- ❖ lower bound $\rightarrow$ $m_{\text{eff}}(t^*)$ with $t^* = 1.8$ fm for all ensembles and both regularisation
- ❖ average over $\delta t \simeq 0.25$ fm after optimal t_c



reference volume interpolation $L_{\text{ref}}=5.46$ fm

❖ B64 ensemble $L \simeq 5.1$ fm

❖ B96 ensemble $L \simeq 7.6$ fm

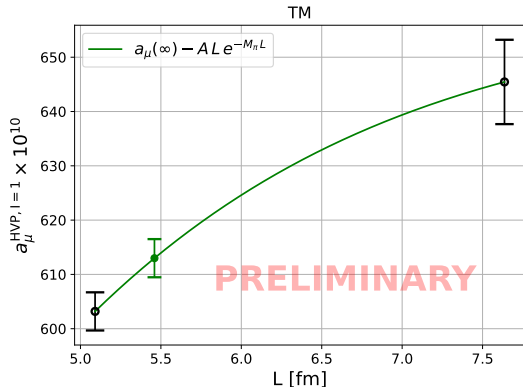


preliminary exponential interpolation
GS and HP inspired

[G.J. Gounaris and J.J. Sakurai – Phys.Rev.Lett. 21 (1968)]

[M.T. Hansen and A. Patella – Phys.Rev.Lett. 123 (2019) 172001]

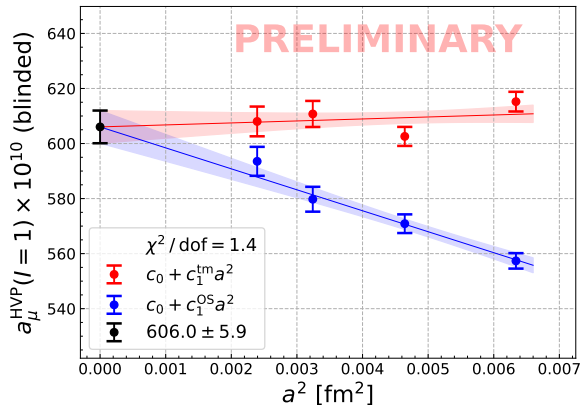
$$a_\mu(L) = a_\mu(\infty) - AL e^{-M_\pi L}$$



final interpolation via data-driven lattice-TM GS model

continuum limit extrapolation

$$a_{\mu}^{\text{HVP,reg}}(l=1) = P_0 + P_1^{\text{reg}} a^2$$



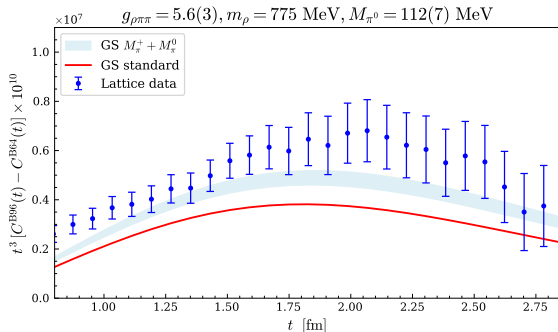
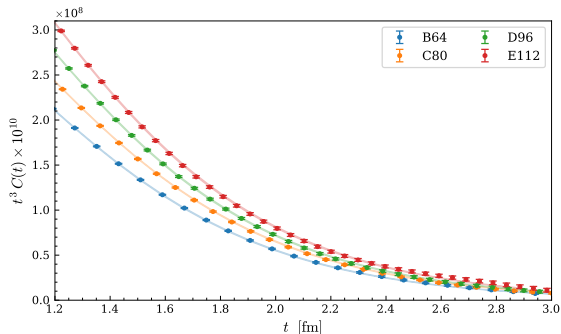
- ❖ $\sim 4 \cdot 10^{-10}$ error at fixed β
- ❖ increase of statistics on D96 and E112
- ❖ TM regularisation data almost flat in a^2
compatible with constant within errors
- ❖ C.L. systematics to be estimated via BAIC
 $a^2 + a^4$ **terms on OS data**
 $a^2 + a^2 / \log(a^2 \Lambda^2), \dots$ **terms**

infinite volume limit

- ❖ lattice-TM GS model using $M_{\pi^\pm}$ and M_{π^0} for 2π energy levels
- ❖ standard GS parameters evaluated from global fit at 4β on TM regularisation
- ❖ validation for FVE between $L \simeq 5.1$ fm (B64) and $L \simeq 7.6$ fm (B96) with a target error $\sim 10\%$ in progress



increase of statistic on B96 in progress



$l = 0$ contribution

- ❖ strange and charm quark-connected contribution already published

[C. Alexandrou *et al.* – Phys.Rev.D 111 (2025) 5, 054502]

no S/N problem at large Euclidean time

- ❖ no bounding

- ❖ no LMA

- ❖ mistuning for strange and charm quark-connected contributions taken into account

- ❖ mistuning on all quark-disconnected contributions negligible within errors

- ❖ charm quark-disconnected contribution expected to be negligible → check in progress at a single beta

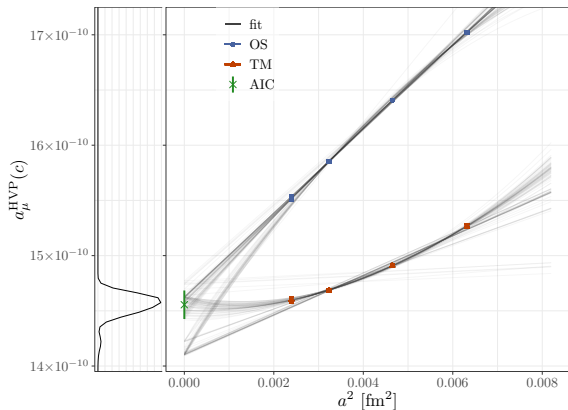
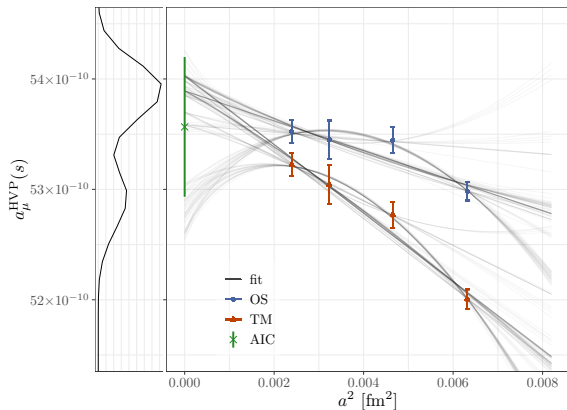
only light and strange contributions taken into account (for now)

$$a_{\mu}^{\text{HVP}}(I=0) = \frac{1}{10} a_{\mu}^{\text{HVP}}(\ell) + a_{\mu}^{\text{HVP}}(s) + a_{\mu}^{\text{HVP}}(\text{disc.}; \ell - s)$$

strange and charm quark-connected contributions

- ❖ FVE negligible but taken into account as systematic effect
- ❖ continuum limit systematics via Bayesian AIC

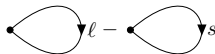
$$a_\mu^{\text{HVP,reg}}(f) = P_0 + P_1^{\text{reg}} a^2 + P_2^{\text{reg}} a^4 \quad f = s, c$$



light-strange disconnected correlators estimation

❖ $B(\ell) - B(s)$ computable via Axial WI

$$\sum_{f=u,d,s} q_f B_f^\mu(z) = Z_V(\mu_\ell + \mu_s) a^4 \sum_x \langle P_{\ell s}(x) A_{\ell s}^\mu(z) \rangle ,$$



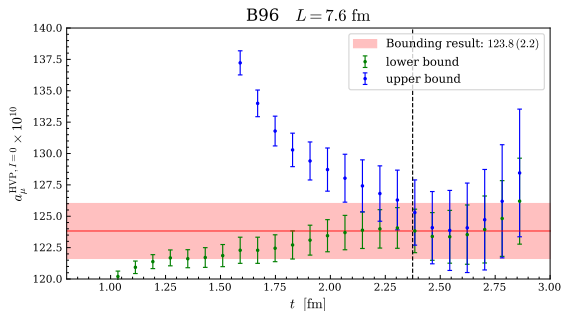
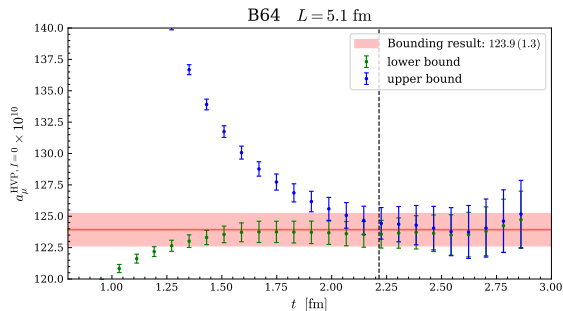
❖ frequency splitting improvement with 4 levels ($N = 3$)

[L. Giusti et al. – Eur.Phys.J.C 79 (2019) 7, 586]

$$\begin{aligned} B^\mu(\mu_\ell) - B^\mu(\mu_s) &= [B^\mu(\mu_\ell) - B^\mu(\mu_1)] \\ &\quad + [B^\mu(\mu_1) - B^\mu(\mu_2)] \\ &\quad + \dots + [B^\mu(\mu_N) - B^\mu(\mu_s)] . \end{aligned}$$

$I = 0$ bounding procedure

- ❖ lower bound $\rightarrow$ correlator positivity
- ❖ upper bound $\rightarrow E_0 = 2E_{\pi^\pm}(\mathbf{p}) + m_{\pi^0}$ with $\mathbf{p} = \pm \frac{2\pi}{L} \mathbf{n}$ (state relevant only at $\mathcal{O}(a)$ in tm-LQCD)
- ❖ average over $\delta t \simeq 0.25$ fm after optimal t_c



FVE negligible within statistical errors as expected

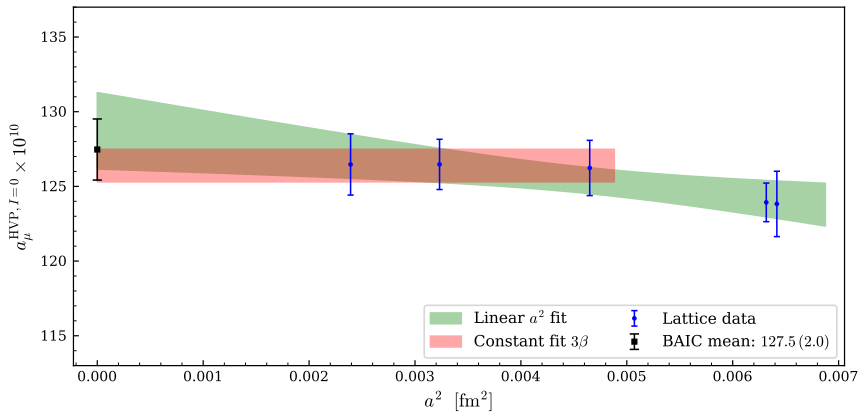
very small lattice artifacts



◆ constant fit with 3β

◆ linear fit in a^2 with 4β

blinded preliminary result



Leading isospin breaking corrections

towards LIB effects from ETMC



computing $a_\mu^{\text{HVP,LO}}$ corrections via RM123 method

$$\Delta a_\mu^{\text{HVP}}(f) = 2\alpha^2 \lim_{t_c \mapsto +\infty} \int_0^{t_c} dt \, t^2 K(m_\mu t) Z_V^2 \Delta C^f(t) + 2 \frac{\Delta Z_{V,f}}{Z_V} \left(a_\mu^{\text{HVP}}(f) \right)^{\text{iso}}$$

- ❖ valence contributions $\longrightarrow$ started and well advanced
- ❖ sea and sea-valence contribution $\longrightarrow$ starting now

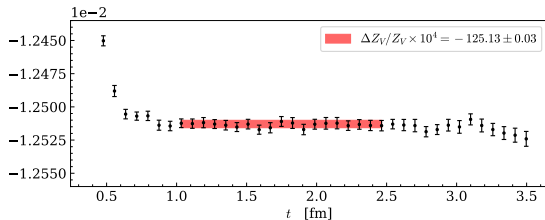
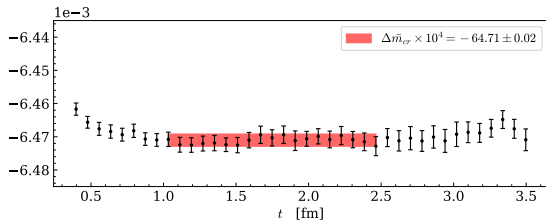
valence quark-connected leading isospin breaking corrections

$$\Delta C^f(t) = e_f^2 \frac{\partial C^f}{\partial e^2}(t) + \Delta m_{cr}^f \frac{\partial C^f}{\partial m_{cr}^f}(t) + \Delta \mu_f \frac{\partial C^f}{\partial \mu_f}(t)$$

❖ e^2 → discrete derivative

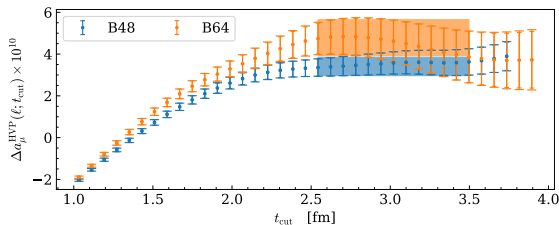
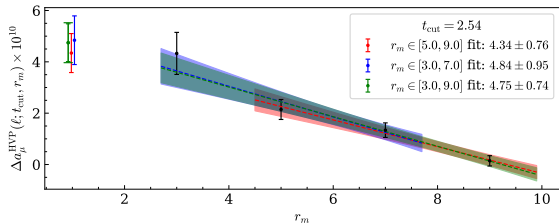
❖ μ_f, m_{cr} → operator insertion

- ❖ Δm_{cr} from parity restoration in QCD+QED
- ❖ $\Delta \mu_f$ from matching physical $\pi^\pm, K^\pm, K^0$ and D_s meson masses
- ❖ $\Delta Z_V^f/Z_V$ from axial WI as in tmQCD by using mixed action setup

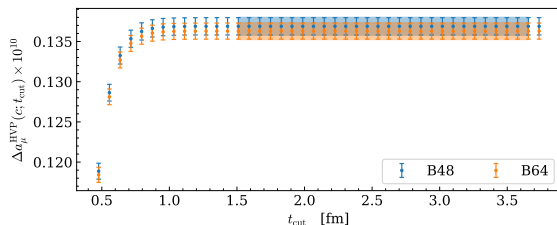
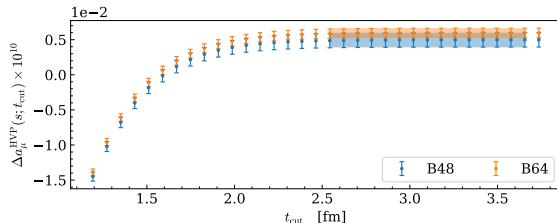


2 volumes: $L_1 \simeq 3.8$ fm and $L_2 \simeq 5.1$ fm at fixed lattice spacing $\simeq 0.08$ fm

❖ $u + d$ contribution needs chiral extrapolation



❖ strange and charm show no S/N problem



isoQCD results

- ❖ $a_\mu^{\text{HVP}}(s)$ and $a_\mu^{\text{HVP}}(c)$ published with $\sim 1\%$ accuracy
- ❖ $a_\mu^{\text{HVP}}(I=0)$ preliminary blinded result with competitive error ($\sim 2 \times 10^{-10}$)
- ❖ $a_\mu^{\text{HVP}}(I=1)$ almost completed
 - increase of statistics on the two finest lattice spacings and on the large volume
 - careful continuum extrapolation study to establish systematics (target error $\sim 5 \times 10^{-10}$)
 - TM-lattice GS model for FVE estimation under validation (target error $\sim 3 \times 10^{-10}$)



final accuracy on a_μ^{HVP} close or below 1% achievable

leading valence isospin breaking effects

- ❖ valence quark-connected data ready to be analysed on 2 finer lattice spacings
- ❖ third volume ($L \simeq 7.6$ fm) at coarsest lattice spacing to check FVE

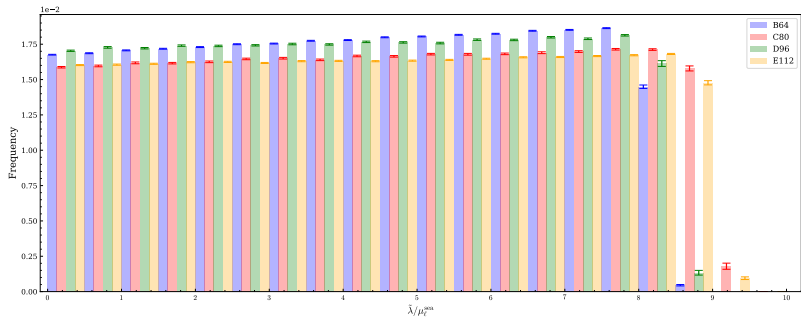
future directions

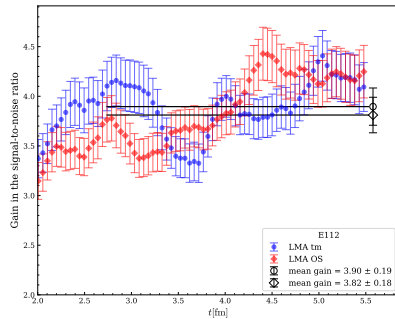
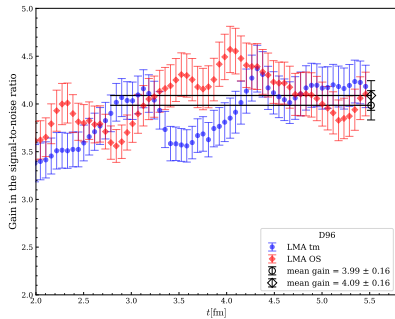
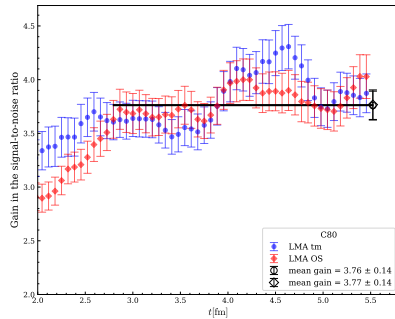
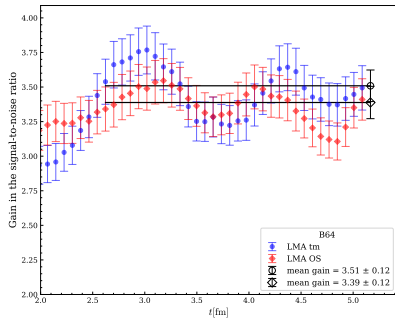
- ❖ **isoQCD** simulations at a bigger volume $L \simeq 10.2$ fm just started
 - ❖ including **electro-unquenched** leading isospin breaking effects at several lattice spacings
 - ❖ **smeared R -ratio** with energy resolution down to ~ 200 MeV just started
- interesting in view of experimental $e^+e^- \mapsto$ hadrons updates

Thank you for the attention!

Backup

$$\frac{\Delta \mathcal{N}}{\Delta \tilde{\lambda}}(0) \propto \Lambda_{\text{QCD}}^3 L^3 \mathcal{T}.$$





μ_ℓ valence mistuning

