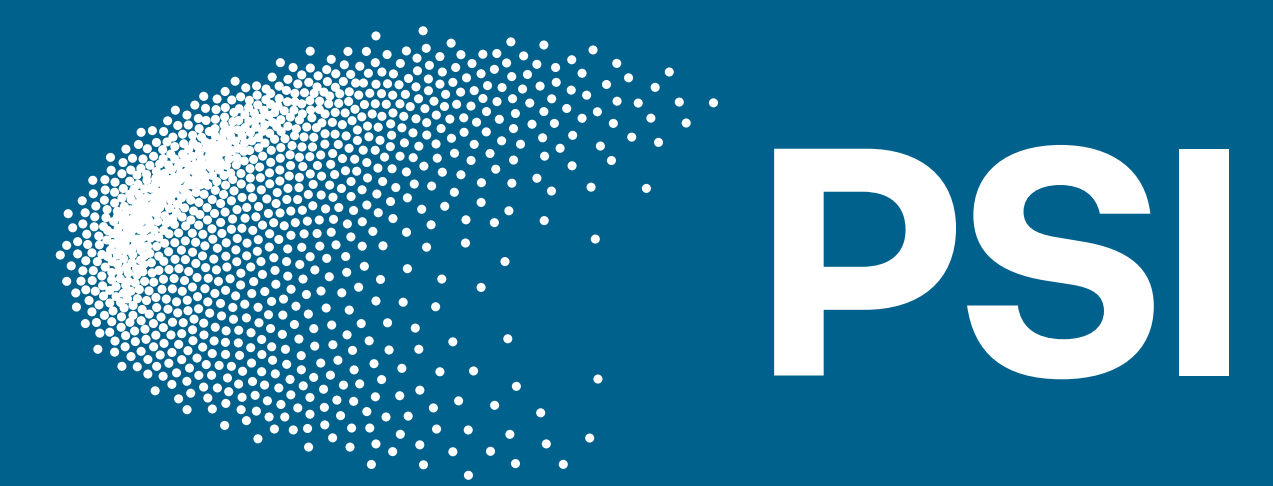


A MODEL FOR LIGHT TENSOR MESON TRANSITION FORM FACTORS



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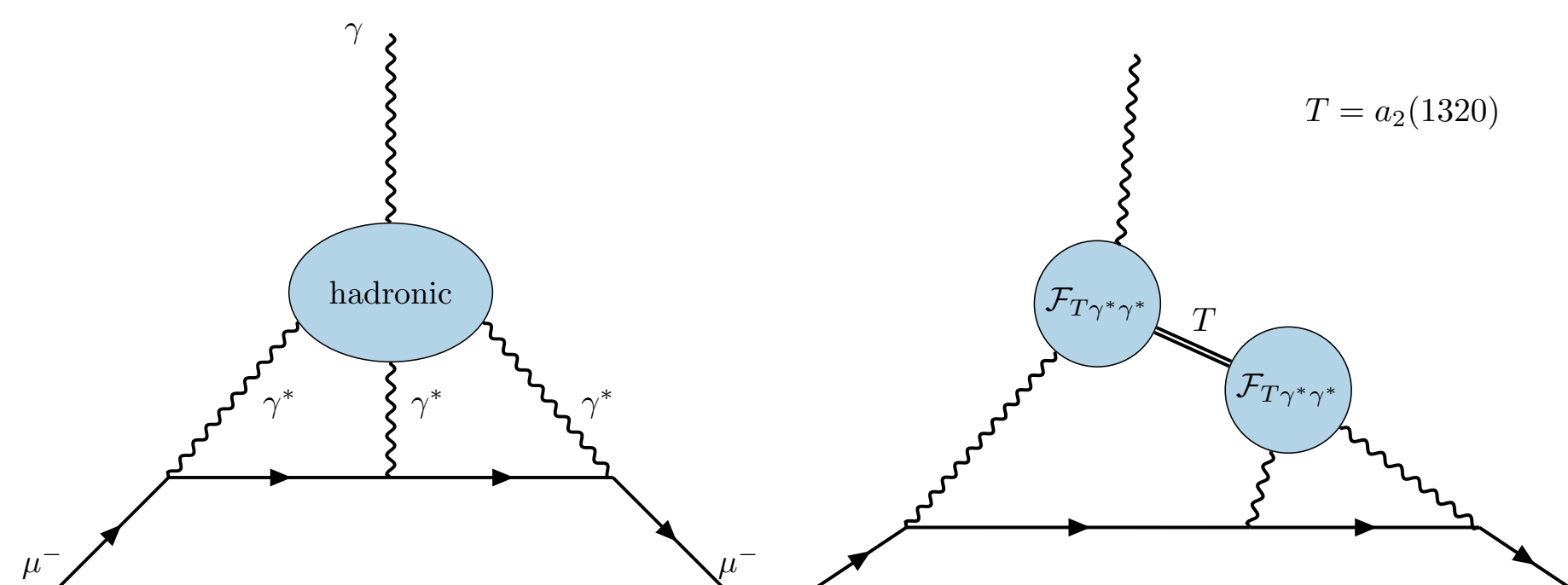
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MOTIVATION : $(g - 2)_\mu$

Anomalous magnetic moment of the muon: hadronic contributions a_μ^{had}

- data-driven dispersive approach for hadronic contributions, hadronic vacuum polarization (HVP) and **hadronic light-by-light** (HLbL)

Focus on **tensor meson** resonance contribution to the HLbL



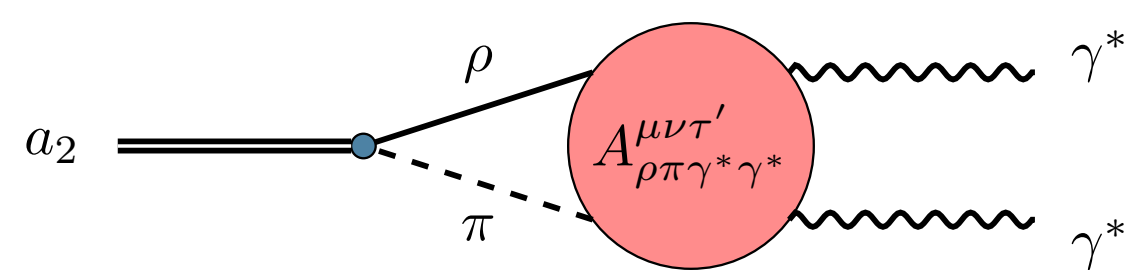
- description for the transition form factors (TFFs) $\mathcal{F}_{T\gamma^*\gamma^*}$, involving vector-meson dominance and dispersive treatment

Tensor mesons of $J^{PC} = 2^{++}$

- multiplet of $\{a_2(1320), K_2^*(1430), f_2'(1525), f_2(1270)\}$
- specific case of $T = a_2(1320)$
 - ↪ utilize the main decay channel of 3π (or $\rho\pi$ via isobar model)
 - ↪ difficulty for phenomenological parameterization or modeling of $\mathcal{F}_{T\gamma^*\gamma^*}$ due to scarcity of data
- need 5 TFFs to describe the $T \rightarrow \gamma^*\gamma^*$ process

FRAMEWORK

Model the dynamical behavior of $\mathcal{F}_{a_2\gamma^*\gamma^*}$ via $\rho\pi$ intermediate state

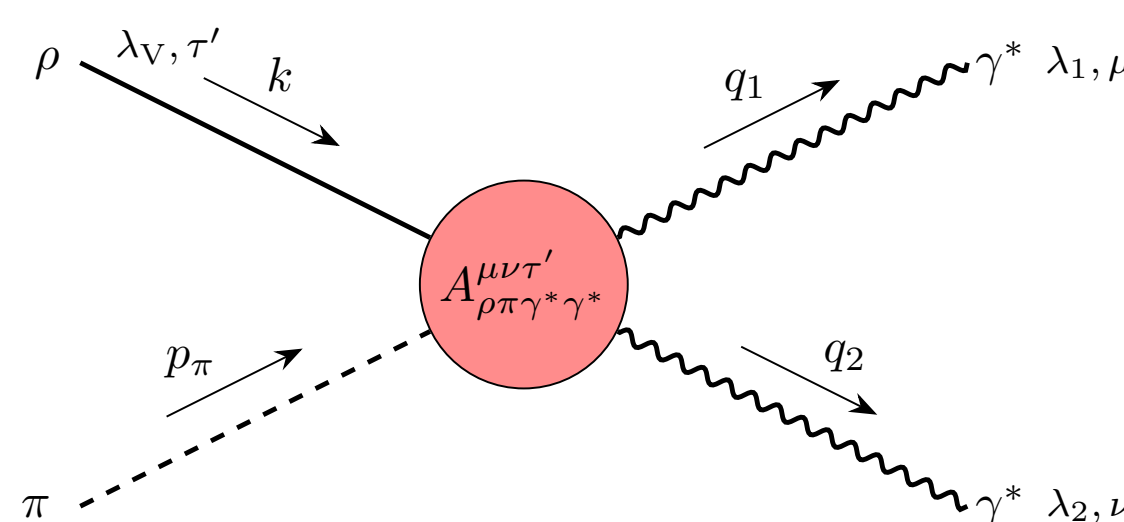


- framework based on ChPT including field representations for the pseudoscalar (P), vector (V_μ) and tensor ($T_{\mu\nu}$) multiplets [1] (containing only the relevant d.o.f. considering $a_2^0 \rightarrow \rho^\pm \pi^\mp$)

$$P = \sqrt{2} \begin{pmatrix} 0 & \pi^+ & 0 \\ \pi^- & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad V_\mu = \sqrt{2} \begin{pmatrix} 0 & \rho_\mu^+ & 0 \\ \rho_\mu^- & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad T_{\mu\nu} = \begin{pmatrix} a_{2\mu\nu}^0 & 0 & 0 \\ 0 & -a_{2\mu\nu}^0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

- phenomenological description of $a_2\rho\pi$ interaction vertex $A_{a_2\rho\pi}^{\mu\nu\tau'}$ with proper spin-2 tensor field representation \rightarrow tracelessness, transversality

Connect to $\rho\pi \rightarrow \gamma^*\gamma^*$ amplitude



- Lorentz decomposition of amplitude as $A_{\rho\pi\gamma^*\gamma^*}^{\mu\nu\tau'} = \sum_i \mathcal{F}_i T_i^{\mu\nu\tau'}$
- employ BTT procedure [2–4] for the initial set of structures
 - ↪ ensure gauge invariance and freedom from kinematic singularities and zeros for the Lorentz structures $T_i^{\mu\nu\tau'}$
 - ↪ form linear combinations in order to remove poles
 - ↪ use Schouten identities to find a minimal basis
- dispersion relations for left-hand cuts (*)
 - ↪ consider only one-particle ρ/π contributions
 - ↪ imaginary part of amplitude described in terms of $\{F_\pi^V, F_{\rho\pi}, G_1, G_2, G_3\}$
- derived a minimal basis of structures $\{\tilde{T}_i^{\mu\nu\tau'}\}_{i=1}^{13}$ and the corresponding form factors $\tilde{\mathcal{F}}_i$

References

- [1] S. Jafarzade et al. “Phenomenology of $J^{PC} = 3^{--}$ tensor mesons”. In: *Physical Review D* (2021).
- [2] W. A. Bardeen et al. “Invariant amplitudes for photon processes”. In: *Physical Review* (1968).
- [3] R. Tarrach. “Invariant amplitudes for virtual Compton scattering off polarized nucleons free from kinematical singularities, zeros and constraints”. In: *Nuovo Cimento A* (1975).
- [4] G. Colangelo et al. “Dispersion relation for hadronic light-by-light scattering: theoretical foundations”. In: *JHEP* (2015). arXiv: 1506.01386 [hep-ph].
- [5] M. Hoferichter et al. “Asymptotic behavior of meson transition form factors”. In: *JHEP* (2020). arXiv: 2004.06127 [hep-ph].

FORM FACTOR DERIVATION

Combine everything and calculate the amplitude

$$i\mathcal{M}_{a_2\gamma^*\gamma^*}^{\alpha\beta\mu\nu} = \sum_{i=1}^{13} \int \frac{d^4k}{(2\pi)^4} A_{a_2\rho\pi}^s{}_{\alpha\beta\tau} \cdot \frac{(g_{\tau\tau'} - k_\tau k_{\tau'}/M_\rho^2)}{(k^2 - M_\rho^2)} \frac{1}{[(k - p_T)^2 - M_\pi^2]} \cdot \tilde{T}_i^{\mu\nu\tau'} \tilde{\mathcal{F}}_i$$

- solve integrals using Passarino–Veltman decomposition

Calculate helicity amplitudes

$$H_{\lambda_T; \lambda_1, \lambda_2} = \mathcal{M}_{a_2\gamma^*\gamma^*}^{\mu\nu\alpha\beta} \cdot \epsilon_{\alpha\beta}^{\lambda_T}(p_T) \epsilon_\mu^{\lambda_1*}(q_1) \epsilon_\nu^{\lambda_2*}(q_2)$$

- non-vanishing helicity amplitudes satisfy : $\lambda_T = \lambda_1 - \lambda_2$

Results expressed $\propto \{f_i\}_{i=1}^4$, with

$$\begin{aligned} f_1 &= G_1(q_1^2) F_{\rho\pi}(q_2^2) & f_2 &= G_1(q_2^2) F_{\rho\pi}(q_1^2) \\ f_3 &= F_\pi^V(q_1^2) F_{\rho\pi}(q_2^2) & f_4 &= F_\pi^V(q_2^2) F_{\rho\pi}(q_1^2) \end{aligned}$$

Use the BTT result for $\mathcal{M}_{T\gamma^*\gamma^*}^{\mu\nu\alpha\beta}$ [5] as a basis

$$\mathcal{M}_{T\gamma^*\gamma^*}^{\mu\nu\alpha\beta} = \sum_{i=1}^5 T_i^{\mu\nu\alpha\beta} \frac{1}{m_T^{n_i}} \mathcal{F}_i^T$$

and project our results onto the five \mathcal{F}_i^T

- **Invariant amplitude approach**

↪ Define projectors that isolate each TFF

$$P_{i\mu\nu\alpha\beta} \mathcal{M}_{T\gamma^*\gamma^*}^{\mu\nu\alpha\beta} = \mathcal{F}_i^T$$

↪ Contract with our amplitude to find $\mathcal{F}_i^T = \sum_{j=1}^4 c_{ij} \cdot f_j$

- **Helicity amplitude approach**

↪ Calculation of helicity amplitudes in both cases

$$\vec{H} = M_2 \cdot \vec{f}, \quad \vec{H}^T = M_1 \cdot \vec{\mathcal{F}}^T$$

↪ Inverse matrix to find the relations

$$\vec{\mathcal{F}}^T = M_1^{-1} M_2 \cdot \vec{f} \rightarrow \mathcal{F}_i^T = \sum_{j=1}^4 \tilde{c}_{ij} \cdot f_j$$

FURTHER WORK

Further adaptations and ideas for future implementation:

- Consideration of structure basis that includes heavy ρ in LHCs
- Solve integrals and express results w.r.t. Feynman integrals
- Consider rescattering effects of the $\rho\pi$ -state for the dispersive description

Final steps:

- Evaluate and finalize the transition form factor functions
- Treat potential UV-divergencies
- Study whether the form factors satisfy the required asymptotic behavior

