

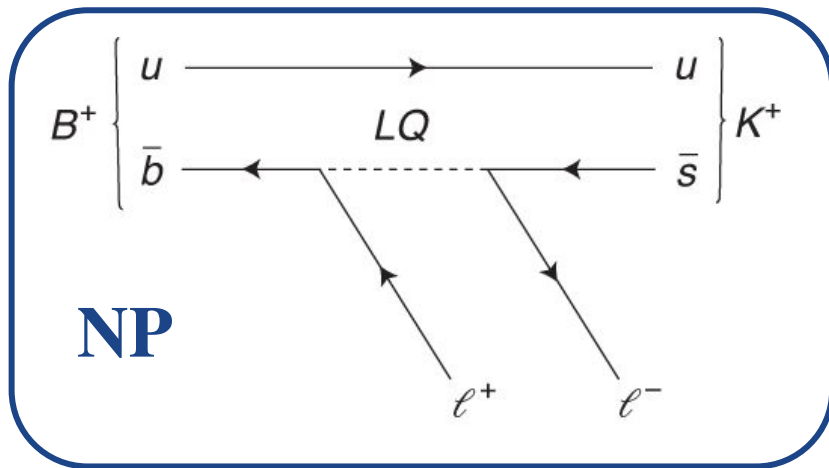
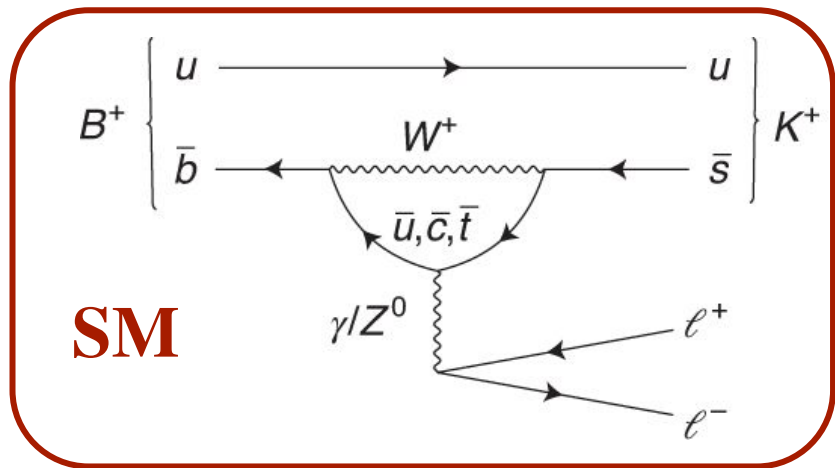
Probing lepton flavour universality with semileptonic b -decays at LHCb

Tommaso Fulghesu¹,
on behalf of LHCb collaboration

¹Aix Marseille Univ, CNRS/IN2P3, CPPM, Marseille, France

Lepton flavour universality in the Standard Model

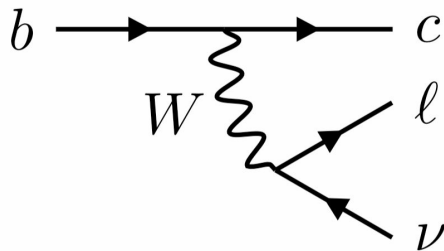
- ❖ **Standard Model (SM)** → 3 lepton flavors (e^\pm, μ^\pm, τ^\pm) with **same couplings** to gauge bosons (γ, W^\pm, Z^0), not Yukawa ($\mathcal{B}(H \rightarrow \mu^+ \mu^-) \neq \mathcal{B}(H \rightarrow \tau^+ \tau^-)$)
 - Differences in decay rates can arise from phase space or long-distance hadronic effects
- ❖ **Couplings can be affected by New Physics (NP) contributions** - leptoquarks [[JHEP 07 \(2019\) 168](#)], W' [[Phys. Rev. D 92, 054018](#)], Z' [[Phys. Rev. D 89, 015005](#)], or extended Higgs sectors [[Phys. Rev. D 86, 054014](#)].



Complementary probes in b -decays at LHCb

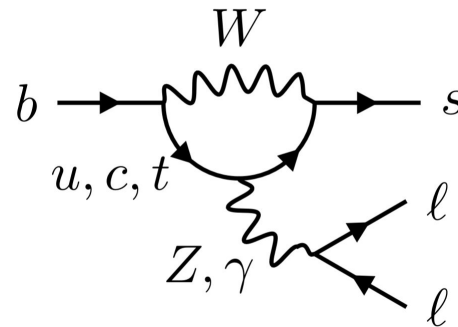
- ❖ Lepton flavour universality measurements in **tree-level** $b \rightarrow c(u)l\nu$ and **loop-level** $b \rightarrow s(d)l^+l^-$ transitions provide sensitive null-tests to New Physics

Flavour changing charged current (FCCC)



High signal yields $\mathcal{B} \sim 1-10\%$
Neutrinos not reconstructed \rightarrow
Challenging analyses

Flavour changing neutral current (FCNC)

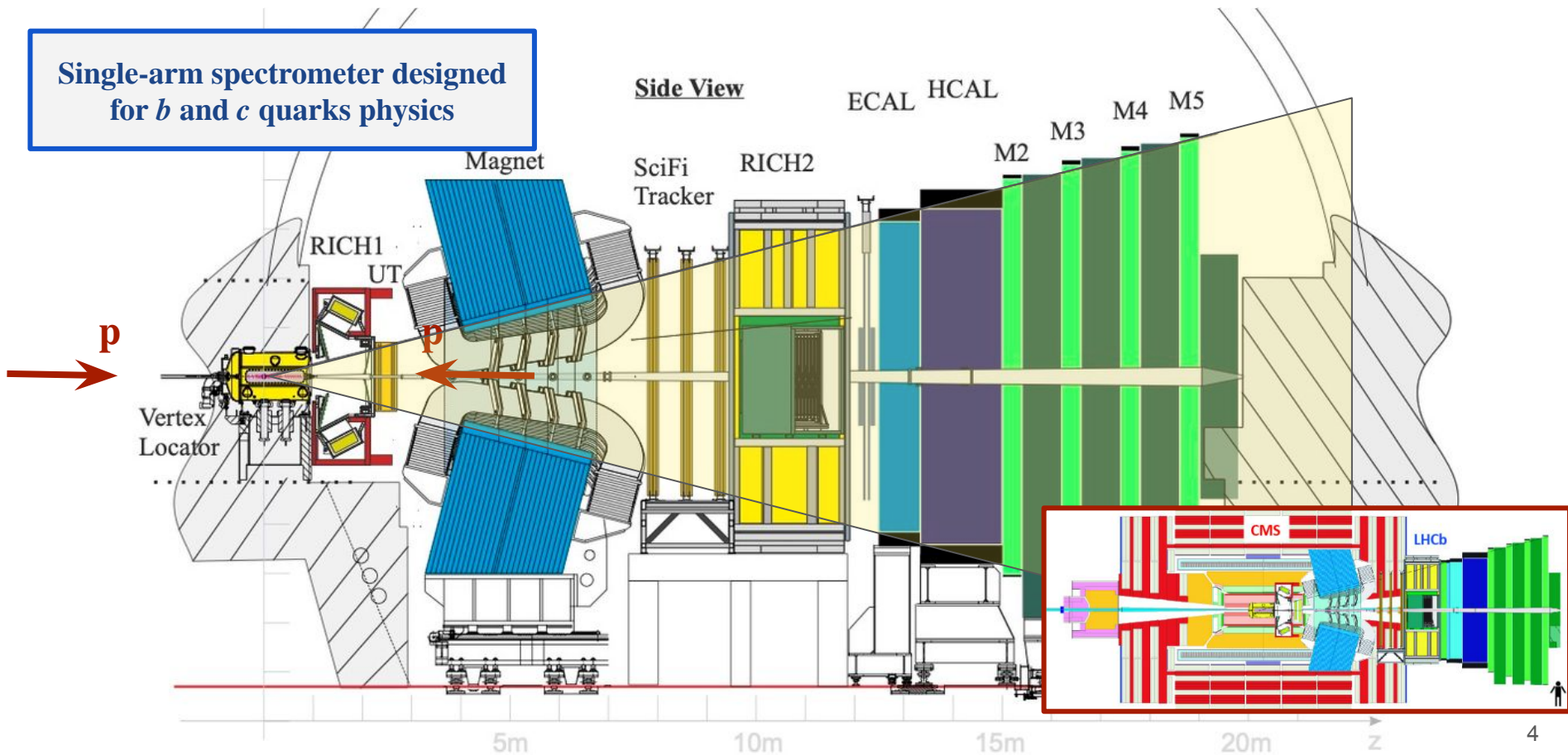


Fully-reconstructed final states
Low signal yields $\mathcal{B} \sim 10^{-6}-10^{-8}$

- ❖ LHCb has access to **all heavy-flavoured b -hadron species** including $B_c, \Lambda_b, \Xi_b, \Omega_b$.

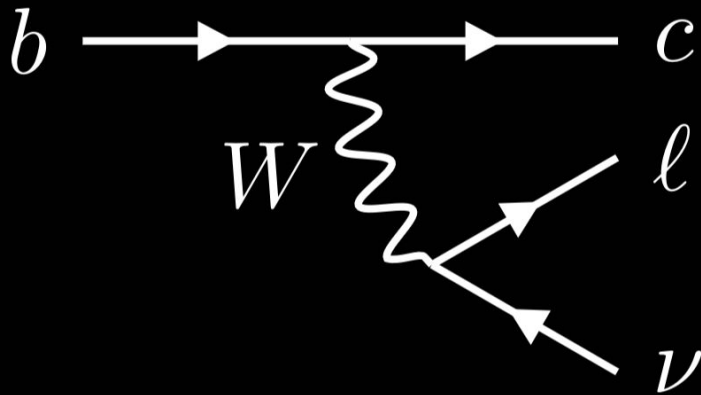
The LHCb detector

Single-arm spectrometer designed for b and c quarks physics



I. Flavour changing charged current

$$b \rightarrow c(u) \ell^- \nu_\ell$$

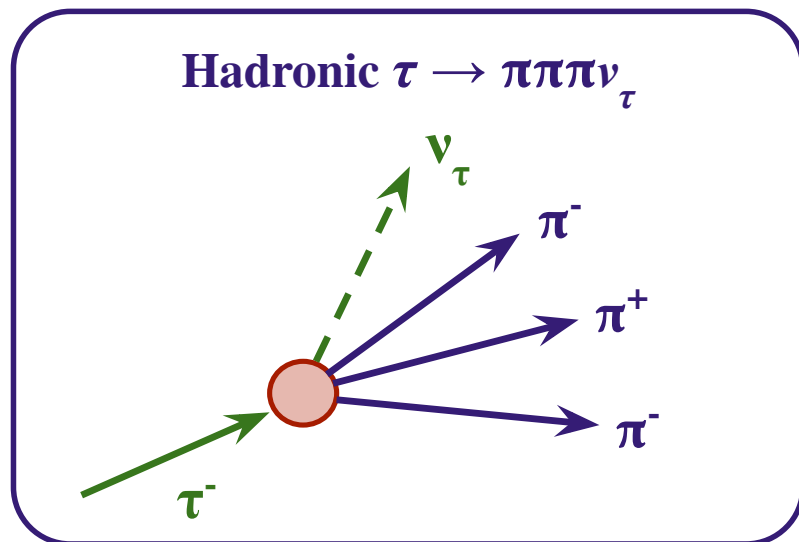
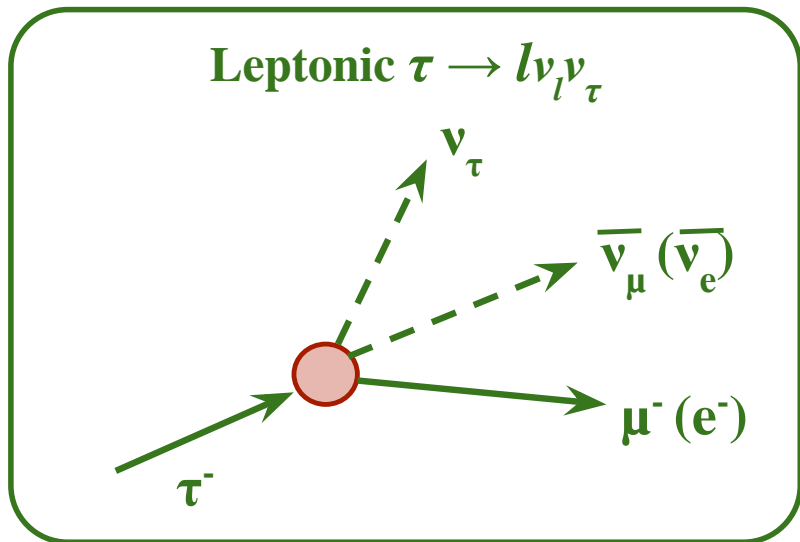


In this section:

- Missing-neutrino problem
- Tau reconstruction at LHCb
- Most recent LHCb measurements for observables $R(D)$ and $R(D^*)$
- Recent results on $R(J/\psi)$
- World average agreement and tension

Experimental strategies and challenges in $b \rightarrow c l \nu$ decays

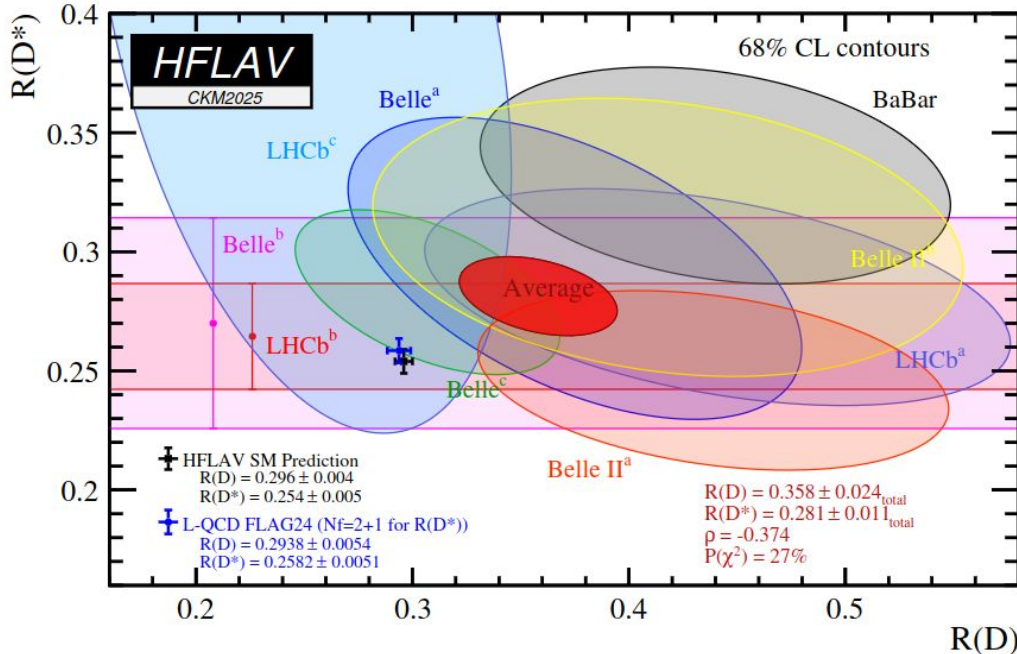
- ❖ **Missing neutrino momentum in the final states**
 - **Large partially reconstructed background** contribution
 - **b -hadron momentum** needs to be **approximated from visible particles**
 - **Large simulated samples needed** for modelling signal and background
- ❖ In $b \rightarrow c \tau \nu$, two decay topologies for tau decays are used:



R(D) and R(D*)

$$R(D^{(*)}) = \frac{\mathcal{B}(B \rightarrow D^{(*)}\tau\nu)}{\mathcal{B}(B \rightarrow D^{(*)}\ell\nu)}$$

Long standing tension of 3.8σ from different experiments!



LHCb^a 2023, muonic
 $R(D^0)$, $R(D^{*+})$, $D^{*+} \rightarrow D^0 X$
[\[PRL 131 \(2023\) 111802\]](#)

LHCb^b 2023, hadronic
 $R(D^{*+})$, $D^{*+} \rightarrow D^0 \pi^+$
[\[PRD 108 \(2023\) 012018\]](#)

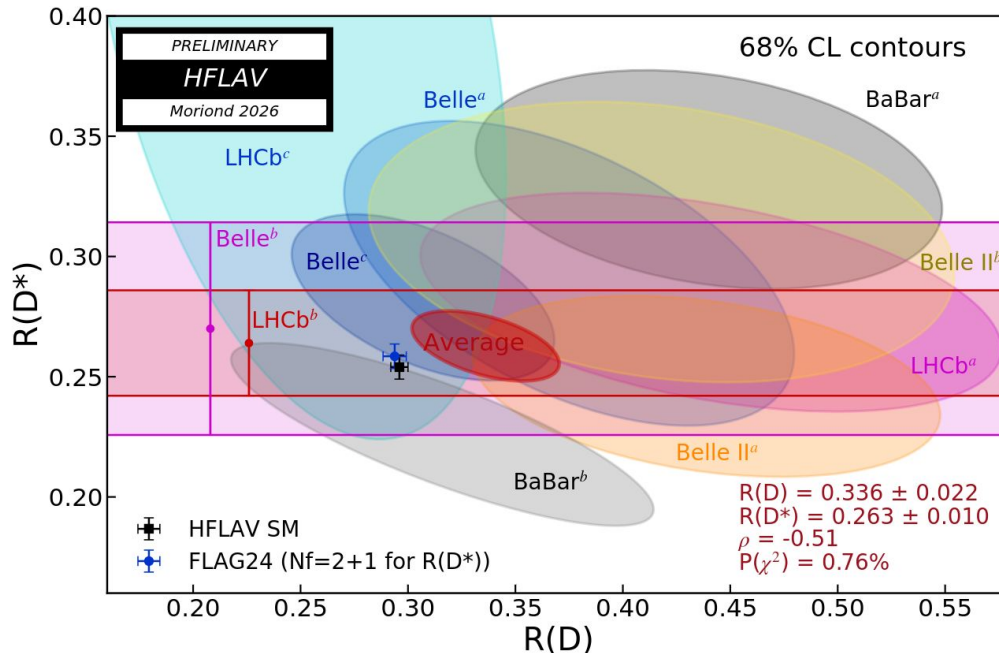
LHCb^c 2025, muonic
 $R(D^+)$, $R(D^{*+})$, $D^{*+} \rightarrow D^+ \pi^0$
[\[PRD 134 \(2025\) 061801\]](#)

[\[PRD 113 \(2026\) 012008\]](#)

R(D) and R(D*)

$$R(D^{(*)}) = \frac{\mathcal{B}(B \rightarrow D^{(*)}\tau\nu)}{\mathcal{B}(B \rightarrow D^{(*)}\ell\nu)}$$

New world average consistent with SM expectations at $\sim 2\sigma$!



LHCb^a 2023, muonic
 $R(D^0)$, $R(D^{*(+)})$, $D^{*(+) \rightarrow D^0 X}$
[\[PRL 131 \(2023\) 111802\]](#)

LHCb^b 2023, hadronic
 $R(D^{*+})$, $D^{*+ \rightarrow D^0 \pi^+}$
[\[PRD 108 \(2023\) 012018\]](#)

LHCb^c 2025, muonic
 $R(D^+)$, $R(D^{*+})$, $D^{*+ \rightarrow D^+ \pi^0}$
[\[PRD 134 \(2025\) 061801\]](#)

[Presented at [Moriond QCD 2026](#)]

R(J/ψ)

$$R(J/\psi) = \frac{\mathcal{B}(B_c^+ \rightarrow J/\psi \tau^+ \nu_\tau)}{\mathcal{B}(B_c^+ \rightarrow J/\psi \mu^+ \nu_\mu)}$$

LHCb 2018 (Run 1 data)

[[PRL 120, 121801 \(2018\)](#)]

LHCb 2026 (Run 2 data)

[LHCb-PAPER-2026-018

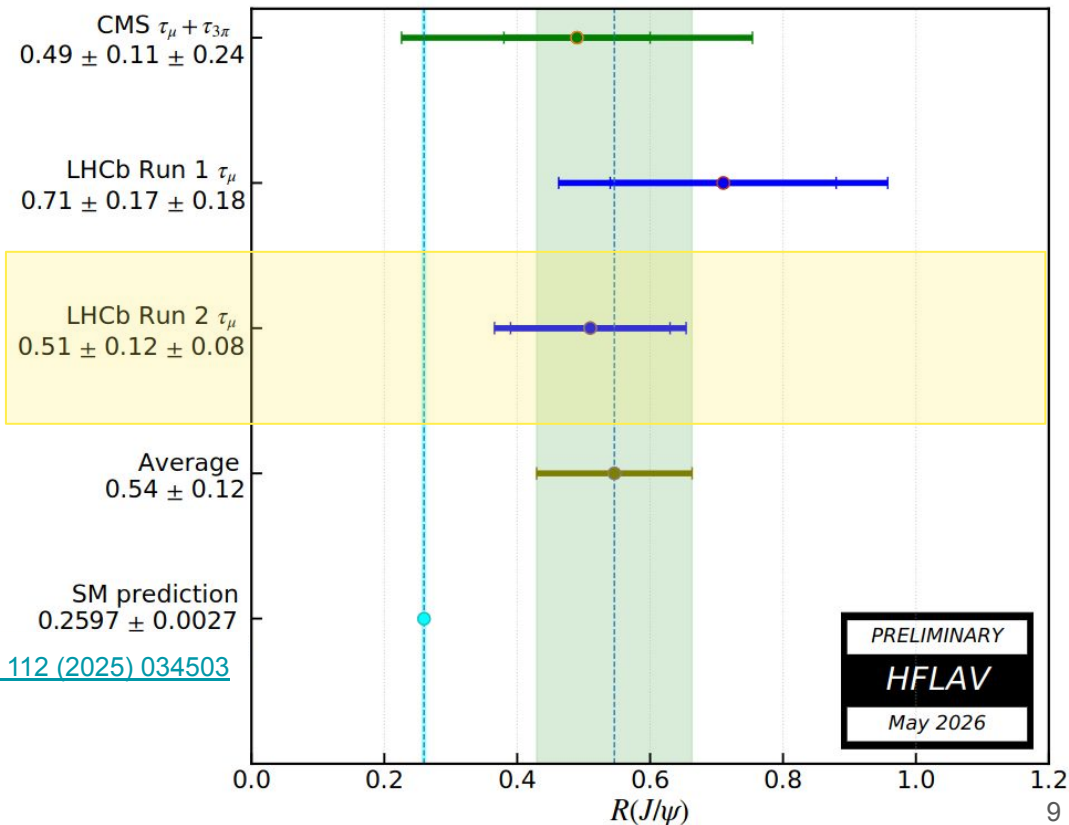
(in preparation)]

NEW

**Compatible with SM at 1.8σ
and consistent with LHCb
Run 1 measurement!**

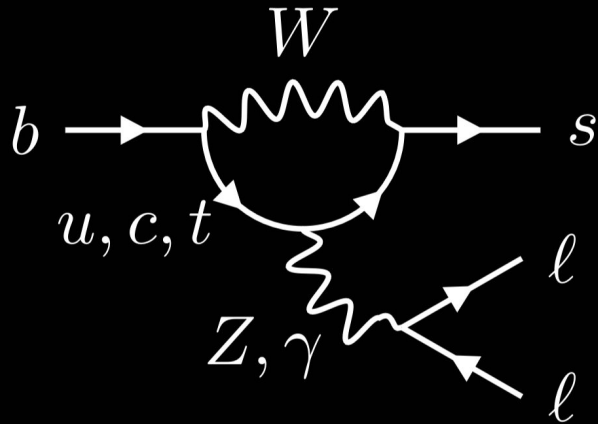
[[PRD 112 \(2025\) 034503](#)]

[[PRD 113 \(2026\) 012008](#)]



II. Flavour changing neutral current

$$b \rightarrow s(d) l^- l^+$$



In this section:

- Electrons vs muons at LHCb
- Sensitivity to New Physics effects for $b \rightarrow sll$ decays
- Invariant di-lepton mass squared q^2 spectrum
- R_X observables

Sensitivity to New Physics in $b \rightarrow sl^+l^-$ decays

❖ Effective field theory

$$H_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i (C_i O_i + C'_i O'_i)$$

Wilson coefficients

Operators

❖ New Physics enter through shift of the Wilson coefficients: $C_i^{(\prime)} = C_i^{(\prime)\text{SM}} + C_i^{(\prime)\text{NP}}$

Operators in $b \rightarrow sl^+l^-$ transitions

$$O_7^{(\prime)} \propto (\bar{s} \sigma_{\mu\nu} P_{R(L)} b) F^{\mu\nu}$$
$$O_9^{(\prime)} \propto (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{l} \gamma_\mu l)$$
$$O_{10}^{(\prime)} \propto (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{l} \gamma_\mu \gamma_5 l)$$

Sensitivity to New Physics in $b \rightarrow sl^+l^-$ decays

❖ Effective field theory

$$H_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i (C_i O_i + C_i' O_i')$$

Wilson coefficients

Operators

Operators in $b \rightarrow sl\ell$ transitions

$$O_7^{(\ell)} \propto (\bar{s} \sigma_{\mu\nu} P_{R(L)} b) F^{\mu\nu}$$

$$O_9^{(\ell)} \propto (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{l} \gamma_\mu l)$$

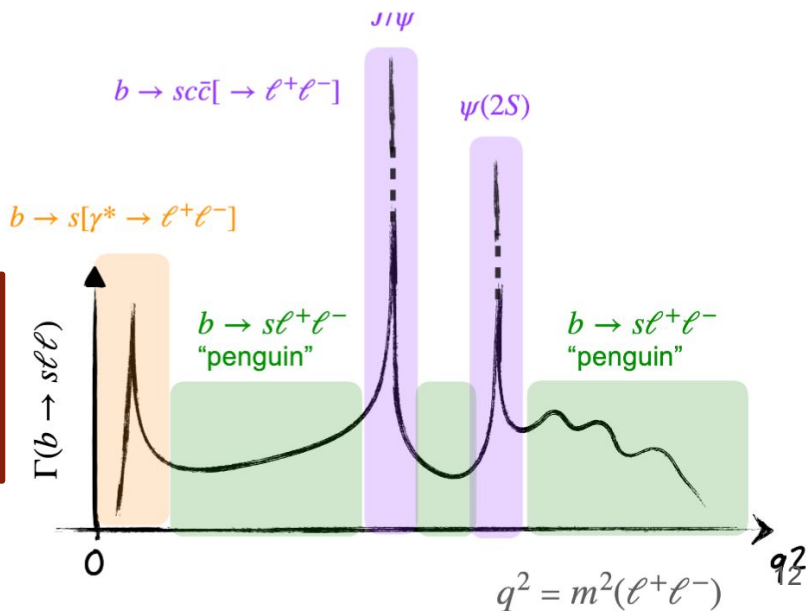
$$O_{10}^{(\ell)} \propto (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{l} \gamma_\mu \gamma_5 l)$$

❖ New Physics enter through shift of the Wilson coefficients: $C_i^{(\ell)} = C_i^{(\ell)\text{SM}} + C_i^{(\ell)\text{NP}}$

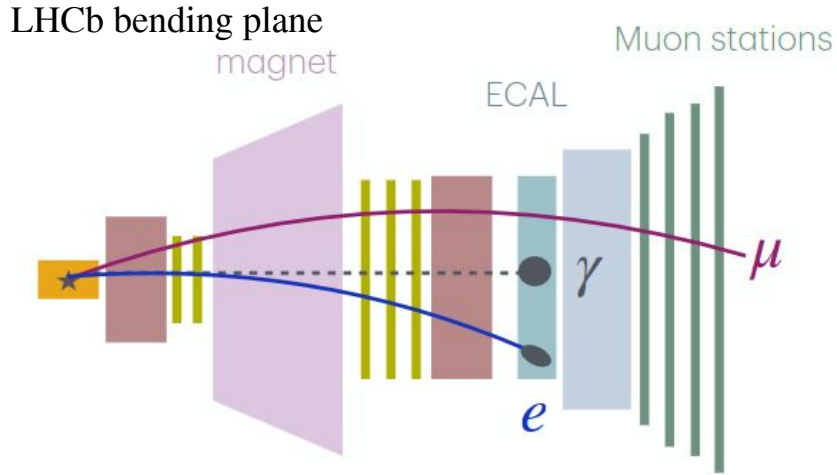
❖ Measurement of double ratio using $c\bar{c}$ resonances as control mode

$$R_{K,K^*}(q_a^2, q_b^2) = \frac{\int_{q_a^2}^{q_b^2} \frac{d\Gamma(B^{(+,0)} \rightarrow K^{(+,*)} \mu^+ \mu^-)}{dq^2} dq^2}{\int_{q_a^2}^{q_b^2} \frac{d\Gamma(B^{(+,0)} \rightarrow K^{(+,*)} e^+ e^-)}{dq^2} dq^2} \times \frac{\Gamma(J/\psi \rightarrow e^+ e^-)}{\Gamma(J/\psi \rightarrow \mu^+ \mu^-)}$$

Measured to be 1



Electrons and muons at LHCb

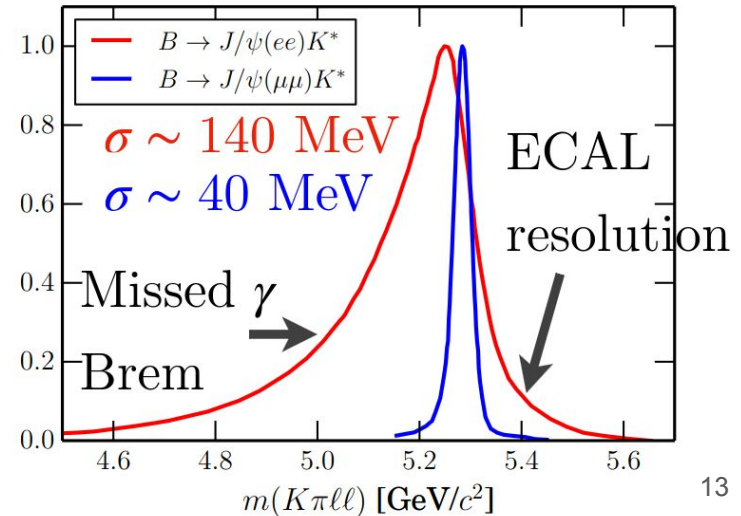


Clear signature in muon stations

Electrons have lower trigger efficiency than muons
Bremsstrahlung emission → before magnet,
recovery procedure $O(50\%)$ efficient (limited by
ECAL resolution)

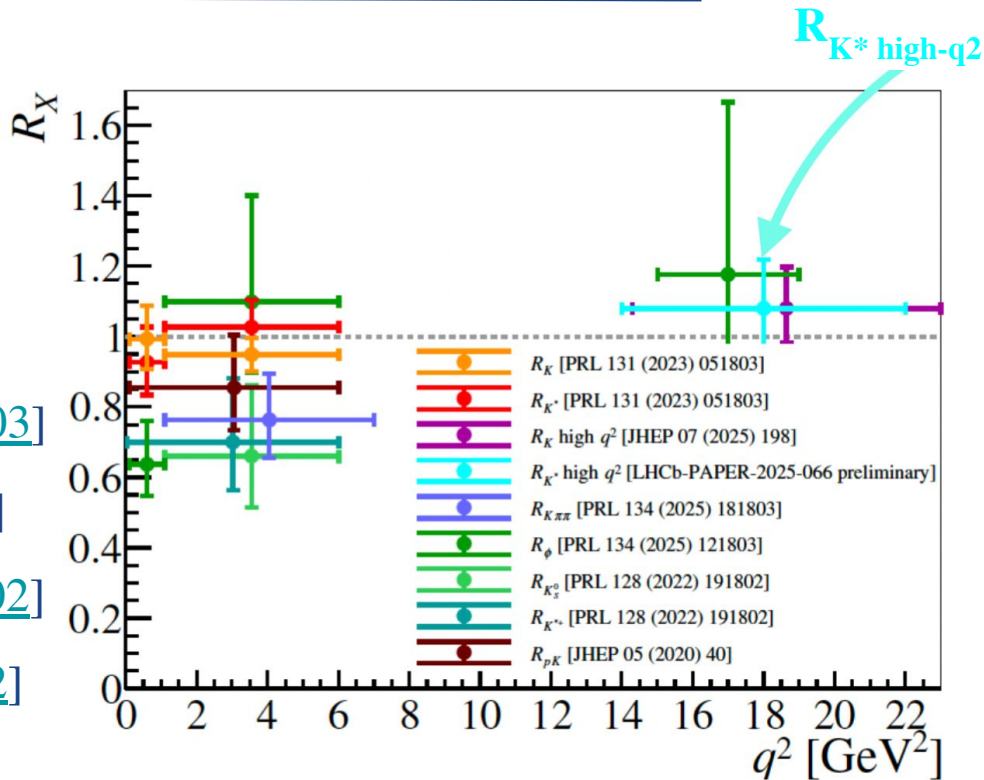
Wider fit range for **electrons** than muons

- more background
- more sensitive to peaking structures
- lineshapes are brem-dependent



Lepton flavour universality tests in $b \rightarrow sl^+l^-$ decays

- R_K [[Phys. Rev. D 108 \(2023\) 032002](#)]
- R_{K^*} [[Phys. Rev. D 108 \(2023\) 032002](#)]
- $R_{K \text{ high-}q^2}$ [[JHEP 07 \(2025\) 198](#)]
- $R_{K^* \text{ high-}q^2}$ [[arxiv:2604.08631](#)] **NEW**
- $R_{K_{\text{int}}}$ [[Phys. Rev. Lett 134 \(2025\) 181803](#)]
- R_ϕ [[Phys. Rev. Lett 134 \(2025\) 121803](#)]
- R_{K^*+} [[Phys. Rev. Lett 128 \(2022\) 191802](#)]
- R_{K_S} [[Phys. Rev. Lett 128 \(2022\) 191802](#)]
- R_{pK} [[JHEP 05 \(2020\) 040](#)]



❖ Overall picture is in good agreement with the SM predictions in $b \rightarrow sl\ell$ LFU tests

Conclusions

Tommaso Fulghesu¹,
on behalf of LHCb collaboration

¹Aix Marseille Univ, CNRS/IN2P3, CPPM, Marseille, France

Where we currently stand

- ❖ Lepton flavour universality tests in $b \rightarrow cl\nu$ and $b \rightarrow sll$ transitions $\rightarrow R(D)$ vs $R(D^*)$ and $R(J/\psi)$ in FCCC, R_X tests in FCNC
 - ongoing exploration $b \rightarrow ul\nu$ transitions and $b \rightarrow cl\nu$ electronic tau decays for FCCC, and $b \rightarrow dll$ decays [[arXiv:2604.26784](https://arxiv.org/abs/2604.26784)] and $b \rightarrow s\tau\tau$ decays [[Phys. Rev. Lett. 136 \(2026\) 181802](https://arxiv.org/abs/2202.05412)] for FCNC

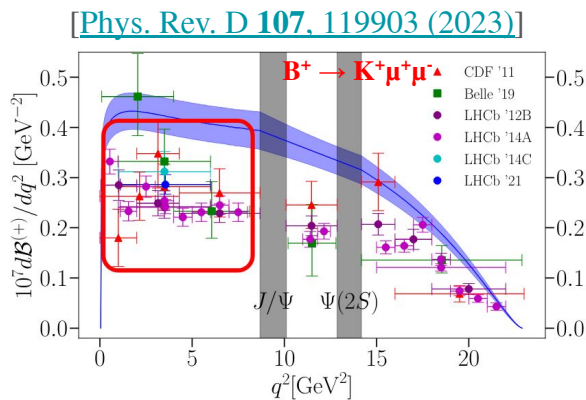
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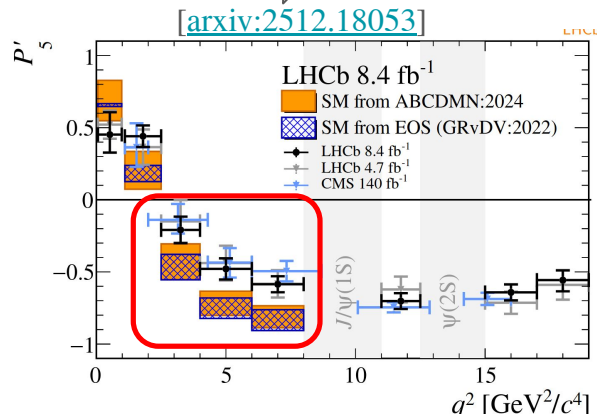
- ❖ Complementary measurements in other observables, see [D. Guadagnoli's talk](#)

INCREASING SM PRECISION

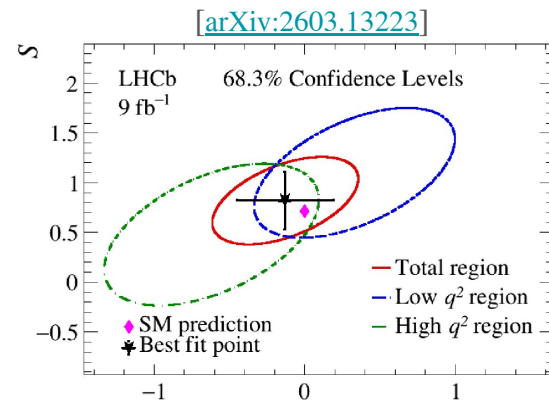
Deviations from SM predictions in $b \rightarrow sll$ FCNC



BRANCHING RATIOS



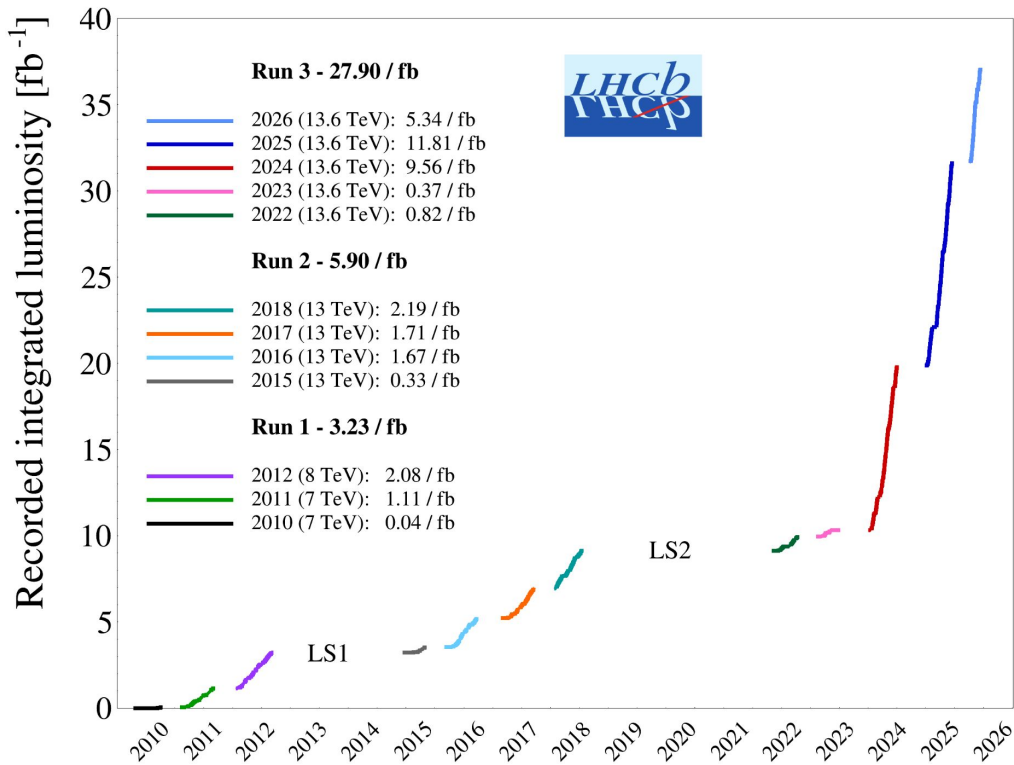
ANGULAR OBSERVABLES



CP-ASYMMETRY

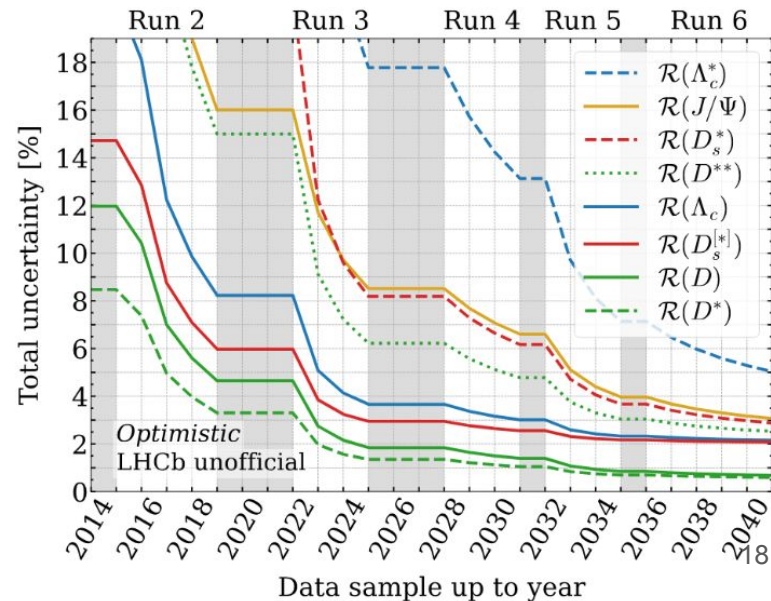
Outlook and future prospects

Total recorded luminosity – pp – 37.0 fb^{-1}



Huge statistics collected during the Run 3 by LHCb will lead to an unprecedented precision on lepton flavour universality measurements!

[[Rev. Mod. Phys \(2022\) 94, 015003](#)]

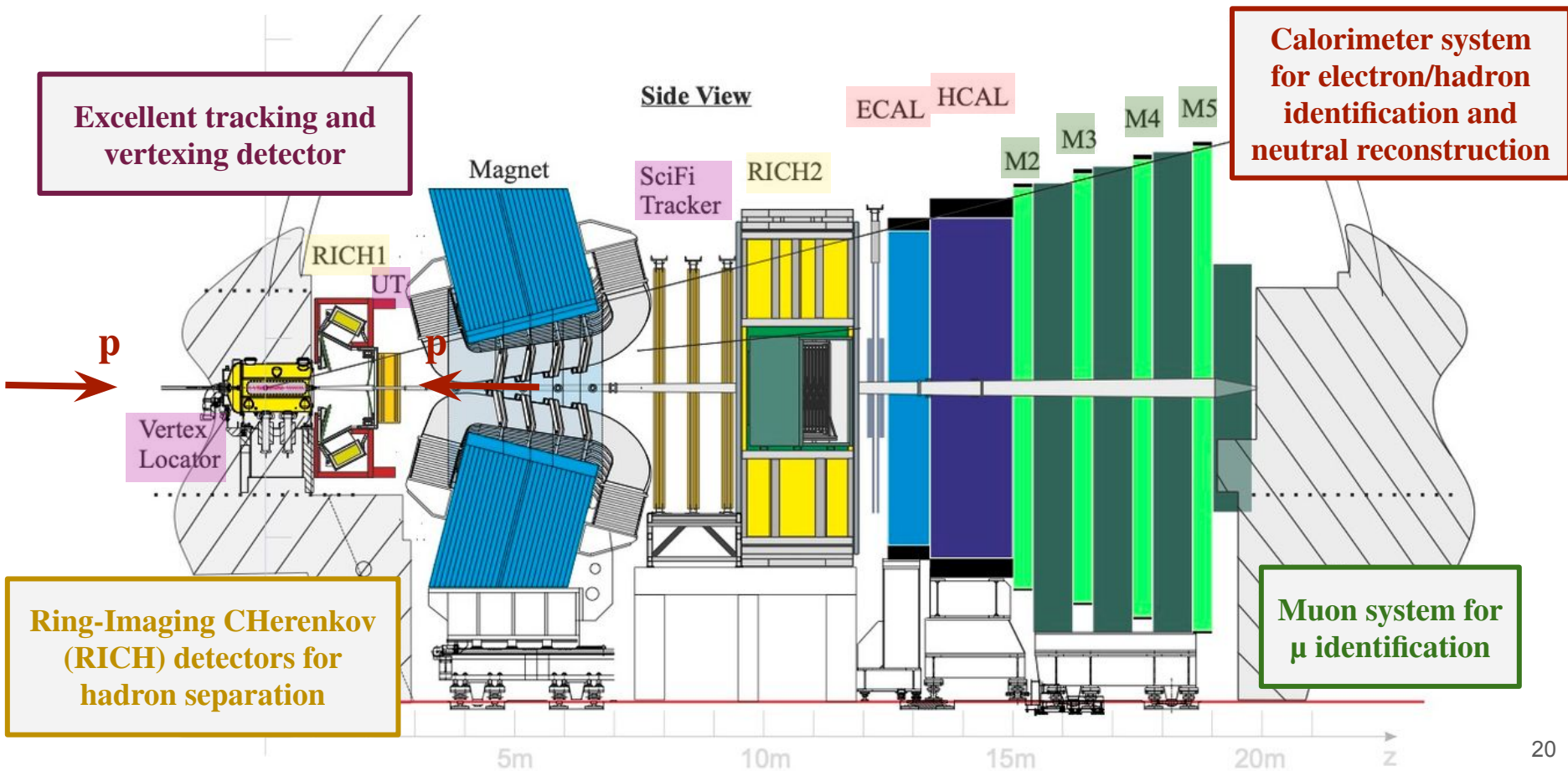


Supplementary material

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on behalf of LHCb collaboration

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The LHCb detector

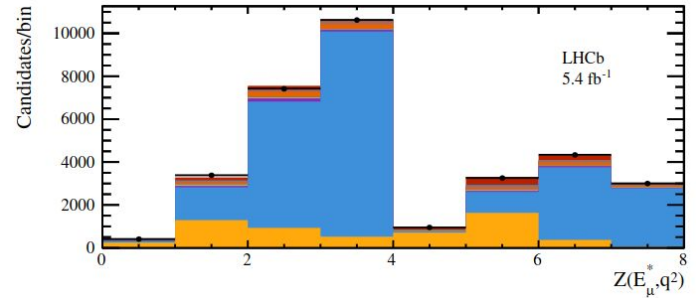
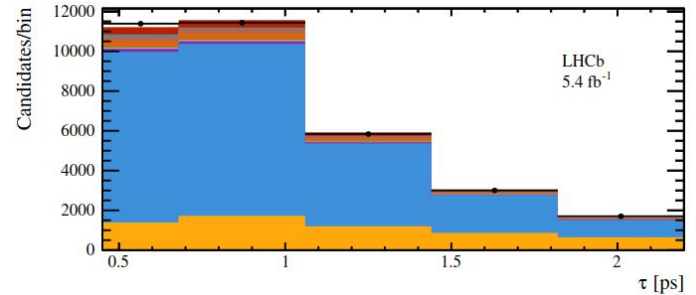
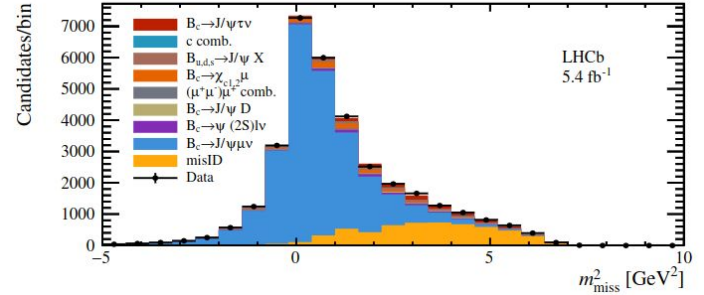
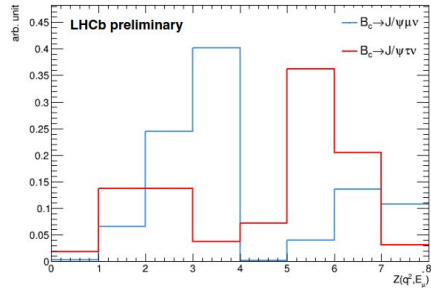
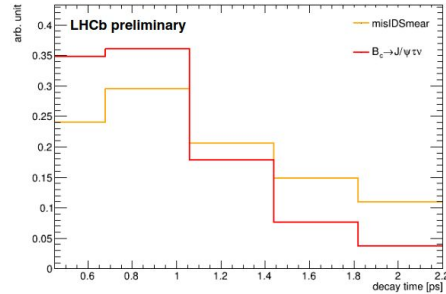
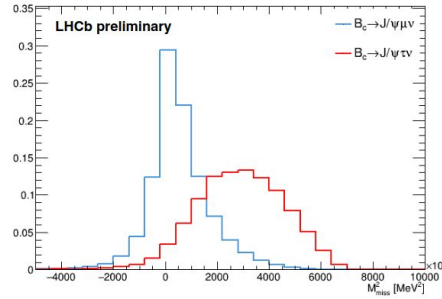
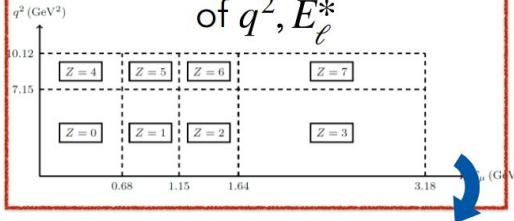


R(J/ψ): strategy

Missing mass to distinguish signal and normalization, 2 extra neutrinos mean larger missing mass

Decay time distinguishes largest background, misID
—light B s have 3x longer lifetime than B_c

$Z(q^2, E_\ell^*)$: binned function of q^2, E_ℓ^*



R(J/ψ): systematics

More reliable checks
on misID modeling,
also significantly
reduced relative to
Run-1

Tighter selection
overall—lose a bit of
statistics but made up
for in systematics!

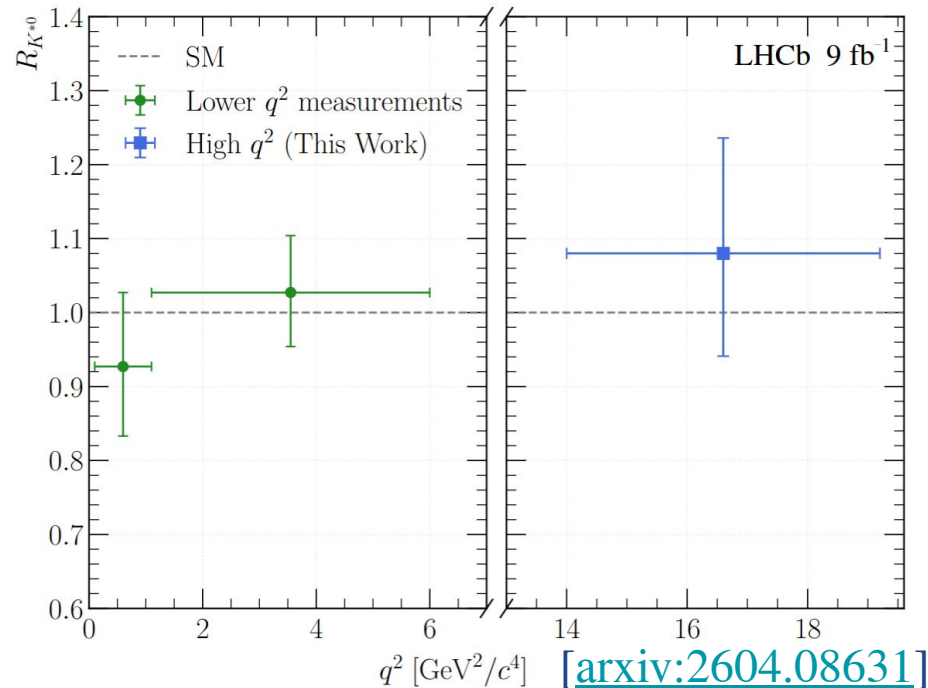
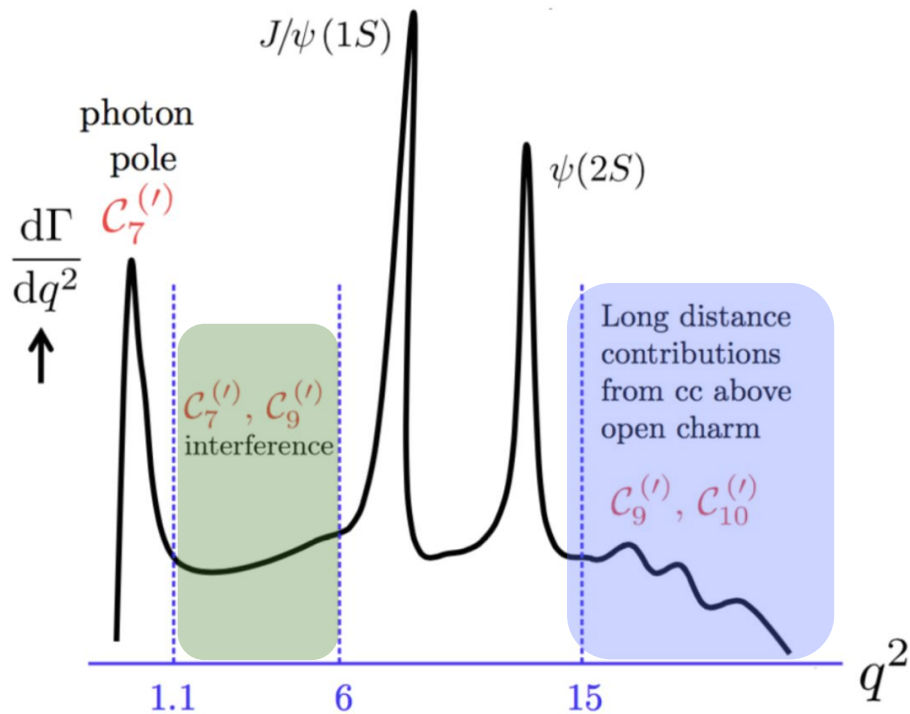
Source of uncertainty on $\mathcal{R}(J/\psi)$	Value ($\times 10^{-2}$)
Simulation statistical uncertainty	5.9
$B_c^+ \rightarrow \psi(2S)\ell^+\nu$ modeling	2.6
MisID composition estimation	2.5
Simulation corrections	2.0
$B_c^+ \rightarrow J/\psi\ell^+\nu$ form factors	1.8
$B_c^+ \rightarrow \chi_c\mu^+\nu$ scaling	1.5
Efficiency ratio	1.3
Fragmentation background modeling	1.3
MisID decay-in-flight correction	1.1
$(J/\psi + \mu^+)$ combinatorial modeling	0.6
MisID decay-in-flight smearing	0.3
$B_c^+ \rightarrow \psi(2S)\ell^+\nu$ scaling	0.3
$\mathcal{B}(\tau^+ \rightarrow \mu^+\nu_\mu\bar{\nu}_\tau)$	0.2
Trimuon effect	0.1
Systematic uncertainty	7.9
Statistical uncertainty	11.9
Total uncertainty	14.3

Lattice QCD
calculations of form
factors enable huge
reduction of largest
systematic
uncertainty in Run-1
measurement

$$\sigma_{\text{stat,Run-2}} \approx 0.65\sigma_{\text{stat,Run-1}}$$

$$\sigma_{\text{sys,Run-2}} \approx 0.5\sigma_{\text{sys,Run-1}}$$

$R(K^*)$ at high q^2



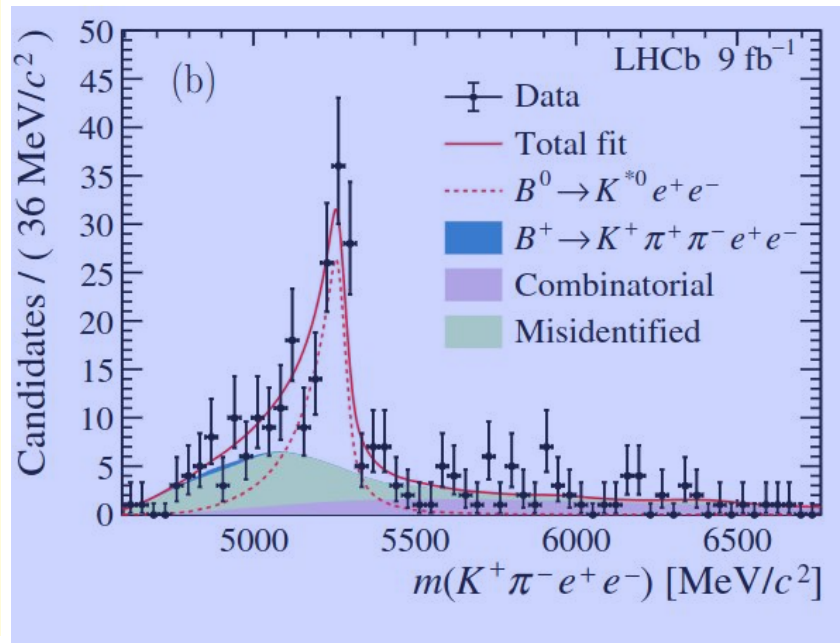
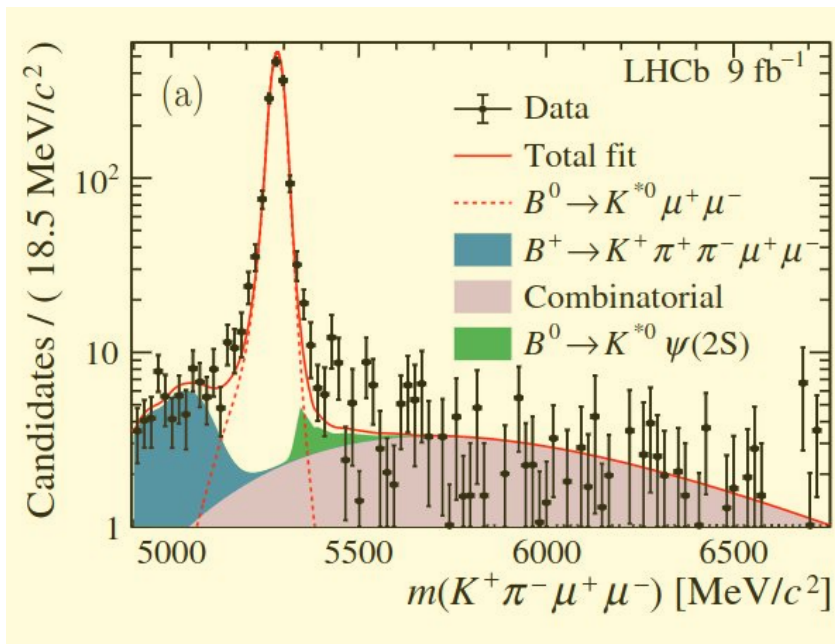
$$R_{K^{*0}} = 1.08^{+0.14}_{-0.12} \text{ (stat.)} \pm 0.07 \text{ (syst.)}$$

in agreement with SM predictions

R(K*) at high q²: strategy

from simulation

$$R_{K^{*0}} = \frac{N/\epsilon(B^0 \rightarrow K^{*0} \mu^+ \mu^-)}{N/\epsilon(B^0 \rightarrow K^{*0} J/\psi (\rightarrow \mu^+ \mu^-))} \bigg/ \frac{N/\epsilon(B^0 \rightarrow K^{*0} e^+ e^-)}{N/\epsilon(B^0 \rightarrow K^{*0} J/\psi (\rightarrow e^+ e^-))}$$



R(K*) at high q^2 : systematics

Source	Uncertainty [%]
Efficiency correction	1.2
Kinematic and topological dependence	0.7
Model dependence	2.0
q^2 smearing	0.9
Integrated luminosity	0.9
Misidentified background	5.2
Signal lineshape	1.7
Combinatorial background	1.3
Partially reconstructed background	2.0

→ **Misidentified $B \rightarrow K^* h^+ h^-$ contributions**, relevant only for electron

mode