

Bayesian calibration, uncertainty quantification, and extrapolation of nuclear reaction models

Kyle Beyer

Collaborators: Filomena M Nunes

Seminar, Orsay 2025



MICHIGAN STATE
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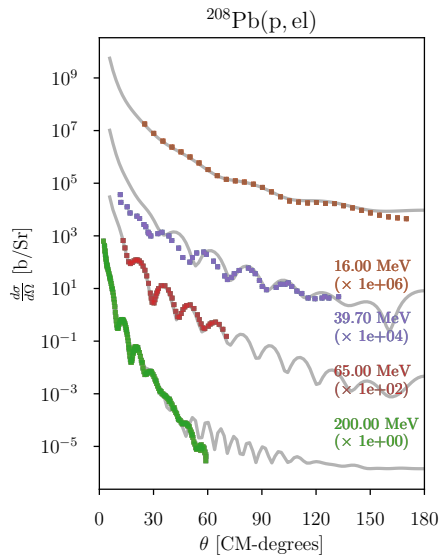
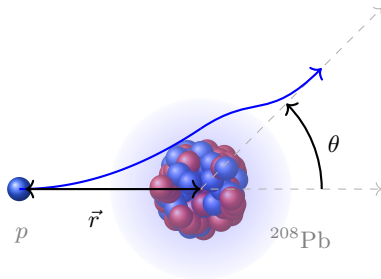
Few Body
Reactions Group

- 1 Introduction
- 2 Calibration (and the software tools enabling it)
- 3 Results
- 4 Extrapolation and uncertainty quantification

Introduction

Nuclear reactions and the optical potential

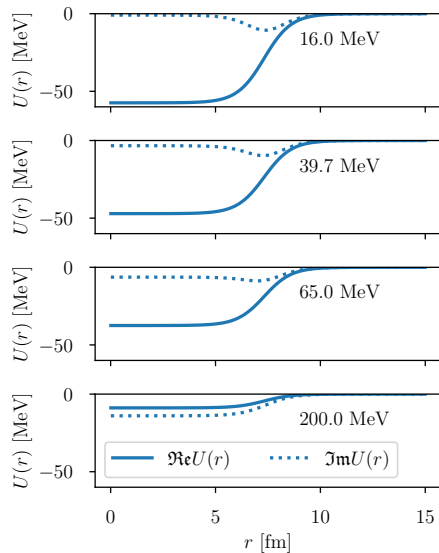
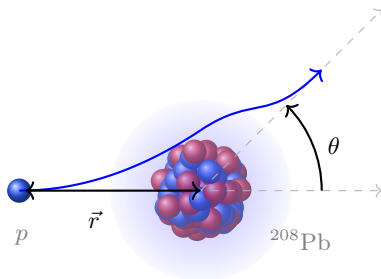
Can you reduce the many-body degrees of freedom of a nuclear reaction problem to few?



Nuclear reactions and the optical potential

Can you reduce the many-body degrees of freedom of a nuclear reaction problem to few?

Yes, with the optical potential!

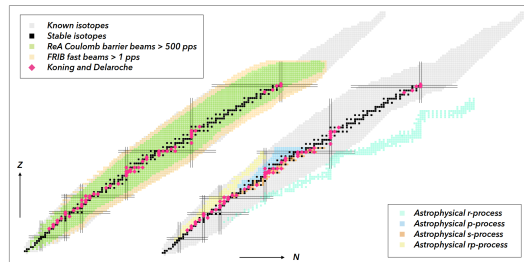
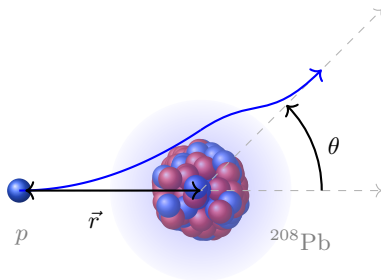


Nuclear reactions and the optical potential

Can you reduce the many-body degrees of freedom of a nuclear reaction problem to few?

Yes, with the optical potential!

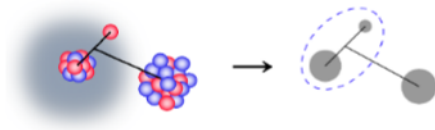
...but it requires fitting and extrapolation



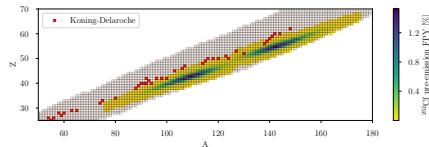
Optical potentials can be fit *locally* or *globally*, but only to stable isotopes^a

^aHebborn et al. 2023, *Journal of Physics G: Nuclear and Particle Physics* 50(6), p. 060501

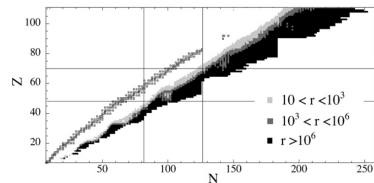
Nuclear reactions on rare isotopes involves extrapolating the optical potential



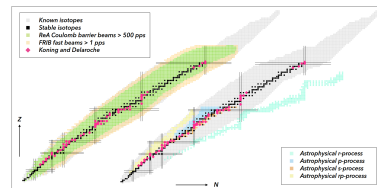
- Extrapolation of global optical potentials are ubiquitous in reaction analyses



de-excitation of fission fragments ¹



(n, γ) rates for r-process ²



astrophysical processes, rare isotope beam experiments (and more!) ³

¹Beyer et al. 2025, *Phys. Rev. C* 112, p. 024604

²Goriely and Delaroche 2007, *Physics Letters B* 653(2-4), pp. 178–183

³Hebborn et al. 2023, *Journal of Physics G: Nuclear and Particle Physics* 50(6), p. 060501

Extrapolation off the line of stability is constrained by the charge-symmetry of the nuclear force

- Reasonable ansatz for isovector dependence of optical potential ⁴:

$$\begin{aligned}
 U(r) &= U_0(r) + 4 \frac{\boldsymbol{\tau} \cdot \mathbf{T}}{A} U_1(r) \\
 &= U_0(r) + \underbrace{4 \frac{\tau_3 T_3}{A} U_1(r)}_{(N-Z)/A \text{ dependence in elastic scattering}} + \underbrace{4 \frac{\tau_+ T_- + \tau_- T_+}{2A} U_1(r)}_{(p, n) \text{ to isobaric analog}}
 \end{aligned}$$

- ~ 10 s of keV mass corrections for higher-order isospin terms suggests they are unimportant for scattering

Goal: improve extrapolation in $(N - Z)/A$ by using $(p, n)_{\text{IAS}}$ as a constraint in a global, uncertainty-quantified optical potential

⁴Lane 1962, *Nuclear Physics* 35, pp. 676–685

(p, n) to the isobaric analog state

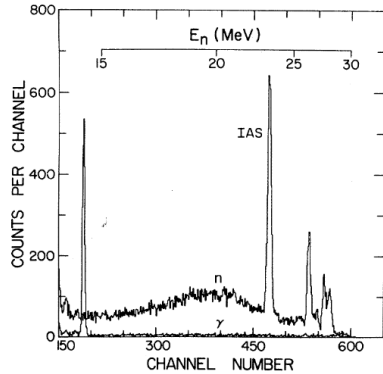
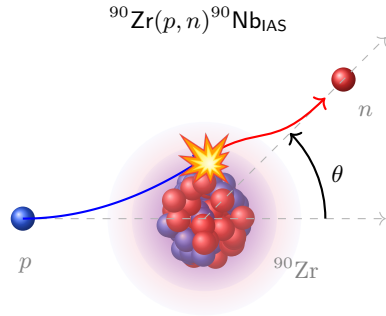


FIG. 1. Neutron (n) and γ -ray (γ) time-of-flight spectra at 0° for 35-MeV protons on ^{90}Zr . Each channel represents 0.1 nsec.

From measurements at MSU Cyclotron Lab ^a

^aDoering et al. 1975, *Physical Review C* 12(2), p. 378

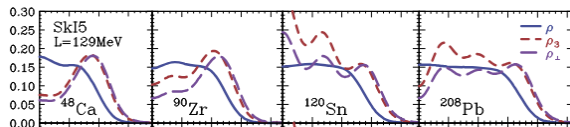


- good single-particle description: replacement of valence neutron orbital with incident proton
- probe of N, Z symmetry energy and neutron skins

The East Lansing Model

$$U(r; E, A, Z) = U_0(r; E, A) + \left(\frac{N - Z}{N + Z} \right) U_1(r; E, A)$$

- 24 parameters (compared to 22 for CHUQ, 48 for KDUQ)
- Lane-consistent, spherical, global, energy dependent
- Coulomb correction to all orders
- relationship between (p, n) , isovector/neutron skins, and symmetry energy by Danielewicz et al.⁵



visitlearn.msu.edu/resources/facility-for-rare-isotope-beams
photos.msu.edu

github.com/beykyle/elm

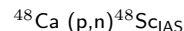
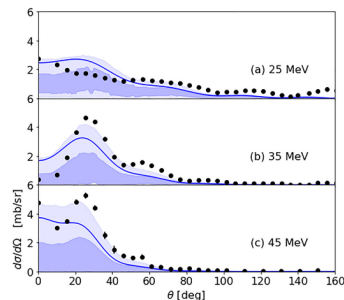
⁵Danielewicz et al. 2017, *Nuclear Physics A* 958, pp. 147–186

Calibration (and the software tools enabling it)

Progress on uncertainty quantification and extrapolation of the optical potential

Statistical approach to optical potentials

- **1950s on:** Frequentist, least-squares fitting, a few parameters at a time. Parameter ambiguities observed. ⁶
- **2003:** Koning-Delaroche ⁷ is state-of-the-art: 50 parameter fit with computational steering, local, global hierarchy
- **2017:** Bayesian calibration and UQ with local potentials (Lovell et al. ⁸)
- **2023:** First global Bayesian calibration (Pruitt et al. ⁹)



- systematic errors, model discrepancy, errors in angle, identifiability, etc., have all been ignored
- Stable-target reactions poorly constrain U_1 ; large uncertainties for (p,n) ¹⁰ and for reactions on highly asymmetric isotopes ¹¹

⁶Satchler et al. 1964, *Physical Review* 136(3B), B637

⁷Koning and Delaroche 2003, *Nuclear Physics A* 713(3-4), pp. 231–310

⁸Catacora-Rios et al. 2021, *Phys. Rev. C* 104, p. 9; Lovell et al. 2017, *Physical Review C* 95(2), p. 024611

⁹Pruitt et al. 2023, *Physical Review C* 107(1), p. 014602

¹⁰Smith et al. 2024, *Physical Review C* 110(3), p. 034602

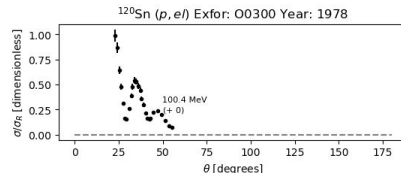
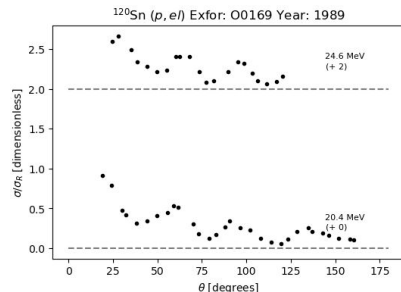
¹¹Goriely and Delaroche 2007, *Physics Letters B* 653(2-4), pp. 178–183

Curation of a body of experimental data

- ~ 100 EXFOR entries, $\sim 10\text{K}$ data points
- $40 \leq A \leq 208$
- $10 \text{ MeV} < E_{\text{lab}} < 200 \text{ MeV}$
- $d\sigma/d\Omega$ and A_y from elastic (p, p) , (n, n)
- $d\sigma/d\Omega$ with clean $\Delta J^\pi = 0^+$ from $(p, n)_{\text{IAS}}$
- manual ID & cleaning of outliers
- types of error reported inconsistent

github.com/beykyle/exfor_tools

```
all_entries_pp = exfor_tools.get_exfor_differential_data(
    target=target,
    projectile=proton,
    quantity="dXS/dA",
    product="EL",
    energy_range=[10, 200], # MeV
)
```



jit $\mathcal{R}$ allows for global calibration without needing emulators

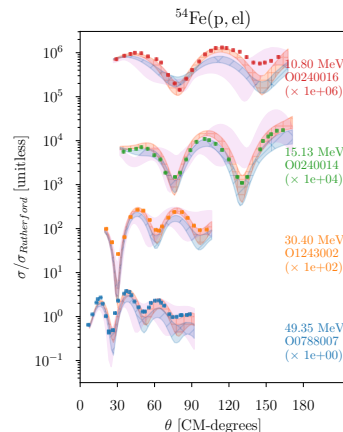
philosophy:

Modular library for building workspaces which take in interaction models and spit out reaction observables

- calculable $\mathcal{R}$ -matrix on Lagrange mesh ^a
- modular python library
- precompute everything you can
- numba just-in-time compilation, heavy numpy vectorization
- **2 orders of magnitude faster than comparable fortran codes**, just as fast as emulators ^b in UQ setting

^aBaye 2015, *Physics reports* 565, pp. 1–107

^bOdell et al. 2024, *Physical Review C* 109(4), p. 044612



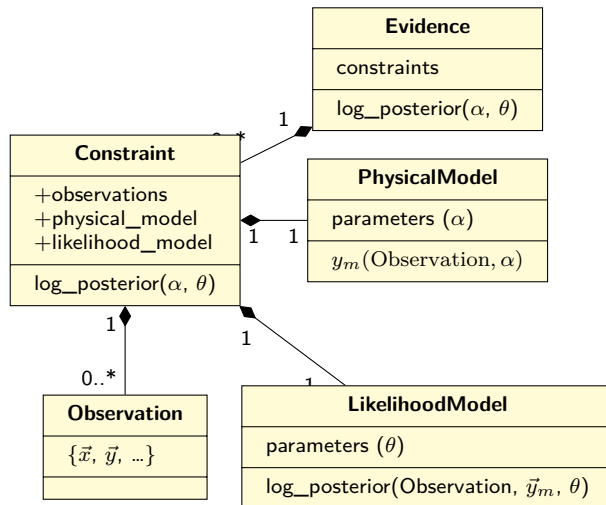
Open source, part of BAND framework ^a

^aBeyer et al. 2023, (Version 0.3.0)



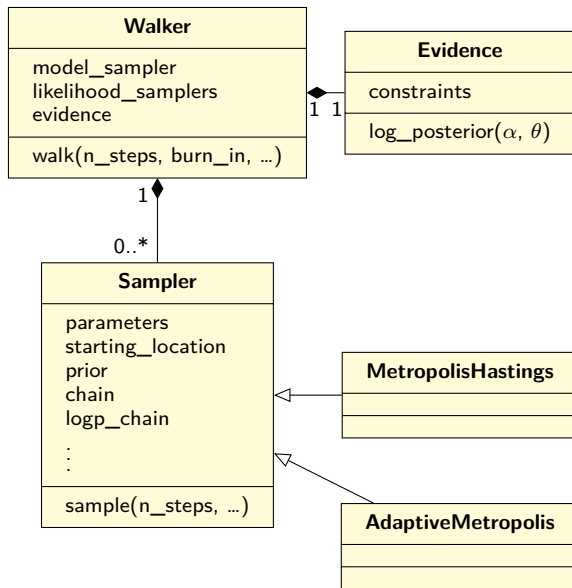
rxmc: spotlight on the body of evidence and the likelihood

- curating a body of evidence from multiple constraints
- Bayesian calibration with flexible likelihood models
- designed with `jitR` and `exfor-tools` in mind
- built-in MCMC, or interface with external MCMC package
- handles parametric likelihood models with batched Metropolis-in-Gibbs sampling
- alpha stage:
<https://github.com/beykyle/rxmc/>

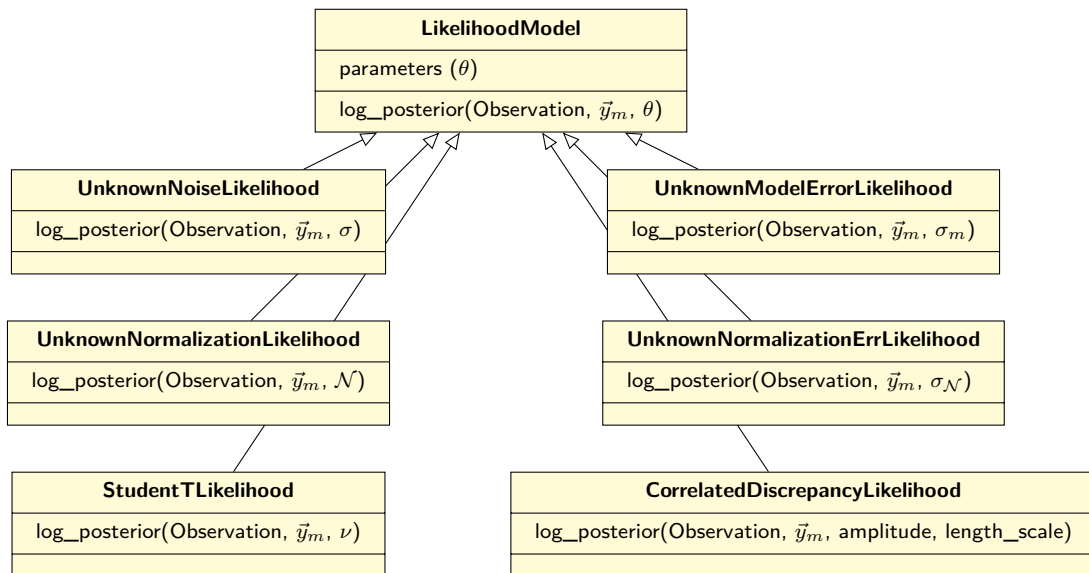


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rxmc: spotlight on the body of evidence and the likelihood



Aside: systematic error in the likelihood

Common approach is to introduce a normalization parameter $\mathcal{N}$:

$$\chi^2(\alpha) = \sum_i \frac{(y_i - \mathcal{N}y_m(x_i; \alpha))^2}{\sigma_i^2} + \frac{(\mathcal{N} - 1)^2}{\sigma_{\mathcal{N}}^2} \quad (1)$$

This is incorrect! The correct approach is actually¹²:

$$\chi^2(\alpha) = \sum_i \frac{(\mathcal{N}y_i - y_m(x_i; \alpha))^2}{(\mathcal{N}\sigma_i)^2} + \frac{(\mathcal{N} - 1)^2}{\sigma_{\mathcal{N}}^2} \quad (2)$$

One can implicitly marginalize over $\mathcal{N}$, which requires use of a non-diagonal covariance¹³. This relates to Peelle's Pertinent Puzzle in nuclear data evaluation¹⁴.

¹²Giulio D'Agostini (1994). "On the use of the covariance matrix to fit correlated data". In: *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 346.1-2, pp. 306–311; Ian J Thompson and Filomena M Nunes (2009). *Nuclear reactions for astrophysics: principles, calculation and applications of low-energy reactions*. Cambridge University Press.

¹³Roger Barlow (2021). "Combining experiments with systematic errors". In: *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 987, p. 164864.

¹⁴R Fröhwrth et al. (2012). "Peelle's pertinent puzzle and its solution". In: *EPJ Web of Conferences*. Vol. 27. EDP Sciences, p. 00008; Denise Neudecker et al. (2012). "Peelle's pertinent puzzle: A fake due to improper analysis". In: *Nuclear science and engineering* 170.1, pp. 54–60.

Handling multiple constraints with systematic errors

- statistical only:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2$$

- statistical error is fit parameter:

$$\Sigma_{ij} = \delta_{ij} (\eta y_m(x_i; \alpha))^2$$

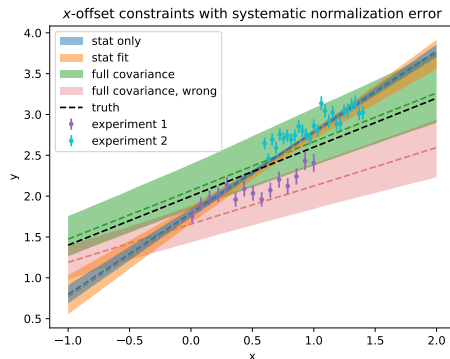
- full covariance:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2 + \sigma_N^2 y_m(x_i; \alpha) y_m(x_j; \alpha)$$

- full covariance, wrong:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2 + \sigma_N^2 y_i y_j$$

We have a high-schooler working on calibrating optical potentials accounting properly for systematic error!



https://github.com/beykyle/rxmc/blob/main/examples/systematic_err_demo.ipynb

Statistical model for ELM

$$\underbrace{p(\boldsymbol{\alpha}|\mathcal{D}, \mathcal{M})}_{\text{posterior}} \propto \underbrace{p(\mathcal{D}|\boldsymbol{\alpha}, \mathcal{M})}_{\text{likelihood}} \underbrace{p(\boldsymbol{\alpha})}_{\text{prior}}$$

Likelihood model ¹⁵

$$y_{i,\text{exp}} + \delta_{i,\text{exp}} = y_{\mathcal{M}}(x_i; \boldsymbol{\alpha}) + \delta_{MD}(x; \boldsymbol{\alpha})$$

$$\ln p(\mathcal{D}|\boldsymbol{\alpha}, \mathcal{M}) = -\frac{1}{2} \left(\boldsymbol{\Delta}^\top \cdot \boldsymbol{\Sigma} \cdot \boldsymbol{\Delta} + \ln \left((2\pi)^k |\boldsymbol{\Sigma}| \right) \right)$$

$$\Delta_i(\boldsymbol{\alpha}) = y_{\mathcal{M}}(x_i, \boldsymbol{\alpha}) - y_i$$

$$\boldsymbol{\Sigma} \equiv \boldsymbol{\Sigma}_{\text{exp}} + \boldsymbol{\Sigma}_{MD}$$

Prior model

- estimated from posterior of prior calibrations w/out (p, n) ¹⁶

Calibration

- Adaptive Metropolis-Hastings in Gibbs with `rxmc`
- Goodman-Weare from `emcee` ¹⁷
- massively parallel computation

¹⁵Kennedy and O'Hagan 2001, *Journal of the Royal Statistical Society: Series B (Statistical Methodology)* 63(3), pp. 425–464

¹⁶Pruitt et al. 2023, *Physical Review C* 107(1), p. 014602

¹⁷Foreman-Mackey et al. 2013, *PASP* 125(925), p. 306; Goodman and Weare 2010, *Communications in applied mathematics and computational science* 5(1), pp. 65–80

Likelihood model

$$\ln p(\mathcal{D}|\boldsymbol{\alpha}, \mathcal{M}) = -\frac{1}{2} \left(\boldsymbol{\Delta}^\top \cdot \boldsymbol{\Sigma} \cdot \boldsymbol{\Delta} + \ln \left((2\pi)^k |\boldsymbol{\Sigma}| \right) \right) \quad \Delta_i(\alpha) \equiv y_{\mathcal{M}}(x_i, \boldsymbol{\alpha}) - y_i$$

- learn model (uncorrelated) discrepancy from data

$$\Sigma_{\text{exp},ij}(\alpha) \equiv \delta_{ij} \sigma_{\text{exp},i}^2$$

$$\Sigma_{\text{MD},ij}(\eta) \equiv \delta_{ij} [\epsilon (y_{\mathcal{M}}(x_i; \boldsymbol{\alpha}) + y_{i,\text{exp}}) / 2]^2$$

- ϵ is a free parameter

To re-scale or not to re-scale the log-likelihood?

- Pruitt et al^a re-scaled the log likelihood by k/N (# of free parameters / # of experimental data points):

$$\ln p(\mathcal{D}|\boldsymbol{\alpha}, \mathcal{M}) \rightarrow -\left(\frac{k}{N}\right) \frac{1}{2} \left(\boldsymbol{\Delta}^\top \cdot \boldsymbol{\Sigma} \cdot \boldsymbol{\Delta} + \ln \left((2\pi)^k |\boldsymbol{\Sigma}| \right) \right)$$

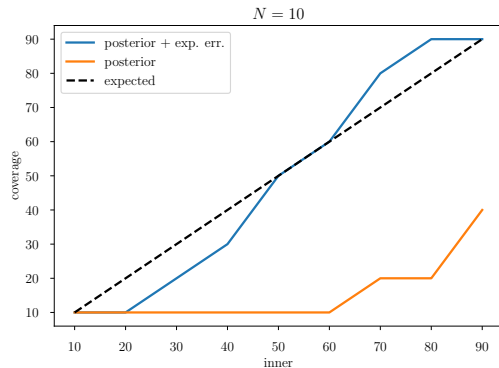
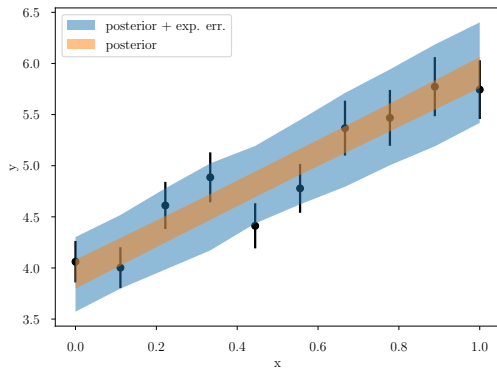
- Are we overconfident if posterior credible intervals are narrow?
- re-scaling is equivalent to “cold” posterior with $\lambda = k/N < 1$ (post-Bayesian approach)^b

$$p_\lambda(\boldsymbol{\alpha}|\mathcal{D}, \mathcal{M}) \propto p(\mathcal{D}|\boldsymbol{\alpha}, \mathcal{M})^\lambda p(\boldsymbol{\alpha})$$

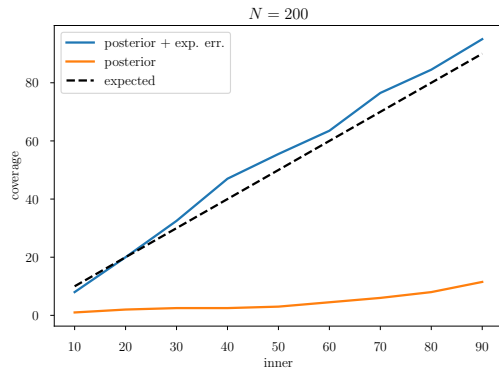
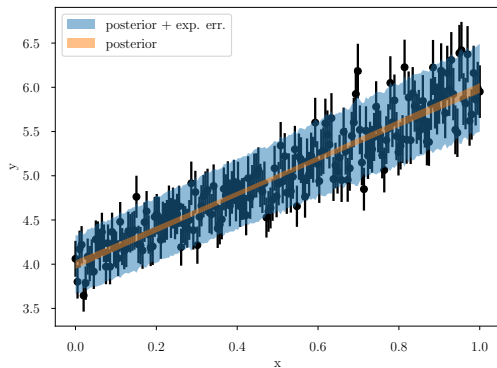
^aPruitt et al. 2023, *Physical Review C* 107(1), p. 014602

^bGrünwald 2011, vol. 19, pp. 397–420

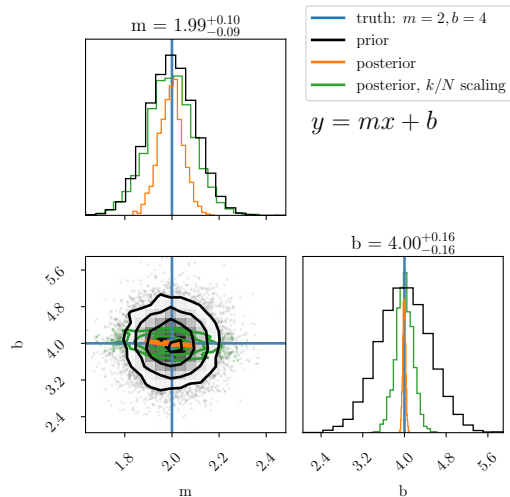
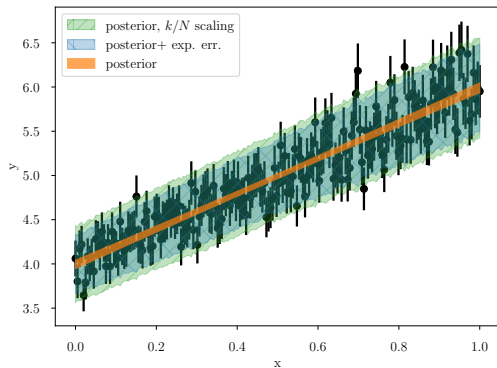
Empirical coverage



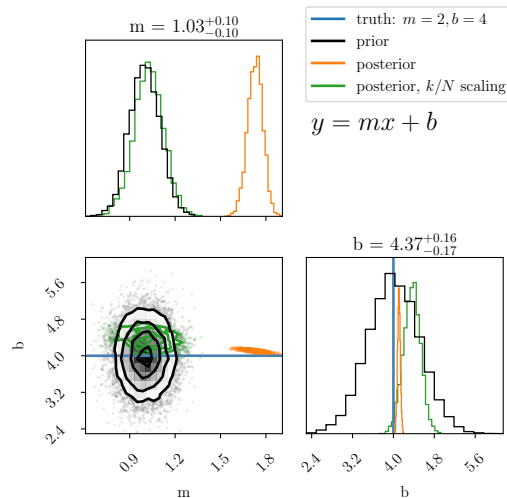
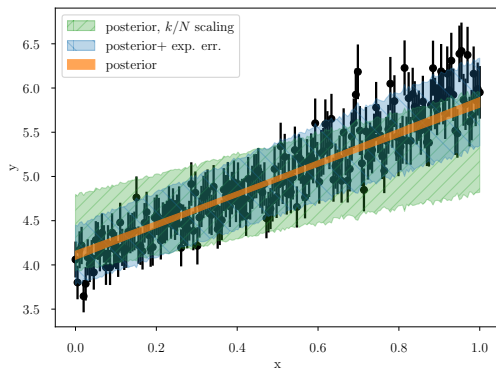
Empirical coverage



“Cold” posterior

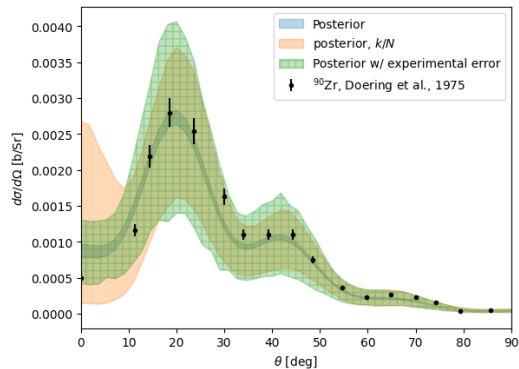
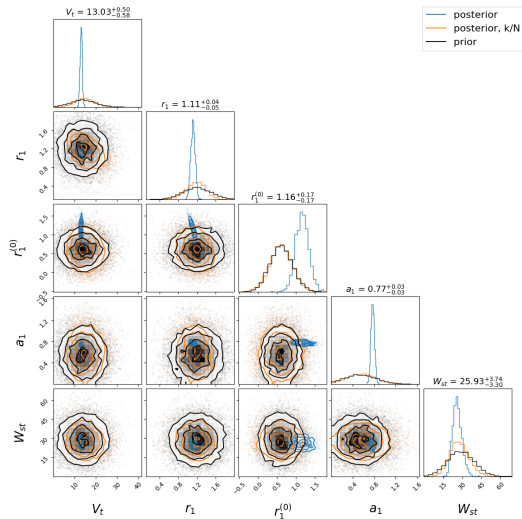


“Cold” posterior: over-emphasis on the prior

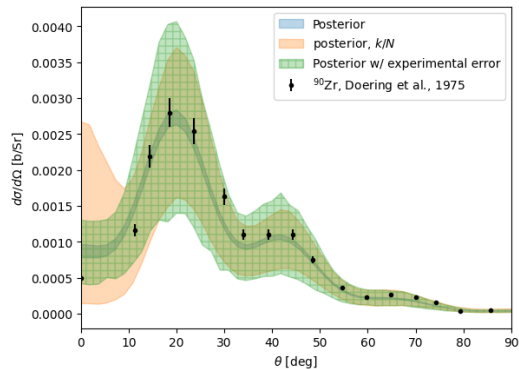
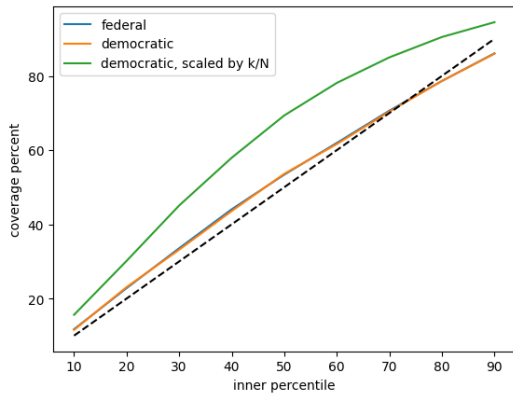


Results

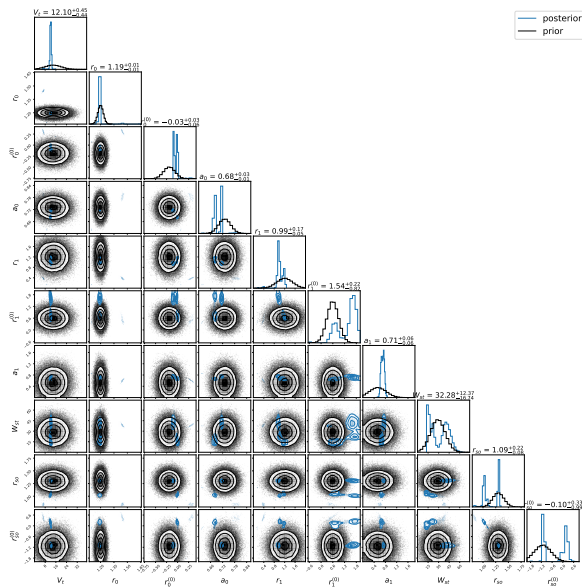
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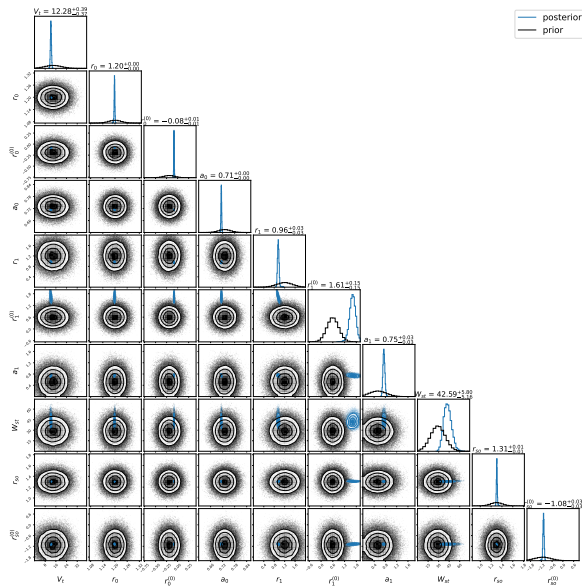
Predictive posteriors and empirical coverage



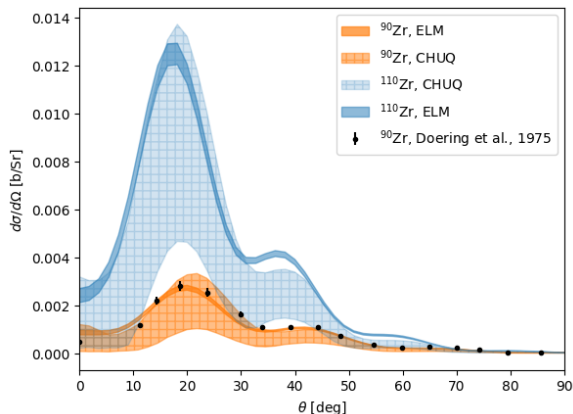
Multimodality, parameter ambiguity in radii



Multimodality, parameter ambiguity in radii



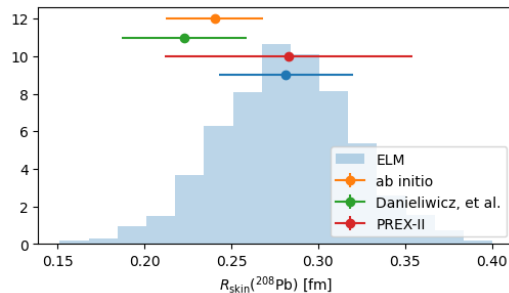
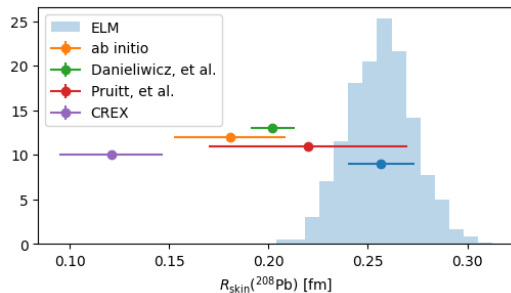
Extrapolation and uncertainty quantification

Extrapolation in $(N - Z)/A$ 

Problem: Model is function of $(N - Z)/A$, but model discrepancy model has no dependence on input space at all!

Bayesian inference of neutron skins

Preliminary ...



Summary and outlook

ELM coming soon!

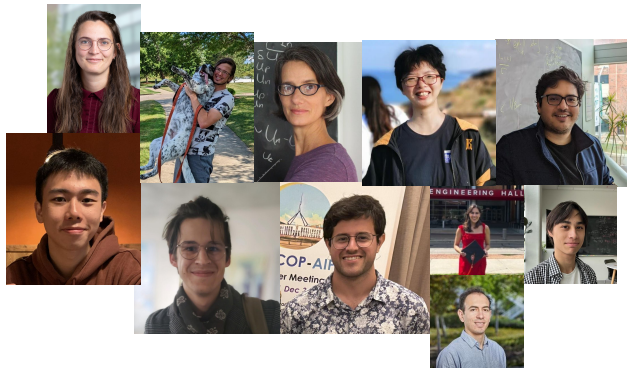
- reasonable posteriors without re-scaling, prediction of $R_{\text{skin}}(^AZ)$
- an open-source software ecosystem for reaction phenomenology has been developed

Outlook and stuff I'd like to do (or am doing)

- correlated experimental error, model discrepancy, errors in angle
- multi-modality in full model calibration? What orthogonal constraints can we fold in?
- model mixing
- microscopic (Skyrme-based) and ab-initio (NCSM) optical potentials for (p, n)
- ELM dispersive

Thank you for your time!

In collaboration with Filomena Nunes and the Few Body Reactions Group



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Reactions Group

References I

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-  Baye, Daniel (2015). “The Lagrange-mesh method”. In: *Physics reports* 565, pp. 1–107.
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-  Beyer, Kyle A. et al. (2025). “Uncertainty quantification of optical models in fission fragment deexcitation”. In: *Phys. Rev. C* 112 (2), p. 024604. DOI: [10.1103/4kny-mj7g](https://doi.org/10.1103/4kny-mj7g). URL: <https://link.aps.org/doi/10.1103/4kny-mj7g>.
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References II

-  Danielewicz, Paweł et al. (2017). “Symmetry energy III: Isovector skins”. In: *Nuclear Physics A* 958, pp. 147–186.
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Extra slides: ELM physical model

$$U(r, E, A, Z) = U_0(r, E, A) \pm \frac{N - Z}{N + Z} U_1(r, E, A) + V_C$$

$$U_0(r, E, A) = V_0(E) f((r - R_0(A))/a_0) + iW_0(E) f((r - R_w(A))/a_w) \\ - 4ia_w W_{s,0}(E) \frac{d}{dr} f((r - R_w(A))/a_w) + 2(\boldsymbol{\sigma} \cdot \boldsymbol{\ell}) V_{so} \frac{1}{r} \frac{d}{dr} f((r - R_{so}(A))/a_{so})$$

$$U_1(r, E, A) = V_1(E) f((r - R_1(A))/a_1) - 4ia_w W_{s,1}(E) \frac{d}{dr} f((r - R_w(A))/a_w)$$

$$f(x) = \frac{1}{1 + \exp(x)}$$

$$R_i(A) = r_{i,0} + r_{i,A} A^{1/3}$$

$$E_C = 6Z\alpha\hbar c/5R_C$$

$$V_{0(1)}(E) = V_{0(1)} - \alpha(E - E_C)$$

$$W(E) = W \left/ \left(1 + \exp\left(\frac{\eta_W - (E - E_C)}{\lambda_W}\right) \right) \right.$$

$$W_{s,0(1)}(E) = W_{s,0(1)} \left/ \left(1 + \exp\left(\frac{(E - E_C) - \eta_s}{\lambda_s}\right) \right) \right.$$

Extra slides: ELM statistical model and calibration

- Goodman-Weare¹⁸ affine-invariant stretch move sampler from emcee¹⁹, using default step-size ($a = 2$)
- 256 chains, 200,000 samples each
- massively parallel calibration requires ~ 3 days to run
- priors taken from CHUQ posterior
- priors for r_1^0 , r_1^A and a_1 taken from Danielewicz²⁰
- each data sector re-scaled to be equally important
- fractional uncertainties reasonably consistent with CHUQ:
 $\frac{d\sigma^{(n,n)}}{d\Omega}$: 35%, $\frac{d\sigma^{(p,p)}}{d\Omega}$: 44%, $A_y^{(n,n)}$: 140%, $A_y^{(p,p)}$: 155%, $\frac{d\sigma^{(p,n)}}{d\Omega}$: 55%

¹⁸Goodman and Weare 2010, *Communications in applied mathematics and computational science* 5(1), pp. 65–80

¹⁹Foreman-Mackey et al. 2013, *PASP* 125(925), p. 306

²⁰Danielewicz et al. 2017, *Nuclear Physics A* 958, pp. 147–186

To re-scale or not to re-scale (the likelihood)

Pruitt et al²¹ re-scaled the log likelihood by k/N (# of free parameters / # of experimental data points):

$$\ln p(\mathcal{D}|\boldsymbol{\alpha}, \mathcal{M}) = -\frac{k}{N} \frac{1}{2} \left(\boldsymbol{\Delta}^\top \cdot \boldsymbol{\Sigma} \cdot \boldsymbol{\Delta} + \ln \left((2\pi)^k |\boldsymbol{\Sigma}| \right) \right)$$

We do not re-scale for ELM

²¹CD Pruitt et al. (2023). “Uncertainty-quantified phenomenological optical potentials for single-nucleon scattering”. In: *Physical Review C* 107.1, p. 014602.

Charge exchange to isobaric analog states

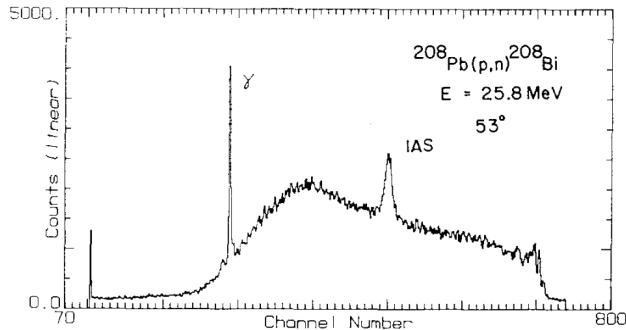
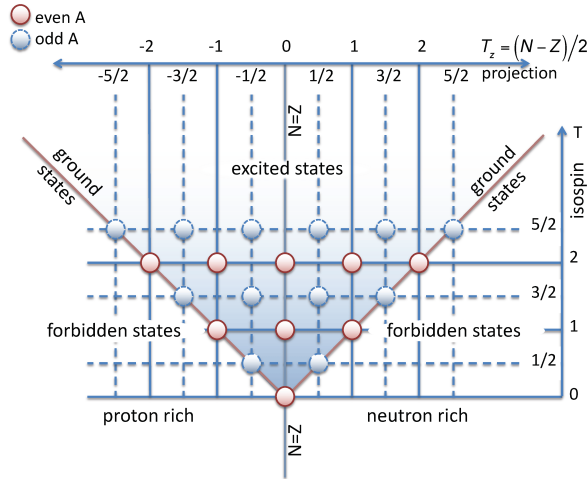


Fig. 1. Time-of-flight neutron spectrum for the reaction $^{208}\text{Pb}(p,n)^{208}\text{Bi}$ at 25.8 MeV showing the broad peak of the isobaric analogue state (IAS) and a narrow peak caused by incomplete rejection of γ -rays. The abscissa gives the time with units of 0.50 nsec per channel; higher energy neutrons are to the right in the spectrum. The flight path was approximately 9 m.

Reprinted from²²

²²SD Schery et al. (1974). "The (p, n) reaction to the isobaric analogue state of high-Z elements at 25.8 MeV". In: *Nuclear Physics A* 234.1, pp. 109–129.

Charge exchange to isobaric analog states



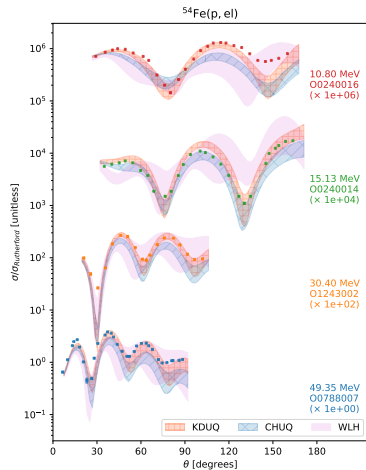
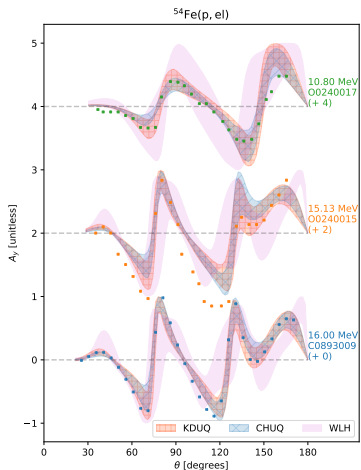
Reprinted from²³

²³Michael A Bentley (2022). "Excited States in Isobaric Multiplets—Experimental Advances and the Shell-Model Approach". In: *Physics* 4.3, pp. 995–1011. ISSN: 2624-8174. DOI: [10.3390/physics4030066](https://doi.org/10.3390/physics4030066). URL: <https://www.mdpi.com/2624-8174/4/3/66>

jit $\mathcal{R}$ demos and tutorials: built-in optical potentials

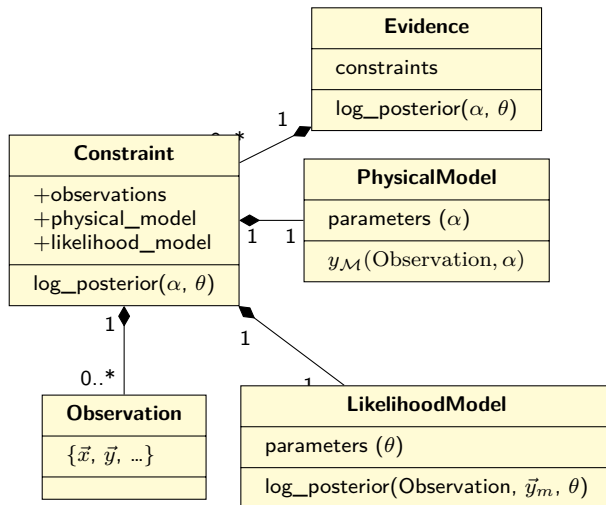
```
[2]: from jitR.optical_potentials import chuq, kduq, wlh
```

https://github.com/beykyle/jitr/blob/main/examples/notebooks/builtin_omps_uq.ipynb



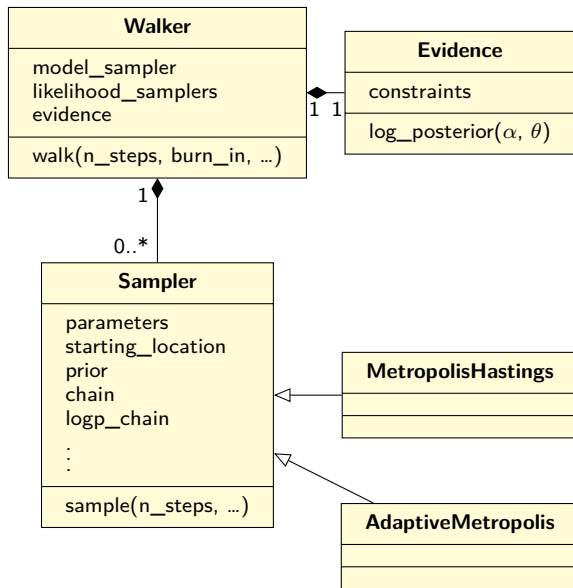
rxmc: spotlight on the body of evidence and the likelihood

- curating a body of evidence from multiple constraints
- Bayesian calibration with flexible likelihood models
- designed with `jitR` and `exfor-tools` in mind
- built-in MCMC, or just calculate the log posterior as input to another MCMC package (e.g. `surmise`)
- handles parametric likelihood models with Gibbs sampling
- alpha stage:
<https://github.com/beykyle/rxmc/>

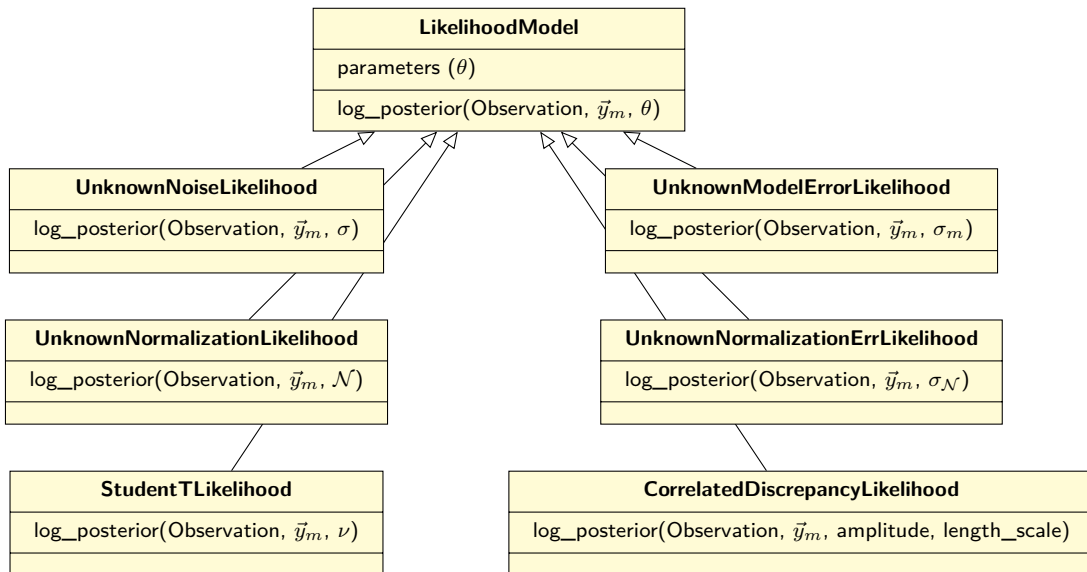


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- alpha stage:
<https://github.com/beykyle/rxmc/>



rxmc: spotlight on the body of evidence and the likelihood



Aside: systematic error in the likelihood

Common approach is to introduce a normalization parameter $\mathcal{N}$:

$$\chi^2(\alpha) = \sum_i \frac{(y_i - \mathcal{N}y_{\mathcal{M}}(x_i; \alpha))^2}{\sigma_i^2} + \frac{(\mathcal{N} - 1)^2}{\sigma_{\mathcal{N}}^2} \quad (3)$$

This is incorrect! The correct approach is actually²⁴:

$$\chi^2(\alpha) = \sum_i \frac{(\mathcal{N}y_i - y_{\mathcal{M}}(x_i; \alpha))^2}{(\mathcal{N}\sigma_i)^2} + \frac{(\mathcal{N} - 1)^2}{\sigma_{\mathcal{N}}^2} \quad (4)$$

One can implicitly marginalize over $\mathcal{N}$, which requires use of a non-diagonal covariance²⁵. This relates to Peelle's Pertinent Puzzle in nuclear data evaluation²⁶.

²⁴D'Agostini 1994, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 346(1-2), pp. 306–311; Thompson and Nunes 2009

²⁵Barlow 2021, *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 987, p. 164864

²⁶Frühwirth et al. 2012, vol. 27, p. 00008; Neudecker et al. 2012, *Nuclear science and engineering* 170(1), pp. 54–60

rxmc demos and tutorials: handling multiple constraints with systematic errors

- statistical only:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2$$

- fit (fractional) statistical error:

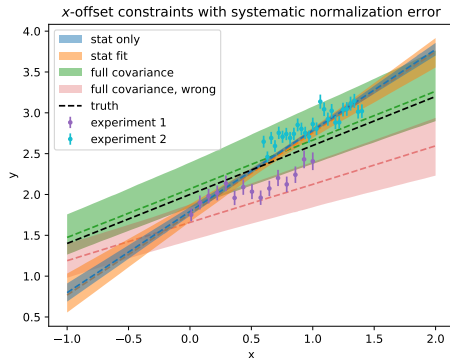
$$\Sigma_{ij} = \delta_{ij} (\eta y_{\mathcal{M}}(x_i; \alpha))^2$$

- full covariance:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2 + \sigma_N^2 y_{\mathcal{M}}(x_i; \alpha) y_{\mathcal{M}}(x_j; \alpha)$$

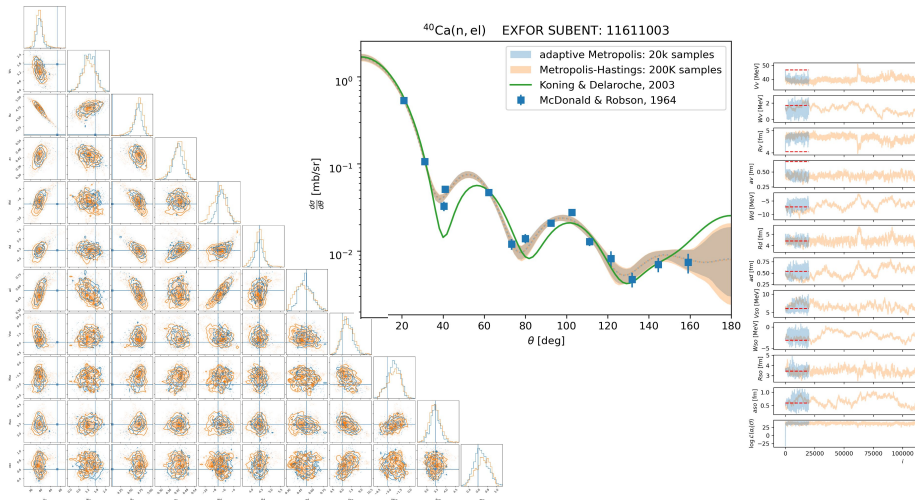
- full covariance, wrong:

$$\Sigma_{ij} = \delta_{ij} \sigma_i^2 + \sigma_N^2 y_i y_j$$



https://github.com/beykyle/rxmc/blob/main/examples/systematic_err_demo.ipynb

rxmc demos and tutorials: compare sampling algorithms for fitting an optical potential



https://github.com/beykyle/rxmc/blob/main/examples/fitting_an_optical_potential.ipynb

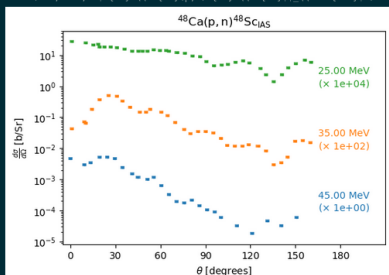
Curating a body of evidence

```
entry_pn = ExforEntry(
    entry="00178",
    reaction=ca48_pn,
    quantity="dXS/dA",
    vocal=True,
    parsing_kwargs={
        "statistical_err_labels": ["DATA-ERR"],
        "systematic_err_labels": ["ERR-1"],
    },
)

Found subentry 00178006 with the following columns:
['ERR-1', 'ERR-2', 'ERR-3', 'ERR-4', 'EN', 'Q-VAL', 'ANG-ERR', 'DATA-ERR', 'ANG-CM', 'DATA', 'ERR-T']
Found subentry 00178007 with the following columns:
['ERR-1', 'ERR-2', 'ERR-3', 'ERR-4', 'EN', 'Q-VAL', 'ANG-ERR', 'DATA-ERR', 'ANG-CM', 'DATA', 'ERR-T']
Found subentry 00178008 with the following columns:
['ERR-1', 'ERR-2', 'ERR-3', 'ERR-4', 'EN', 'Q-VAL', 'ANG-ERR', 'DATA-ERR', 'ANG-CM', 'DATA']

fig, ax = plt.subplots(1, 1, figsize=(6, 4))
entry_pn.plot(
    ax, offsets=100, label_kwargs={"label_xloc_deg": 165, "label_offset_factor": 0.0001}
)
ax.set_xlim([-5, 210])
ax.set_title(ax.title.get_text() + r"${\text{IAS}}$")

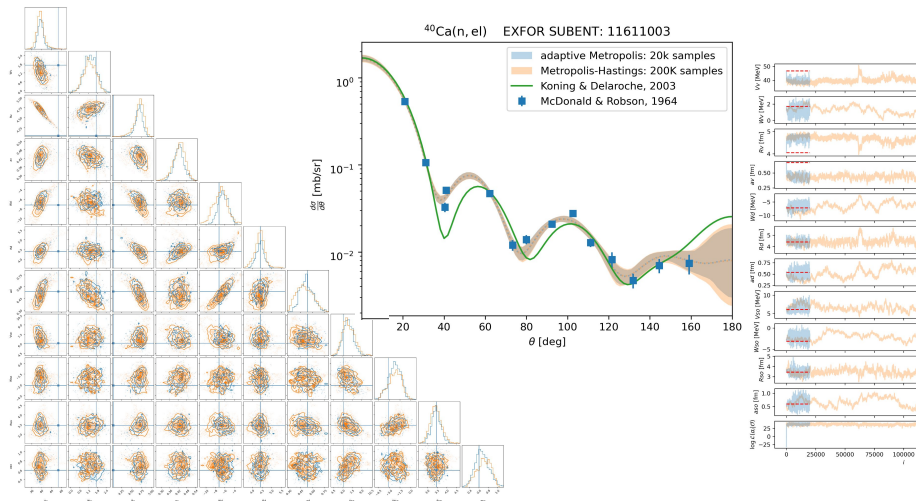
Text(0.5, 1.0, '$^{48}$Ca(p,n)$^{48}$ScIAS')
```



https://github.com/beykyle/exfor_tools

- reproducible curation
- systematically improvable with new evaluations

Comparison of sampling algorithms for fitting an optical potential



https://github.com/beykyle/rxmc/blob/main/examples/fitting_an_optical_potential.ipynb