

Introduction to High Vacuum in Accelerators

Jonathan YEMANE

Le 18/06/2026

Introduction to High Vacuum in accelerators.

Part1: vacuum basis

- 1-1 Kinetic behavior of gas molecules
- 1-2 Mean free path / residence time
- 1-3 outgassing/desorption

Part3: Technology evolution

- 3-1 Pressure distribution in accelerators
- 3-2 Distributed pumping
 - 3-2-a getter pump
 - 3-2-b NEG coating

Part2: vacuum in accelerators

- 2-1 interaction residual gas - beam
 - 2-1-a Coulomb scattering
 - 2-1-b Bremsstrahlung
 - 2-1-c Ionization energy loss
 - 2-1-d ion accumulation
- 2-2 Interaction particles surface
 - 2-2-a Synchrotron radiation an Photon stimulated desorption
 - 2-2-b Ion stimulated Desorption
 - 2-2-c Electron stimulated Desorption an electron-cloud

What does vacuum mean?

Vacuum: an absence of matter ?

- Space where the molecules are strongly rarefied compared to atmospheric pressure

10^{-9} mbar $\sim$ 25 millions/cm³ at 20°C

Pollution source, background noise,
parasitic interactions

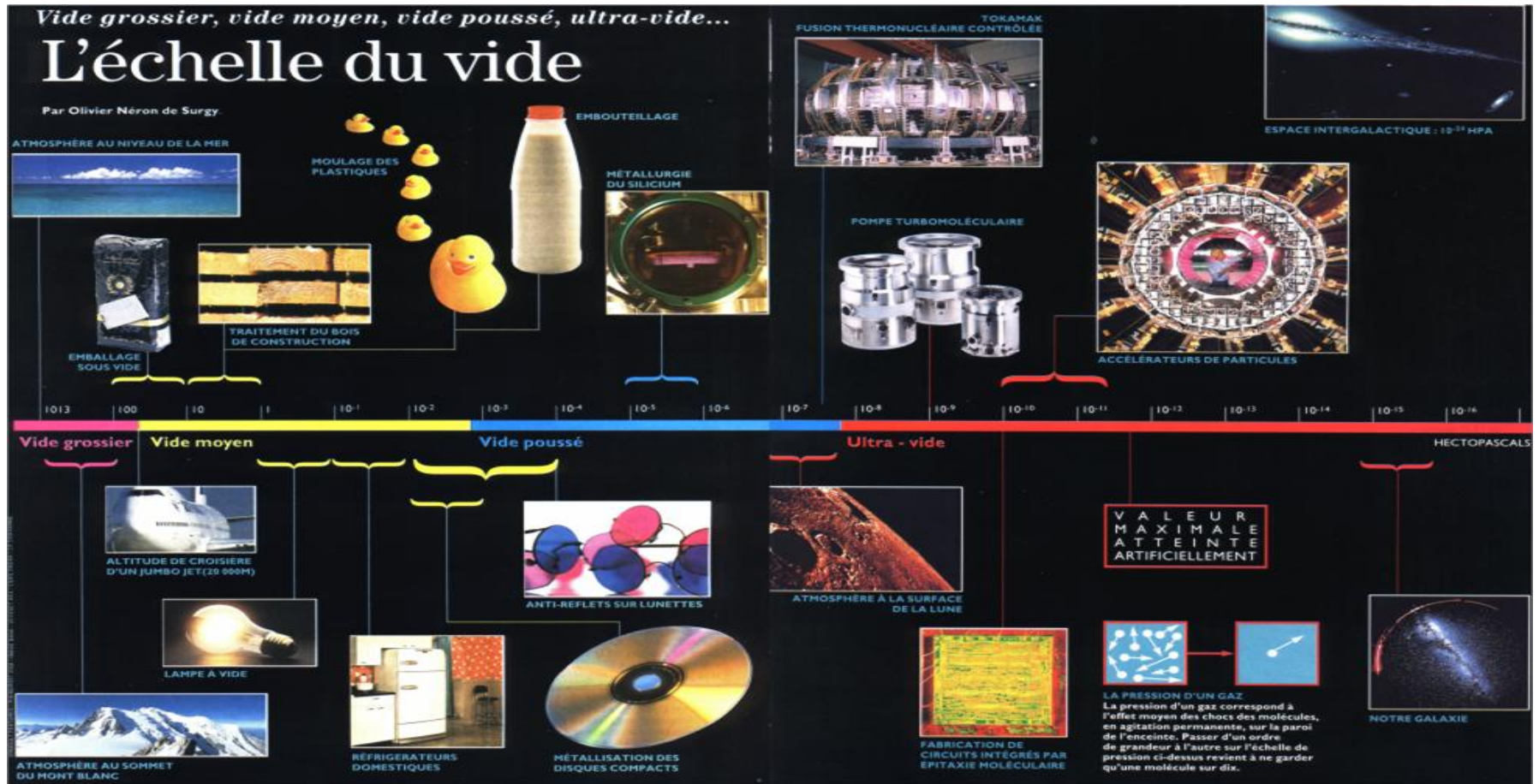
- quantum vacuum:

State minimum energy, vacuum fluctuation, ...

a vacuum: full of matter !!

The scale

Pressure $100 \text{ Pa}(\text{N/m}^2) = 1 \text{ mbar} = 0,75 \text{ Torr}$



Maxwell-Boltzmann velocity distribution

$$\bar{V} = V_m = \sqrt{\frac{8RT}{\pi M}} \quad V_{rms} = c = \sqrt{\frac{3RT}{M}}$$

At 20°C

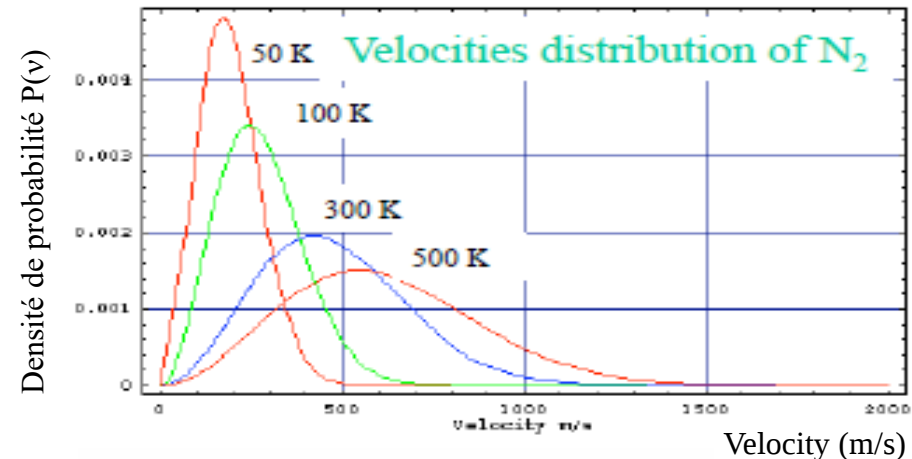
Ar ~ 394 m/s

N₂ ~ 471 m/s

He ~ 1246 m/s

Dependent upon composition of gas and temperature

$$P(v)dv = \frac{dn}{n} = 4\pi \left(\frac{m}{2\pi kT}\right)^{3/2} v^2 e^{-\frac{mv^2}{2kT}} dv$$



residence time on the surface

Frenkel formula

$$\tau_0 \sim 10^{-13} \text{ s}$$

$$\tau = \tau_0 \cdot e^{\frac{E}{RT}}$$

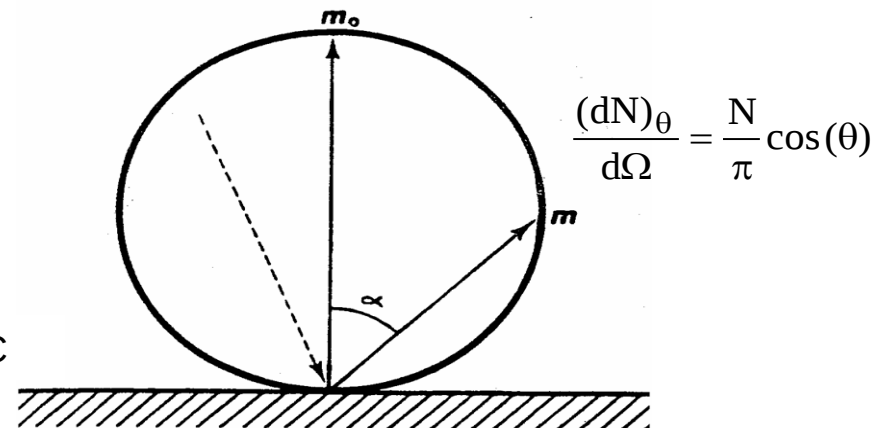
E desorption energy

$$20 < E < 25 \text{ kcal/mol}$$

$$0,65 < E < 1,08 \text{ eV/molecule}$$

$$80 \text{ s} < \tau < 5 \text{ days at } 20^\circ \text{C}$$

Increase pumping time



Emission under cosine law

Monolayer formation time

And considering the molecule as spherical, for perfect atomic arrangement:

$$\tau = \frac{4}{d^2 \cdot n \cdot v_m}$$

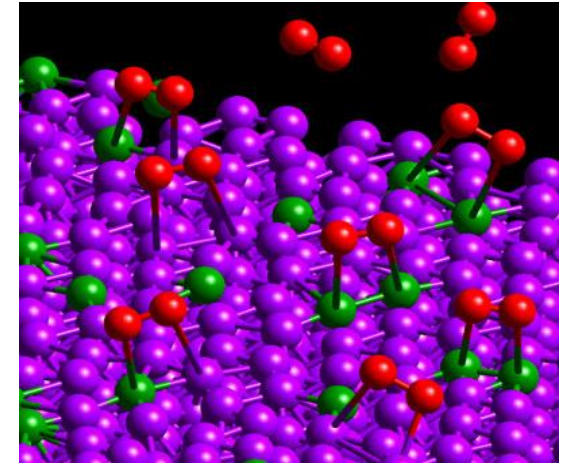
(very approximative)

n molecular density
d molecular diameter
 v_m mean velocity

Order of magnitude one monolayer $\sim 10^{15}$ molécules/cm²

$\tau \sim 40$ min at $P=10^{-9}$ mbar (N_2)

$\tau \sim 3$ sec at $P=10^{-6}$ mbar (N_2)



The concept of monolayer can appreciate the cleanliness of the surface

Mean Free Path

It is the path length that a molecules traverse between two succesives impacts with other molecules.

$$l_m = \frac{1}{\sqrt{2} \pi \cdot d^2 \cdot n}$$

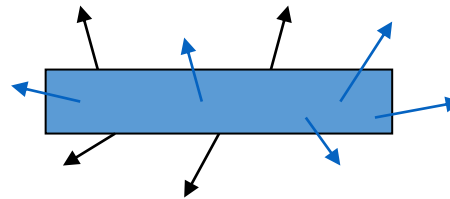
$l_m \sim 64$ km at $P=10^{-9}$ mbar (N_2)

$l_m \sim 65$ m at $P=10^{-6}$ mbar (N_2)

$l_m \sim 0,65$ mm at $P=10^{-1}$ mbar (N_2)

outgassing

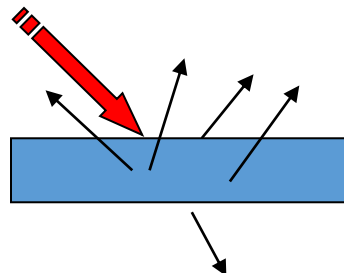
Outgassing is the **spontaneous** evolution of gas from solid or liquid.



desorption /degassing

Desorption is the release of adsorbed chemical species from the surface of a solid or liquid.

Particles (ions, photons, electrons,...)



Desorption: gas flow

At the thermodynamic equilibrium

Gas flow

$$Q = P \cdot \frac{dV}{dt} = \mathbf{P \cdot S} = \frac{dN}{dt} \cdot K \cdot T$$

S effectif pumping speed
N molecule number

$$\text{Pa.m}^3/\text{s} (=10 \text{ mbar.l/s})$$

The outgassing rate τ : a gas flow per unit of surface

$$\tau = Q / A$$

A desorption area

$$\text{Pa.m}^3.\text{s}^{-1}.\text{m}^{-2} = \text{Pa.m.s}^{-1} (=10^{-3} \text{ mbar.l.s}^{-1}.\text{cm}^{-2})$$

Importance of the gas flow

At the thermodynamic equilibrium, the pressure in the system is expressed by:

$$P = \frac{\sum Q}{S} + P_0$$

P_0 limit pressure of pumping system

S effective pumping speed

Decrease of the work pressure P

⇒ **Increase the pumping speed**

Approximately 2 to 3 orders of magnitude with a high cost and limited effectiveness conductances

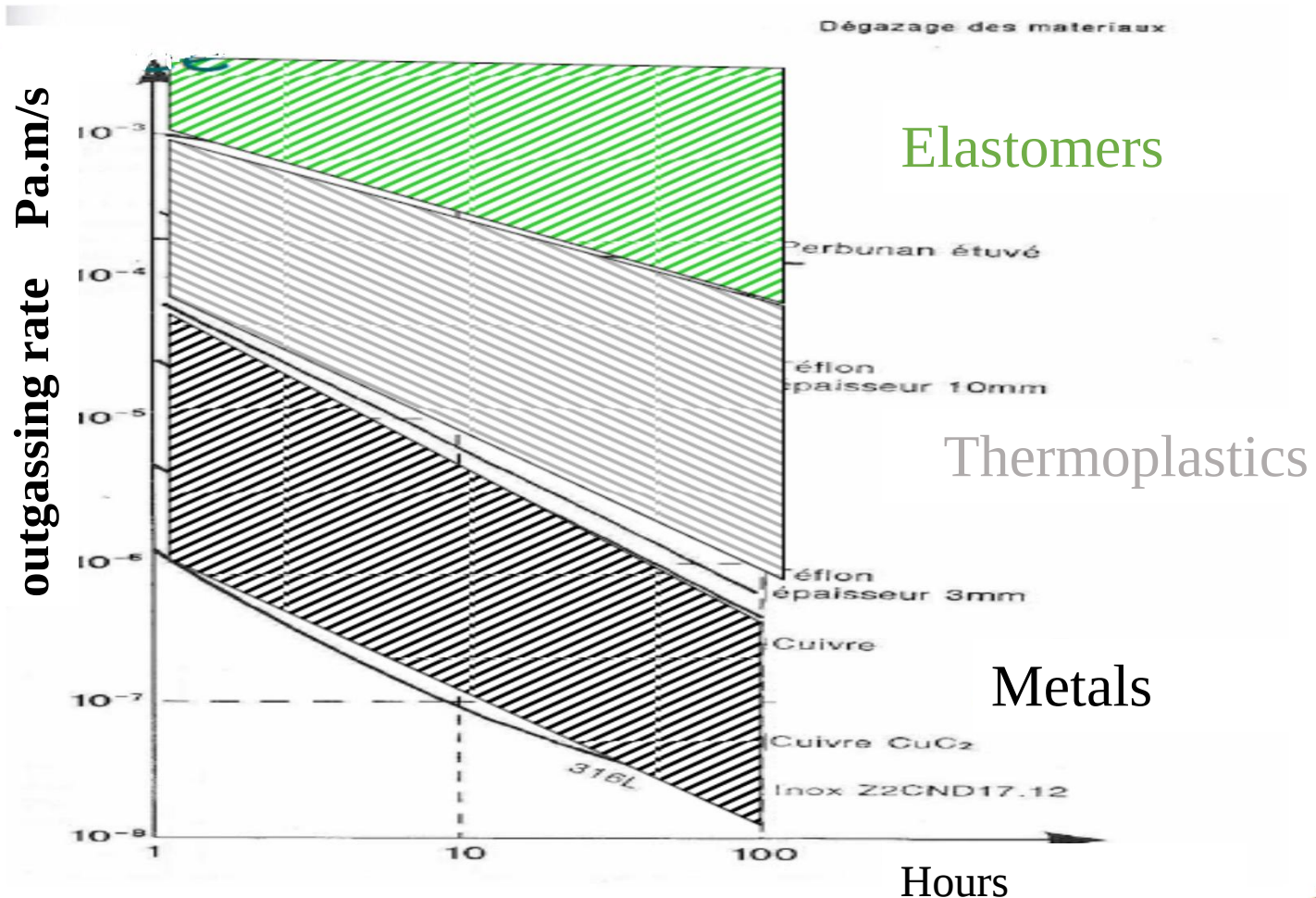
⇒ **Decrease the flow**

About 10 orders of magnitude depending on the choice of materials and treatments

Attention to minority surfaces

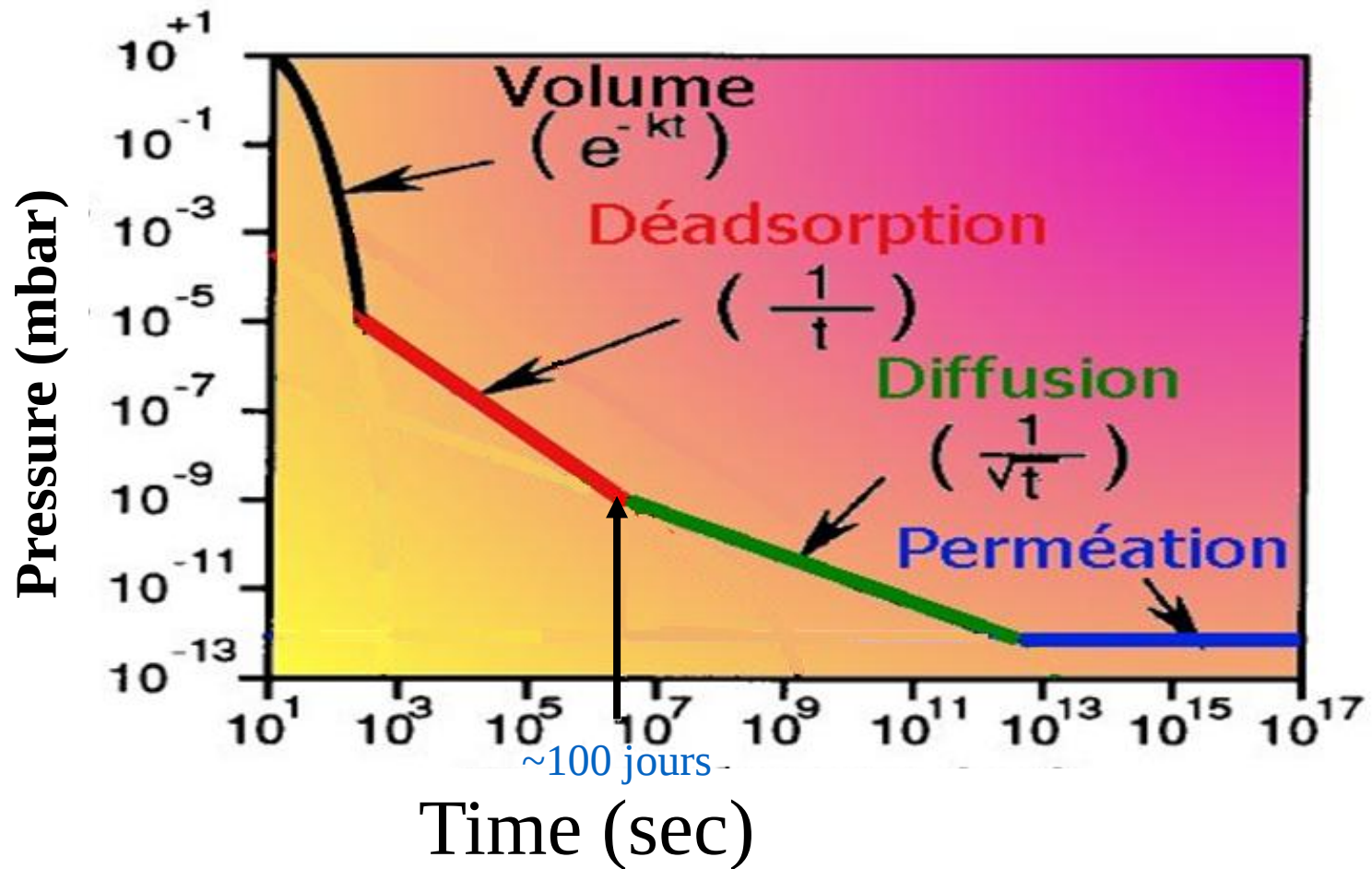
First of all decrease the outgassing rates

Time dependance of the outgassing flow rate for différent materials

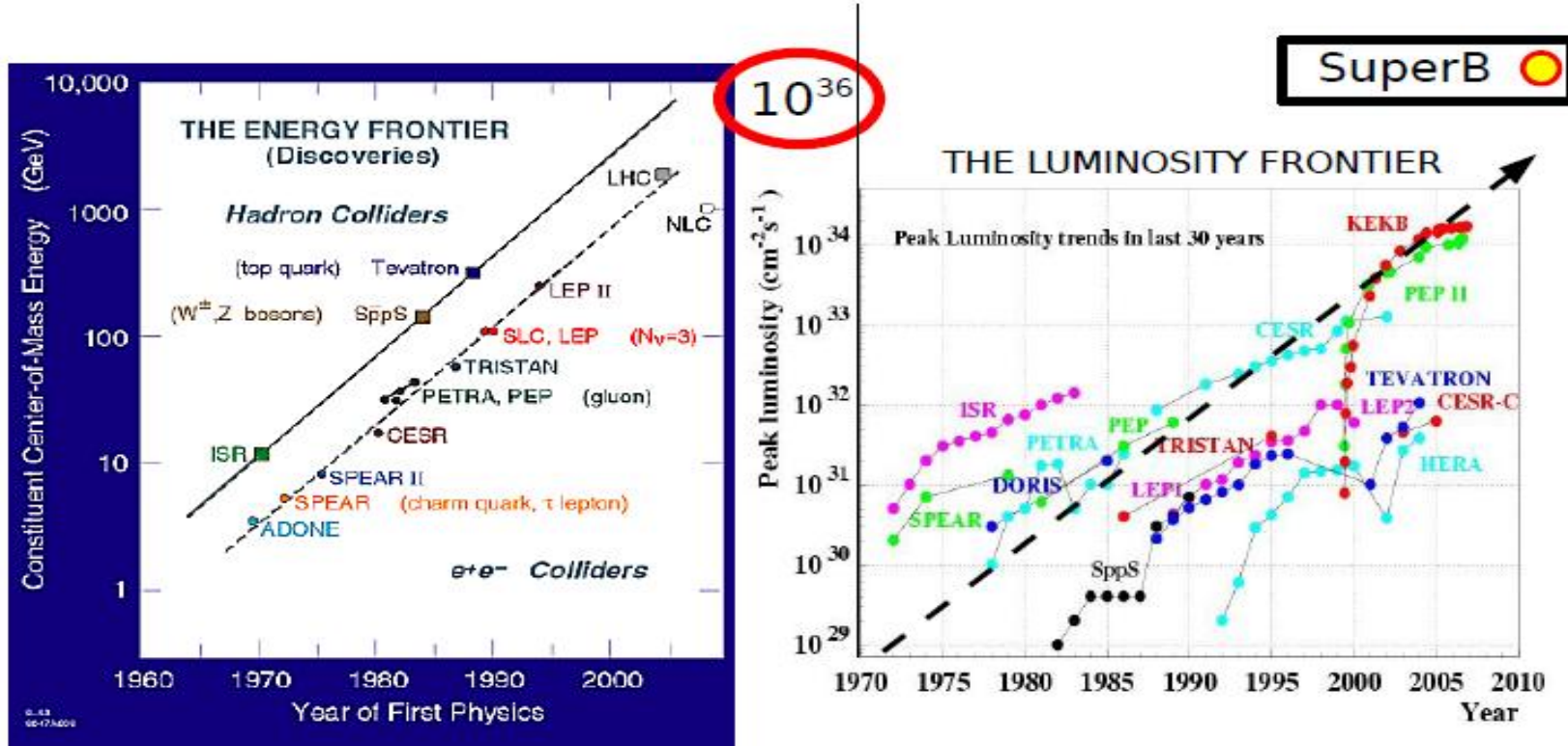


$$1 \text{ Pa.m/s} = 10^{-3} \text{ mbar.l.s}^{-1}.\text{cm}^{-2}$$

The pressure vs. Time behavior of a typical vacuum system (metal chamber)



Vacuum in accelerators :



$$\mathcal{L} \sim f_{\text{coll}} \frac{N^+ N^-}{4\pi \sigma_x \sigma_y} = 10^{36} \text{ cm}^{-2} \text{ s}^{-1}$$

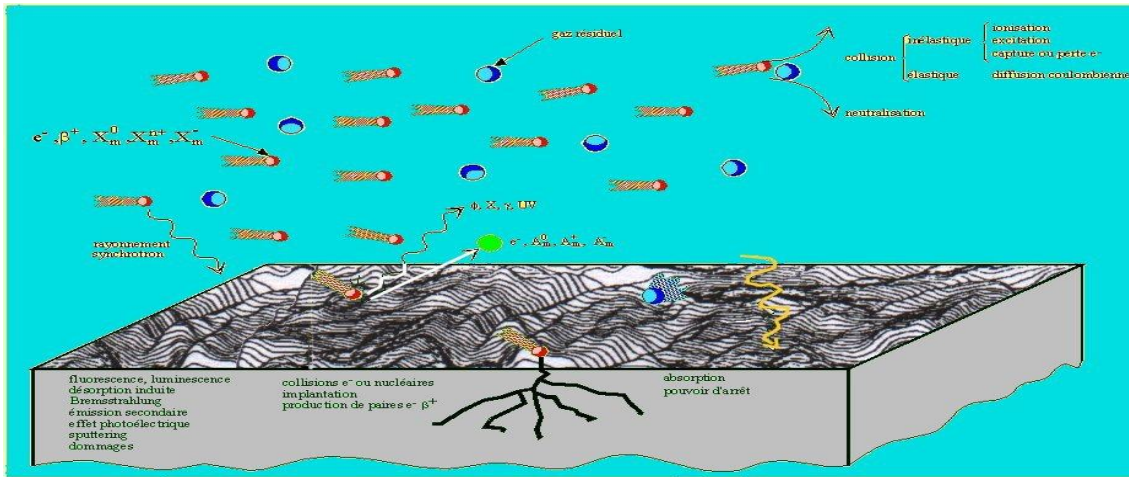
High luminosity

small beam, high intensity and high frequency



Level and quality of the residual gas (dynamic pressure $\sim 10^{-9}$ mbar)

Vacuum in accelerators :



**Beam & Vacuum & Surface:
a difficult coexistence!**

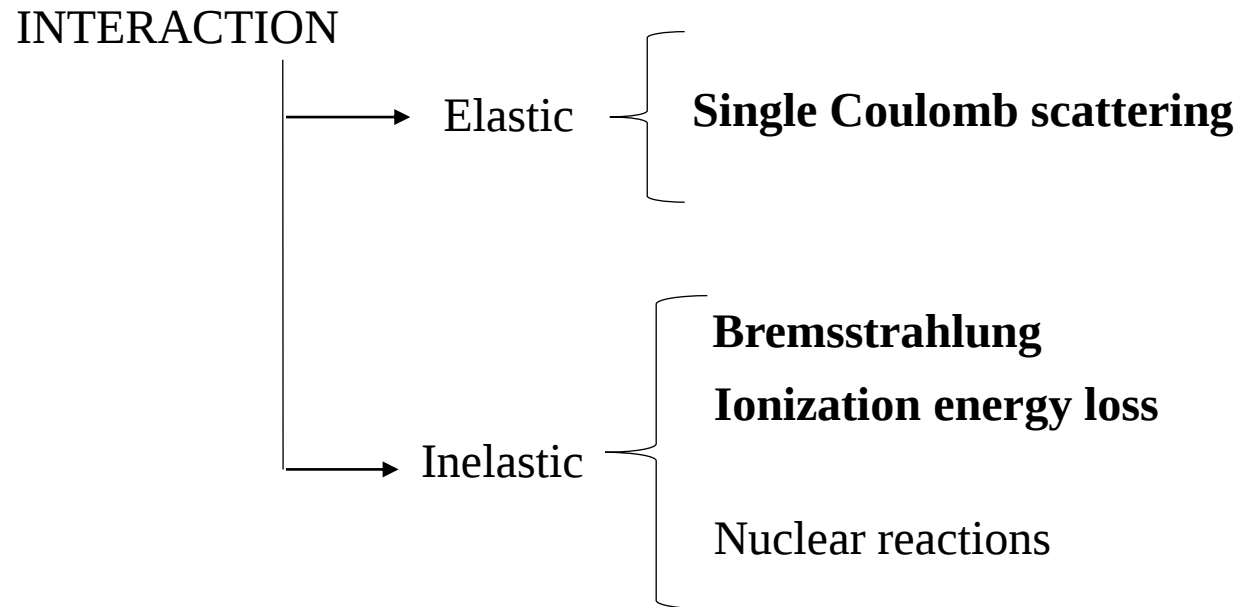
Interaction: Beam - résiduel gas

- ➡ Changing the shape and the beam energy
- ➡ Secondary particle emission (ion, e-, photons)
- ➡ Loss of beam particles

Interaction: Particules (secondary, primary) - surface

- ➡ Desorption by electron, photon, ion impact
- ➡ Secondary emission

Interaction: Beam- residual gas



Interaction: Beam- residual gas

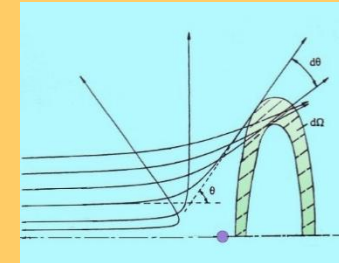
Elastic interaction $\left\{ \begin{array}{l} \text{Single Coulomb scattering} \end{array} \right.$

$$\frac{d\varepsilon}{dt} \propto Z^2 \cdot P$$

ε Emittance is proportional to residual gas pressure and depends square residual gas atomic number (Z)

La Diffusion Coulombienne

$$F = \frac{1}{4\pi\epsilon_0} \xi_i \xi_c e^2 \frac{\exp(-d/a)}{d^2}$$



$$\sigma_{\theta \max} = 2\pi \left(\frac{\xi_i Z_c e^2}{8\pi\epsilon_0 m_i v_i^2} \right)^2 \cot^2(\theta_{\max} / 2)$$

Nuclear scattering LHC

Life time limit $\sim 100h$ $P \sim 10^{-8}$ mbar H_2 equivalent

Synchrotron Soleil

Life time limit $\sim 60h$ $P \sim 10^{-9}$ mbar N_2 equivalent $E = 2,15$ GeV

Interaction: Beam- residual gas

→ Inelastic interaction

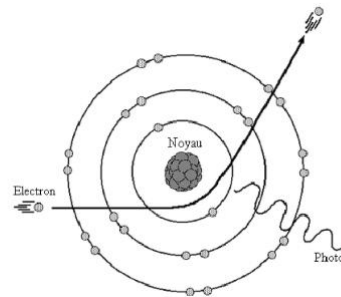
Bremsstrahlung

Ionization energy loss
Electron capture
Electron loss
Nuclear reactions

Bremsstrahlung

Material	Z	A	Radiation length		Density [g/cm ³] () is for gas [g/l]
			X_0 [g/cm ²]	[cm]	
H ₂	1	1.01	61.28	865	0.0708(0.090)
D ₂	1	2.01	122.6	757	0.162(0.177)
He	2	4.00	94.32	755	0.125(0.178)
Li	3	6.94	82.76	155	0.534
Be	4	9.01	65.19	35.3	1.848

Valid for light particles (e⁺,e⁻)



$$\frac{dE}{E} = -\frac{dx}{\chi_0}$$

$$\chi_{0[g/cm^2]} = \frac{714.6A_c}{Z_c(Z_c + 1) \ln(287/\sqrt{Z_c})}$$

Loss beam enregy

the radiation interacts with the surface

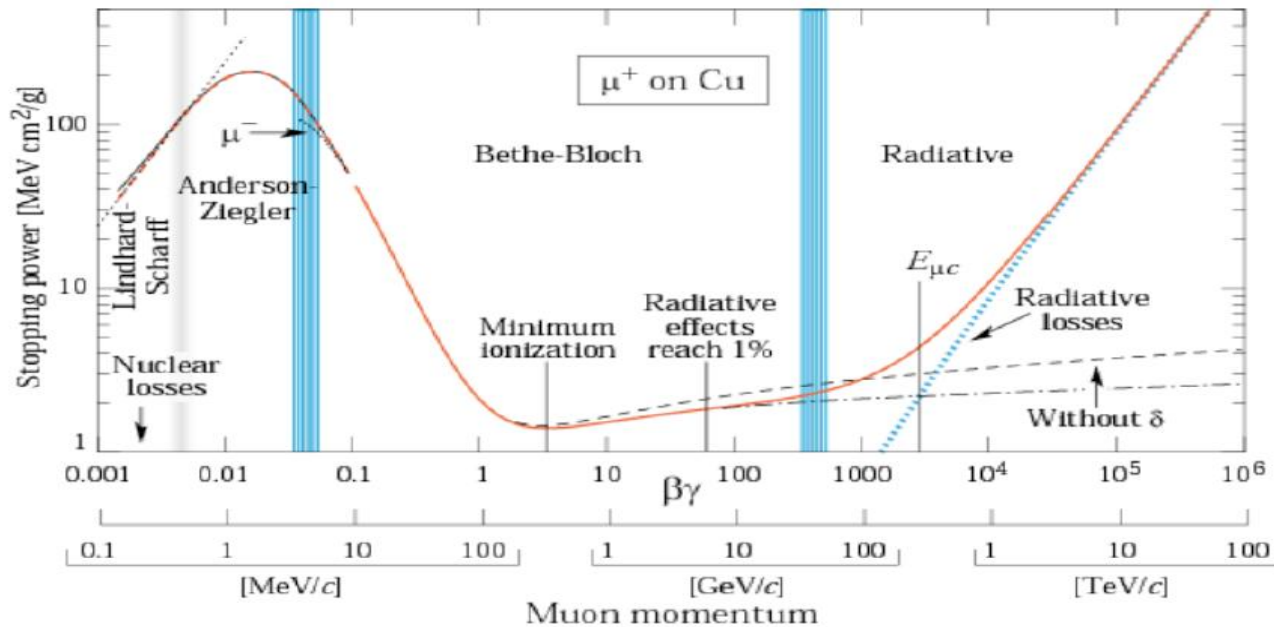


Desorption

A very strong dependence on the target atomic number is observed, which means that even a small amount of a heavy gas may have a higher influence on the lifetime than the usually dominating hydrogen content.
about 4.10⁶ times lower for proton

Interaction: Beam – residual gas

Ionization energy loss



Inelastic interaction {
 Bremsstrahlung
Ionization energy loss
 Electron capture
 Electron loss
 Nuclear reactions

Loss beam energy

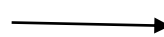
LHC

$\Delta E \sim 1350 \text{ eV}$ per LHC revolution sur 7 TeV with 10-8 mbar (H₂)

« SOLEIL » synchrotron

Life time limit $\sim 60 \text{ h}$ With $P \sim 10^{-9} \text{ mbar N}_2 \text{ eq.}$ $E = 2,15 \text{ GeV}$ RF acceptance $\epsilon_{\text{RF}} = 1,8 \%$

Secondary particule emission: **electrons, ions**



Interaction on the surface and desorption

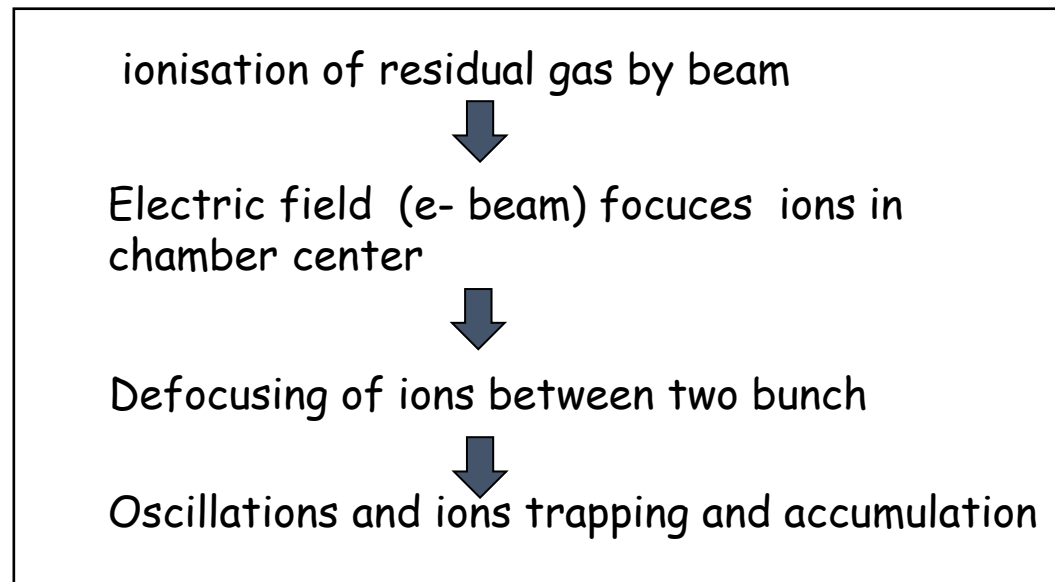
Ion accumulation

negative charge beam

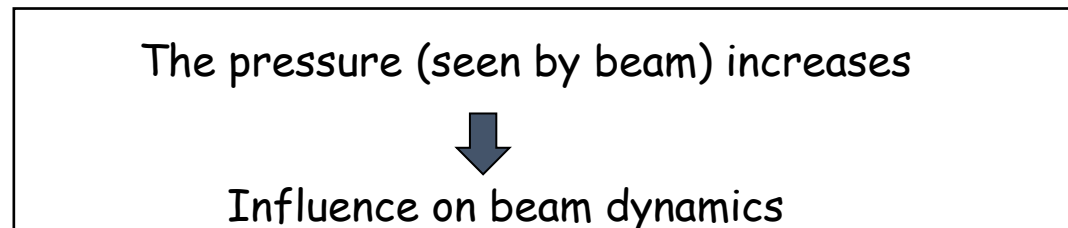
An important phenomenon for negative charges beam, for high current and for short bunch spacing

Pressure due to the ion containment

mechanism:



consequence:



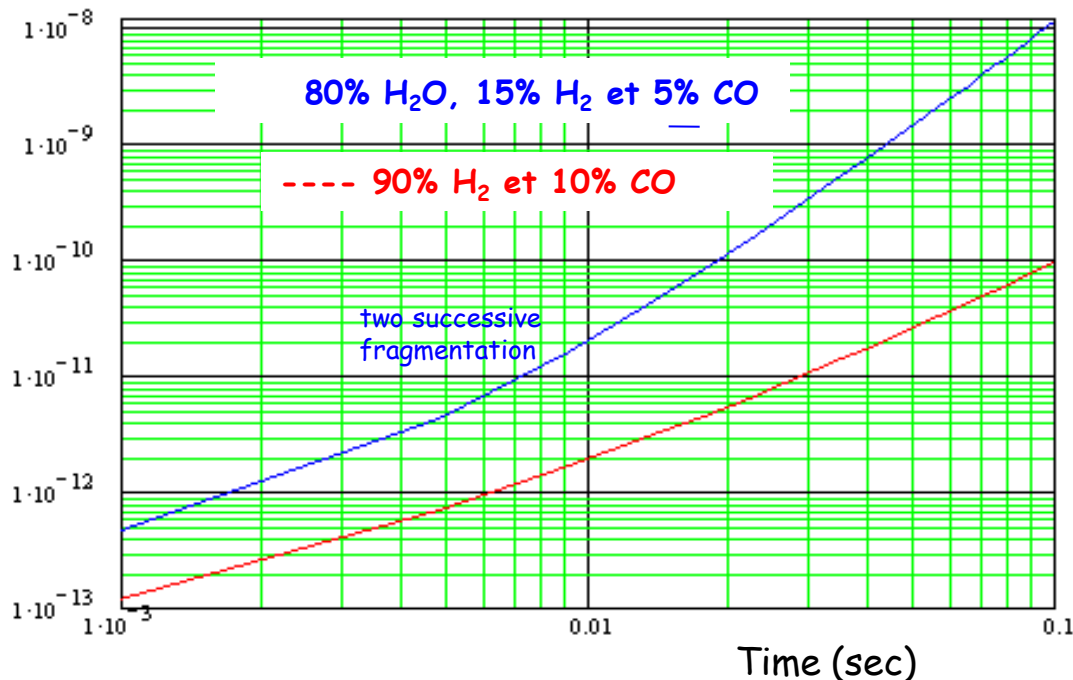
Ion accumulation

Ionic pressure seen by beam

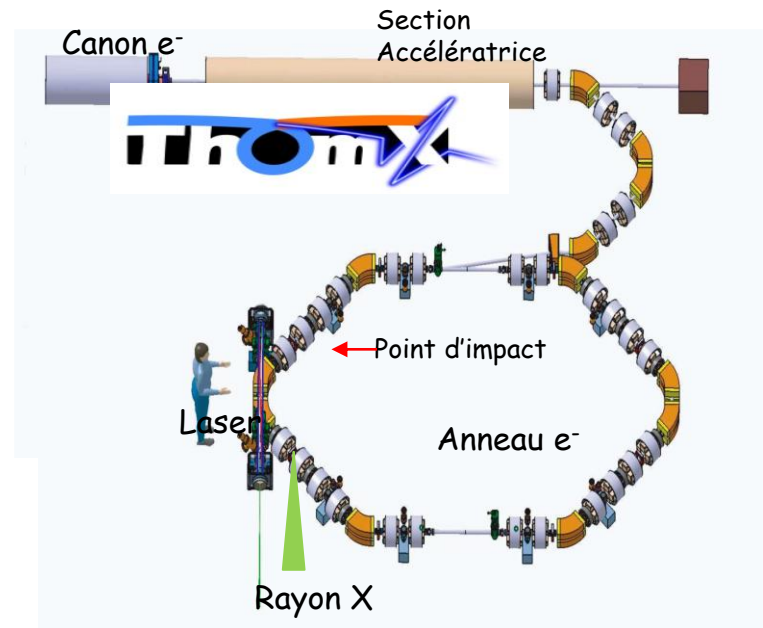
Increased pressure due to the ion density.

$$P_{\text{ion}} = \frac{I}{e} \cdot \frac{\eta^* \cdot k \cdot T}{\beta \cdot C \cdot \pi \cdot 2 \cdot \sigma_x \cdot \sigma_y}$$

Ionic pressure (mbar)



At l' IP $\sigma_x = \sigma_y = 0.07 \text{ mm}$ $P \approx 2 \cdot 10^{-9} \text{ mbar}$



Solutions:

Electrodes + gap between two injections ($5 \mu\text{s}$)

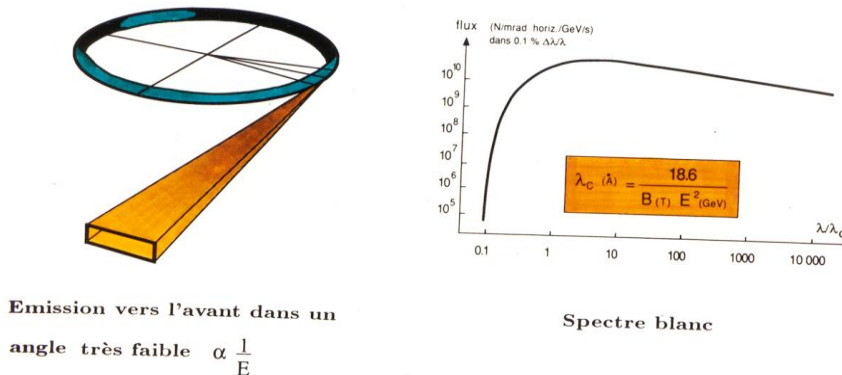
C. Bruni, J. Haissinski, T. Demma

Synchrotron radiation

Any charged particle undergoing acceleration, centripetal or longitudinal, produces radiation.

Main source on circular colliders

RAYONNEMENT SYNCHROTRON : DIPOLES



Power loss

electron $P_0 \text{ [W/m]} = 88.57 \frac{E[\text{GeV}]^4}{2\pi \rho[\text{m}]^2} I[\text{mA}]$

proton $P_0 \text{ [W/m]} = 7.79 \cdot 10^{-12} \frac{E[\text{GeV}]^4}{2\pi \rho[\text{m}]^2} I[\text{mA}]$

Critical energy

Electrons: $\varepsilon_c [\text{eV}] = 2.218 \cdot 10^3 \frac{E[\text{GeV}]^3}{\rho[\text{m}]}$

Protons: $\varepsilon_c [\text{eV}] = 3.5835 \cdot 10^{-7} \frac{E[\text{GeV}]^3}{\rho[\text{m}]}$

Flux

• $\Gamma [\text{photons} \cdot \text{m}^{-1} \cdot \text{s}^{-1}] = 1.288 \cdot 10^{17} \frac{E[\text{GeV}]}{\rho[\text{m}]} I[\text{mA}]$

• $\Gamma [\text{photons} \cdot \text{m}^{-1} \cdot \text{s}^{-1}] = 7.017 \cdot 10^{13} \frac{E[\text{GeV}]}{\rho[\text{m}]} I[\text{mA}]$

synchrotron radiation

some examples

		Soleil	KEK-B		LEP			LHC	
			LER	HER	Inj.	1	2	Inj.	Col.
Particule		e ⁻	e ⁺	e ⁻	e ⁻	e ⁻	e ⁻	p	p
Courant	mA	500	2600	1100	3	3	7	584	584
Energie	GeV	2.75	3.5	8	20	50	96	450	7000
R. courbure	m	5.36	16.31	104.46	2962.96			2784.302	
Puissance	W/m	14 030	20 675	5 820	0.8	30	955	0	0.2
En. critique	eV	8 600	5 800	11 000	6 000	94 000	660 000	0	44
Flux	photons/m/s	3 10¹⁹	7 10¹⁹	1 10¹⁹	3 10¹⁵	7 10¹⁵	3 10¹⁶	7 10¹⁵	1 10¹⁷
Dose a 3000 h	photons/m	4 10 ²⁶	8 10 ²⁶	1 10 ²⁶	3 10 ²²	7 10 ²²	3 10 ²³	7 10 ²²	1 10 ²⁴

Photon stimulated desorption

photodesorption Yield η

$$\eta = \frac{\text{desorbed molecule}}{\text{Incident photons}}$$

material type, surface finish, energy and particle type, incidence angle, particle dose

Experimental determination

photodesorption Yield η at machine start-up

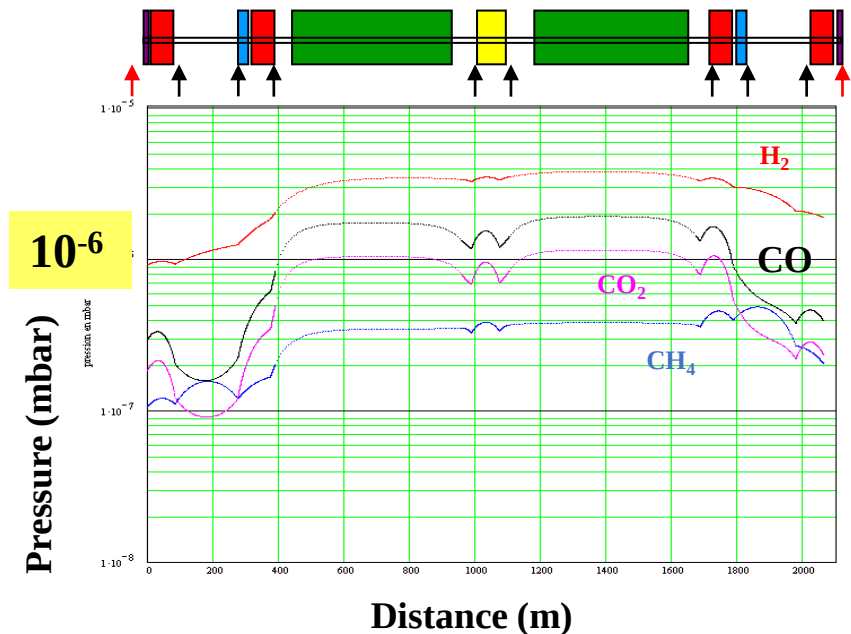
gas	η (mol/ph)
H ₂	$\sim 1.10^{-3}$
CH ₄	$\sim 1.10^{-4}$
Co	$\sim 5.10^{-4}$
Co ₂	$\sim 3.10^{-4}$

For baked Cu 150°C and Ec = 4,5 Kev

Gröbner and all, JVSTA 12(3) 1994

The static pressure increases by several orders of magnitude due to dynamic effects

the pressure distribution in Cell HER (SUPERB) with synchrotron radiation at machine start-up

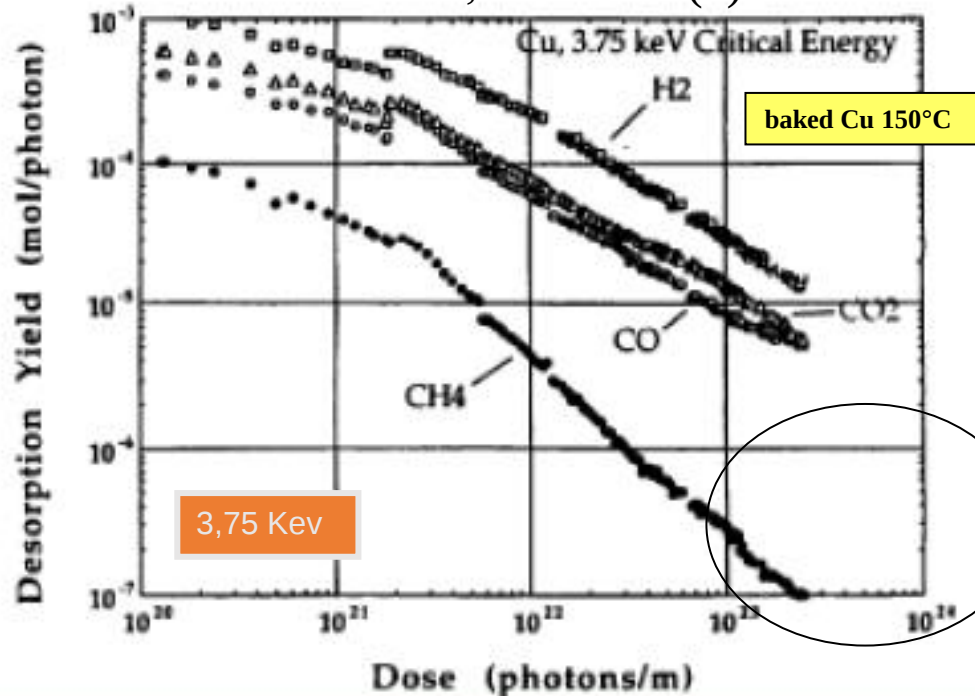


photodesorption

Conditioning - Scrubbing

Surface conditioning according to the received dose

Gröbner and all, JVSTA 12(3) 1994



After conditionning

D = photon dose 10^{25} photons/m ~ 360 A.h

gas	η (mol/ph)
H ₂	$\sim 6 \cdot 10^{-7}$
CH ₄	$\sim 2.5 \cdot 10^{-8}$
Co	$\sim 2.5 \cdot 10^{-7}$
Co ₂	$\sim 2.5 \cdot 10^{-7}$

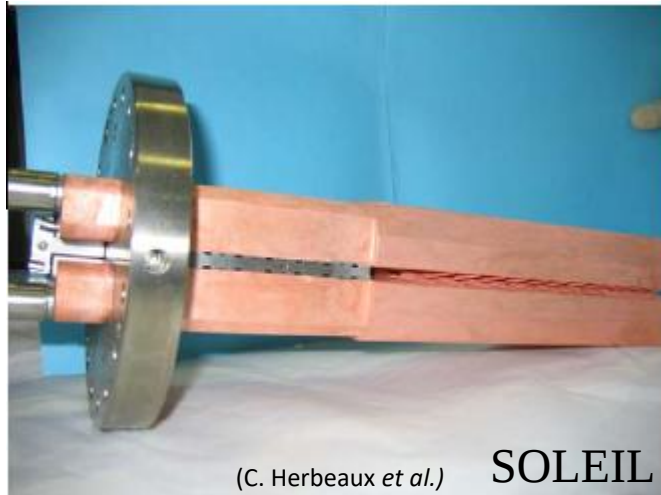
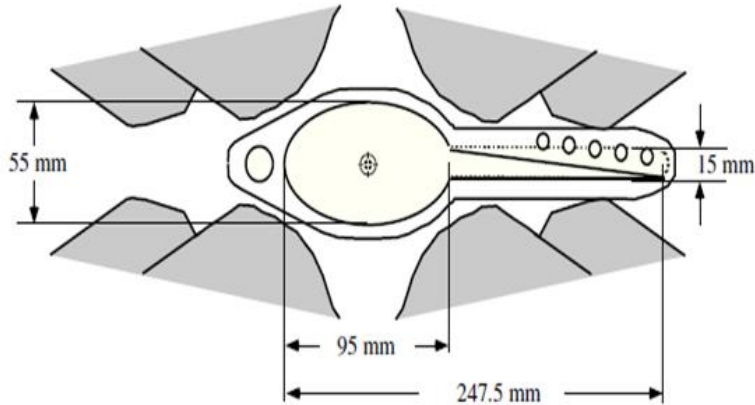
$$\eta = \eta_0 \cdot \left(\frac{D}{D_0} \right)^{-\alpha} \quad \alpha = 0.84$$



Scrubbing but needed time

Photon stimulated desorption

Important implications for the design



(C. Herbeaux *et al.*) **SOLEIL**

Absorbing power "crotch" : GLIDCOP copper cooled by water(256 W / mm²)

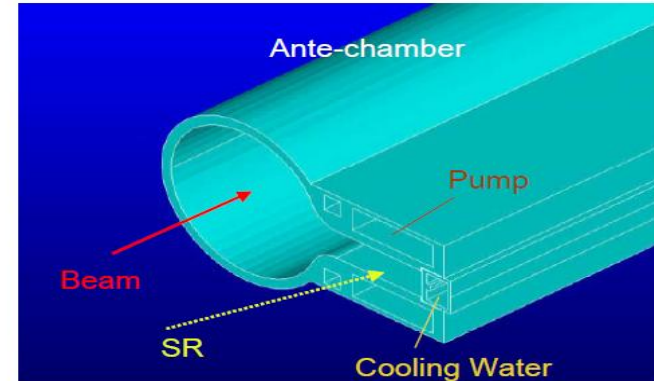


Figure 1: Typical structure of ante-chamber for the LER of Super-KEKB.

SOLEIL

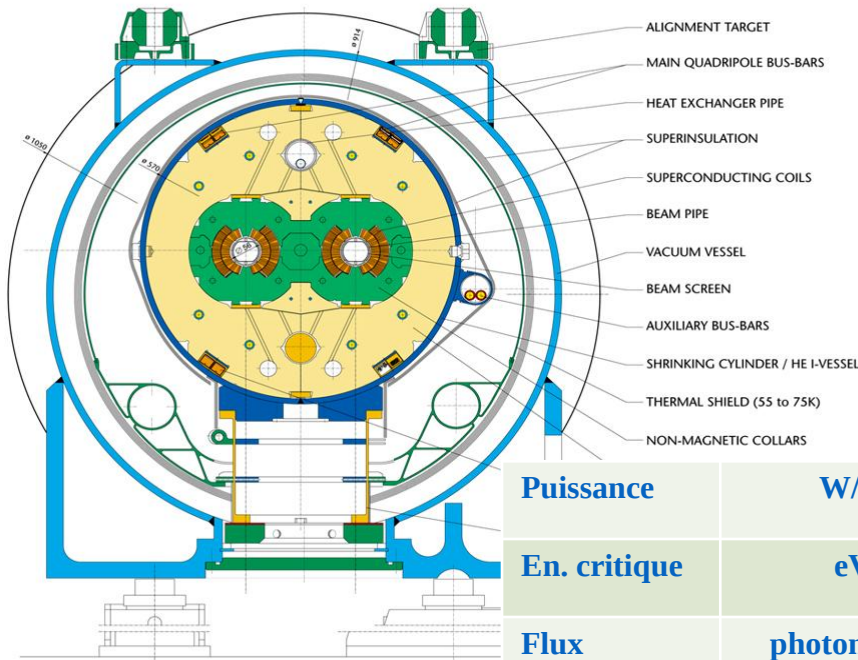
Puissance	W/m	14 030
En. critique	eV	8 600
Flux	photons/m/s	3 10 ¹⁹

Photon stimulated desorption

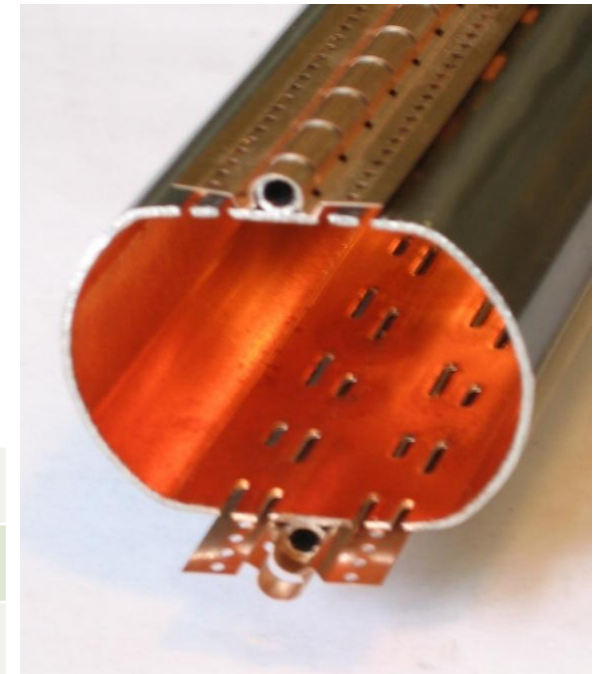
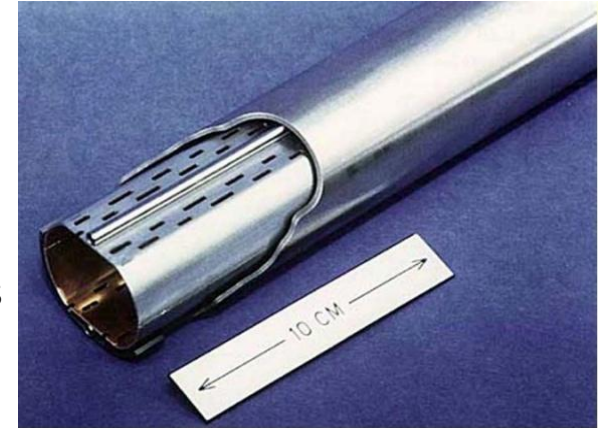
LHC Dipole Vacuum System

- Cold bore (CB) at 1.9 K which ensures **leak tightness**
- Beam screen (BS) at 5-20 K which intercepts **thermal loads** and acts as a screen

LHC DIPOLE : STANDARD CROSS-SECTION



Puissance	W/m	0.2
En. critique	eV	44
Flux	photons/m/s	$1 \cdot 10^{17}$



Courtesy N. Kos CERN AT/VAC

V. BAGLIN

Ion stimulated desorption

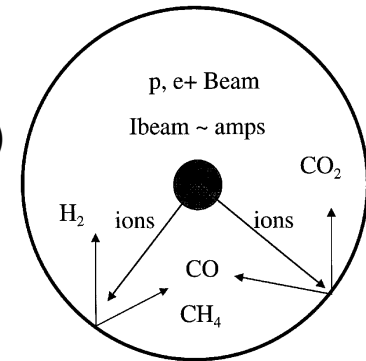
Positive charge beam

The mechanism:

ionization of the residual gas by the beam

Acceleration of these ions by the beam electric field (e +, p)

The ions/surface impact creates a molecular desorption

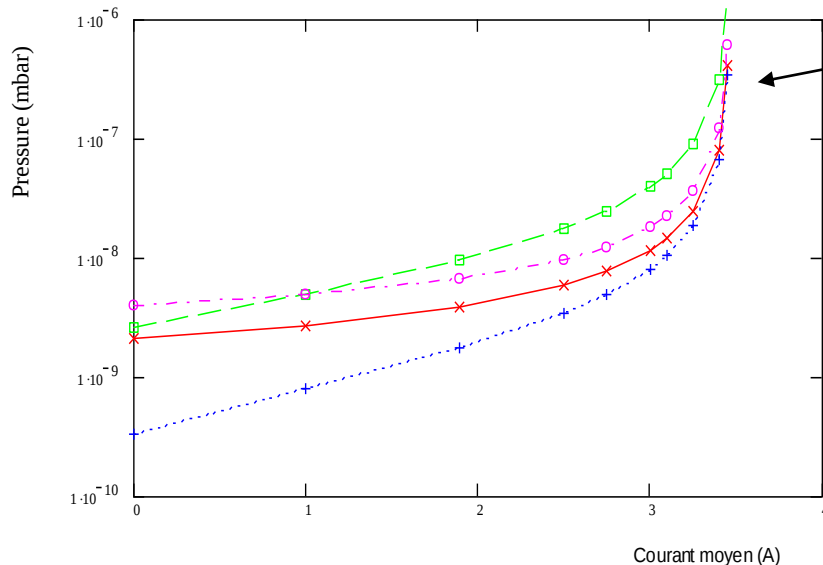


The important parameter

Rendement η = $\frac{\text{desorption molecular}}{\text{Incident ions}}$

$$\text{flux: } \Gamma_{\text{ion}} \approx \sigma \frac{I}{e} \cdot \frac{\beta \cdot P}{k \cdot T}$$

HER (SuperB) simulation



pressure divergence

$$I_{\text{crit}} \approx \frac{1}{\eta} \frac{eS}{\sigma}$$

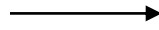
Effective Pumping speed

Cross section
(Avoid heavy gas)

• High current accelerators : ISR, SNS, LHC...

Ion stimulated desorption

Ionodesorption Yield η

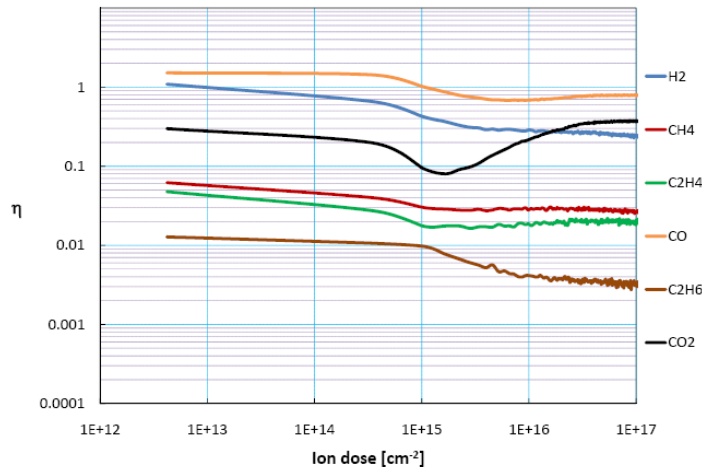


material type, surface finish, energy and particle type, incidence angle, particle dose

Experimental determination

η versus ion dose

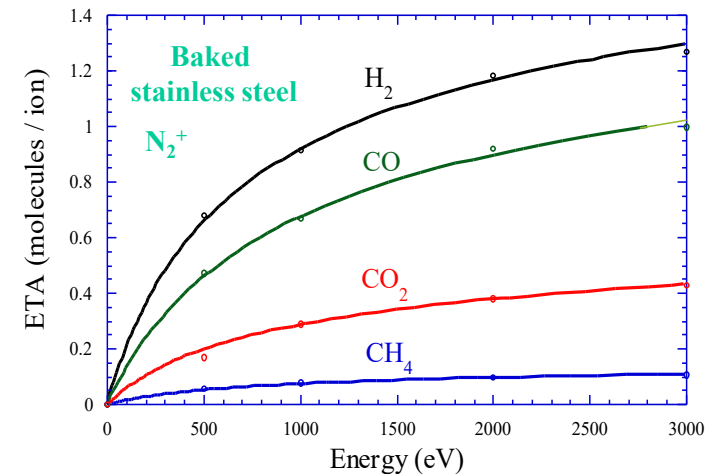
7 keV, Cu Baked, CO^+



G. Hulla, PhD Thesis, Vienna Tech. U, 2009

LHC : ion dose $\sim 3 \cdot 10^{15}$ ions/($\text{cm}^2 \cdot \text{an}$)

η versus energy



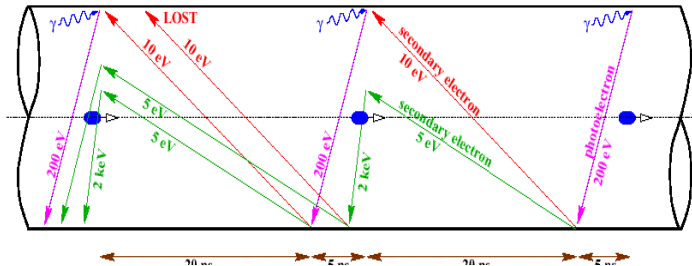
A.G. Mathewson, CERN ISR-VA/76-5

Inefficient scrubbing !!

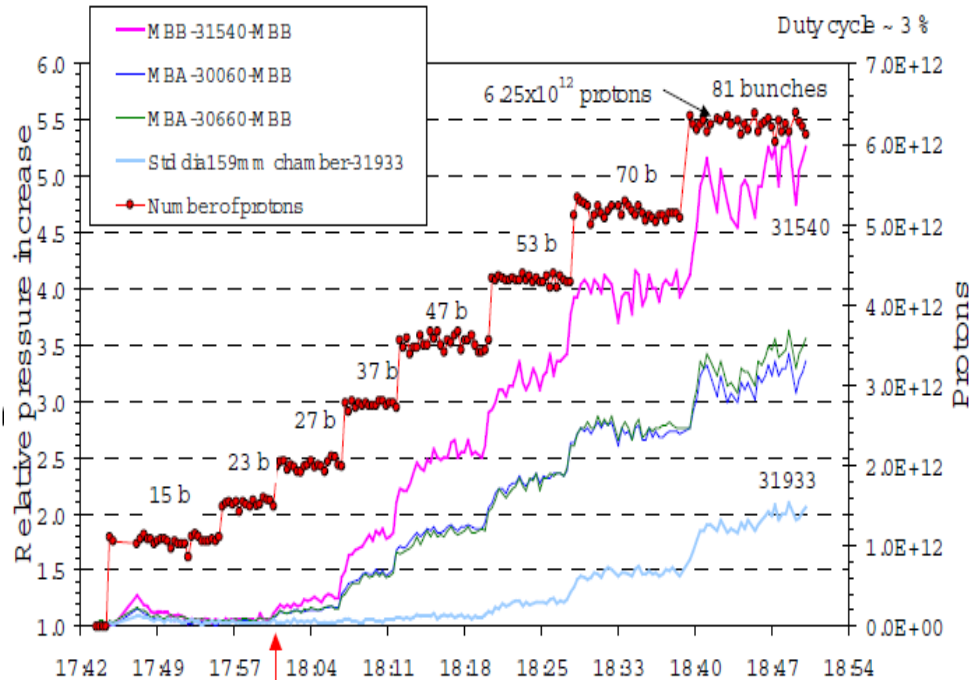
Electron-Cloud / Electron Stimulated Desorption

KEKB, PEP II, SPS ~ année 2000
RHIC, LHC

The mechanism:



F. Ruggiero *et al.*, LHC Project Report 188 1998, EPAC 98



J.M. Jimenez *et al.*, EPAC 2000, Vienna, Austria

photoelectrons acceleration (with electron coming from ionization) by the beam field E

Interaction photoelectron – surface

Molecular desorption

Secondary electron

Accélération E faisceau

Interaction electron surface

Increase pressure

e-Cloud

Increase emittance and beam loss, heat load, decrease luminosity

Electron-Cloud / Electron Stimulated Desorption

An important phenomenon for positive charges beam, for high current and for short bunch spacing

Main parameters:

- Beam structure(size,time between bunch)
 - bunch/charge
- chamber size and geometry
- secondary electron yeld (SEY)
 - photo-electrons rate
 - photons /electrons reflected
- electron stimulated desorption (ESD)



reduction

Some Solutions

material (coating), surface finish, scrubbing,
Cleaning electrodes, solenoidal field,
Geometry chamber (antichamber)

Electron-Cloud / Electron Stimulated Desorption Surface conditioning by electron impingement (Scrubbing)

Electron Stimulated Desorption (ESD)

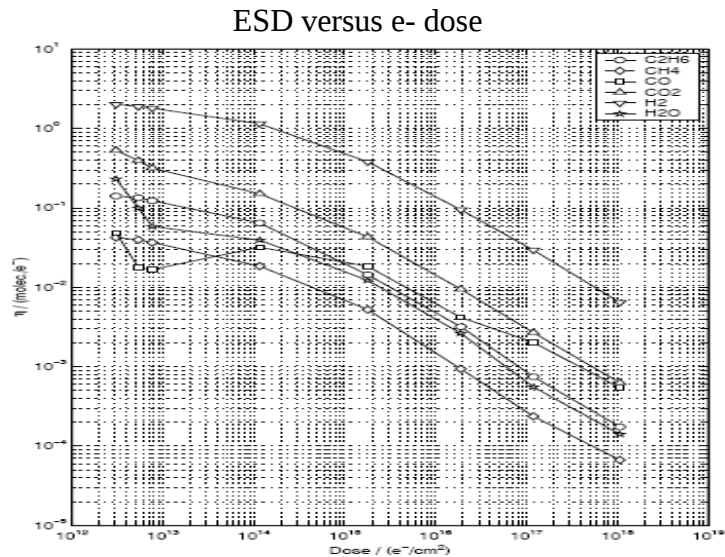


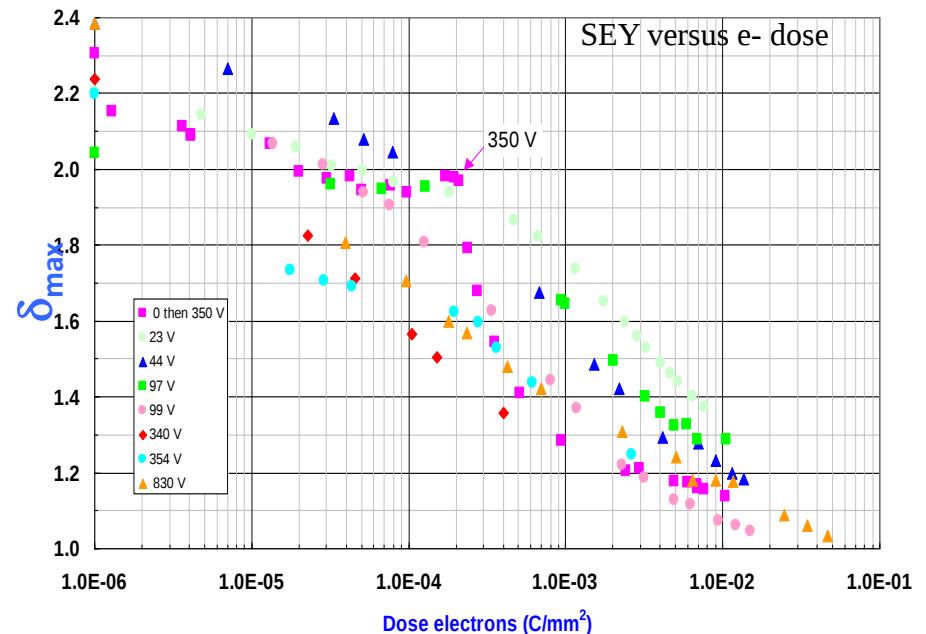
Figure 3: Effect of the electron dose on the electron induced desorption yield of an unbaked copper sample. The electron energy during bombardment and measurement was 300 eV.

G. Vorlaufer *et al.*, CERN VTN, 2000

Reduction ESD and SEY

Secondary electron yield (SEY)

$$\text{SEY} = \frac{\text{number of emitted electrons (secondary)}}{\text{number of impinging electrons (primary)}}$$



V. Baglin *et al.*, Chamonix, 2001

Carbone layer on surface (AES, XPS,...)

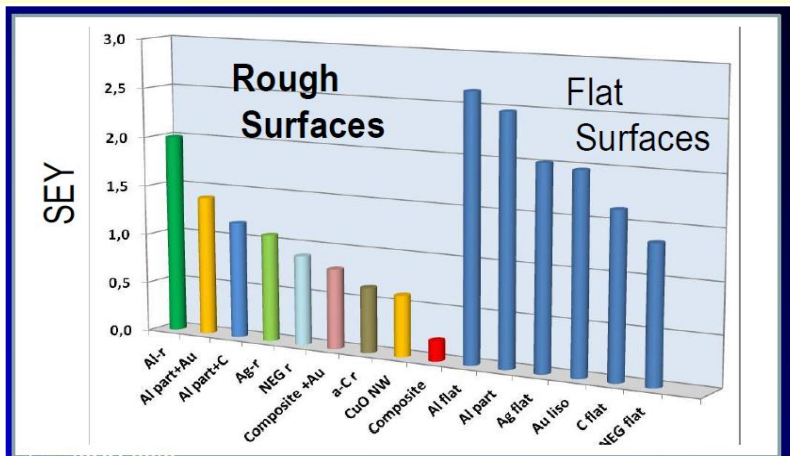
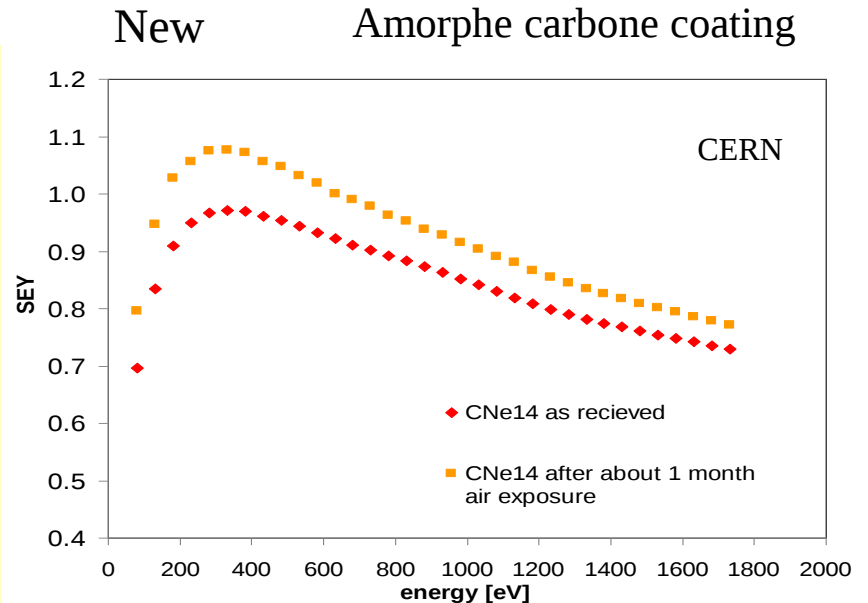
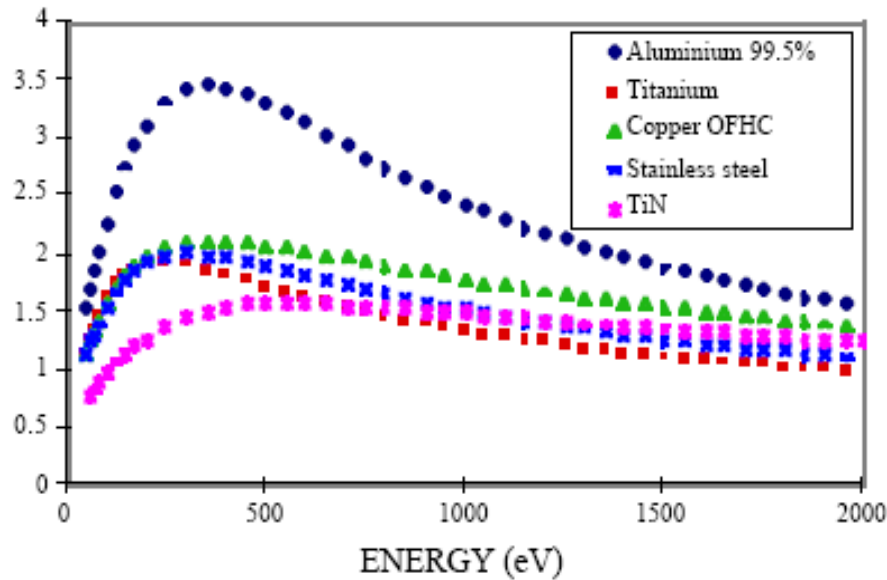


Scrubbing but needed time

Electron-Cloud / Electron Stimulated Desorption

SEY Reduction

Material and relief surface



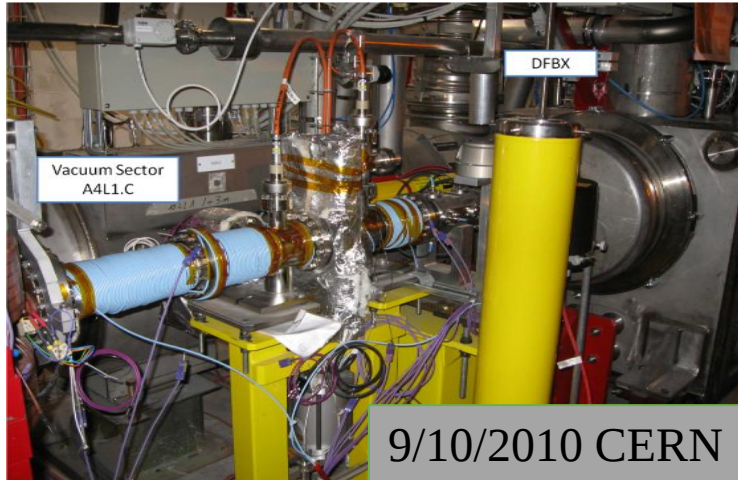
Coating TiN, TiZrV, amorphe C
coating on vacuum chamber

Isabel Montero, ICMM-CSIC (AEC 2012)

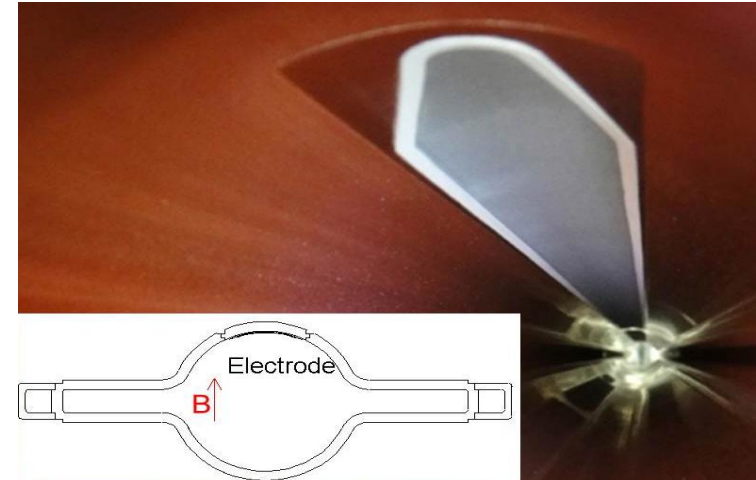
Electron-Cloud / Electron Stimulated Desorption

Other solution for e-Cloud decrease

External solenoid on chamber

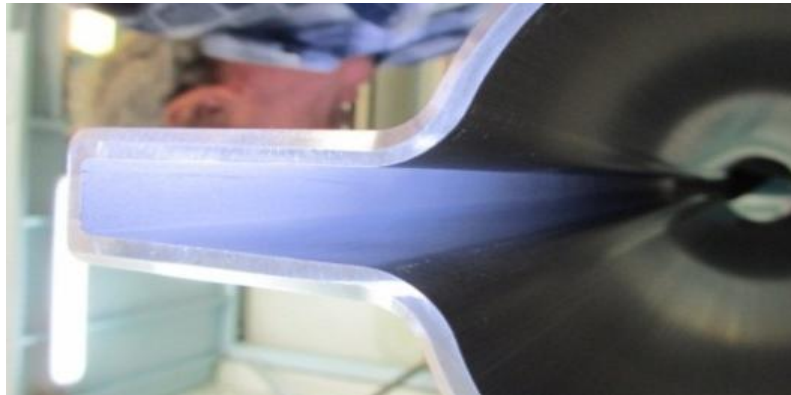


Clearing electrodes

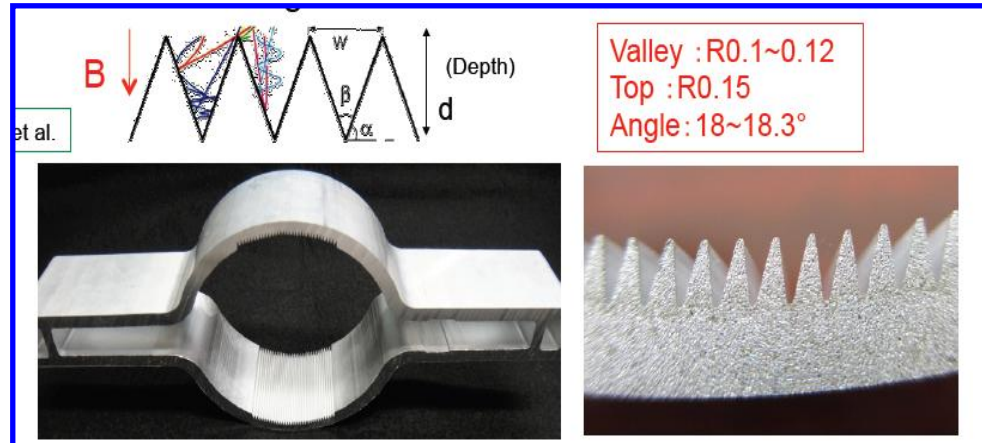


Y. Suetsugu SuperKEKB

Antichamber and Rough surface at the side wall (Ra ~20) reduces the photon reflection.



antichamber, grooves and coating TiN



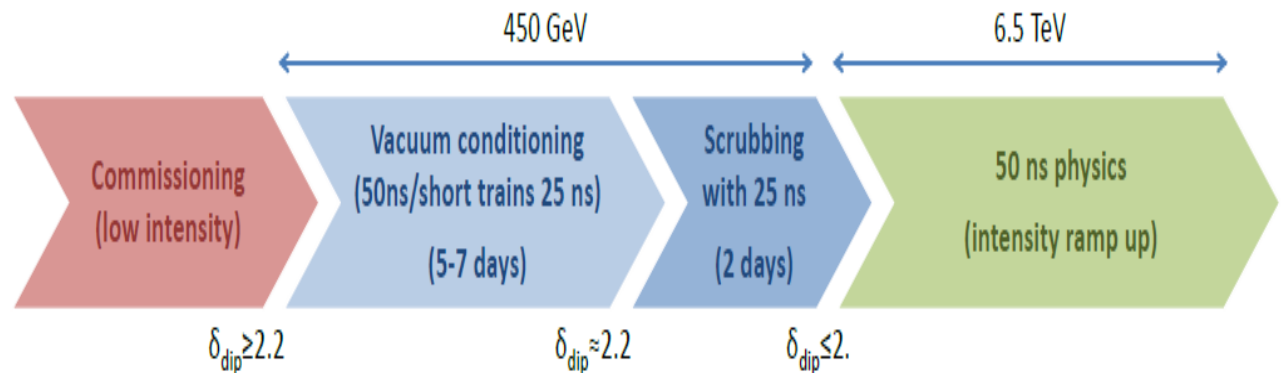
Y. Suetsugu SuperKEKB SuperKEKB Dipole Chamber Extrusion

Electron-Cloud / Electron Stimulated Desorption

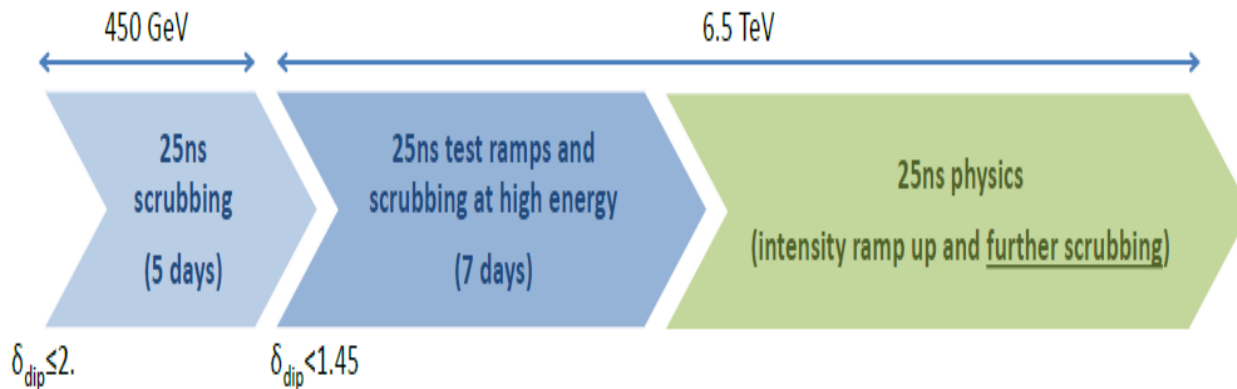
Scrubbing strategy for operation restart in 2015

Requirements for operation with 50 ns beams:

After the 2013-14 Long Shutdown, the Secondary Electron Yield of the beam screen in the arcs will likely be reset to values higher than 2.3, as it was before the 2011-2012 machine scrubbing.



Further requirements for operation with 25 ns beams:



According to the 2012 experience, the proposed scrubbing run(s) will not completely suppress the EC in the LHC. Further scrubbing will have to be achieved during the physics run implying a slow down of the intensity ramp-up process due to heat load, emittance blow up and poor lifetime.

Electron-Cloud / Electron Stimulated Desorption

Frank Zimmermann ECLLOUD'18

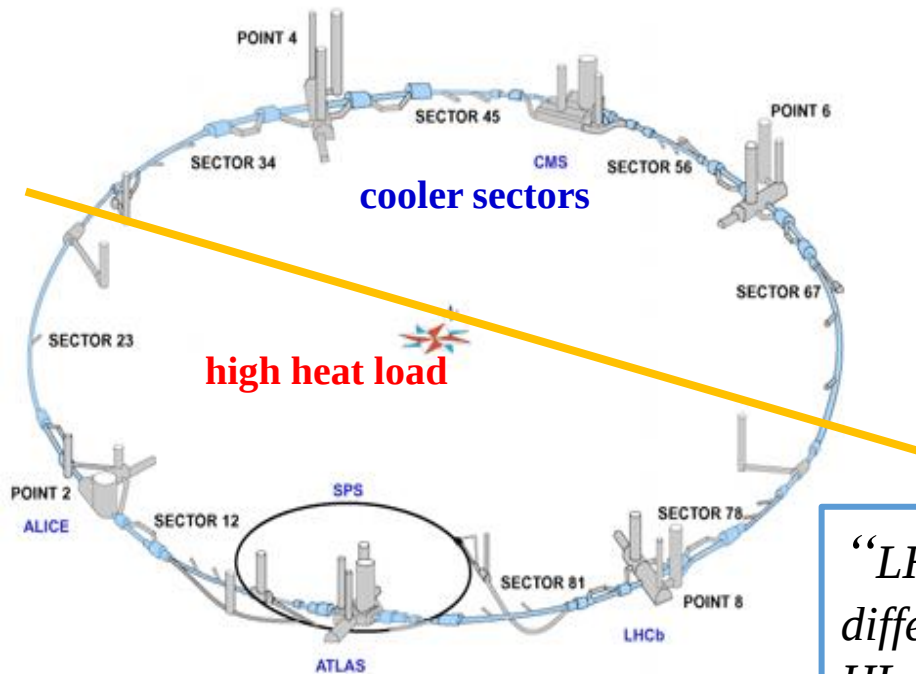
G. Iadarola, L. Tavian, et al,
2018

what is different in the high-load sectors?
lowest heat in sectors around CMS

arc heat loads already challenge for LHC Run 2 with 25 ns

- close to cryogenics limits in some of the arcs
- heat load much **larger** than expected from **impedance and synchrotron radiation**
- large differences observed between sectors
- **e-cloud** most plausible explanation

extrapolated HL-LHC heat loads for high-load sectors **not acceptable!**



highest heat in sectors
around ATLAS

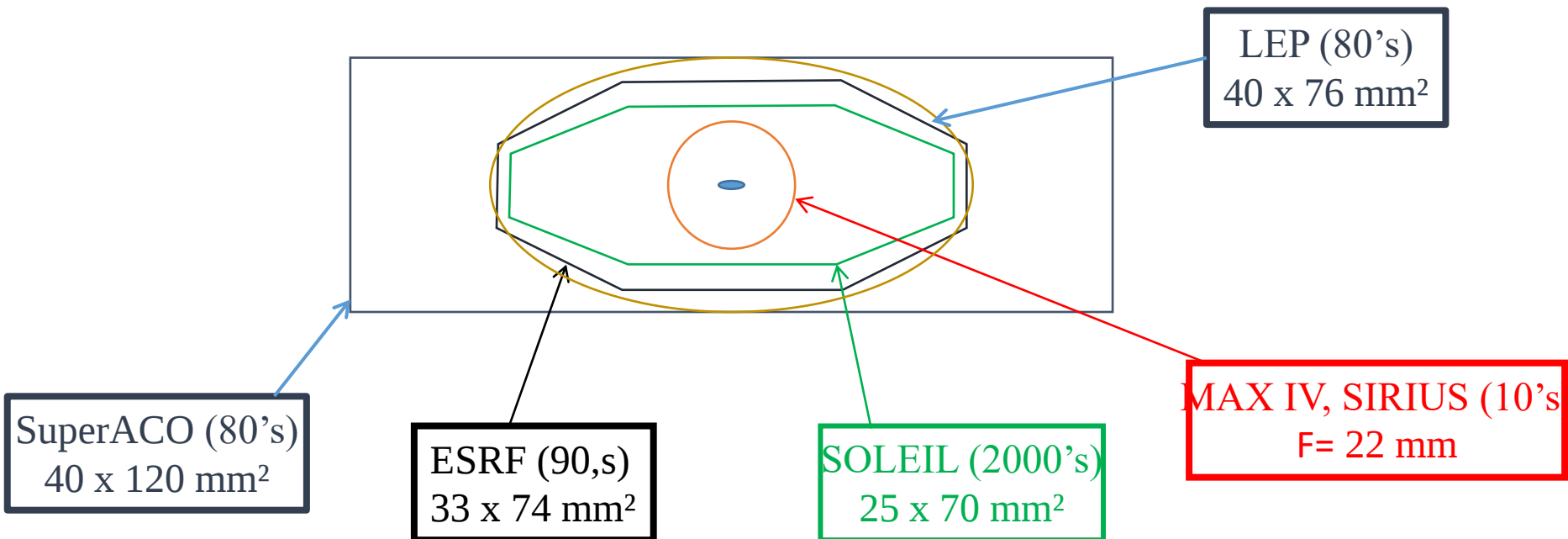
“LHC electron cloud is top priority. If e-cloud heat differences between sectors are not understood and HL-LHC not mastered, CERN may be shut down!”

Norbert Holtkamp, Chair CERN MAC, 28
January 2018

size evolution of vacuum chambers on Accelerators

Size vacuum chambers are constrained by the beam performance expected by the magnetic system including:

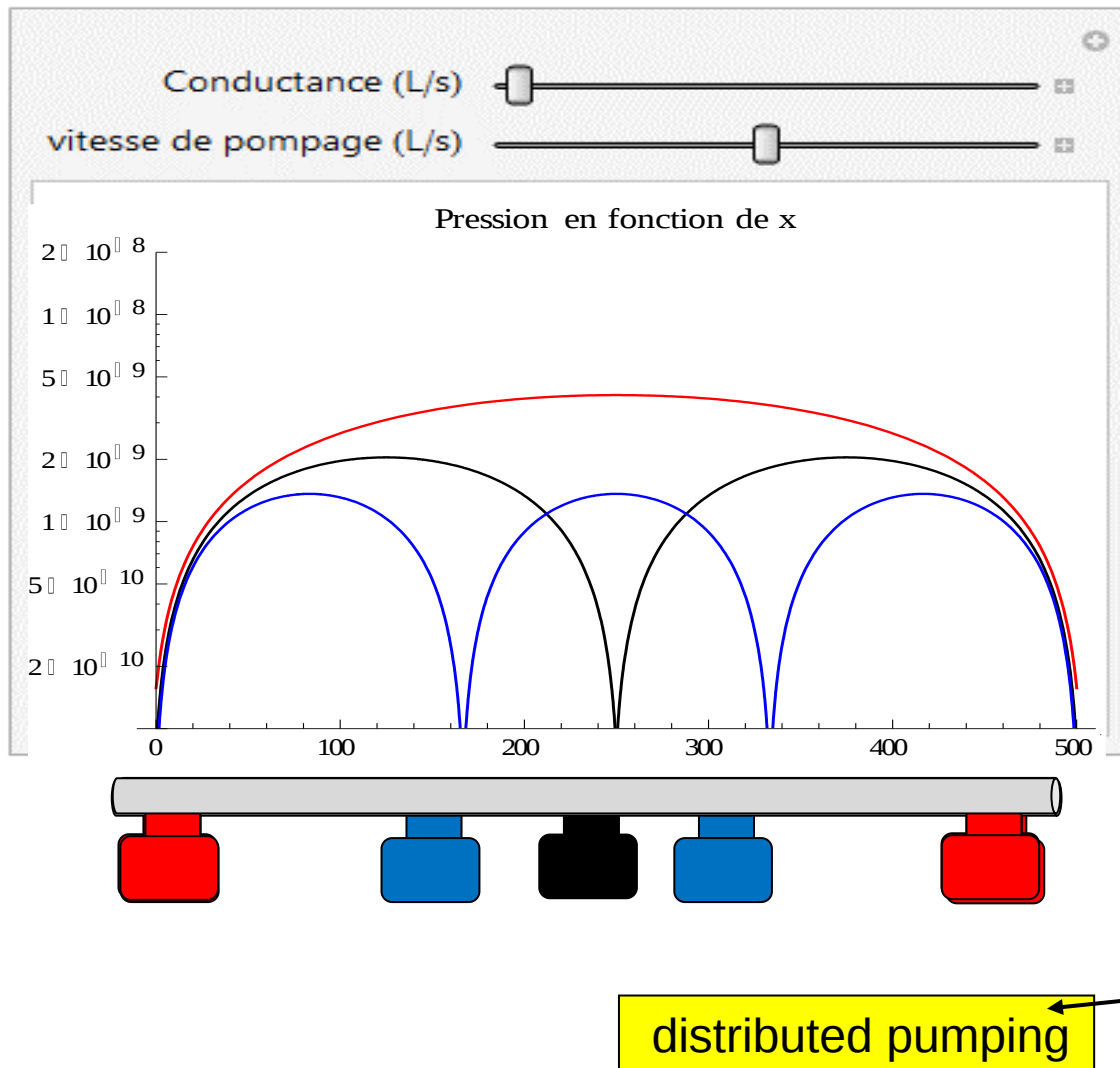
- Strong focus therefore reduces throat circle
- Low distance between the magnets.



Evolution vacuum technology

Consequences : size room more and more small

Significant decrease in conductance and increasing the maximum pressure



$$P_{mean} = A \cdot \tau \cdot L \left(\frac{1}{12C} + \frac{1}{S} \right)$$

$$P_{max} = A \cdot \tau \cdot L \left(\frac{1}{8C} + \frac{1}{S} \right)$$

τ outgassing rate C conductance

A circonference of chamber

if $S \gg 8C$ then $P_{max} \sim A\tau L / (8C)$

Decrease of P_{mean} and P_{max} is limited by the conductance C :

- Reduce outgassing rate (τ) and the stimulated desorption
- Reduce length L

Evolution vacuum technology

distributed pumping

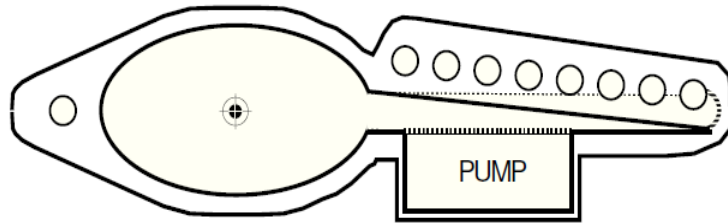
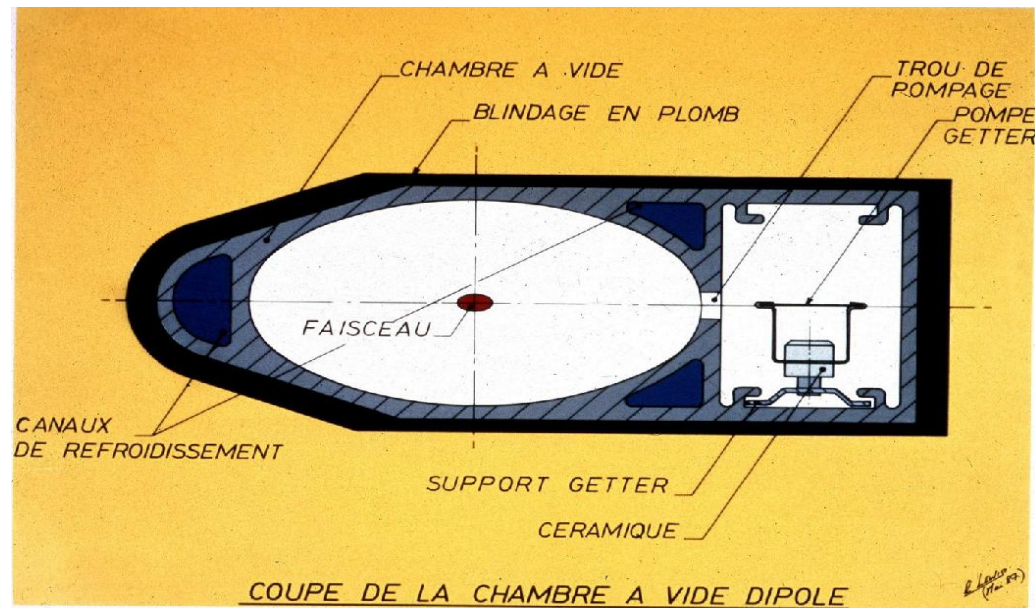
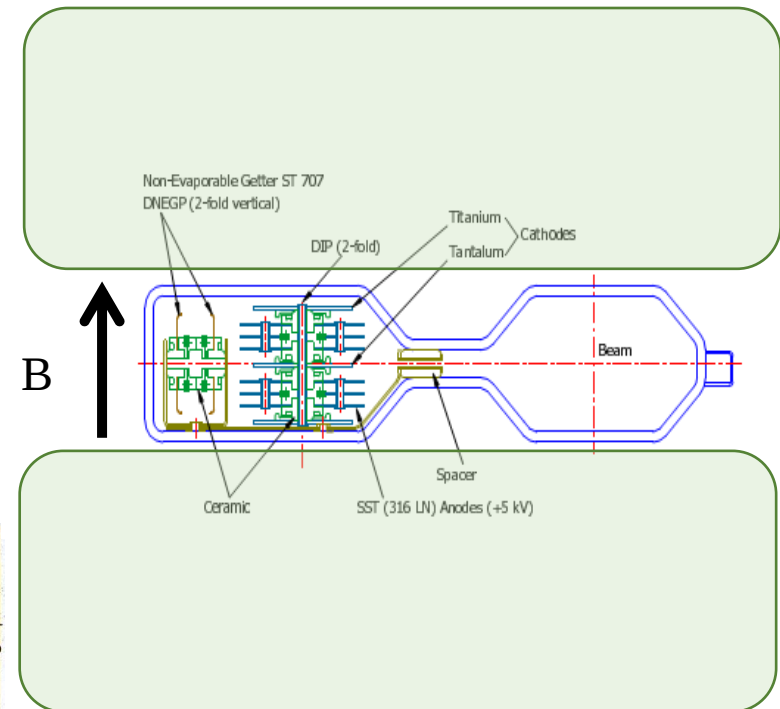
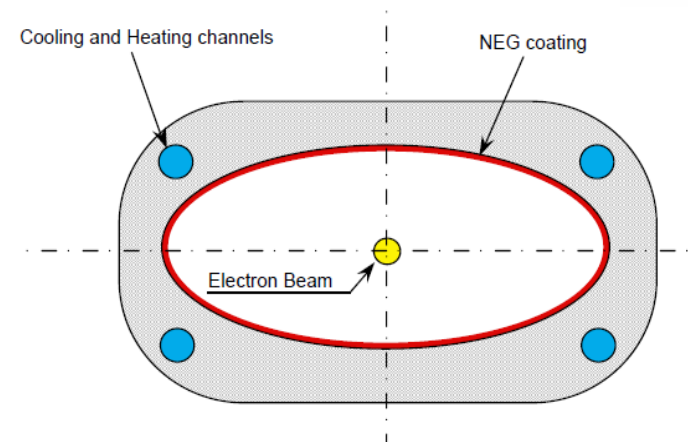


Figure 3-108. LER dipole vacuum chamber concept. Inside dimensions match those of Fig. 3-107.



LEP Design

(CERN LEP Vacuum group)



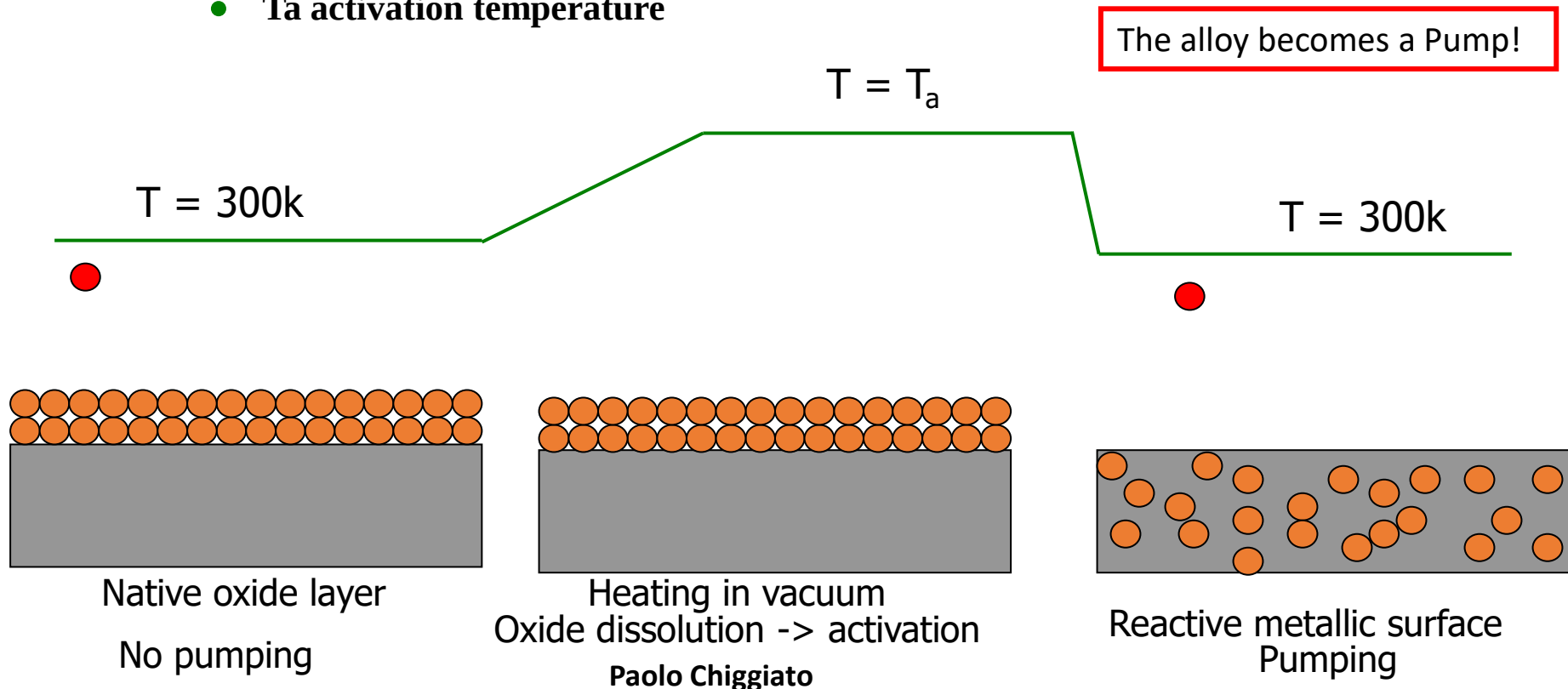
LHC

Non evaporable Getter (NEG)

Principle

coating of metal alloy (NEG) on the walls of vacuum chambers or on a strip

- coating Zr-Al, Zr-V-Fe, Ti-Zr-V,
- Ta activation temperature

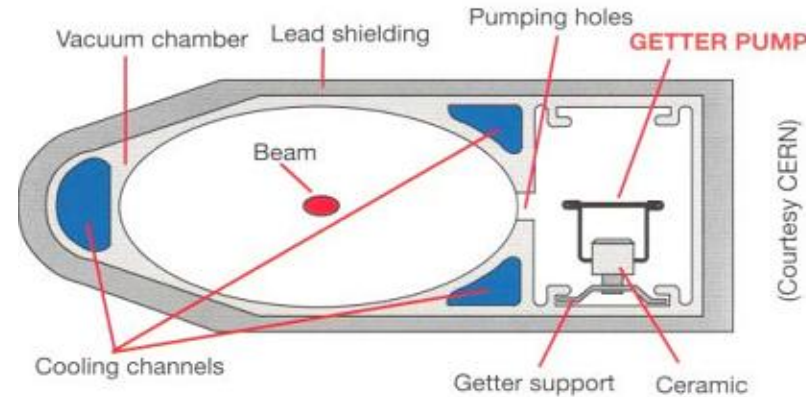


NEG strips

St101 NON-EVAPORABLE GETTER (Zr Al)
activation temperature at 700°C

**St707 NON-EVAPORABLE GETTER (Zr 70-V
24.6-Fe 5.4 wt%)**
activation temperature at 450°C

NEG materials are embedded in a constantan ribbon that is heated by the Joule effect



LEP – (CERN) 24 km de ruban installés
 $P < 10^{-11}$ mbar

The high temperature activation NEG can generate mechanical stresses in the vacuum chamber.

Need to have an antechamber

NEG strips

Design and construction of the SuperKEKB vacuum system

Yusuke Suetsugu,^{a)} Ken-ichi Kanazawa, Kyo Shibata, Takuya Ishibashi, Hiromi Hisamatsu, Mitsuru Shirai, and Shinji Terui
High Energy Accelerator Research Organization (KEK), Tsukuba, Ibaraki 305-0801, Japan

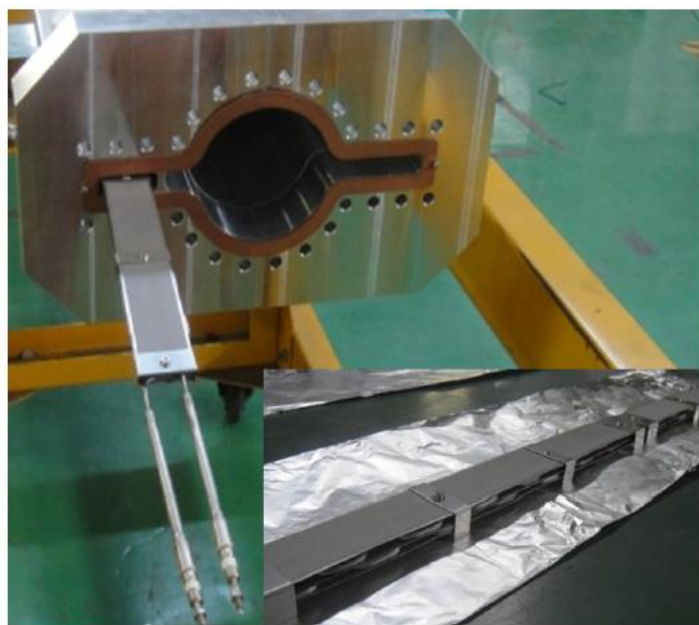


FIG. 4. (Color online) Multilayered NEG strips for the antechamber.

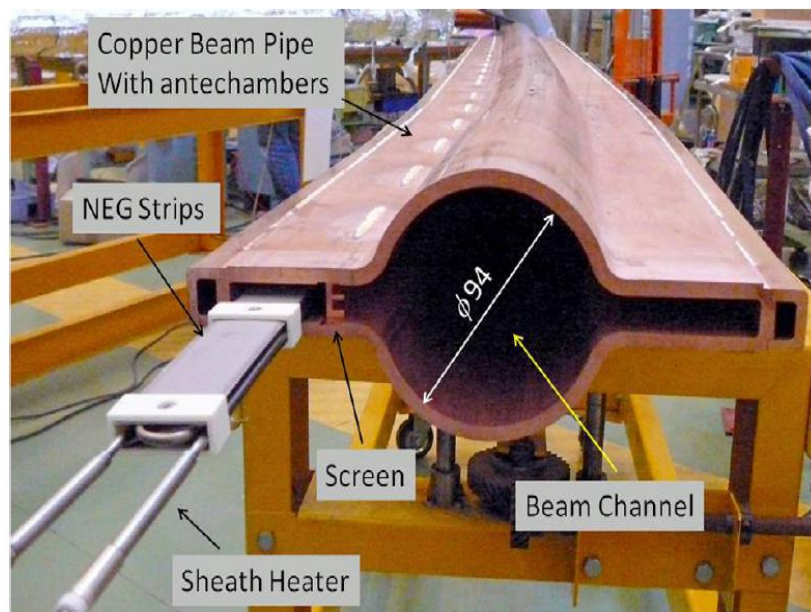


Fig. 1. Example of a beam pipe with antechambers and distributed pumps.

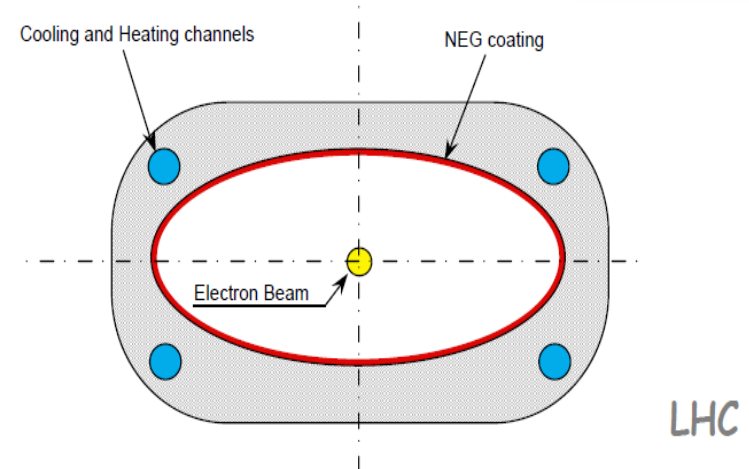
St707 NON-EVAPORABLE GETTER

J. Vac. Sci. Technol. A, Vol. 30, No. 3, May/June 2012

NEG coating in vacuum chamber

- Titane Zirconium Vanadium coating
- Low activation temperature $T_a = 180^\circ\text{C}$

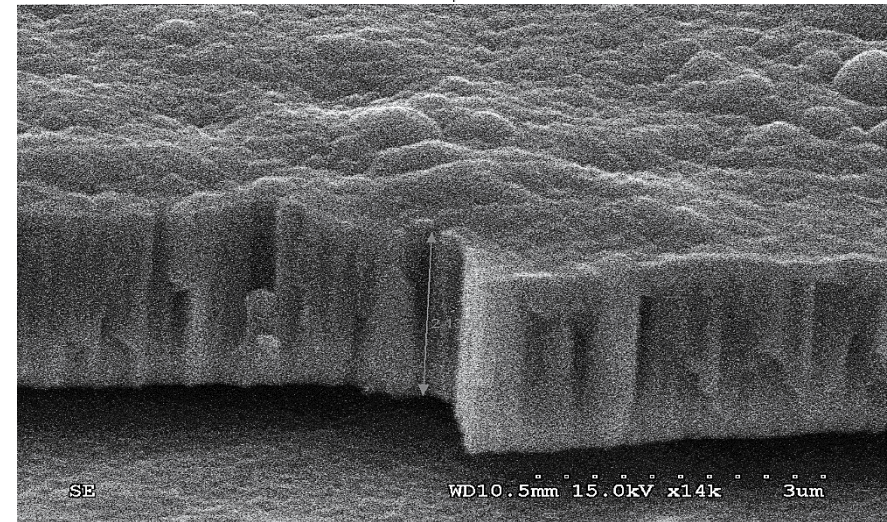
surface of vacuum chamber becomes a pump!!!



SOLEIL: NEG 60 % of the circumference
Upgrade : 100% of the ring!

LHC : 6 km in warm part

MAX IV: 95% of the circumference



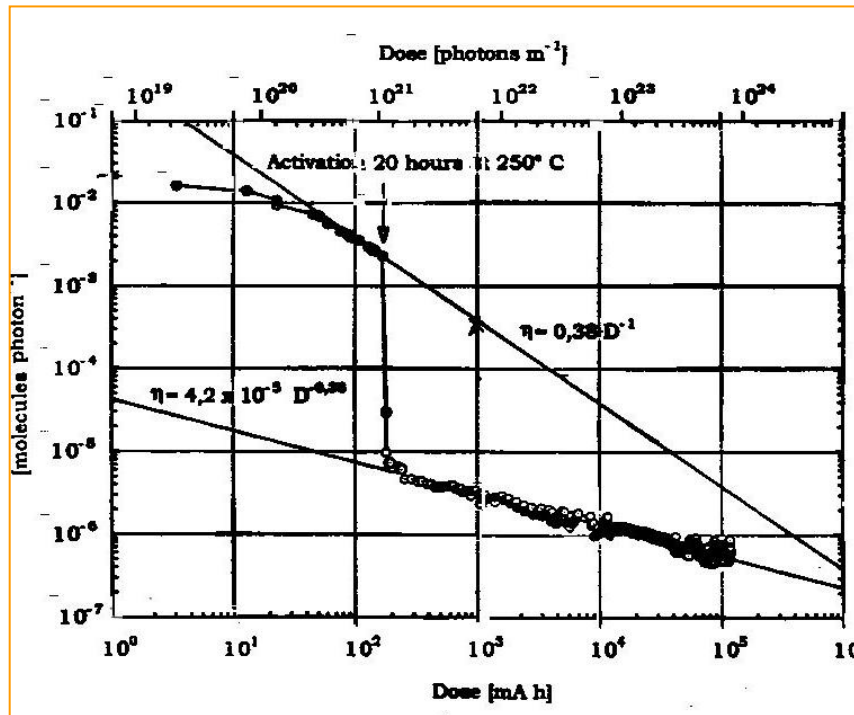
Coating at LAL ($\sim 2 \mu\text{m}$)

NEG coating-vacuum performances

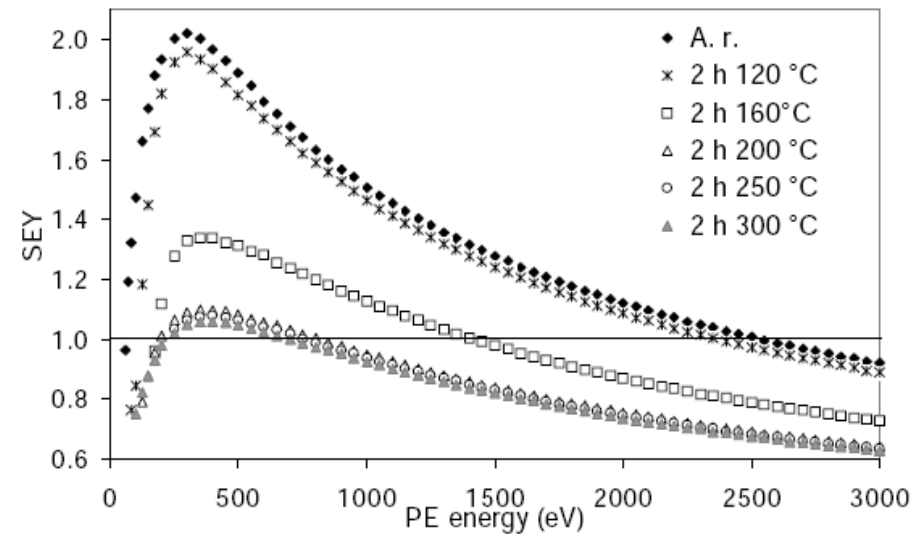
Titane ZirconiumVanadium coating

Electron and photon stimulated desorption photons (ESD and PSD) are reduced

A low secondary electron yield (SEY)



NEG SEY



C. Scheuerlein et al. Appl.Surf.Sci 172(2001)

P. Chiggiato, R. Kersevan / Vacuum 60 (2001) 67-72

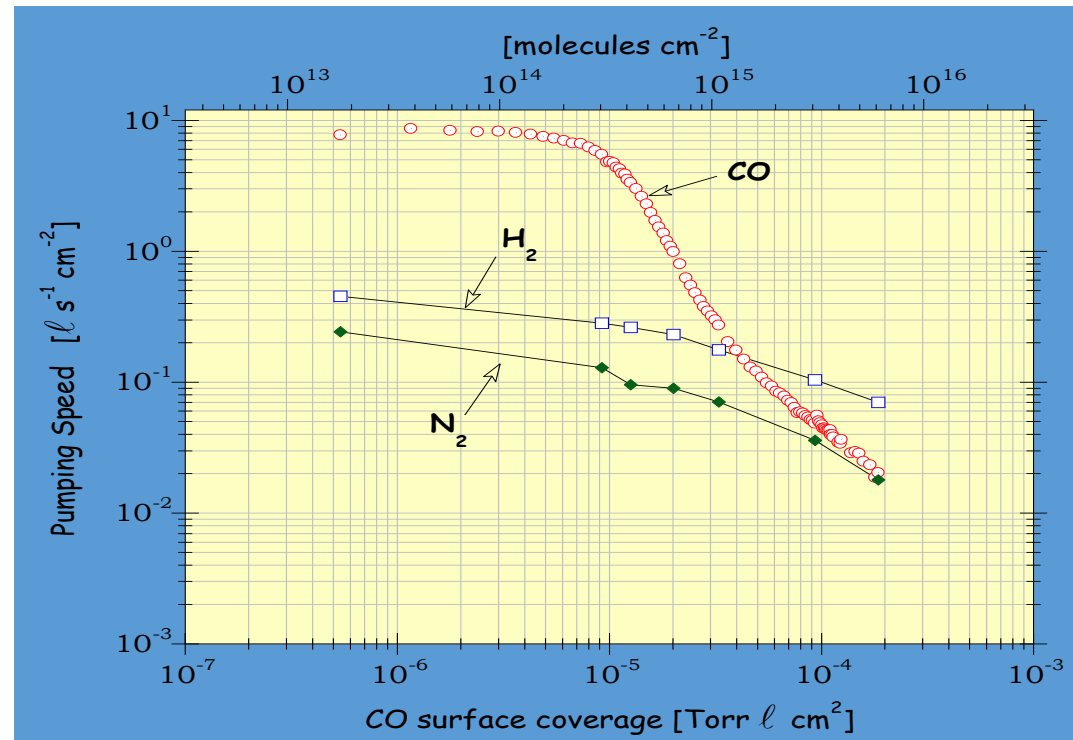
NEG coating-vacuum disadvantage

Titane Zirconium Vanadium coating

Selective pumping speed

(No pumping of rare gas, CH_4)

saturation effect



P. Chiggiato, R. Kersevan / Vacuum 60 (2001) 67-72

Conclusion

Trend: High luminosity accelerators (involves small-sized beams, high current and a high repetition frequency).

Many beams interactions, residual gas and surface will generate induced outgassing sources and impact the shape and the beam energy

Reduction of induced outgassing sources

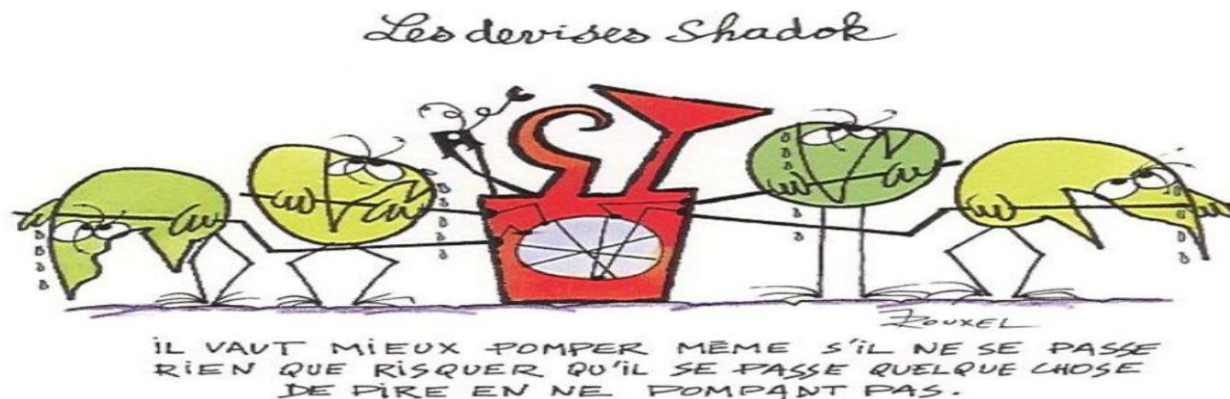
Geometry, coating, surface condition, **scrubbing**,.....

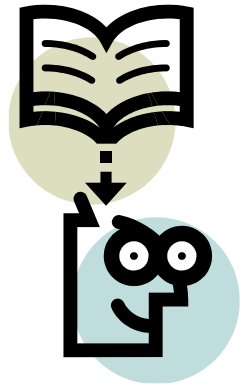
Increase the effective pumping speed

Distributed pumping

NEG strips

Neg coating (TiZrV)





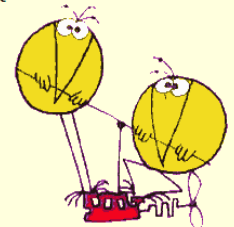
Some references

- Cern Accelerator School, Vacuum technology, CERN 99-05
- Cern Accelerator School, Vacuum in accelerators, CERN 2007-03
- Scientific foundations of vacuum technique, S. Dushman, J.M Lafferty. J. Wiley & sons. Elsevier Science.
- Vacuum Technology, A. Roth. Elsevier Science.
- Jacow
- L'ultra vide des grands instruments – V. Baglin Gif-sur-Yvette - 22 septembre 2009

Journals

- Journal of vacuum science and technology
- Vacuum
- Physical review special topics Accelerators and Beam

C'EST POURQUOI TOUTE LA POPULATION SHADOK.....



POMPAIT, POMPAIT, POMPAIT....

D'où, un des fameux adages Shadok :

Je pompe donc je suis

Interaction: Beam- residual gas

Elastic interaction { Single Coulomb scattering

$$\frac{d\varepsilon}{dt} \propto Z^2 \cdot P$$

ε Emittance is proportional to residual gas pressure and depends square residual gas atomic number (Z)

Diffusion sur les noyaux du gaz résiduel

$$\frac{1}{\tau_{DN}} := \frac{2\pi r_0^2}{\gamma^2} Z^2 \rho c \left[\frac{\langle \beta_z \rangle \beta_{zmax}}{a^2} + \frac{\langle \beta_x \rangle \beta_{xmax}}{b^2} \right]$$

β_z (m) oscillation betatron verticale

a (m) demi ouverture verticale

β_z max(m) oscillation betatron verticale

b (m) demi ouverture horizontale

β_x (m) oscillation betatron horizontale

Z n° atomique moyen du gaz

β_x max(m) oscillation betatron horizontale

ρ densité des noyaux du gaz

γ Facteur gamma

Nuclear scattering LHC

Life time limit ~100h P~10-8 mbar H₂ equivalent E=7 Tev

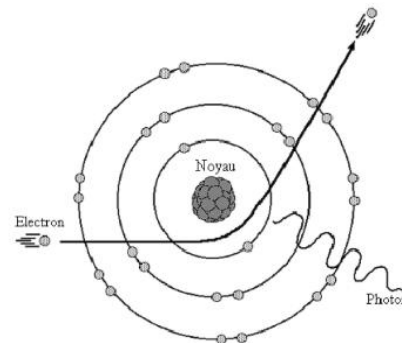
Synchrotron Soleil

Life time limit ~60h P~10-9 mbar N₂ equivalent E=2,15 Gev

Interaction: Beam- residual gas

Bremsstrahlung

Material	Z	A	Radiation length X_0 [g/cm ²] [cm]	Density [g/cm ³] () is for gas [g/l]
H ₂	1	1.01	61.28 865	0.0708(0.090)
D ₂	1	2.01	122.6 757	0.162(0.177)
He	2	4.00	94.32 755	0.125(0.178)
Li	3	6.94	82.76 155	0.534
Be	4	9.01	65.19 35.3	1.848



→ Inelastic interaction

Bremsstrahlung

Ionization energy loss
Electron capture
Electron loss
Nuclear reactions

$$\frac{dE}{E} = -\frac{dx}{\chi_0}$$

Valid for light particles (e⁺, e⁻)

$$\chi_{0[g/cm^2]} = \frac{714.6A_c}{Z_c(Z_c + 1) \ln(287/\sqrt{Z_c})}$$

➤ Loss beam energy

Synchrotron Soleil

$$\frac{1}{\tau_{BN}} := \frac{16}{411} r_0^2 Z^2 \rho c \left[\ln \left(\frac{183}{Z^{1/3}} \right) \left(-\ln(\epsilon_{RF}) - \frac{5}{8} \right) + \frac{1}{12} [-\ln(\epsilon_{RF}) - 1] \right]$$

Life time limit ~60h

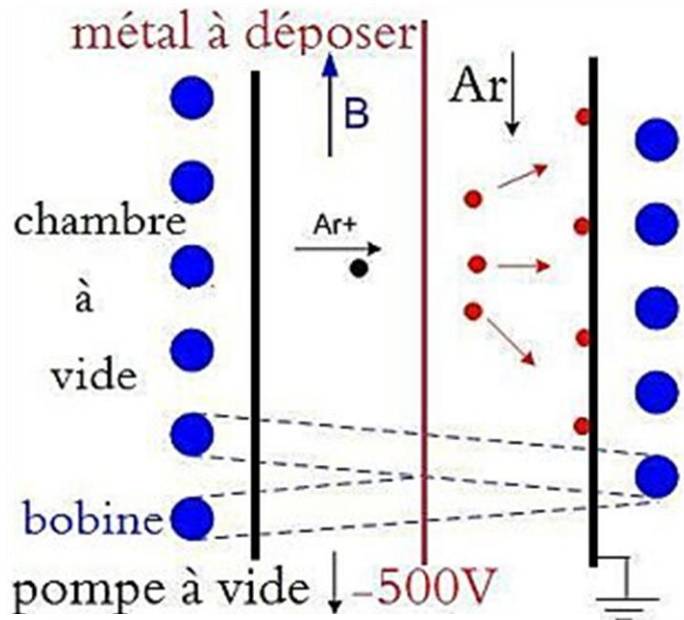
P~10-9 mbar N₂eq. E=2,15 GeV RF acceptance $\epsilon_{RF}=1,8 \%$

➤ the radiation interacts with the surface → Desorption

A very strong dependence on the target atomic number is observed, which means that even a small amount of a heavy gas may have a higher influence on the lifetime than the usually dominating hydrogen content.
about 4.10⁶ times lower for proton

Le dépôt NEG

La pulvérisation cathodique magnétron



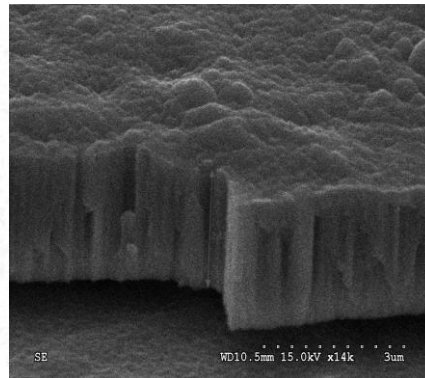
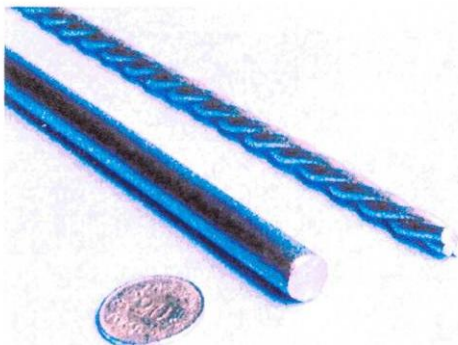
Paramètres de dépôt

$B = 100 \text{ G}$

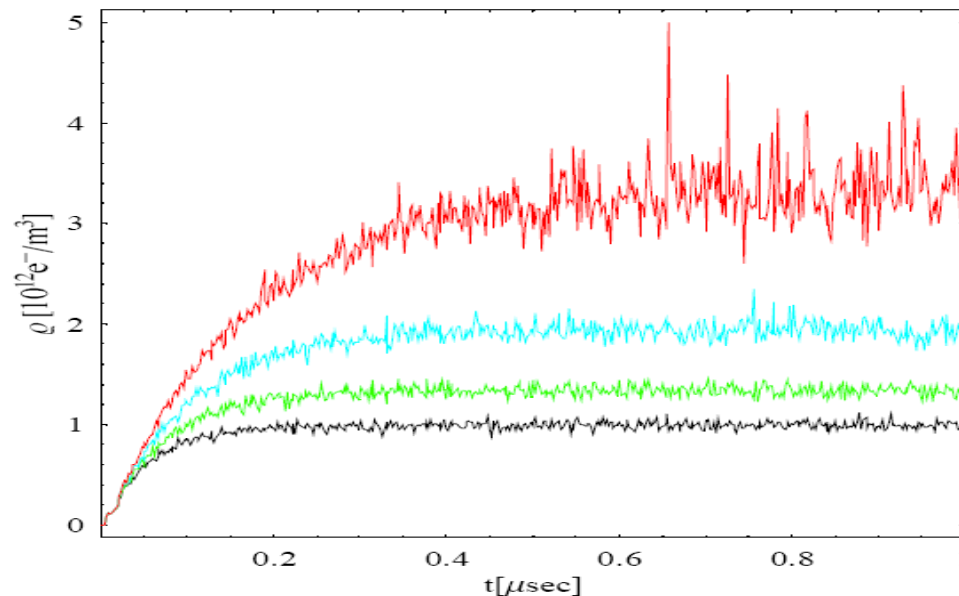
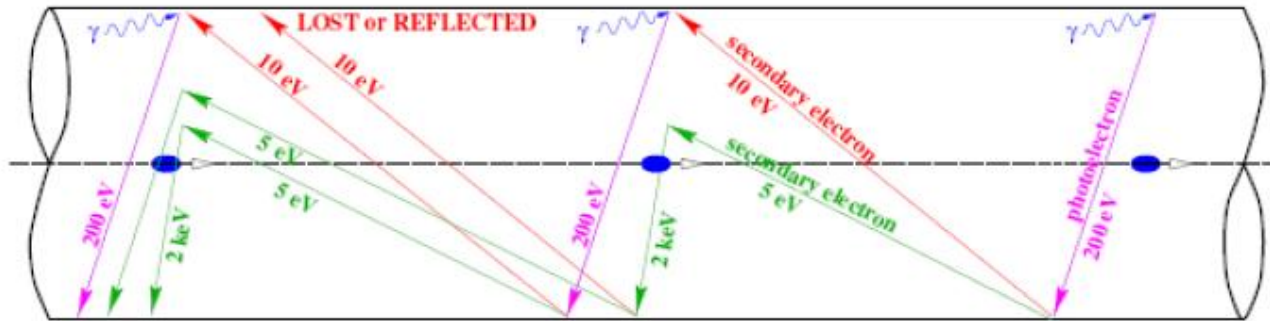
$V = - 500 \text{ V}$

Pression Krypton : $2 \cdot 10^{-2} \text{ mbar}$

Vitesse de dépôt : 1 Å/s



Electron Cloud Effects



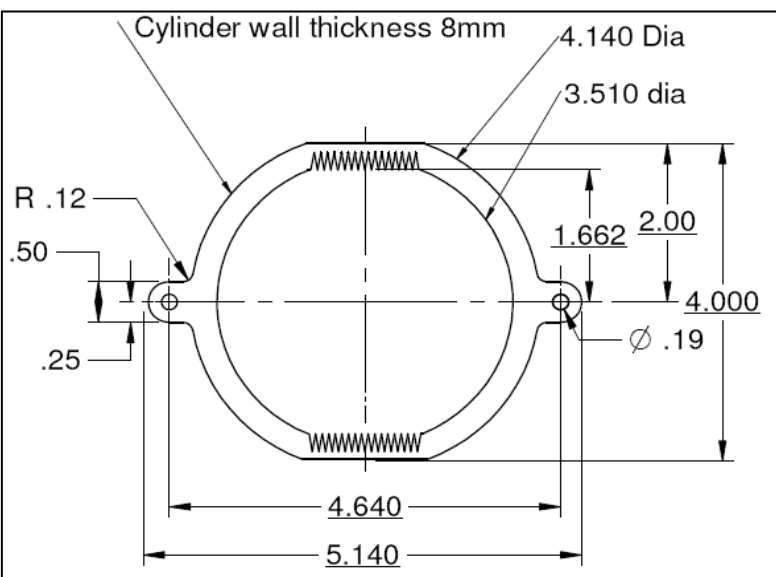
95% prot.

$\delta_{\max}=1$

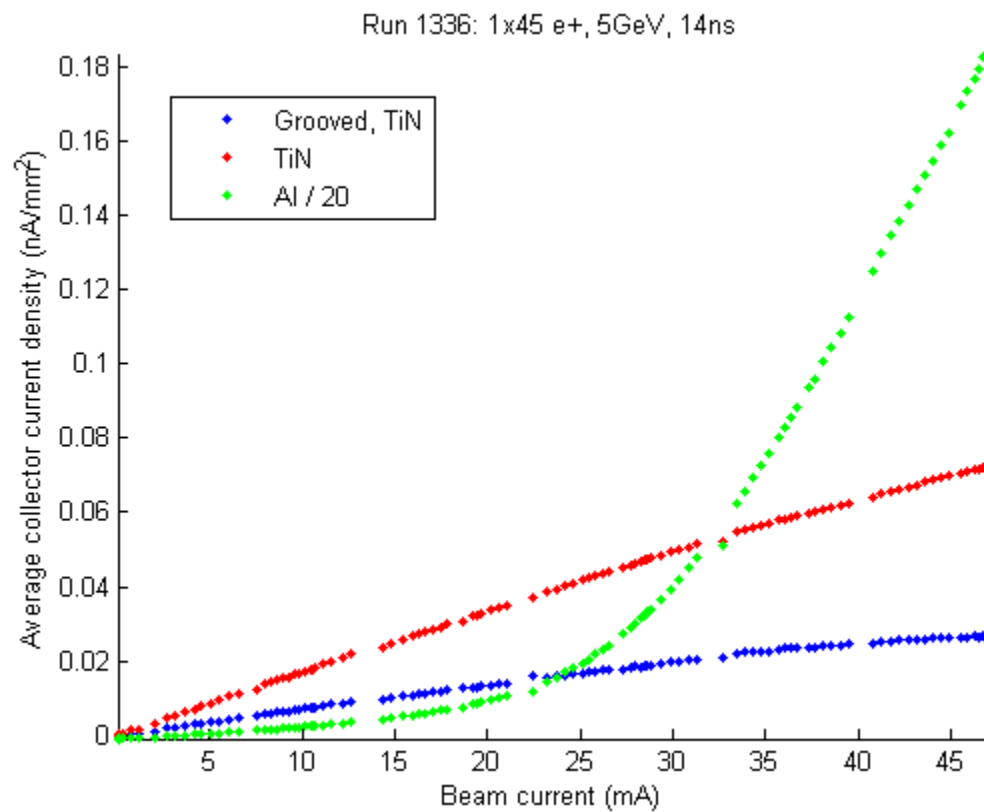
$\delta_{\max}=1.1$

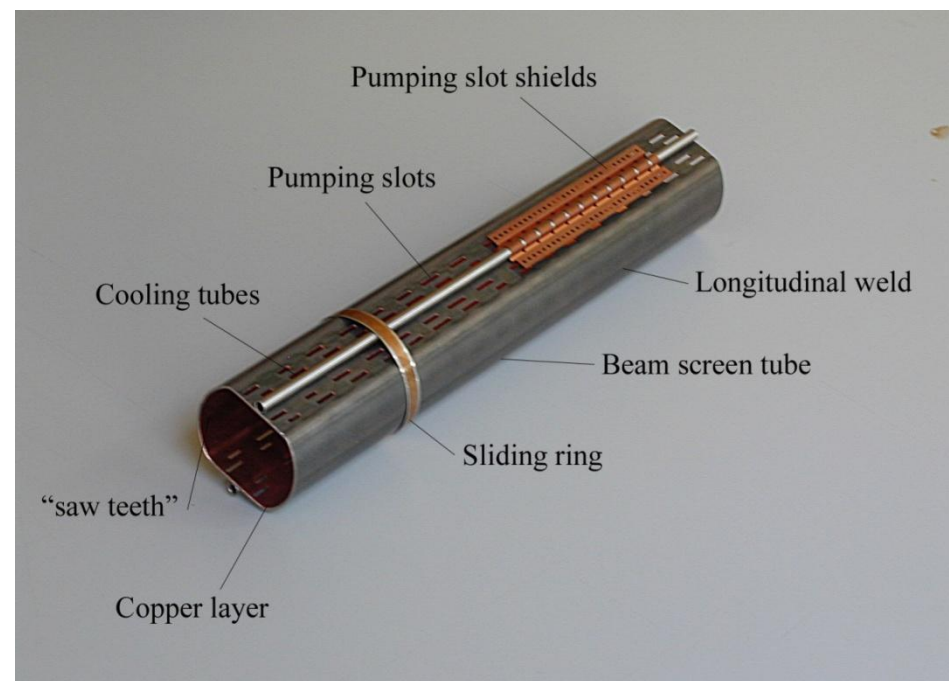
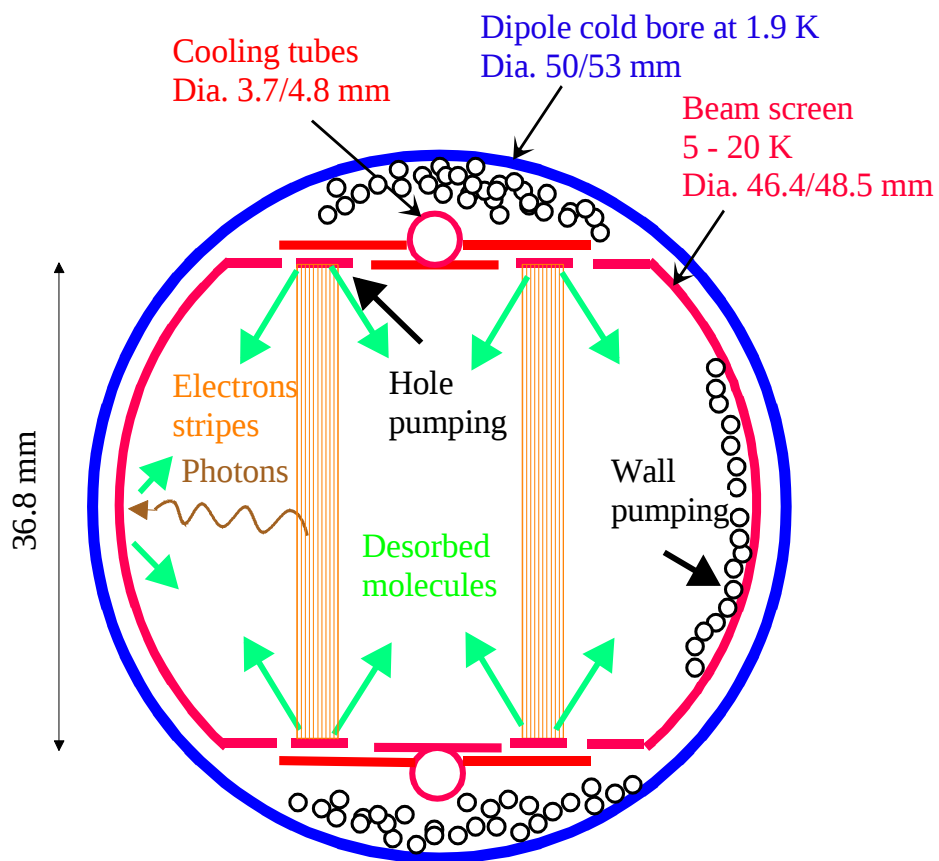
$\delta_{\max}=1.2$

$\delta_{\max}=1.3$



M. Pivi and L. Wang design, SLAC





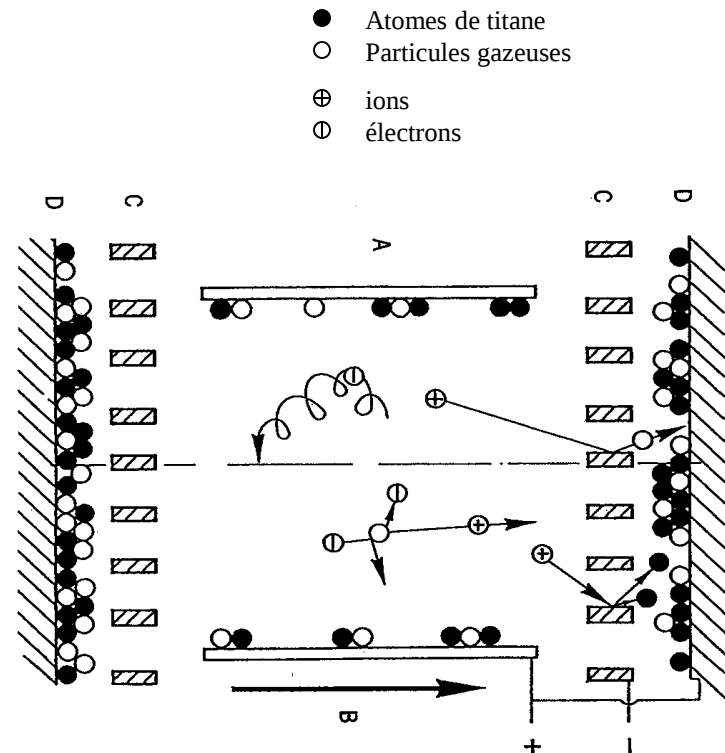
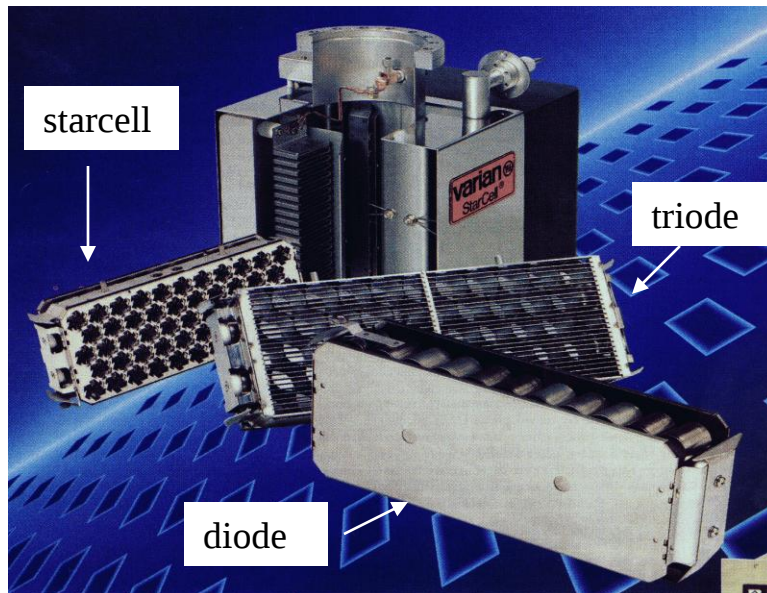
POMPES IONIQUES

Trois phénomènes de pompage :

- par chimisorption → effet Getter (gaz actifs, CO, CO₂, H₂, N₂, O₂,...)
- par diffusion (principalement H₂)
- par physisorption → enterrement (unique pompage pour les gaz rares)

Principalement deux configurations:

- Pompes ioniques Diodes (et noble diode)
- Pompes ioniques triodes (et starcell)



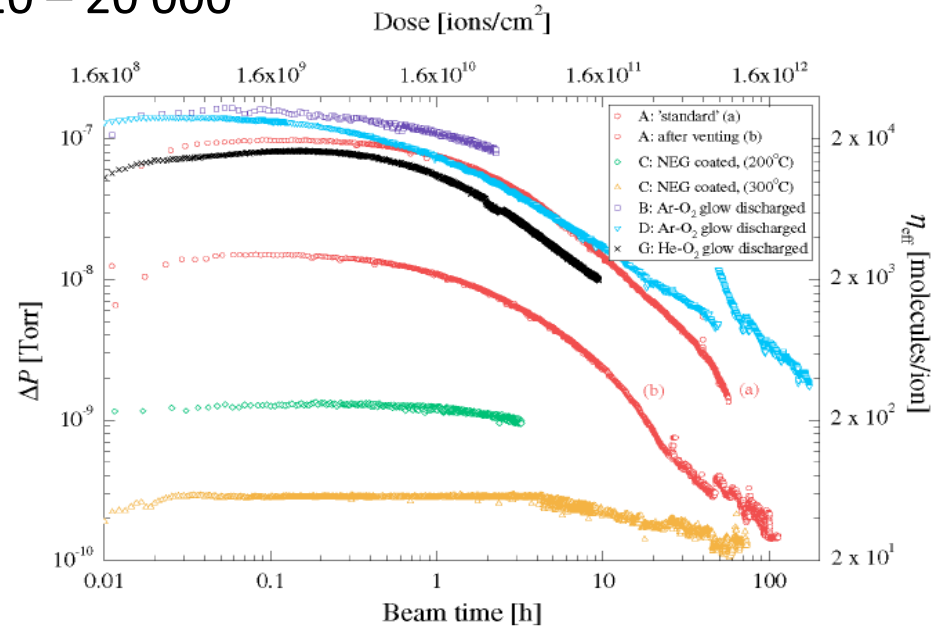
Pompe ionique triode

Ions à très haute énergie

- Des taux de desorption de 20 – 20 000

Pb^{53+} , 4.2 MeV/u, 14 mrad

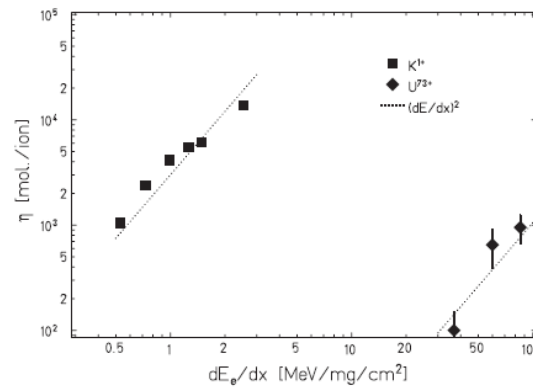
E. Mahner *et al.*, Phys. Rev. ST Accel. Beams 6, 013201 (2003)



- La désorption est déterminée par l' énergie donnée aux électrons (pouvoir d'arrêt électronique)

$$\eta_{\text{ion}} \propto \left(\frac{dE_e}{dx} \right)^2$$

L. Prost *et al.*, PRL 98, 064801 (2007)



- La désorption induite par les électrons est le mécanisme responsable

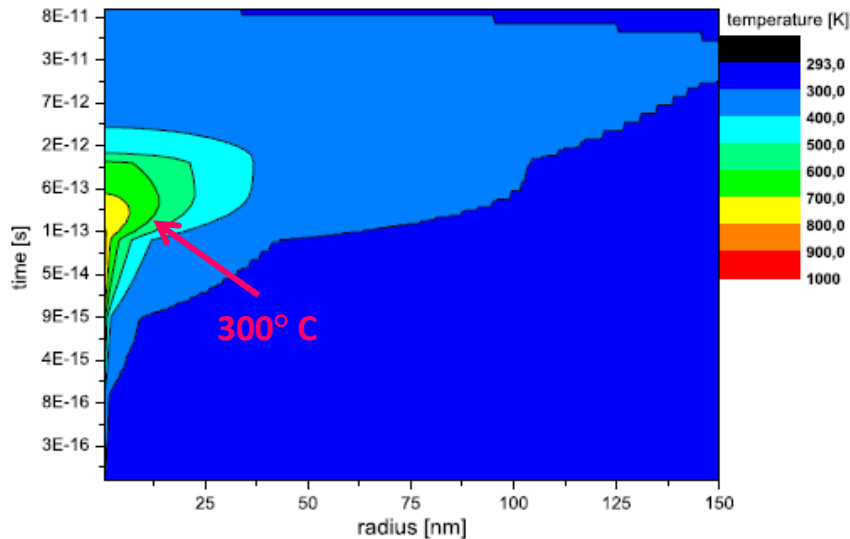
Mécanisme

- Effet de surface (sauf diffusion H₂) dû à une activation thermique
- « Inelastic thermal spike model » : carte de température couplée au modèle de désorption thermique

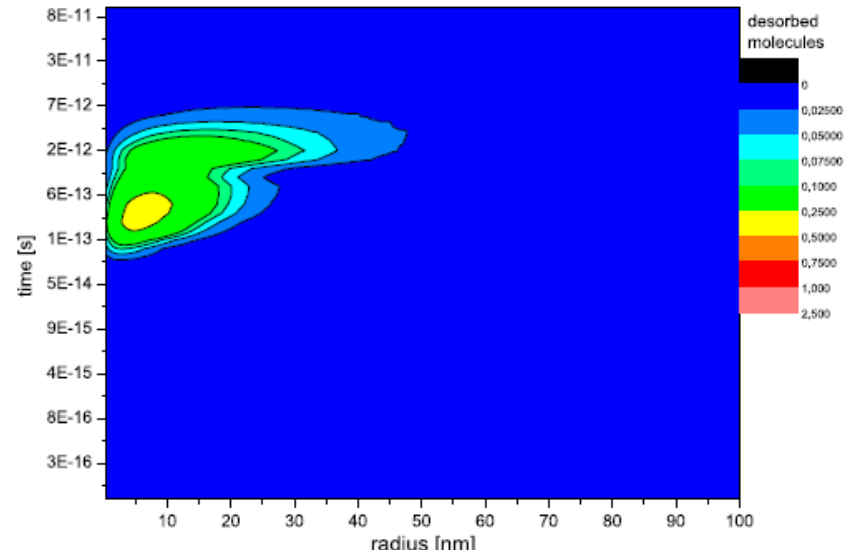
$$\eta = \int_0^{t_{\max}} \int_0^{r_{\max}} v_0(T(r, t)) \cdot \tilde{n}(r, t) \cdot \exp\left(-\frac{E_{\text{des}}}{k_B \cdot T(r, t)}\right) \cdot 2\pi \cdot r dr dt,$$

M. Bender *et al.*, NIM B 267 (2007) 885-890

Xe²⁹⁺, 1.4 MeV/u, Perpendiculaire



Temperature of atomic Cu subsystem after Xe impact



Desorbed particles per Xe per dt

$$\eta_{\text{calculated}} = 185$$

Remèdes

- Utilisation de revêtements NEG (LEIR, RHIC, GSI)
- Intercepter les ions « perdus » sur des collimateurs dédiés :
 - LEIR : plaquage d'or de 30 μm sur de l'acier 316 LN, incidence perpendiculaire
 - GSI : plaquage d'or de 0.1 μm , incidence perpendiculaire. Absorbeur introduit dans une chambre secondaire. Revêtement NEG

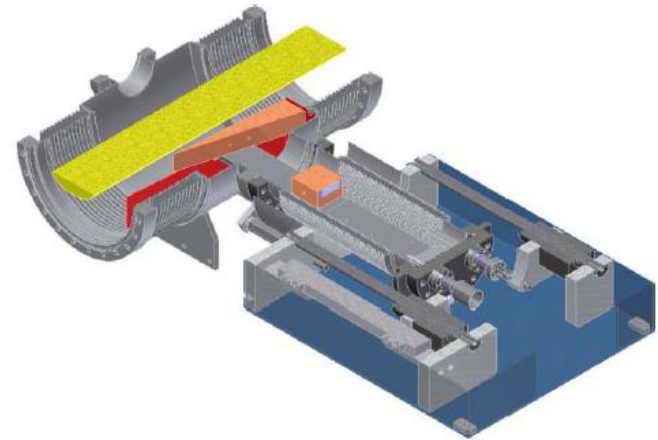
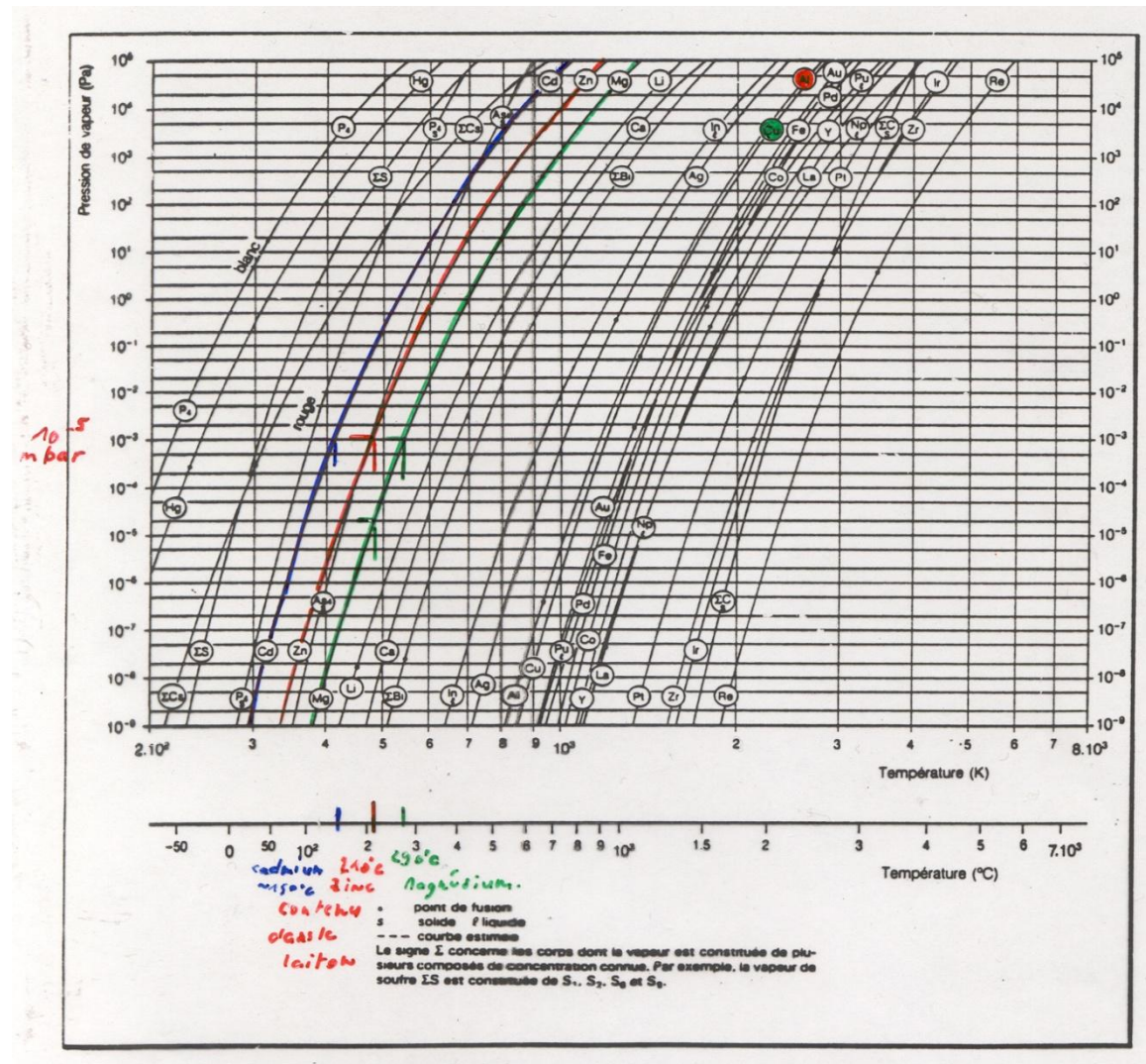


Figure 2: Horizontal cut through the installed SIS18 ion catcher prototype. Yellow: beam, red: secondary chamber, brown: beam absorbers.

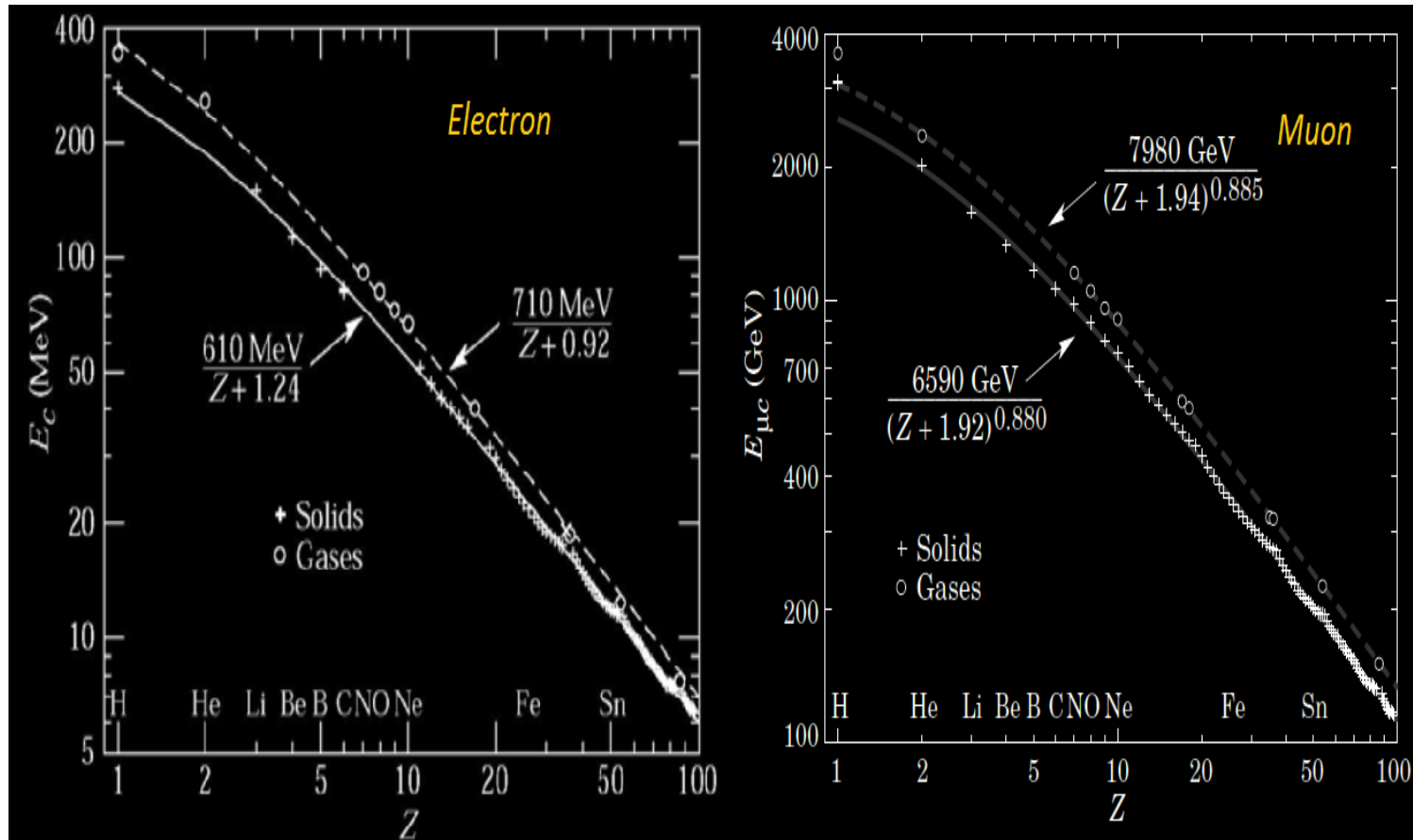
C. Omet *et. al.* EPAC 2008, Genoa, Italy

- Utiliser le conditionnement par le faisceau

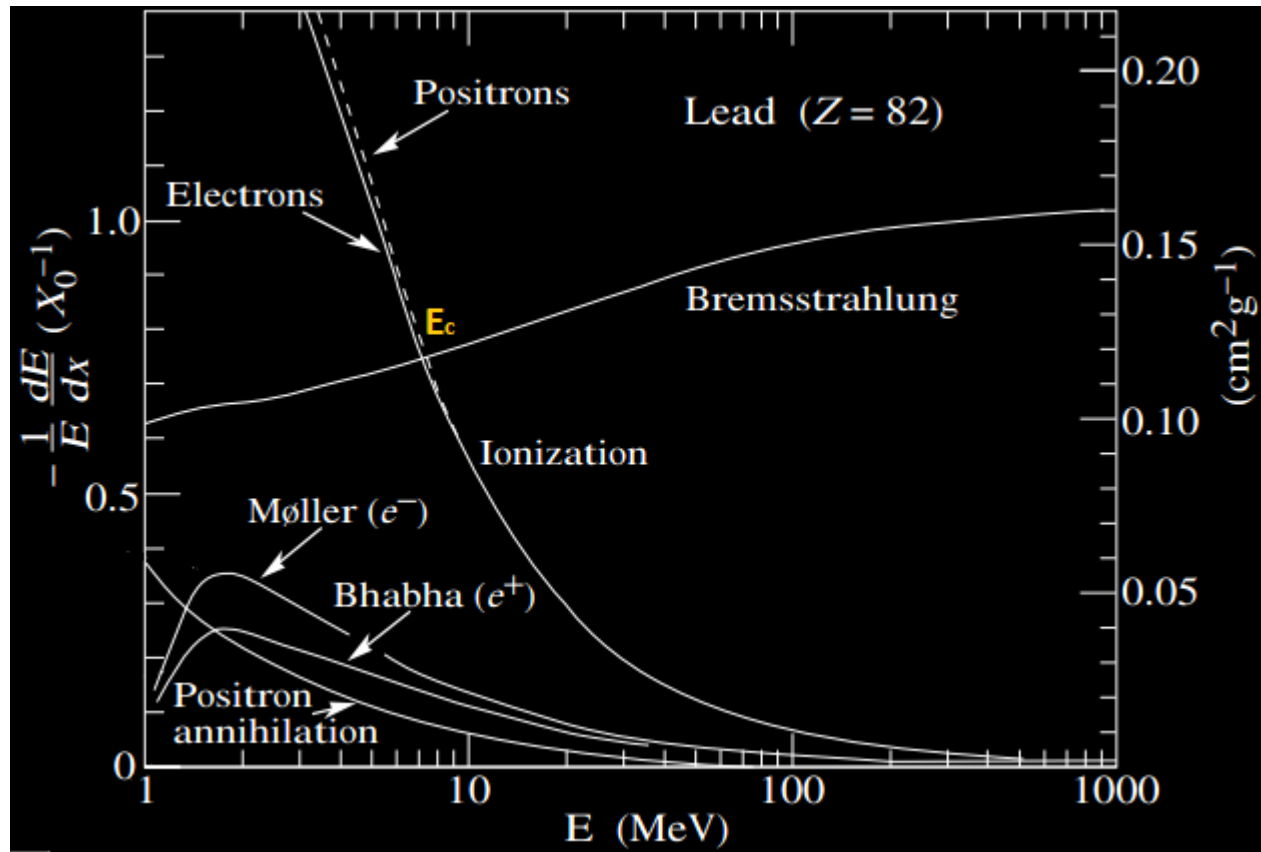
Vapor pressure



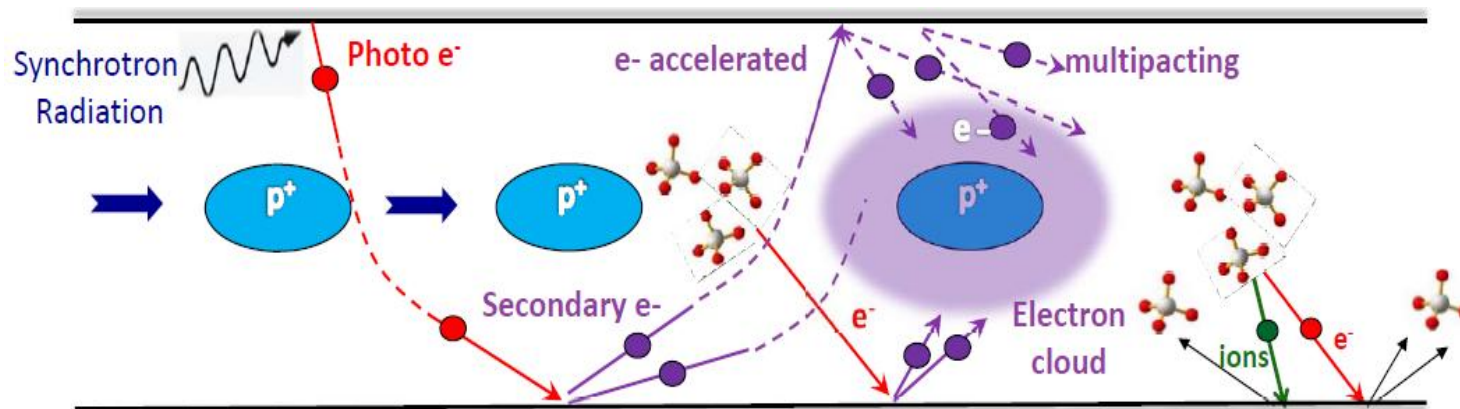
Energie critique 'collision élastique/ émission radiative) en fonction de Z



Perte d'énergie des électrons dans du plomb par unité de radiation X0.
L'énergie critique est de [7,43 MeV](#).



Plomb 11,35 g/cm³ (20°C)



2018 LHC
26.7 km, 7.74 T
13 TeV, 2556 b, 1.1×10^{11} ppb
RUN 2

LHC
27 km, 8.33 T
14 TeV, 2808 b, 1.2×10^{11} ppb
NOMINAL PARAMETERS

HL-LHC (2026)
26.7 km, 8.33 T
14 TeV, 2808 b, 2.2×10^{11} ppb

FCC-hh
100 km, 16 T
100 TeV, 10600 b

FCC-ee
100 km
45 GeV, 16640 b

To stay operational during 2018, the LHC was limited to fewer than the nominal number of bunches due to **dynamic pressure phenomena**

Can these nominal parameters be reached for future colliders ?

Ion accumulation

Ions stability criterion :

Calculation of critical mass ¹ :

All ions of mass greater than the critical mass will be accumulated.

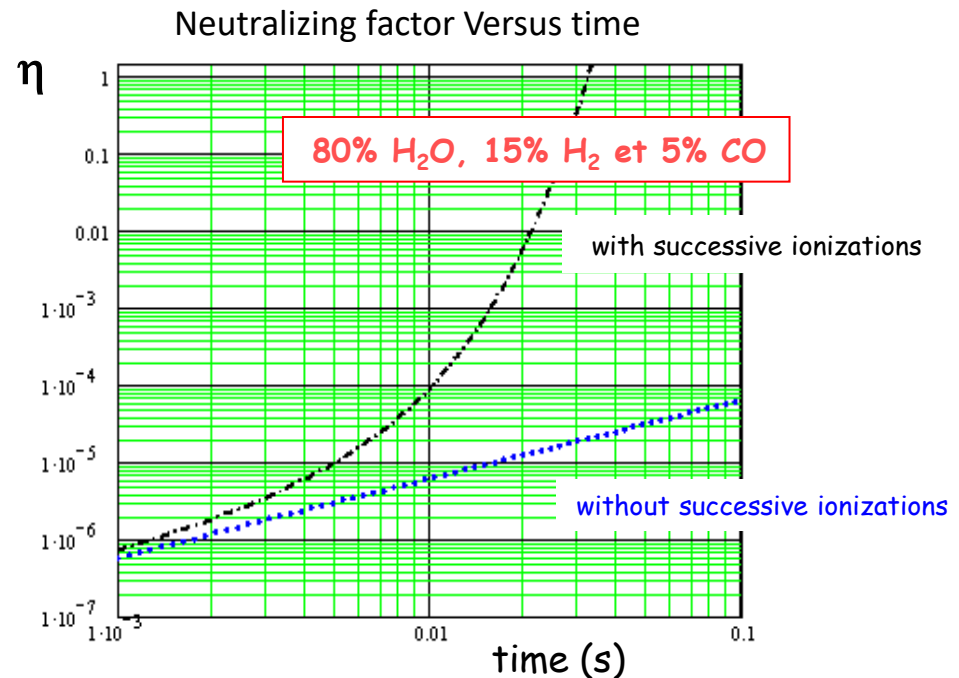
$$A_c = \frac{N_e}{N_b^2} \cdot r_p \cdot \frac{2 \cdot \pi \cdot R}{\beta \cdot b^2 \cdot (1 + a/b)}$$

avec N_e e- number, N_b bunch number, r_p proton radius,
 $2 \cdot \pi \cdot R$ ring circumference, β ratio of v to the speed of light c , a and b beam size .

[1] Neutralisation of accelerator beams by ionisation of the residual gas Y. Baconnier, A. Poncet and P.F. Tavares CERN

Neutralizing Factor

$$\eta = \frac{\text{number of positive charge (ions)}}{\text{Number of negative charge (beam electron)}}$$



summary Tables

Pour N₂ at 20°C

Pressure (hPa)	Strength per unit area	density molecular (molecule/cm ³)	Mean free path	Impingement rate (per cm ² and per sec)	Monolayer formation time
1013	1 kg/cm ²	2,50x10 ¹⁹	0,1 μm	2,9x10 ²³	3,4x10 ⁻⁹ s
1	1 g/cm ²	2,47x10 ¹⁶	0,1 mm	2,9x10 ²⁰	3,4x10 ⁻⁶ s
10 ⁻³	1 mg/cm ²	2,47x10 ¹³	10 cm	2,9x10 ¹⁷	3,4x10 ⁻³ s
10 ⁻⁶	1 μg/cm ²	2,47x10 ¹	100 m	2,9x10 ¹⁴	3,4 s
10 ⁻⁹	10 μg/m ²	2,47x10 ⁷	100 km	2,9x10 ¹¹	1 heure
10 ⁻¹²	0,01 μg/m ²	2,47x10 ⁴	10 ⁵ km	2,9x10 ⁸	40 jours
10 ⁻¹⁶	1 μg/km ²	2,47	10 ⁹ km	2,9x10 ⁴	1000 ans
10 ⁻²³	10 ⁻⁷ μg/km ²	2,47x10 ⁻⁷	10 ¹⁶ km	2,9x10 ⁻³	10 ¹⁰ ans