

# The Ellis and Baldwin test and the cosmic dipole anomaly

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Under the supervision of Roya Mohayaee

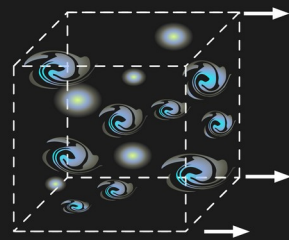


# I. The Cosmological Principle

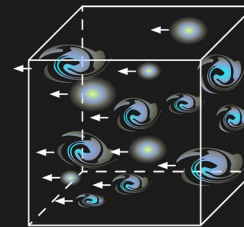
« There exist a cosmic rest frame in which the universe is, at sufficiently large scale, homogeneous and isotropic. »

(idea from Edward Milne, 1935)

The Cosmological Principle is only true within the CRF !



Rest frame



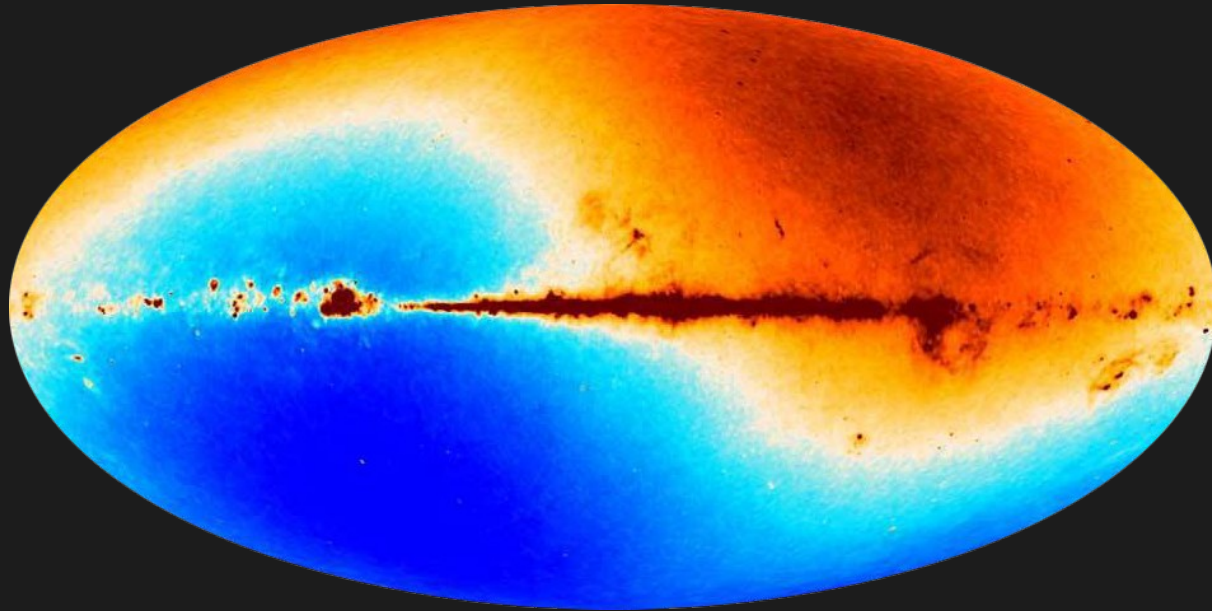
Moving frame

—▶ everything has a velocity  
—▶ anisotropic

What is this Cosmological Rest Frame (CRF) ?

## II. The CRF and the CMB

The CMB should be at rest in the CRF !

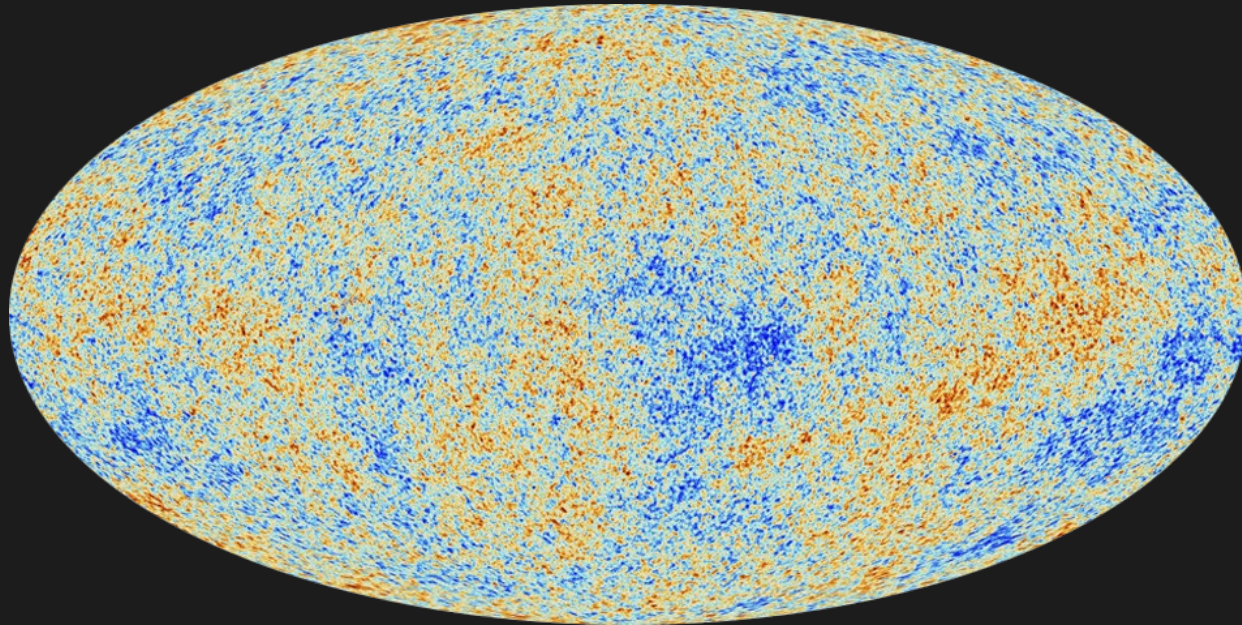


Main interpretation : CMB dipole is dominated by kinematic effects

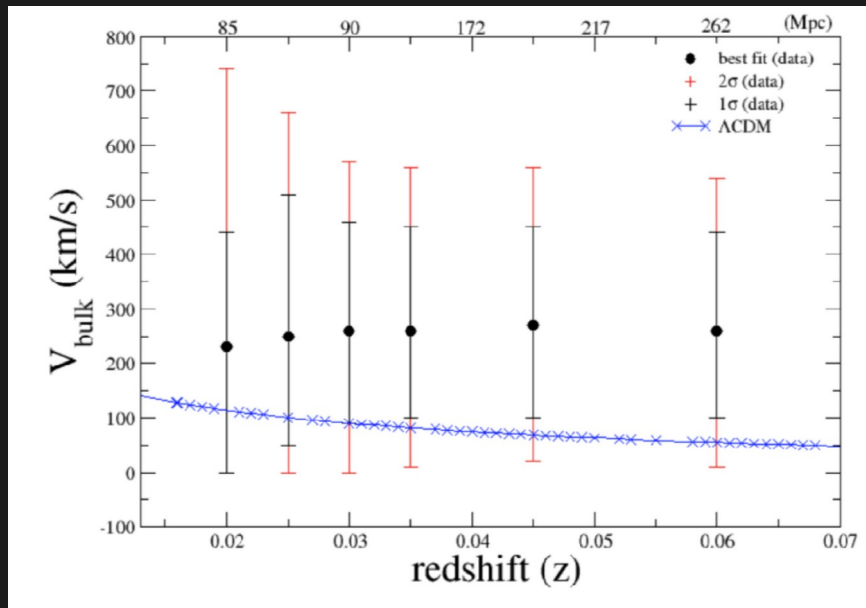
$$v_{CMB} = 369.82 \pm 0.11 \text{ km} \cdot \text{s}^{-1} \text{ toward the Virgo constellation}$$

## II. The CRF and the CMB

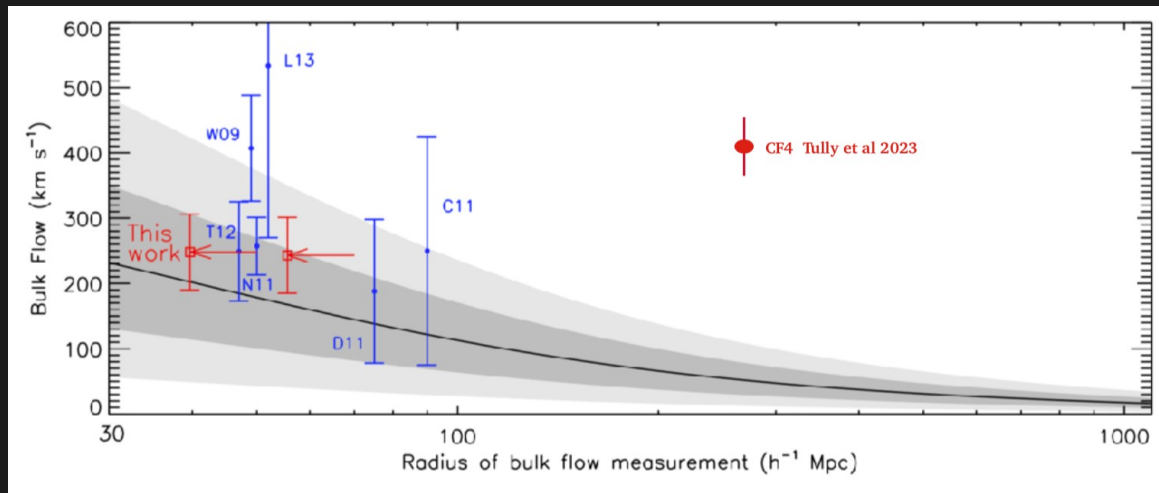
The CMB should be at rest in the CRF !



# III. Is this CMB rest frame a Cosmological rest frame ?



Bulk velocity of SN Ia, with respect to the CMB [Colin et al. 2011]



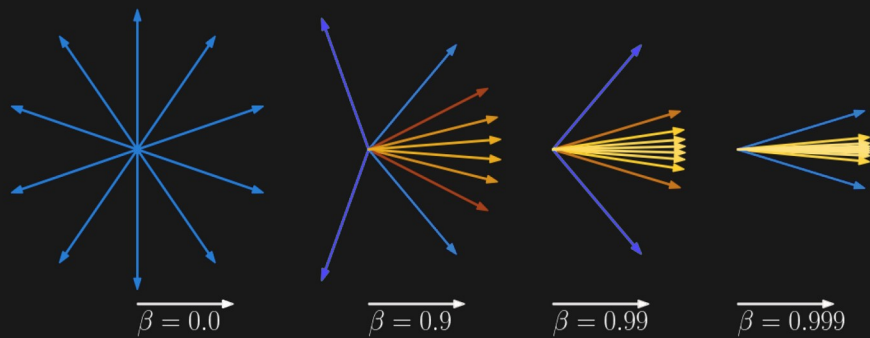
Summary of several measurement of the bulk flow velocity

It looks like the bulk-flow may be bigger than what we expect within LCDM, but it's not that clear

### III. The Ellis and Baldwin test

Our own velocity has an impact on our observations :

1) Relativistic aberrations : our field of view is distorted



$$\delta(\theta) = \gamma^{-1} (1 - \beta \cos(\theta))^{-1}$$

$$d\Omega_{obs} = \delta(\theta)^2 d\Omega_{rest}$$

2) Doppler boosting : some sources go above/bellow the detection threshold

$$dN(>F_{min}) \propto F_{min}^{-x}$$

$$F \propto \nu^{-\alpha}$$

$$F_{obs} = \delta(\theta)^{1+\alpha} F_{rest}$$

→ we model our source spectra and luminosity distribution with power laws

There is an (intrinsic) anisotropy in our observations of the universe !

$$\begin{aligned}\frac{dN}{d\Omega_{obs}}(>F_{min}, \hat{n}) &= \frac{dN}{d\Omega_{rest}}(>F_{min}, \hat{n}) \delta(\theta)^{2+x(1+\alpha)} \\ &= \bar{N} (1 + d_{kin} \cdot \hat{n})\end{aligned}$$

$$d_{kin} = (2 + x(1 + \alpha)) \beta \quad \text{(Ellis and Baldwin 1984)}$$

Model independant way to test the cosmological principle !

$$d_{kin} = (2 + \chi(1 + \alpha))\beta \quad [\text{Ellis and Baldwin 1984}]$$

Model independent way to test the cosmological principle !

Only need :

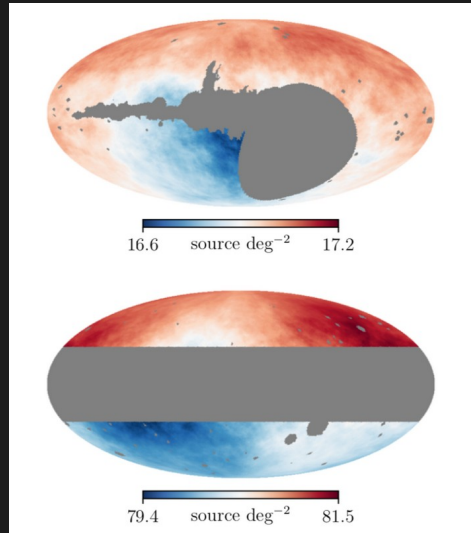
- (almost) full-sky survey → good dipole fit + prevent leakage
- *sufficiently* far-away sources → no clustering/LSS
- *minimal* knowledge of spectra/luminosity distribution → to have  $\chi, \alpha$
- above  $10^5$  sources → prevent shot noise
- consistent scanning pattern

Quite a lot of constraints actually....

## IV. What are the results for now ?

It doesn't coincide with the CMB rest frame...

- Measurement of the dipole in NVSS since 2002 [Blake and Wall 2002]
- Coherent results with other radio surveys (e.g TGSS, SUMSS, WENSS) (for ex. [Rubart Schwarz 2013], [Colin et al. 2017], etc.)
- [Secret et al. 2022] study of both CatWISE quasars and NVSS radio sources,  $5\sigma$  anomaly



NVSS radio galaxies :  
( $5.0 \times 10^5$  sources)

$$v_{NVSS} = 1100 \pm 200 \text{ km} \cdot \text{s}^{-1}$$

CatWISE quasars :  
( $1.6 \times 10^6$  sources)

$$v_{WISE} = 750 \pm 80 \text{ km} \cdot \text{s}^{-1}$$

The number count dipole amplitude is far too big for both (independent) catalogue !  
+ there is a  $\sim 40^\circ$  difference in the direction between the CMB and the matter dipole

## V. Is there something not taken into account in the EB formula ?

- Real number count : [Friedman-Shaw et al. 2024]

$$\begin{aligned}
 \Delta_g = & \left\{ b_1 \delta_m + \mathcal{H}^{-1} \partial_r v_{\parallel} \right\} + \left\{ \frac{5s-2}{2} \int_0^{\chi} d\chi' \frac{\chi - \chi'}{\chi \chi'} \Delta_{\Omega}(\Psi + \Phi) \right\} \\
 & + \left\{ \mathcal{R}_v (v_{\parallel} - v_{\parallel,o}) - (2 - 5s) v_{\parallel,o} \right\} \\
 & + \left\{ \left( \mathcal{R}_v - \frac{2-5s}{\mathcal{H}_0 \chi} \right) \mathcal{H}_0 V_o + (\mathcal{R}_v + 1) \Psi - \mathcal{R}_v \Psi_o + (5s - 2) \Phi + \dot{\Phi} \mathcal{H}^{-1} + (b_e - 3) \mathcal{H} V \right\} \\
 & + \left\{ \frac{2-5s}{\chi} \int_0^{\chi} (\Psi + \Phi) d\chi' + \mathcal{R}_v \int_0^{\chi} (\dot{\Psi} + \dot{\Phi}) d\chi' \right\}
 \end{aligned}$$

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- Real number count : [Friedman-Shaw et al. 2024]

Real matter  
distribution

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 & + \left\{ \mathcal{R}_v(v_{\parallel} - v_{\parallel,o}) - (2 - 5s)v_{\parallel,o} \right\} \quad \text{Kinematic effects} \\
 & + \left\{ \left( \mathcal{R}_v - \frac{2-5s}{\mathcal{H}_0 \chi} \right) \mathcal{H}_0 V_o + (\mathcal{R}_v + 1)\Psi - \mathcal{R}_v \Psi_o + (5s-2)\Phi + \dot{\Phi} \mathcal{H}^{-1} + (b_e - 3)\mathcal{H}V \right\} \\
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Real matter distribution

Redshift-space distortion

$$\Delta_g = \left\{ b_1 \delta_m + \mathcal{H}^{-1} \partial_r v_{\parallel} \right\} + \left\{ \frac{5s-2}{2} \int_0^{\chi} d\chi' \frac{\chi - \chi'}{\chi \chi'} \Delta_{\Omega}(\Psi + \Phi) \right\}$$

Kinematic effects

$$+ \left\{ \mathcal{R}_v (v_{\parallel} - v_{\parallel,o}) - (2 - 5s)v_{\parallel,o} \right\}$$

$$+ \left\{ \left( \mathcal{R}_v - \frac{2-5s}{\mathcal{H}_0 \chi} \right) \mathcal{H}_0 V_o + (\mathcal{R}_v + 1)\Psi - \mathcal{R}_v \Psi_o + (5s-2)\Phi + \dot{\Phi} \mathcal{H}^{-1} + (b_e - 3)\mathcal{H}V \right\}$$

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Real matter distribution      Redshift-space distortion      Weak-lensing

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+  $\left\{ \mathcal{R}_v (v_{\parallel} - v_{\parallel,o}) - (2 - 5s)v_{\parallel,o} \right\}$       Kinematic effects

+  $\left\{ \left( \mathcal{R}_v - \frac{2-5s}{\mathcal{H}_0 \chi} \right) \mathcal{H}_0 V_o + (\mathcal{R}_v + 1)\Psi - \mathcal{R}_v \Psi_o + (5s-2)\Phi + \dot{\Phi} \mathcal{H}^{-1} + (b_e - 3)\mathcal{H}V \right\}$

+  $\left\{ \frac{2-5s}{\chi} \int_0^{\chi} (\Psi + \Phi) d\chi' + \mathcal{R}_v \int_0^{\chi} (\dot{\Psi} + \dot{\Phi}) d\chi' \right\}$

GR

1) The dominant effect here should be the clustering

= the intrinsic anisotropy of the matter distribution, due to LSS

How can we constrain it ?

Expected power spectrum within LCDM, for a source distribution  $n_s(r)$  :

$$\langle C_l^{cl} \rangle = \frac{2}{\pi} \int dk k^2 P(k, z=0) W_l^{cl}(k)^2$$

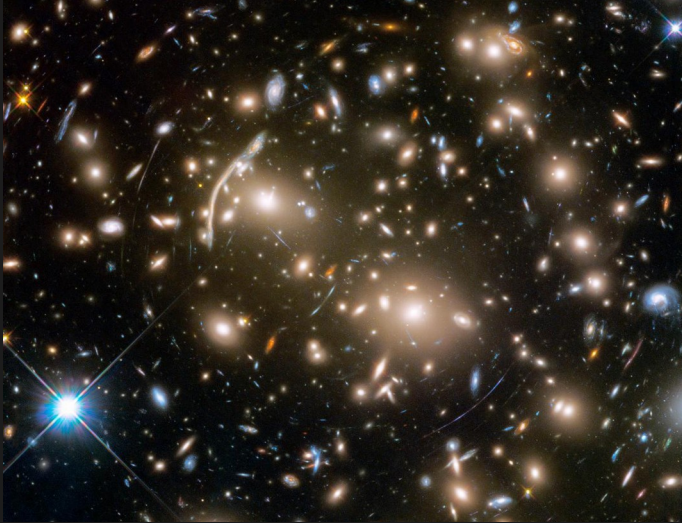
$$W_l^{cl}(k) = \int_0^{r_{max}} dr n_s(r) b_s(r) D_+(r) j_l(kr)$$

$$\langle d_{cl}^2 \rangle = \frac{9 \langle C_1^K \rangle}{4 \pi}$$

At far enough redshift, this clustering dipole can be as small as we want, but dominates at close distance.

In CatWISE  $d_{cl} \approx 2.4 \times 10^{-4}$  compared to  $d_N \approx 1.5 \times 10^{-2}$

## 2) What about Weak-lensing ?



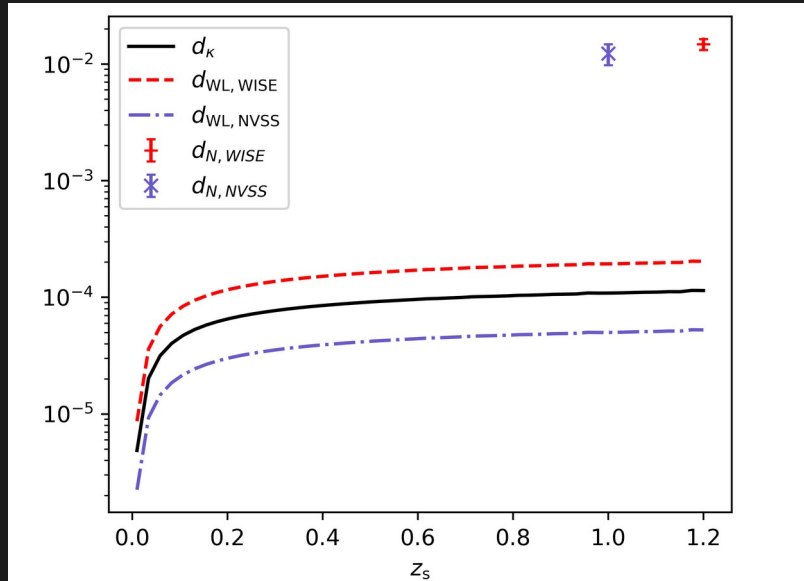
The local universe deforms observations along the line of sight, through the convergence field :

$$\kappa(r_s, \hat{n}) = \frac{3H_0\Omega_m}{2c^2} \int_0^{r_s} dr \frac{r(r_s - r)}{r_s a(r)} \delta_m(r, \hat{n})$$

- WL impact should be *stronger* in the same direction as the dipole
- Should be stronger at large redshift
- Impact of the convergence dipole on the number count dipole

$$d_{WL} = 2(x-1)d_\kappa$$

[Mohayaee and Bonnefous 2026]  $\longrightarrow$  Its also negligible !



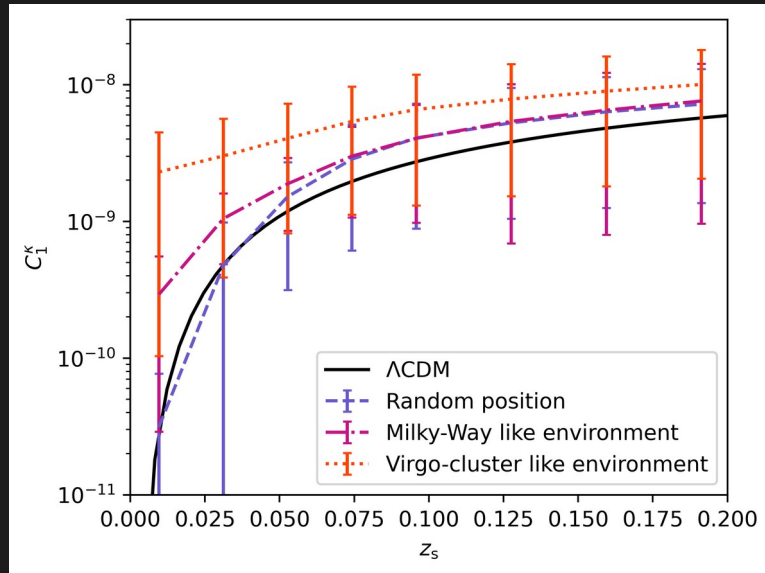
Expected convergence dipole in LCDM,  
compared with the anomalous dipole in NVSS  
and WISE

- Same as clustering, we can obtain expected LCDM convergence dipole
- How much that expectation is impacted by the local matter distribution ?

$$\langle C_l^K \rangle = \frac{2}{\pi} \int dk k^2 P(k, z=0) W_l^K(k)^2$$

$$W_l^{cl}(k) = \frac{3 H_0^2 \Omega_m}{2 c^2} \int_0^{r_{max}} dr \frac{r(r_{max}-r)}{r_{max} a(r)} D_+(r) j_l(kr)$$

- Using simulation from the Quijote suite [F. Villaescusa-Navarro et al. 2020], to account for the particular position of the MW, we can constrain the value of the convergence dipole



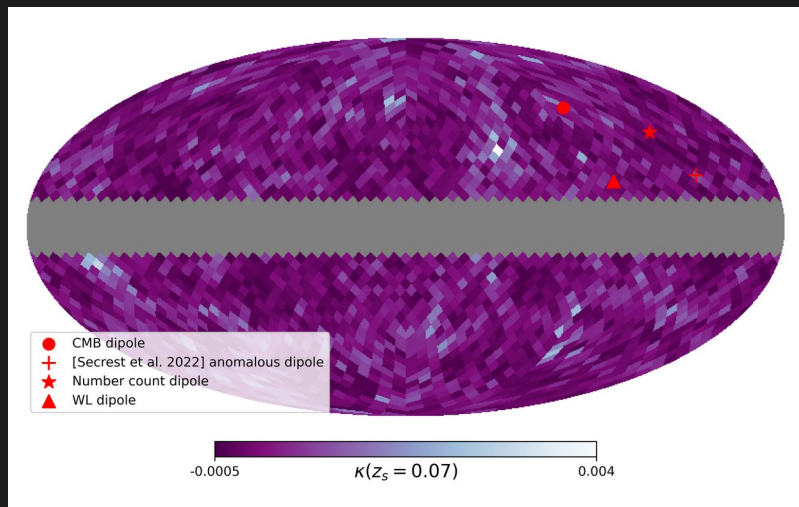
- Defined 3 types of observer (Random, Virgo-cluster-like, MW-like)
  - defined according to halo mass, distance to cluster and velocity
- For each type, take 50 of them
- Reconstruct the convergence field and extract low multipoles

$l=1$  angular power spectrum of the convergence field in a Quijote simulation for different types of observer

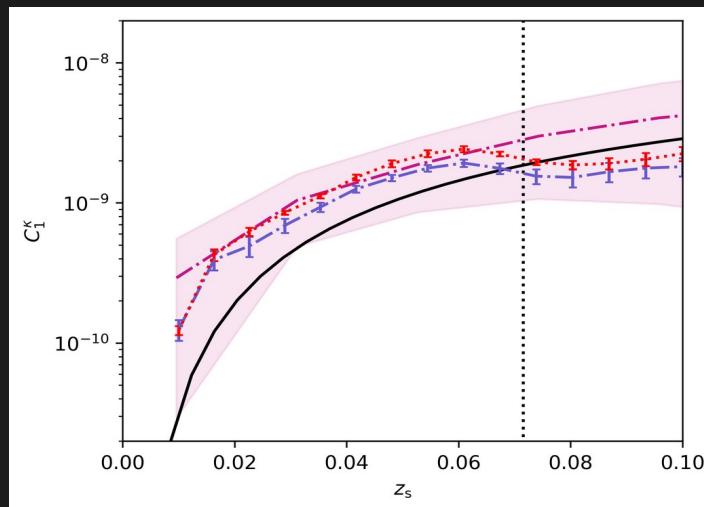
Local universe has a huge impact on the convergence field !

→ Only at low redshift (up to  $z \sim 0.1$ )

- Using 2MRS, we can reconstruct the (local) matter density field, then the local convergence field (up to  $z \sim 0.07$ )



Reconstructed convergence field using 2MRS



Reconstructed dipole of the convergence field at different redshift compared with the LCDM expectation

→ align with the CMB !

Conclusion : For now, not the beginning of an explanation from the cosmological perspective within LCDM model

### 3) What if we don't have a perfect power law ?

[Bonnetfous 2026]

—————▶ it doesn't change anything, we just have to redefine the coefficients

$$\alpha_{eff} = \frac{\int d\nu T_X(\nu) \nu \frac{\partial S_\nu}{\partial \nu}}{\int d\nu T_X(\nu) S_\nu} \quad \chi_{eff} = \frac{-\partial \log N}{\partial \log L}$$

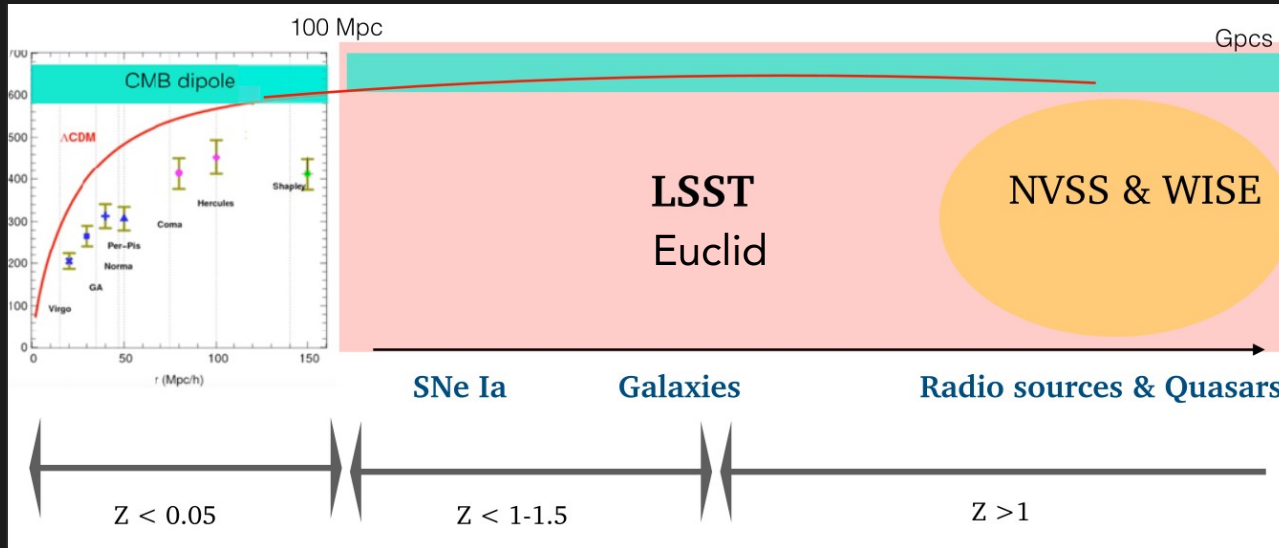
We can generalize the EB test to (pretty much) any large scale survey !

—————▶ especially NIR/visible galactic survey

# VI. Future application of this test

Future large scale surveys : LSST, Euclid, SPHEREx

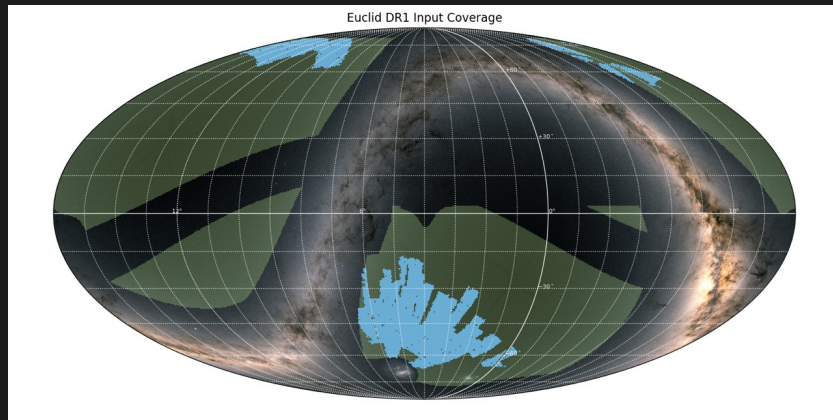
Goal : bridging the local universe to the distant radio sources and quasars



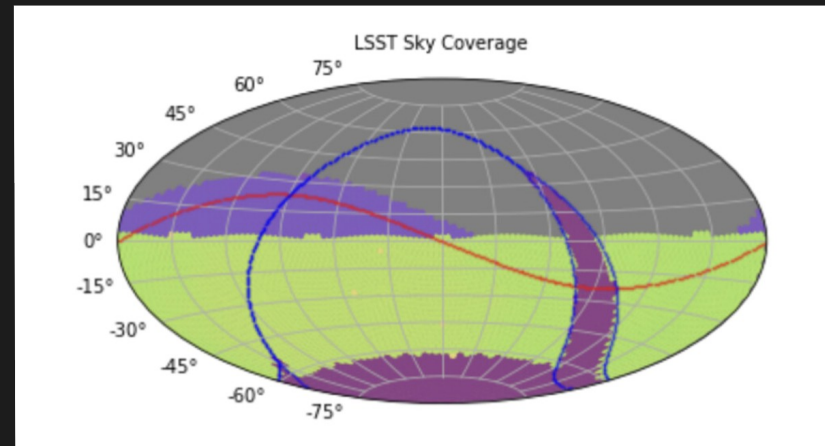
# VI. Future application of this test

Completely new challenges :

- Smaller cover = more leakage
- Visible/NIR galaxies = no power laws, obtaining  $\alpha$  is tricky
- REDENNING, extinction



Euclid sky coverage (in blue : DR1)



LSST sky coverage

## VI. Conclusion :

- There is a  $5\sigma$  discrepancy between the CMB kin. dipole, and the matter dipole
- Weak lensing is a too weak effect to have to be taken into account, still no explanation of this anomaly **within** LCDM
- We can generalize the EB test to any spectral profile/any magnitude distribution !
- We're working with LSST, Euclid (and SPHEREx ?), to apply this test to new sets of data
- Current work :
  - 1) Creating mock catalogue
  - 2) Obtaining  $\alpha$  for galaxies in the visible/NIR