



# Statistical Nuclear Reactions & Quantum Pre-Equilibrium Models

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- Microscopic Projectors
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# Statistical Nuclear Reactions

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# Single Resonances

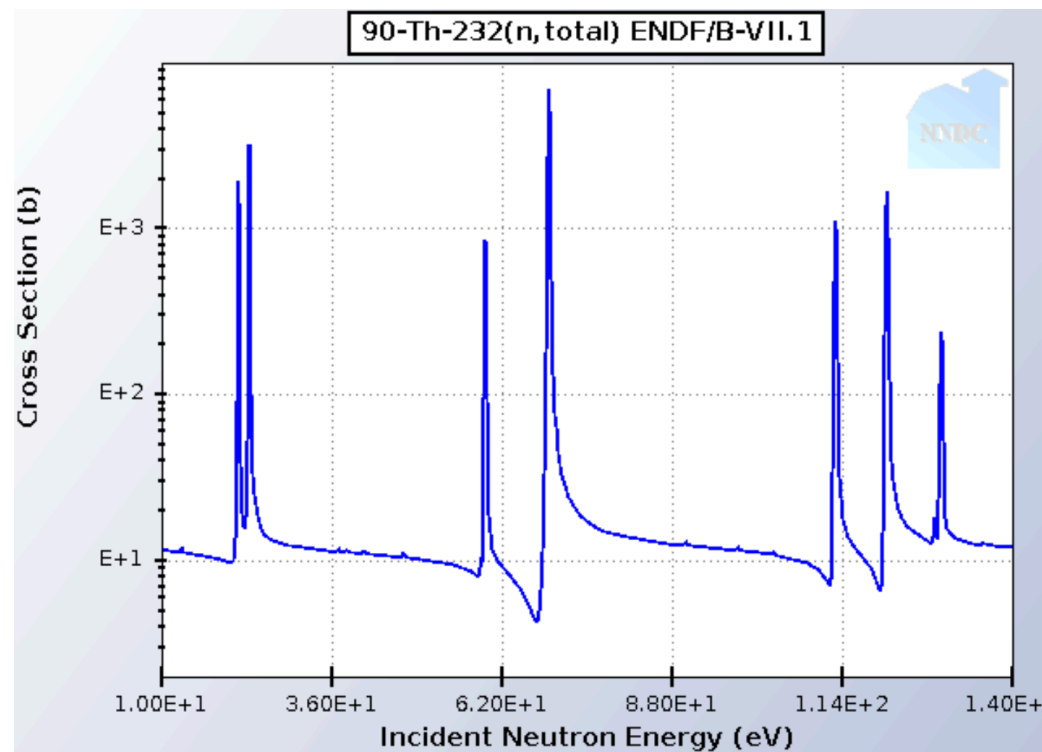
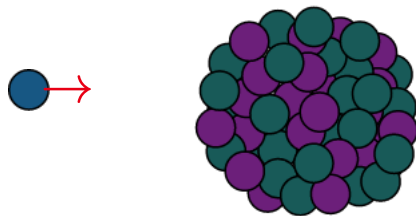
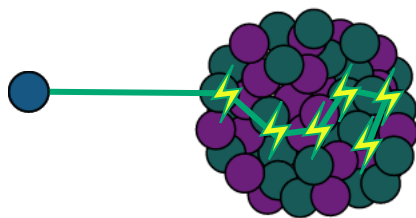


Figure 1:  $n + {}^{232}\text{Th}$  (from NNDC)

# Single Resonances



Compound Nucleus Formation

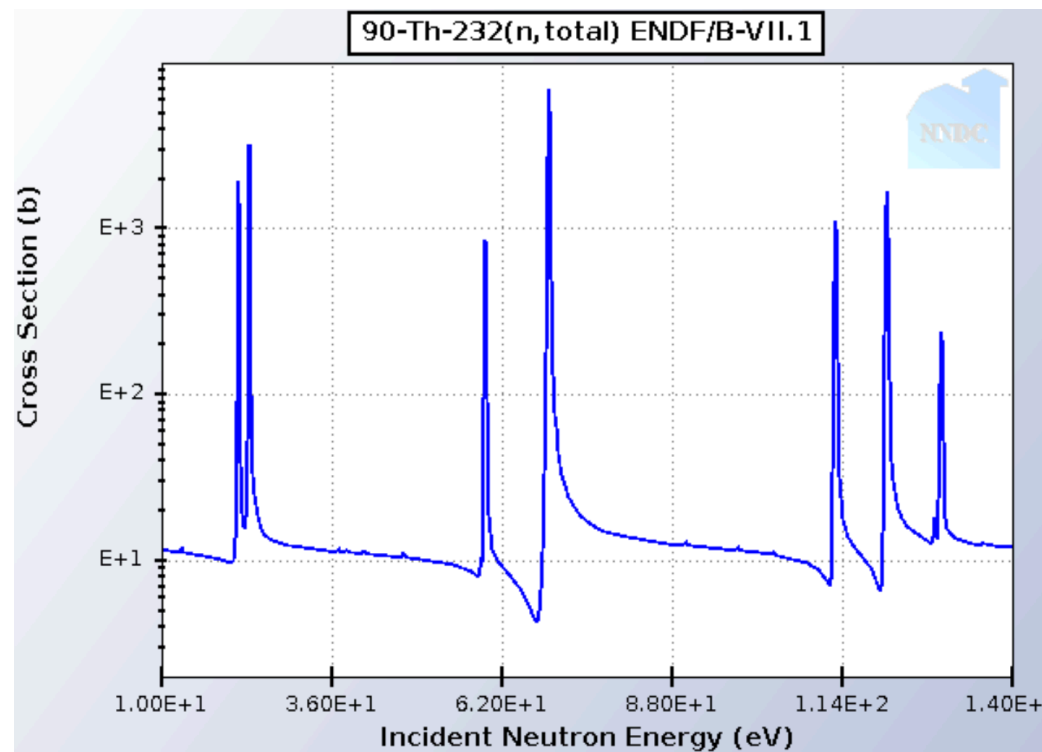
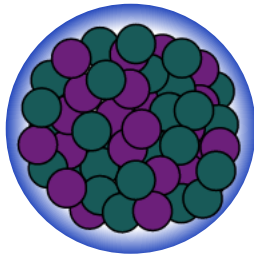


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# Single Resonances



Compound Nucleus Formation  
form Quasibound state of lifetime  $\sim \frac{1}{\Gamma}$

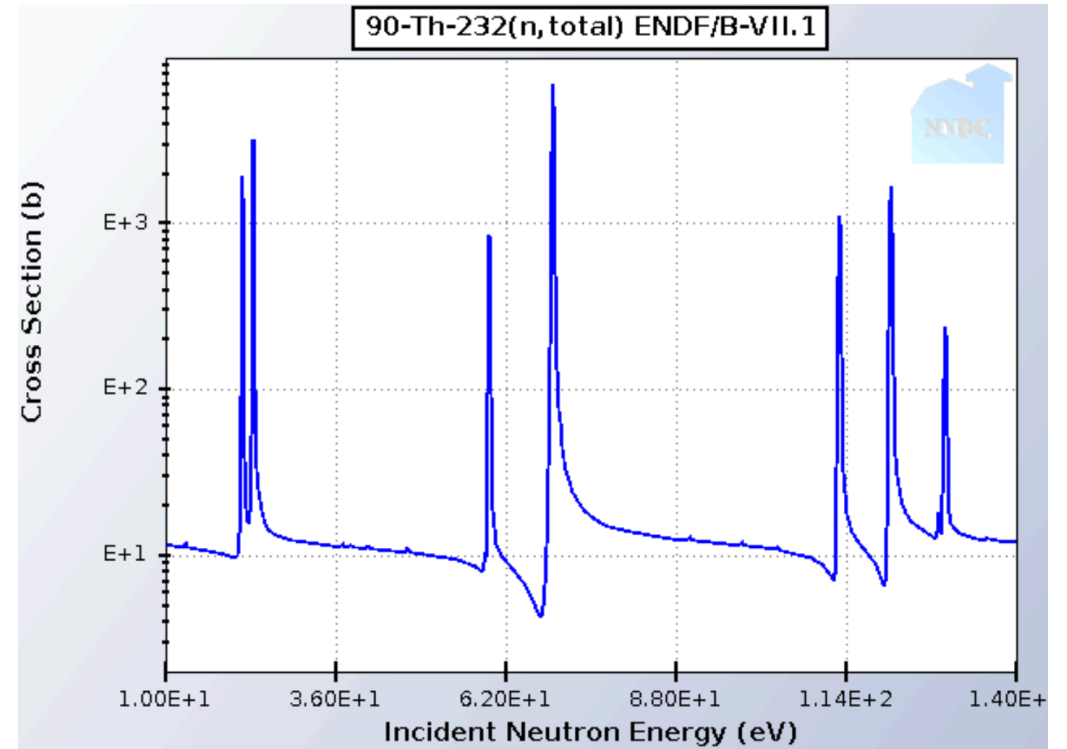
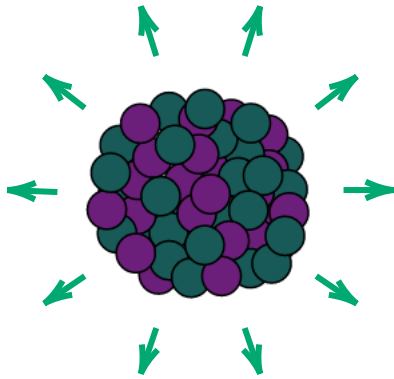


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# Single Resonances



Compound Nucleus Formation  
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Compound Nucleus Decay

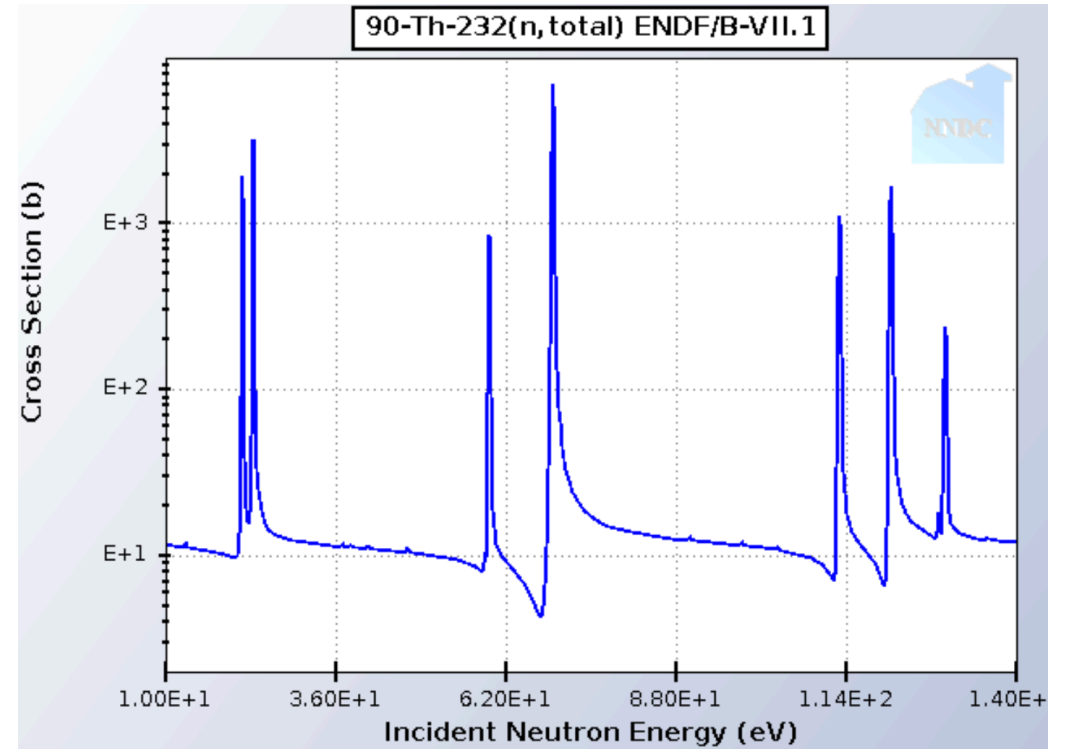
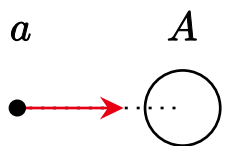
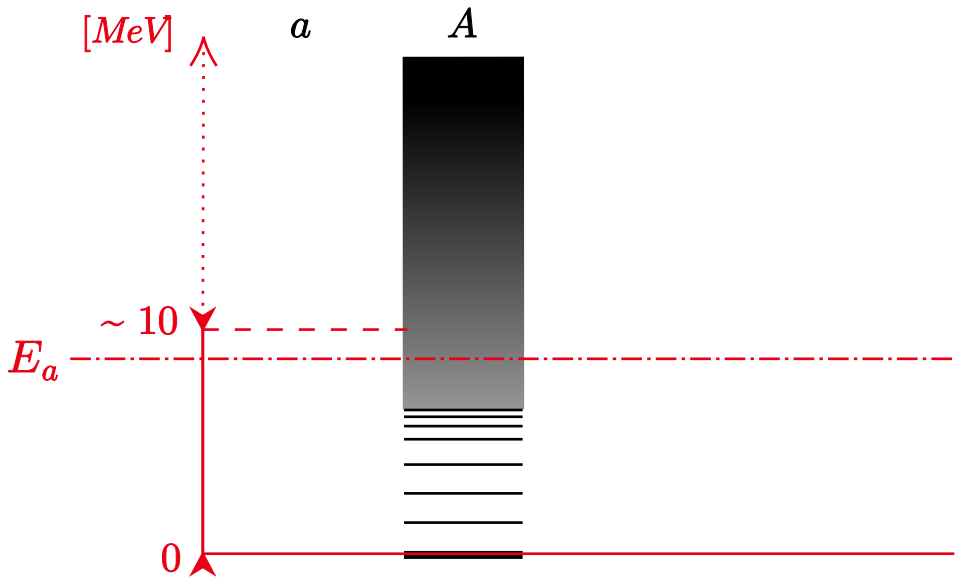
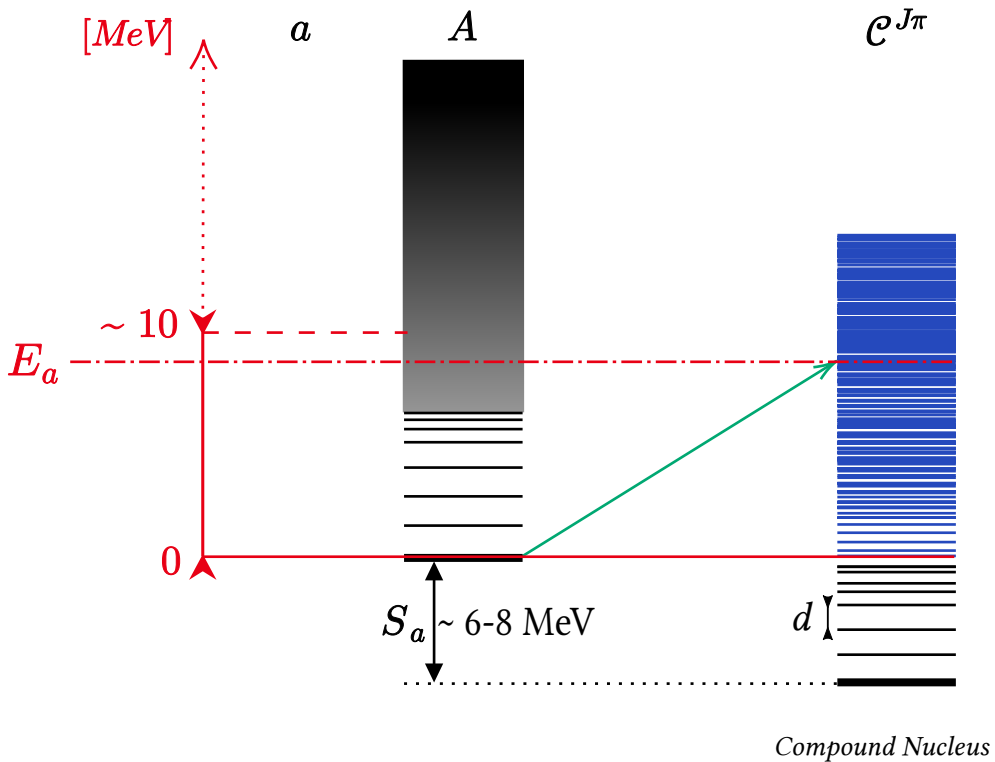


Figure 1:  $n + {}^{232}\text{Th}$  (from NNDC)

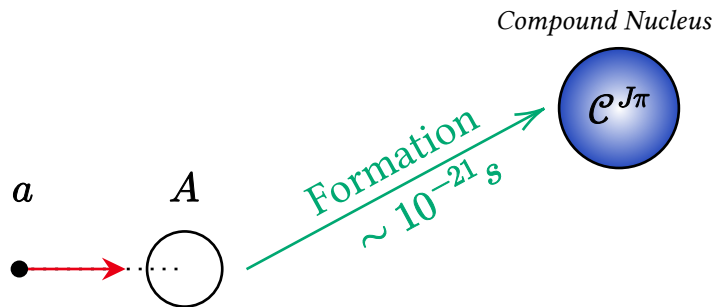
# Overlapping Resonances



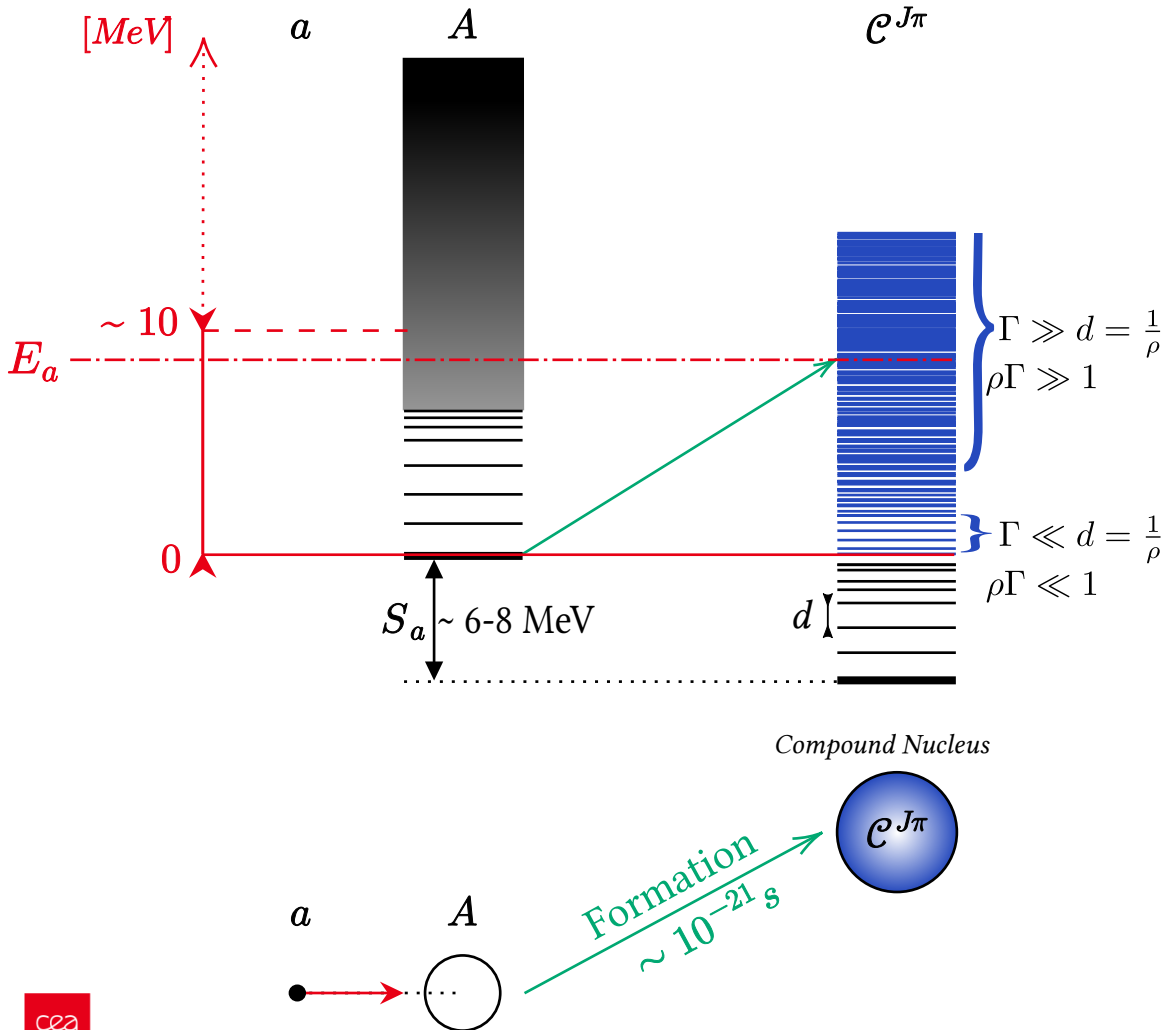
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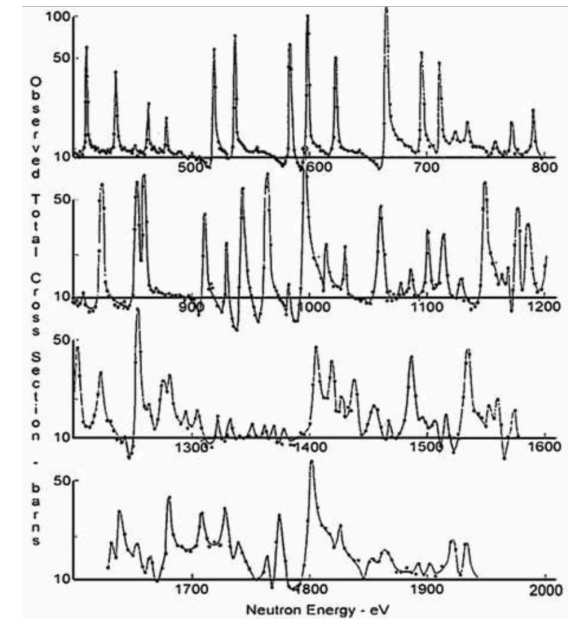
- $\rho \sim e^{\sqrt{cE}}$



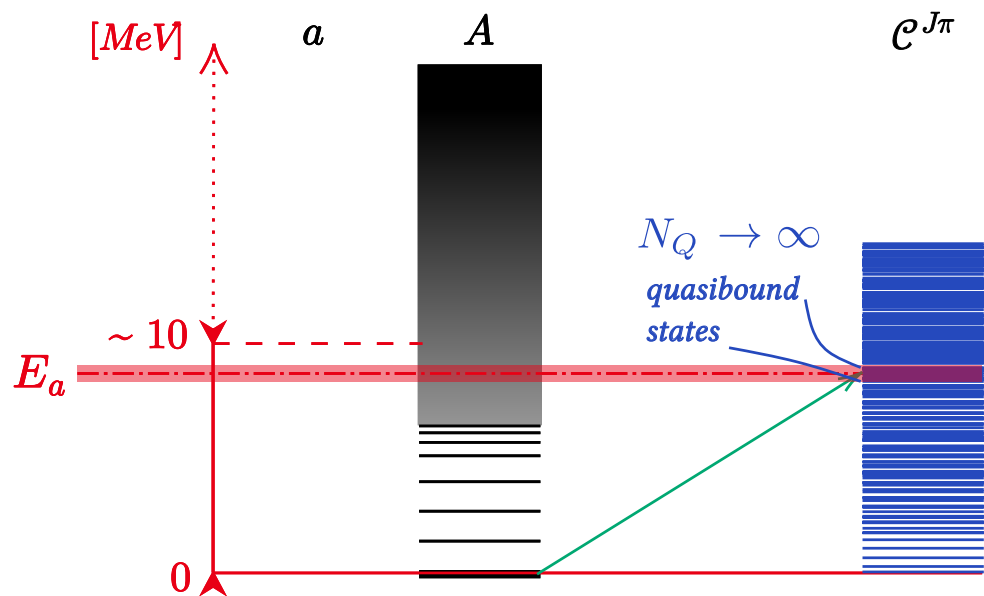
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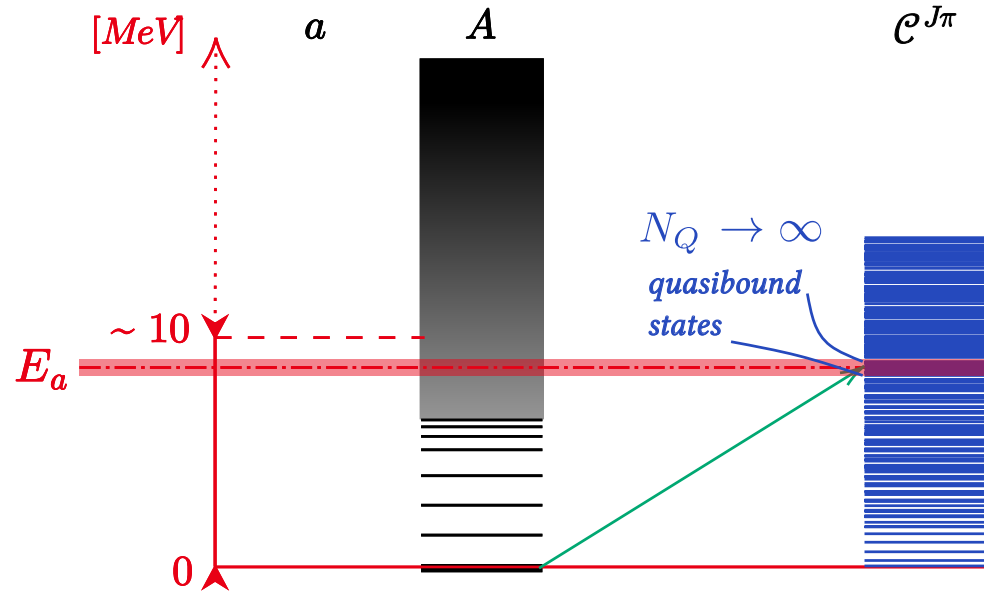
- $\rho \sim e^{\sqrt{cE}}$   
 $\hookrightarrow$  Number of open channels  $\Lambda$  grows similarly
- $\rho\Gamma \propto \sum_{\beta=1}^{\Lambda} T_{\beta} = \Lambda^{\text{eff}}$   
 $\hookrightarrow$  semiclassical expansion in  $(\Lambda^{\text{eff}})^{-1}$
- From single to overlapping resonances



# Average & Fluctuations



# Average & Fluctuations



## Feshbach projectors

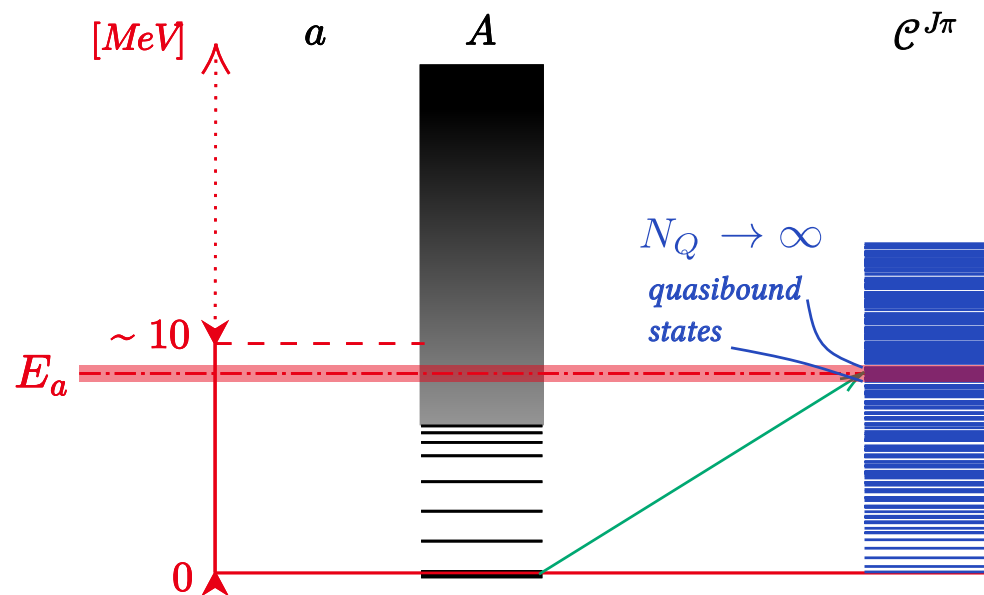
$P(E) = \sum_{\beta=1}^{\Lambda(E)} P_{\beta}$  onto the  $\Lambda(E)$  open channels  $\beta$

$Q(E) = 1 - P(E)$  onto the  $N_Q$  quasibound states  
& the closed channels

$$|\Psi_{\alpha}^{+}\rangle = P|\Psi_{\alpha}^{+}\rangle + Q|\Psi_{\alpha}^{+}\rangle = |\Psi_{\alpha}^{+}\rangle + |\Psi_{\alpha}^{+}\rangle$$

$$H^N = H_{PP}^N + H_{QQ}^N + H_{PQ}^N + H_{QP}^N$$

# Average & Fluctuations



## Coupled equations

$$\begin{aligned} \left[ E - H_{QQ}^N - H_{QP}^N \frac{1}{E^+ - H_{PP}^N} H_{PQ}^N \right] |\Psi_\alpha^+\rangle &= H_{QP}^N |\psi_\alpha^+\rangle \\ \left[ E - H_{PP}^N - H_{PQ}^N \frac{1}{E^+ - H_{QQ}^N} H_{QP}^N \right] |\Psi_\alpha^+\rangle &= 0 \end{aligned}$$

## Feshbach projectors

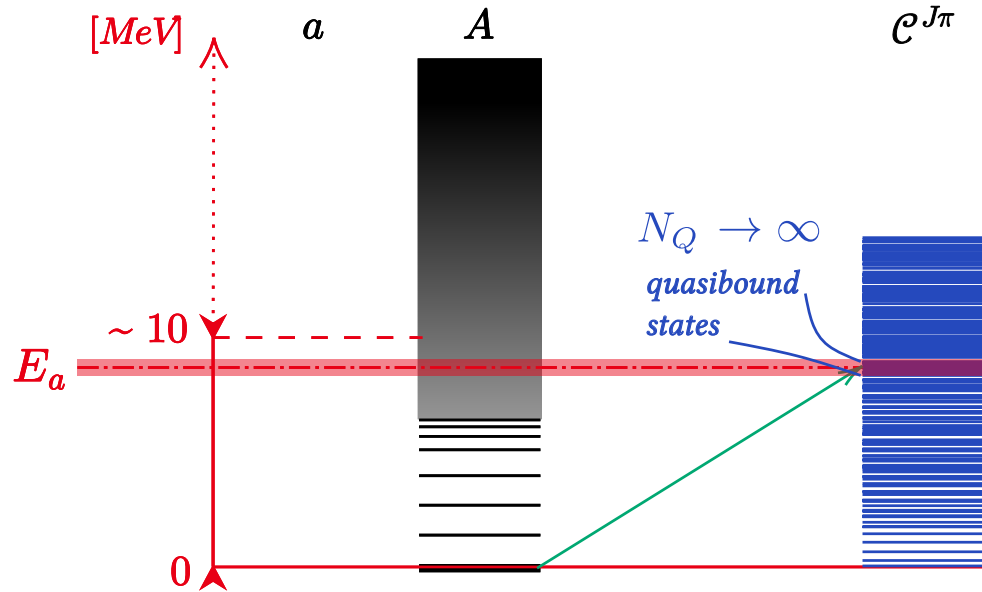
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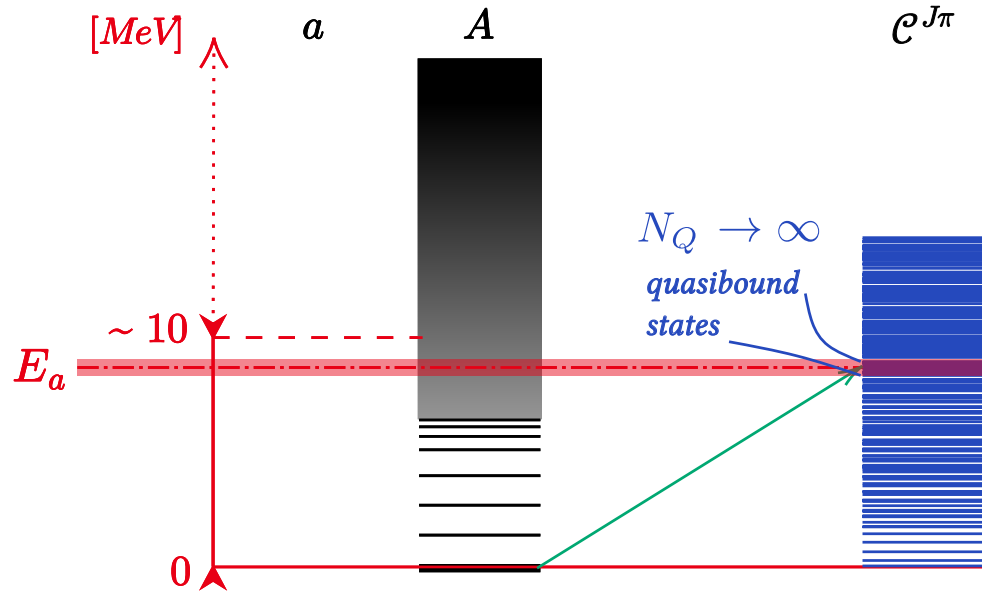
## S-matrix

$$S_{\alpha\beta} = \langle \psi_{\beta}^{-} | \psi_{\alpha}^{+} \rangle - i2\pi \langle \psi_{\beta}^{-} | H_{PQ}^N \frac{1}{\mathcal{D}(E)} H_{QP}^N | \psi_{\alpha}^{+} \rangle$$

where  $\mathcal{D} = E - H_{QQ}^N - \xi(E) + i\frac{1}{2}\Gamma(E)$

$$= \boxed{\langle S_{\alpha\beta} \rangle^{\Delta E}} + \boxed{S_{\alpha\beta}^{\text{fl}}}$$

# Average & Fluctuations



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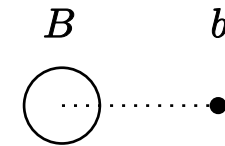
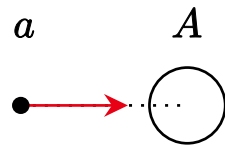
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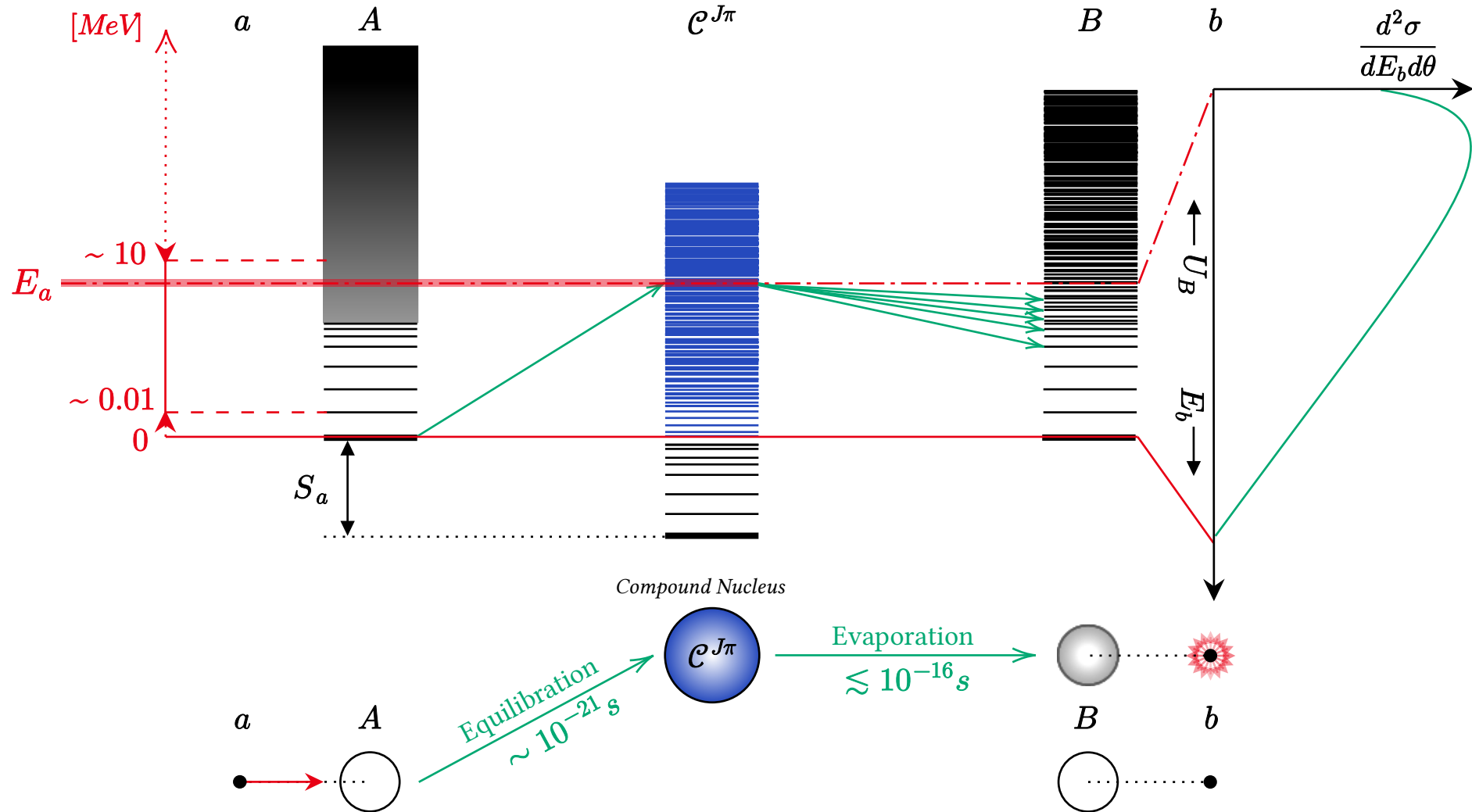
## Cross section

$$\begin{aligned} \langle \sigma_{\alpha\beta} \rangle^{\Delta E} &\propto \langle |S_{\alpha\beta} - \delta_{\alpha\beta}|^2 \rangle^{\Delta E} \\ &\propto \underbrace{\langle |S_{\alpha\beta} \rangle^{\Delta E} - \delta_{\alpha\beta}|^2 }_{\sigma_{\alpha\beta}^{\text{D}}} + \underbrace{\langle |S_{\alpha\beta}^{\text{fl}}|^2 \rangle^{\Delta E}}_{\sigma_{\alpha\beta}^{\text{CN}}} \end{aligned}$$

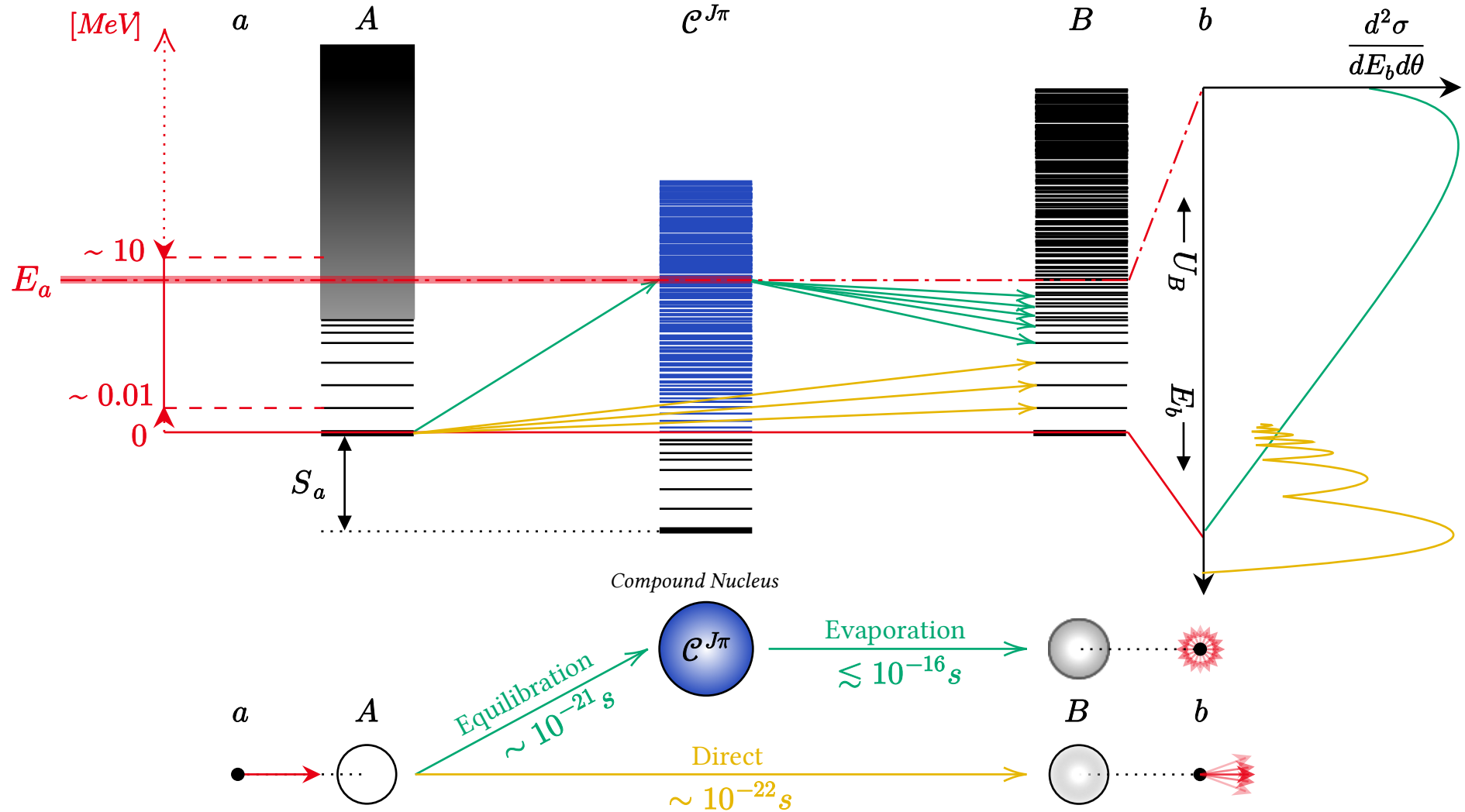
# Nucleon-channels reactions at $E_a \lesssim 10$ MeV



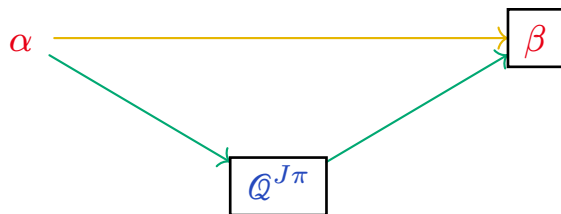
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# Statistical Reaction Theory



## S-matrix

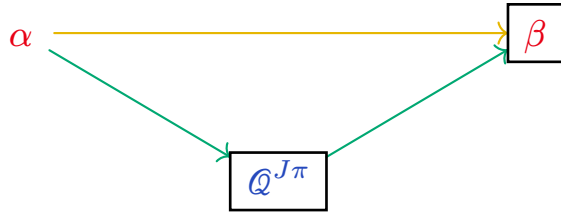
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where  $\mathcal{D} = E - H_{QQ}^N - \xi(E) + i\frac{1}{2}\Gamma(E)$

## Why $J\pi$ ?

$$\langle \sigma_{\alpha\beta} \rangle^{\Delta E} \propto \underbrace{\left| \langle S_{\alpha\beta} \rangle^{\Delta E} - \delta_{\alpha\beta} \right|^2}_{\sigma_{\alpha\beta}^{\text{D}}} + \underbrace{\left\langle \left| S_{\alpha\beta}^{\text{fl}} \right|^2 \right\rangle^{\Delta E}}_{\sigma_{\alpha\beta}^{\text{CN}}}$$

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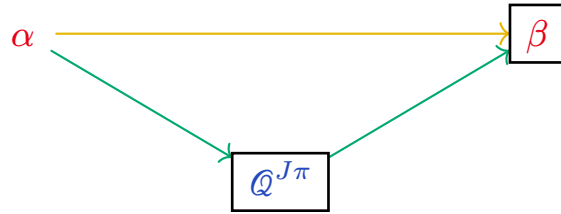
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**Chaos in nuclei**

#constraints  $\ll$  #dof  $\Rightarrow$   $\begin{cases} \text{regular motion} \\ \text{chaotic motion} \end{cases}$

$$H_{QQ}^N = \begin{pmatrix} H_{QQ}^{N J_1 \pi_1} & & \\ & H_{QQ}^{N J_2 \pi_2} & \\ & & \dots \end{pmatrix}$$

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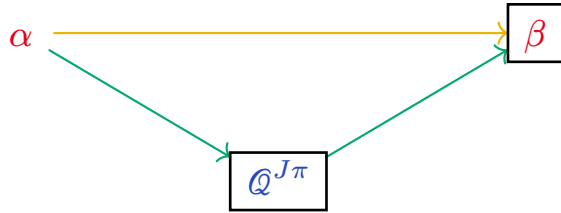
$$\langle \sigma_{\alpha\beta} \rangle^{\Delta E} \propto \underbrace{\left| \langle S_{\alpha\beta} \rangle^{\Delta E} - \delta_{\alpha\beta} \right|^2}_{\sigma_{\alpha\beta}^{\text{D}}} + \left\langle \left| \sum_{J\pi} S_{\alpha\beta}^{\text{fl } J\pi} \right|^2 \right\rangle^{\Delta E}$$

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↪ uncorrelated  $S^{\text{fl}}$ -matrix for different  $J, \pi$

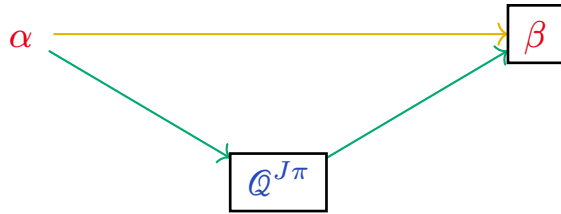
↪  $90^\circ$  symmetrical emission 

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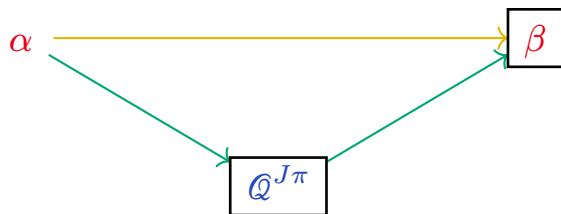
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**Random Hamiltonians**

$$H_{QQ}^{NJ\pi} \rightarrow H^{\text{GOE}} \in \left\{ \begin{pmatrix} H_{11} & \dots & H_{1N_Q} \\ & \ddots & \vdots \\ & & H_{N_Q N_Q} \end{pmatrix} \right\}$$

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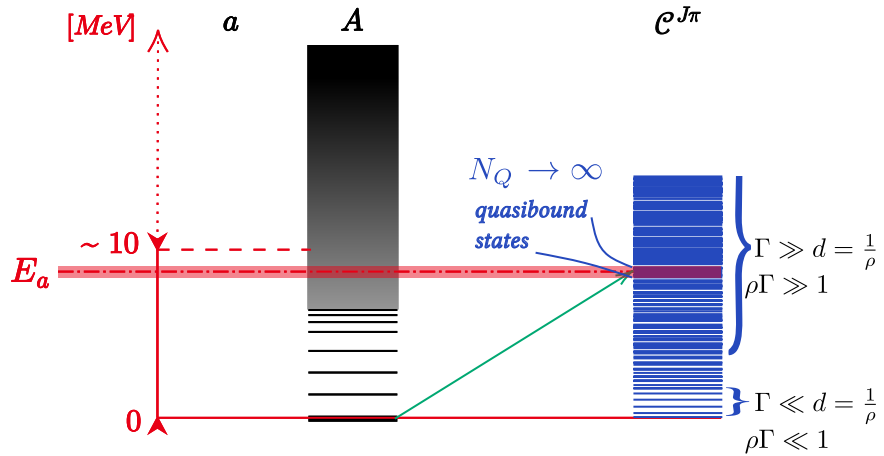
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**Ergodicity of GOE**

$$\langle f(H^{\text{GOE}}) \rangle^{\Delta E} \xrightarrow{N_Q \rightarrow \infty} \overline{f(H^{\text{GOE}})}^N (= \mathbb{E}(f(H^{\text{GOE}})))$$

$$\langle |S_{\alpha\beta}^{\text{fl} J\pi}|^2 \rangle^{\Delta E} \rightarrow \overline{|S_{\alpha\beta}^{\text{fl} J\pi}|^2}^N \sim \frac{1}{\mathcal{D}} \frac{1}{\mathcal{D}}^N$$

# Statistical Reaction Theory



$$\left\langle |S_{\alpha\beta}^{\text{fl } J\pi}|^2 \right\rangle^{\Delta E} \rightarrow \overline{|S_{\alpha\beta}^{\text{fl } J\pi}|^2}^N \sim \frac{1}{\mathcal{D}} \frac{1}{\mathcal{D}}^N$$

- Diagrammatic expansion & Partial Resummation  
 ✓ leading terms in  $(\Lambda^{\text{eff}} \propto \rho\Gamma)^{-1}$
- Generating functional approach  $\frac{1}{\mathcal{D}} \frac{1}{\mathcal{D}}^N \propto \frac{\partial^2}{\partial J^2} \ln Z(E, J)|_{J=0}$   
 $\hookrightarrow$  “replica trick”  $\ln(Z) = \lim_{n \rightarrow 0} \frac{1}{n} (Z^n - 1)$  ✓  
 $\hookrightarrow$  ✓ SUSY integration method to normalize  $Z(E, 0) = 1$   
 $\Rightarrow \frac{\partial^2}{\partial J^2} \ln Z(E, J)|_{J=0} = \frac{\partial^2}{\partial J^2} Z(E, J)|_{J=0}$
- SUSY 1985 : ✓ CN (VWZ) [2]



# Random Matrix Theory & Nuclear Physics

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# Symmetries & Chaos



- Which matrix ensemble for which symmetries
  - **3-fold way** by Wigner, Dyson (1950s, 1960s)
  - **10-fold way** (~ Cartan) by Altland, Zirnbauer (1997)
  - **non-Cartan way** for non-Hermitian Hamiltonians (RPA...)

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  - **non-Cartan way** for non-Hermitian Hamiltonians (RPA...)
- **no microscopic information required !!!**
  - ↪ universal ( $c \sim d \sim \text{GOE}$ )
  - ↪ ~ thermodynamics for spectral statistics

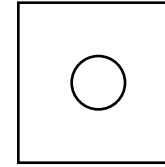


Figure 1: Sinai Billiard

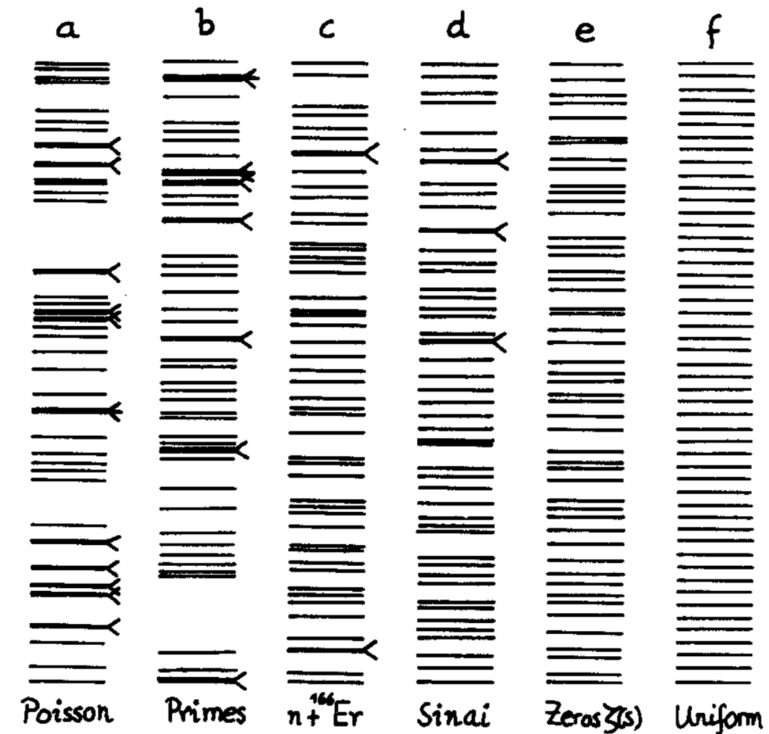


Figure 2: Various unfolded spectra (from [3])

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  - ↪ ~ thermodynamics for spectral statistics
- A **quantum** system is **chaotic** if ?
  1. its classical counterpart is chaotic (based on periodic orbits)
  2. it displays random matrix spectral statistics

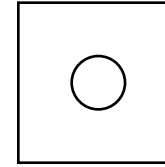


Figure 1: Sinai Billiard

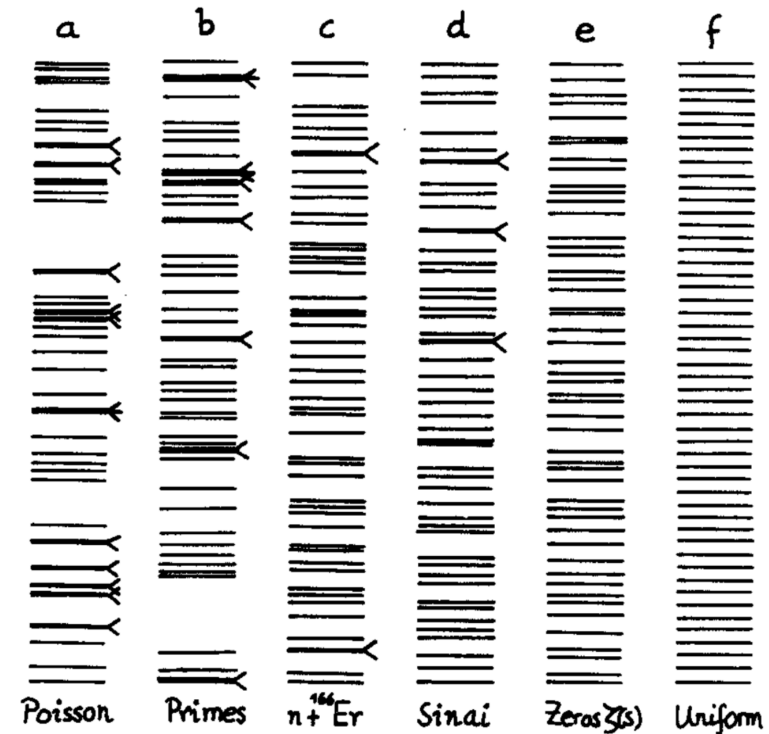


Figure 2: Various unfolded spectra (from [3])

# Randomness in Shell-Model



- GOE vs shell-model

$$H_{QQ}^{NJ\pi} \rightarrow H^{\text{GOE}} \in \left\{ \left( \begin{array}{ccc} H_{11} & \dots & H_{1N_Q} \\ & \ddots & \vdots \\ & & H_{N_Q N_Q} \end{array} \right); \text{iid } H_{\mu\nu} \hookrightarrow \mathcal{N}(0, (1 + \delta_{\mu\nu})\sigma^2) \right\}$$

every pair of quasibound states is equally coupled  
equivalent ‘1, ...,  $N$ -body’ interaction  $\rightarrow \kappa$

- 1970-1971 : TBRE (Two-Body Random Ensemble)
- 1975 : EGOE(k) (Embedded k-body interaction GOE)
- 2015 : SYK model (Dirac  $\rightarrow$  Majorana fermions)

$$\text{AdS}_2 \leftrightarrow \text{CFT}_1 = \text{QM}$$

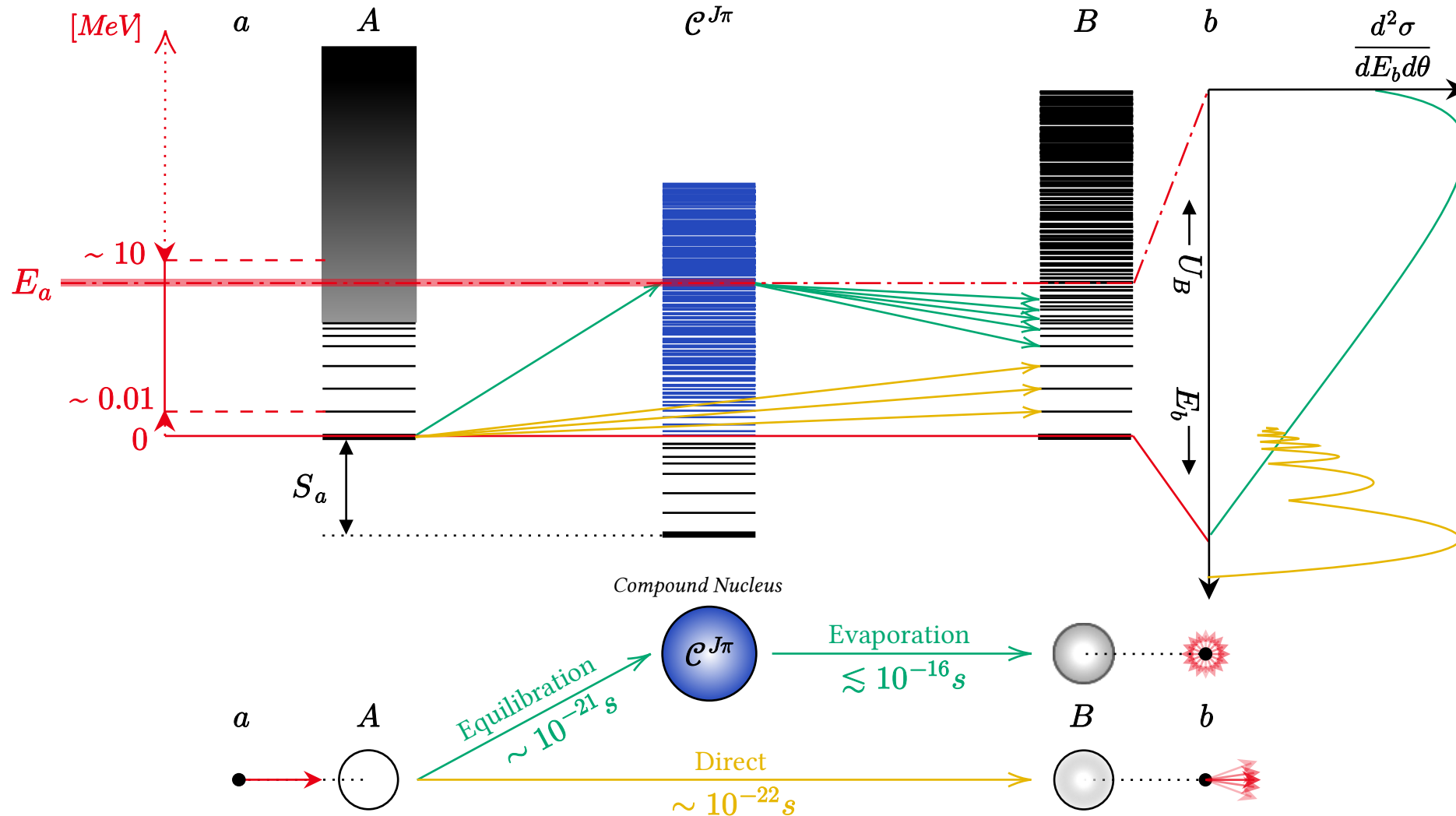
$$\text{black-hole} \leftrightarrow \text{SYK model} \sim \mathcal{C}^{J\pi}$$



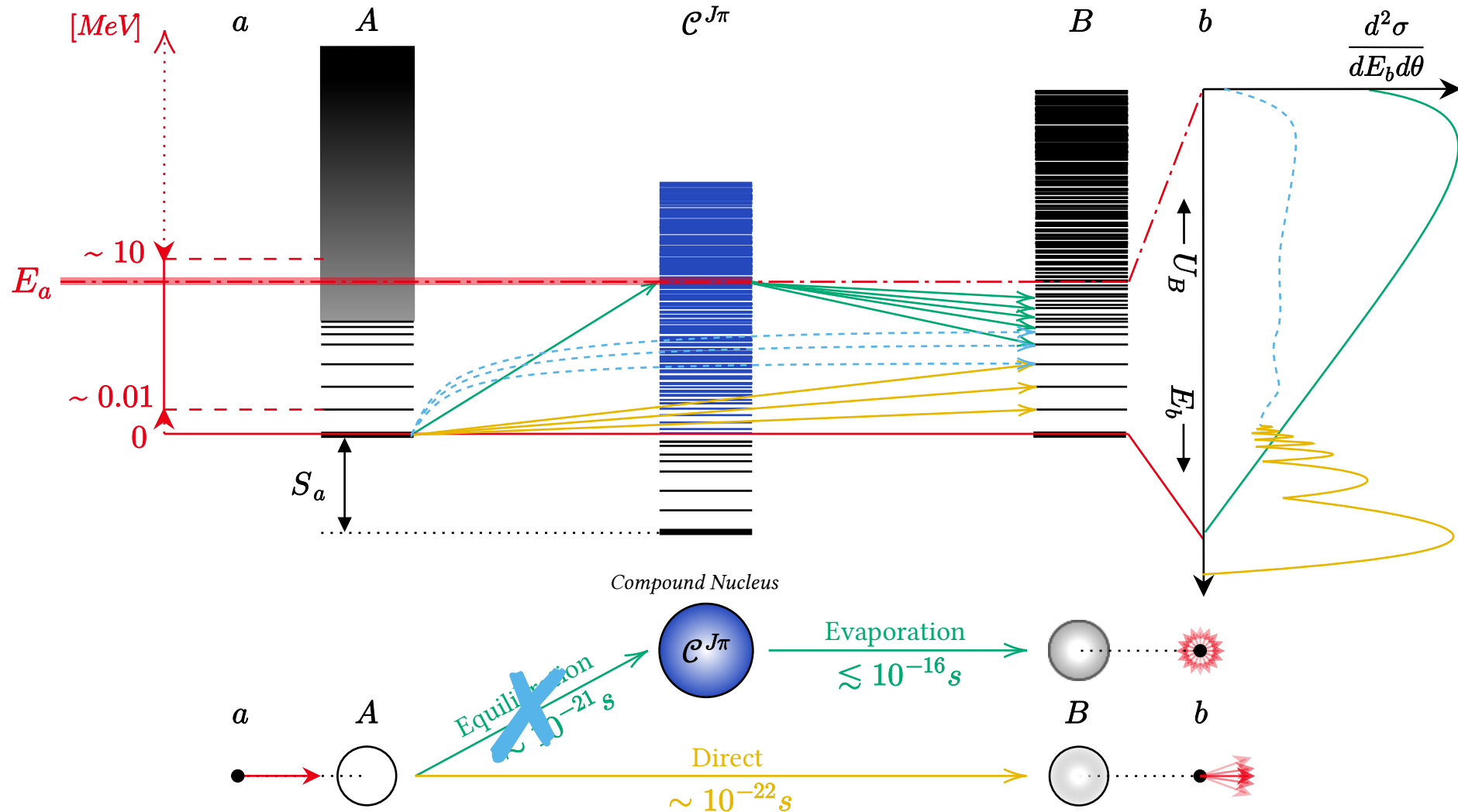
# Quantum Pre-Equilibrium Models

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# Intermediate Structure & Pre-Equilibrium Emission



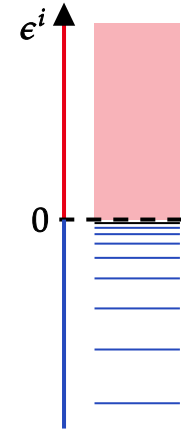
# Intermediate Structure & Pre-Equilibrium Emission



# Microscopic Projectors

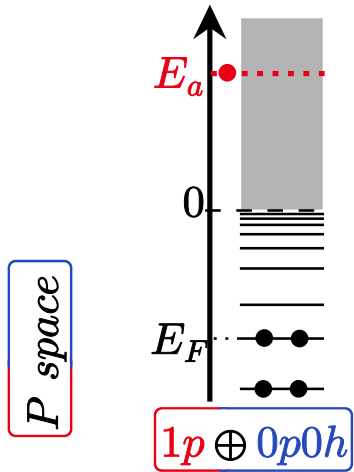
$$\begin{aligned}
 H^N &= \sum_{i=1}^N t(x_i) + V^N \\
 &= \underbrace{\sum_{i=1}^N \overbrace{[t(x_i) + v_0(x_i)]}^{h_0(x_i)}}_{\text{IPM term}} + \underbrace{V^N - \sum_{i=1}^N v_0(x_i)}_{\text{residual interaction}} \\
 &= H_{\text{IPM}}^N + V_{\text{IPM}}^N
 \end{aligned}$$

distribute  $N$  single-particles  
in  $\text{Sp}(h_0(x_i)) \rightarrow$



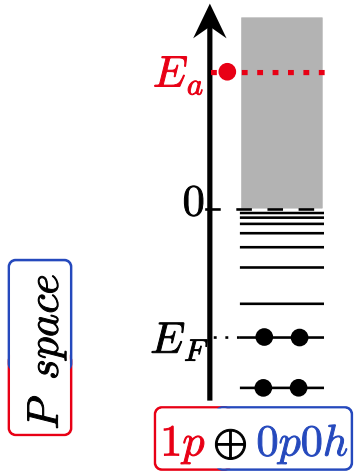
$$\begin{aligned}
 \left( \begin{array}{c} \mathcal{P}(E) \\ \xrightarrow{+ \text{ closed}} \end{array} \right) \mathcal{P} &= \{ \text{scattering states of } H_{\text{IPM}}^N \} \\
 \left( \begin{array}{c} \mathcal{Q}(E) \\ \xrightarrow{- \text{ closed}} \end{array} \right) \mathcal{Q} &= \{ \text{bound states of } H_{\text{IPM}}^N \} = \{ \text{quasibound states} \}
 \end{aligned}$$

# Pre-equilibrium emission examples



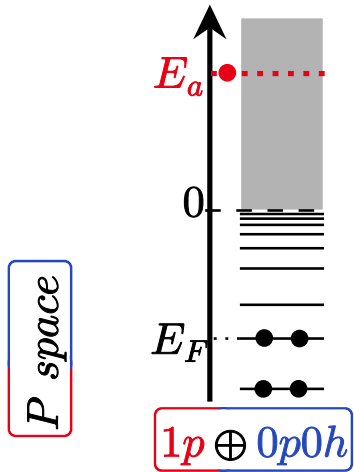
$$H^{N-1} + t(N)$$

# Pre-equilibrium emission examples



$$H^{N-1} + t(N) + V(N, \leq (N-1))$$

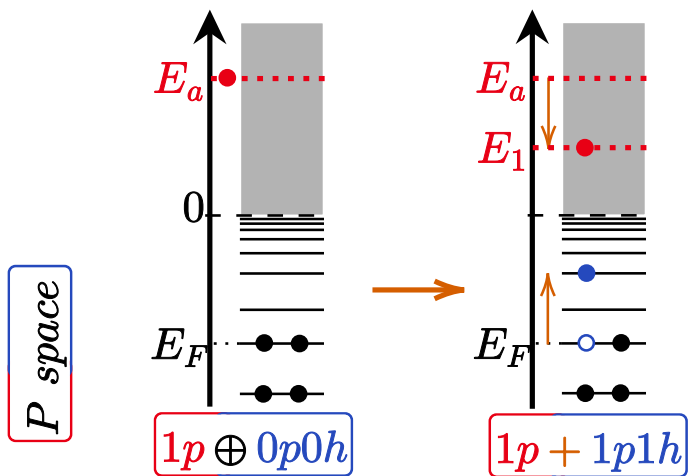
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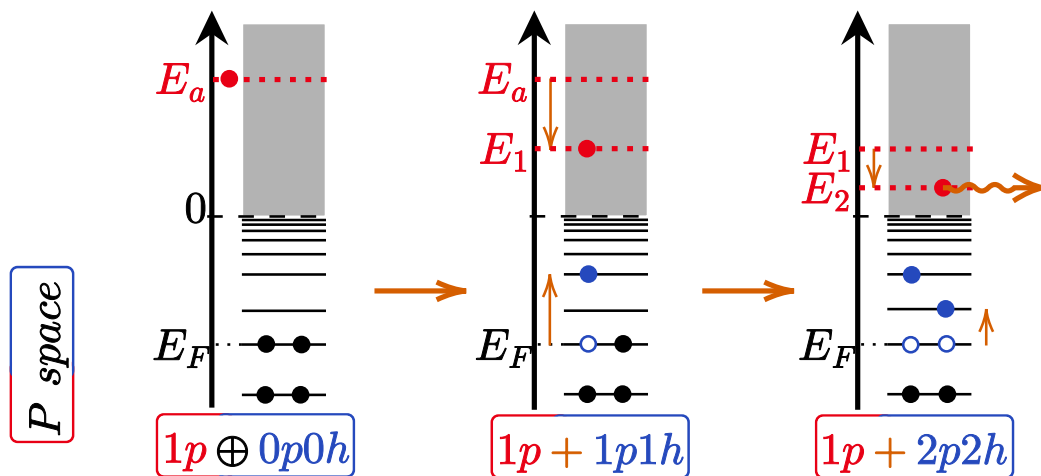
$$+ u_{\text{opt}}(N) - u_{\text{opt}}(N)$$

# Pre-equilibrium emission examples



$$H^{N-1} + \mathcal{H}_{\text{opt}}(N) + \mathcal{V}_{\text{opt}}$$

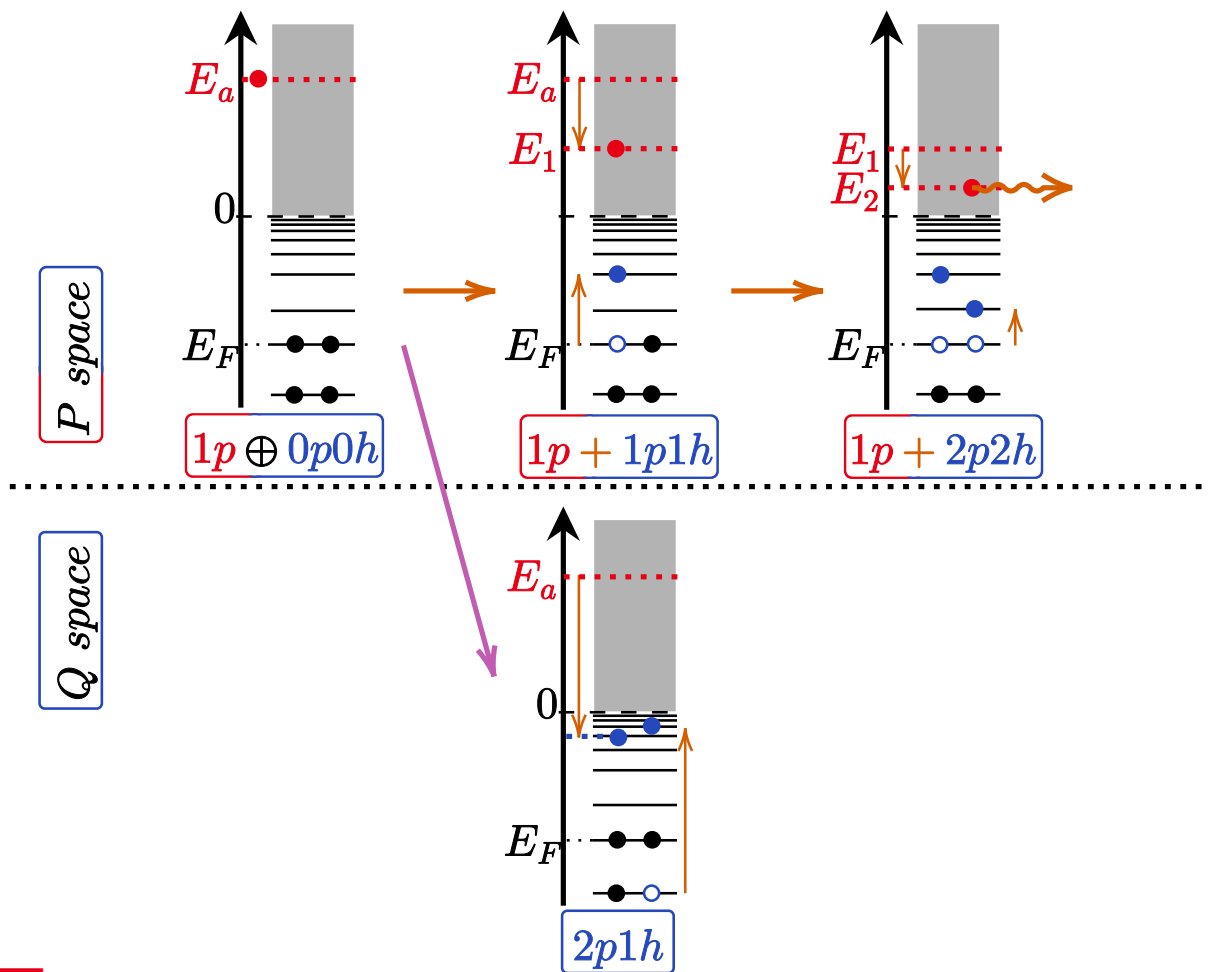
# Pre-equilibrium emission examples



$$H^{N-1} + \mathcal{H}_{\text{opt}}(N) + \mathcal{V}_{\text{opt}}$$

- only through P-space emission { } :  
 $\hookrightarrow$  Multi-Step Direct emission  
 MSD

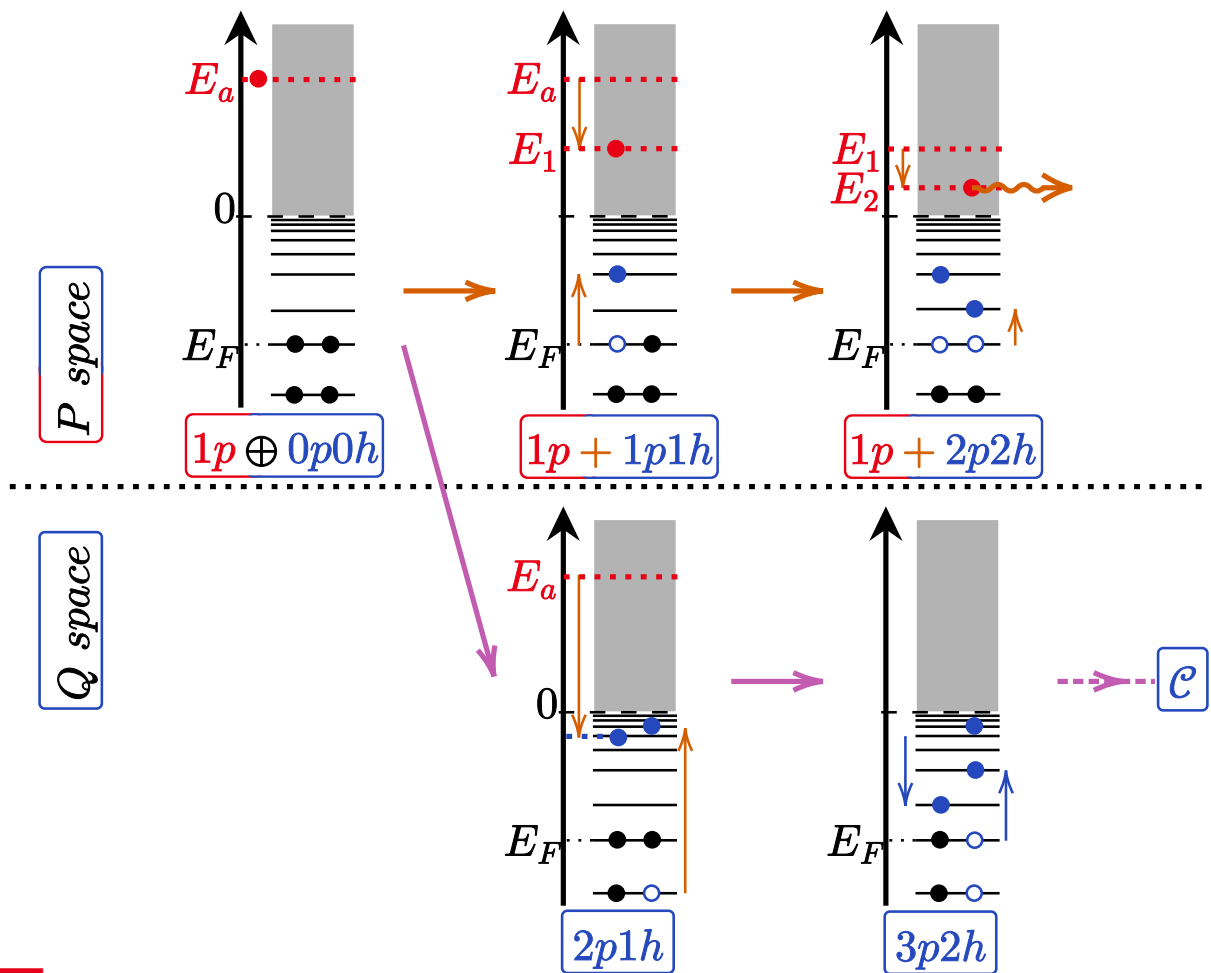
# Pre-equilibrium emission examples



$$H^{N-1} + \mathcal{H}_{\text{opt}}(N) + \mathcal{V}_{\text{opt}}$$

- only through P-space emission {wavy}:  
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**MSD**

# Pre-equilibrium emission examples

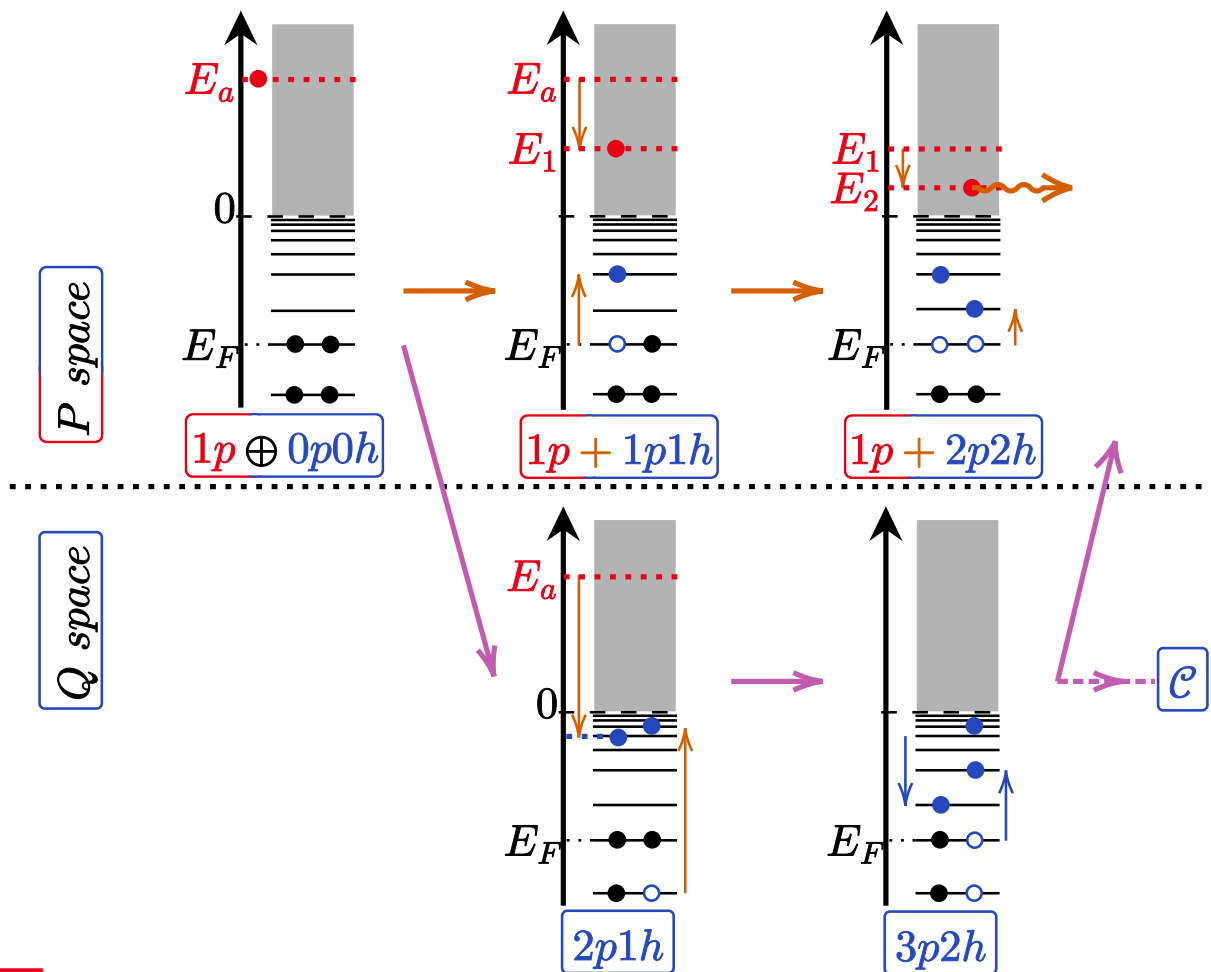


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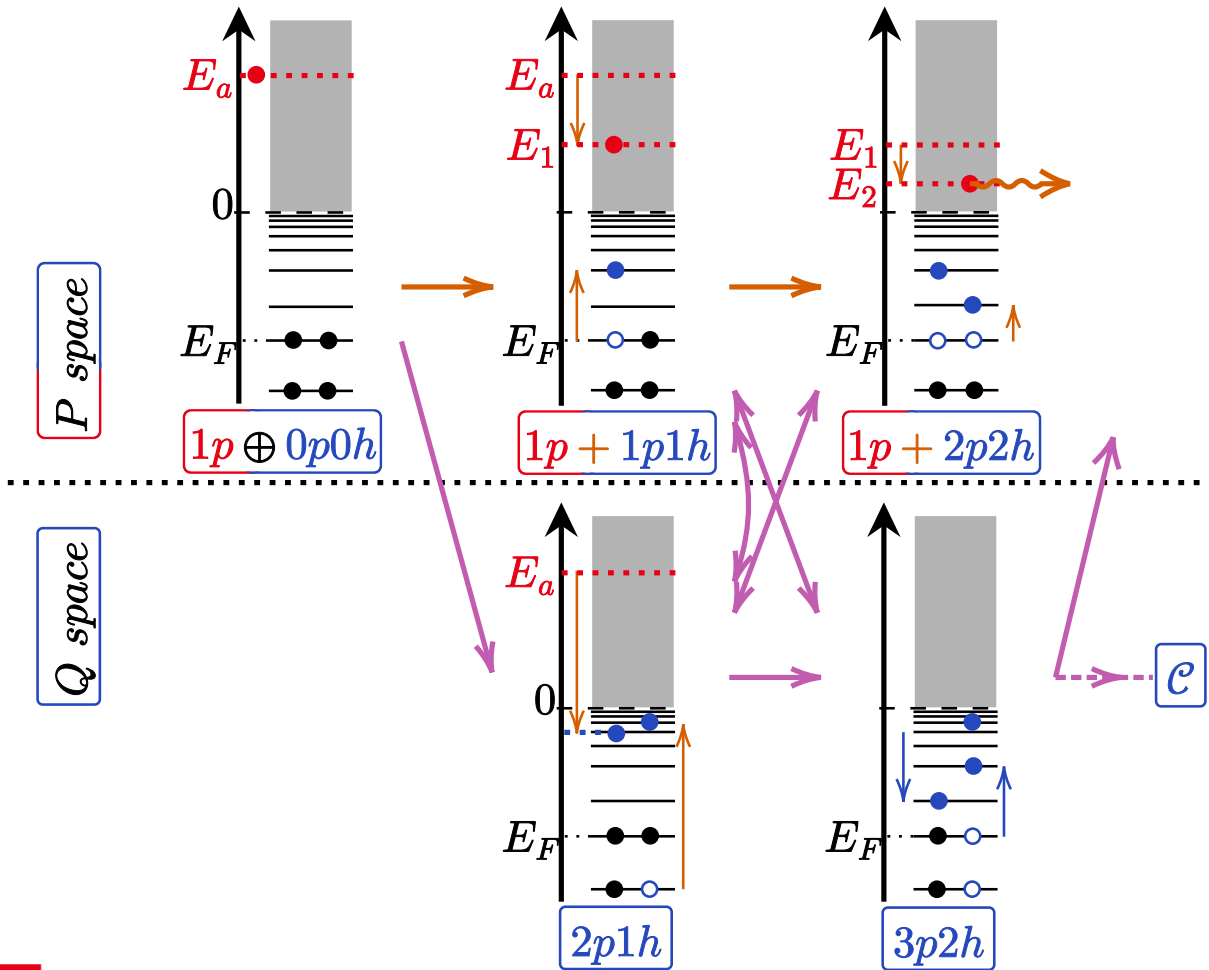
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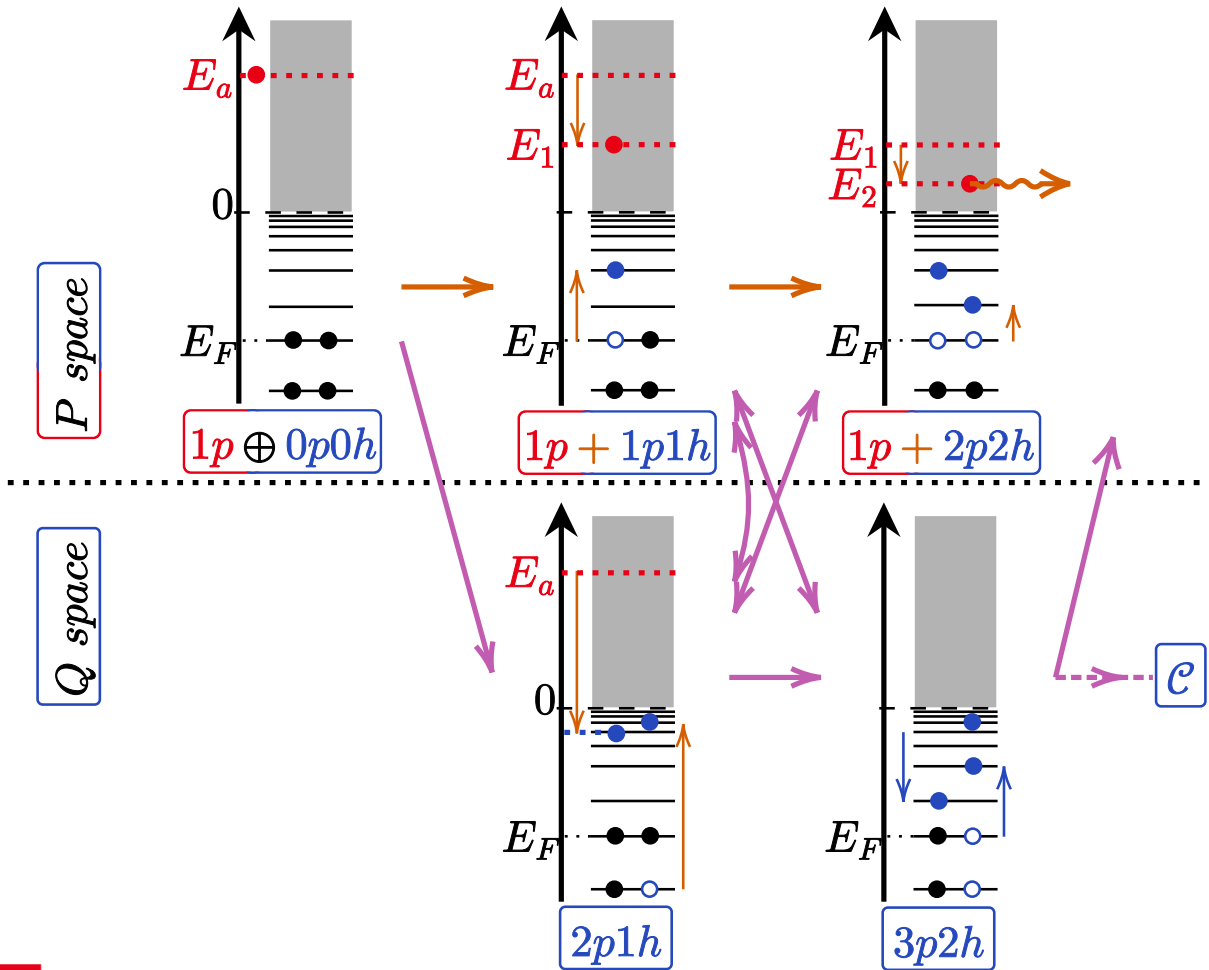
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- $(P \rightarrow Q) = \{\swarrow, \downarrow, \searrow\}$  :  
 $\hookrightarrow$  **gradual absorption**

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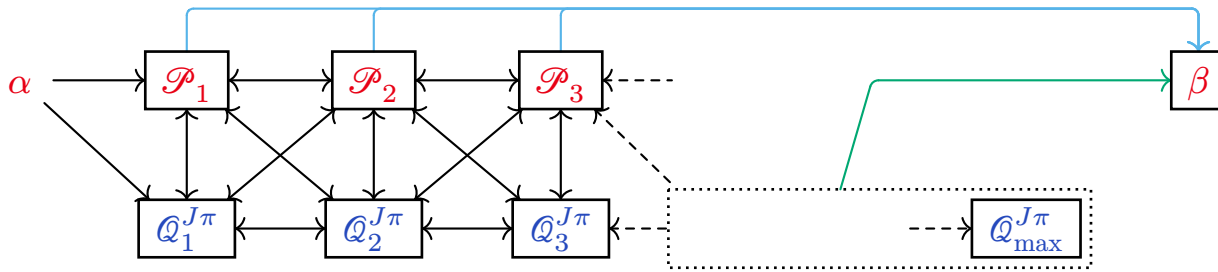
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# Microscopic Projectors



one GOE  $\rightarrow$  many GOEs : how do they couple ?

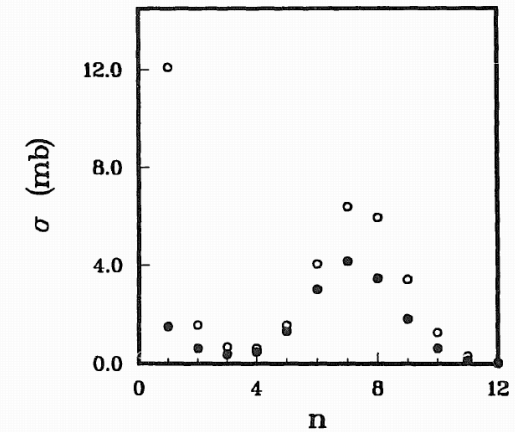
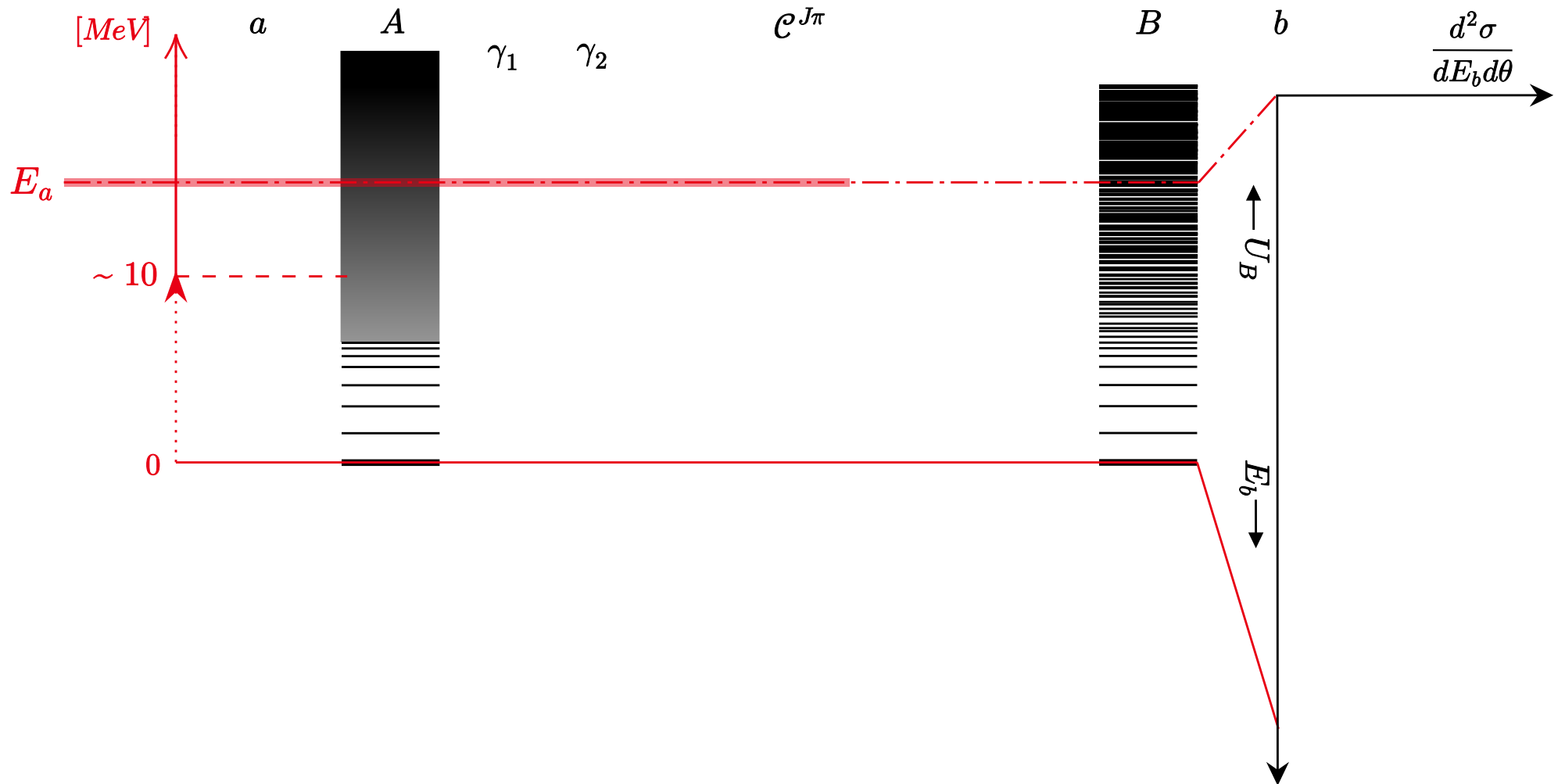
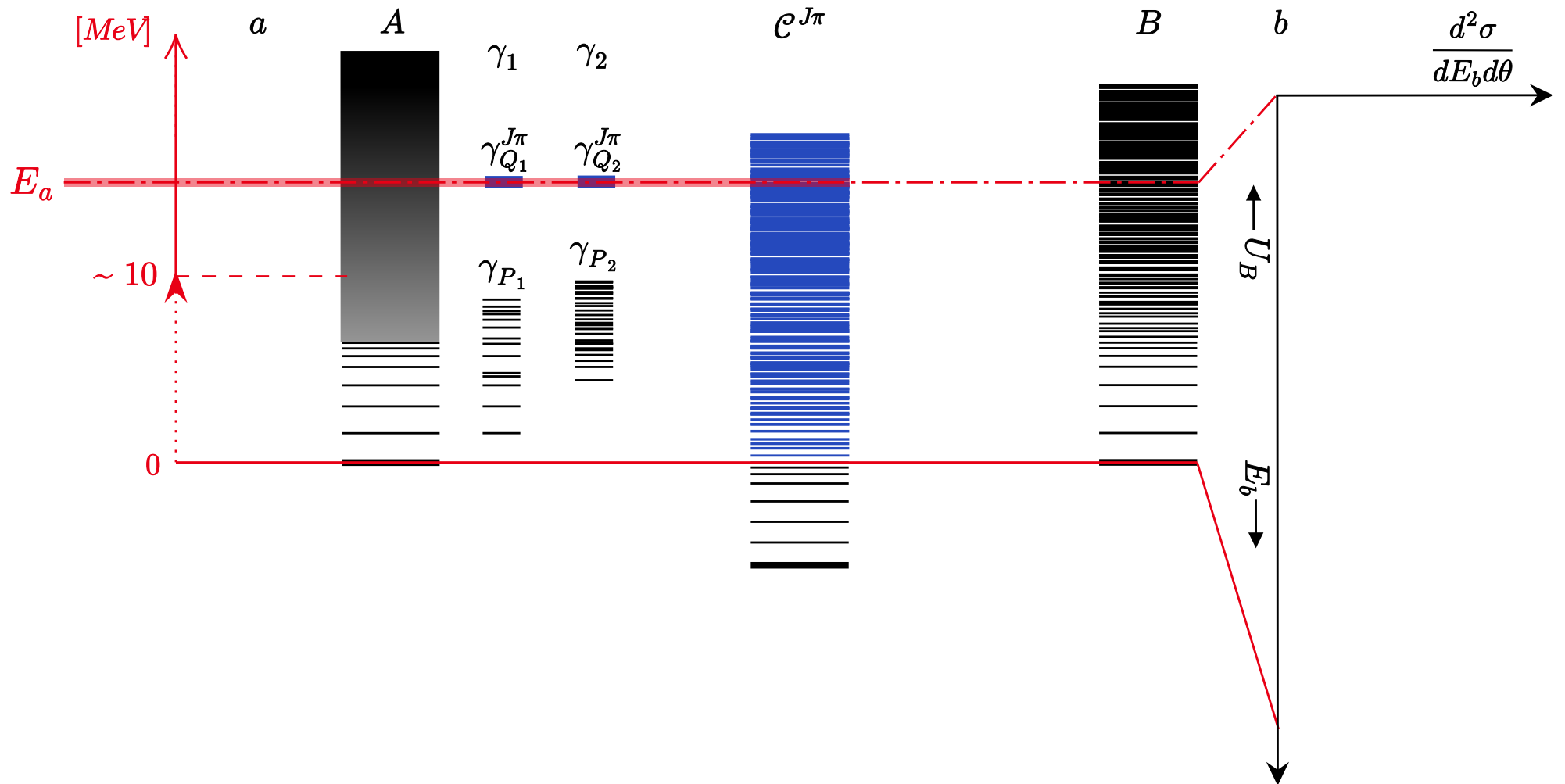


Fig. 4. Total (energy integrated) neutron emission as a function of the class number for two different values of the composite nucleus spin ( $J^\pi = 2^+$  - solid dots,  $J^\pi = 12^+$  - open dots).

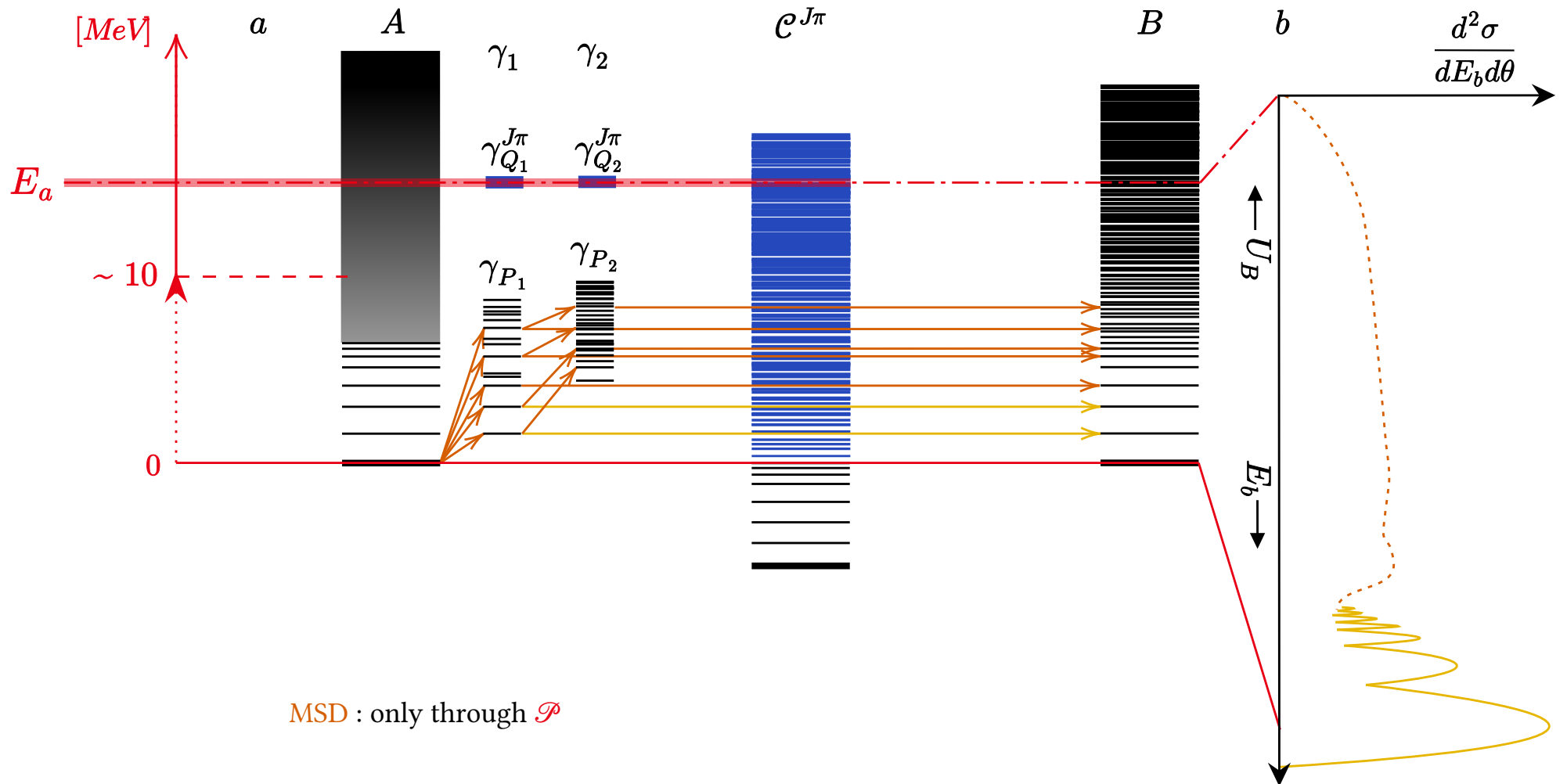
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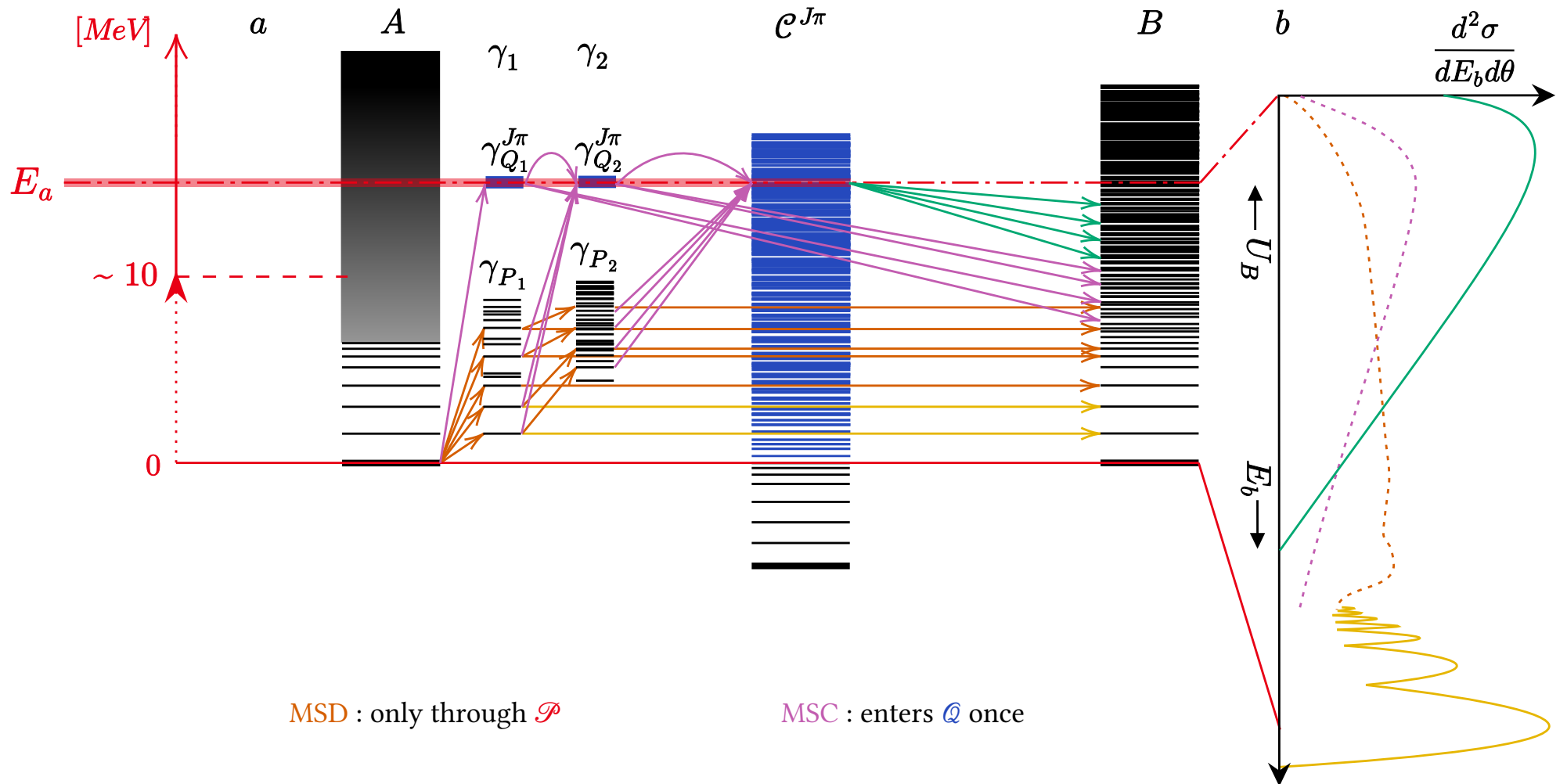


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MSD : only through  $\mathcal{P}$

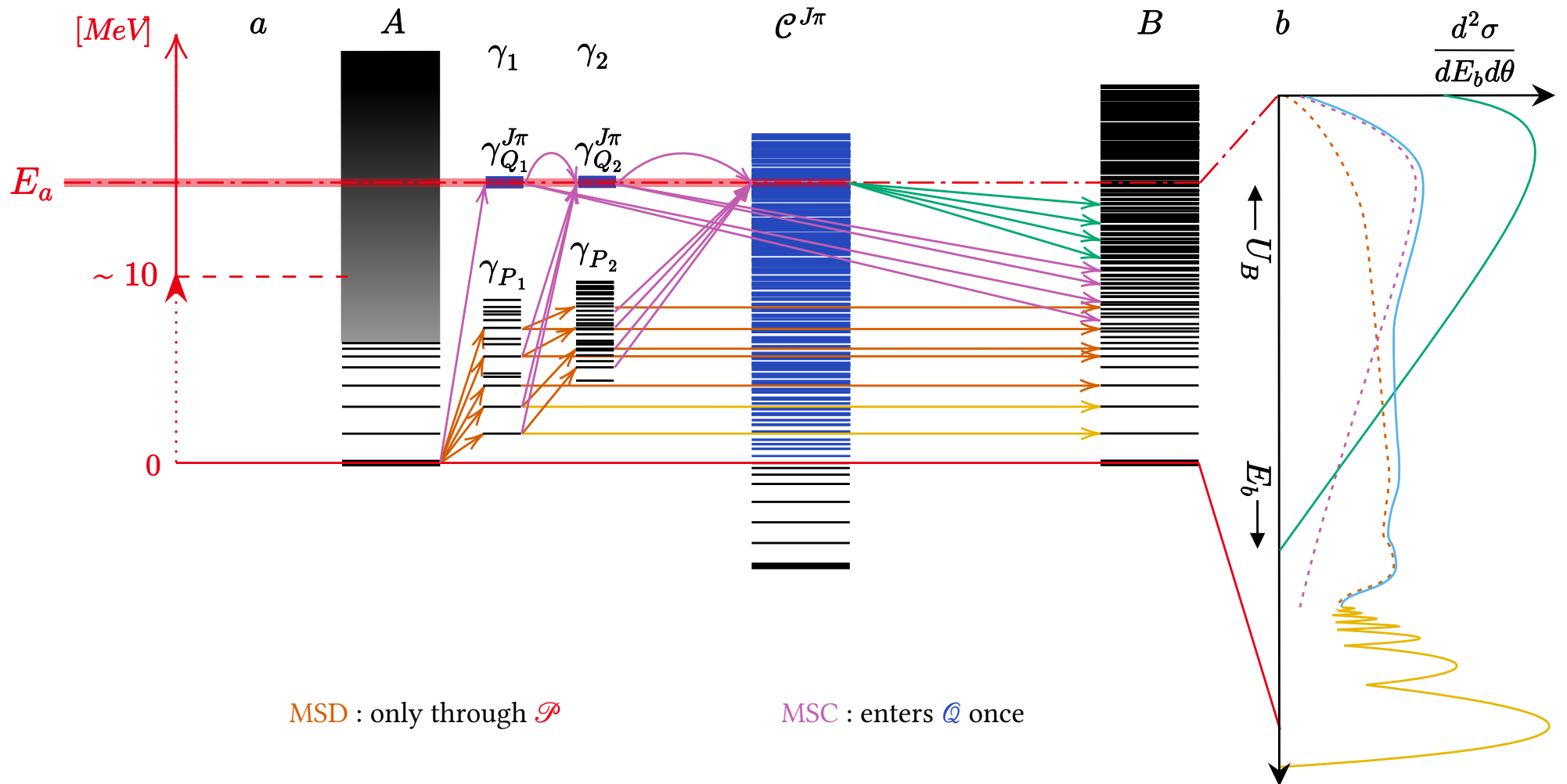
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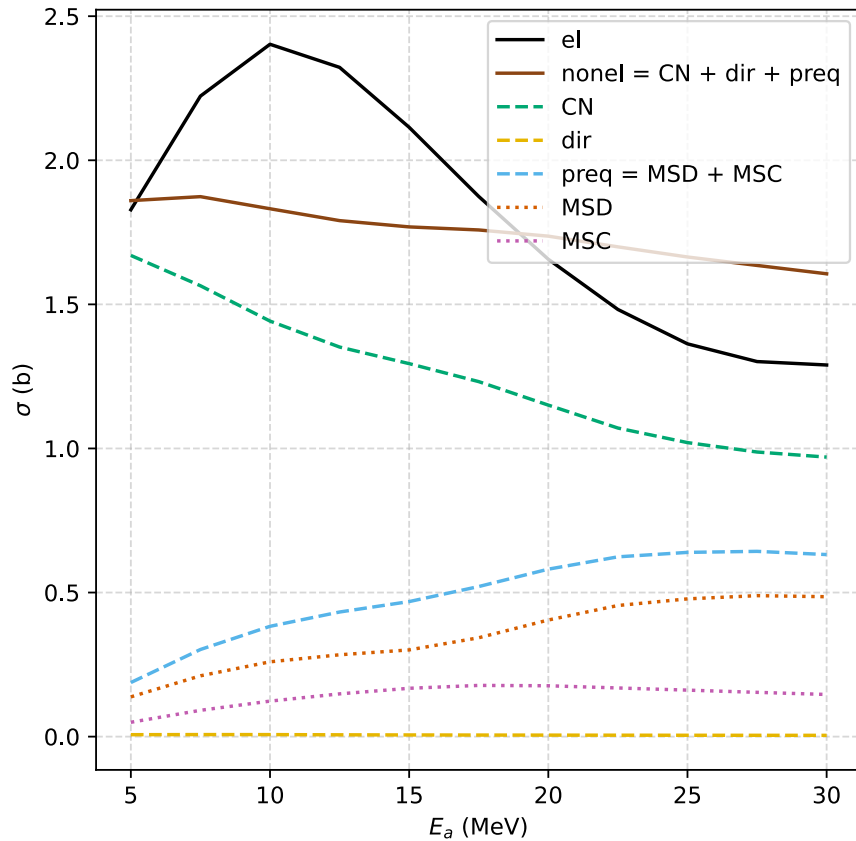
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# MSD & MSC contribution



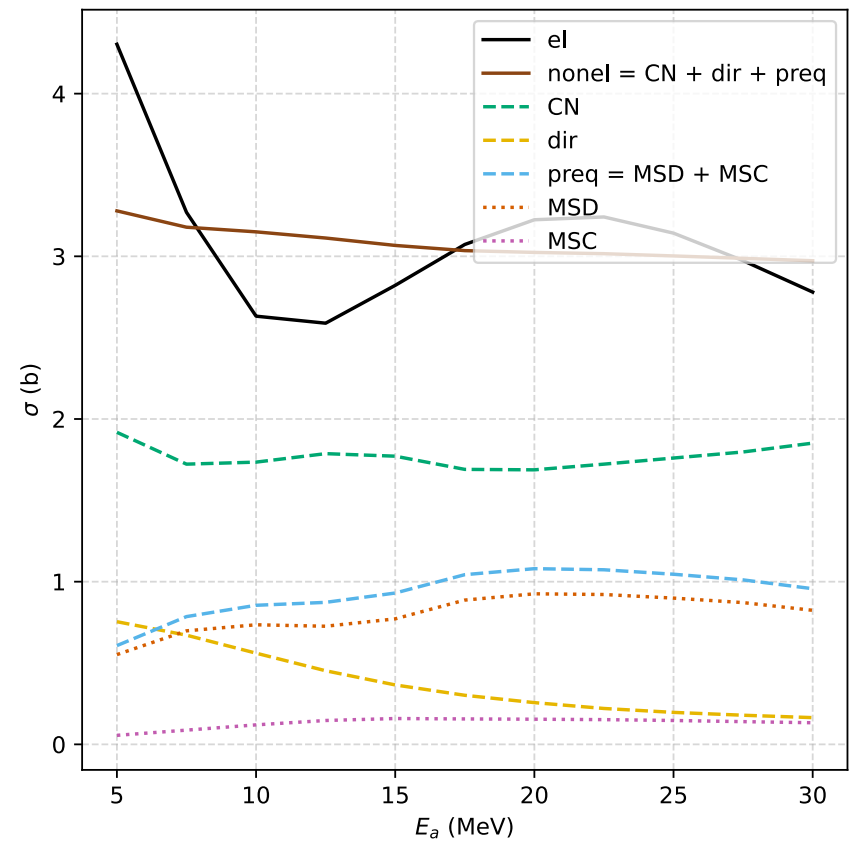
## $^{93}\text{Nb}(n,xn)$ pre-equilibrium and direct first chance emission

Results of EMPIRE with : MSD & MSC pre-equilibrium model



## $^{238}\text{U}(n,xn)$ pre-equilibrium and direct first chance emission

Results of EMPIRE with : MSD & MSC pre-equilibrium model



# “Heidelberg Group” Theory of Pre-Equilibrium

- CN (VWZ) [2]
- MSC (NVWY) [4]
  - weak-coupling limit (equivalent to AWM [5])
    - ↪ numerical implementation [6] in the weak-coupling limit
  - and the more general **strong-coupling** limit → systematically emphasized as more physical
    - ↪ **no implementation** (requires enhanced level densities)
- MSD (NWY) [7]
  - assuming  $\mathcal{V}_{\text{opt}}$  is 2-body type → **same description of  $N - 1$  nucleon system as  $N$**  → common inputs with NVWY [8]

$$\langle \sigma_{\alpha\beta} \rangle^{\Delta E} \propto \underbrace{\left| \langle S_{\alpha\beta} \rangle^{\Delta E} - \delta_{\alpha\beta} \right|^2}_{\sigma_{\alpha\beta}^{\text{MSD}}} + \sum_{J,\pi} \underbrace{\left\langle |S_{\alpha\beta}^{J\pi \text{fl}}|^2 \right\rangle^{\Delta E}}_{\sigma_{\alpha\beta}^{J\pi \text{MSC}}}$$

$$\langle X \rangle^{\Delta E} \rightarrow \bar{X}^N \quad \& \quad \langle X \rangle^{\Delta U_B} \rightarrow \bar{X}^{N-1}$$

# Summary & Challenges

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    - 1-SD : (Q)RPA for spherical & deformed nuclei
    - 2-SD :
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- **MSD challenges**

- $\overbrace{H_0^{N-1} + V_0^{N-1}}^{H^{N-1}} + \mathcal{R}_{\text{opt}}(N) + \mathcal{V}_{\text{opt}}$   
 $\hookrightarrow V_0^{N-1}$  vs  $\mathcal{V}_{\text{opt}}$  : MSD equilibration competition
- extended (Q)RPA models

Any question ?



# References

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- [1] F. Firk, J. Lynn, and M. Moxon, “Parameters of neutron resonances in U238 below 1.8 keV,” in *Proc. Kingston Intern. Conf. on Nuclear Structure*, University of Toronto Press, Toronto, 1960, pp. 757–759.
- [2] J. Verbaarschot, H. Weidenmüller, and M. Zirnbauer, “Grassmann integration in stochastic quantum physics: The case of compound-nucleus scattering,” *Physics Reports*, vol. 129, no. 6, pp. 367–438, 1985, doi: [https://doi.org/10.1016/0370-1573\(85\)90070-5](https://doi.org/10.1016/0370-1573(85)90070-5).
- [3] O. Bohigas and M.-J. Giannoni, “Chaotic motion and random matrix theories,” in *Mathematical and Computational Methods in Nuclear Physics*, J. S. Dehesa, J. M. G. Gomez, and A. Polls, Eds., Berlin, Heidelberg: Springer Berlin Heidelberg, 1984, pp. 1–99.
- [4] H. Nishioka, J. Verbaarschot, H. Weidenmüller, and S. Yoshida, “Statistical theory of precompound reactions: The multistep compound process,” *Annals of Physics*, vol. 172, no. 1, pp. 67–99, 1986, doi: [https://doi.org/10.1016/0003-4916\(86\)90020-5](https://doi.org/10.1016/0003-4916(86)90020-5).
- [5] D. Agassi, H. Weidenmüller, and G. Mantzouranis, “The statistical theory of nuclear reactions for strongly overlapping resonances as a theory of transport phenomena,” *Physics Reports*, vol. 22, no. 3, pp. 145–179, 1975, doi: [https://doi.org/10.1016/0370-1573\(75\)90028-9](https://doi.org/10.1016/0370-1573(75)90028-9).
- [6] M. Herman, G. Reffo, and H. Weidenmüller, “Multistep-compound contribution to precompound reaction cross section,” *Nuclear Physics A*, vol. 536, no. 1, pp. 124–140, 1992, doi: [https://doi.org/10.1016/0375-9474\(92\)90249-J](https://doi.org/10.1016/0375-9474(92)90249-J).
- [7] H. Nishioka, H. Weidenmüller, and S. Yoshida, “Statistical theory of precompound reactions: The multistep direct process,” *Annals of Physics*, vol. 183, no. 1, pp. 166–187, 1988, doi: [https://doi.org/10.1016/0003-4916\(88\)90250-3](https://doi.org/10.1016/0003-4916(88)90250-3).
- [8] T. Kawano and S. Yoshida, “Interference effect in the scattering amplitudes for nucleon-induced two-step direct process using the sudden approximation,” *Phys. Rev. C*, vol. 64, no. 2, p. 24603, Jun. 2001, doi: [10.1103/PhysRevC.64.024603](https://doi.org/10.1103/PhysRevC.64.024603).