



UNIVERSITAT DE  
BARCELONA

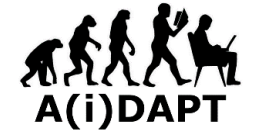


EXCELENCIA  
MARIA  
DE MAEZTU  
04/2025-03/2031

beatriu  
de pinós bp'



Jefferson Lab



AI for Data Analysis and Preservation

# Reconstructing scattering amplitudes with physics-informed generative models

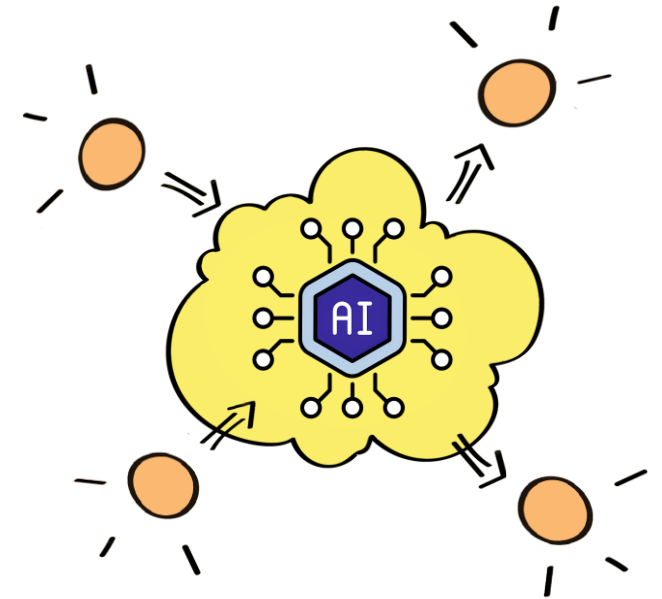
Glòria Montaña

University of Barcelona & ICCUB

In collaboration with J. Xu, Y. Li,  
A. Piloni, M. Battaglieri and others

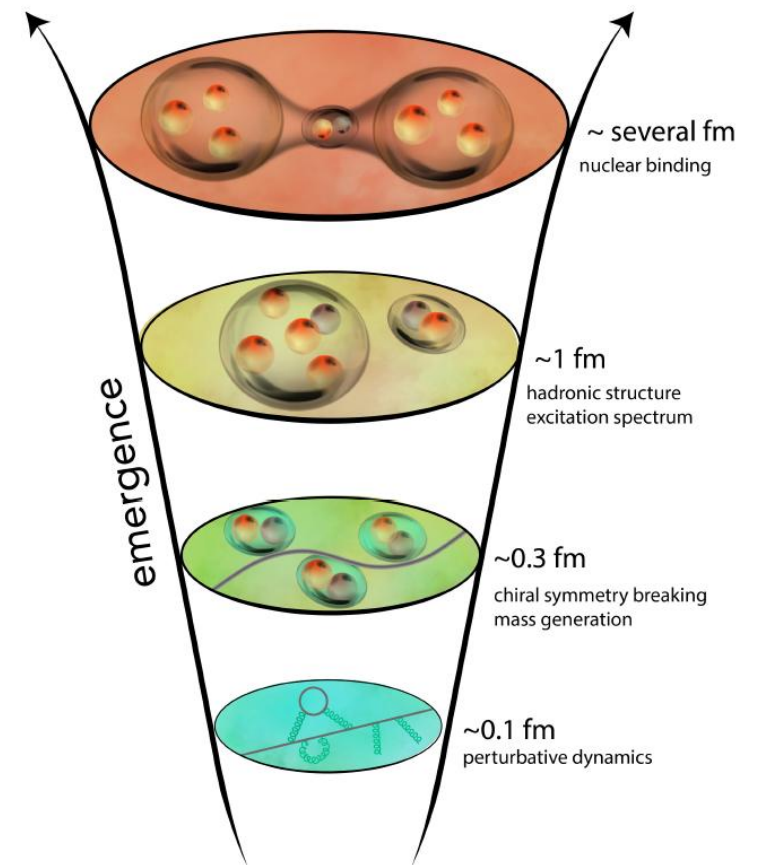
**Progress in algorithms and numerical tools for QCD**

Orsay, 9-10 June 2026



# The hadron spectrum

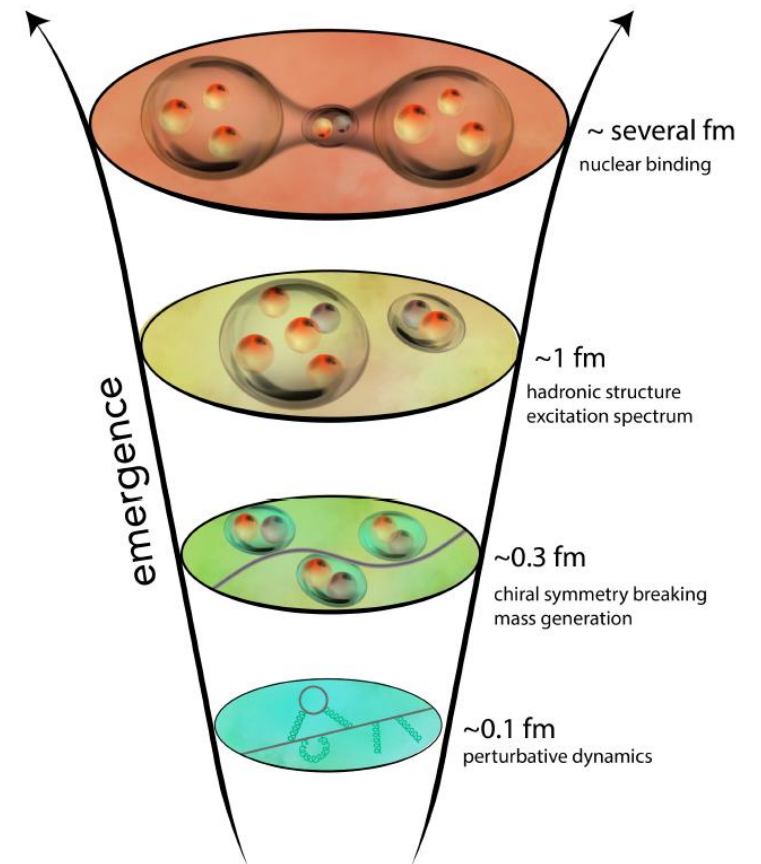
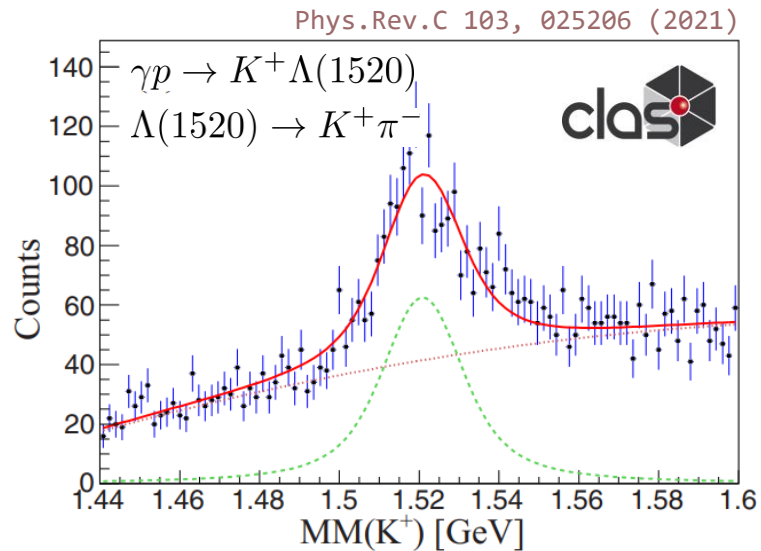
- \* Hadrons are an emergent phenomenon of QCD at low energies (non-perturbative regime)



Eur.Phys.J.A 60, 173 (2024)

# The hadron spectrum

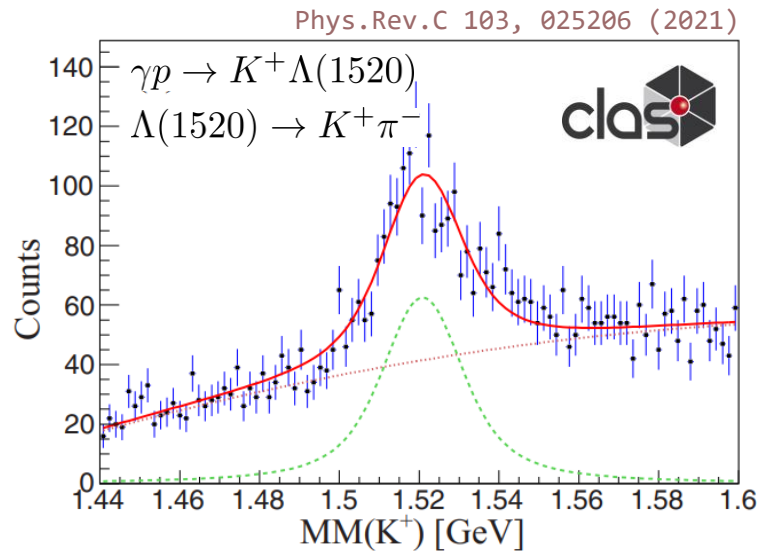
- \* Hadrons are an emergent phenomenon of QCD at low energies (non-perturbative regime)
- \* Most hadrons are short-lived **resonances**



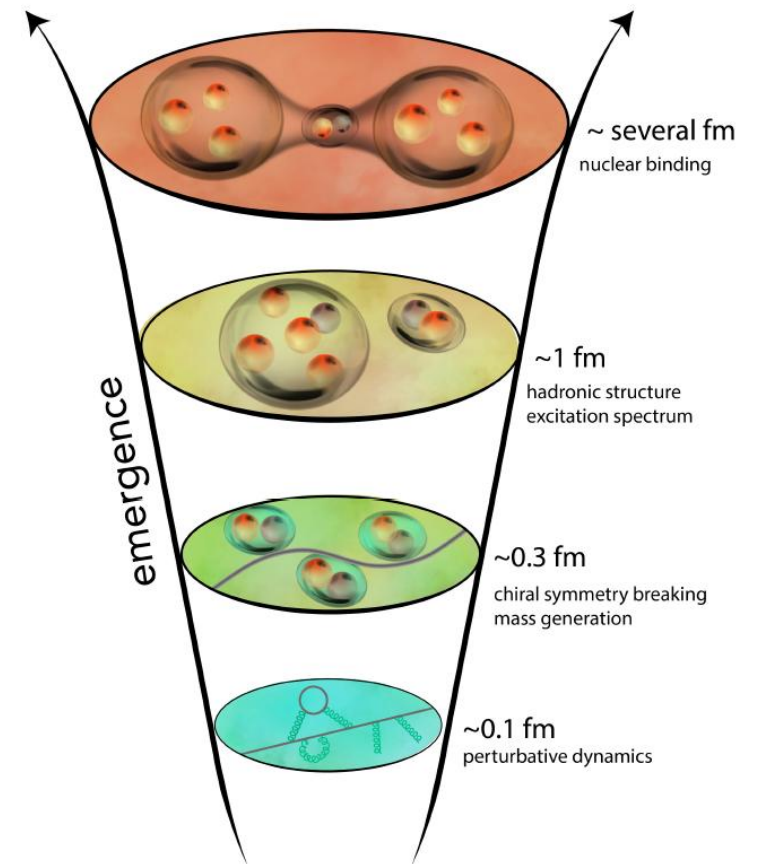
Eur.Phys.J.A 60, 173 (2024)

# The hadron spectrum

- \* Hadrons are an emergent phenomenon of QCD at low energies (non-perturbative regime)
- \* Most hadrons are short-lived **resonances**

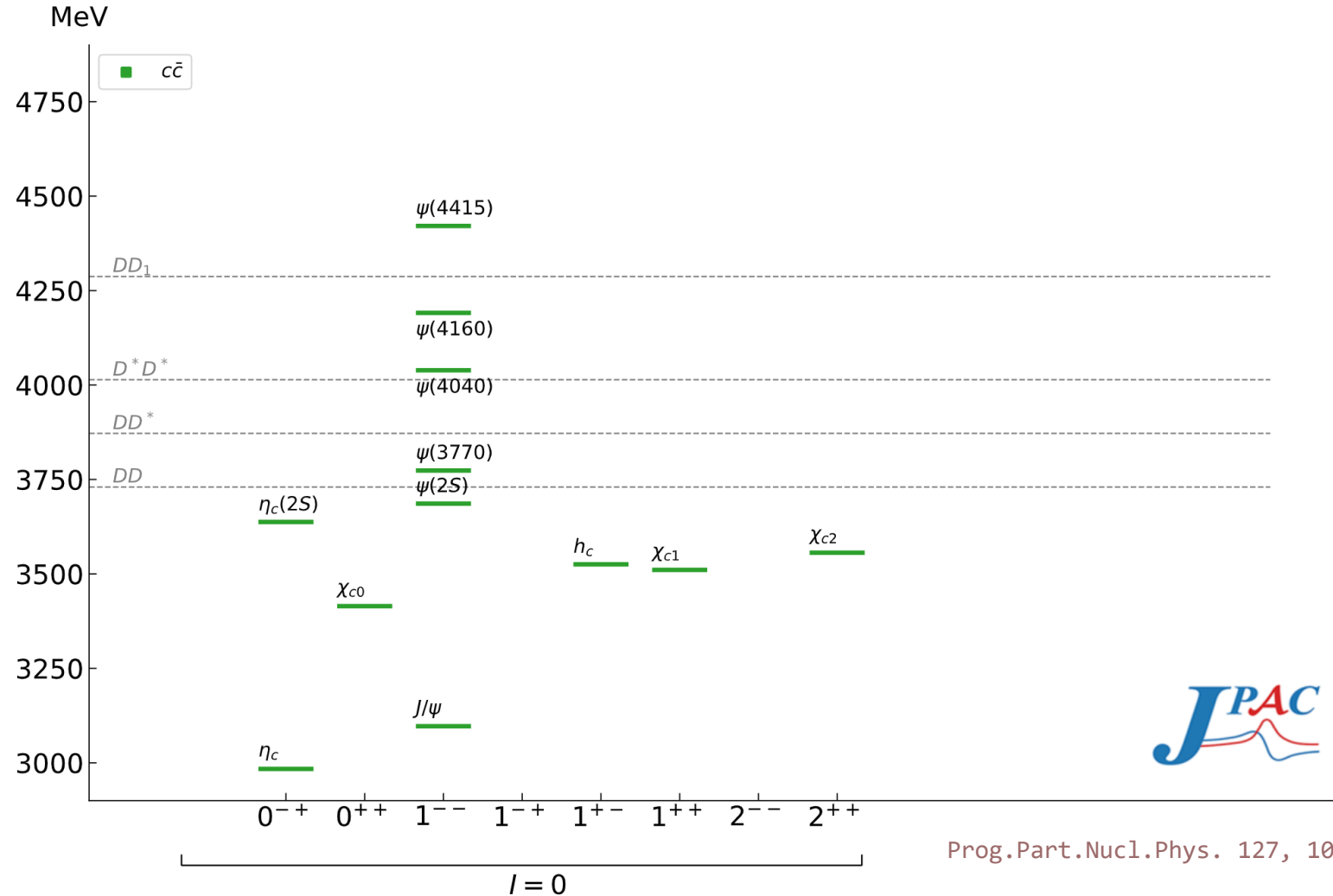


- Ultimate goals: Understand in terms of quarks and gluons  
Learn about QCD

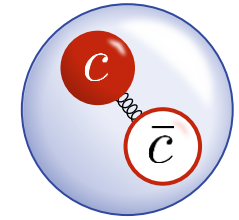


Eur.Phys.J.A 60, 173 (2024)

# The heavy hadron spectrum... before 2003



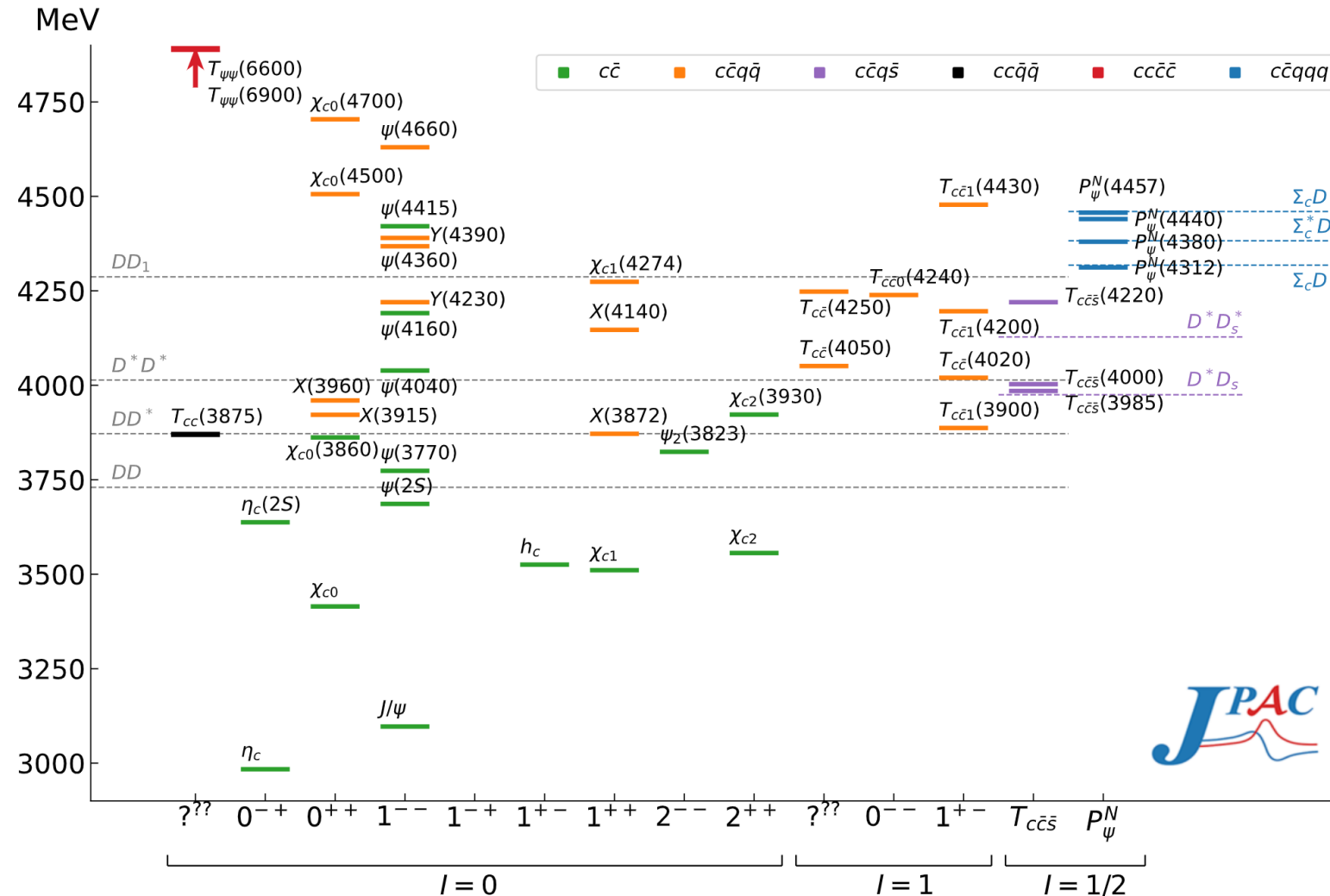
Prog.Part.Nucl.Phys. 127, 103981 (2022)



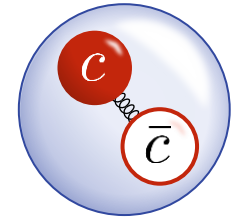
charmonium

The quark model organizes the spectrum of “conventional” hadrons

# The heavy hadron spectrum... after 2003

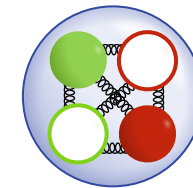


Prog.Part.Nucl.Phys. 127, 103981 (2022)

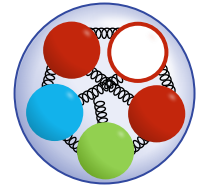


charmonium

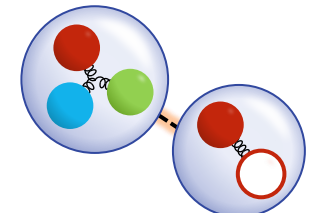
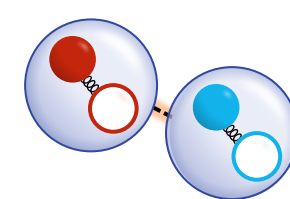
**Exotic hadrons:** don't fit in the "conventional" quark model



tetraquark



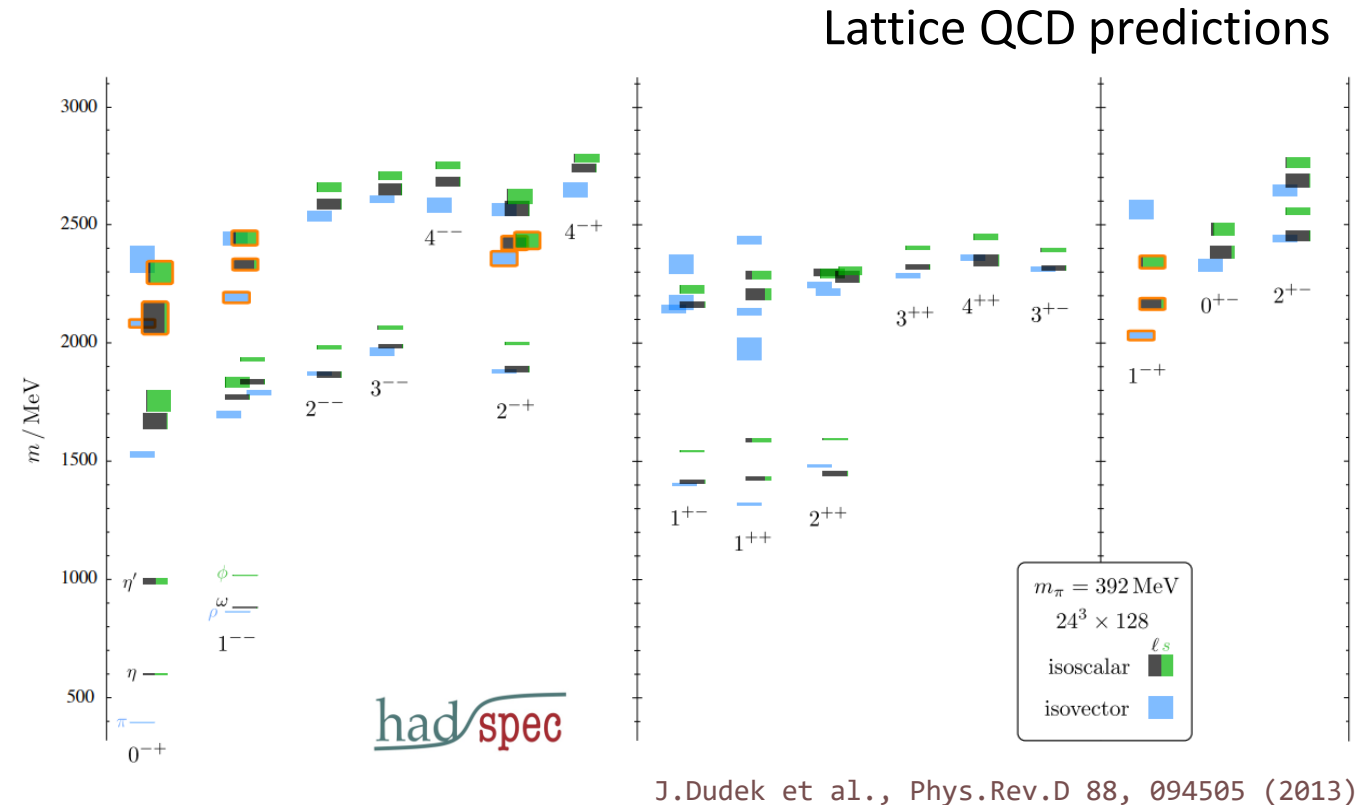
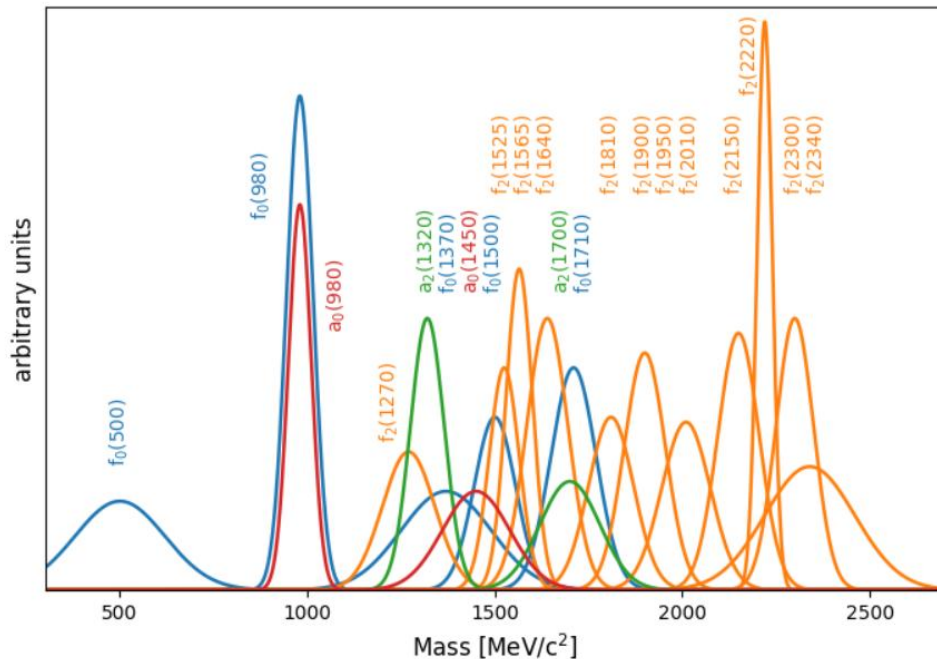
pentaquark



hadronic molecules

# Challenges in the light sector

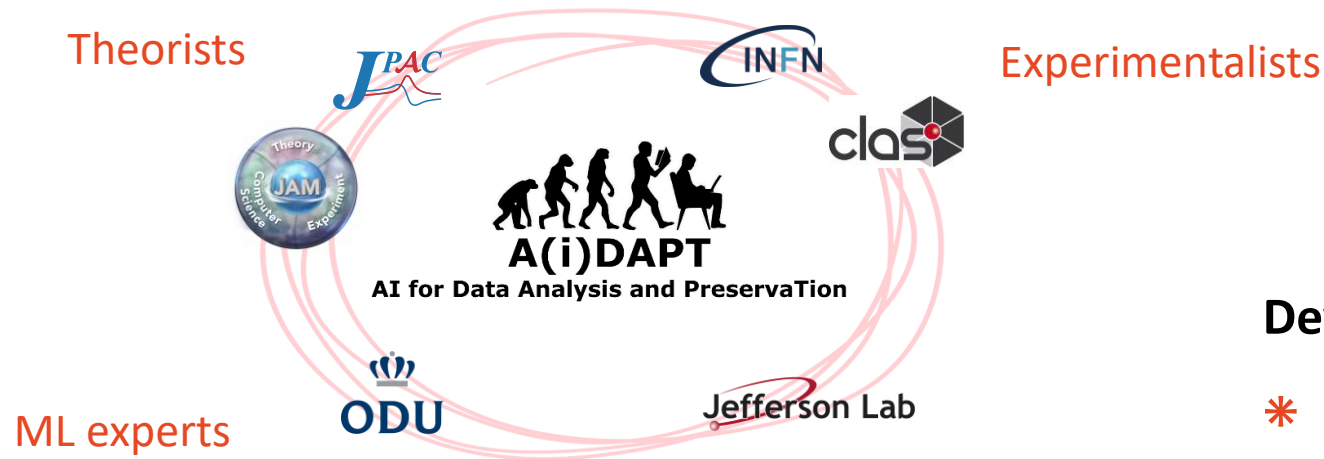
## Experimental observations (PDG)



- \* Wide, overlapping states
- \* Can't use simplistic resonance extraction techniques (e.g. Breit Wigner fits)
- \* Need for amplitude level analyses

# The A(i)DAPT collaborative effort

- Experimental data are distorted by detector effects → must be unfolded to extract physics
- Traditional observables may not be adequate in multidimensional space (correlations)
- Modern experiments produce large datasets, which can be difficult to manipulate/preserve



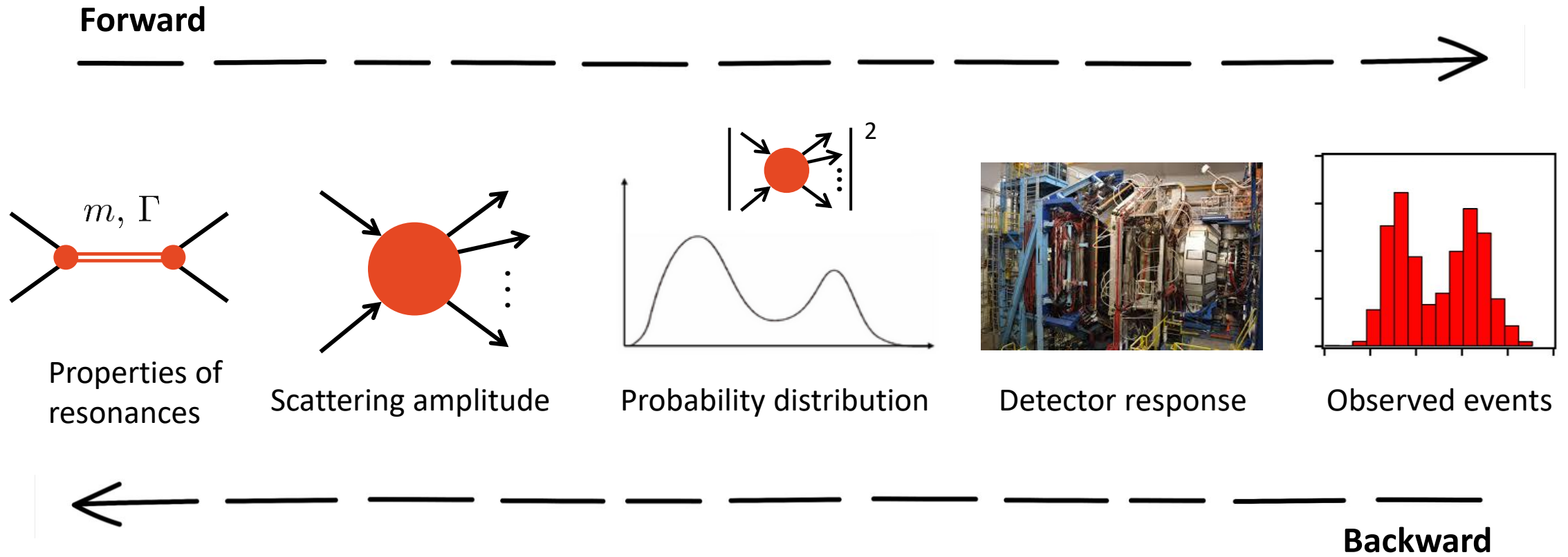
## Develop AI tools to:

- \* Accurately fit data in multiD space
- \* Unfold detector effects
- \* Quantify the uncertainties
- \* **Move from cross sections to amplitudes**

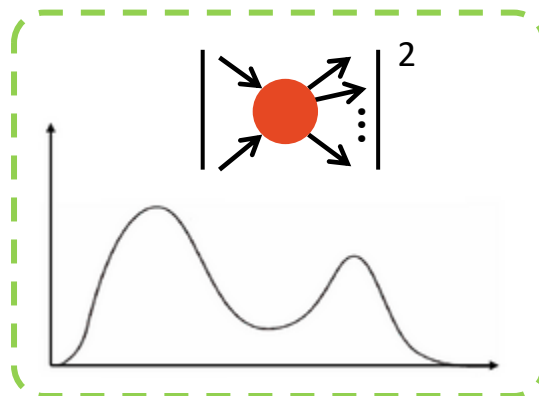
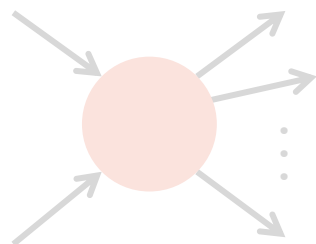
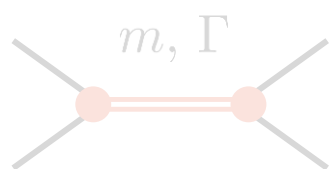
M. Battaglieri, Y. Li, A. Pilloni, N. Sato, A. Szczepaniak,  
D. Glazier, L. Bibrzycki, G. Montana, G. Foti, J. Xu,  
T. Alghamdi, T. Vittorini, and others



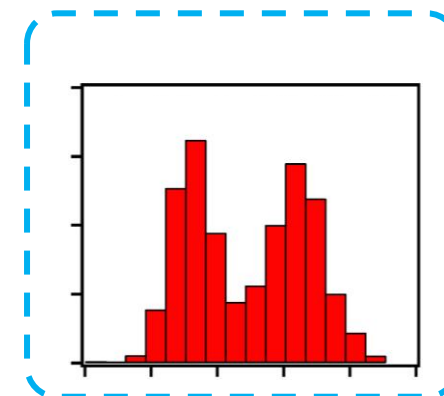
# From data to amplitudes and resonance properties



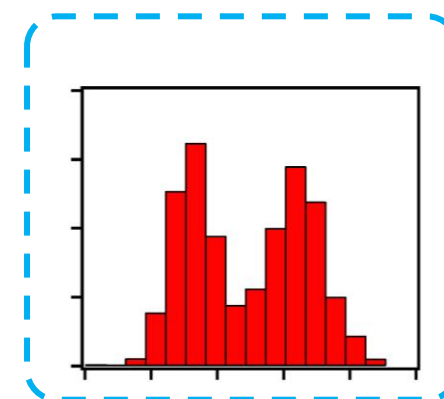
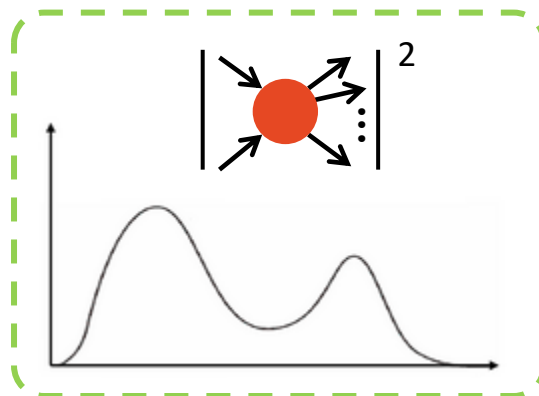
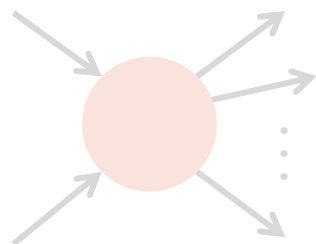
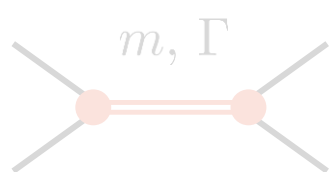
# Unfolding of detector effects



Unfold  
detector effects



# Unfolding of detector effects



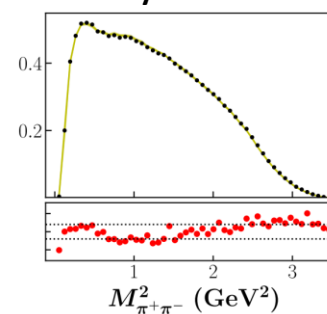
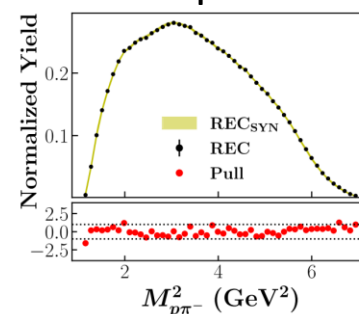
Unfold  
detector effects

\* Smearing (finite detector resolution)

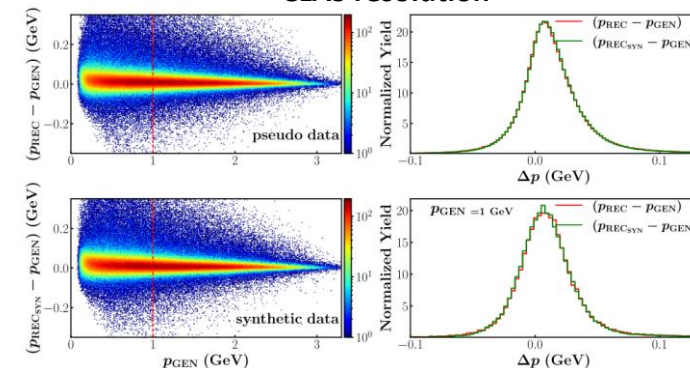
## Closure test on pseudodata of CLAS exclusive $2\pi$ photoproduction

Alghamdi et al., Phys.Rev.D 108, 094030 (2023)

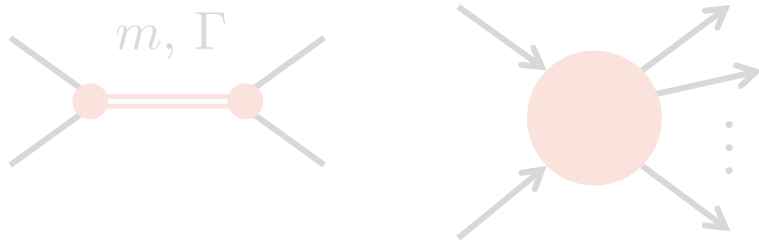
### MC REC pseudodata vs DS-GAN synthetic data



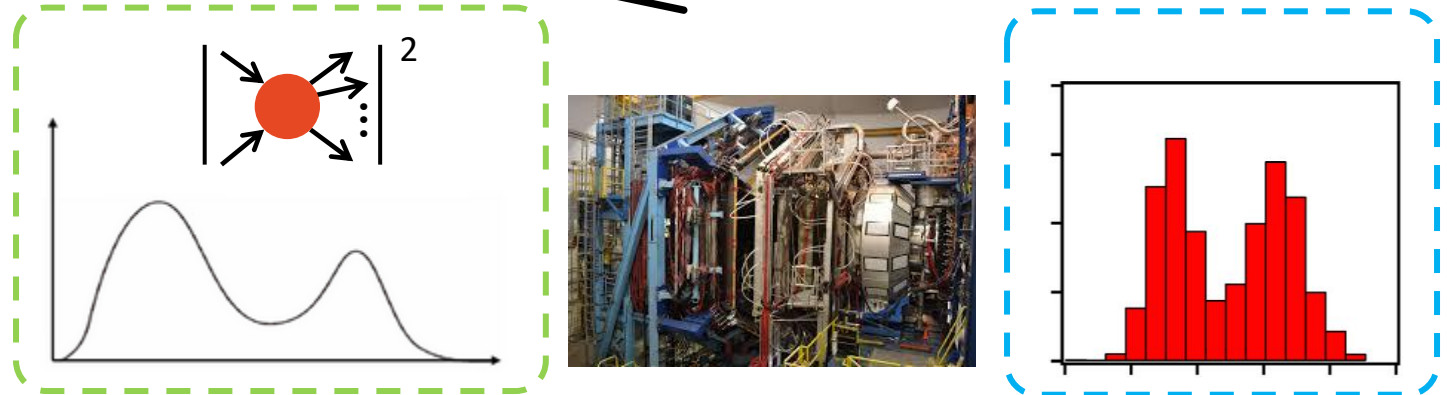
### CLAS resolution



# Unfolding of detector effects



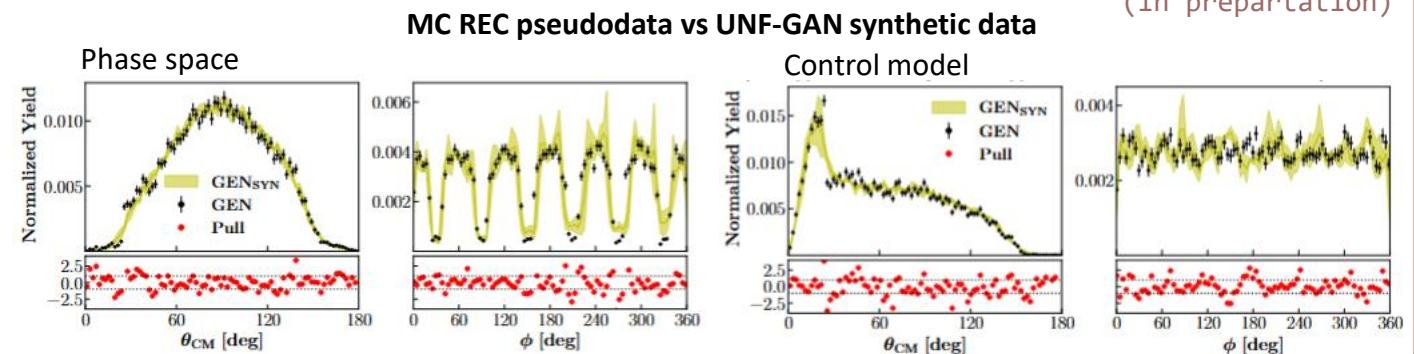
Unfold  
detector effects



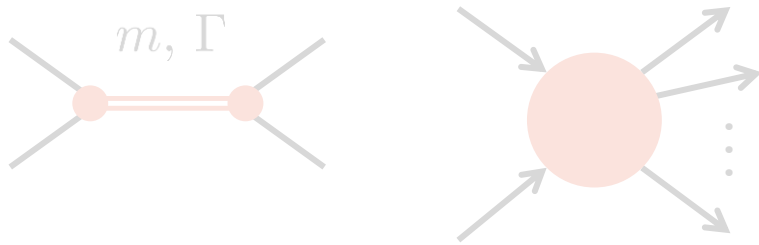
- \* Smearing (finite detector resolution)
- \* Acceptance (detector coverage)
- \* Inefficiency (missed events within acceptance)

## Unfolding of CLAS detector in $\pi^0$ photoproduction

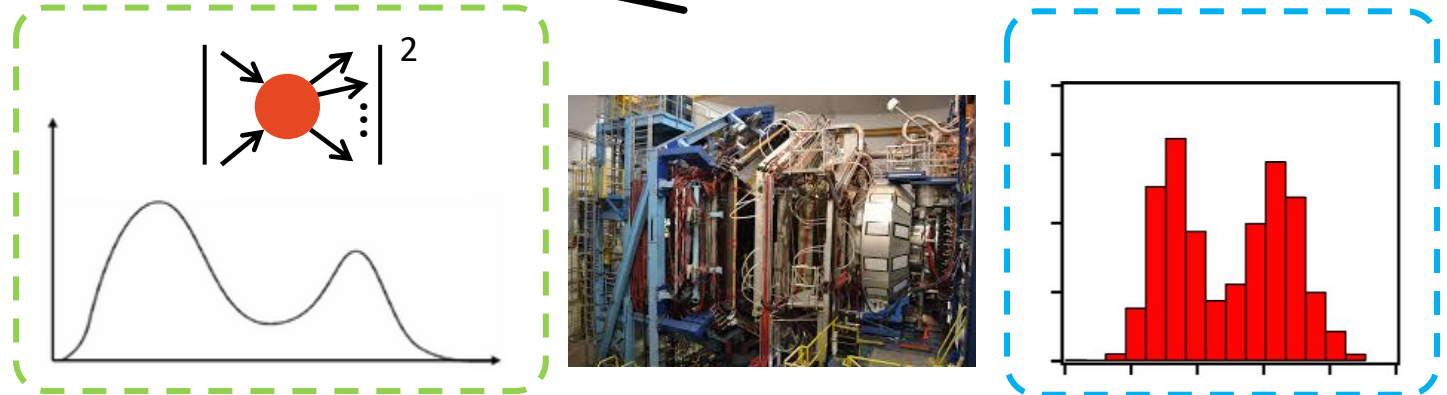
Alghamdi et al.,  
(in preparation)



# Unfolding of detector effects



Unfold  
detector effects

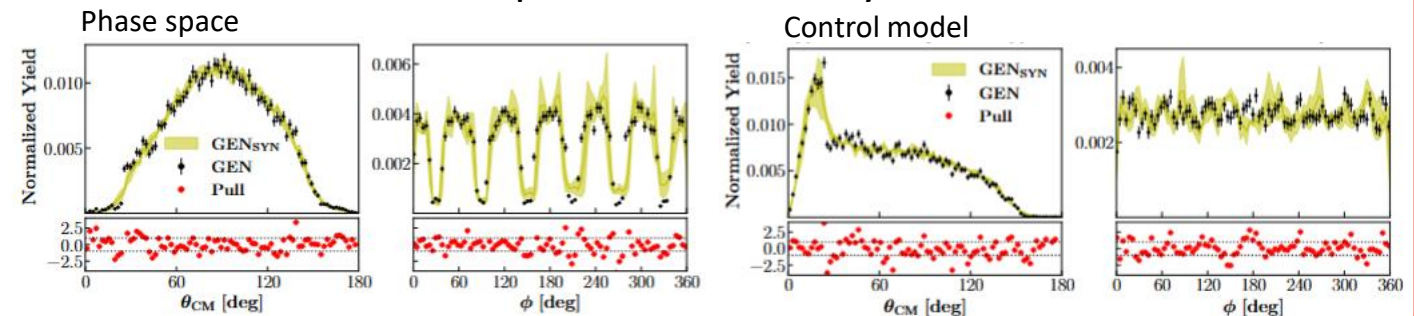


- \* Smearing (finite detector resolution)
- \* Acceptance (detector coverage)
- \* Inefficiency (missed events within acceptance)

## Unfolding of CLAS detector in $\pi^0$ photoproduction

Alghamdi et al.,  
(in preparation)

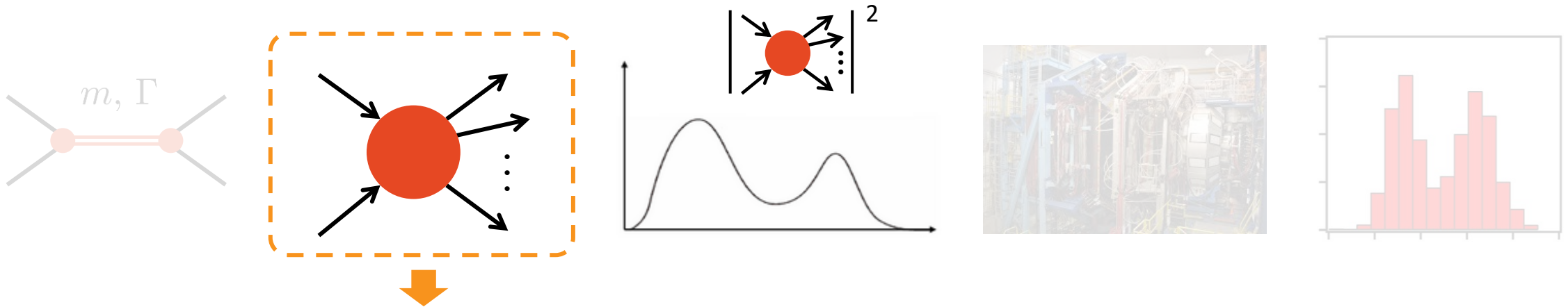
MC REC pseudodata vs UNF-GAN synthetic data



**Diffusion Model** for CLAS detector unfolding in photoproduction of  $\pi^+\pi^-$

G. Foti

# Amplitude level unfolding



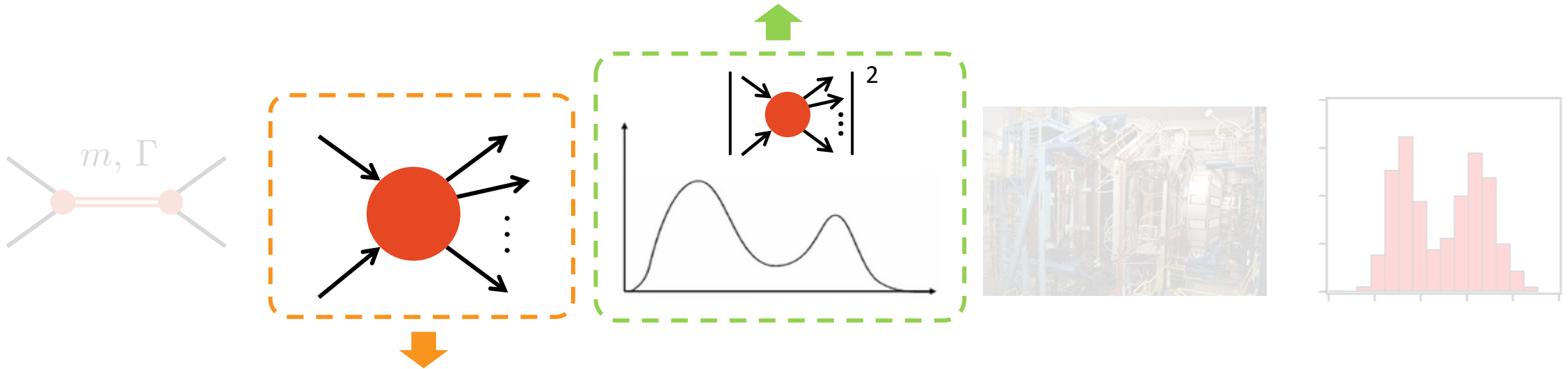
The **scattering amplitude**  $\mathcal{A}$  is the fundamental quantity of interest in hadron reactions:

- ▷ Encodes the underlying dynamics of the interaction
- ▷ Is a complex quantity: magnitude + phase

# Amplitude level unfolding

The **cross section**  $\sigma$  is an experimentally observable quantity:

▷ Related to  $|\mathcal{A}|^2$  (phase information is lost)



The **scattering amplitude**  $\mathcal{A}$  is the fundamental quantity of interest in hadron reactions:

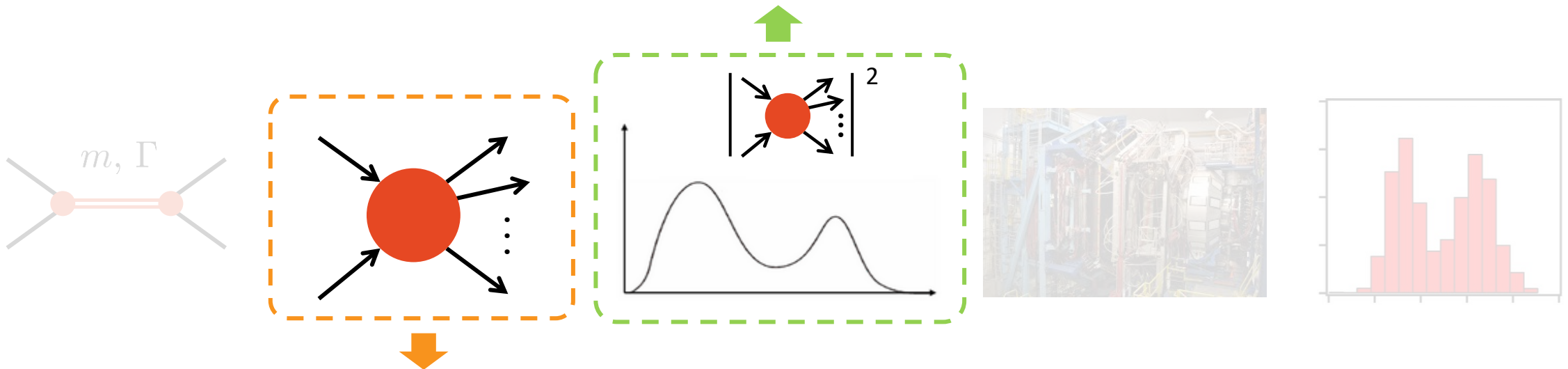
- ▷ Encodes the underlying dynamics of the interaction
- ▷ Is a complex quantity: magnitude + phase



# Amplitude level unfolding

The **cross section**  $\sigma$  is an experimentally observable quantity:

▷ Related to  $|\mathcal{A}|^2$  (phase information is lost)



The **scattering amplitude**  $\mathcal{A}$  is the fundamental quantity of interest in hadron reactions:

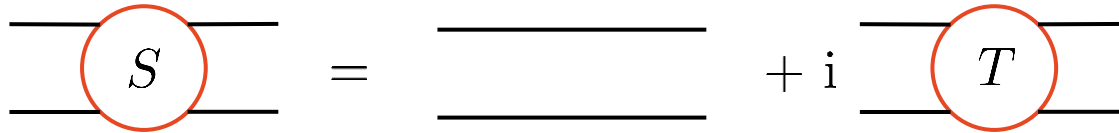
- ▷ Encodes the underlying dynamics of the interaction
- ▷ Is a complex quantity: magnitude + phase

→ Ill-posed inverse problem



# Properties of scattering amplitudes

## S-matrix theory



The diagram shows a horizontal line with two parallel lines above and below it, representing an incoming state. A red circle labeled  $S$  is placed on this line. This is followed by an equals sign, then a horizontal line with two parallel lines above and below it, representing an outgoing state. This is followed by a plus sign, the letter  $i$ , another plus sign, and finally a red circle labeled  $T$  on a horizontal line with two parallel lines above and below it, representing an incoming state.

$$\text{---} \bigcirc_S \text{---} = \text{---} + i \text{---} \bigcirc_T \text{---}$$

# Properties of scattering amplitudes

**S-matrix theory**  $\longrightarrow$  Fundamental principles of scattering

\* Rotational invariance  $\rightarrow$  **Partial wave expansion**

$$\langle \text{out} | T | \text{in} \rangle \equiv A(s, \cos \theta) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}(s) P_{\ell}(\cos \theta)$$

$$\Sigma \quad \begin{array}{c} \diagup \quad \text{---} \quad \diagdown \\ \diagdown \quad \text{---} \quad \diagup \end{array}$$

$$f_{\ell} \sim \frac{1}{s - m_R^2}$$

# Properties of scattering amplitudes

**S-matrix theory**  $\longrightarrow$  Fundamental principles of scattering

\* Rotational invariance  $\rightarrow$  **Partial wave expansion**

$$\langle \text{out} | T | \text{in} \rangle \equiv A(s, \cos \theta) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}(s) P_{\ell}(\cos \theta)$$

\* Probability conservation  $\rightarrow$  **Unitarity**

$$SS^{\dagger} = 1 \rightarrow \text{Im} f_{\ell}(s) = |f_{\ell}(s)|^2$$

# Properties of scattering amplitudes

**S-matrix theory**  $\longrightarrow$  Fundamental principles of scattering

\* Rotational invariance  $\rightarrow$  **Partial wave expansion**

$$\langle \text{out} | T | \text{in} \rangle \equiv A(s, \cos \theta) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}(s) P_{\ell}(\cos \theta)$$

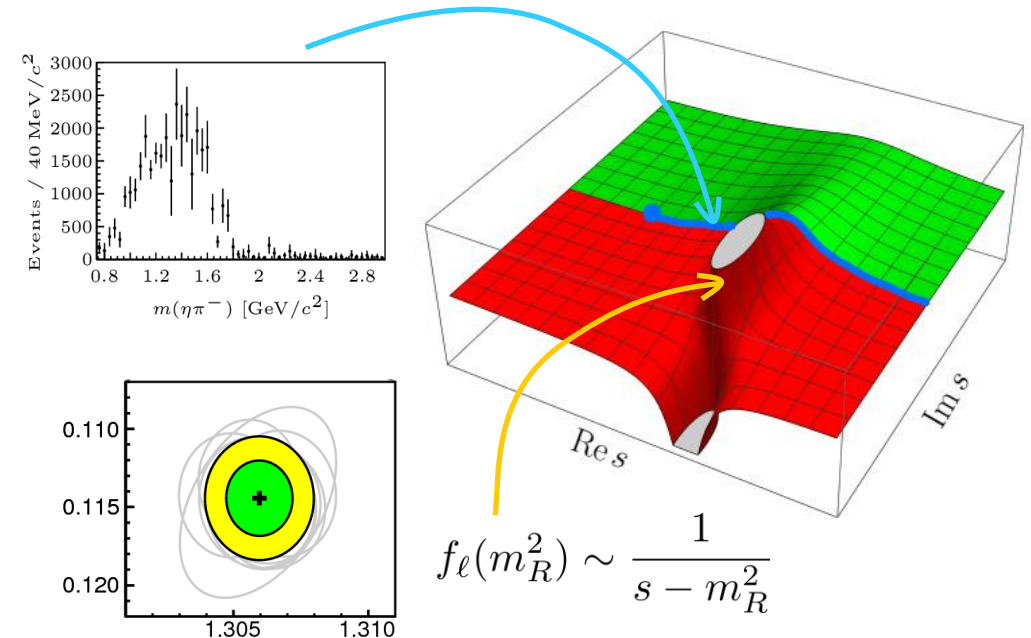
\* Probability conservation  $\rightarrow$  **Unitarity**

$$SS^{\dagger} = 1 \rightarrow \text{Im} f_{\ell}(s) = |f_{\ell}(s)|^2$$

\* Causality  $\rightarrow$  **Analyticity**

$$f_{\ell}(s) = \frac{1}{\pi} \int ds' \frac{\text{Disc } f_{\ell}(s')}{s' - s - i\epsilon}$$

Resonances are characterized by poles in the complex energy plane



# Properties of scattering amplitudes

**S-matrix theory**  $\longrightarrow$  Fundamental principles of scattering

\* Rotational invariance  $\rightarrow$  **Partial wave expansion**

$$\langle \text{out} | T | \text{in} \rangle \equiv A(s, \cos \theta) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}(s) P_{\ell}(\cos \theta)$$

\* Probability conservation  $\rightarrow$  **Unitarity**

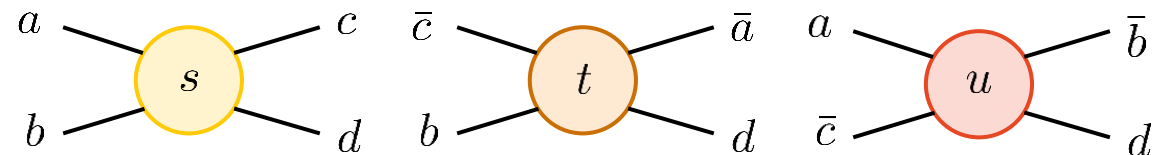
$$SS^{\dagger} = 1 \rightarrow \text{Im} f_{\ell}(s) = |f_{\ell}(s)|^2$$

\* Causality  $\rightarrow$  **Analyticity**

$$f_{\ell}(s) = \frac{1}{\pi} \int ds' \frac{\text{Disc } f_{\ell}(s')}{s' - s - i\epsilon}$$

\* Existence of antiparticles  $\rightarrow$  **Crossing symmetry**

$$A(s, \cos \theta_s) = A(t, \cos \theta_t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}(t) P_{\ell}(\cos \theta_t)$$

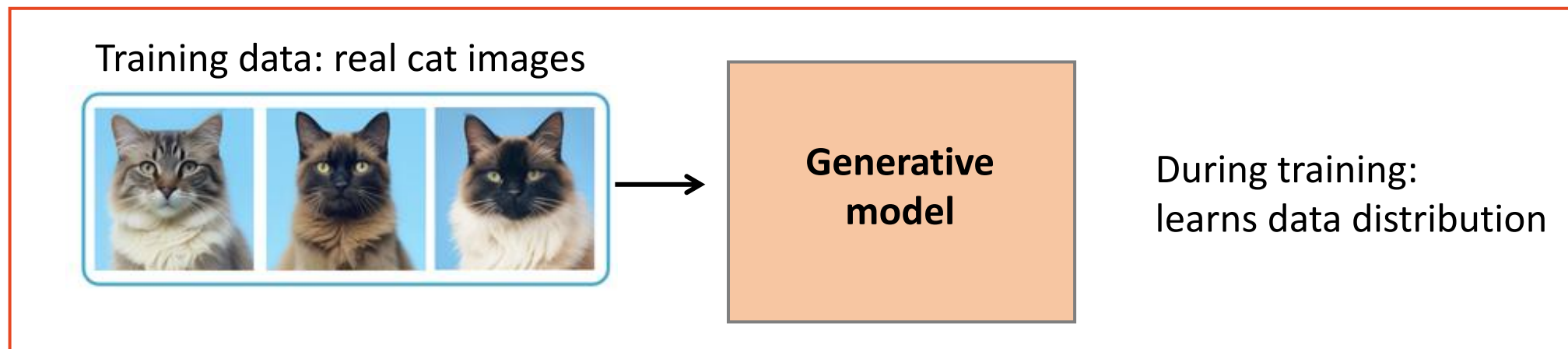


## Build a generative approach informed by S-matrix principles to reconstruct scattering amplitudes

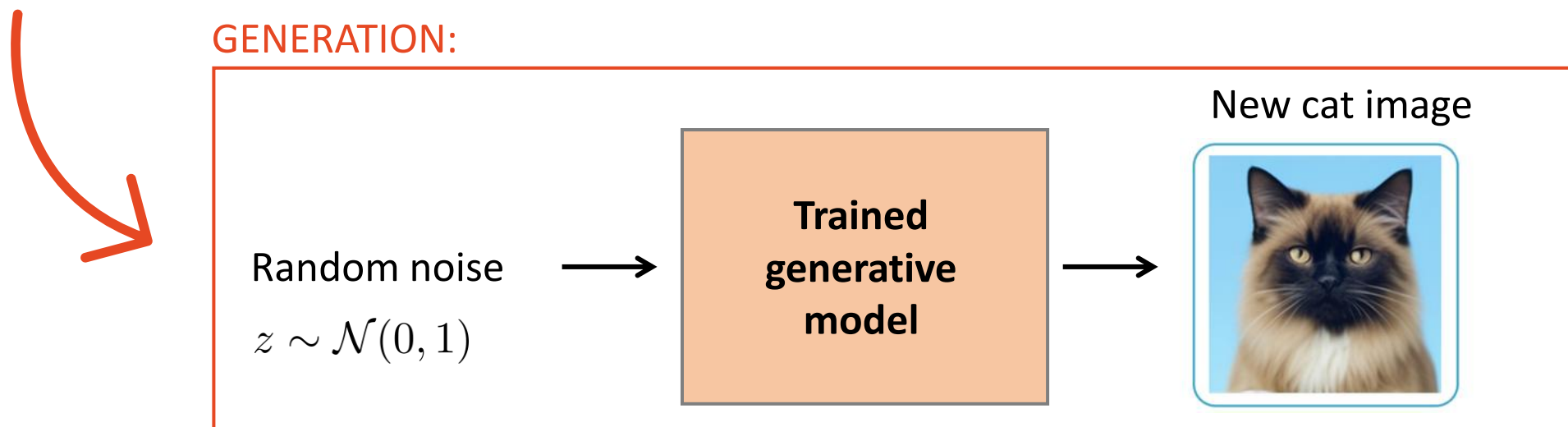
- \* Generative model → Generative adversarial network (GAN)
- \* Physics case → Pion-pion scattering

# Generative models

## TRAINING:



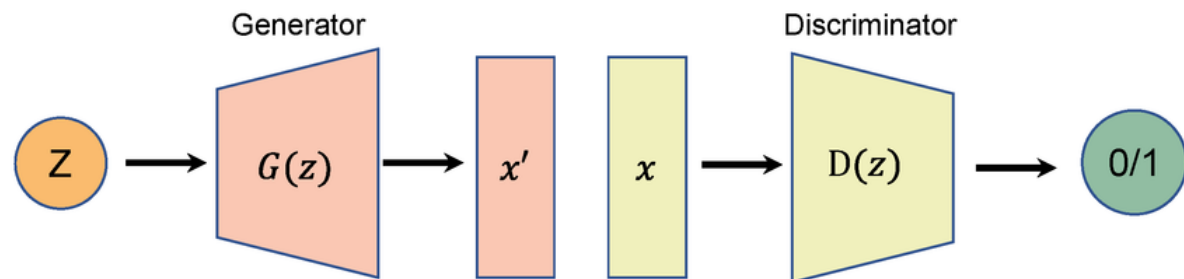
## GENERATION:



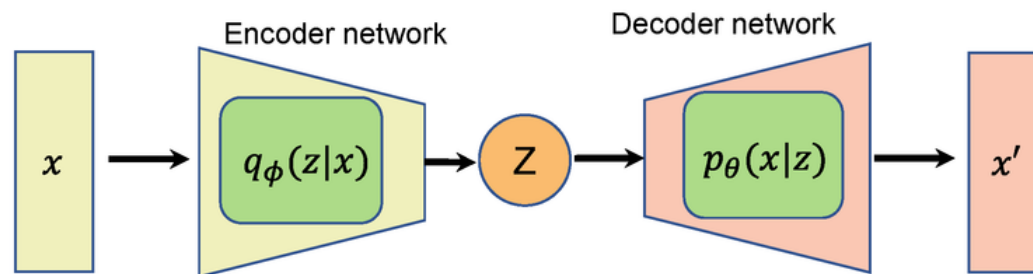


# Generative models

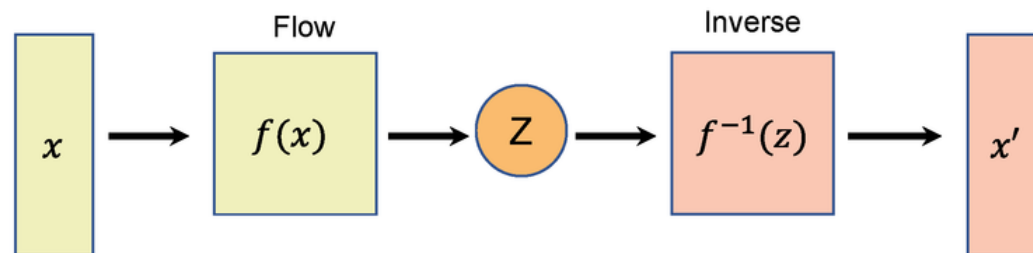
# Generative adversarial network (GAN)



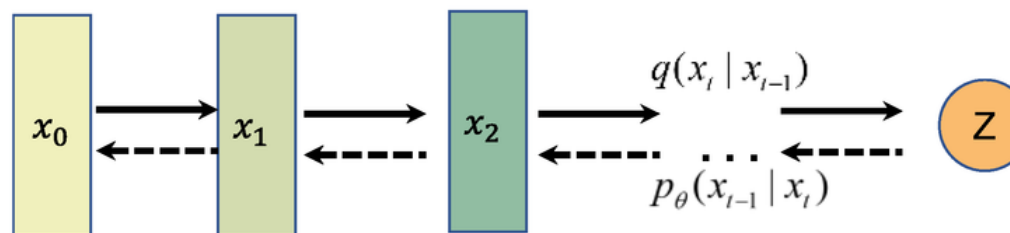
## Variational autoencoder (VAE)



## Normalizing flow (NF)



## Diffusion model (DM)



# $\pi^+\pi^-$ elastic scattering

Simplest physics case: elastic  $2 \rightarrow 2$  of spinless particles

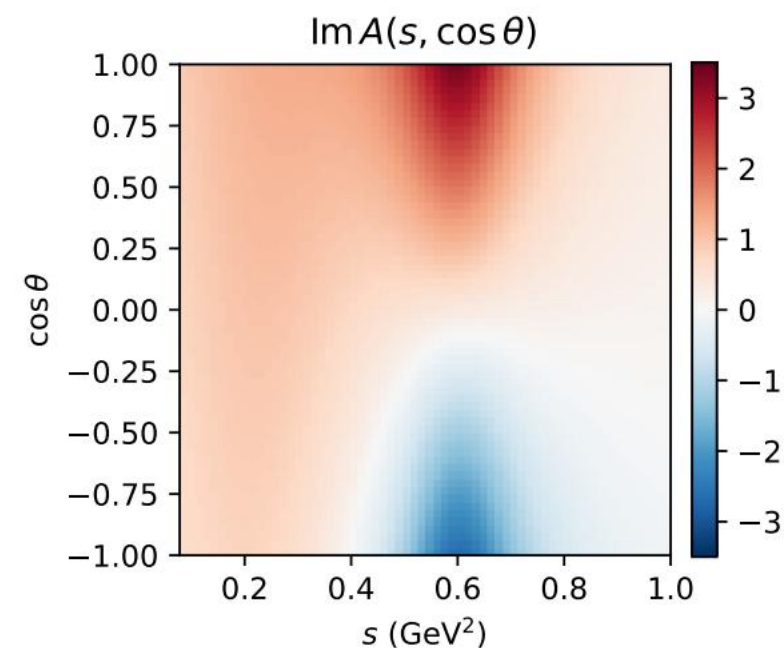
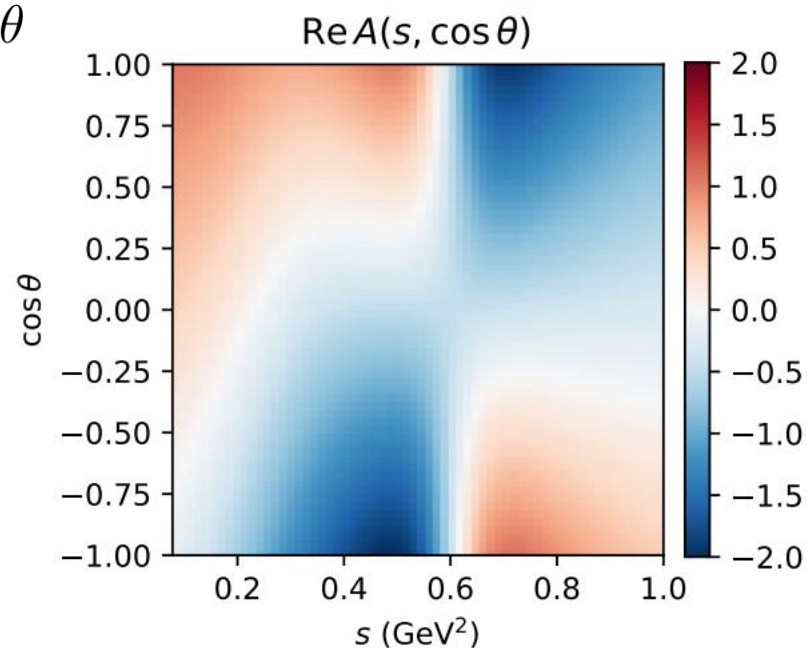
▷ Dominated by  $f_0(500)$  and  $\rho(770)$  resonances

## Model:

Partial-wave decomposition of the **amplitude** truncated to  $n = 1$  
$$\mathcal{A}(s, \cos \theta) = \sum_{\ell=0}^{n=1} (2\ell + 1) f_{\ell}(s) P_{\ell}(\cos \theta)$$

and Breit-Wigner type partial waves: 
$$f_{\ell} = \frac{m_{\ell} \Gamma_{\ell}}{m_{\ell}^2 - s - im_{\ell} \Gamma_{\ell}}$$

$$\mathcal{A}(s, \cos \theta) = f_0(s) + 3f_1(s) \cos \theta$$



# $\pi^+\pi^-$ elastic scattering

Simplest physics case: elastic  $2 \rightarrow 2$  of spinless particles

▷ Dominated by  $f_0(500)$  and  $\rho(770)$  resonances

## Model:

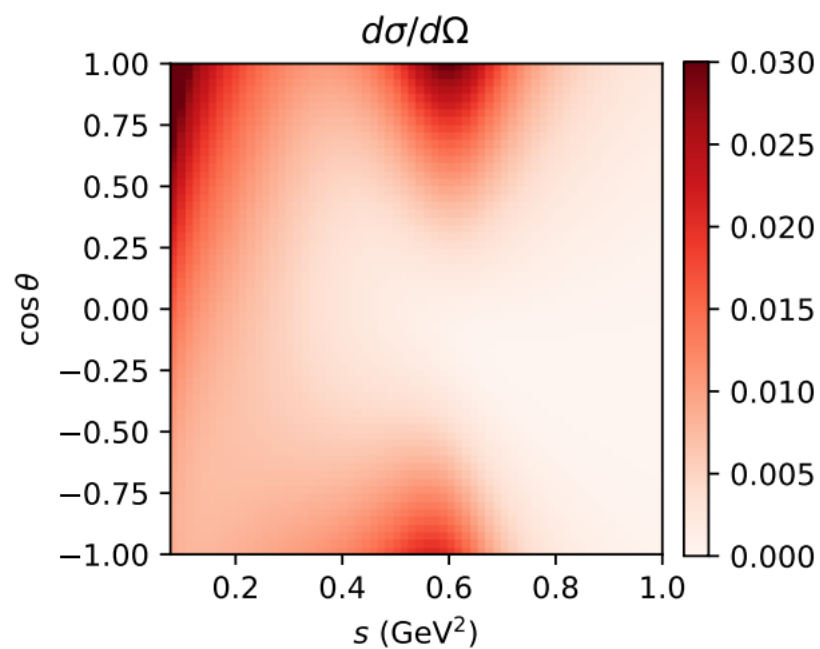
Partial-wave decomposition of the **amplitude** truncated to  $n = 1$   $\mathcal{A}(s, \cos \theta) = \sum_{\ell=0}^{n=1} (2\ell + 1) f_\ell(s) P_\ell(\cos \theta)$

and Breit-Wigner type partial waves:  $f_\ell = \frac{m_\ell \Gamma_\ell}{m_\ell^2 - s - im_\ell \Gamma_\ell}$

$$\mathcal{A}(s, \cos \theta) = f_0(s) + 3f_1(s) \cos \theta$$

## Differential cross section

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \frac{1}{s} |\mathcal{A}(s, \cos \theta)|^2$$



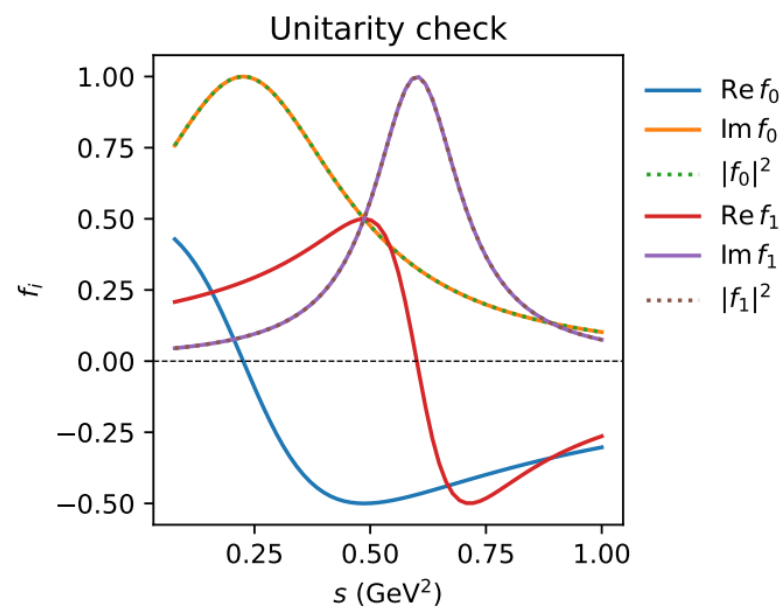
Amplitudes and cross sections as images (discretized functions)

# $\pi^+\pi^-$ elastic scattering

## Unitarity relation

For the partial waves:  $\text{Im } f_\ell(s) = |f_\ell(s)|^2$

$$f_\ell(s) = \frac{1}{2} \int_{-1}^{+1} dz P_\ell(z) \mathcal{A}(s, z)$$

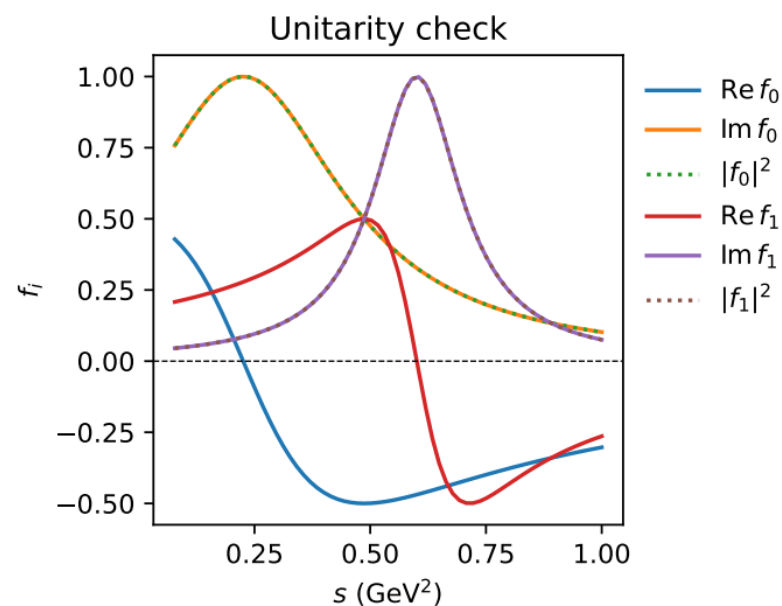


# $\pi^+\pi^-$ elastic scattering

## Unitarity relation

For the partial waves:  $\text{Im } f_\ell(s) = |f_\ell(s)|^2$

$$f_\ell(s) = \frac{1}{2} \int_{-1}^{+1} dz P_\ell(z) \mathcal{A}(s, z)$$



Integral relation for the full amplitude:

$$\text{Im } \mathcal{A}(s, z) = \frac{1}{4\pi} \int_0^{2\pi} d\phi \int_{-1}^{+1} dz' \mathcal{A}(s, z') \mathcal{A}^*(s, z'')$$

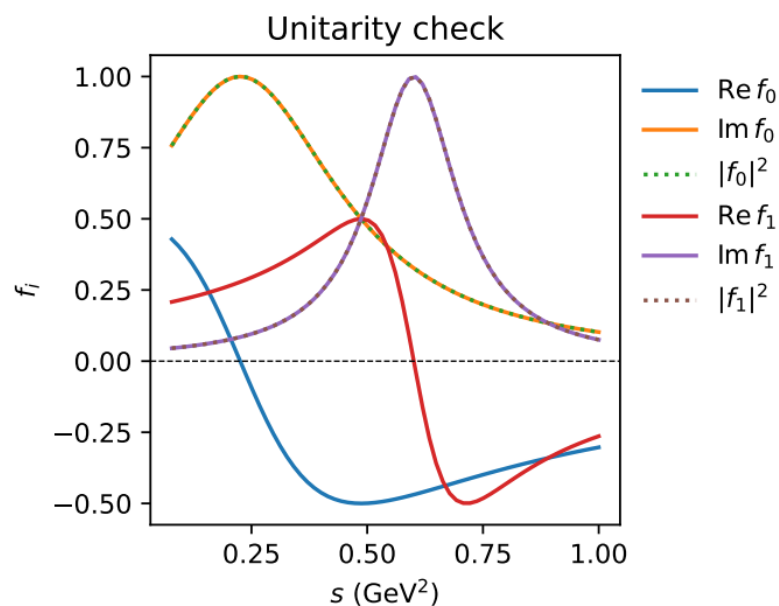
$$z'' = zz' + \sqrt{1-z^2} \sqrt{1-z'^2} \cos \phi$$

# $\pi^+\pi^-$ elastic scattering

## Unitarity relation

For the partial waves:  $\text{Im } f_\ell(s) = |f_\ell(s)|^2$

$$f_\ell(s) = \frac{1}{2} \int_{-1}^{+1} dz P_\ell(z) \mathcal{A}(s, z)$$



Integral relation for the full amplitude:

$$\text{Im } \mathcal{A}(s, z) = \frac{1}{4\pi} \int_0^{2\pi} d\phi \int_{-1}^{+1} dz' \mathcal{A}(s, z') \mathcal{A}^*(s, z'')$$

$$z'' = zz' + \sqrt{1-z^2} \sqrt{1-z'^2} \cos \phi$$

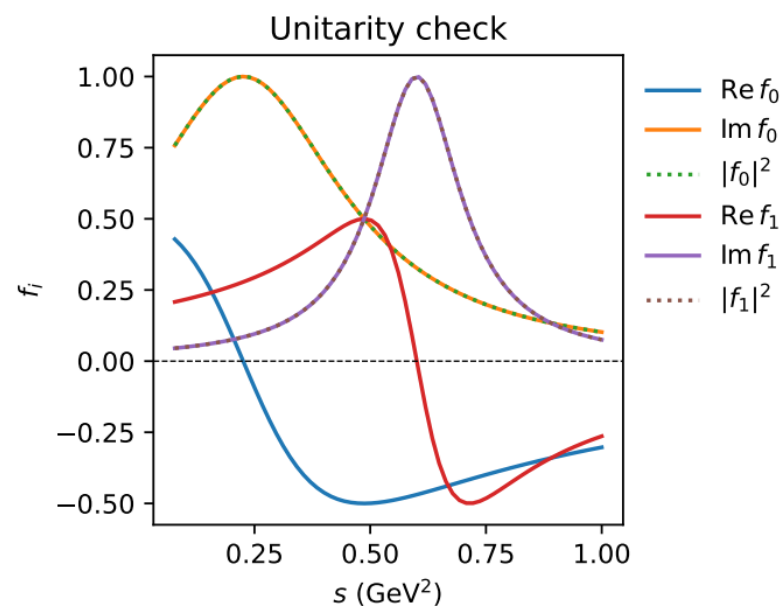
Phase ambiguity:  $\mathcal{A}(s, z) \rightarrow -\mathcal{A}^*(s, z)$

# $\pi^+\pi^-$ elastic scattering

## Unitarity relation

For the partial waves:  $\text{Im } f_\ell(s) = |f_\ell(s)|^2$

$$f_\ell(s) = \frac{1}{2} \int_{-1}^{+1} dz P_\ell(z) \mathcal{A}(s, z)$$



Integral relation for the full amplitude:

$$\text{Im } \mathcal{A}(s, z) = \frac{1}{4\pi} \int_0^{2\pi} d\phi \int_{-1}^{+1} dz' \mathcal{A}(s, z') \mathcal{A}^*(s, z'')$$

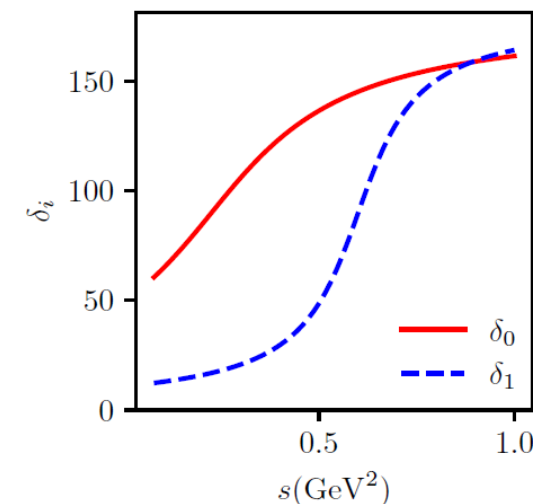
$$z'' = zz' + \sqrt{1-z^2} \sqrt{1-z'^2} \cos \phi$$

Phase ambiguity:  $\mathcal{A}(s, z) \rightarrow -\mathcal{A}^*(s, z)$

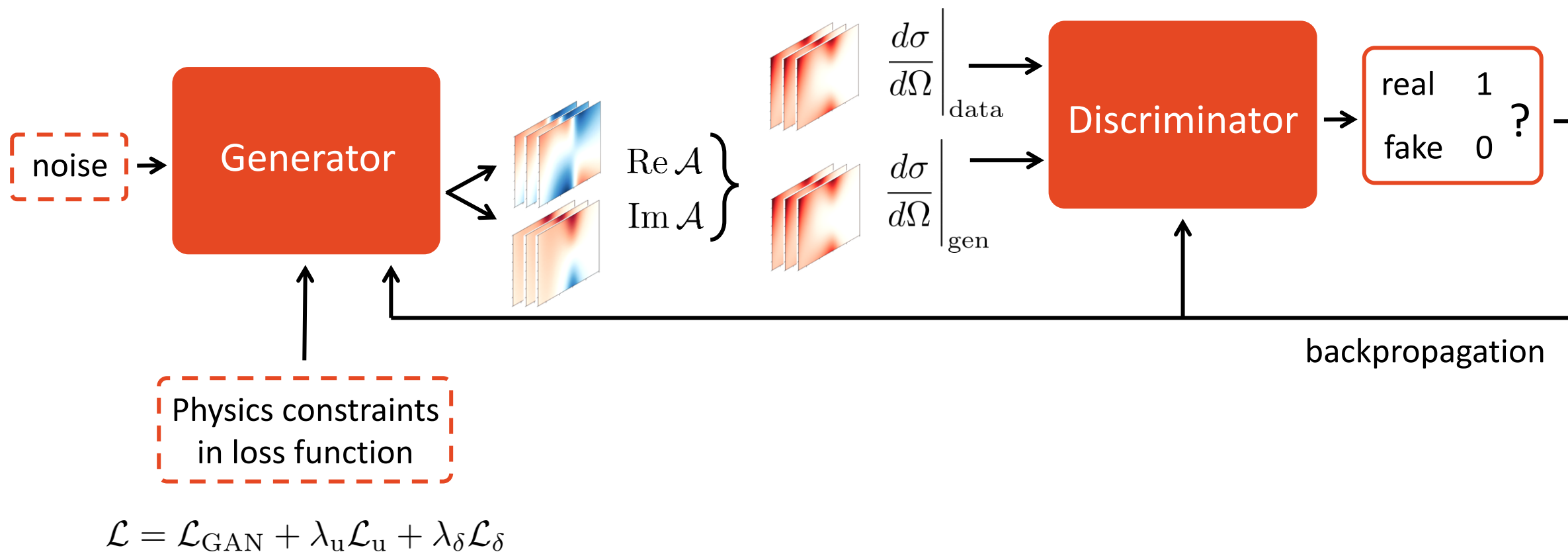
Phase shift monotonicity  
(**causality**)

$$\delta_\ell = \text{atan} \left( \frac{\text{Im } f_\ell(s)}{\text{Re } f_\ell(s)} \right)$$

$$\frac{d\delta_\ell(s)}{ds} \geq 0$$



# GAN with physics constraints





# Loss function

**Adversarial loss:** Mean squared error (LSGAN)

**Unitarity loss:** Enforce unitarity by comparing the modulus squared of the numerical integral of the scattering amplitude over angular variables to the imaginary part

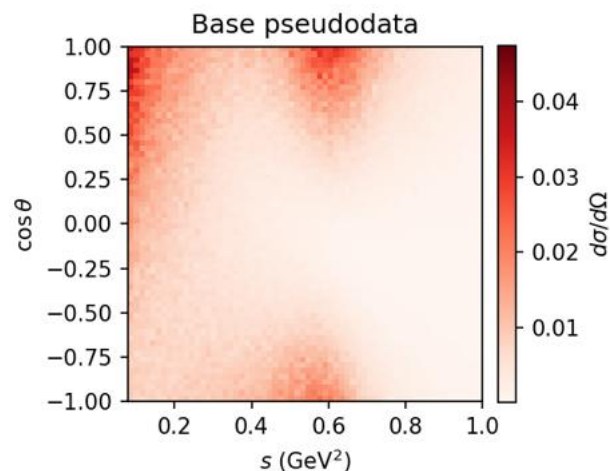
$$\mathcal{L}_u = \frac{1}{N \cdot N_s \cdot N_z} \sum_{i=1}^N \sum_{j=1}^{N_s} \sum_{k=1}^{N_z} \left( |\text{Im } \mathcal{A}(s, z) - \text{Re } \mathcal{I}(s, z)| + |\text{Im } \mathcal{I}(s, z)| \right)$$

$$\mathcal{I}(s, z) = \frac{1}{4\pi} \int_{-1}^1 dz' \int_0^{2\pi} d\phi \left( \mathcal{A}(s, z') \mathcal{A}(s, z''(z, z', \phi)) \right) \quad [\cos \theta \times \phi] = 64 \times 10$$

**Phase shift loss:** Ensure the positive derivative of the phase shift (causality)

$$\mathcal{L}_{D_\ell} = \frac{1}{N} \sum_{i=1}^N \log \left( \max \left( 0, -\Delta \delta_\ell(s) \right) + 1 \right),$$

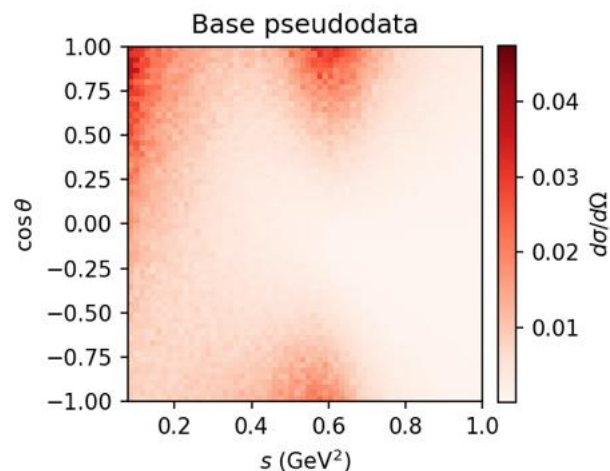
# Training pseudodata



Differential cross section from the model  
with relative Gaussian noise  $\mathcal{N}(0, \sigma_{\text{exp}})$  (mimic experimental data)

- ▷ discretized  $64 \times 64$ ,  $s \in [(2m_\pi)^2, 1 \text{ GeV}^2]$ ,  $\cos \theta \in [-1, 1]$
- ▷ normalized

# Training pseudodata



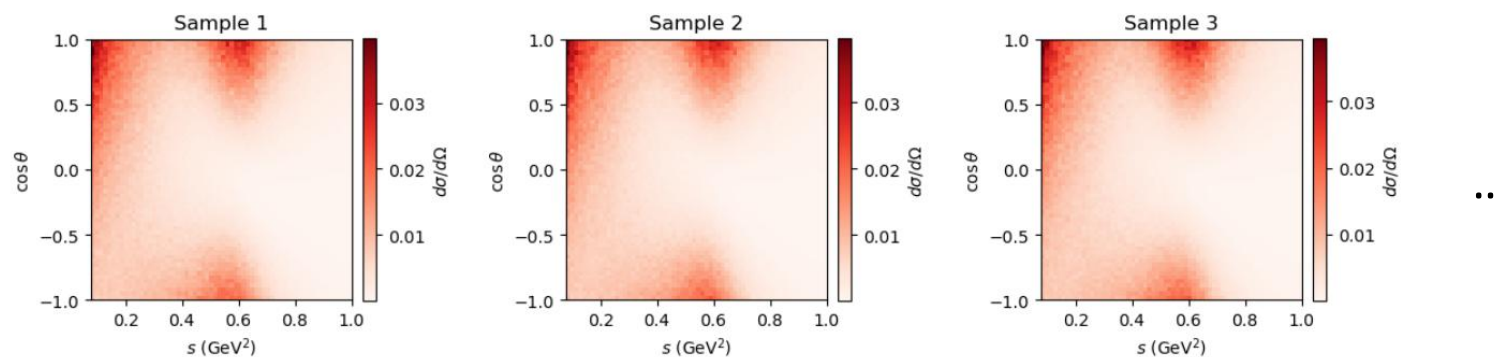
Differential cross section from the model  
with relative Gaussian noise  $\mathcal{N}(0, \sigma_{\text{exp}})$  (mimic experimental data)

▷ discretized  $64 \times 64$ ,  $s \in [(2m_\pi)^2, 1 \text{ GeV}^2]$ ,  $\cos \theta \in [-1, 1]$

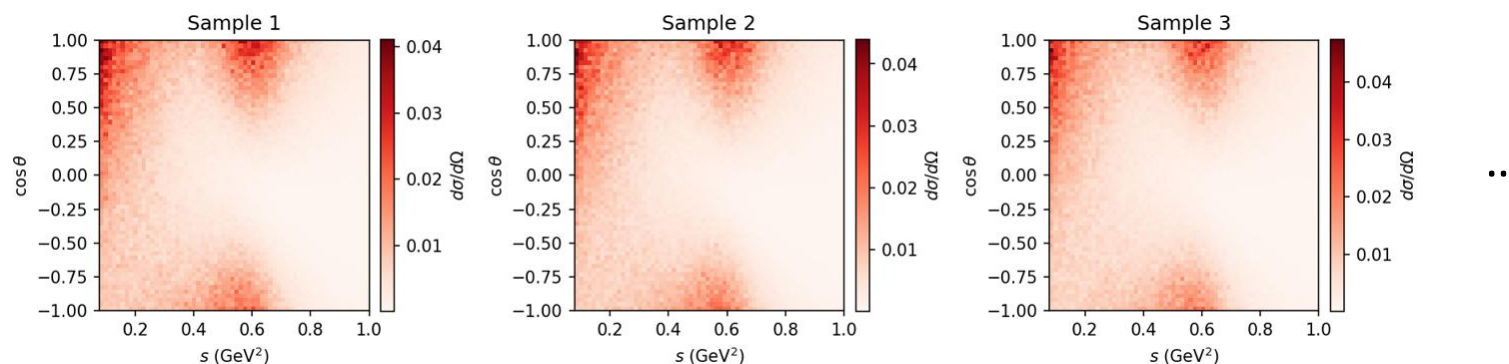
▷ normalized

▷ with additional noise  $\mathcal{N}(0, \sigma_{\text{aug}})$  (augment around base image)

$$\sigma_{\text{exp}} = \sigma_{\text{aug}} = 0.05$$



$$\sigma_{\text{exp}} = \sigma_{\text{aug}} = 0.10$$



# The convergence challenge

## Stabilizing the adversarial game

- ▷ Spectral normalization to prevent runaway gradients
- ▷ Asymmetric learning rates ( $LR_D < LR_G$ ) and label smoothing (real = 0.9) to avoid over-confident discriminator
- ▷ Dropout in D after each conv block
- ▷ LSGAN (MSE) instead of BCE (non-saturating gradients far from the decision boundary)

## Multi-loss optimization

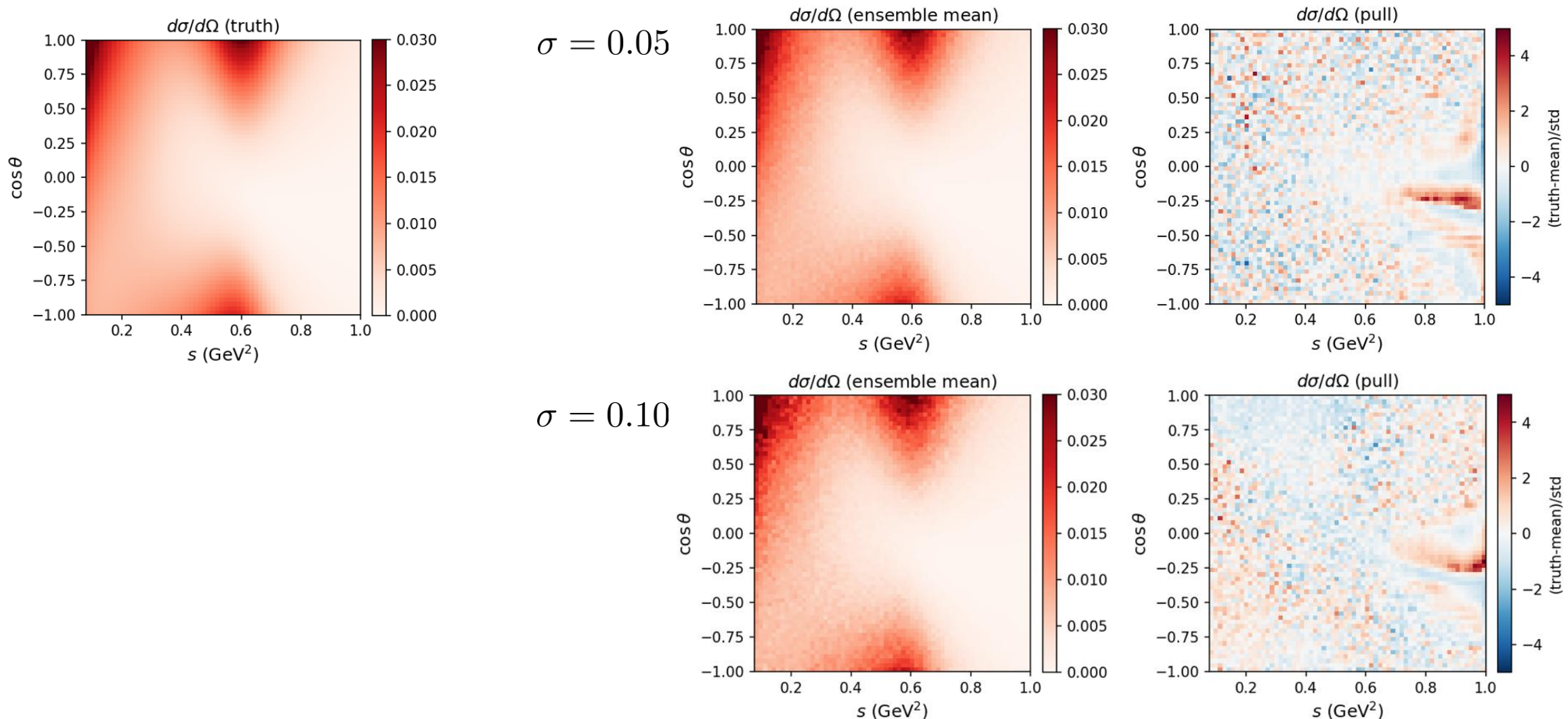
- ▷ PCGrad gradient surgery: physics gradients are projected onto the orthogonal direction when they conflict before optimization

## Convergence detection (based on $\mathcal{L}_{\text{unit}}$ )

- ▷ Early failure detection: abort the trial and restart training (up to 5 restarts before declaring failure)
- ▷ Early stopping
- ▷ Learning-rate reduction near the minimum

# Preliminary results: cross section

Mean over 10 trained GAN, with 10 generated samples per GAN (100 total samples)

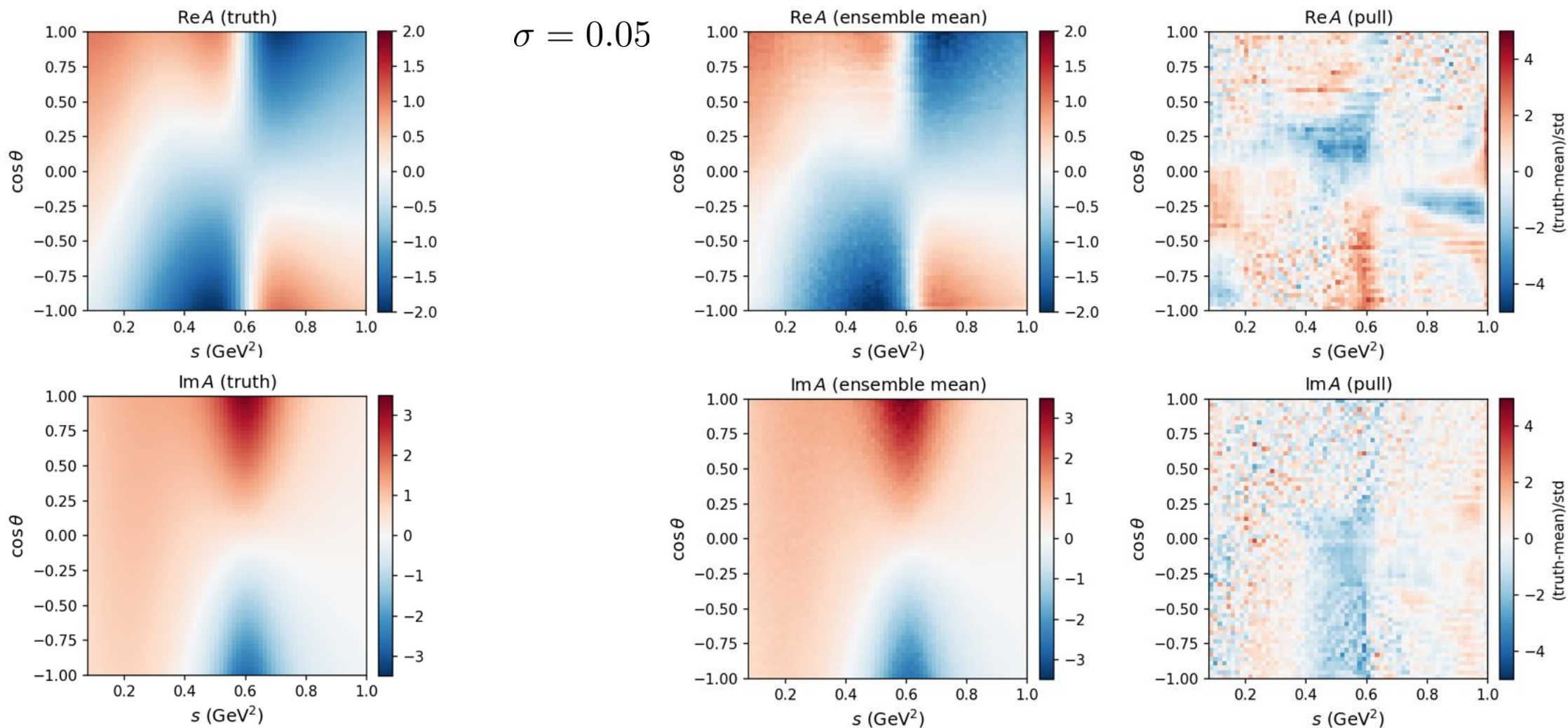




# Preliminary results: amplitude

Mean over 10 trained GAN, with 10 generated samples per GAN (100 total samples)

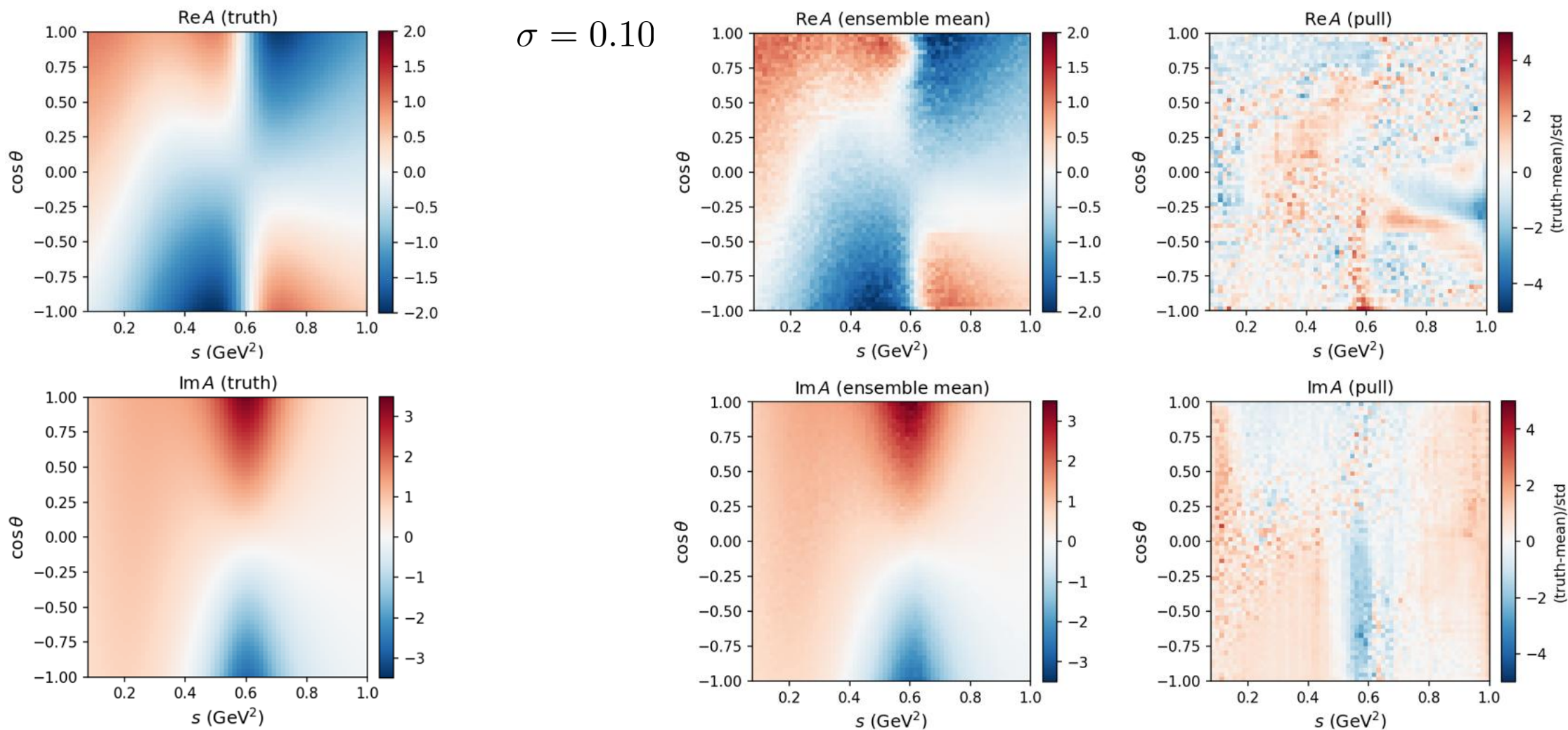
$$\sigma = 0.05$$



# Preliminary results: amplitude

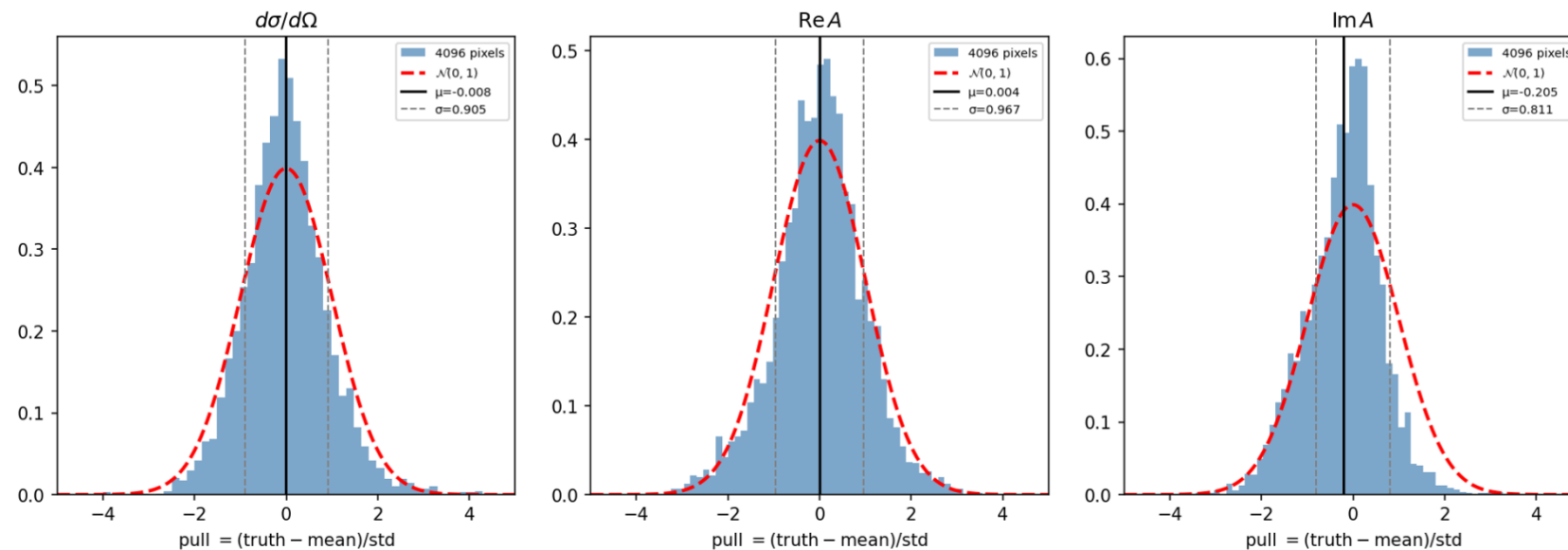
Mean over 10 trained GAN, with 10 generated samples per GAN (100 total samples)

$$\sigma = 0.10$$

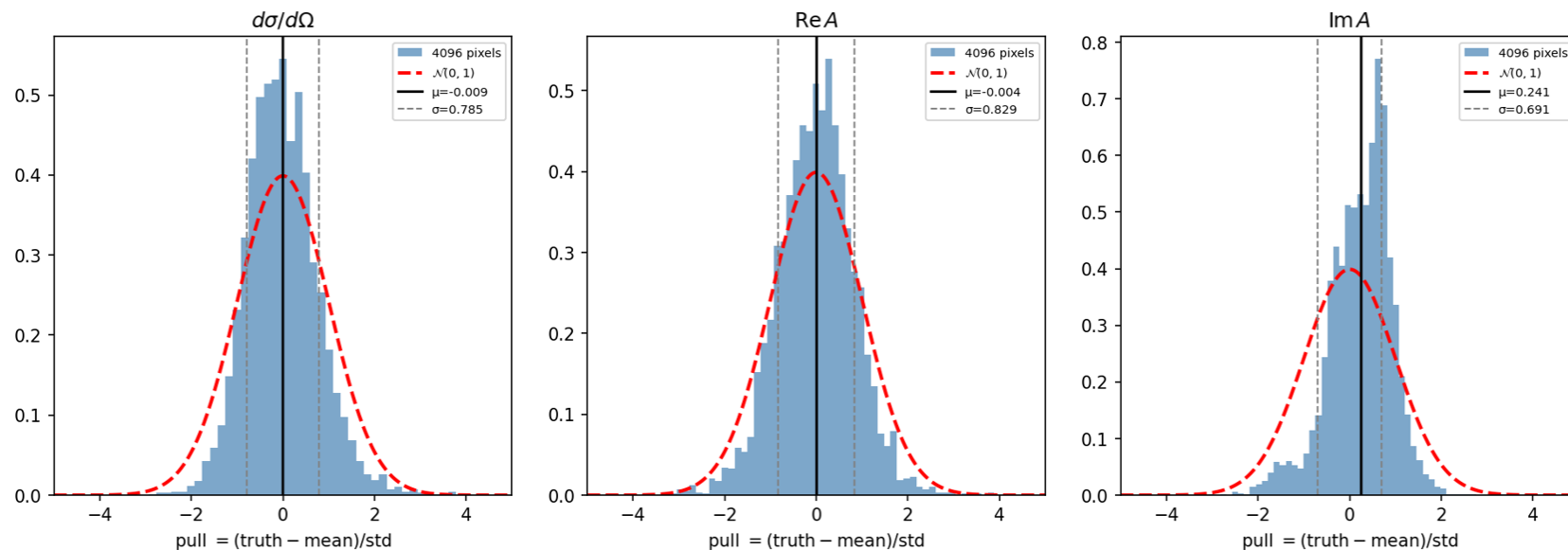


# Preliminary results: pull histograms

$\sigma = 0.05$



$\sigma = 0.10$

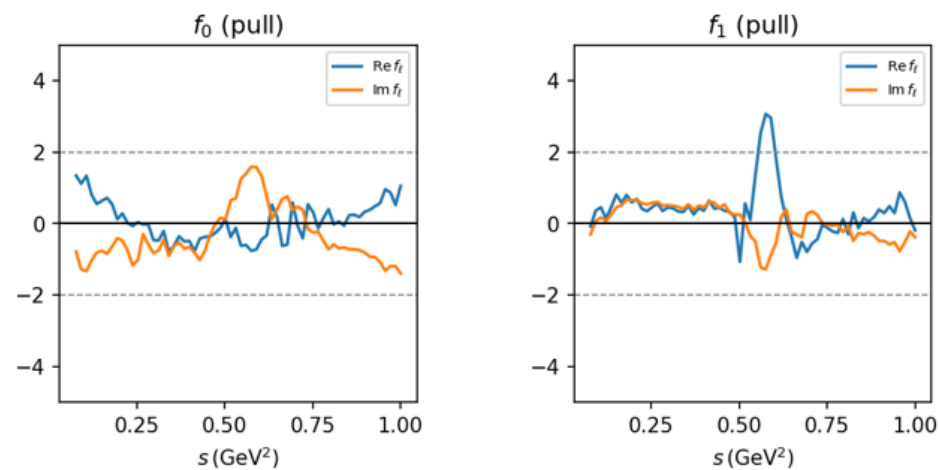
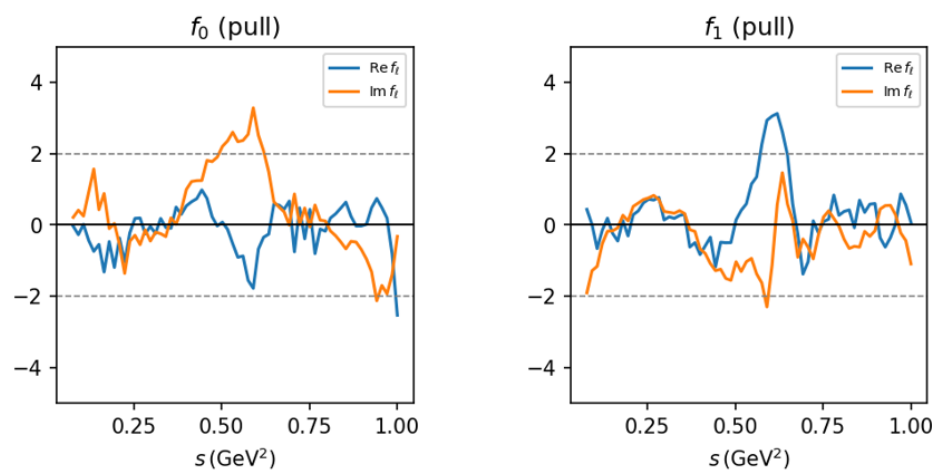
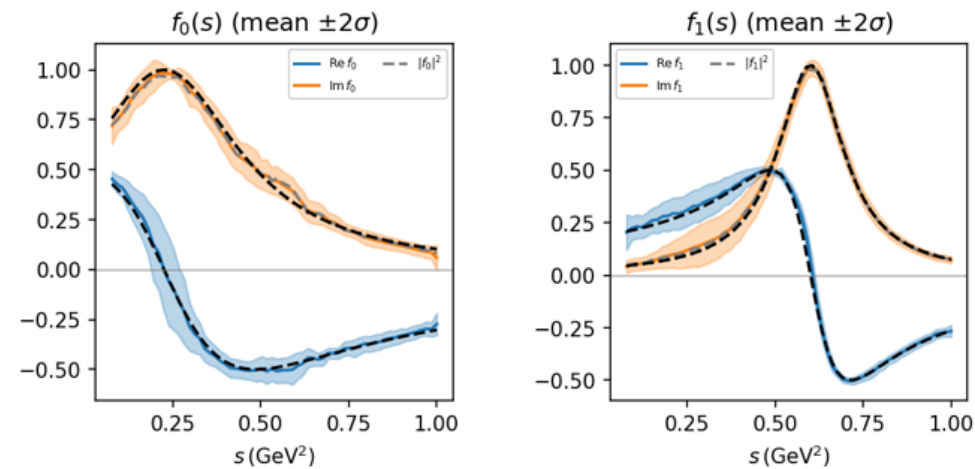
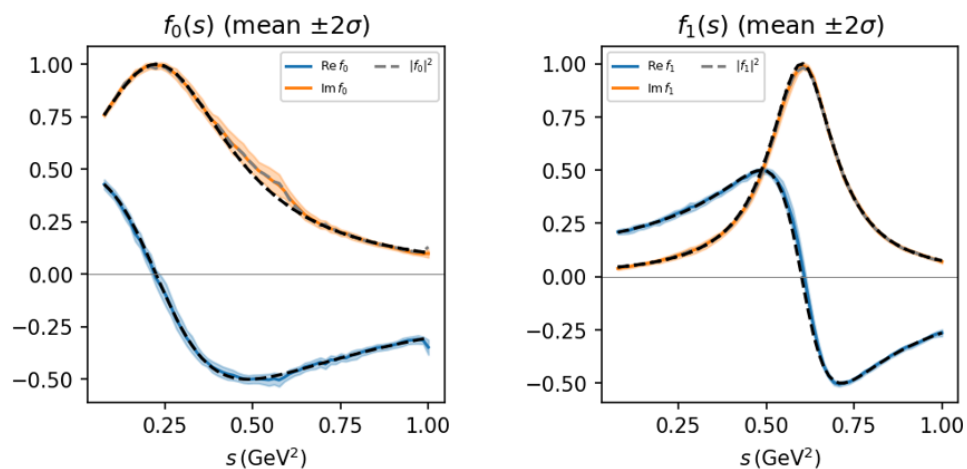




# Preliminary results: partial waves $\ell = 0, 1$

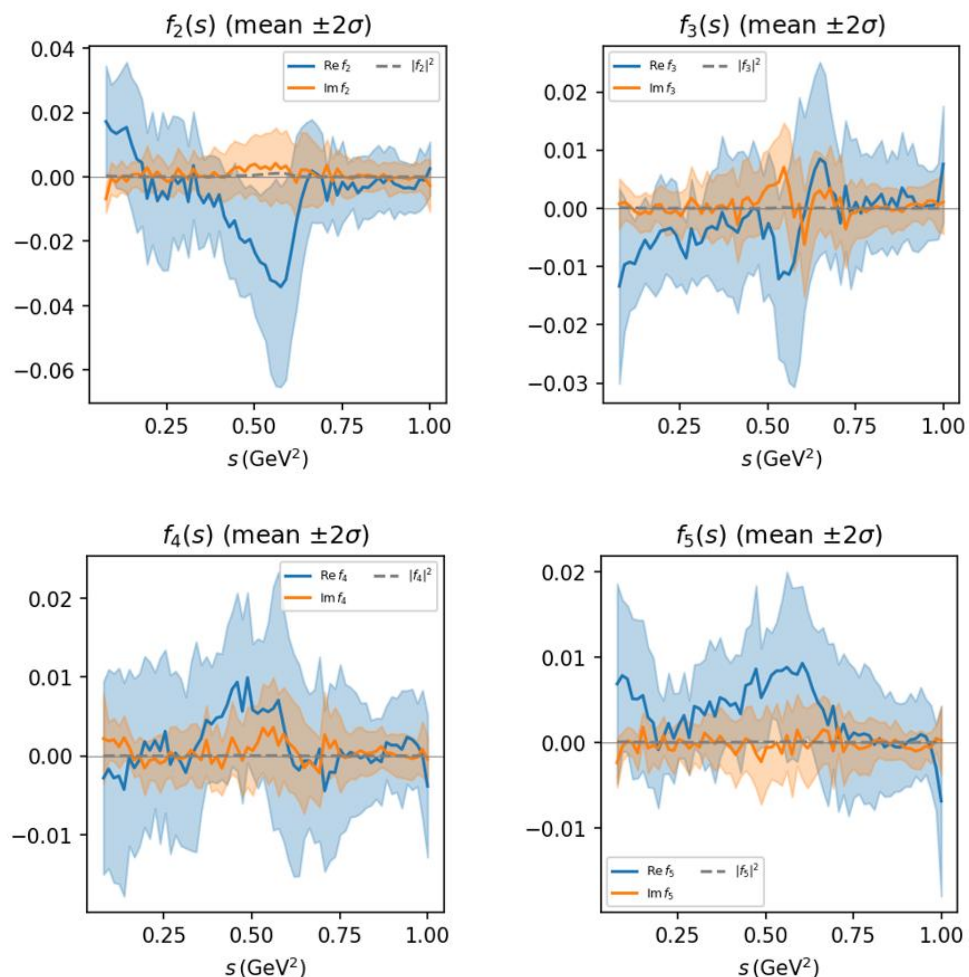
$\sigma = 0.05$

$\sigma = 0.10$

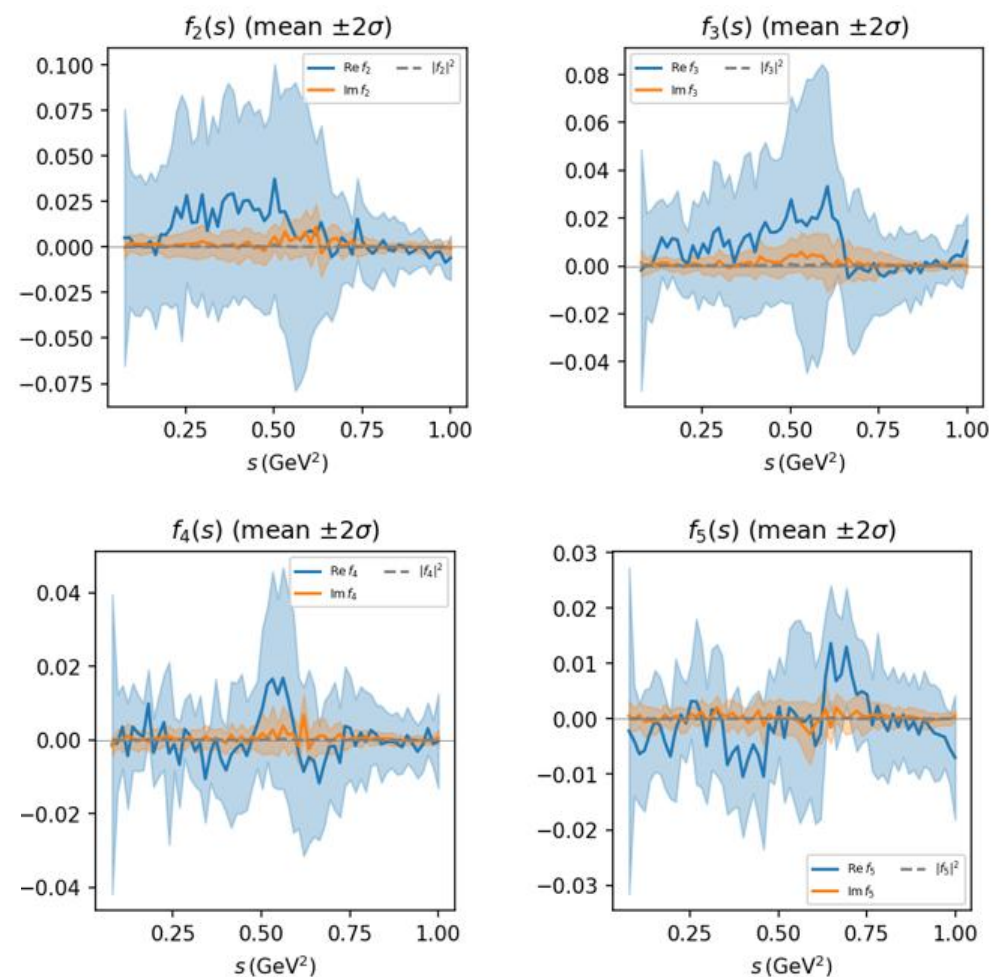


# Preliminary results: higher partial waves

$$\sigma = 0.05$$



$$\sigma = 0.10$$



# Conclusions & outlook

- \* Amplitude reconstruction from cross section is an ill-posed inverse problem
- \* Generative models with S-matrix constraints to solve it
- \* Can handle noisy pseudodata
- \* Questions:
  - ▷ Can we scale to inelastic scattering ( $\pi\pi \leftrightarrow K\bar{K}$ )?  
channels with spin ( $\gamma p \rightarrow \pi N$ )?  
2  $\rightarrow$  3 reactions?
  - ▷ Apply to real experimental data from Jefferson Lab
- \* GANs are difficult to train: consider alternative generative models