

Neutrino Physics – Exercises for the Ecole doctorale PHENIICS

Salvador Urrea and Luighi Leal

March 2026

Contents

1 Sterile neutrinos and non-unitarity	1
2 Simple neutrino mass models	3

1 Sterile neutrinos and non-unitarity

Without resorting to any particular model, one may assume in full generality that n extra fermions mix with the light neutrinos. The leptonic mixing matrix \mathcal{U} that diagonalizes the full neutrino mass matrix and connects flavor and mass eigenstates is then a unitary $(3+n) \times (3+n)$ matrix:

$$\mathcal{U} = \begin{pmatrix} N & \Theta \\ R & S \end{pmatrix}. \quad (1)$$

The 3×3 block N plays the role of an effective PMNS matrix and is in general non-unitary:

$$NN^\dagger \neq \mathbb{I}. \quad (2)$$

It can be parametrized as:

$$N = (I - T)U, \quad (3)$$

where $(I - T)$ is a lower triangular matrix:

$$T = \begin{pmatrix} \alpha_{ee} & 0 & 0 \\ \alpha_{\mu e} & \alpha_{\mu\mu} & 0 \\ \alpha_{\tau e} & \alpha_{\tau\mu} & \alpha_{\tau\tau} \end{pmatrix}, \quad (4)$$

and U is unitary.

Exercise 1

1.1 Why can N be written in the form $N = (I - T)U$ without loss of generality?

1.2 Prove that:

$$\alpha_{\beta\beta} = \frac{1}{2} \left(\Theta \Theta^\dagger \right)_{\beta\beta} = \frac{1}{2} \sum_{i=4}^{3+n} |\mathcal{U}_{\beta i}|^2, \quad (5)$$

$$\alpha_{\gamma\beta} = \left(\Theta \Theta^\dagger \right)_{\gamma\beta} = \sum_{i=4}^{3+n} \mathcal{U}_{\gamma i} \mathcal{U}_{\beta i}^*. \quad (6)$$

Exercise 2

The oscillation probability in vacuum is:

$$P_{\gamma\beta} = \left| \left(\mathcal{U} \mathcal{S} \mathcal{U}^\dagger \right)_{\beta\gamma} \right|^2, \quad (7)$$

with

$$\mathcal{S} = \text{diag} \left(\exp \left(-i \frac{\Delta m_{j1}^2 L}{2E} \right) \right). \quad (8)$$

2.1 Assuming a distance $L = 574$ m and neutrino energy $E = 1$ GeV, what is the expected standard oscillation result?

2.2 Assuming one sterile neutrino with $\Delta m_{41}^2 = 10 \text{ eV}^2$, compute the probability $P_{\gamma\beta}$ for:

- appearance,
- disappearance.

2.3 What happens to this probability in the limit $\Delta m_{41}^2 \rightarrow \infty$?

2.4 Taking into account that detectors have a finite energy resolution, can the previous issue be resolved?

2.5 Express this oscillation probability in terms of the α parametrization.

Exercise 3

3.1 An experiment observes a signal in $P_{\mu e}$ but no signal in P_{ee} or $P_{\mu\mu}$.

Is this consistent with a scenario containing only one sterile neutrino?

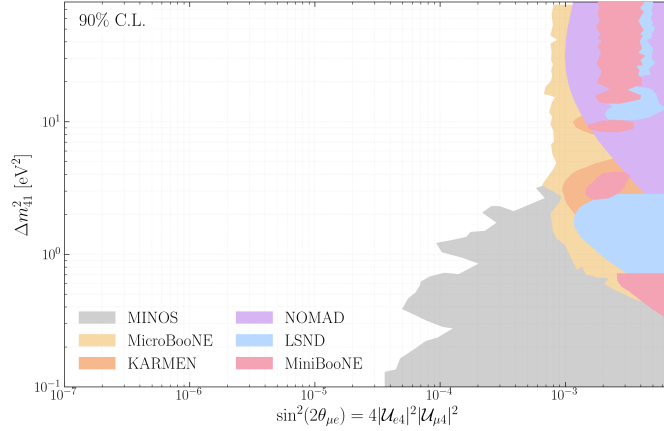


Figure 1: The shaded pink, orange, yellow, and gray areas show current constraints from other experiments, while the shaded purple regions are favored at 99% CL by the LSND and MiniBooNE anomalies.

2 Simple neutrino mass models

The Linear-Inverse Seesaw (LISS) combines elements from both the inverse and linear seesaw mechanisms, providing a framework that generates the required neutrino mass spectrum and mixing patterns.

A minimal realization is obtained by extending the SM with two right-handed neutrinos, N_R^1 and N_R^2 , at a new physics scale Λ , carrying opposite lepton numbers:

$$L(N_R^1) = +1, \quad L(N_R^2) = -1.$$

The Lagrangian is:

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + i\bar{N}_R^1 \not{\partial} N_R^1 + i\bar{N}_R^2 \not{\partial} N_R^2 - \left(Y_\alpha \bar{\ell}_\alpha \tilde{\phi} N_R^1 + \epsilon Y'_\alpha \bar{\ell}_\alpha \tilde{\phi} N_R^2 + \frac{\Lambda}{2} \bar{N}_R^{1c} N_R^2 + \frac{\mu}{2} \bar{N}_R^{2c} N_R^2 \right) + \text{h.c.} \quad (9)$$

After electroweak symmetry breaking, the neutrino mass terms can be written as:

$$-\mathcal{L}_{m_\nu} = \frac{1}{2} n_L^T C \mathcal{M}_{\text{LISS}} n_L + \text{h.c.}, \quad (10)$$

where

$$n_L = (\nu_{eL}, \nu_{\mu L}, \nu_{\tau L}, N_1^c, N_2^c)^T, \quad C = i\gamma^2 \gamma^0.$$

Exercise 5

5.1 Write explicitly the 5×5 neutrino mass matrix $\mathcal{M}_{\text{LISS}}$.

5.2 How many light neutrinos can be massive? Is it compatible with neutrino oscillation data?

Defining $m_D = \mathbf{Y}v$, $m_L = \epsilon \mathbf{Y}'v$, and in the following limit

$$|m_D = \mathbf{Y}v|, |m_L = \epsilon \mathbf{Y}'v|, |\mu| \ll \Lambda,$$

we can perform a block diagonalization:

$$U_B^T \mathcal{M}_{\text{LISS}} U_B = \begin{pmatrix} m_{\text{light}}^{3 \times 3} & 0 \\ 0 & M_{\text{heavy}}^{2 \times 2} \end{pmatrix}, \quad (11)$$

where U_B is a unitary matrix and, at leading order:

$$m_{\text{light}} = \frac{\mu}{\Lambda^2} m_D m_D^T - \frac{1}{\Lambda} (m_D m_L^T + m_L m_D^T).$$

$$M_{\text{heavy}}^{2 \times 2} \simeq \begin{pmatrix} 0 & \Lambda \\ \Lambda & \mu \end{pmatrix}$$

Exercise 7

- 7.1** Identify the contributions to m_{light} proportional to μ and to ϵ . Explain why they are referred to as the *inverse seesaw* and *linear seesaw* terms.

Exercise 8

- 8.1** Obtain the masses of the heavy neutrinos in the model.
- 8.2** What happens in the limit $\mu \rightarrow 0$?
- 8.3** Do you understand why they would be called a pseudo-Dirac pair?