Recent results from E-203 at FACET (SLAC):

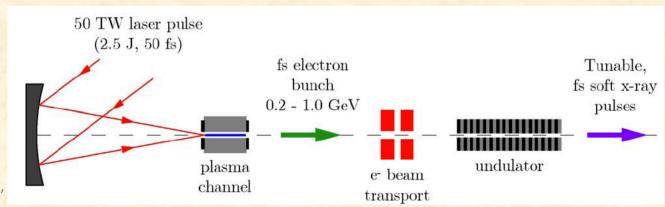
Longitudinal profile measurement of ultra-short electron bunches

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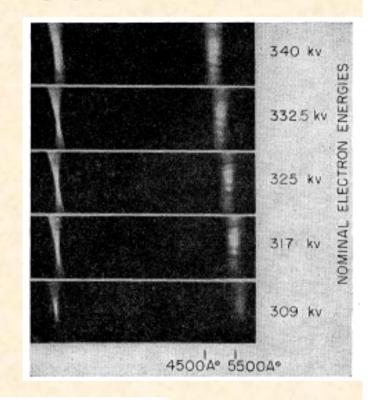
Motivations

- In particle accelerators some effects (eg: coherent effects) depends on the bunch current (charge/time), not the bunch charge.
- Laser-driven plasma accelerators (LPA) use fs laser pulses to produce ultra-short bunches of electrons.
- To use LPA as drivers for FEL it is important to be able to measure the longitudinal profile of the bunch.
- This can be attempted using Smith-Purcell radiation.

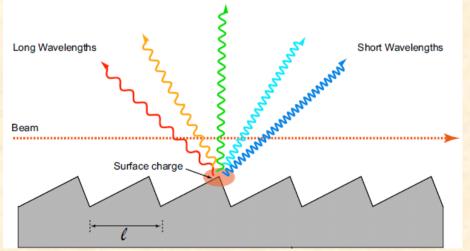


Smith-Purcell effect

- Discovered 1953 by Smith and Purcell.
- Surface charges induced in a grating emit light.
- This can be explained by the accelerating of the induced charges at the teeth of the grating.
- This light is spectrally dispersed by the grating.
- When the charges are induced by a short bunch there can be coherent emission.



$$\lambda = \frac{l}{n} \left(\frac{1}{\beta} - \cos \theta \right)$$



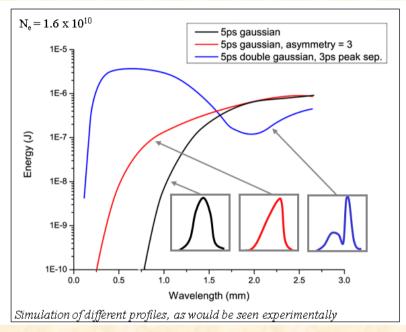
Coherent emission

- Like most radiative processes SPR has an incoherent and a coherent component.
- The coherent component is the product of N² and the Fourier transform of the bunch shape squared.

$$\begin{split} \left(\frac{dI}{d\Omega d\omega}\right)_{N_{e}}(\Omega,\omega) &= \left(\frac{dI}{d\Omega d\omega}\right)_{sp}(\Omega,\omega) \cdot \left[N_{e} + N_{e}(N_{e} + 1) \mid F(\omega) \mid^{2}\right] \\ \left(\frac{dI}{d\Omega}\right)_{sp} &= 2\pi q^{2} \frac{Z}{\ell^{2}} \frac{n^{2}\beta^{3}}{(1 - \beta\cos\theta)^{3}} R^{2} \exp\left(-\frac{2x_{0}}{\lambda_{e}}\right) \\ \lambda_{e} &= \frac{\lambda}{2\pi} \frac{\beta\gamma}{\sqrt{1 + \beta^{2}\gamma^{2}\sin^{2}\theta\sin^{2}\phi}} \end{split}$$

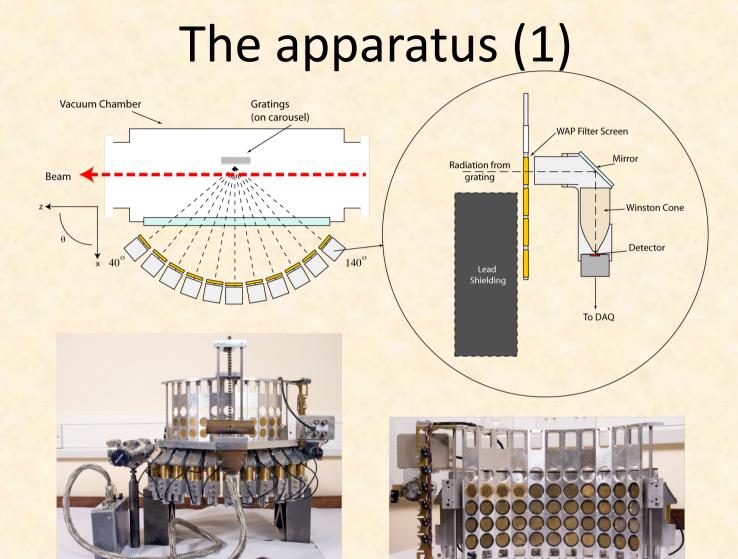
 A measure of the SPR therefore gives a measure of the Fourier transform of the bunch shape.

Effect of the profile on the emitted spectrum



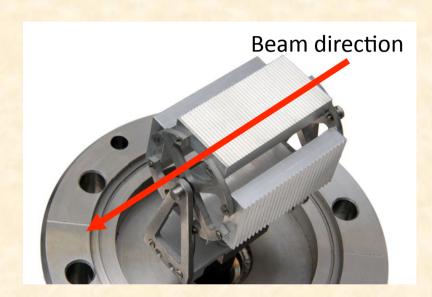
$$\left(\frac{\mathrm{dI}}{\mathrm{d}\Omega\mathrm{d}\omega}\right)_{\mathrm{N_e}}(\Omega,\omega) = \left(\frac{\mathrm{dI}}{\mathrm{d}\Omega\mathrm{d}\omega}\right)_{\mathrm{sp}}(\Omega,\omega)\cdot\left[\mathrm{N_e} + \mathrm{N_e}(\mathrm{N_e} + 1)\,|\,\mathrm{F}(\omega)\,|^2\right]$$

- Pulses with the same "length" but different shapes will produce different SPR spectrum.
- A device able to measure the emitted radiation at different angles is therefore needed.

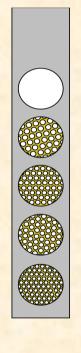


Ability to switch between 3 gratings (sets of wavelengths)

The apparatus (2)





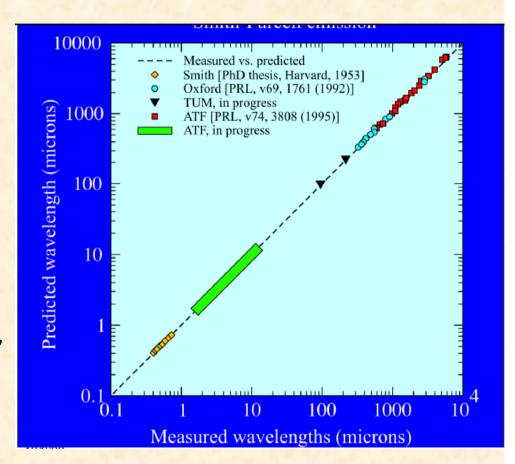




Gratings carrousel, filters, filters ladder and winston cones.

Early experiments

- In the early 1990s Smith-Purcell radiation (SPR) was seen as a possible seed for FELs.
- After that several experiments were done by G. Doucas to assess its potential as a longitudinal profile monitor in the ps range:
 - FELIX (Netherlands) 45 MeV, 5.5 ps
 - End station A, SLAC (US) 28 GeV, 5 ps

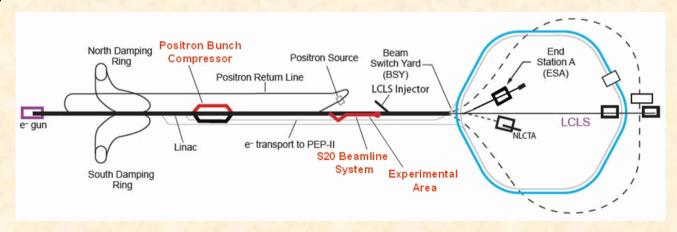


5ps => mm wave

Application to LPA

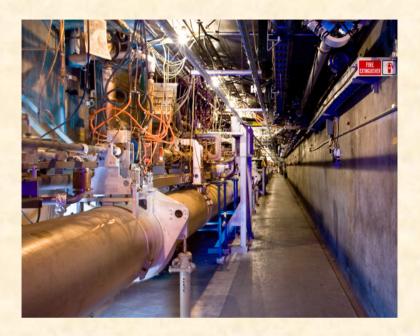
- In 2008, studies started to use SPR at LPA.
- LPA have very short bunches: less than 100fs (some claim just a few fs).
- Such short bunches become coherent near the THz range (50fs =>

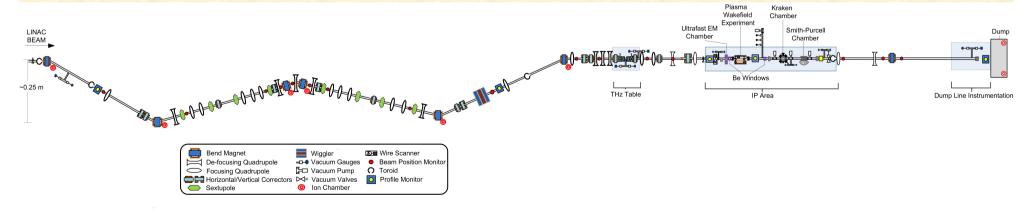
 - In the THz range many materials are absorbing radiation. $\lambda = \frac{l}{n} \left(\frac{1}{\beta} \cos \theta \right)$ => Measurement is more complicated
- Before trying at a LPA where there are many other unknown parameters we applied for beam time at a conventional accelerator: FACET



FACET

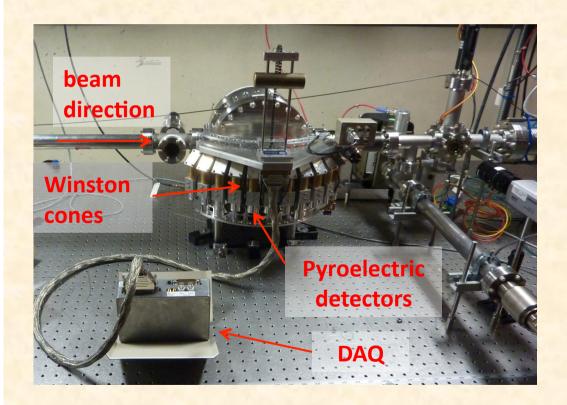
- Facility for Advanced Experimental Tests
- Part of the SLC Linac (2/3) not used by LCLS
- Goal: advanced experiments on beam driven plasma acceleration but open to other R&D.
- Very short bunches (design 60fs) at 20 GeV.
- Up to 3nC/bunch
- 1-30 Hz
- Now under commissioning...



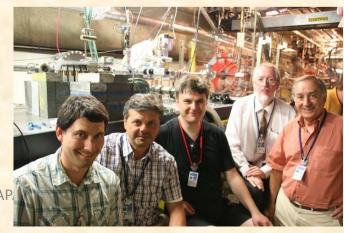


E-203

 Longitudinal profile measurements using Smith-Purcell radiation.

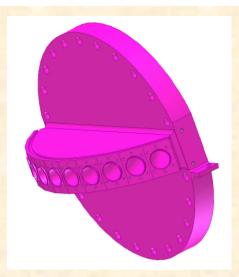




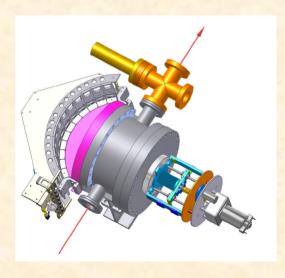


SPR in the THz gap

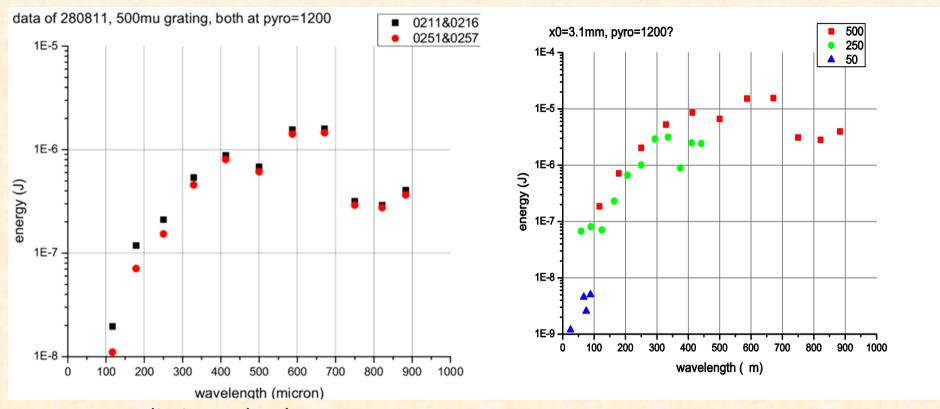
- Moving from mm-wave to micrometric/THz wave was not easy.
- With mm-wave Band pass filters were easy to build: a plate with holes drilled in the workshop! Filters in the THz rang are complicated and much more expensive arrays of mesh.
- The new gratings (50um and 250um pitch) were also more difficult to manufacture (requires high precision).
 Regular optical grating are not suitable as they are made of glass.
- Fused silica absorbs radiation on part of the THz range. We had to find a new window type. Synchrotron beam lines use synthetic diamond, but at £5-10000 a piece it was out of price => We decided to use special Silicium (2 suppliers worldwide).
- However silicium is strongly birefringent
 the signal needs to be orthogonal to the window
 partial redesign of the chamber.
- It was not possible to braze the new windows and silicium is very brittle
 - => Vacuum sealing was complicated







SPR data

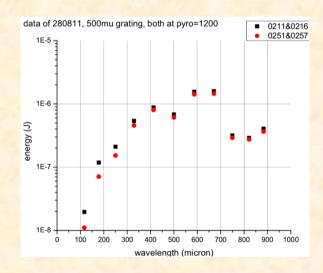


- Very preliminary data!
- We had some experimental issues at the beginning but we observed a nice once we understood our system.
- Data with 3 gratings still under investigation
- Relation between the gratings not as expected.

Going from the data to the bunch profile

$$\left(\frac{\mathrm{dI}}{\mathrm{d}\Omega\mathrm{d}\omega}\right)_{\mathrm{N_e}}(\Omega,\omega) = \left(\frac{\mathrm{dI}}{\mathrm{d}\Omega\mathrm{d}\omega}\right)_{\mathrm{sp}}(\Omega,\omega) \cdot \left[\mathrm{N_e} + \mathrm{N_e}(\mathrm{N_e} + 1) \,|\, \mathrm{F}(\omega) \,|^2\right]$$

- The measured spectrum gives only the square of the Fourier Transform of the bunch profile.
- Some phase information has been lost in the process.
- It is necessary to "reconstruct" this missing information.



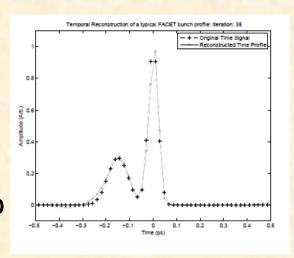
Phase recovery: Kramers-Kronig

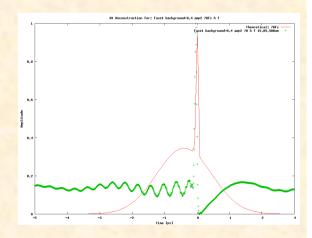
- Several methods exist to recover missing phase information (this is done every day by crystallographers at SOLEIL):
- In the Kramers-Kronig methods one tries to find the minimal phase satisfying all the constraints set by the measurement:

$$G(\omega) \approx N_e^2 |F(\omega)|^2$$
 $G(\omega) = \rho(\omega) \cdot \exp[i\phi(\omega)]$

$$\phi(\omega) = \frac{2\omega}{\pi} \int_{0}^{\omega} \frac{\ln[\rho(\omega)/\rho(\omega_0)]}{\omega_0^2 - \omega^2} d\omega$$

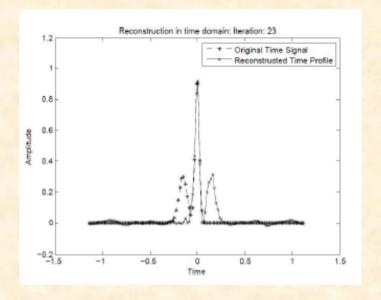
 However the results produced by this methods are sometimes non physical.

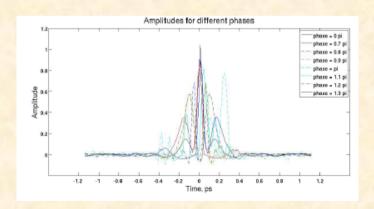




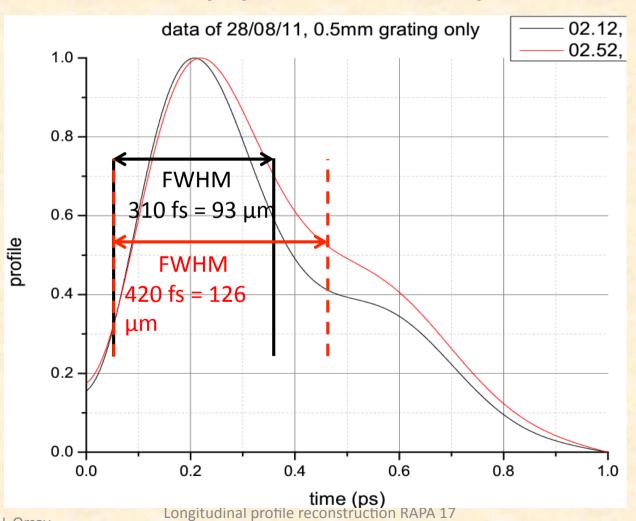
Shrink wrap method

- To address the problems of the KK methods we also use another method inspired from crystallography: the shrink wrap method.
- In that method an initial phase is chosen and used to reconstruct the profile.
- Tests are then applied on the reconstructed profile. If the profile fails the tests, then the phase of one parameter is slightly changed and the tests are applied again.
- However this method may lead to a large number of solutions/profiles.
 => The reconstructed profile is not necessarily unique. (R.Tovey / C.Santamaria)



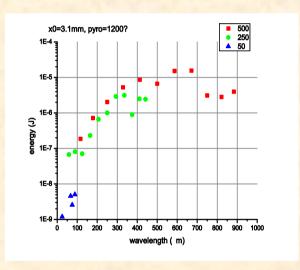


Reconstructed profiles (KK method) Very preliminary



Remaining issues

- At the moment we do not understand the data produced by the shorter grating: the signal intensity is much mower than expected.
- We are still comparing the results obtained with the Shrink wrap method and the KK method.
- We identified several operational problems (data alignment, grating/ filters mismatch, ...) that we need to address.

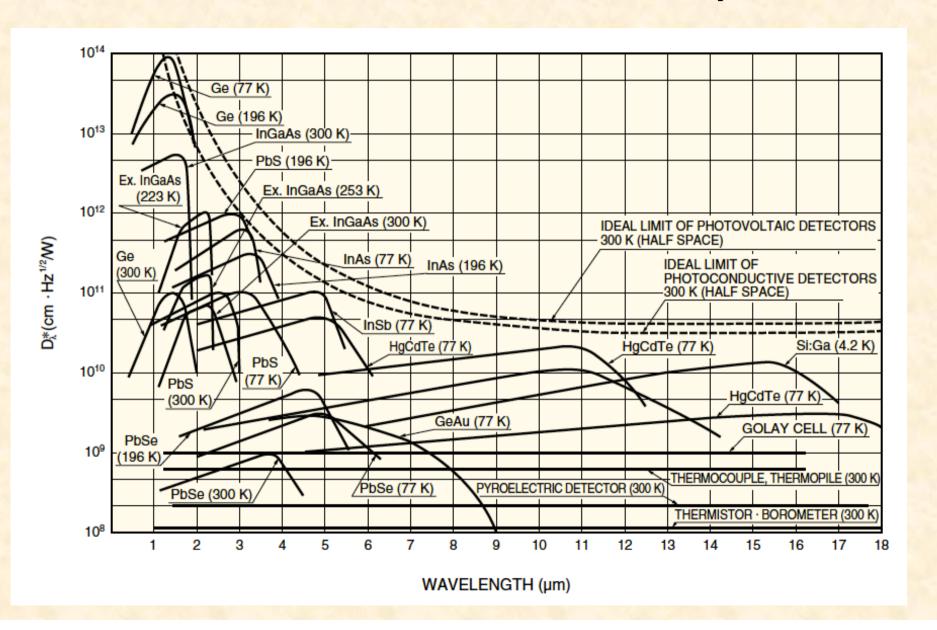


Further developments

- E-203 is currently approved until Summer 2012.
- We have applied this week for an extension until summer 2013.
- Our preliminary results have triggered a lot of interest from the community.
- To reach our goal and deliver a device suitable for a LPA we need several improvements:
 - Make it single shot (3 gratings in a row? Gratings around the beam pipe?...)
 - Increase the sensitivity of the detectors (Bolometers???)
 - Extend the ability to shorter pulses (1-50fs), deeper in the THz gap and reaching the water absorption band.

=> ANR

IR detectors sensitivity

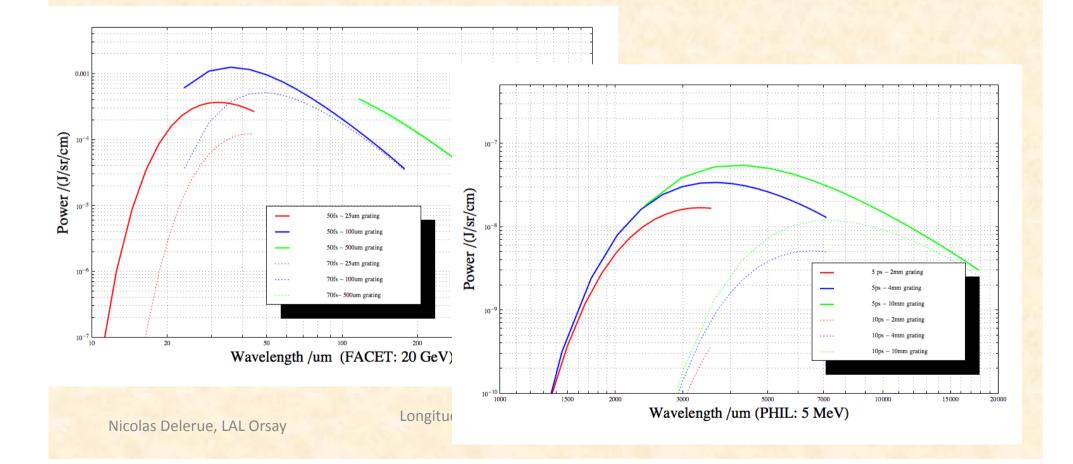


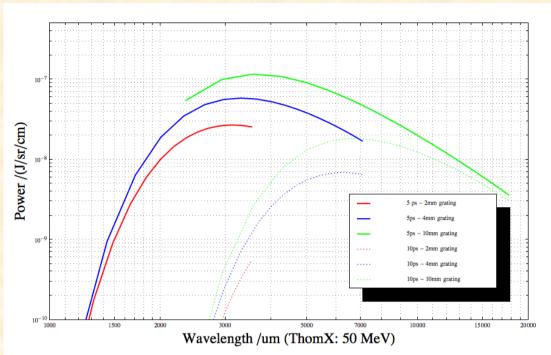
PHIL, ThomX, SOLEIL, LPA

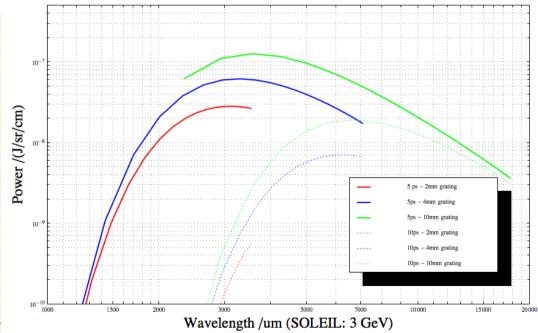
$$\left(\frac{\mathrm{dI}}{\mathrm{d}\Omega}\right)_{\mathrm{sp}} = 2\pi q^2 \frac{Z}{\ell^2} \frac{n^2 \beta^3}{\left(1 - \beta \cos \theta\right)^3} R^2 \exp\left(-\frac{2x_0}{\lambda_e}\right)$$

$$\left(\frac{dI}{d\Omega}\right)_{sp} = 2\pi q^2 \frac{Z}{\ell^2} \frac{n^2 \beta^3}{(1 - \beta \cos \theta)^3} R^2 \exp\left(-\frac{2x_0}{\lambda_e}\right) \qquad \left(\frac{dI}{d\Omega d\omega}\right)_{N_e} (\Omega, \omega) = \left(\frac{dI}{d\Omega d\omega}\right)_{sp} (\Omega, \omega) \cdot \left[N_e + N_e(N_e + 1) \mid F(\omega) \mid^2\right]$$

45 MeV, few ps already tested at FELIX







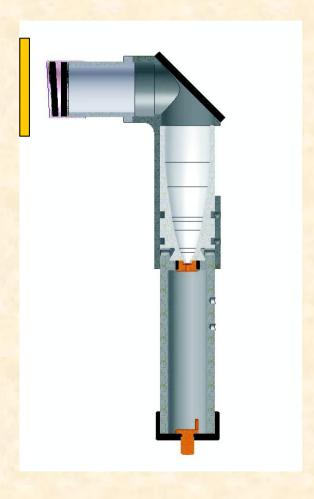
Summary

- We are investigating SPR as a diagnostic in the fs range.
- Most hardware difficulties have been addressed.
- Currently trying to understand the data.
- First results promising.
- Step toward SPR at a LPA.

Thank you for your attention

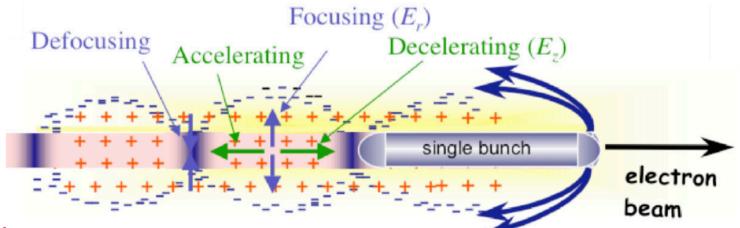
I am grateful to all those at LAL who have helped with the preparation of this experiment and in particular the design office and the vacuum group.

Detector assembly



Plasma Wakefield Acceleration

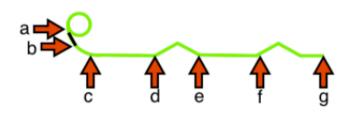
- Continuation of very successful experiments at SLAC
 - Single electron bunch in field ionized lithium plasma
 - Single bunch in field ionized caesium plasma
 - Use Notch Collimator
 - Two electron bunches (drive and witness) in field ionized caesium plasma
- Second stage: pre-ionized plasma, positrons and electrons





Bunch Compression

Three-stage bunch compression:



- On entering linac from damping rings
- Magnetic chicane in S10
- Magnetic chicane in S20

