

# Dark Matter: Candidates, signals and LHC consequences

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Yann Mambrini, LPT Orsay



LAL seminar, february 21th 2012, Paris

# Outline

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- Dark matter evidences : scales and methodology

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- Dark matter detection : principle and experimental status
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- 2 examples :  $Z'$  and Higgs-portal
- Complementarity with the LHC

# Astroparticle

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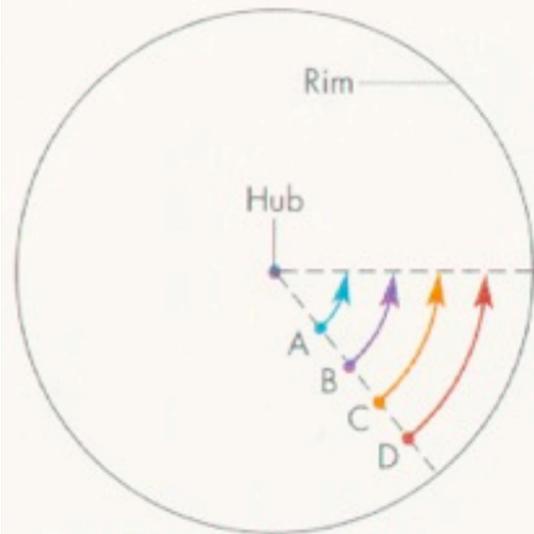
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- LHC provides 600 millions collisions per seconds whereas we just have 1 Universe : we cannot reproduce the experiment to increase the luminosity!!!
- LHC provides his own background, whereas in the Universe, you have no idea of the background as you always measure Signal+Background.

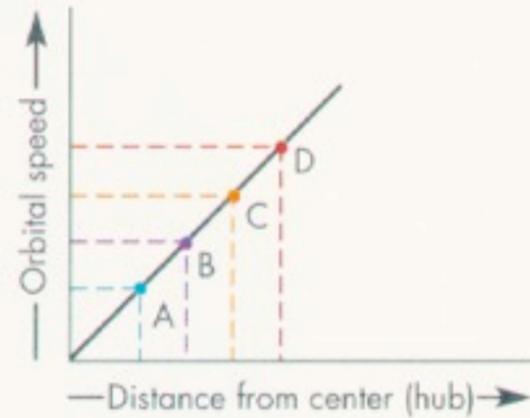
# Dark matter evidence : Galactic scale



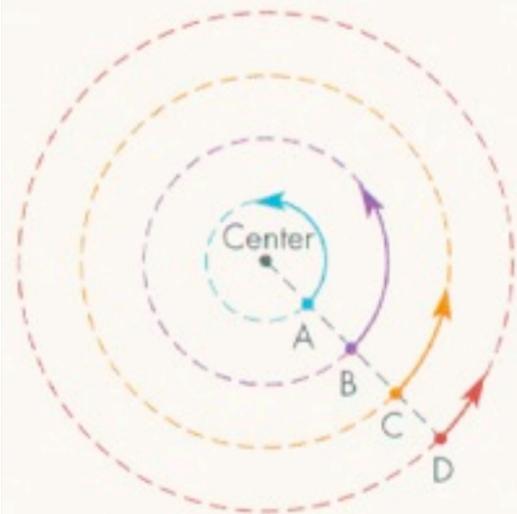
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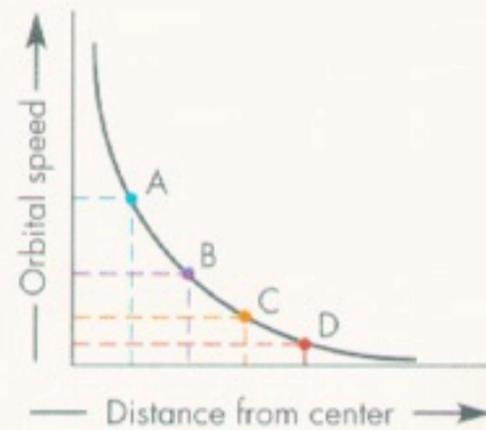
Wheel-like rotation



Rotation curve for wheel-like rotation

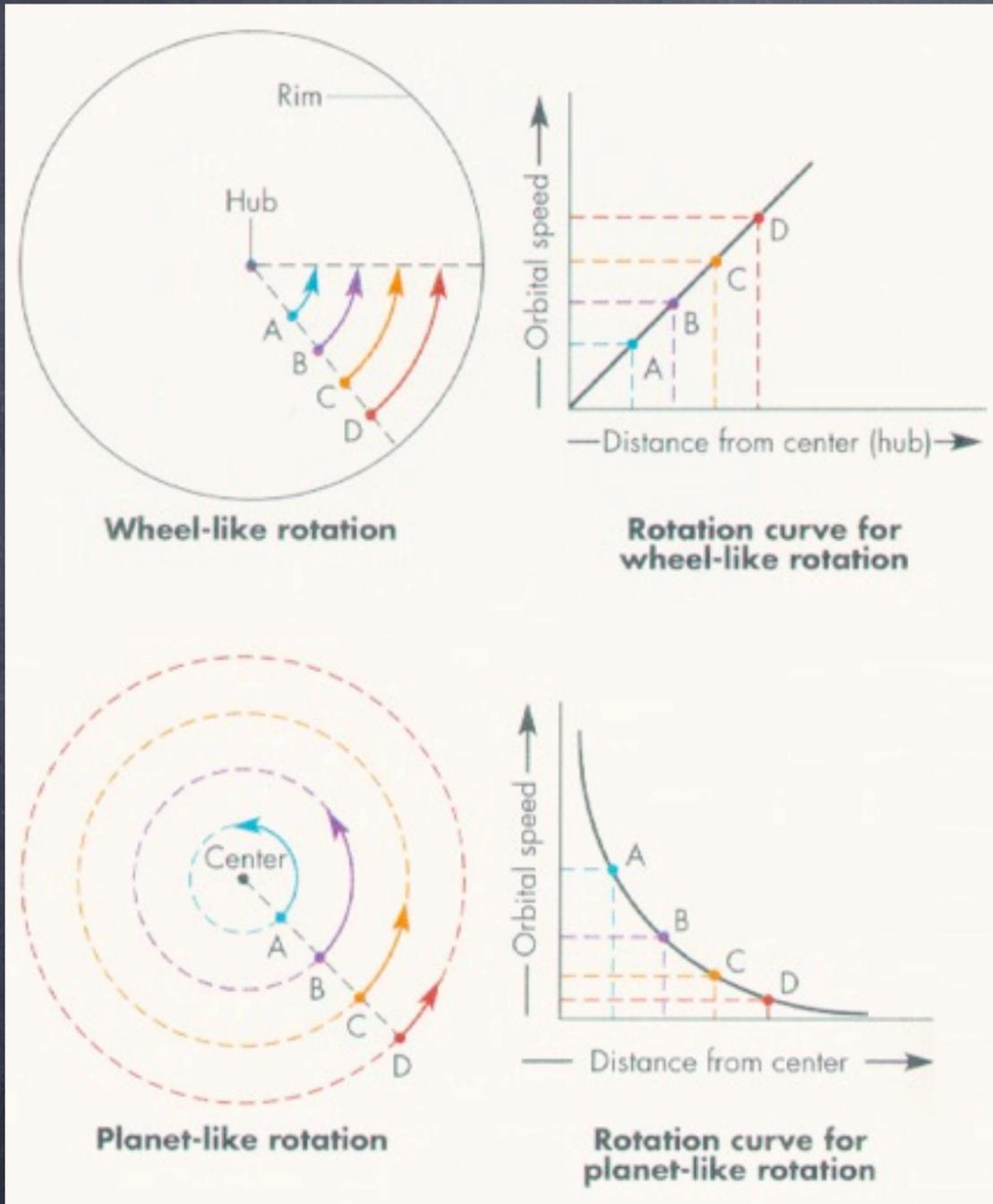


Planet-like rotation



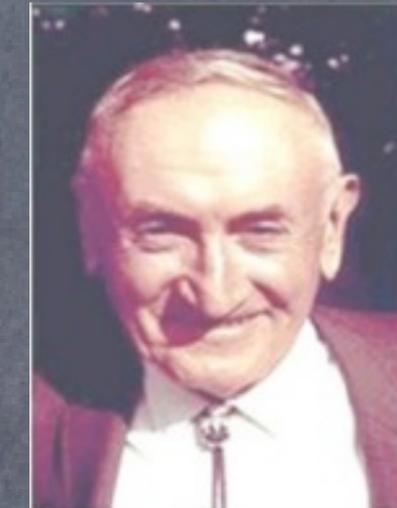
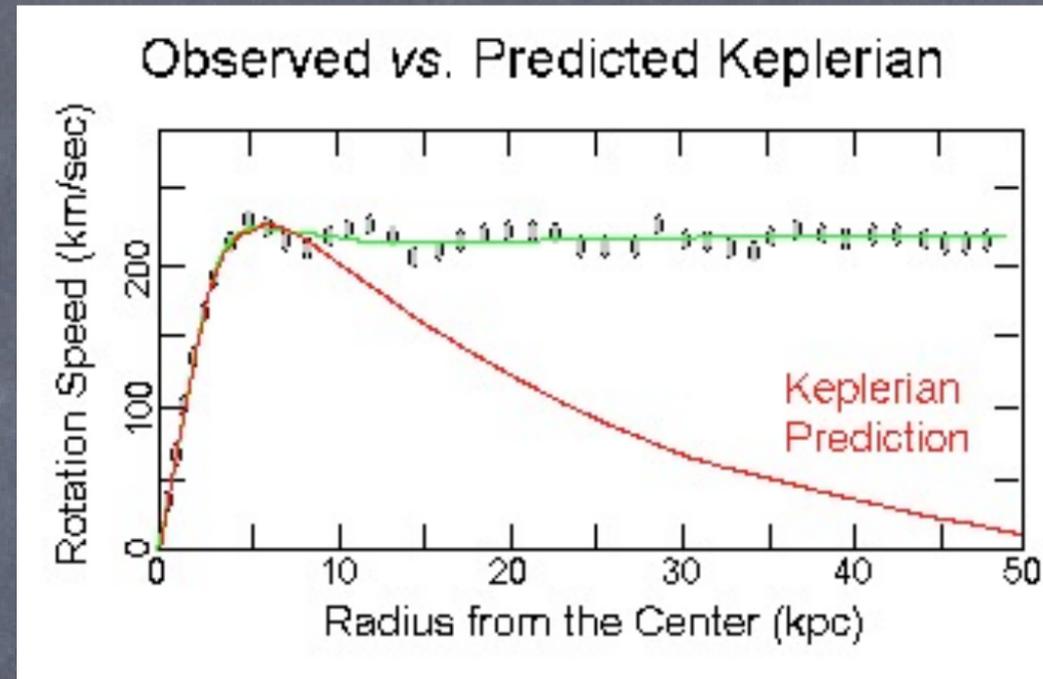
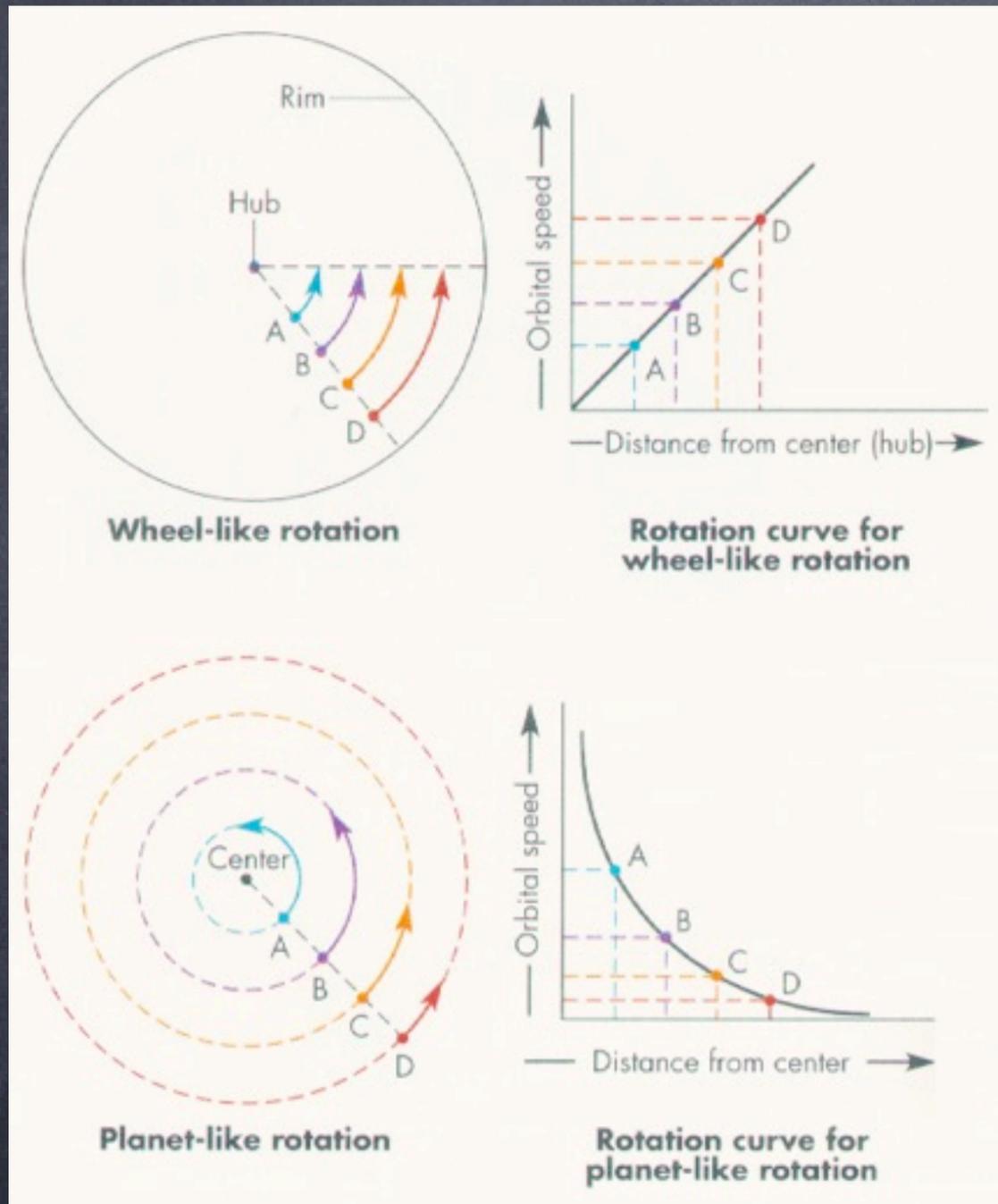
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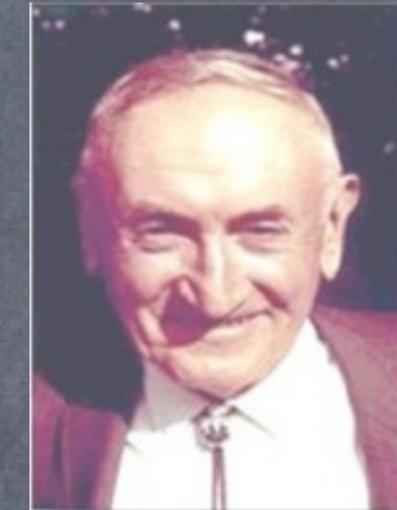
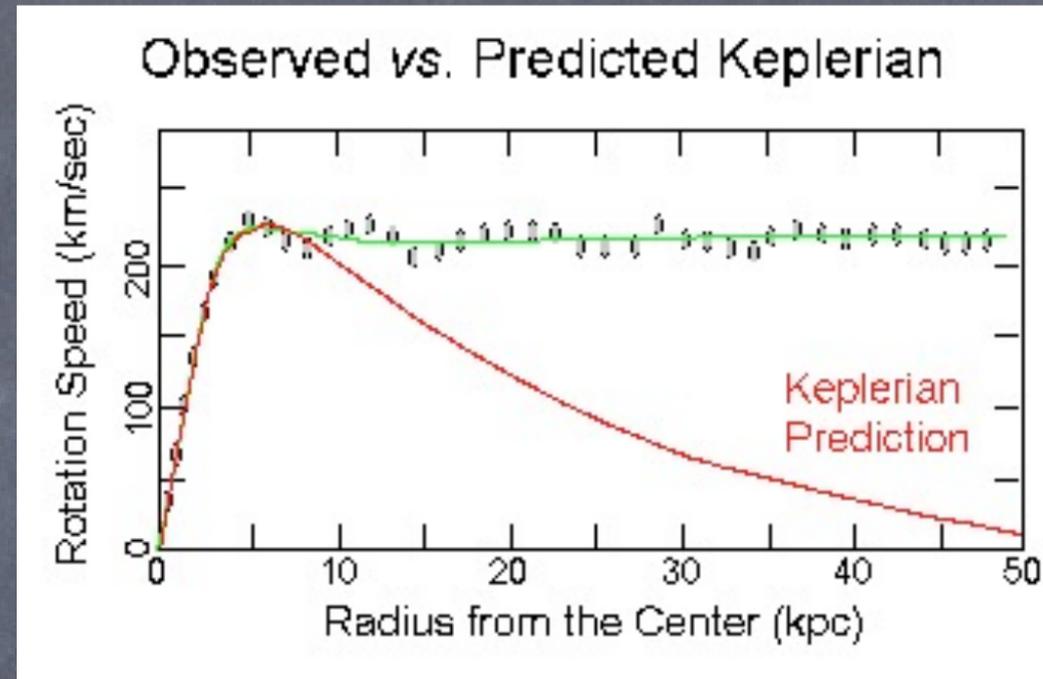
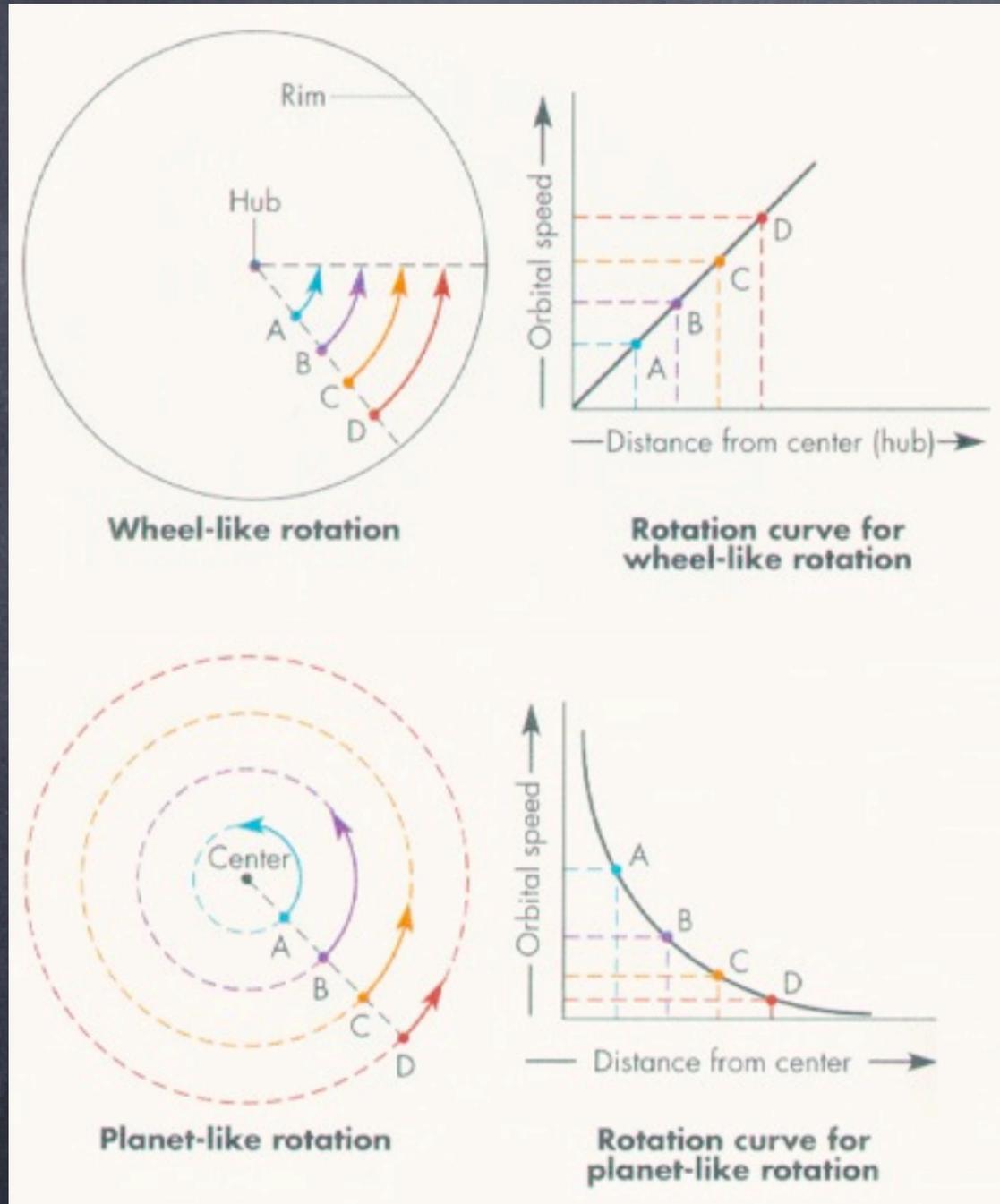
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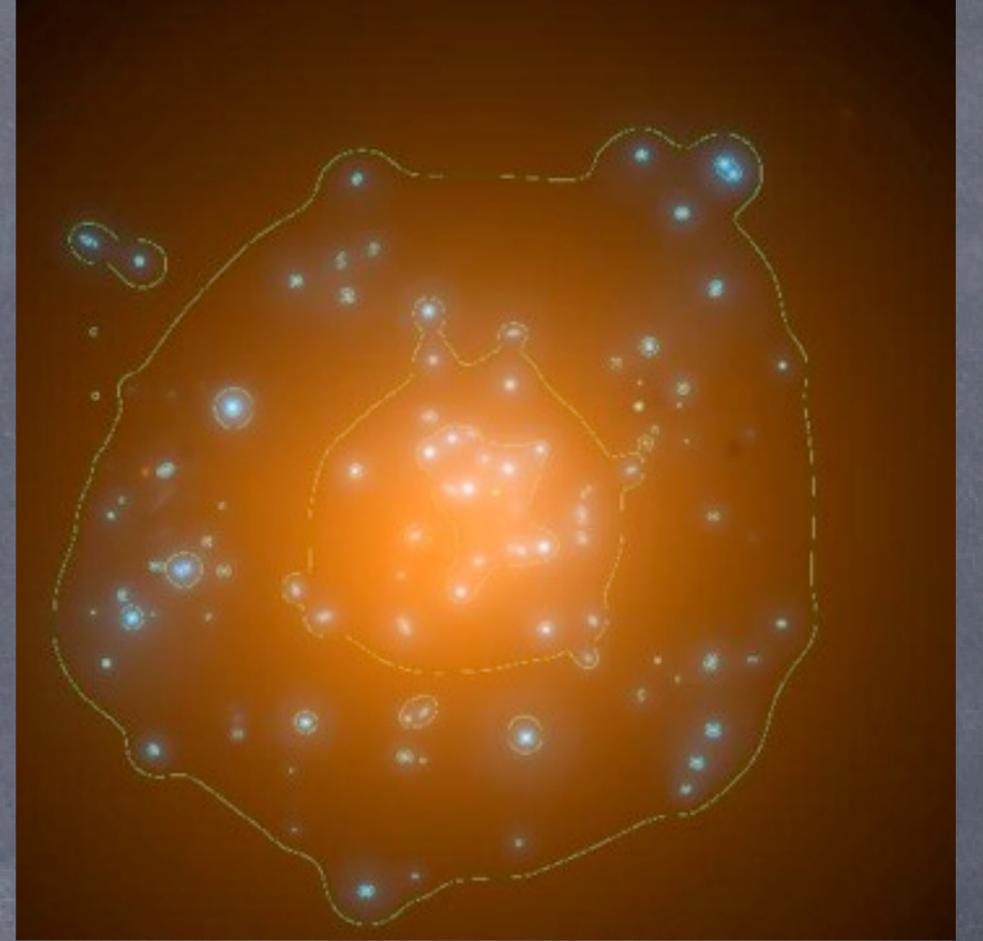
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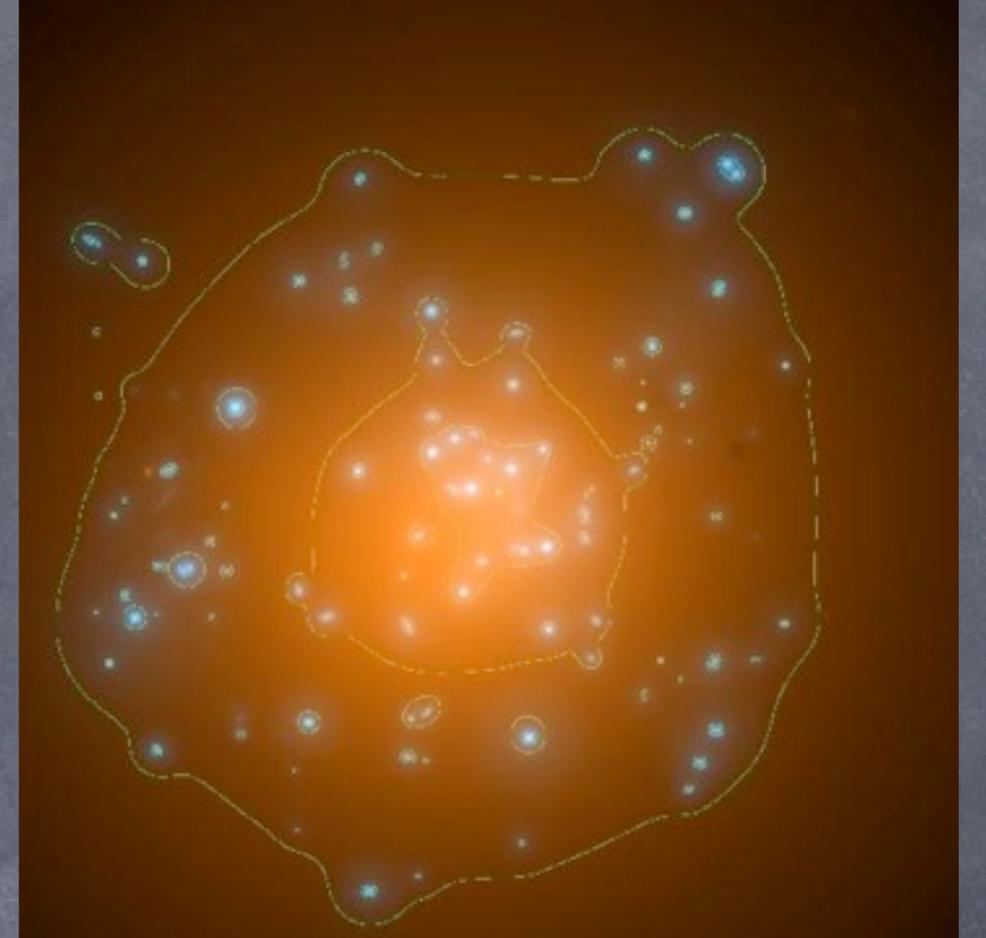
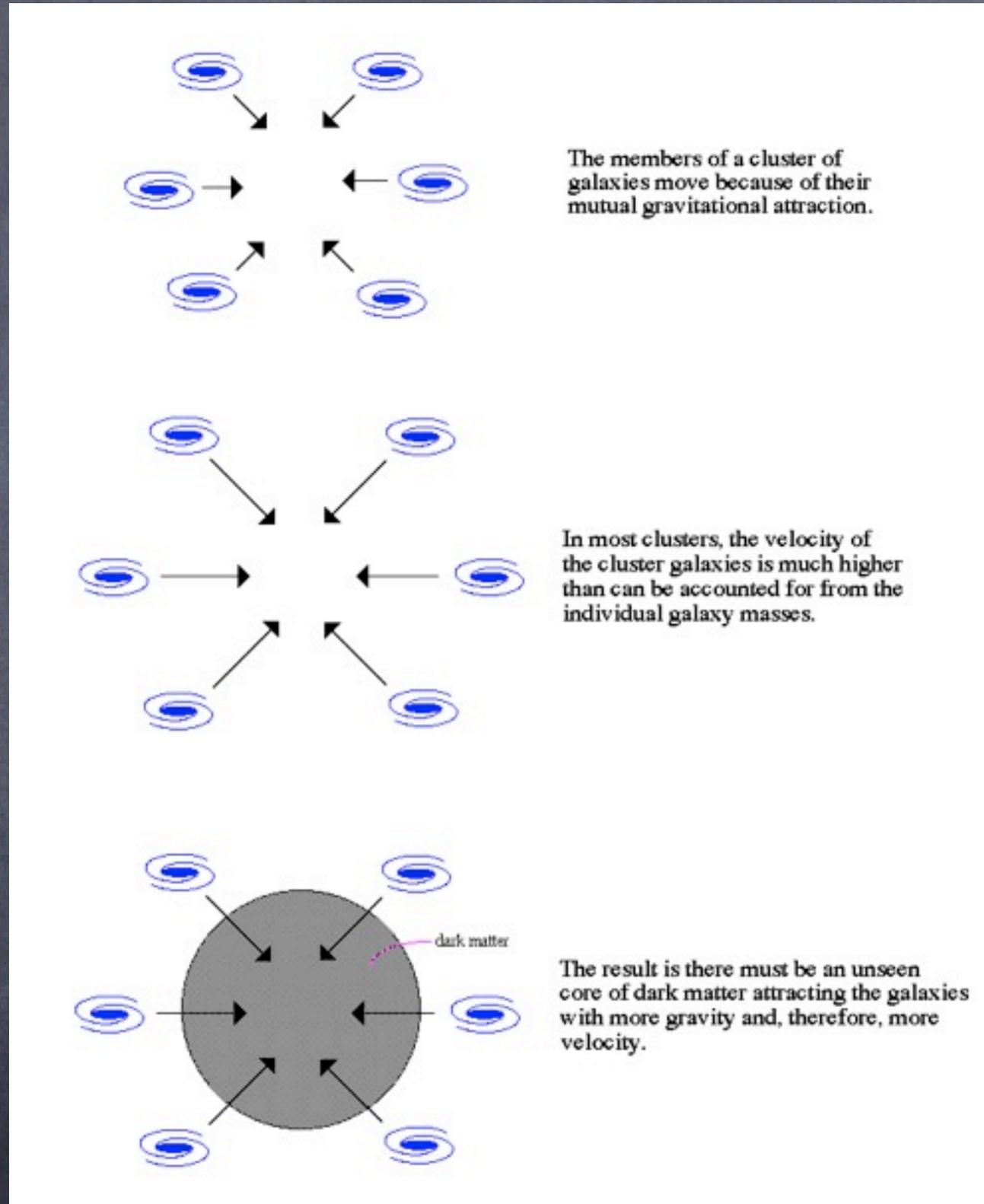
$$\rho_{DM} \propto \frac{1}{r^2} \rightarrow M_{gal} = Volume * \rho \propto r \rightarrow v \sim cte$$

$$\frac{GmM_{gal}}{r} = mv^2 \rightarrow v \propto \frac{1}{\sqrt{r}}$$

# Cluster of Galaxy scale



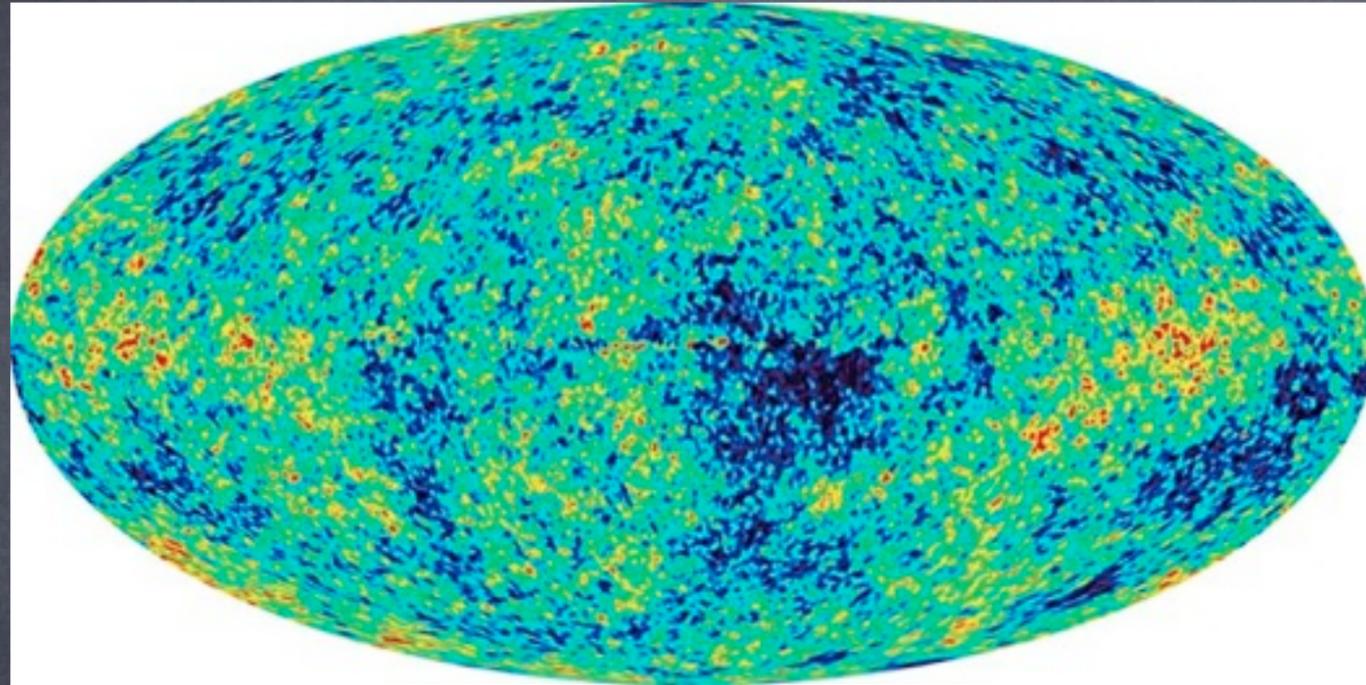
# Cluster of Galaxy scale



# Cosmological scale

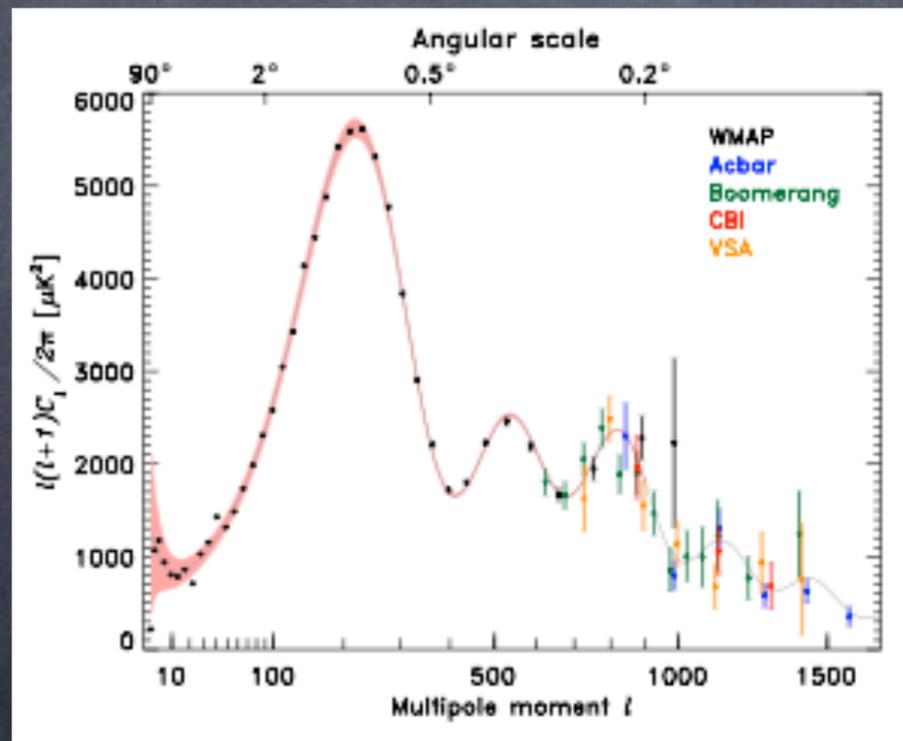
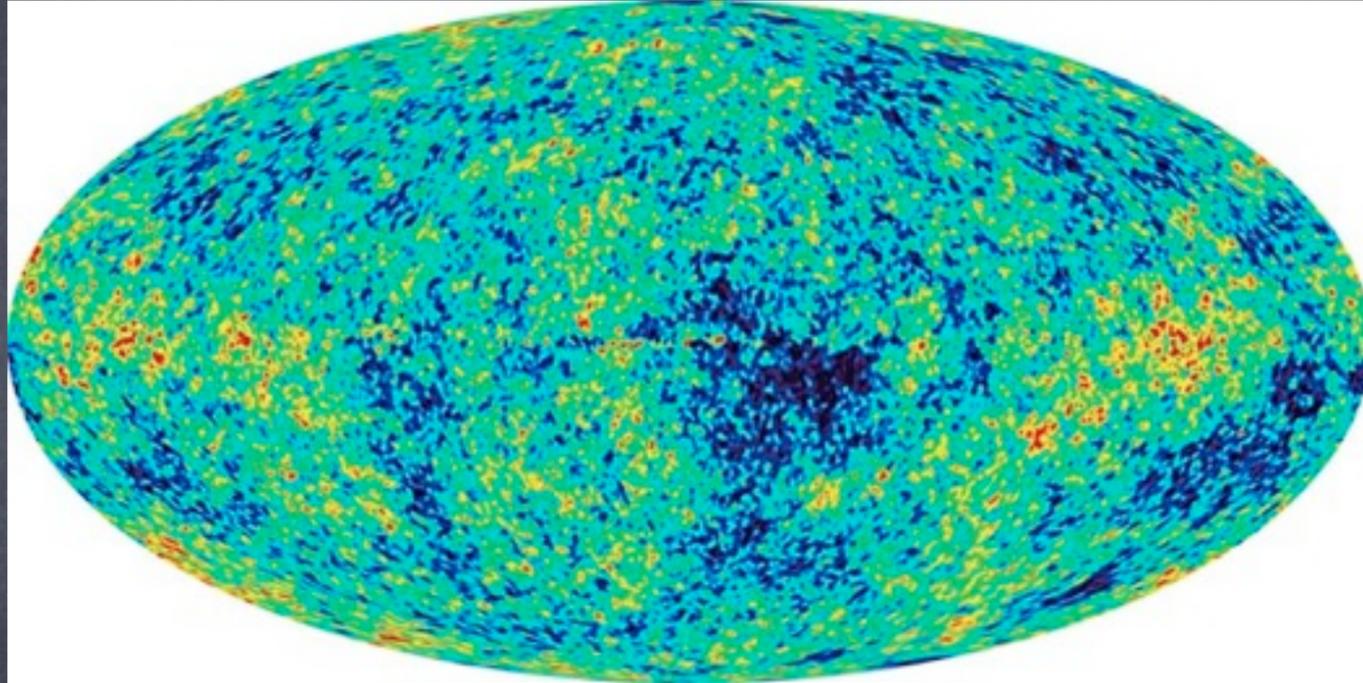
# Cosmological scale

Cosmic Microwave Background (CMB)



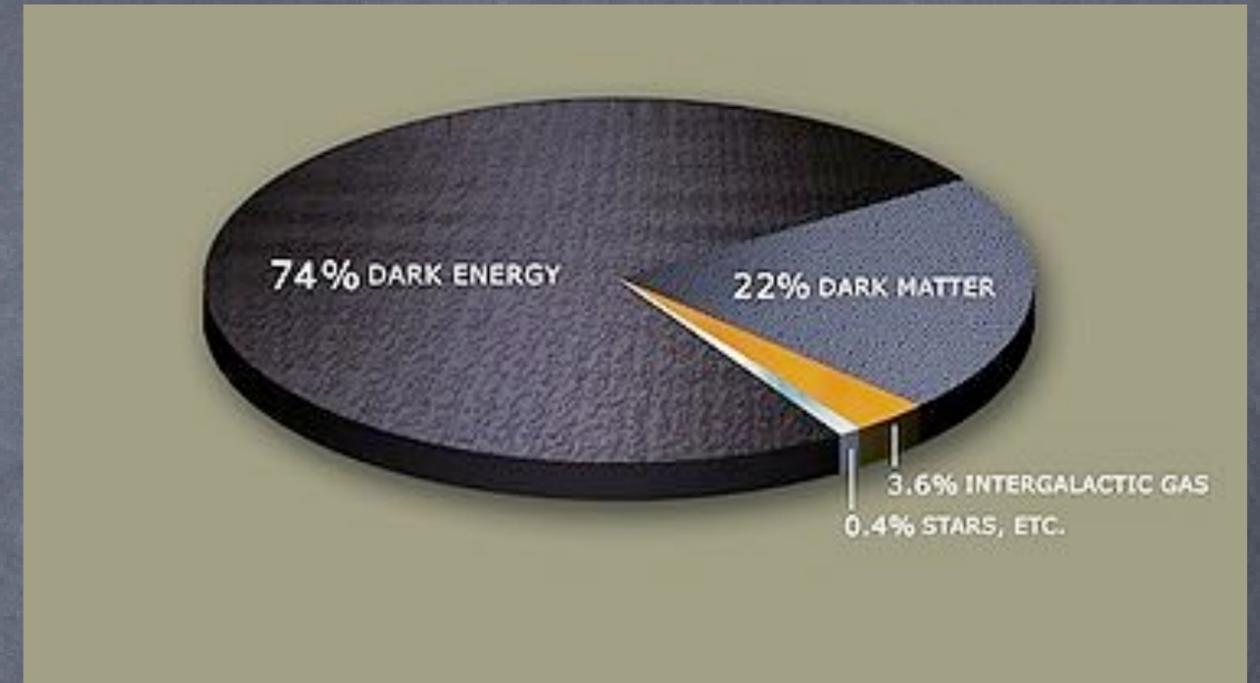
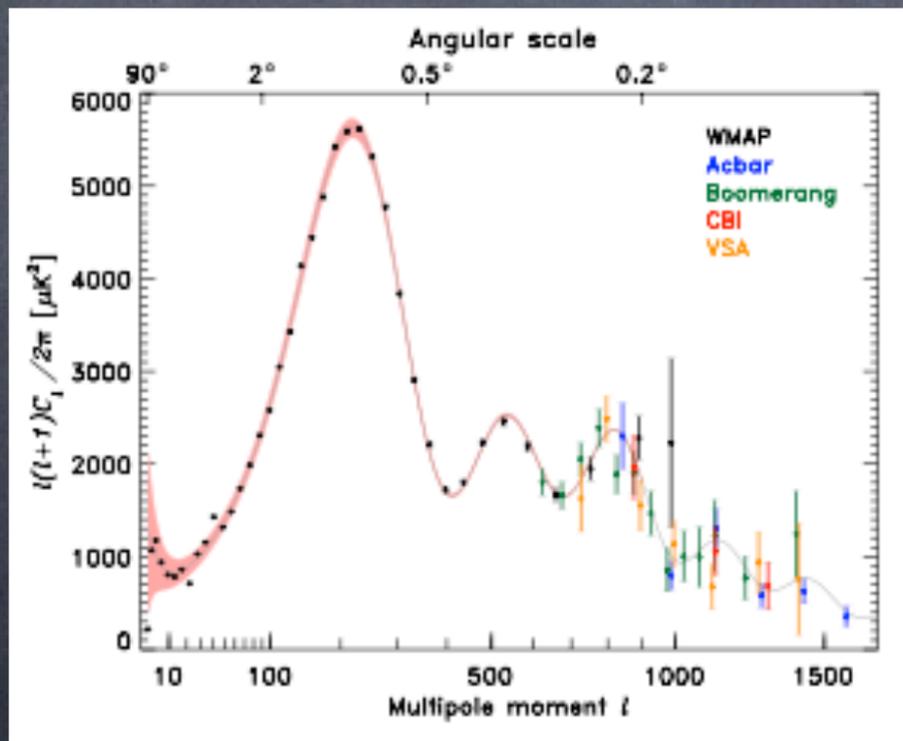
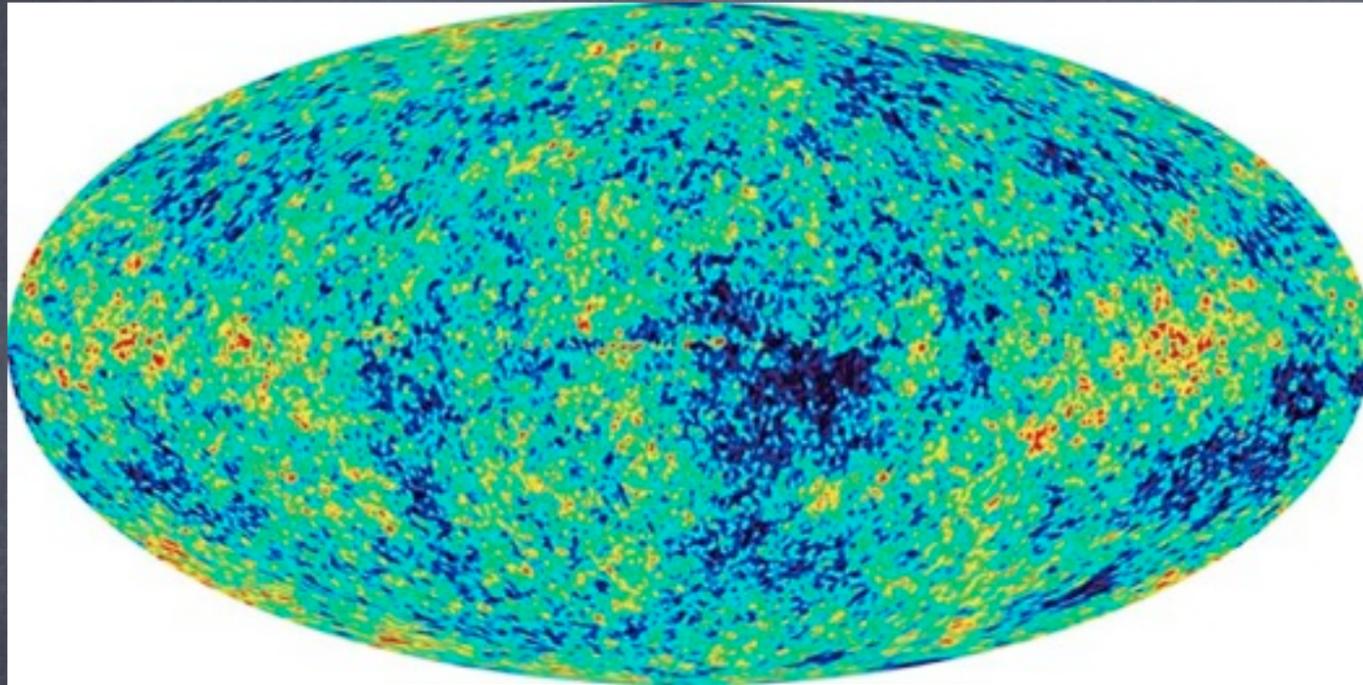
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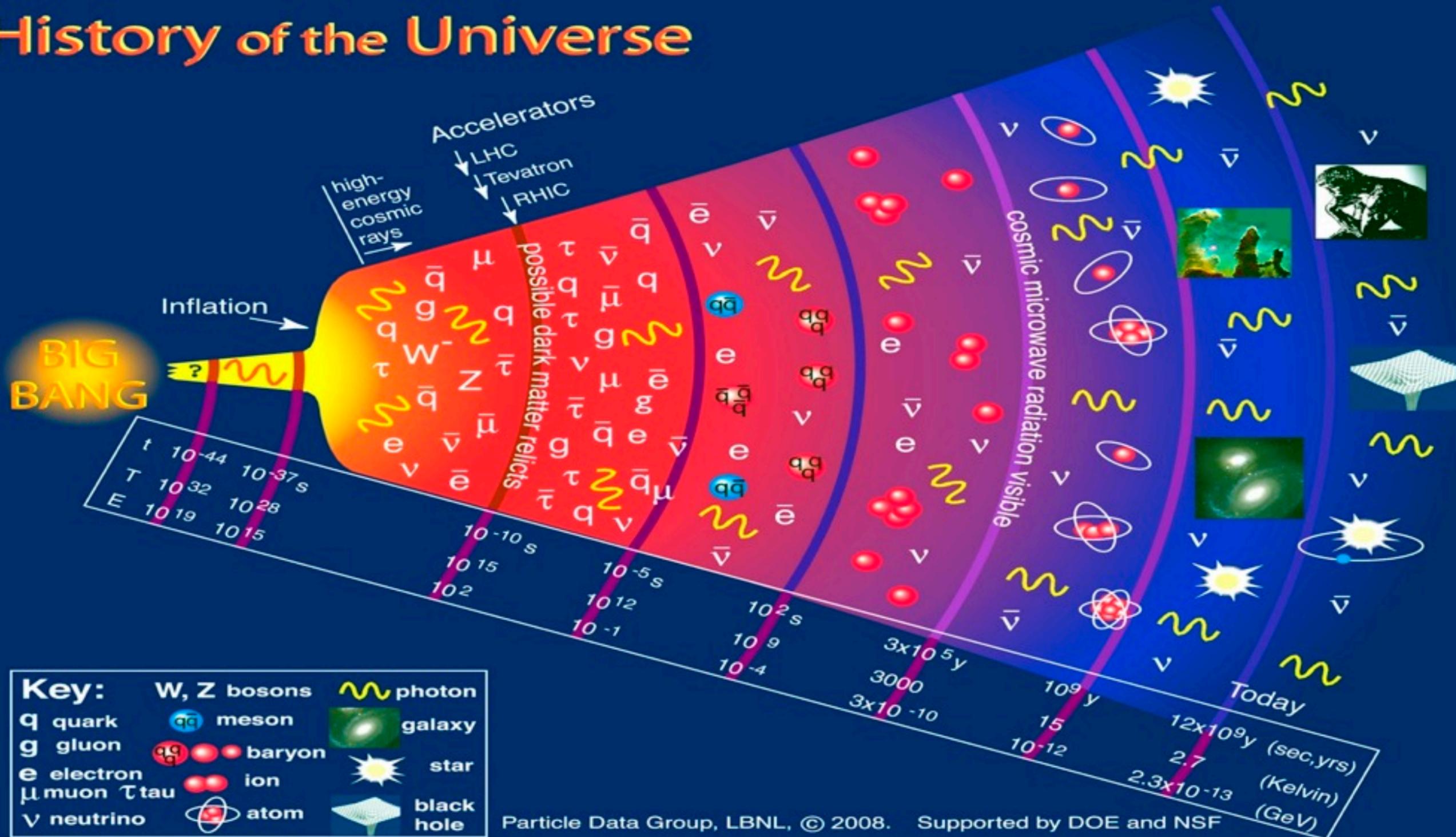
Budget of the Universe

# A Little thermal history of the Universe (I)

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## History of the Universe



# A little thermal history of the Universe (II)

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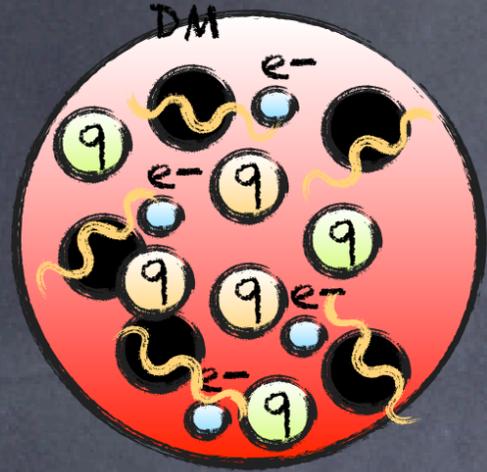
T

time



# A little thermal history of the Universe (II)

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$$f \sim \frac{e^{-\frac{E}{T}}}{1 \pm e^{-\frac{E}{T}}}$$

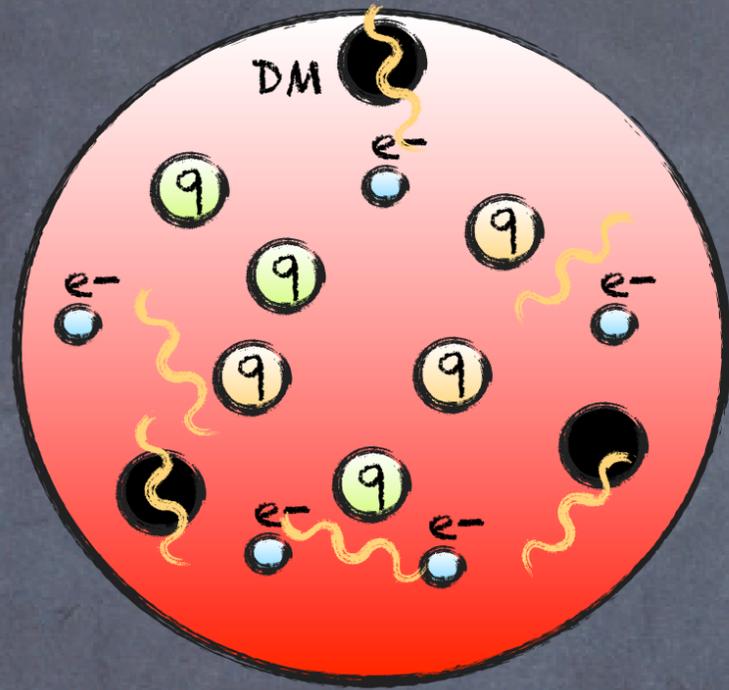
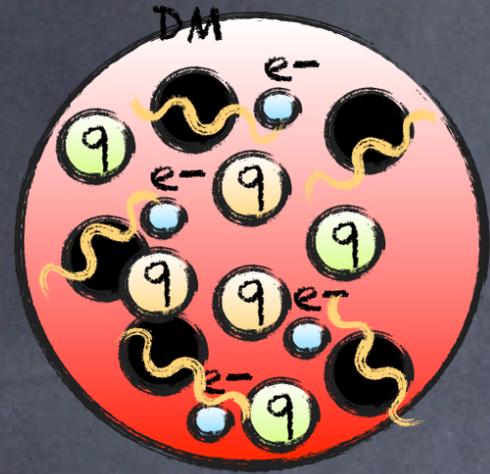
Equilibrium:  
Thermal bath

T

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# A little thermal history of the Universe (II)



$$f \sim \frac{e^{-\frac{E}{T}}}{1 \pm e^{-\frac{E}{T}}}$$

$$\rho_{DM} \sim e^{-\frac{m}{T}}$$

Equilibrium:  
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Dark Matter  
decoupling

T

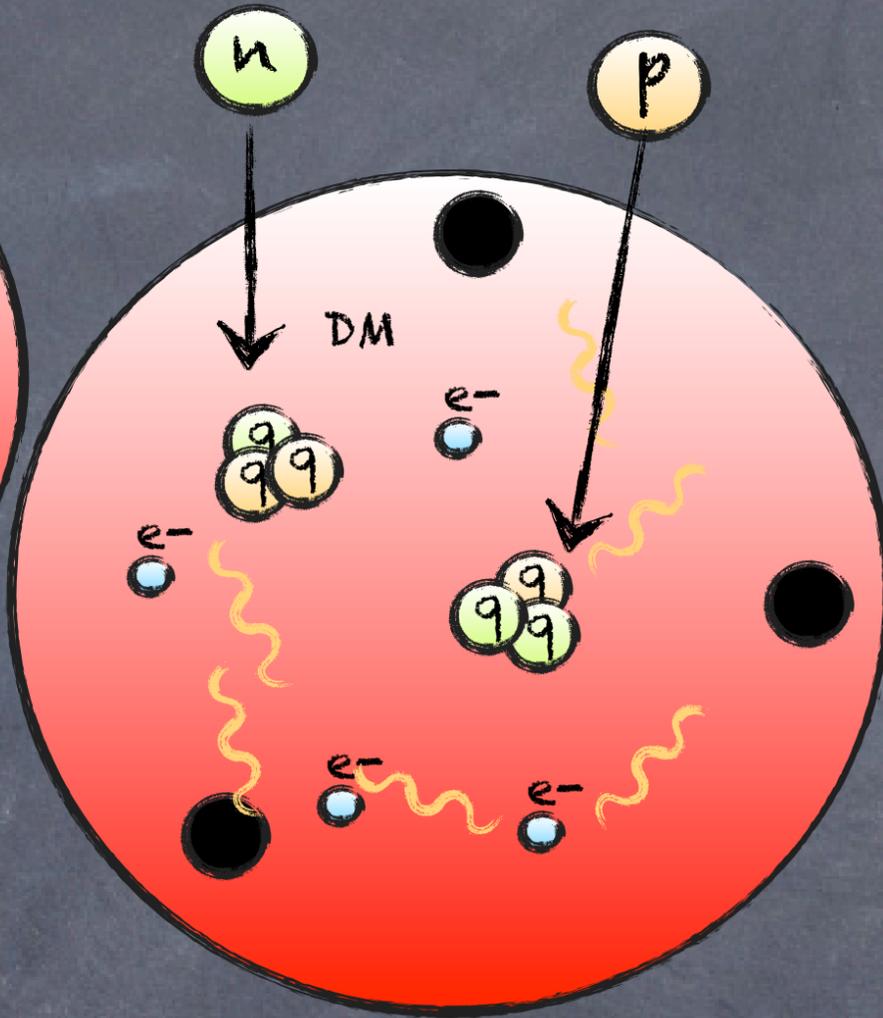
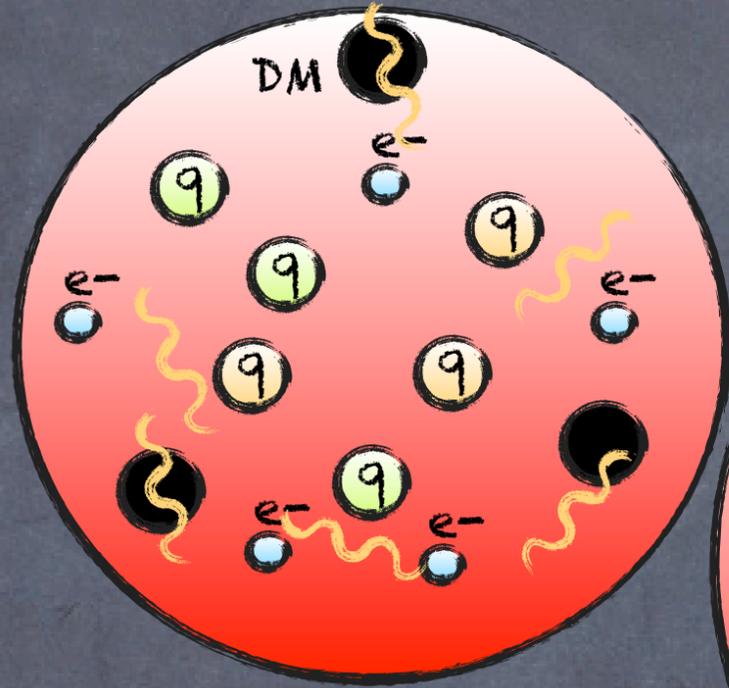
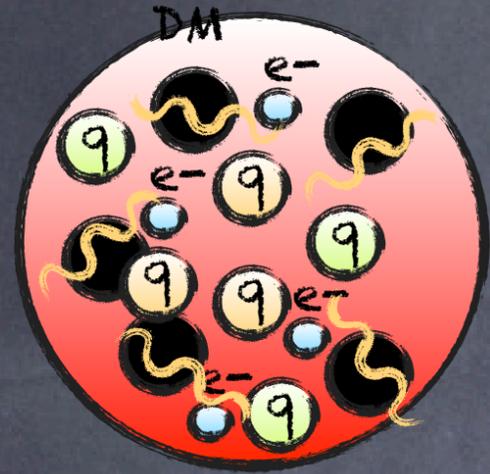
10 GeV

10<sup>-5</sup> sec

time



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Equilibrium:  
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Nucleosynthesis

T

10 GeV

300 keV

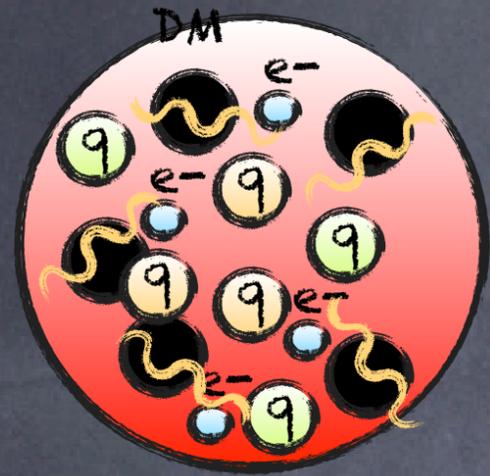
10<sup>-5</sup> sec

1 min

time

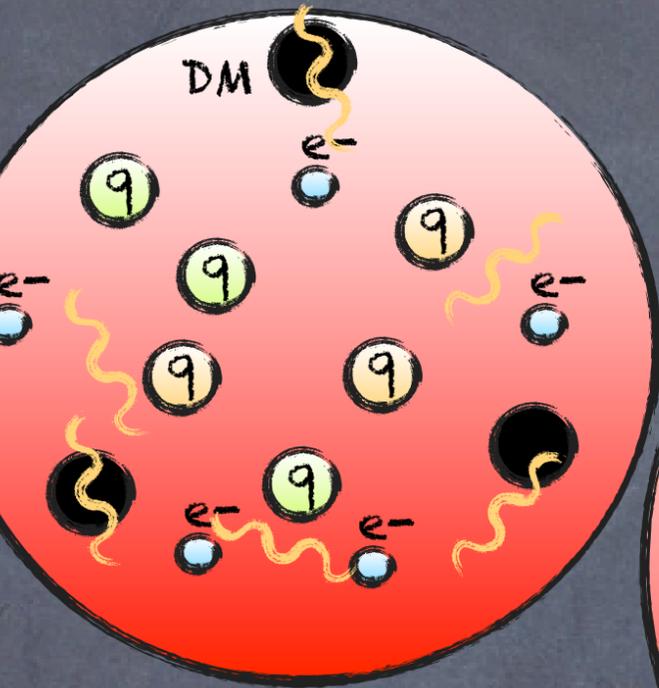


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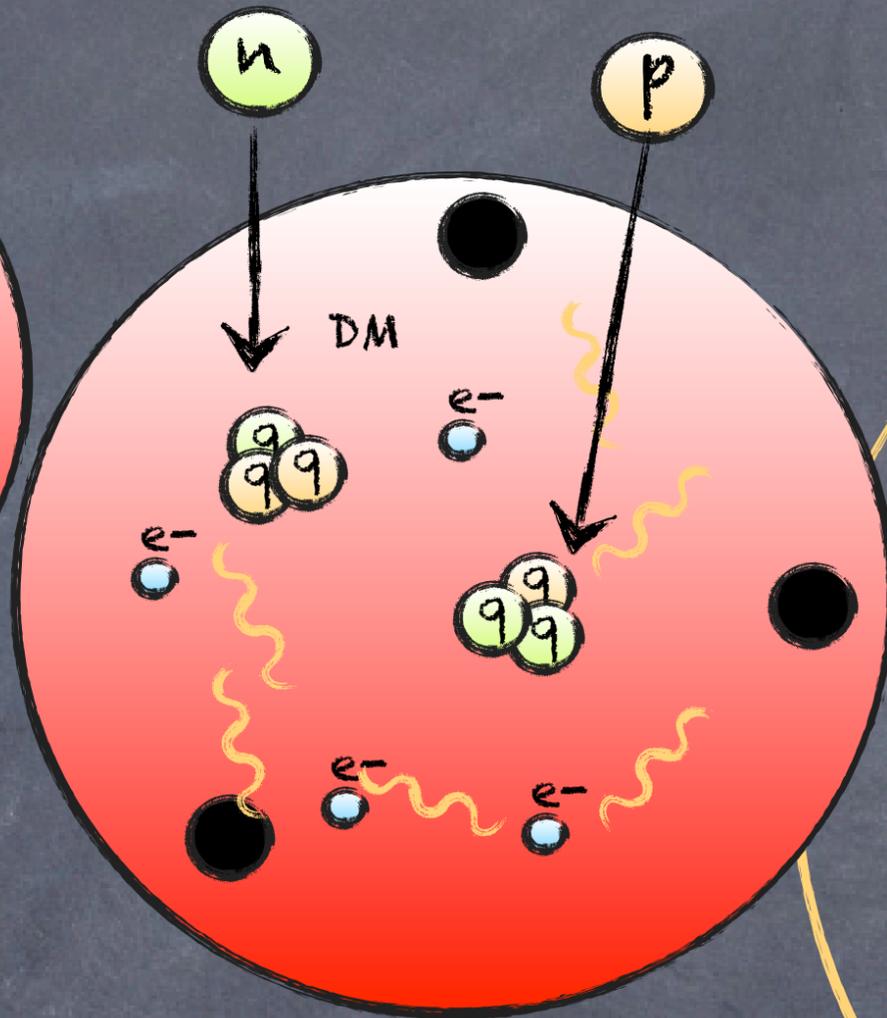
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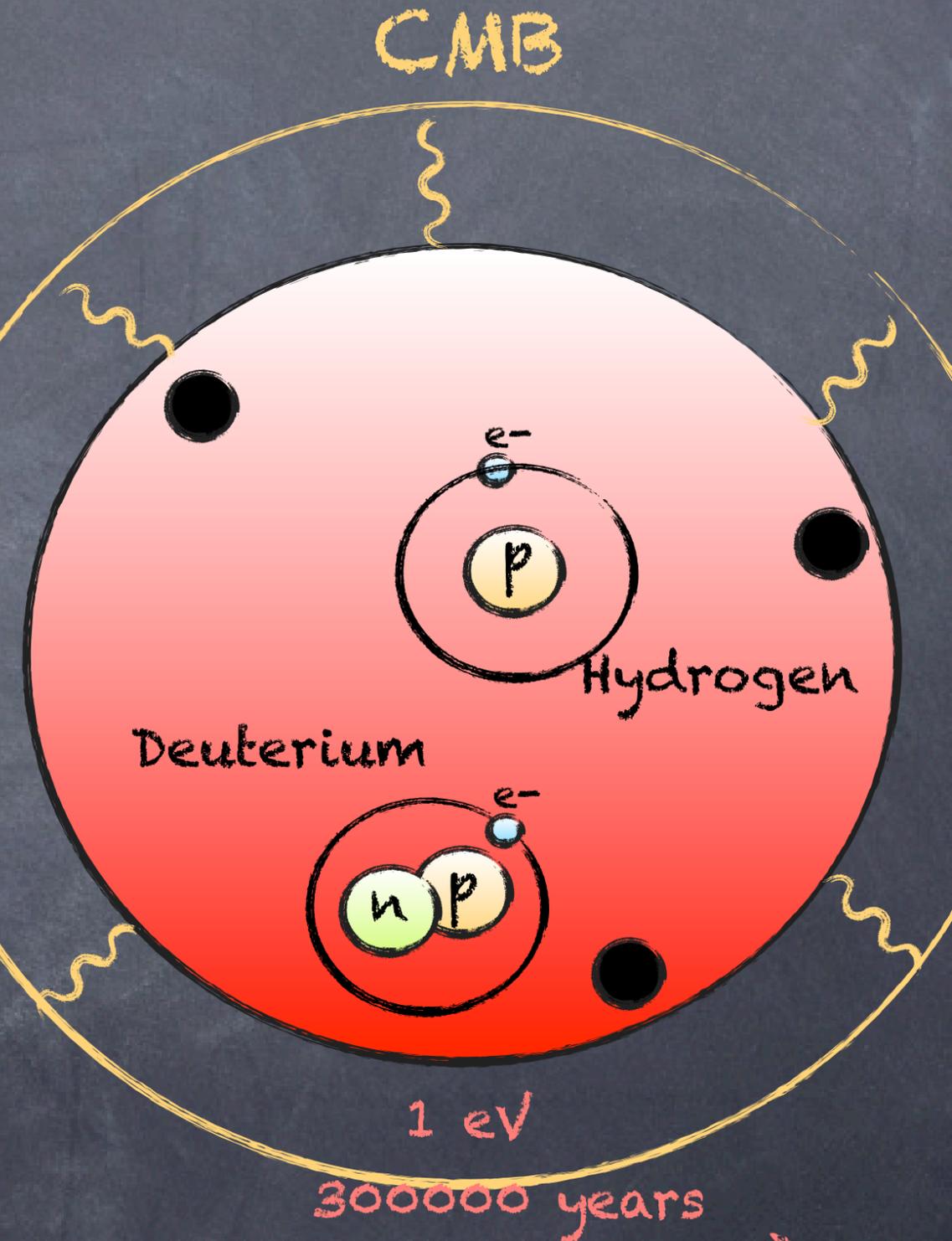
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10 GeV  
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Nucleosynthesis

300 keV  
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1 eV  
300000 years

T  
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An example : the recombination temperature

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Suppose a plasma made  
of electron, proton,  
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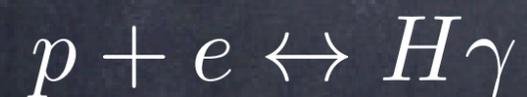


# An example : the recombination temperature

The density of each species can be written

$$n_i = g_i \left( \frac{2\pi}{m_i T} \right)^{3/2} e^{\frac{\mu_i - m_i}{T}}$$

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Chemical equilibrium  
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$$n_H = \frac{g_H}{g_p g_e} n_p n_e \left( \frac{2\pi}{m_e T} \right)^{3/2} e^{\frac{B_H}{T}}$$

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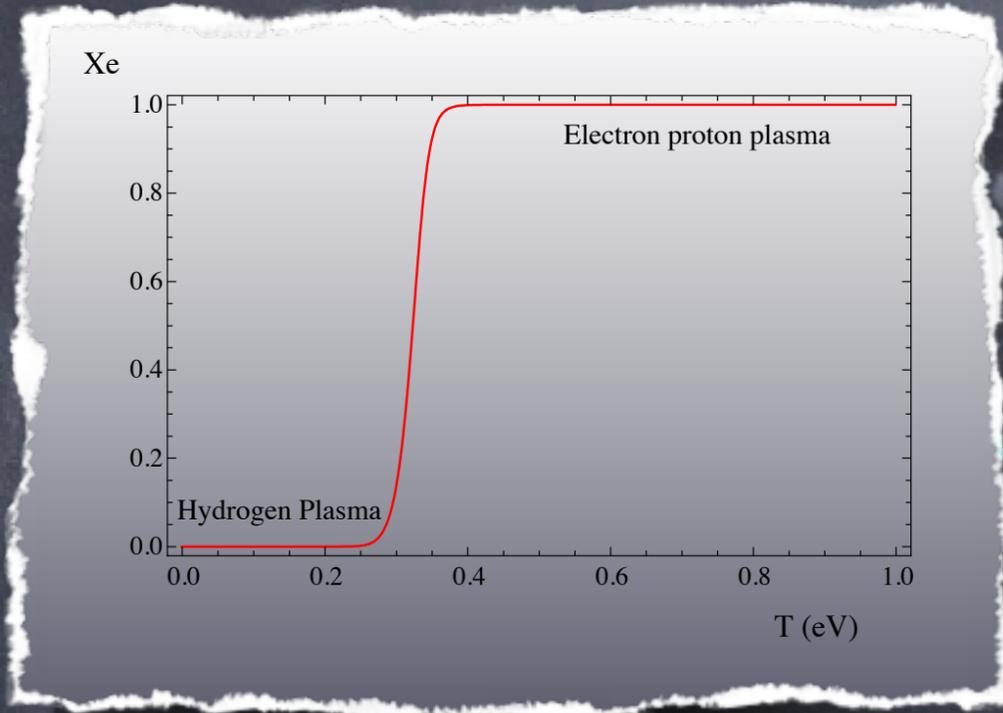
$$X_e = \frac{n_e}{n_B} = \frac{n_e}{n_p + n_H}$$

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$$1 - X_e = X_e^2 \frac{2.68 \times 10^{-8} 4\sqrt{2}\zeta(3)}{\sqrt{\pi}} \left( \frac{T}{m_e} \right)^{3/2} e^{\frac{B_H}{T}} \Omega_B h^2$$

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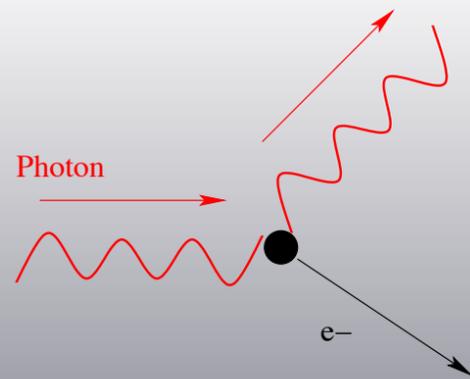
When  $T \sim B_H \Rightarrow$  Exponential dominates : the photon is not able anymore to destroy the atom of Hydrogen



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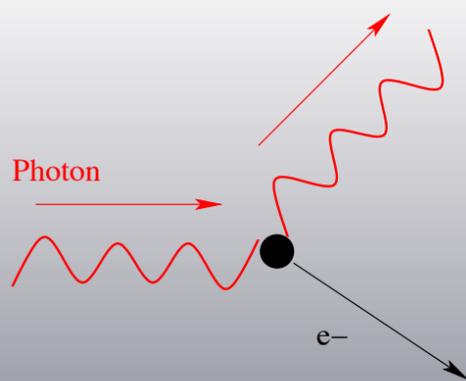


Thomson scattering

$$\Gamma_{\gamma} = \sigma_T c n_e \quad [\text{s}^{-1}]$$

$$\sigma_T = \frac{8\pi}{3} \left( \frac{q^2}{4\pi\epsilon_0 m_e c^2} \right)^2$$

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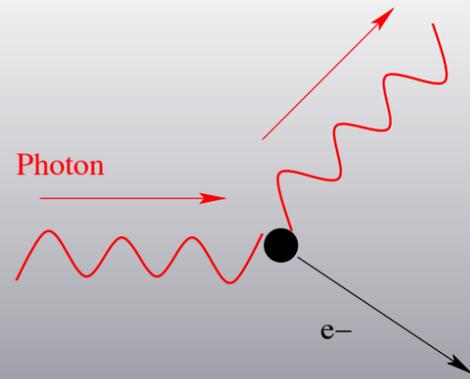
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then the photon travel a distance larger than the Hubble horizon

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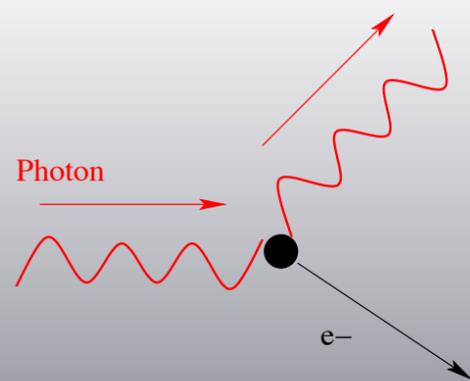
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It happens for  $T_{\text{dec.}} = 0.234 \text{ eV} = 280\,000 \text{ years}$

Remark :  $T_{\text{dec}}$  is lower than BH ( $0.2 < 13.6 \text{ eV}$ ) because even if  $T < \text{BH}$ , still sufficiently numerous photons with  $E_\gamma > 13.6 \text{ GeV}$  to destroy the atom of Hydrogen. (both  $\sigma$  and  $n$  contribute to annihilation rate)

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Boltzmann Equation

$$\frac{dn}{dt}$$

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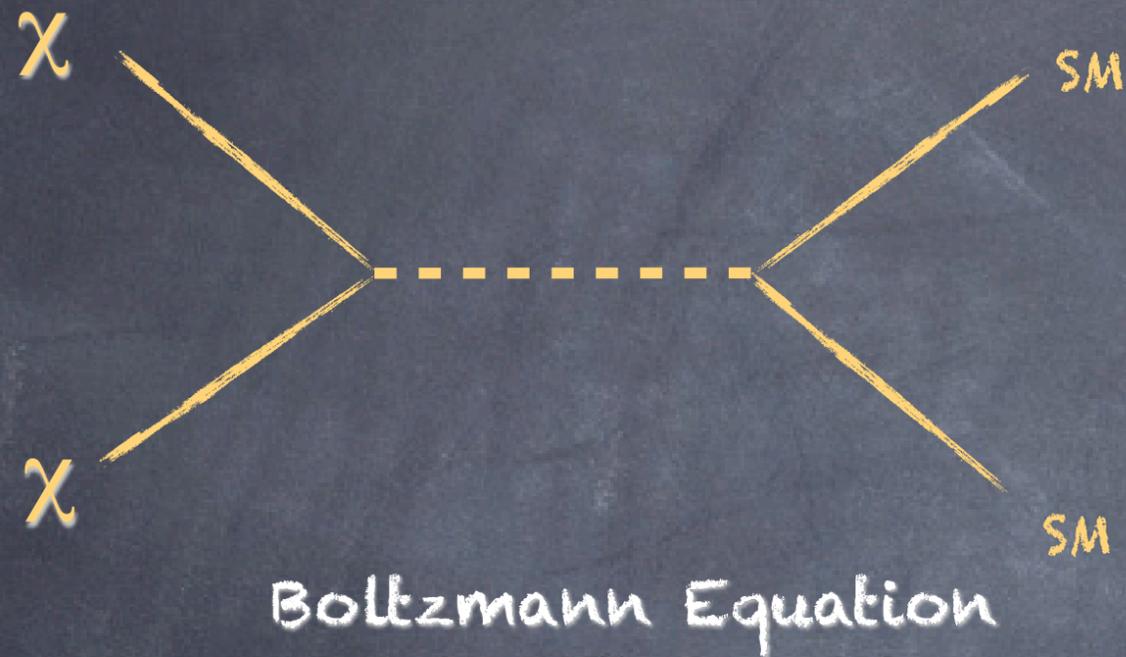
Boltzmann Equation

$$\frac{dn}{dt} = -3 H n$$

$$(H = \dot{R} / R)$$

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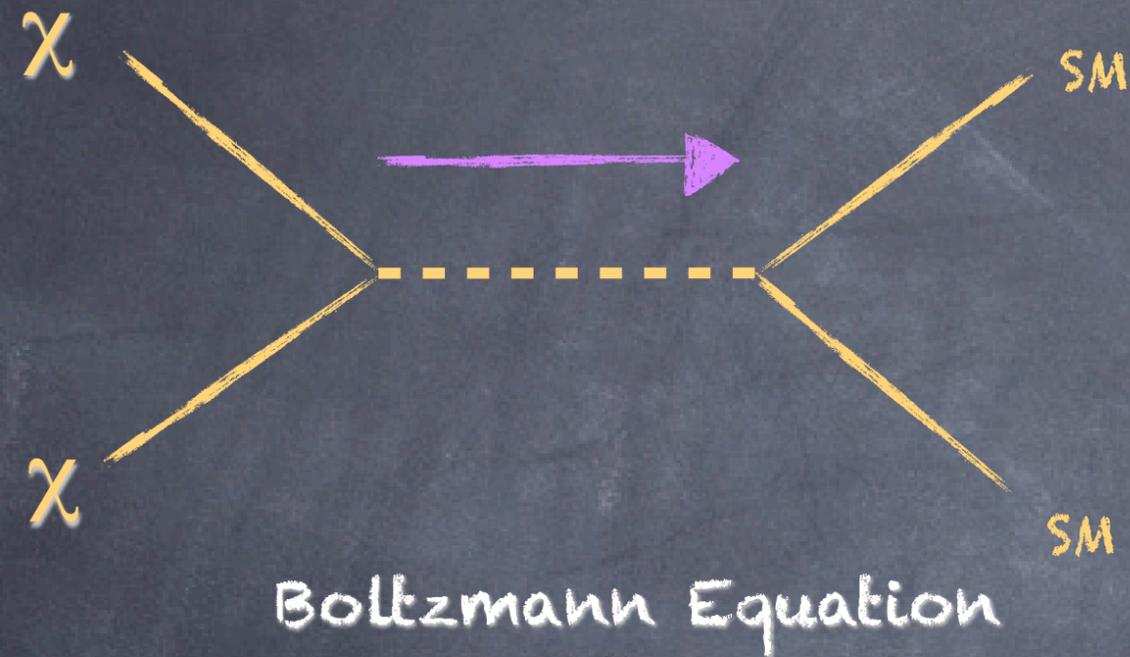


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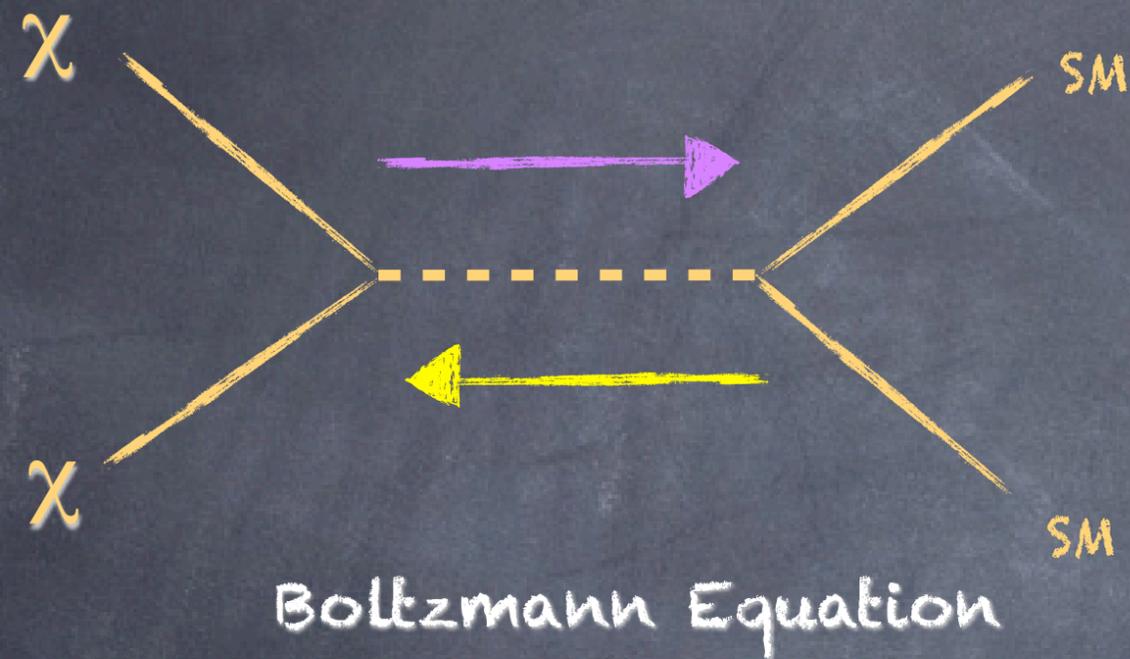


$$\frac{dn}{dt} = -3Hn - \langle \sigma v \rangle n^2$$

$$(H = \dot{R} / R)$$

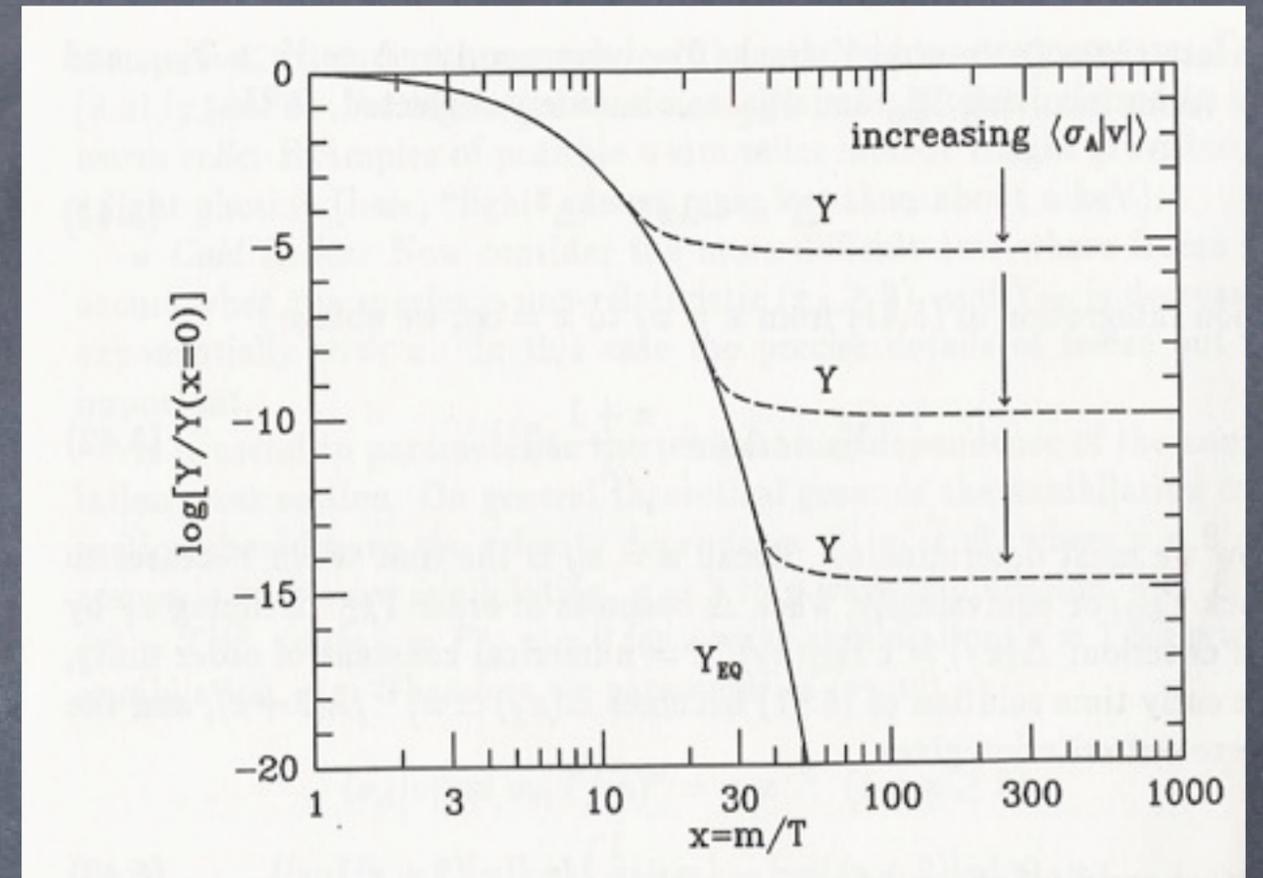
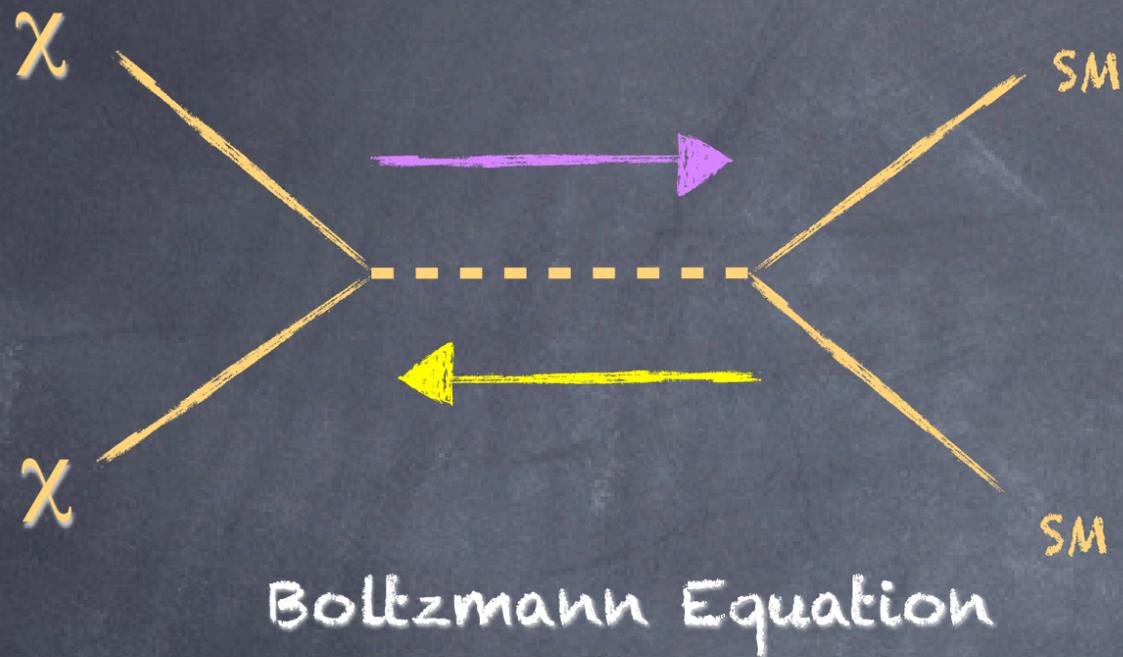
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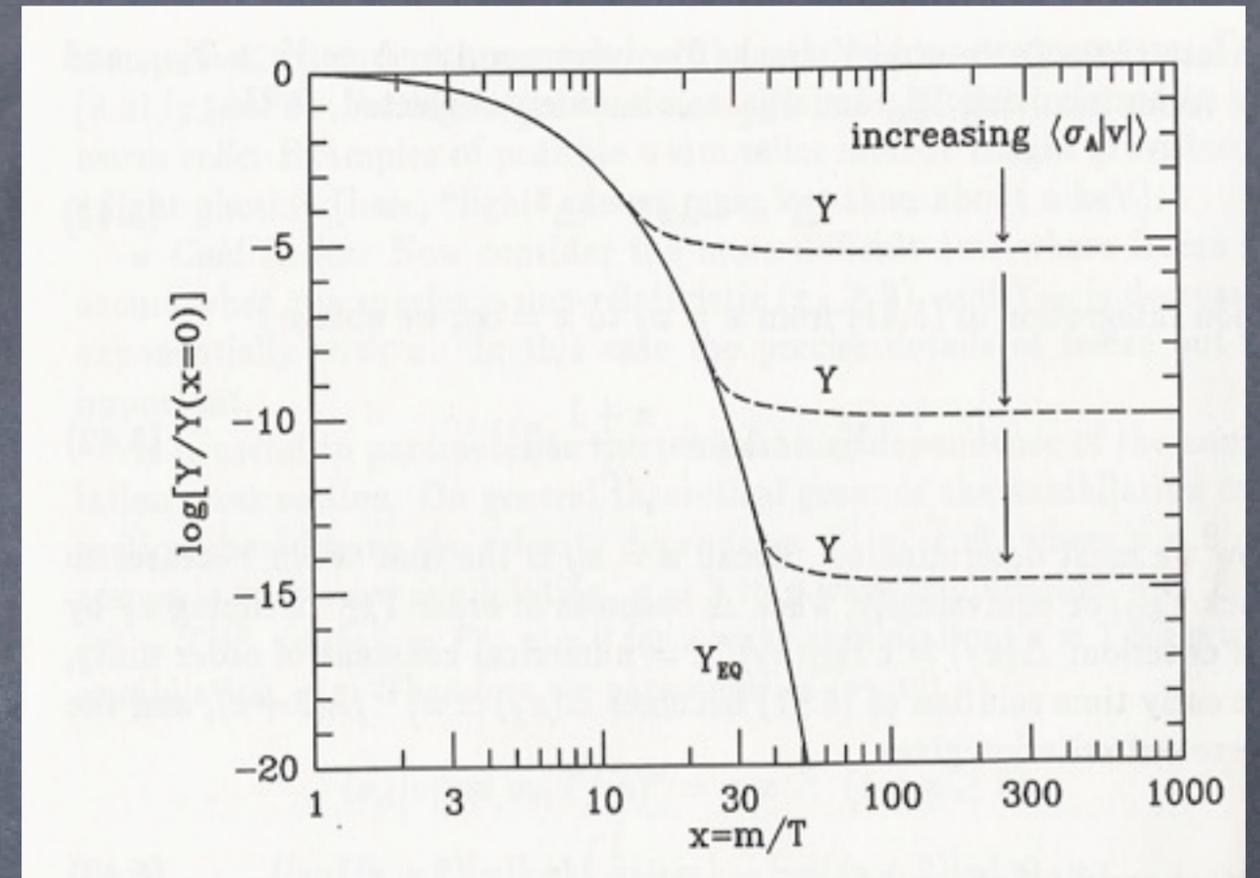
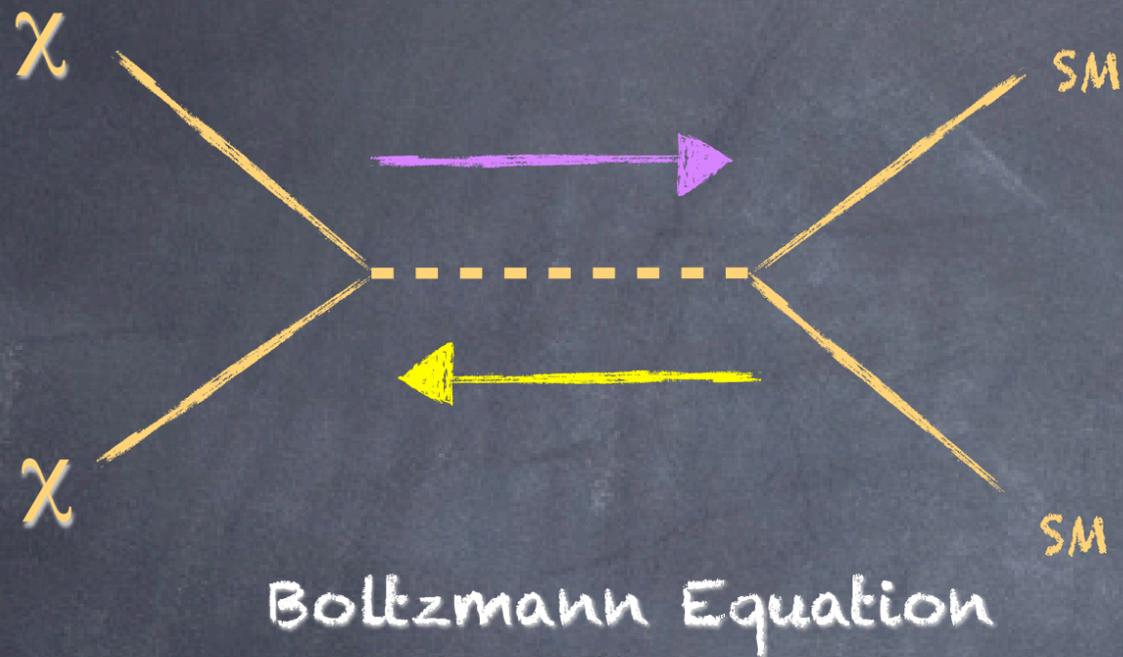
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$$\frac{dn}{dt} = -3Hn - \langle \sigma v \rangle [n^2 - n_{eq}^2] \quad (H = \dot{R}/R)$$

$$\Omega h^2 \sim \frac{3 \cdot 10^{-27} \text{ cm}^3 \text{ s}^{-1}}{\langle \sigma v \rangle} \sim 0.1,$$

decoupling at  $T \sim M_X/20$

# Astroparticle data..

[DD = Direct detection; ID $\gamma$  = Indirect gamma; ID $e^+$  = Indirect antimatter]

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- DAMA + DAMA/LIBRA : DD ; 5-20 GeV WIMP

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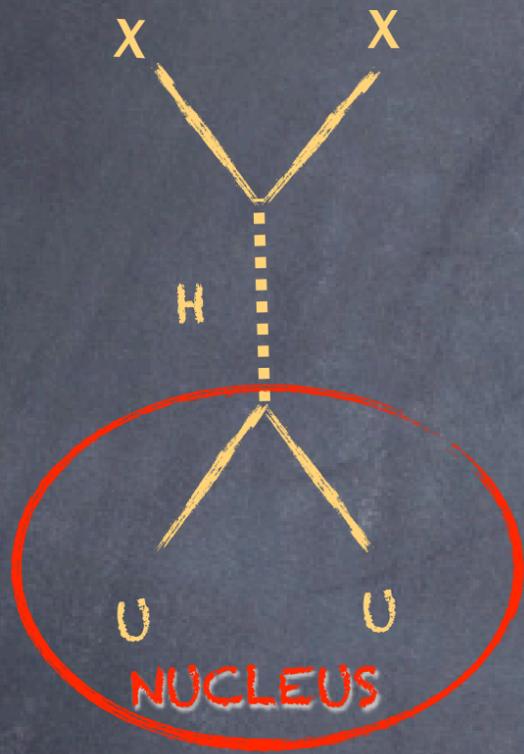
[DD = Direct detection; ID $\gamma$  = Indirect gamma; ID $e^+$  = Indirect antimatter]

- DAMA + DAMA/LIBRA : DD ; 5-20 GeV WIMP
- EGRET : ID $\gamma$  ; 40-100 GeV WIMP
- INTEGRAL : ID $e^+$  ; 2-5 MeV DM
- HEAT : ID $e^+$  ; 100 GeV DM
- HESS : ID $\gamma$  ;  $> 15$  TeV DM
- PAMELA + FERMI : ID $e^+$  ; 100 GeV or  $> 5$  TeV DM
- CDMS : DD ; 10-100 GeV WIMP
- CoGENT : DD ; 6-15 GeV WIMP
- CRESST : DD ; 10-20 GeV WIMP

# Direct detection : principle

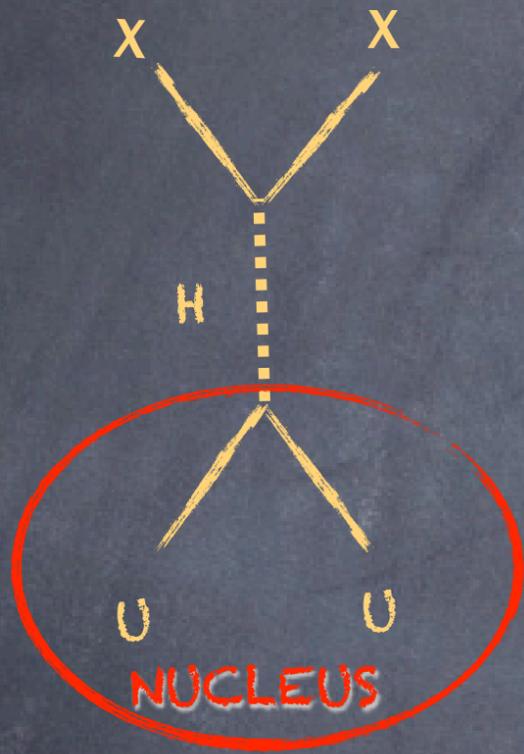
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# Direct detection : principle



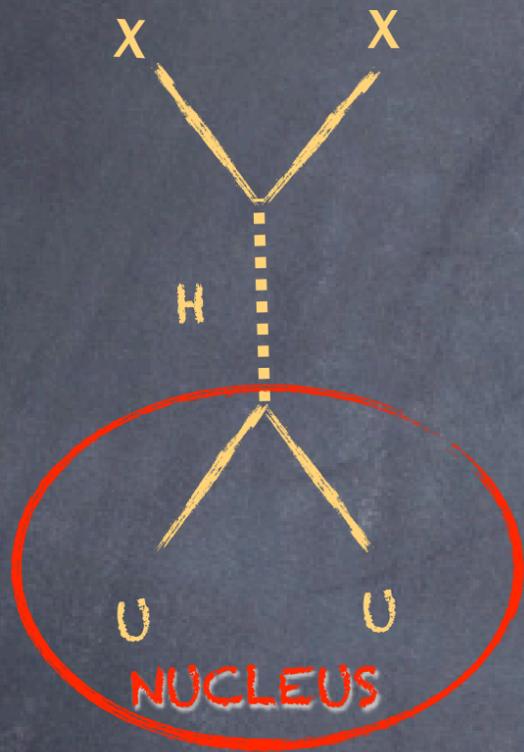
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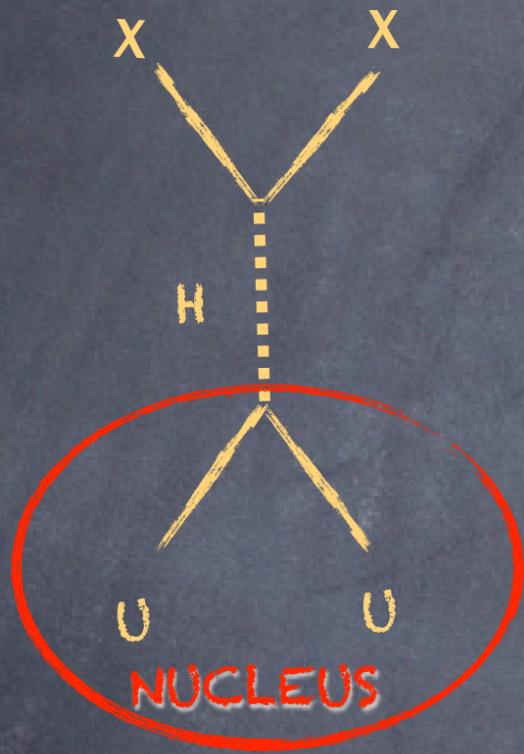
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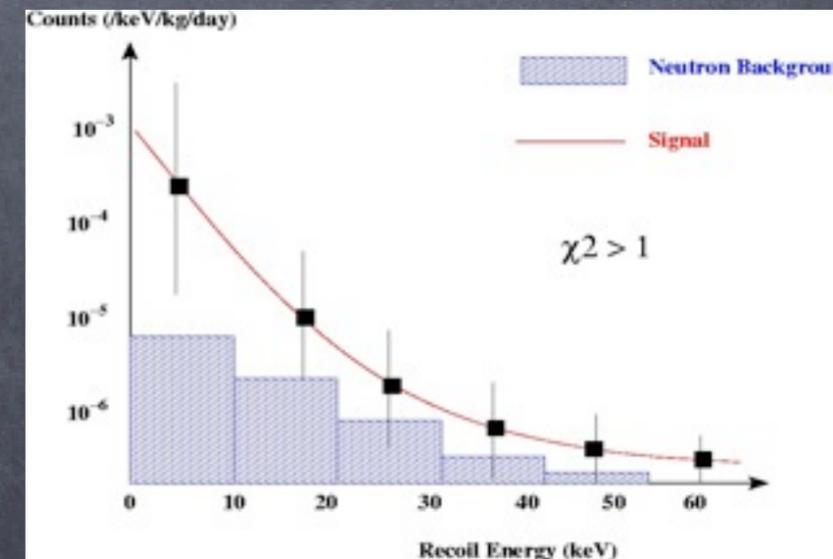
$$\frac{dN}{dE_r} = \frac{\sigma^{SI} \rho}{2m_r^2 m_{dm}} F(E_r)^2 \int_{v_{min}}^{\infty} \frac{f(v)}{v} dv$$

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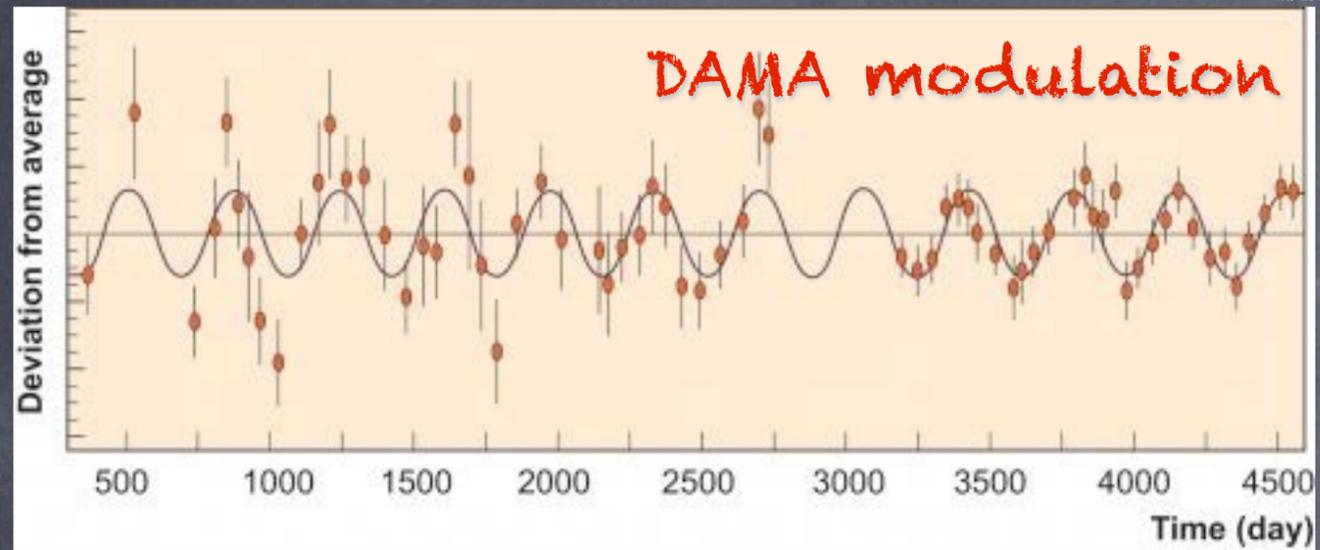


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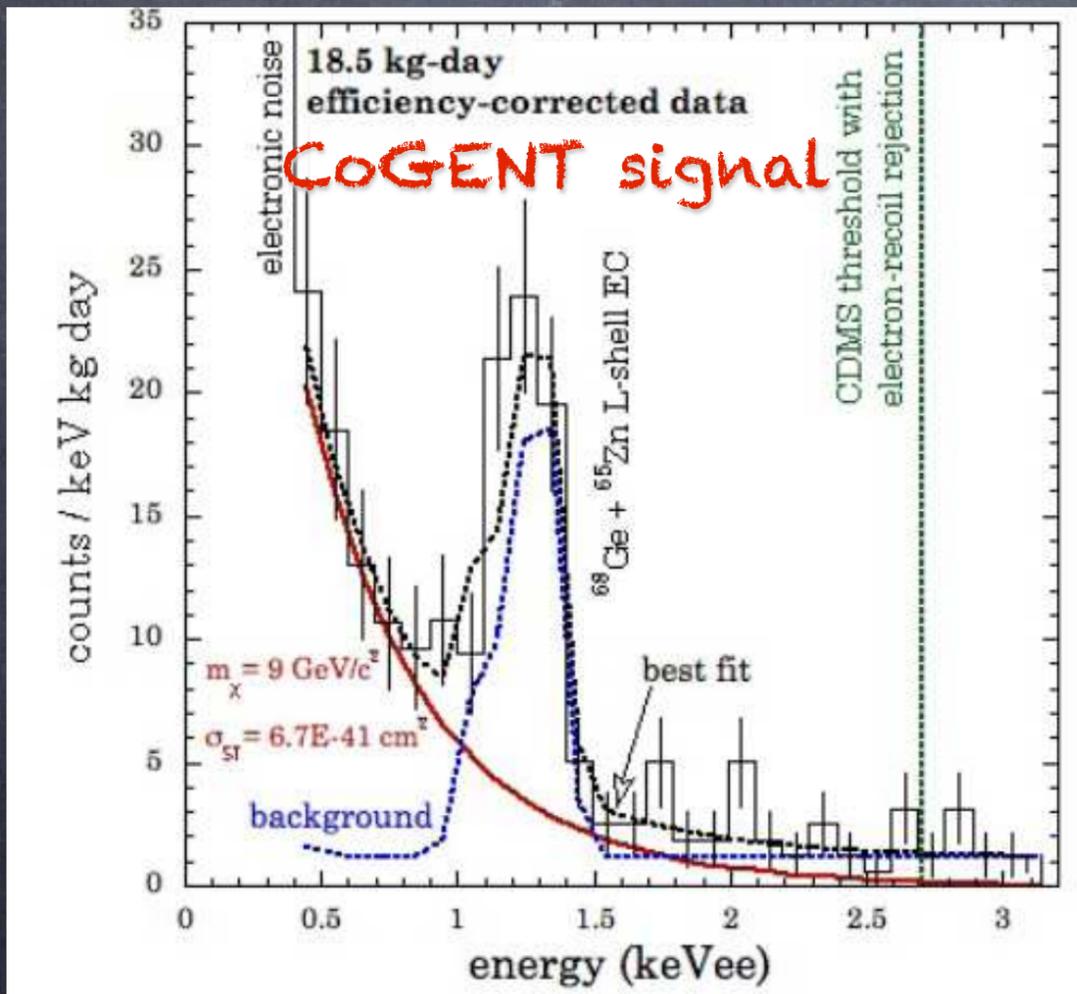
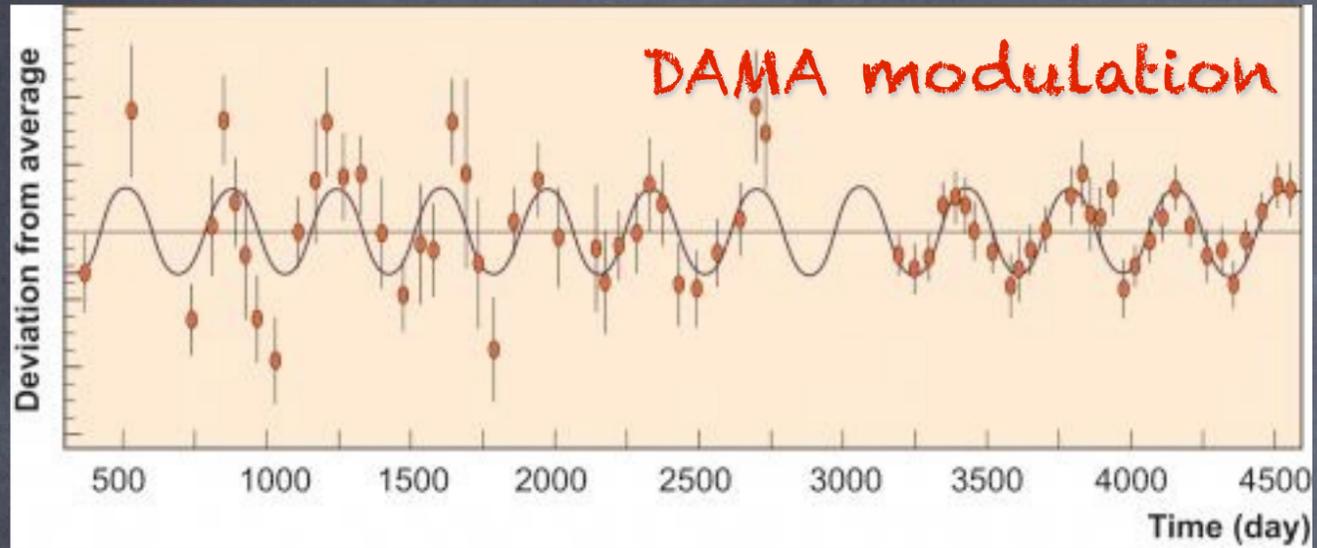


Direct detection : signals

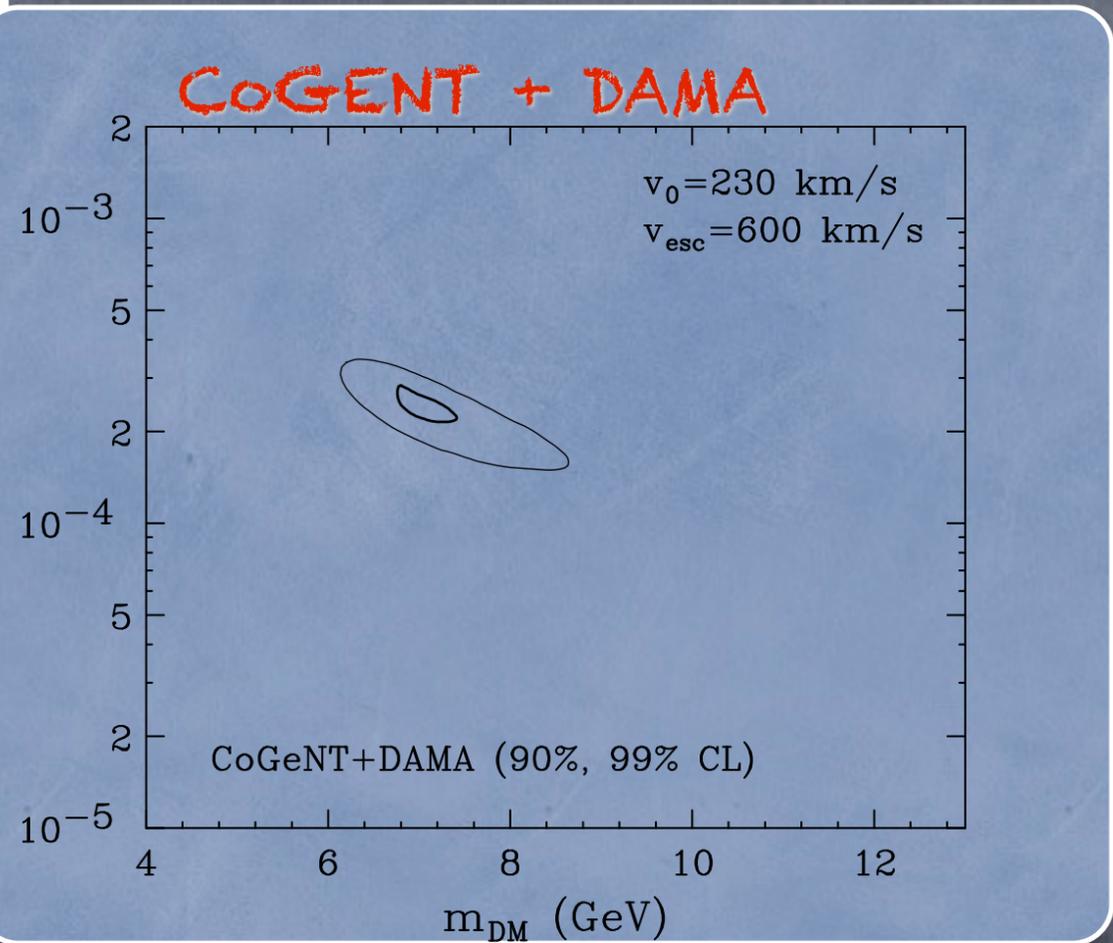
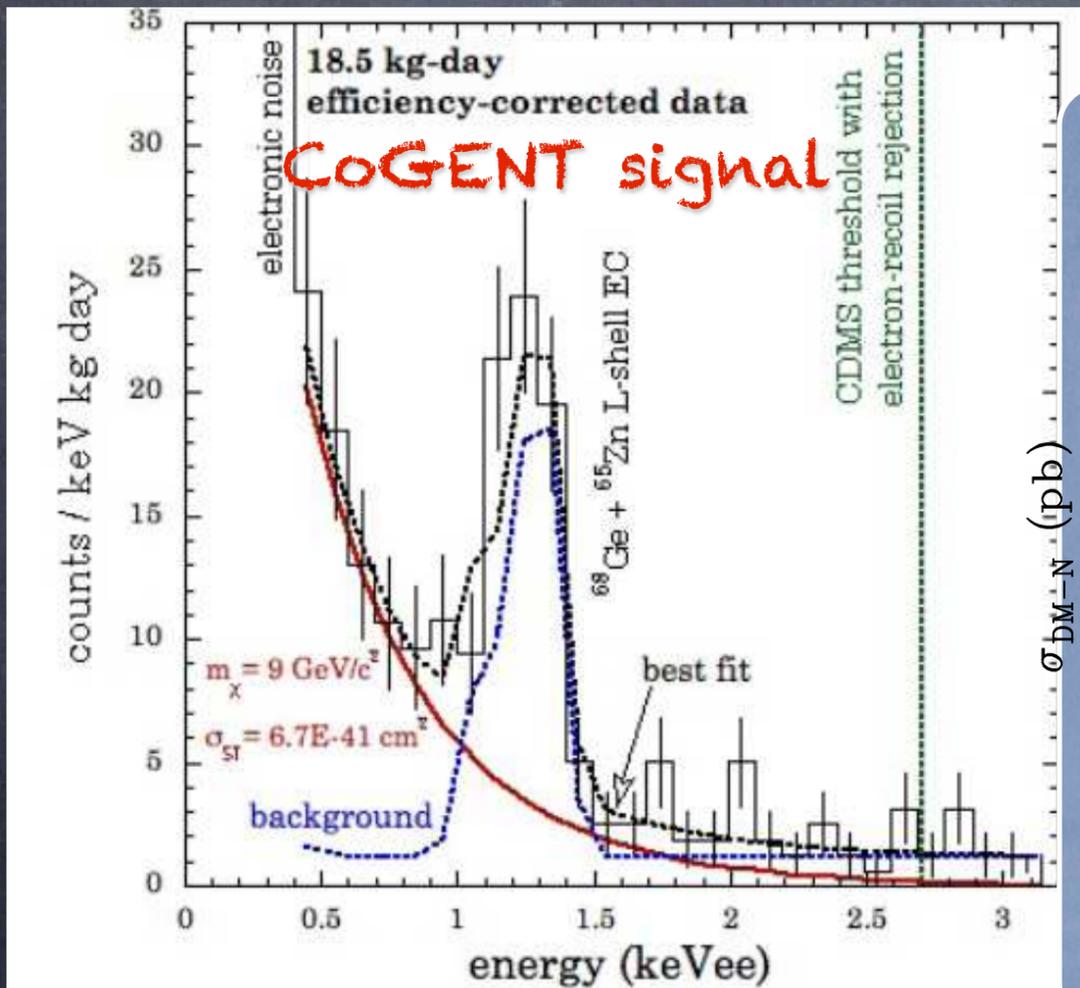
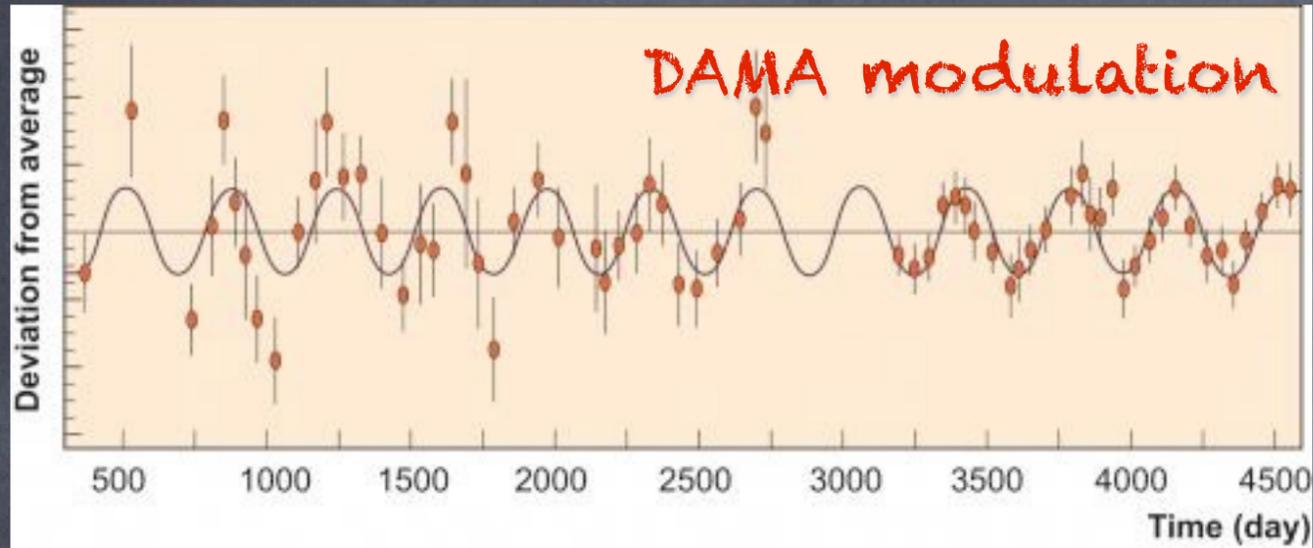
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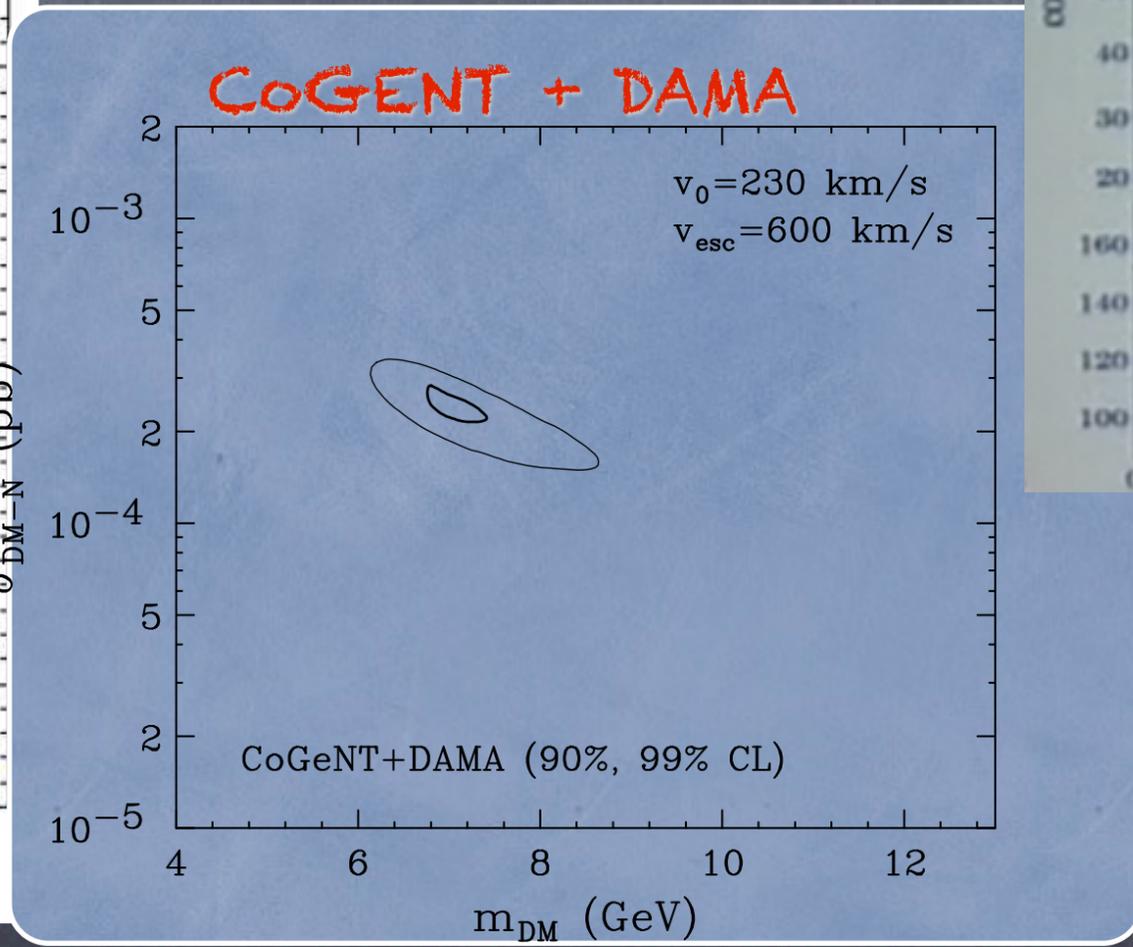
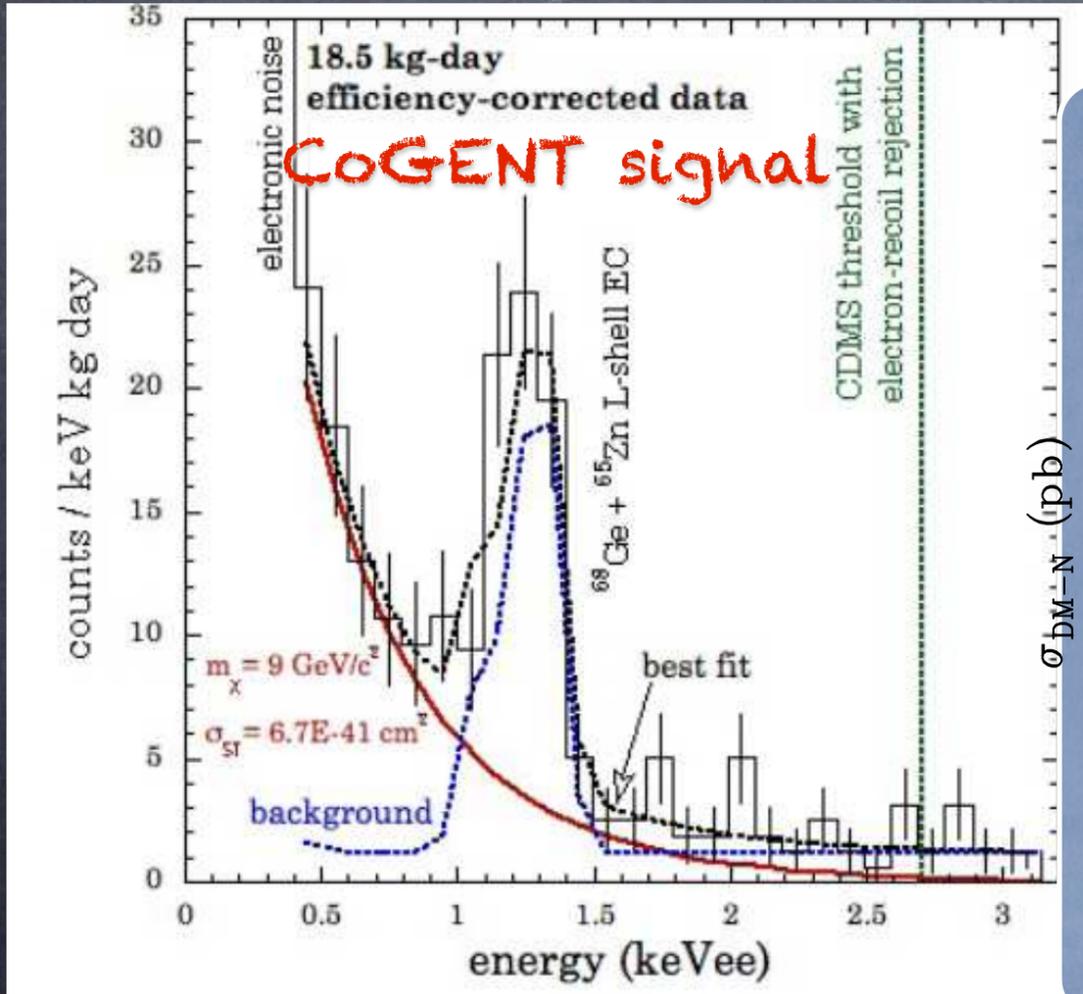
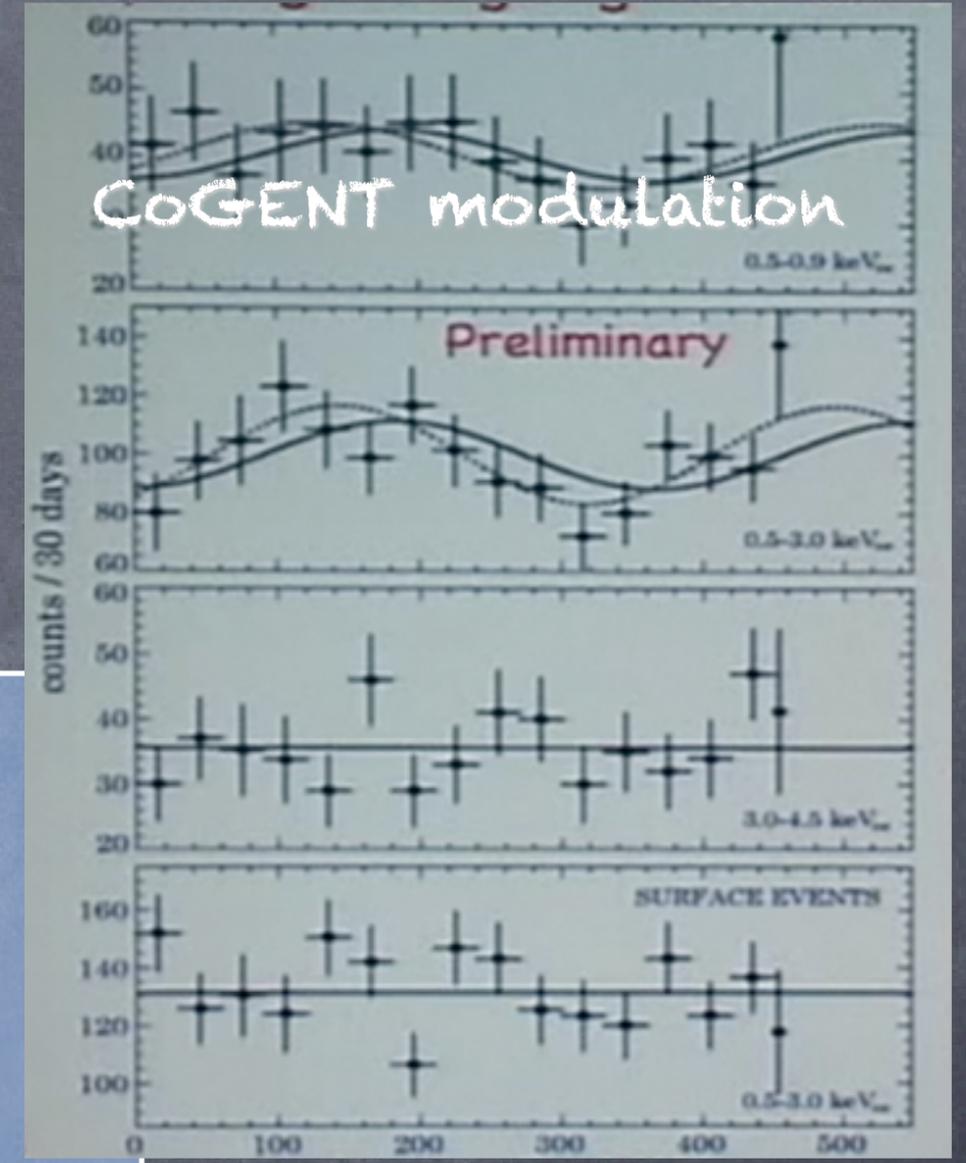
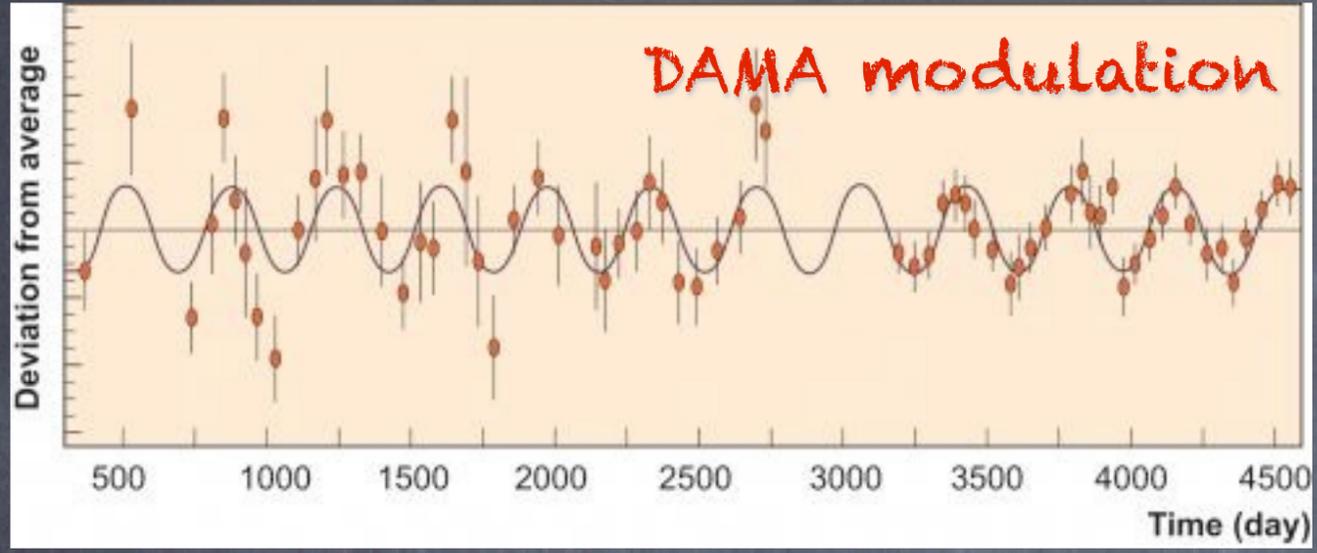
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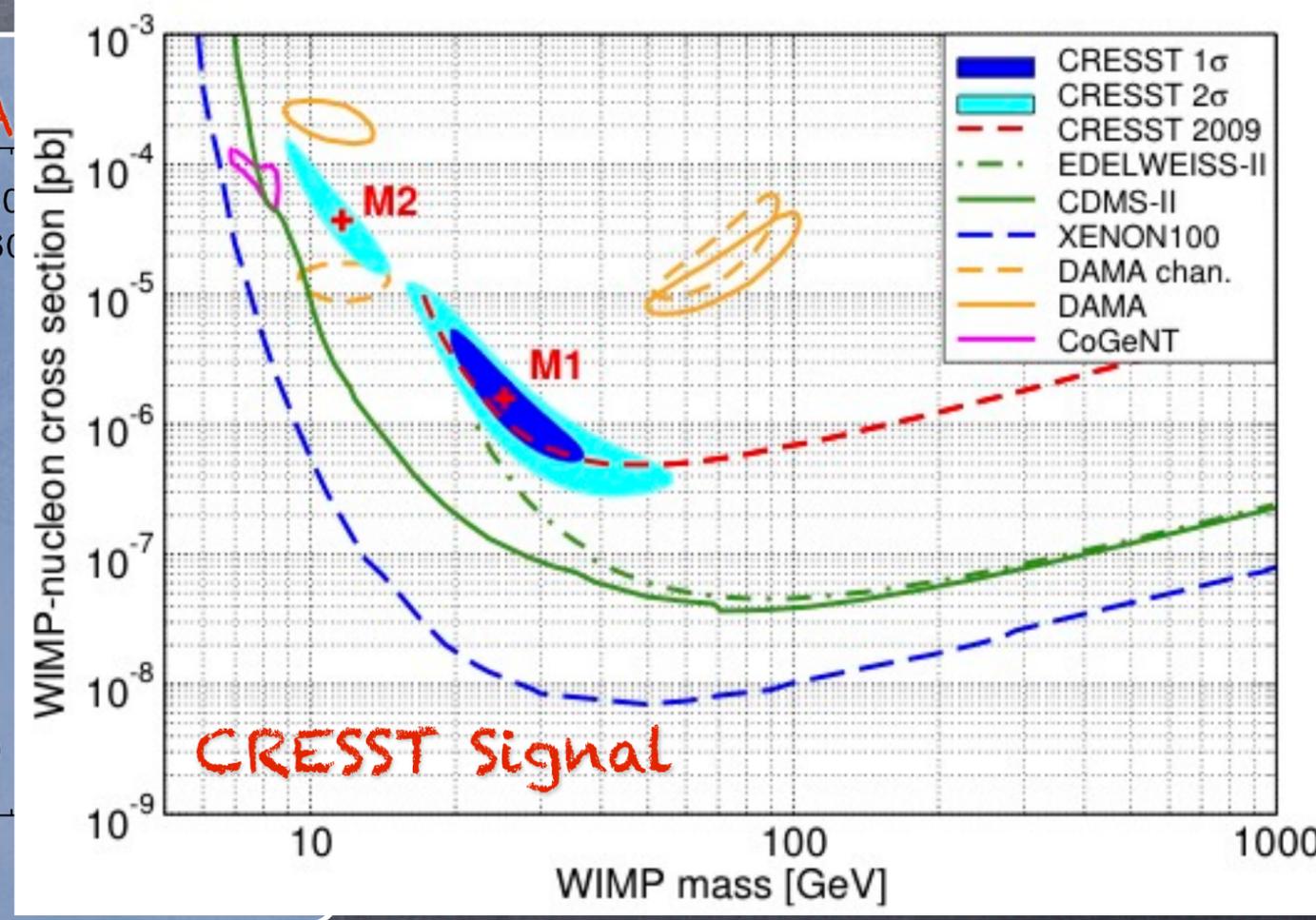
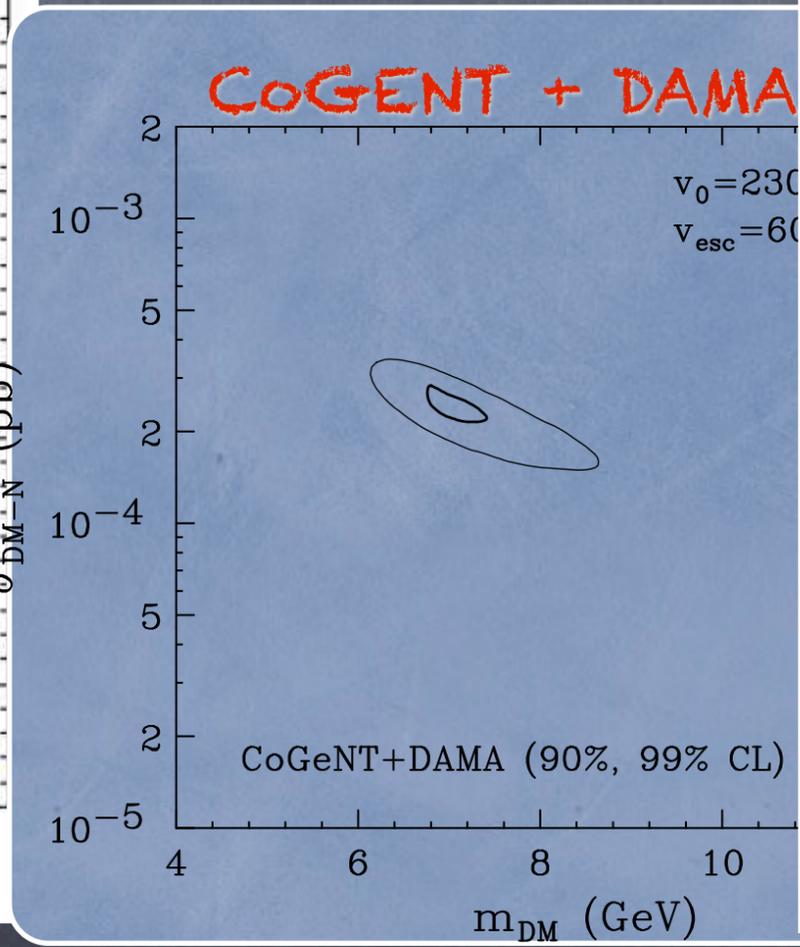
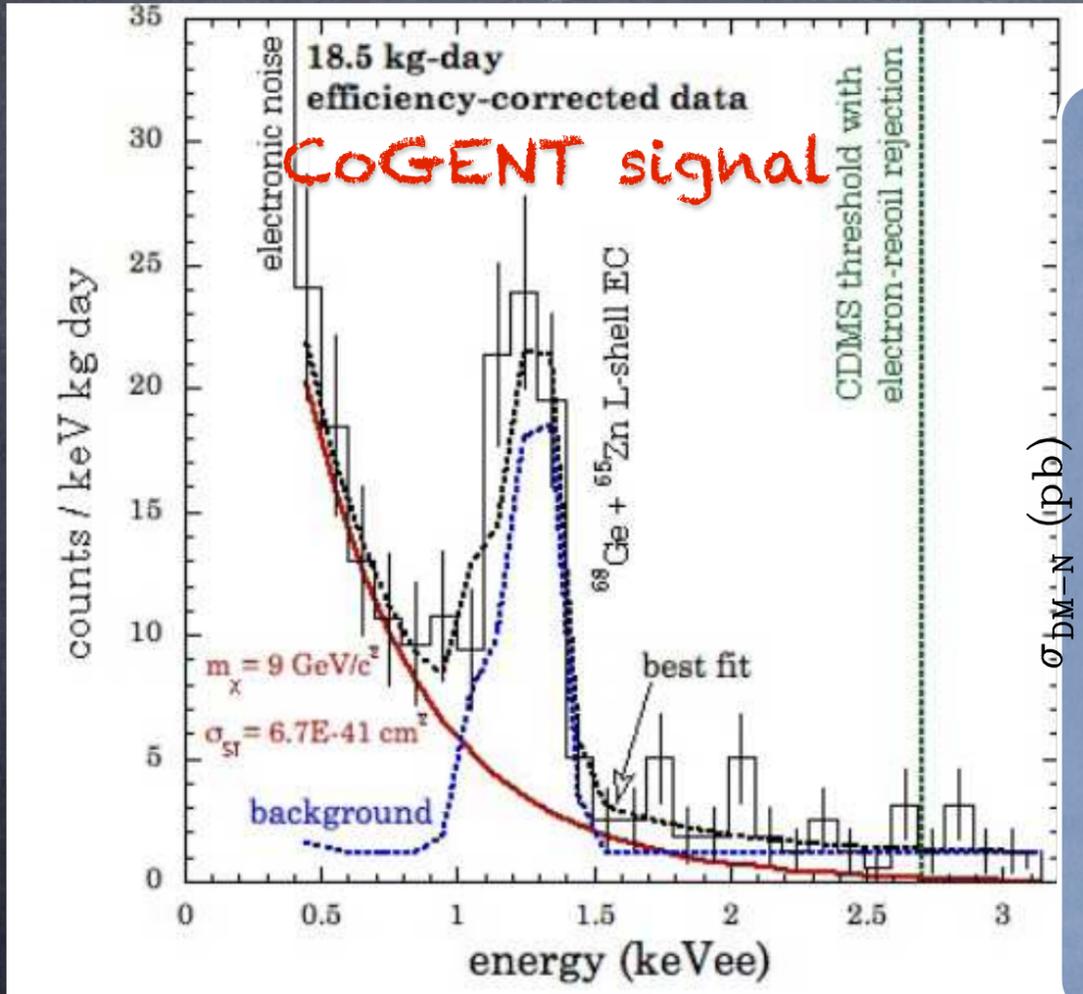
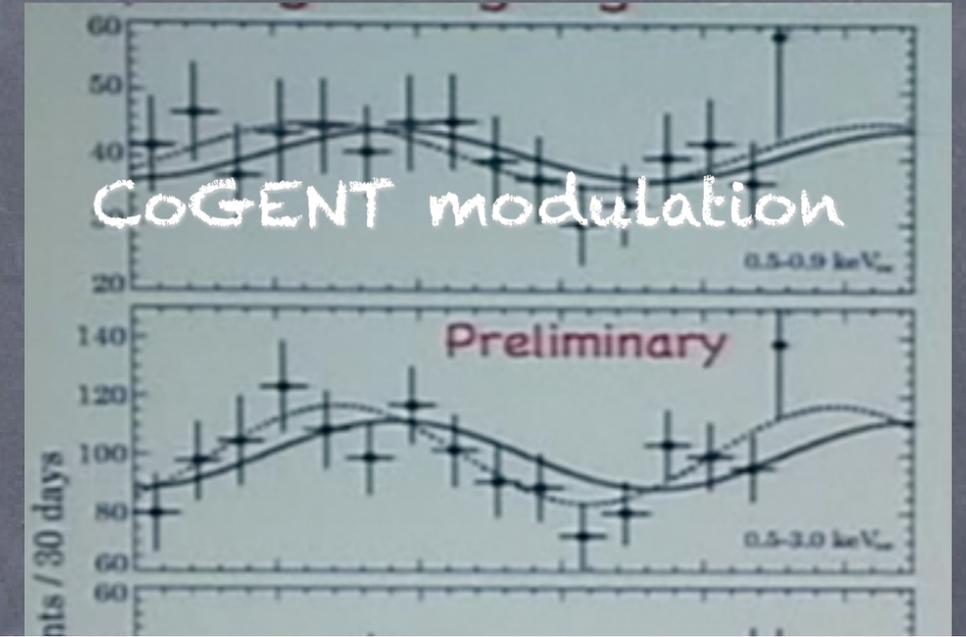
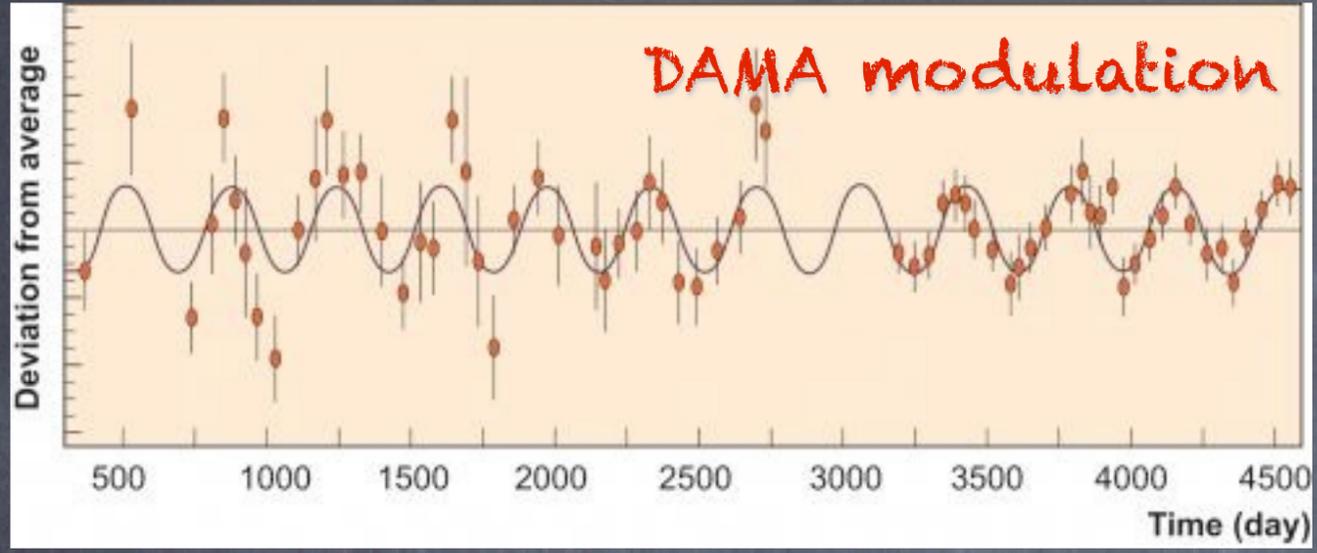
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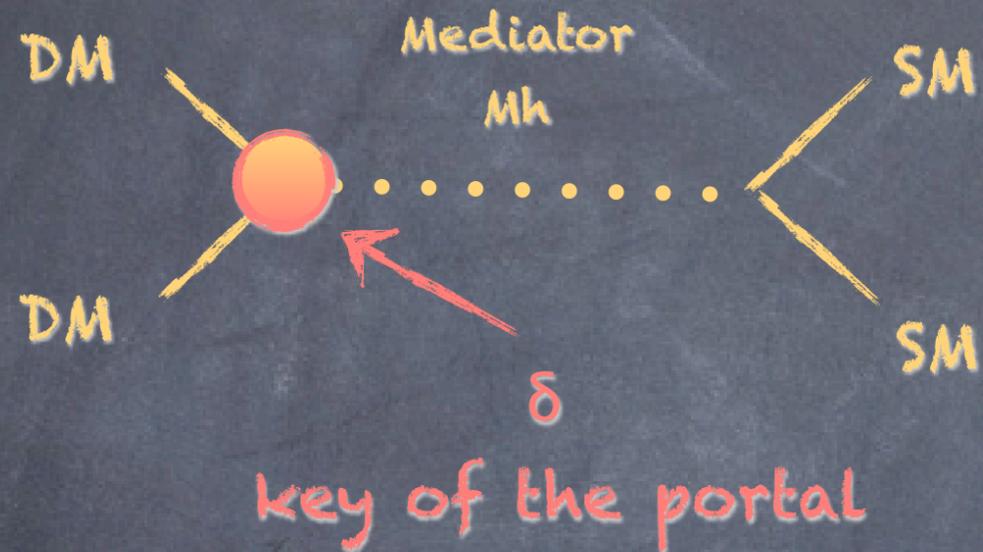
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# Constraints in «portal like» models

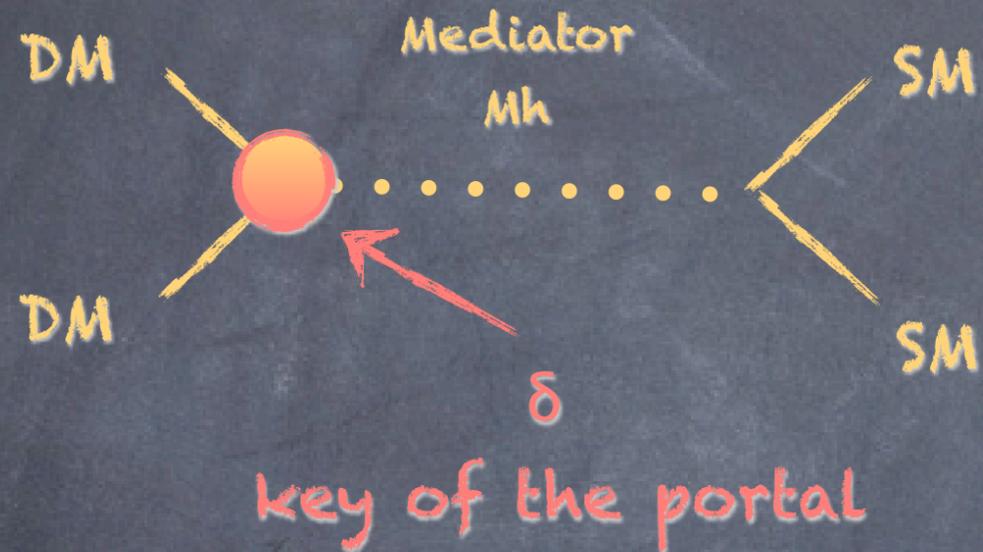
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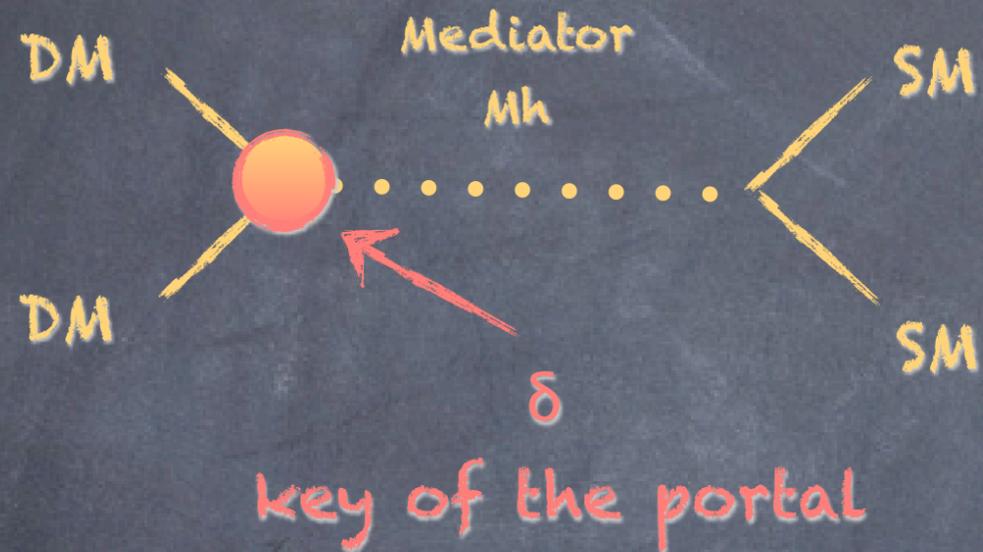
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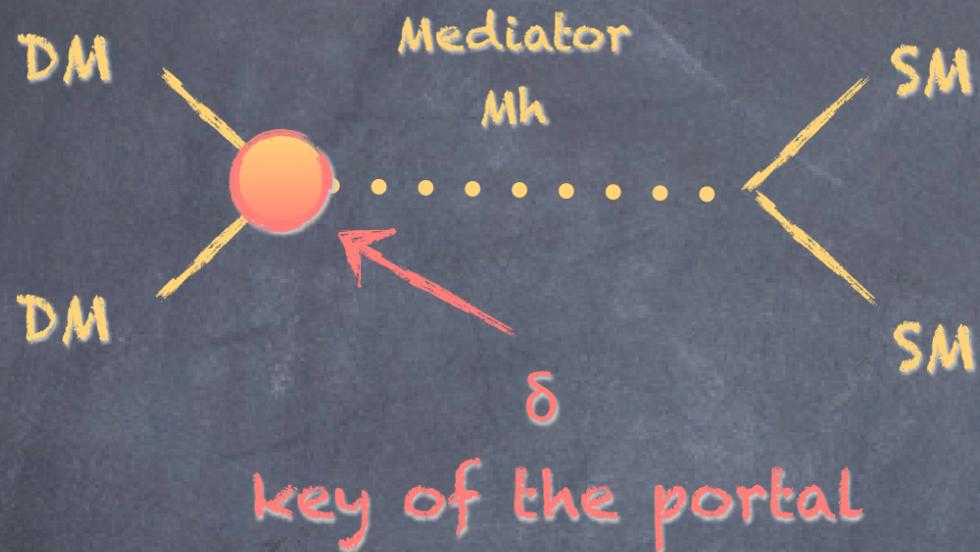
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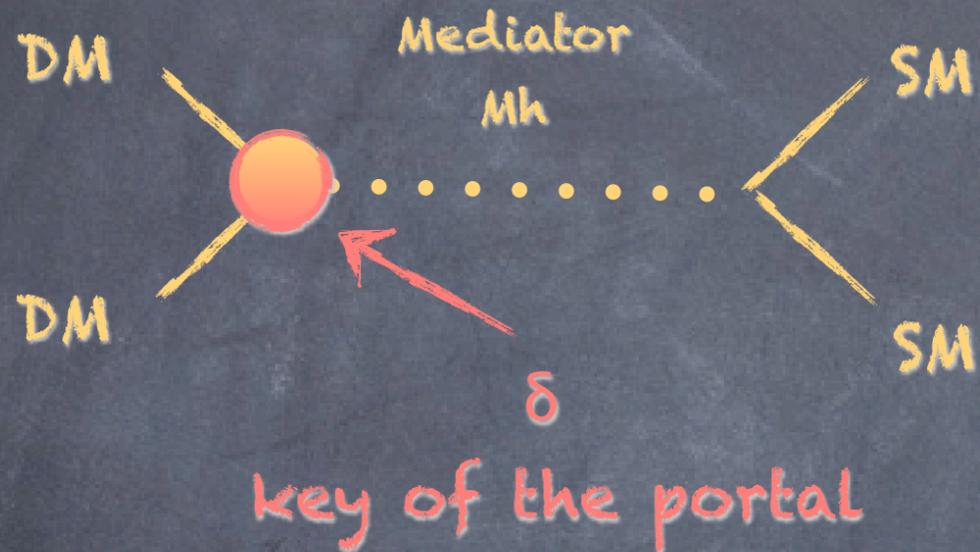


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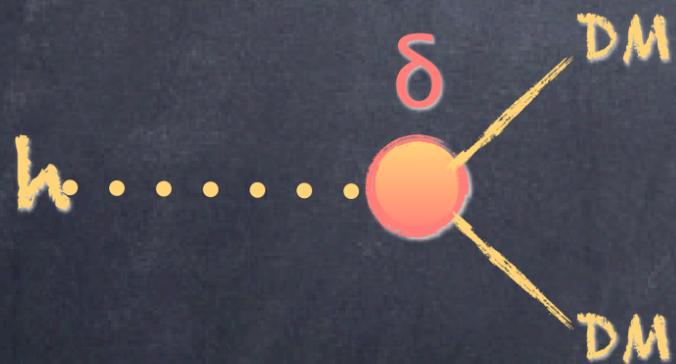
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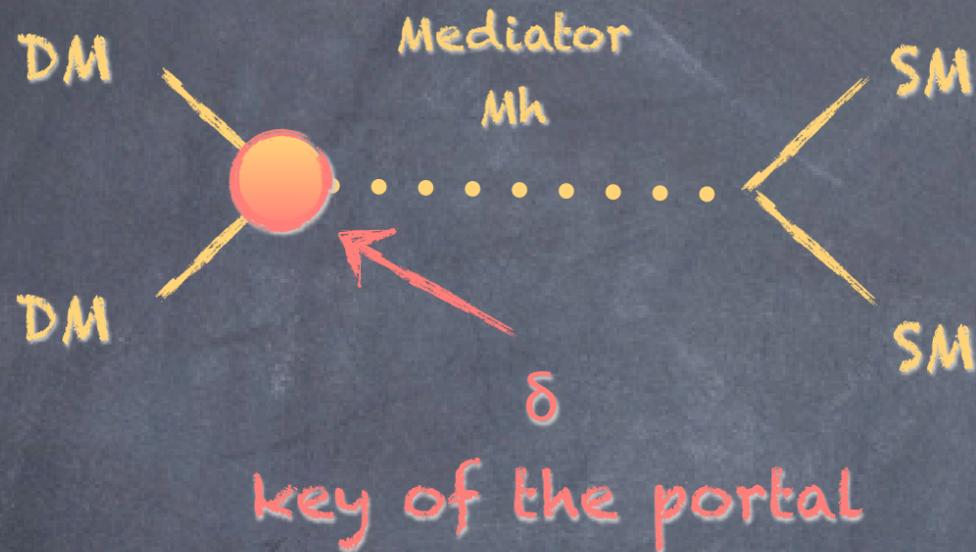
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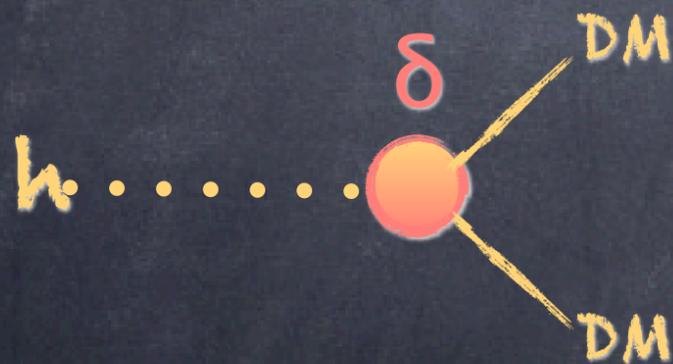
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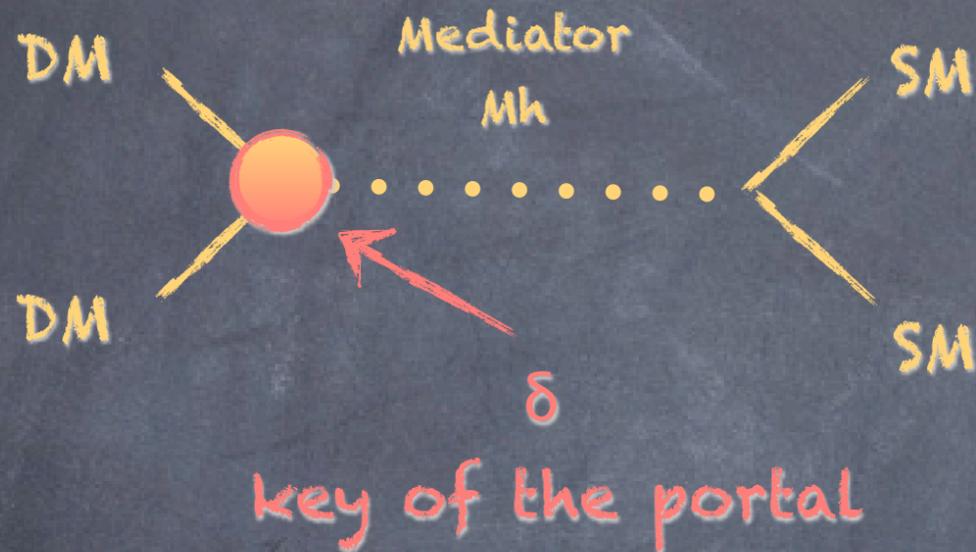
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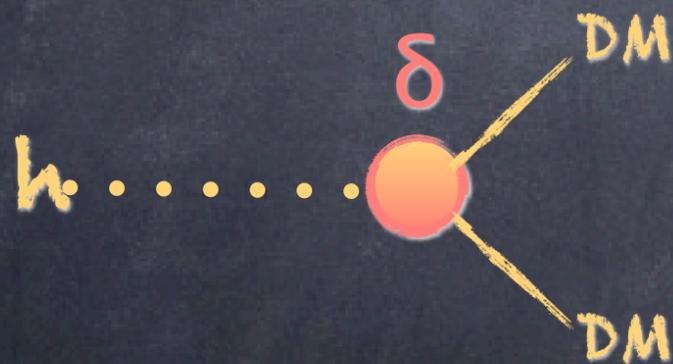
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XENON

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LHC/ILC

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Gauge extension: Extra  $U_D(1)$

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$$SU(3) * SU(2) * U(1) * U_D(1)$$

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$$\begin{array}{cccc} \text{SU}(3) * \text{SU}(2) * \text{U}(1) * \text{U}_D(1) \\ g_\mu & W_\mu & Y_\mu & X_\mu \end{array}$$

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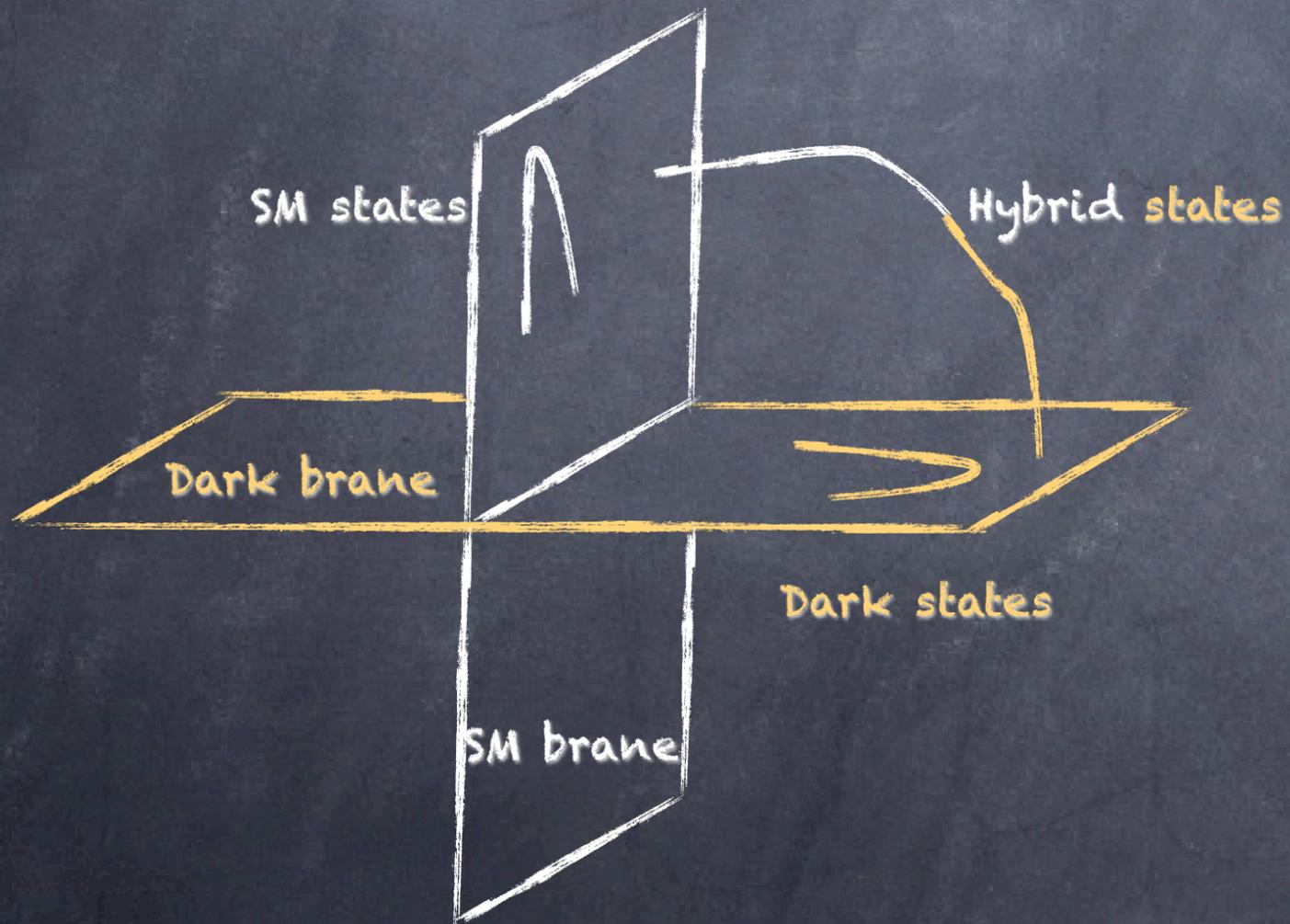
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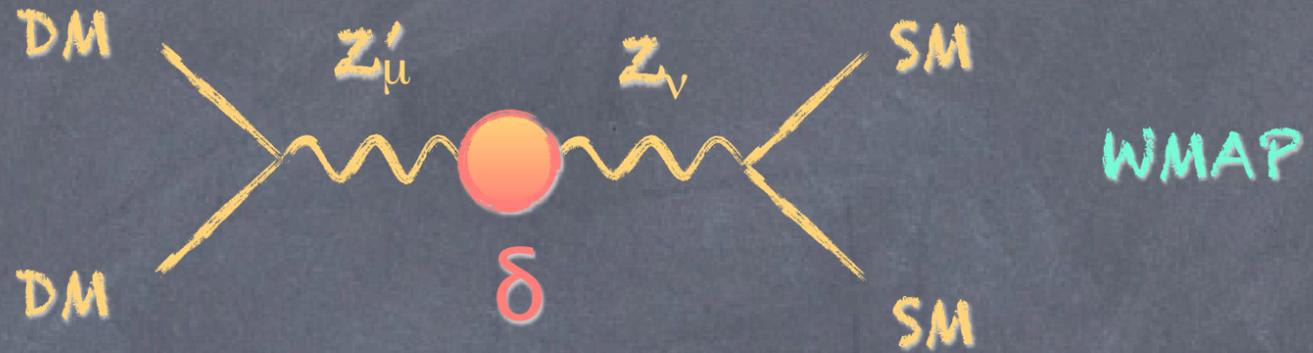
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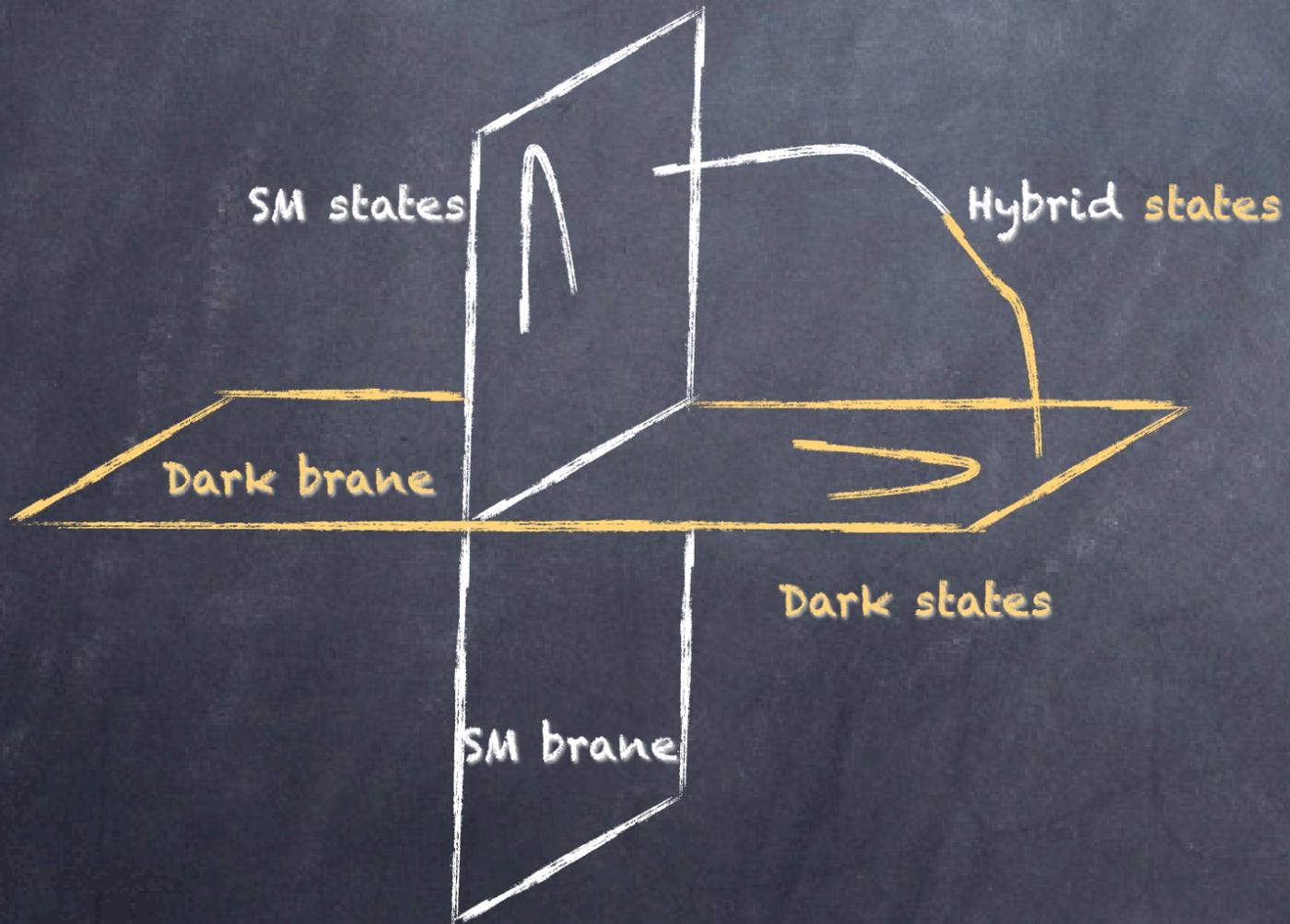
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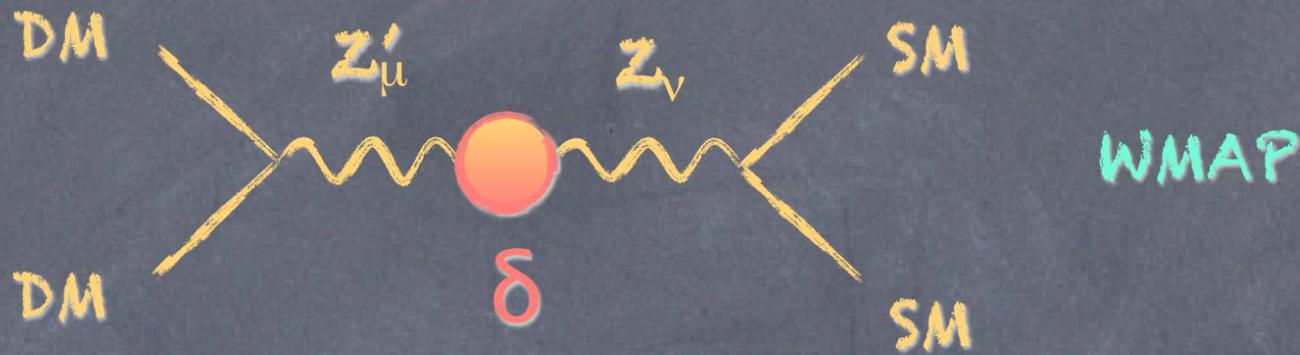
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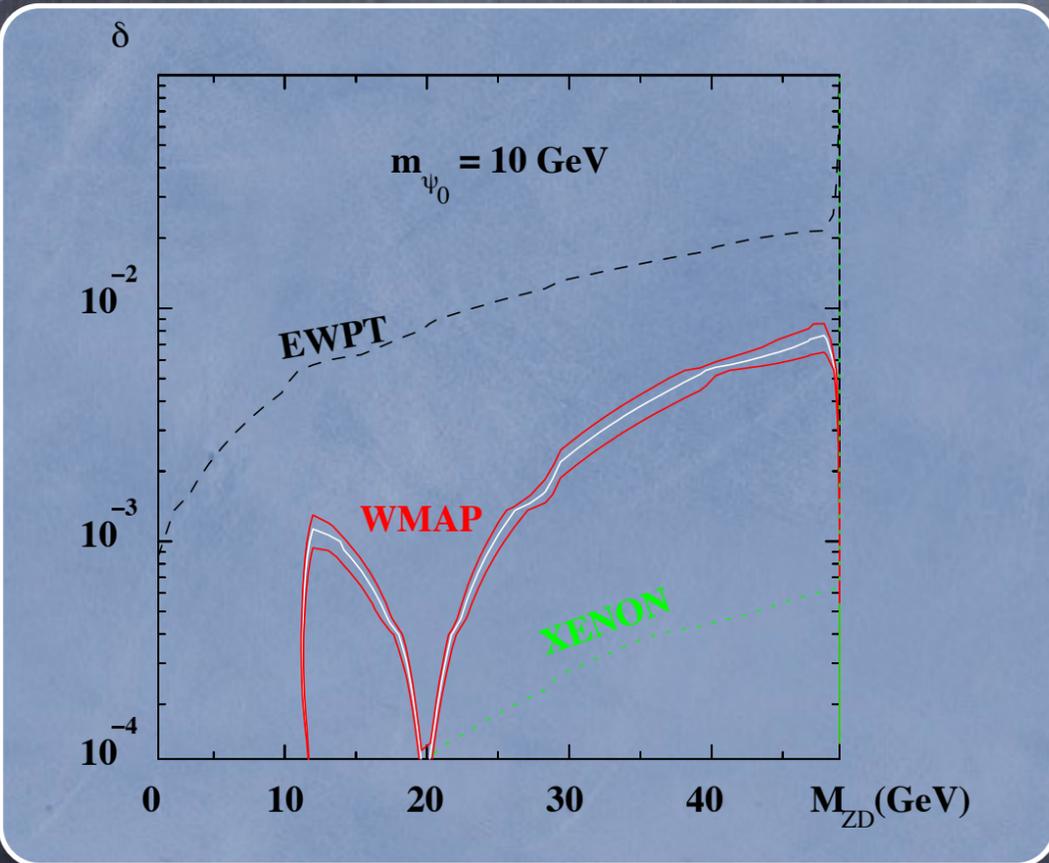
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Hybrid states

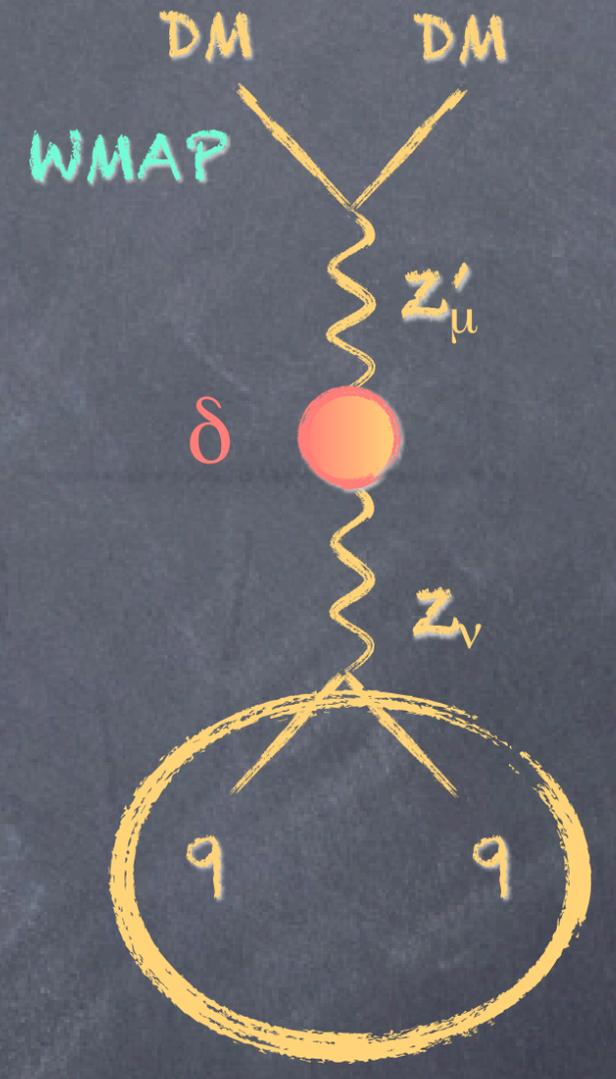
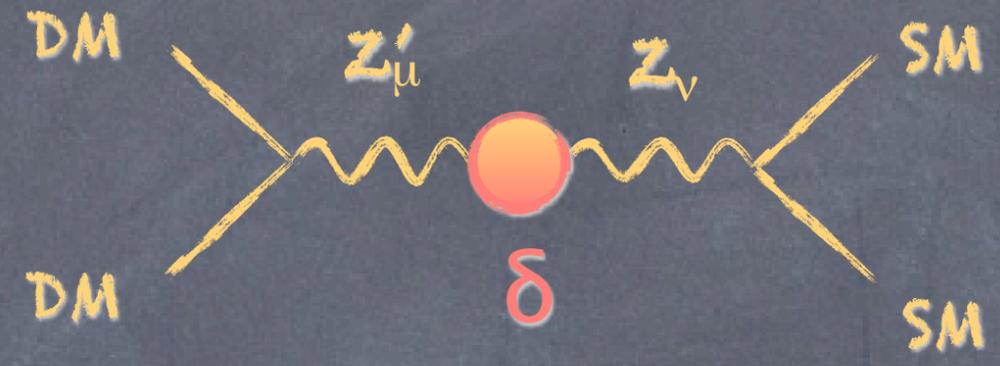
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$M_{DM} = 10 \text{ GeV}$  /  $M_X$

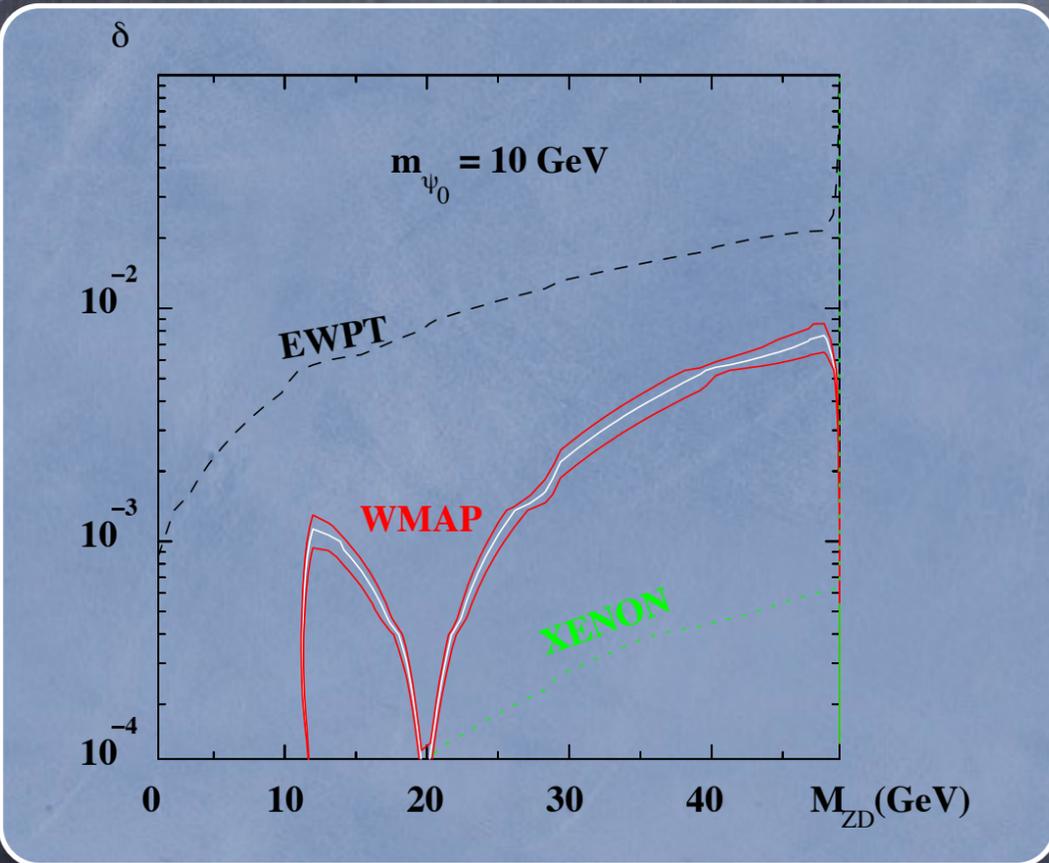
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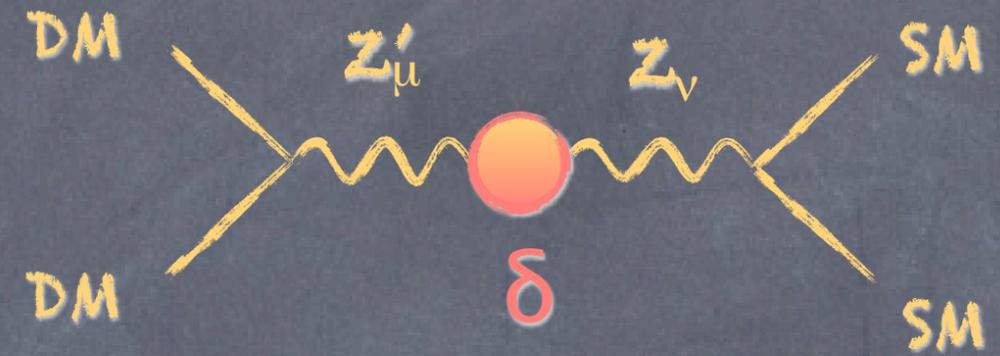
CoGENT

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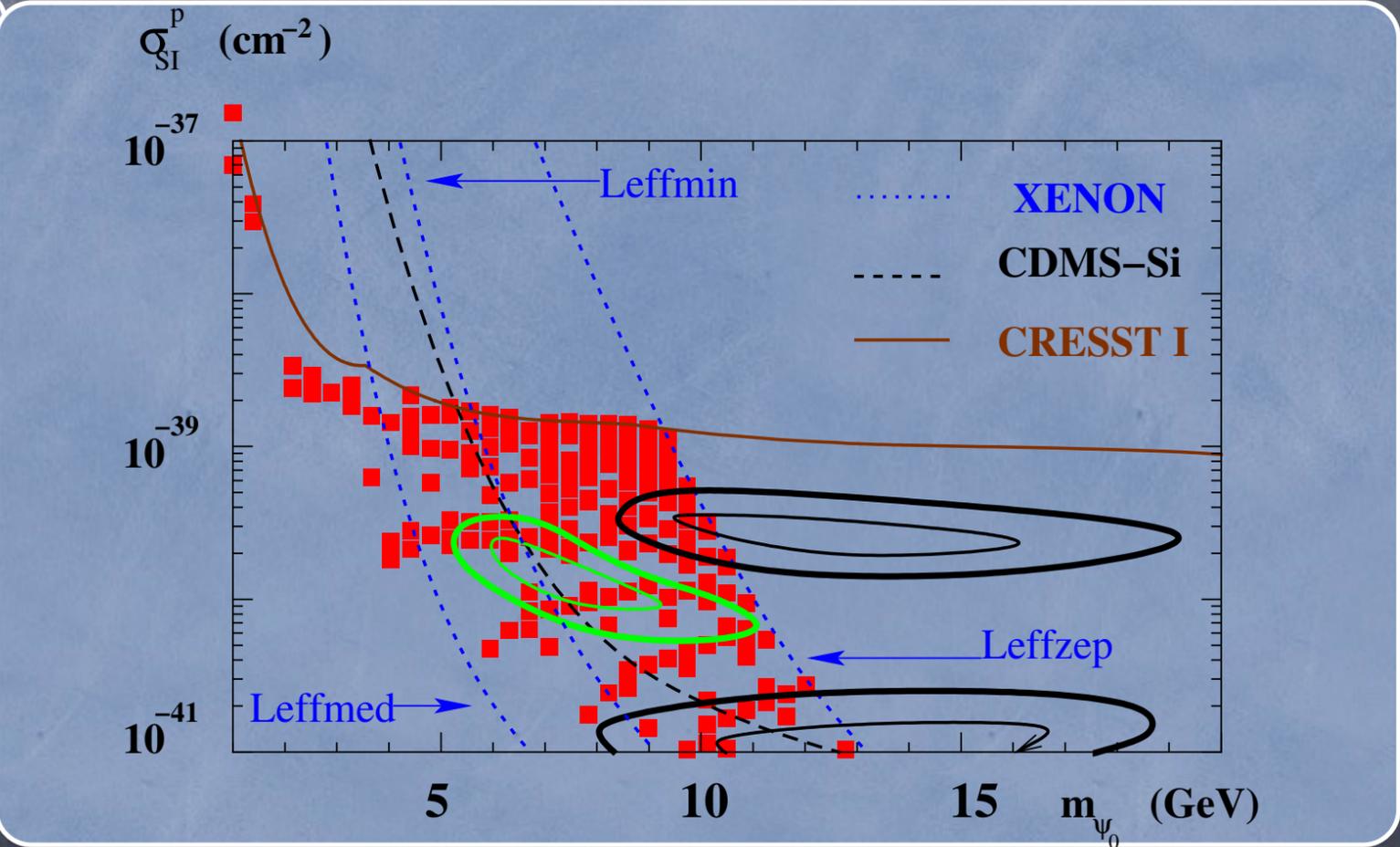
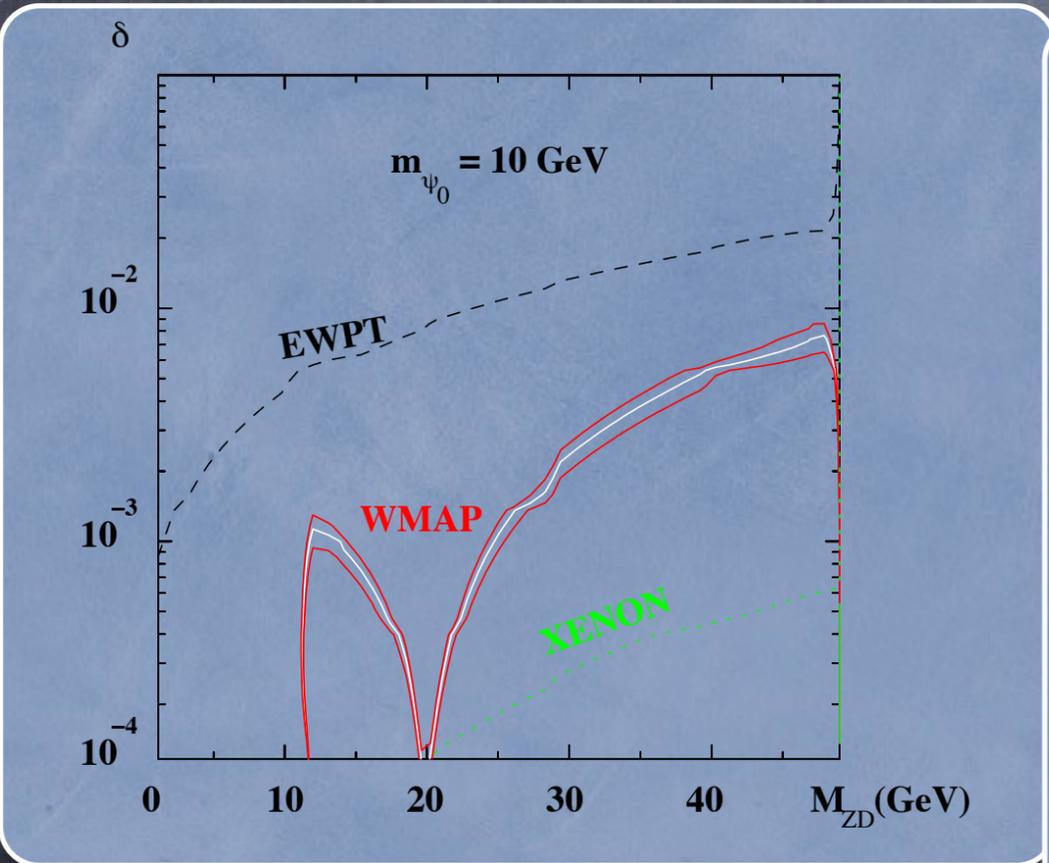
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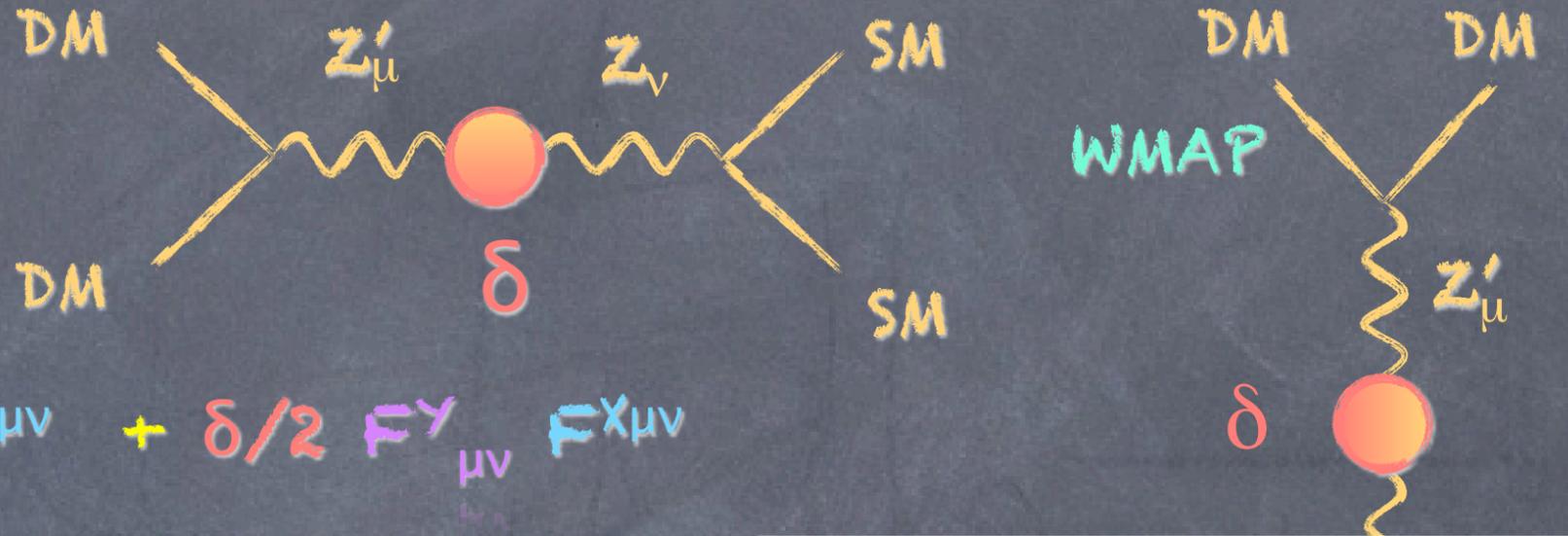
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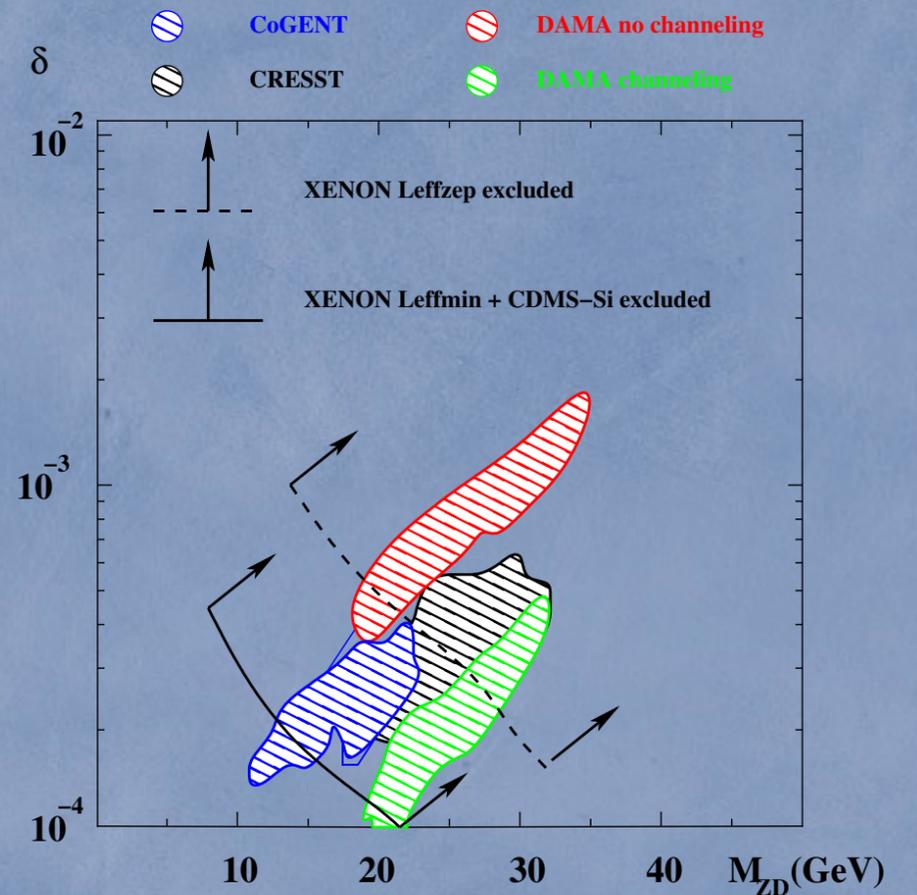
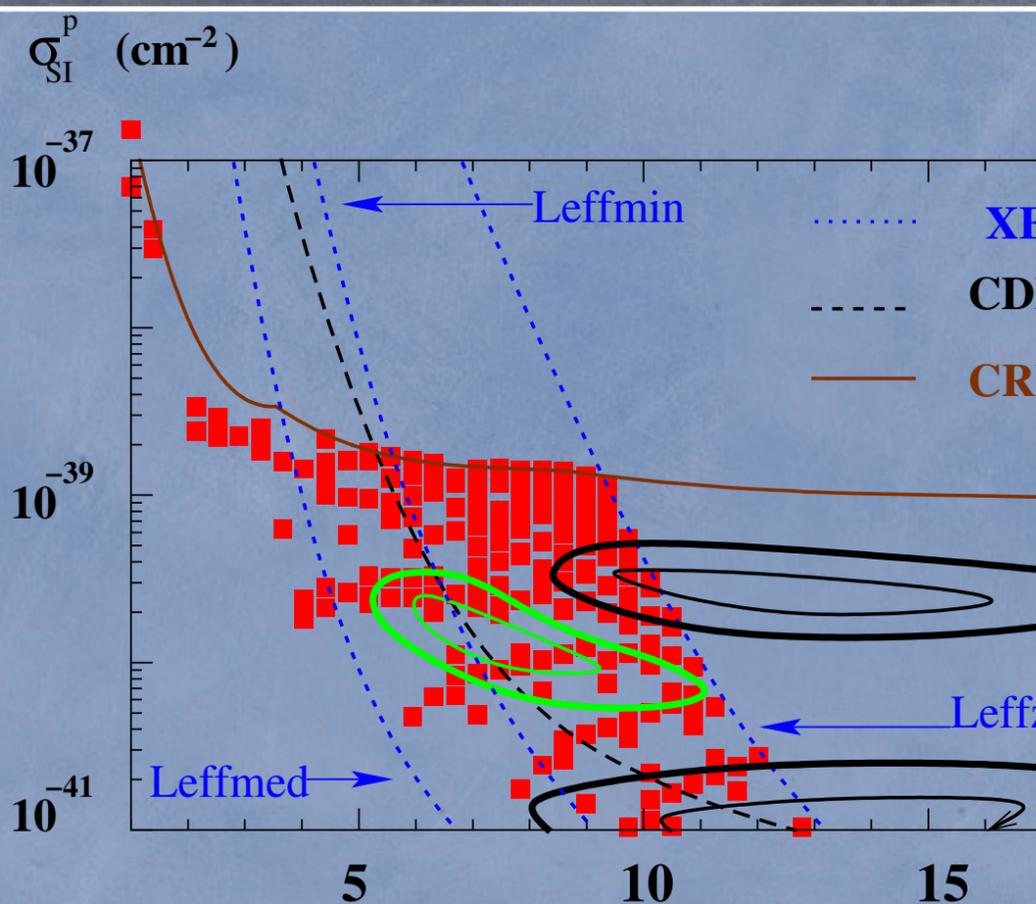
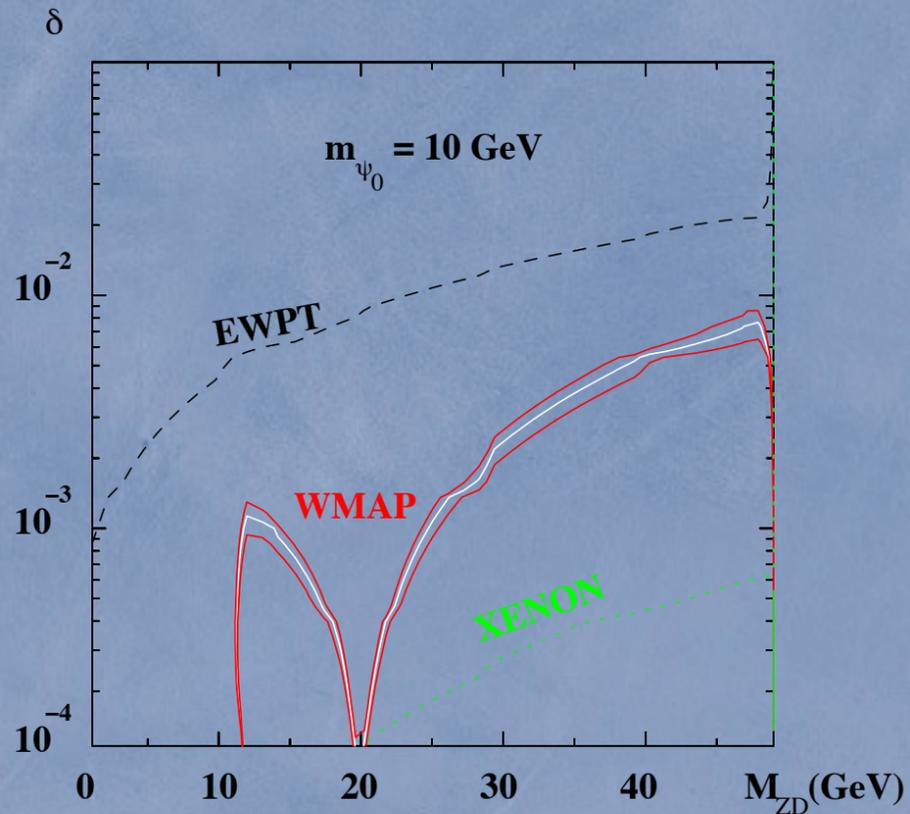
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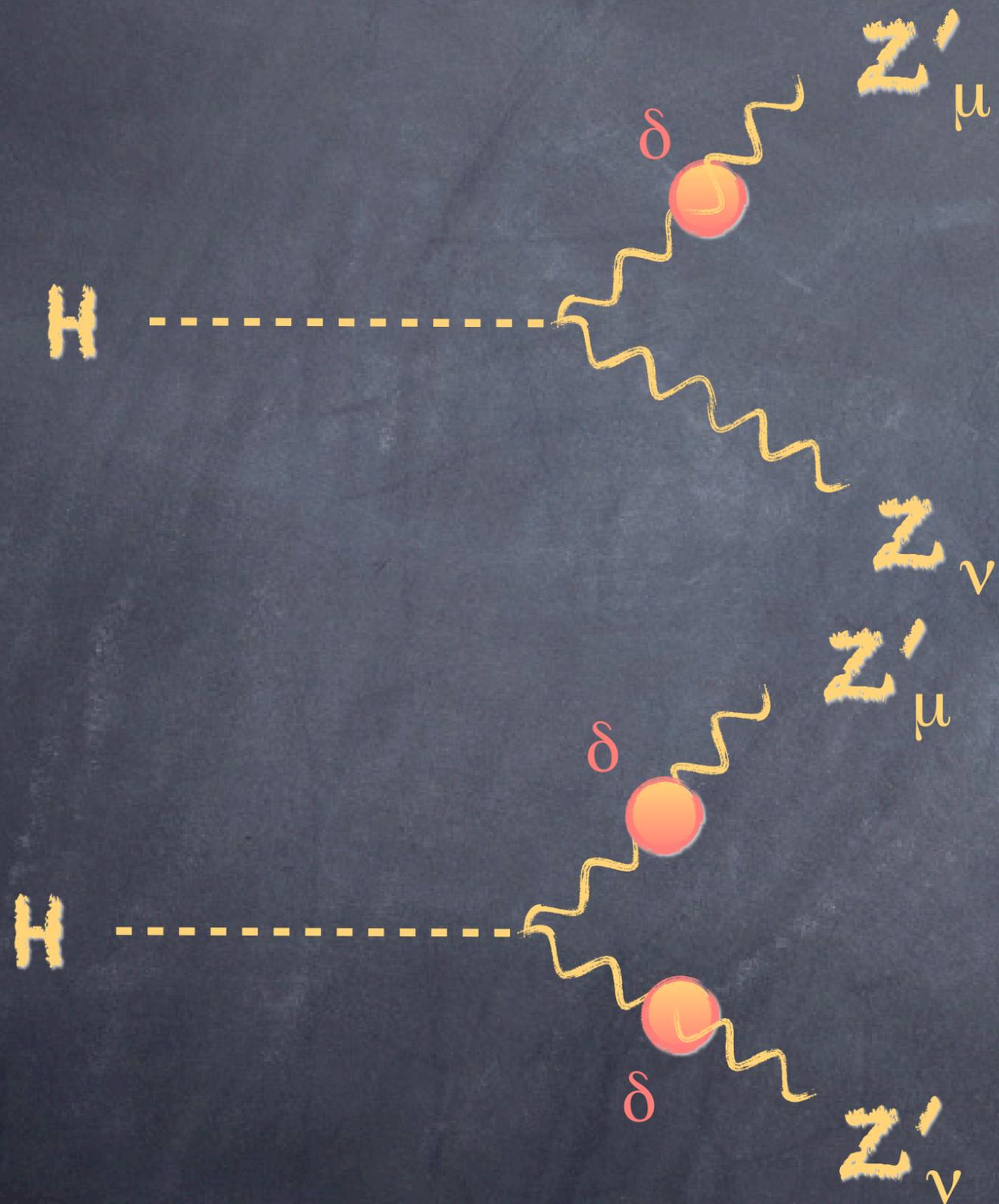
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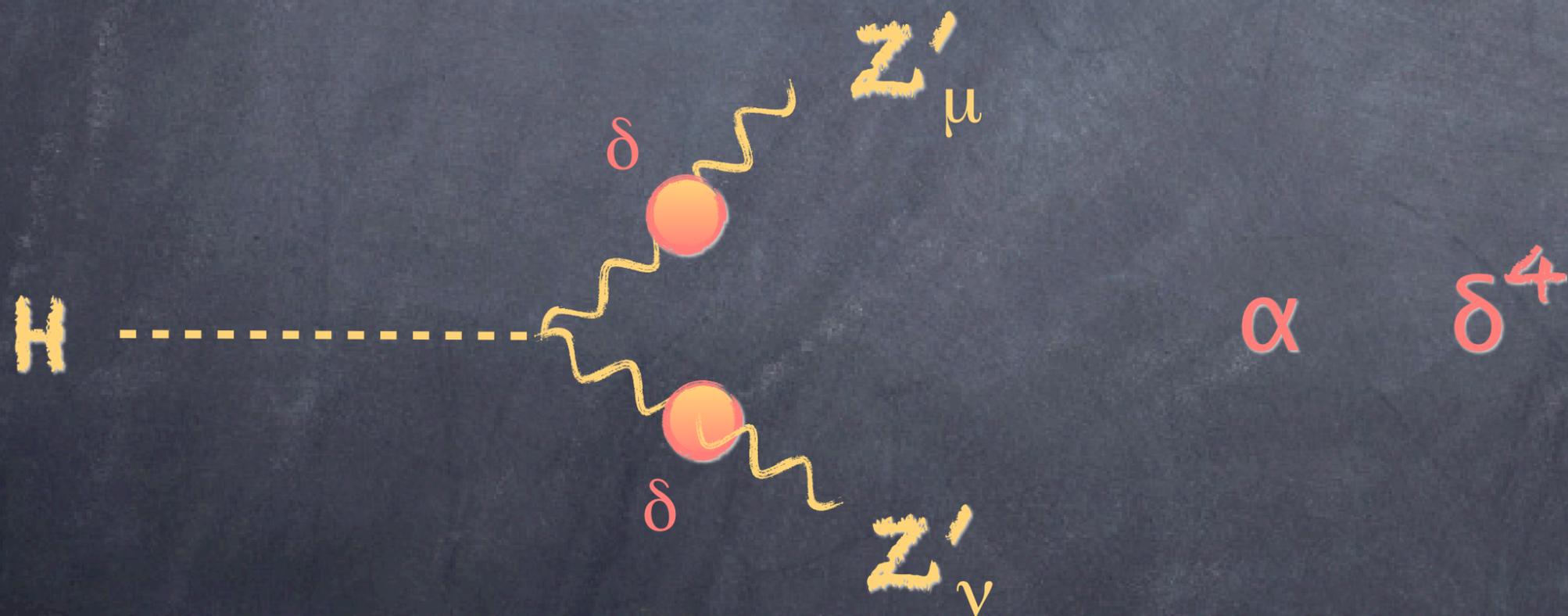
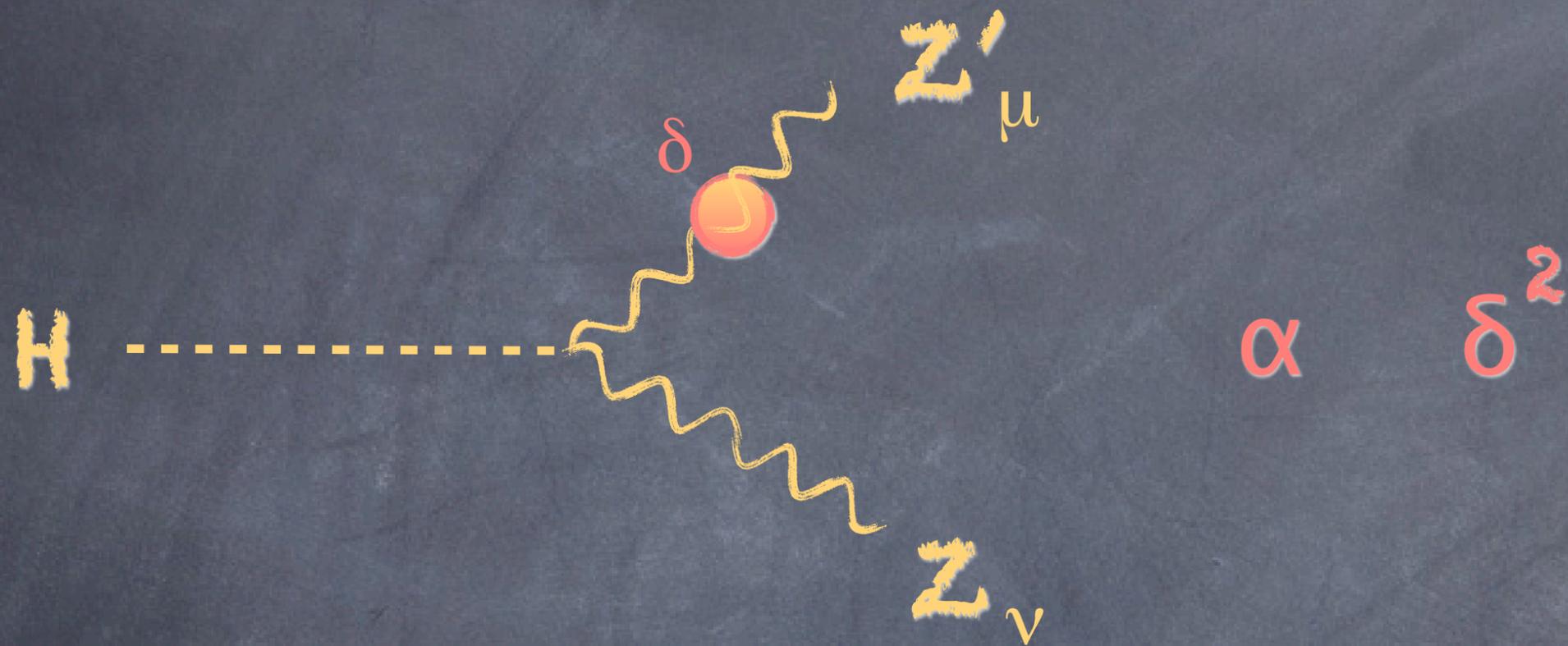
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Constraint from Higgs physics?

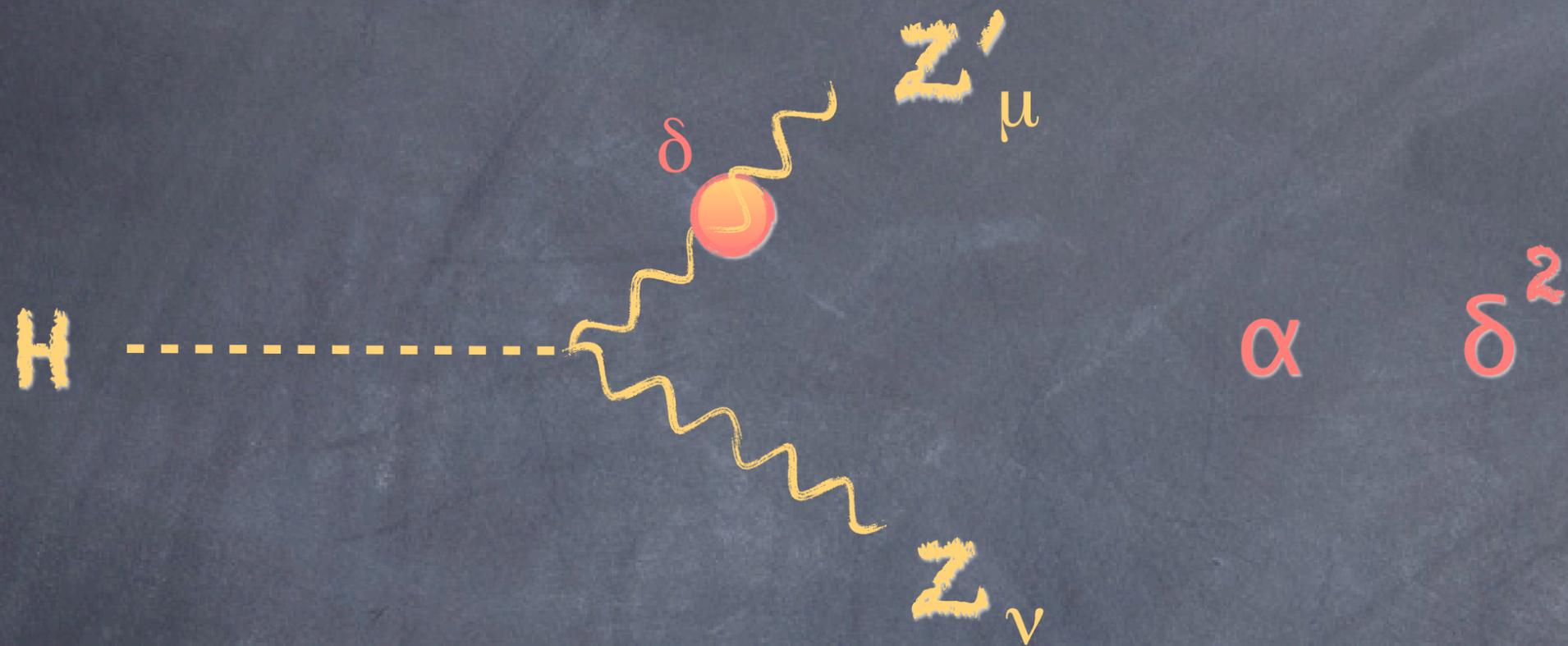
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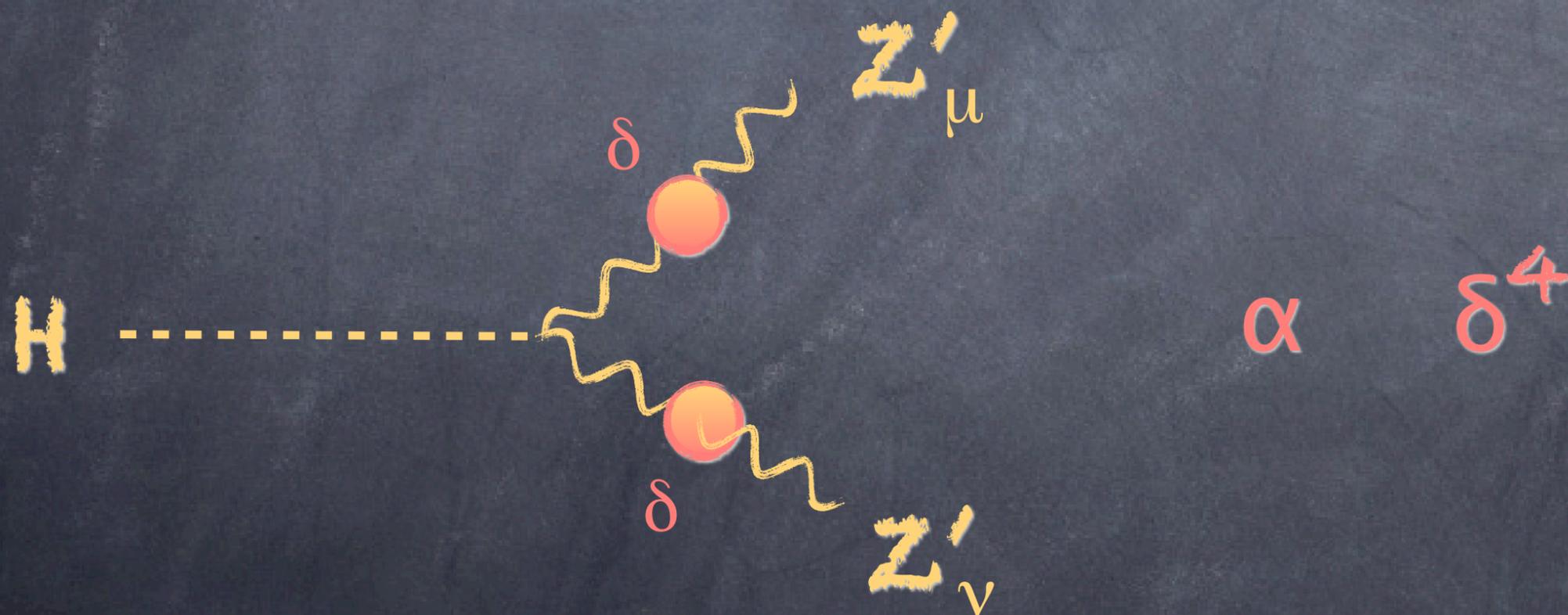
$$M_{Z'} = 5 \text{ GeV}$$

$$\delta = 5 \times 10^{-3}$$

$$m_h = 125 \text{ GeV}$$

$$\Rightarrow \Gamma_{h \rightarrow ZZ'} \simeq 10^{-10} \text{ GeV}$$

$$(\Gamma_h^{SM} = 3.9 \times 10^{-3} \text{ GeV})$$



Y. Mambrini  
2011

# Singlet Extension of the SM: Higgs portal

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Y. Mambrini  
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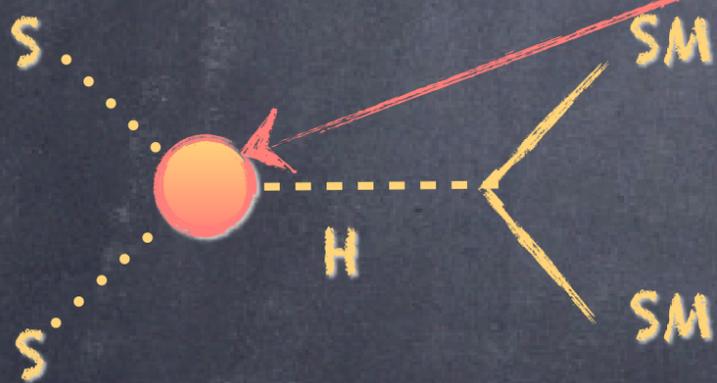
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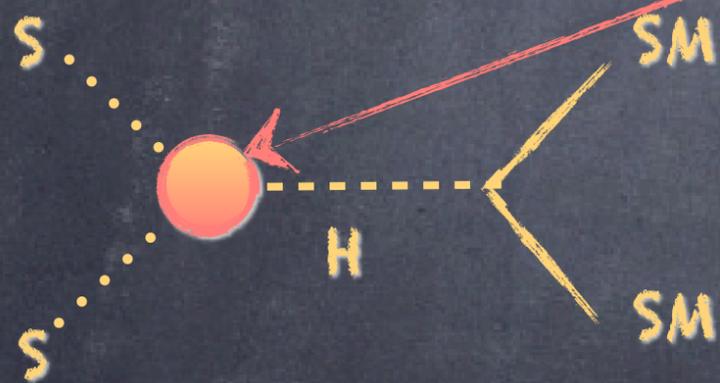
breaking

$\rightarrow$  Higgs mixes with S

$\rightarrow S \rightarrow f\bar{f}$  possible and is thus not a viable DM candidate.

Solved by imposing a  $Z_2$  symmetry

$S \rightarrow -S$



$$\langle \sigma_{f\bar{f}v} \rangle = \frac{\lambda_{HS}^2 (m_S^2 - m_f^2)^{3/2} m_f^2}{16\pi m_S^3 [(4m_S^2 - M_H^2)^2 + M_H^2 \Gamma_H^2]}$$

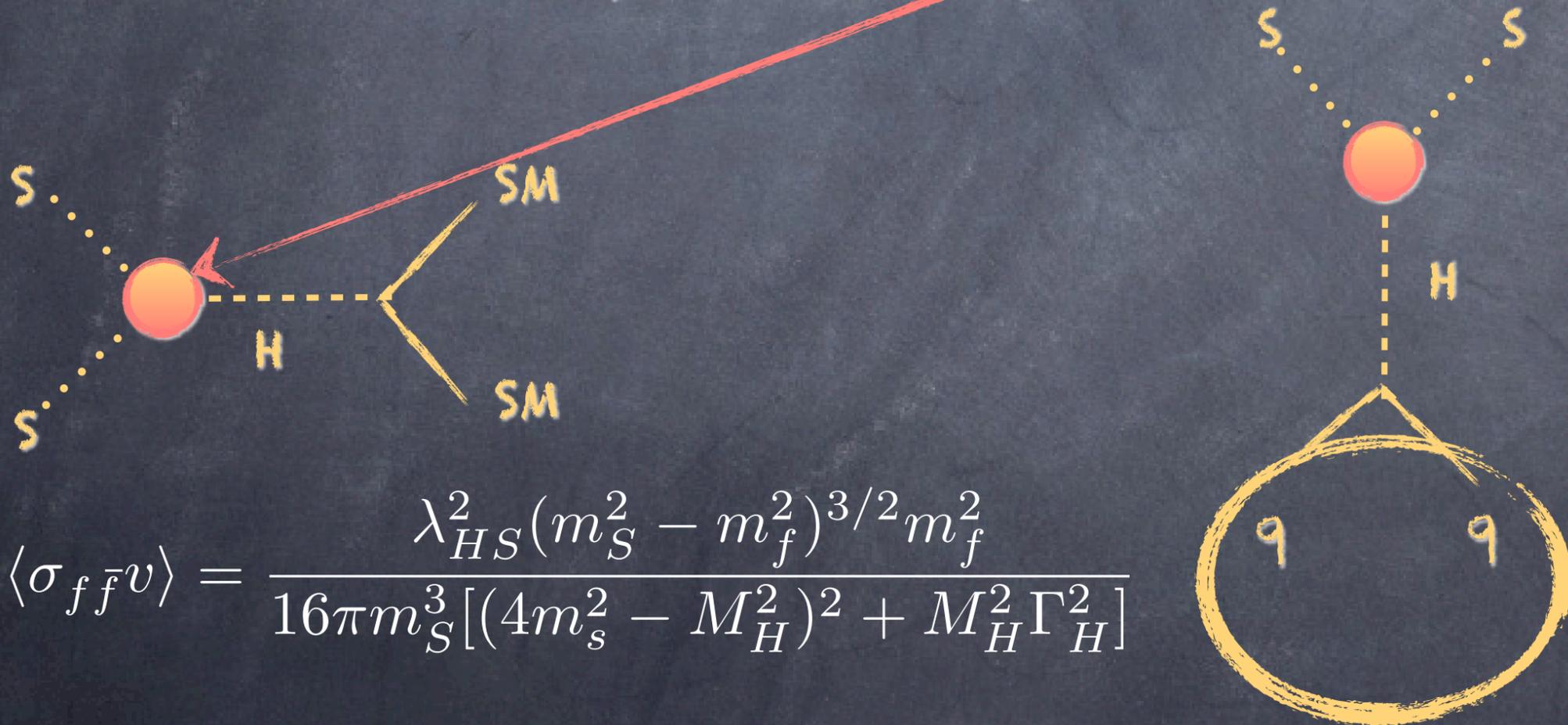
# Singlet Extension of the SM: Higgs portal

Y. Mambrini  
2011

To build the simplest gauge invariant extension of the SM

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \partial_\mu S \partial^\mu S - \frac{\lambda_S}{4} S^4 - \frac{\mu_S^2}{2} S^2 - \frac{\lambda_{HS}}{4} S^2 H^\dagger H - \frac{\kappa_1}{2} H^\dagger H S - \frac{\kappa_3}{3} S^3$$

No phenomenology ( $\langle S \rangle = 0$ )



Stability of  $S$  as DM candidate:  
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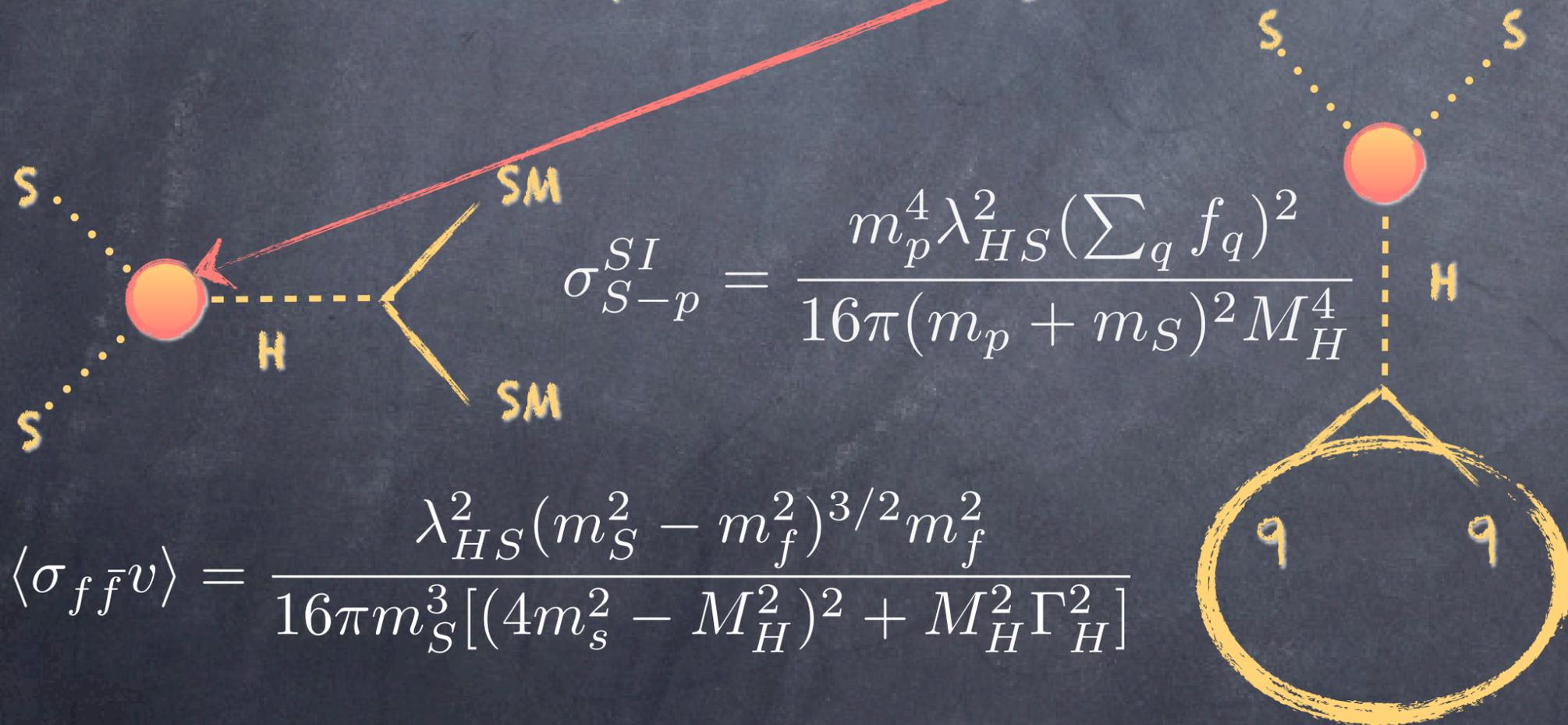
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$$\sigma_{S-p}^{SI} = \frac{m_p^4 \lambda_{HS}^2 (\sum_q f_q)^2}{16\pi (m_p + m_S)^2 M_H^4}$$

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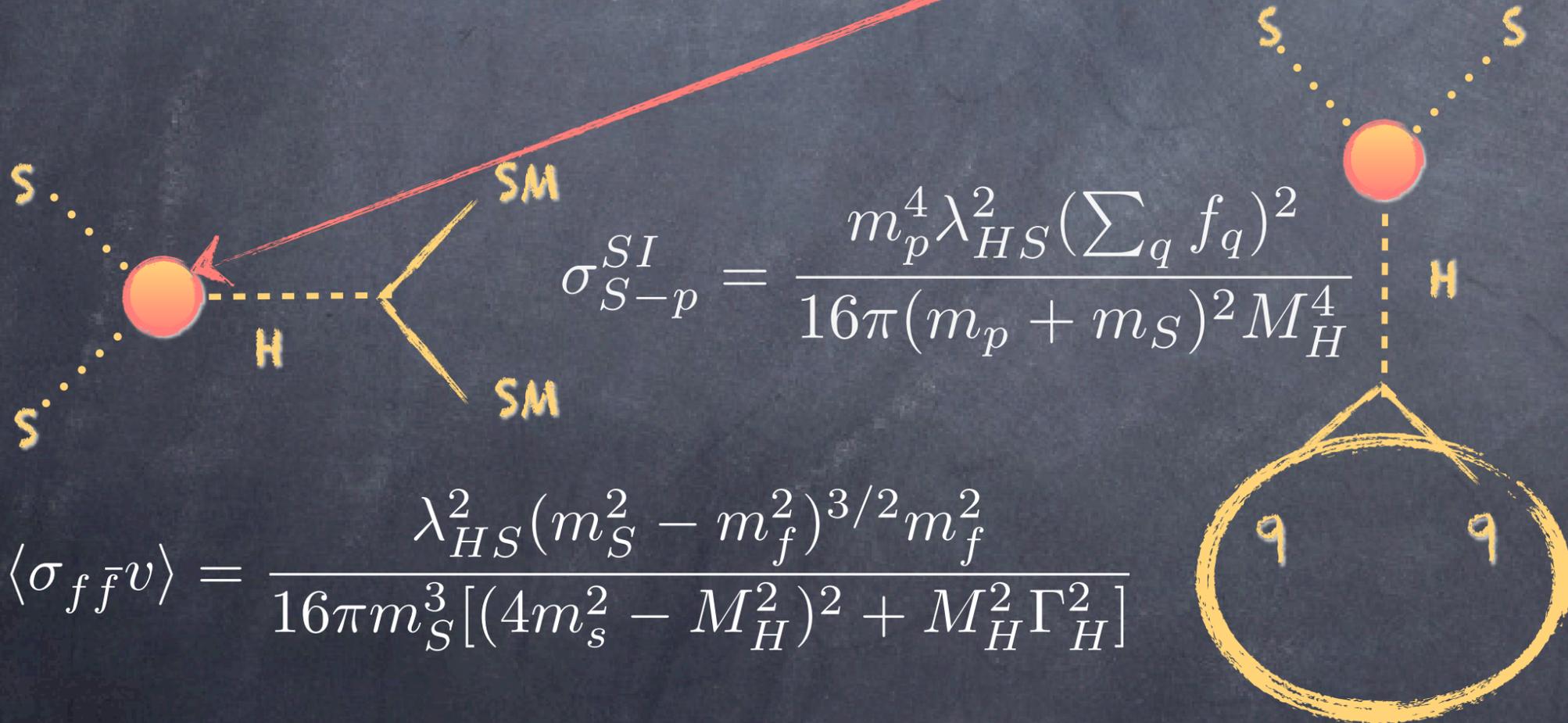
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Y. Mambrini  
2011

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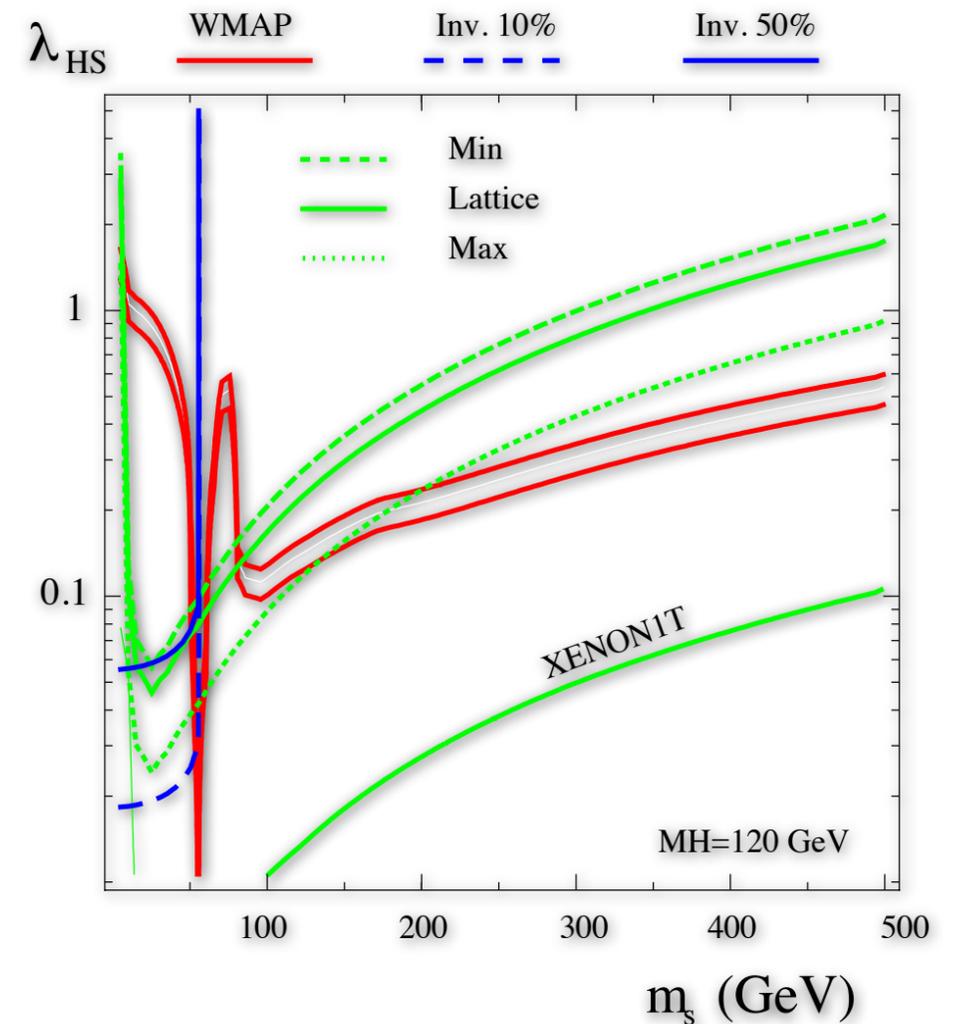
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No phenomenology ( $\langle S \rangle = 0$ )



$$\sigma_{S-p}^{SI} = \frac{m_p^4 \lambda_{HS}^2 (\sum_q f_q)^2}{16\pi (m_p + m_S)^2 M_H^4}$$

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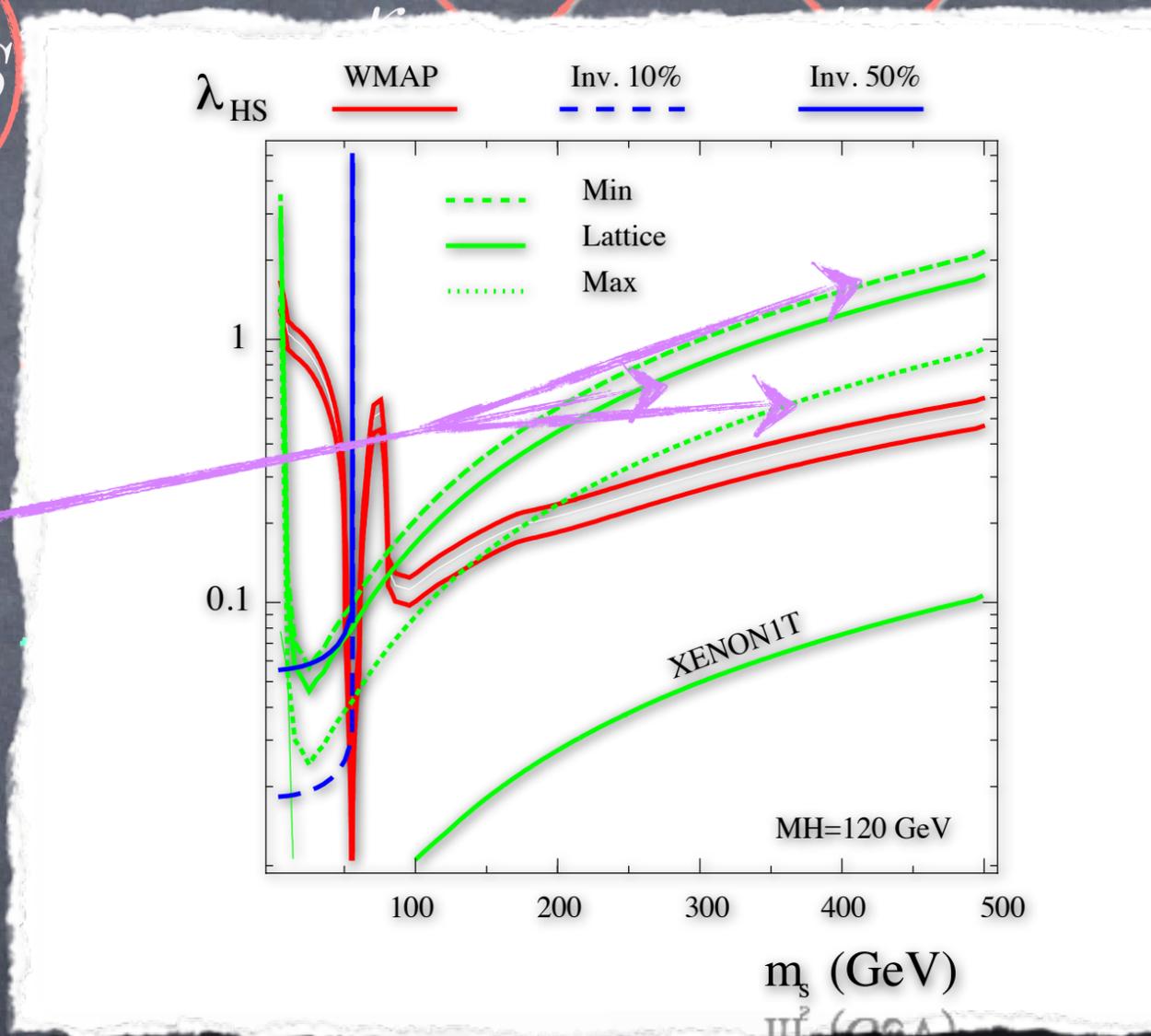
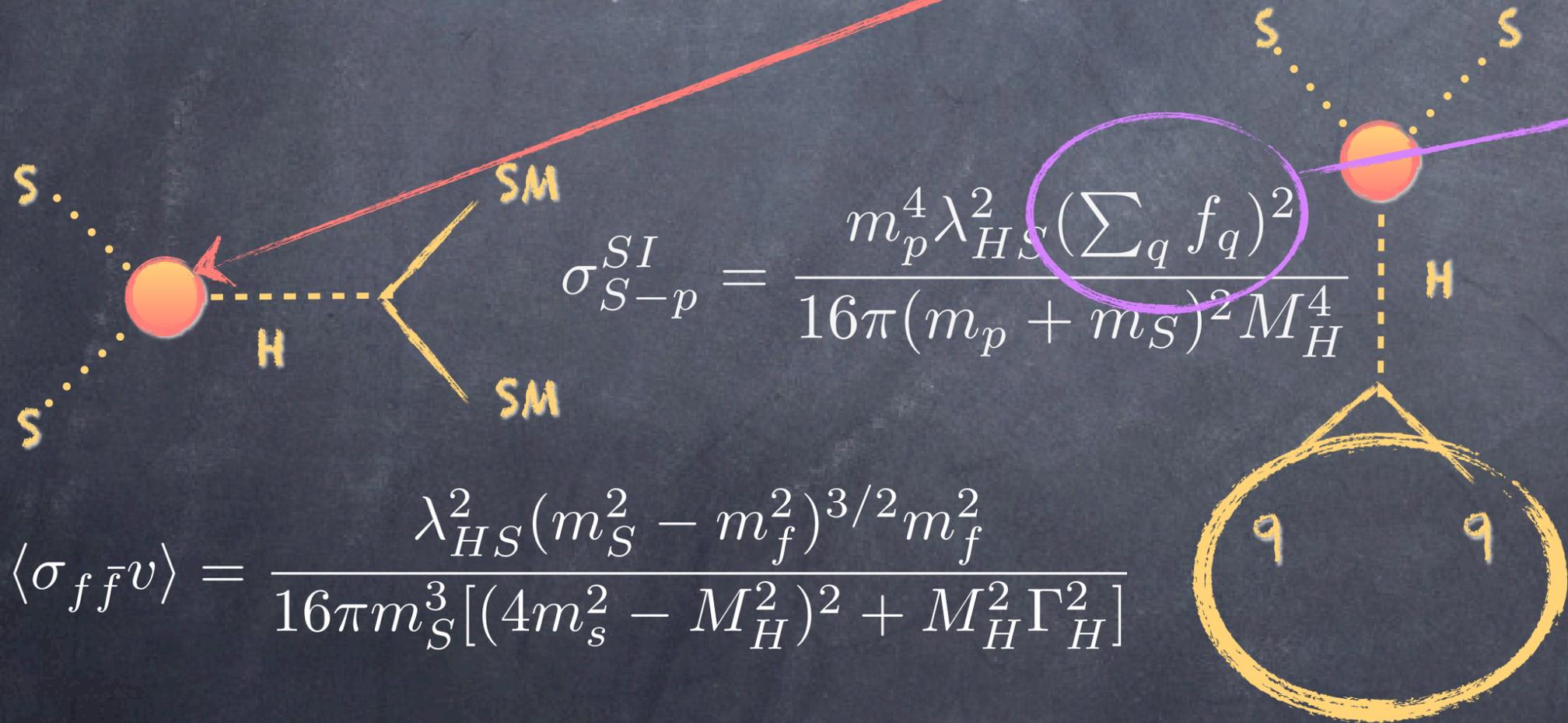
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No phenomenology ( $\langle S \rangle = 0$ )



Y. Mambrini  
1108.0671

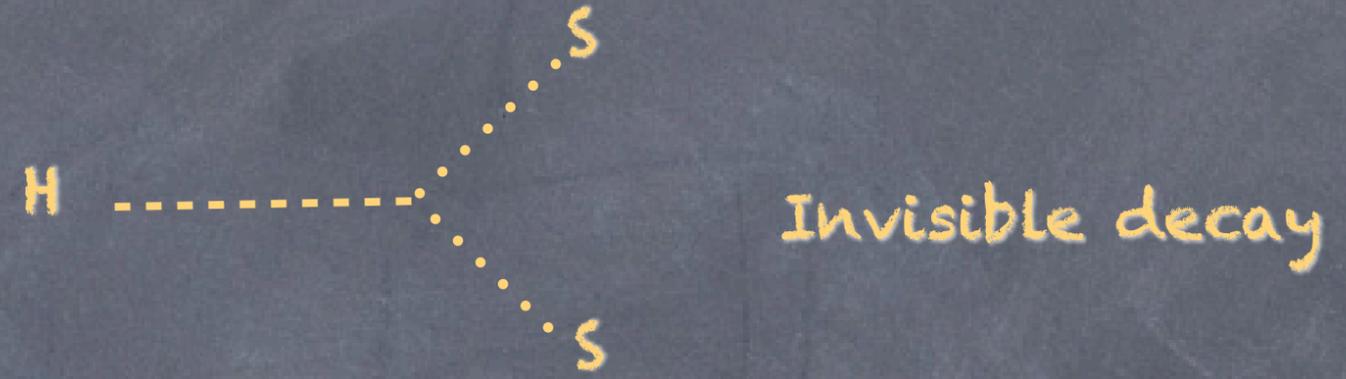
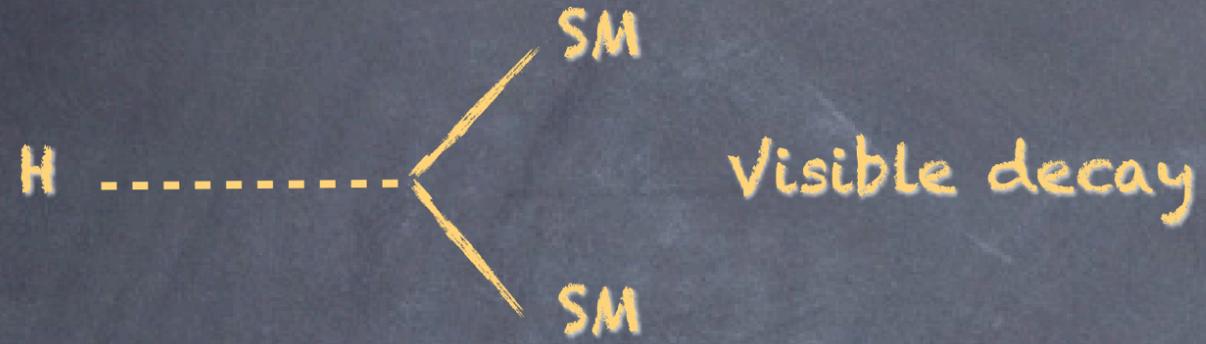
# Invisible width of the Higgs

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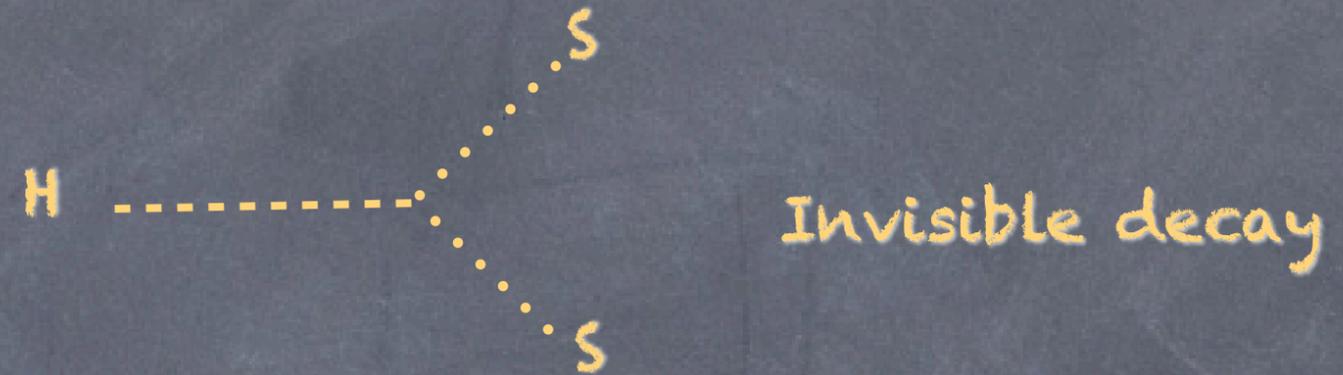
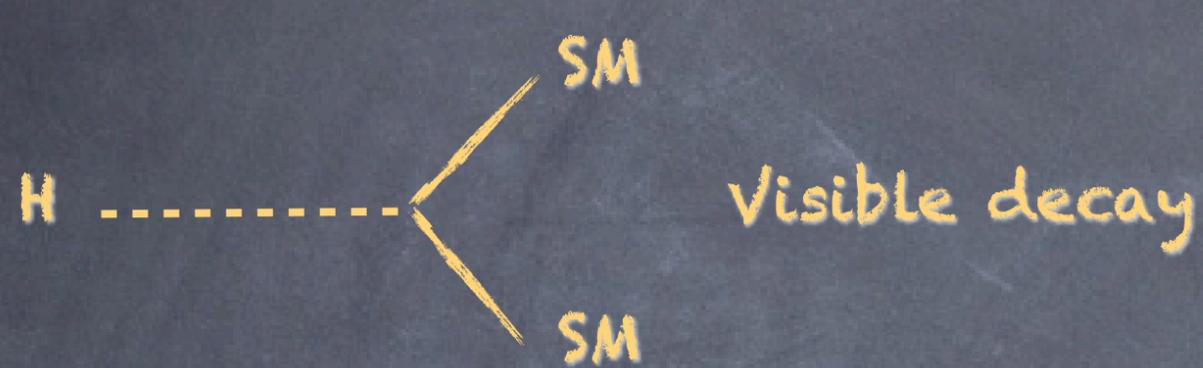
# Invisible width of the Higgs

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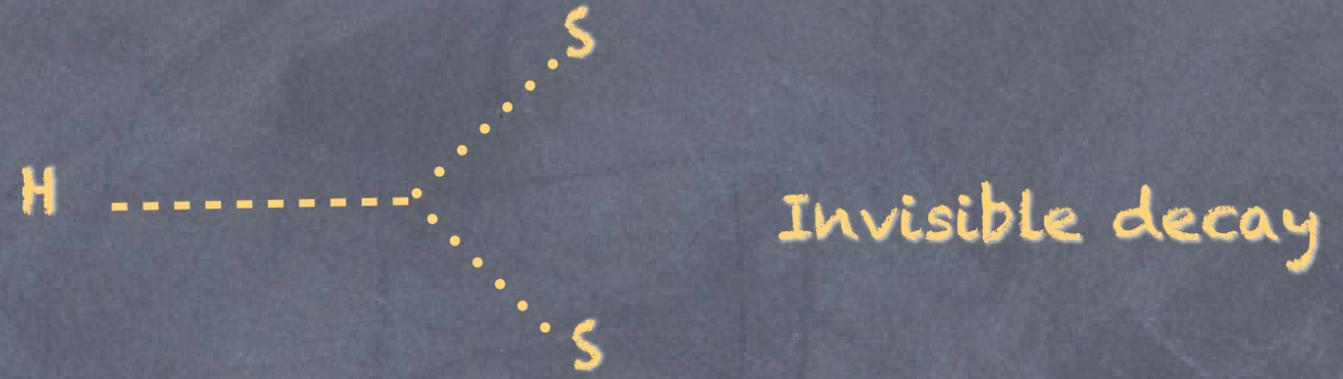
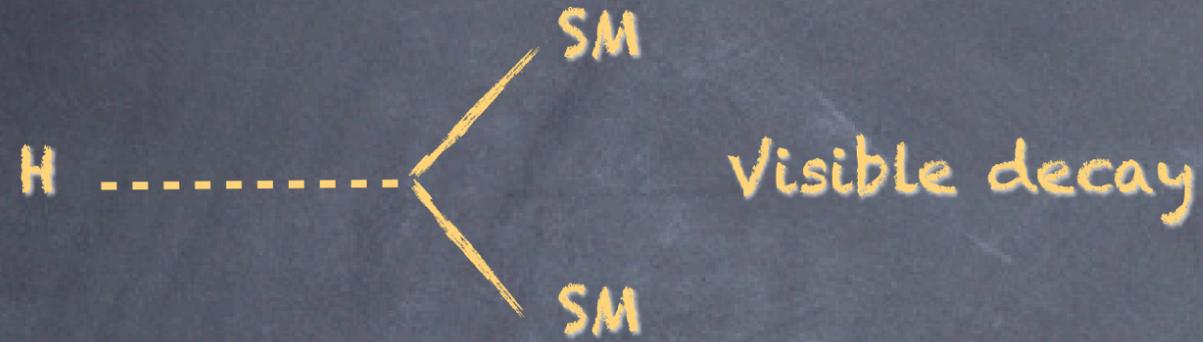
Y. Mambrini  
1108.0671

# Invisible width of the Higgs



$$\Gamma_H(H \rightarrow SS) = \frac{\lambda_{HS}^2 M_W^2}{32\pi g^2 M_H^2} \sqrt{M_H^2 - 4m_S^2}$$

# Invisible width of the Higgs

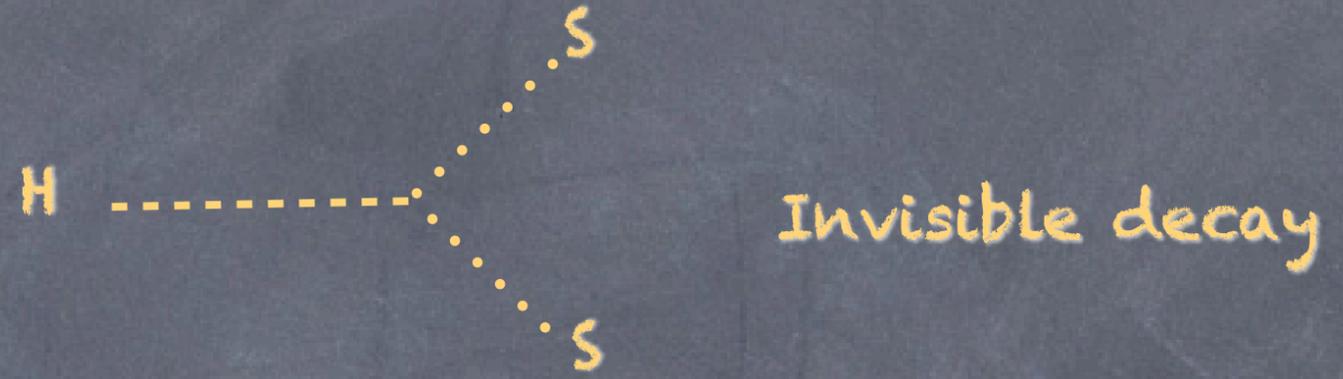
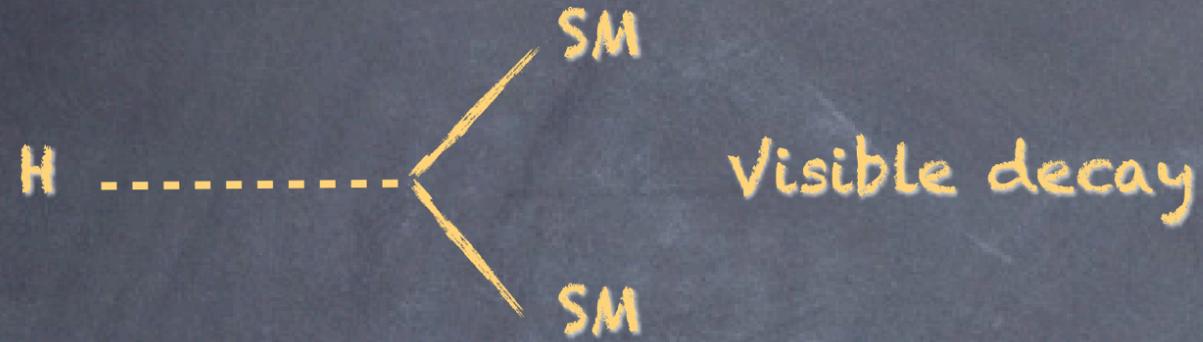


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# Invisible width of the Higgs



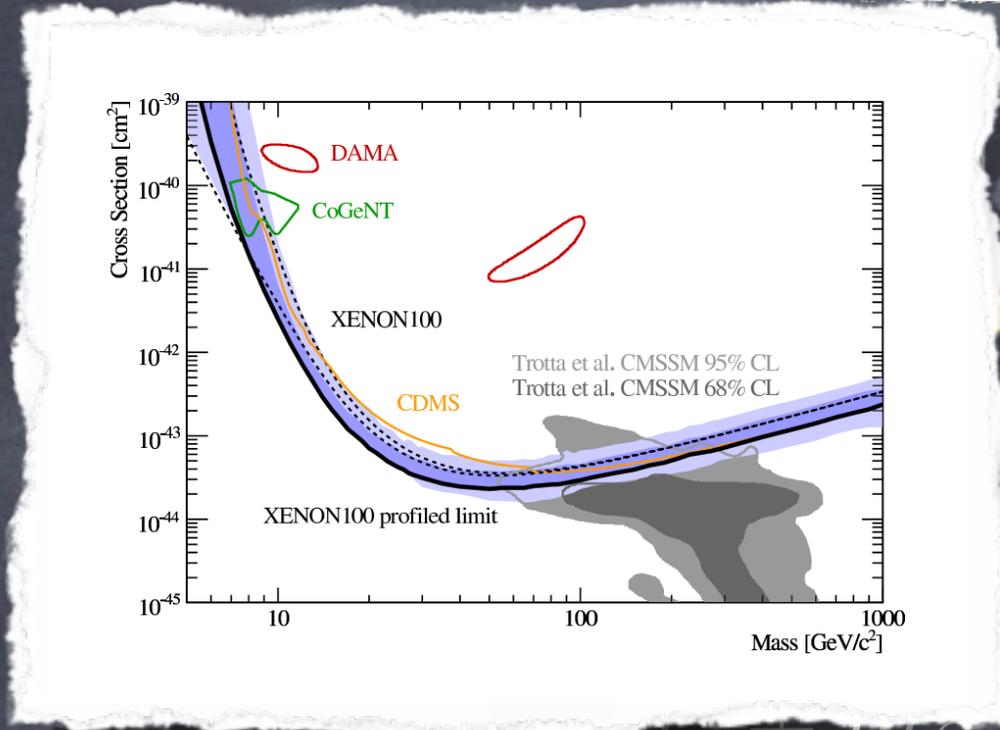
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$$\frac{\Gamma_H^{Inv}}{\sigma_{S-p}^{SI}} = \frac{(m_S + m_p)^2 M_H^2 M_W^2 \sqrt{M_H^2 - 4m_S^2}}{2g^2 f^2 m_p^4}$$

# Invisible width of the Higgs



Visible decay



Invisible decay

$$\Gamma_H(H \rightarrow SS) = \frac{\lambda_{HS}^2 M_W^2}{32\pi g^2 M_H^2} \sqrt{M_H^2 - 4m_S^2}$$

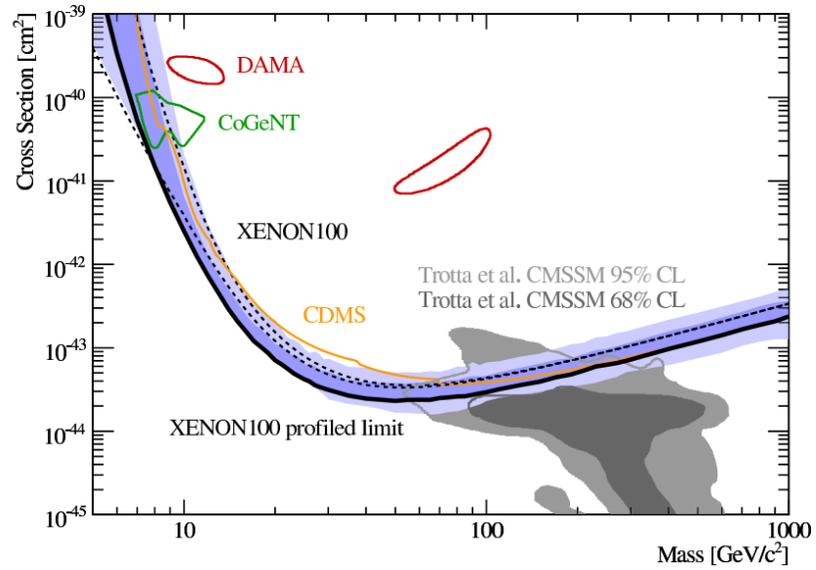
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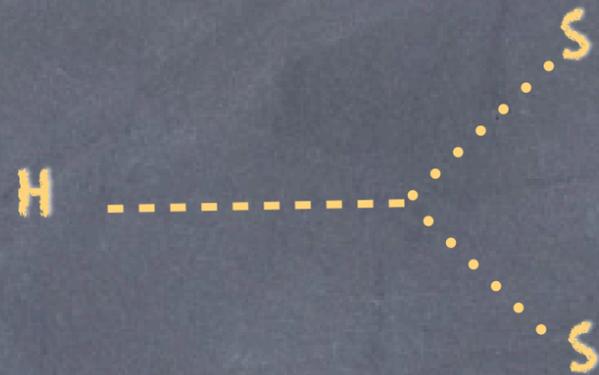
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Y. Mambrini  
1108.0671

# Invisible width of the Higgs



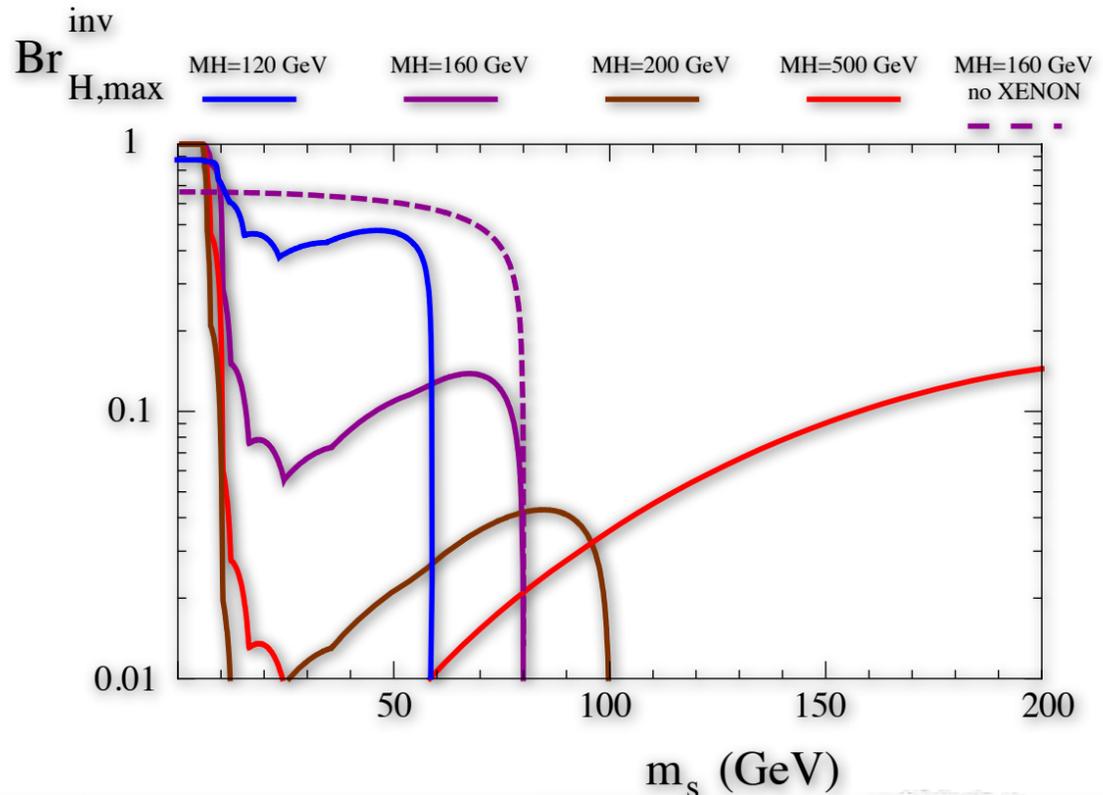
the decay



Invisible decay

$(f_q)^2$

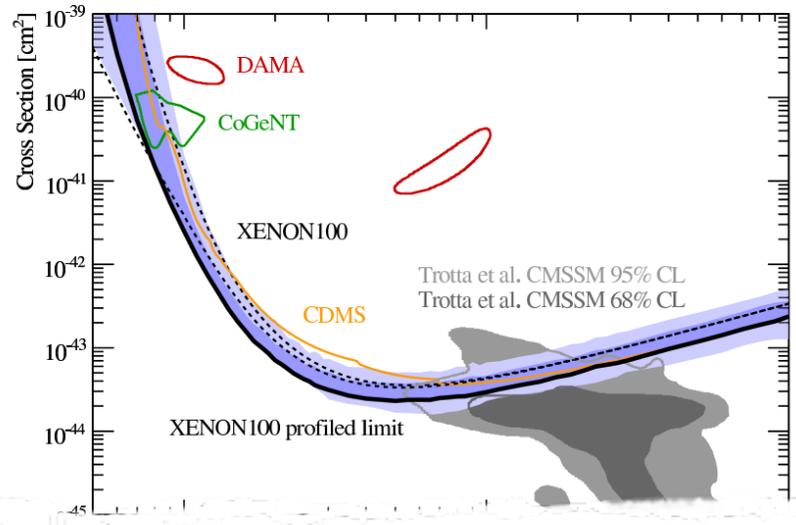
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# Invisible width of the Higgs



Visible decay

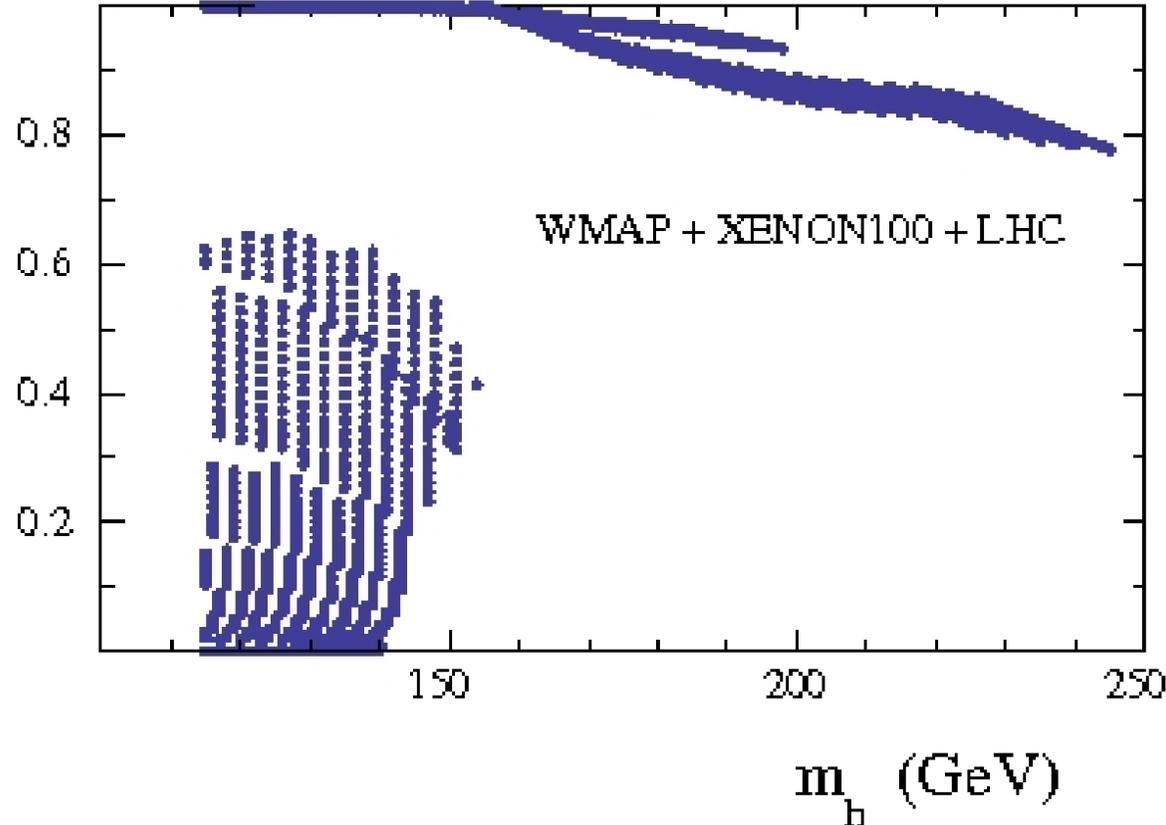
H

S

Invisible decay

$f \sim 2$

Br (h → SS)

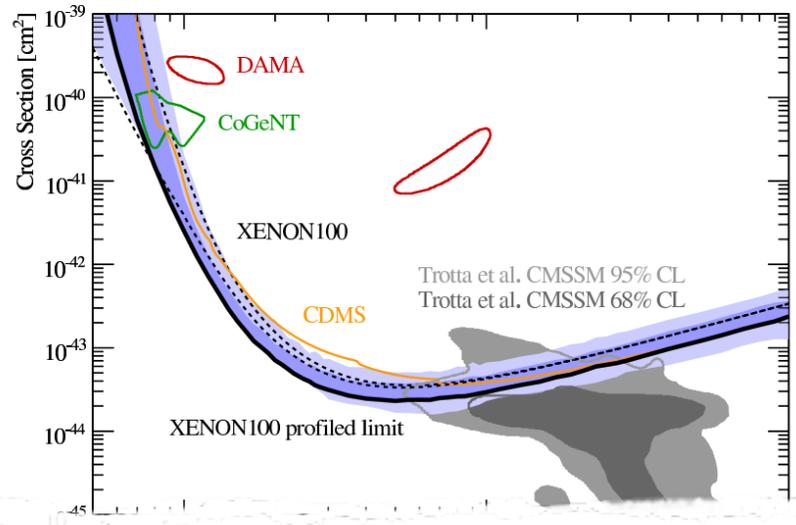


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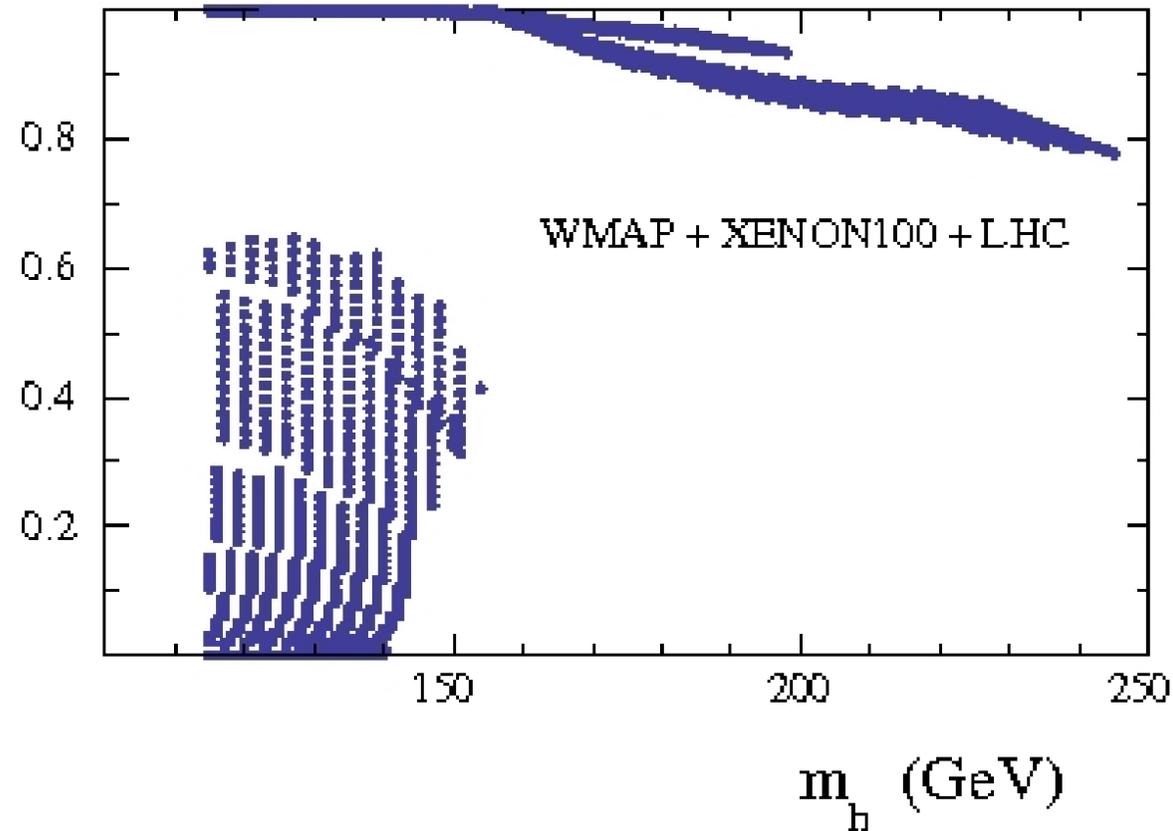
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# Invisible width of the Higgs



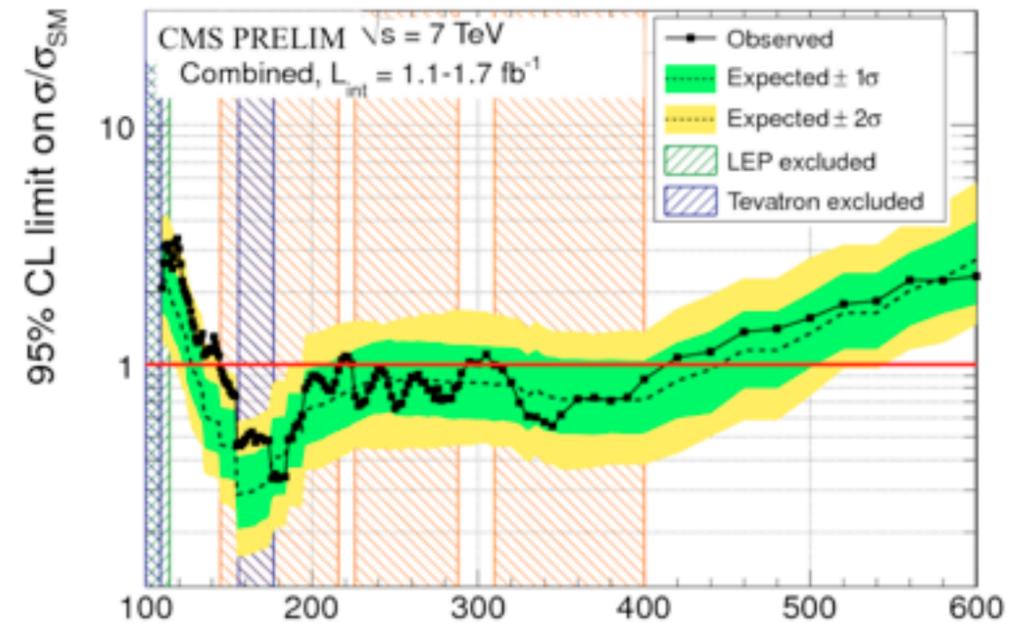
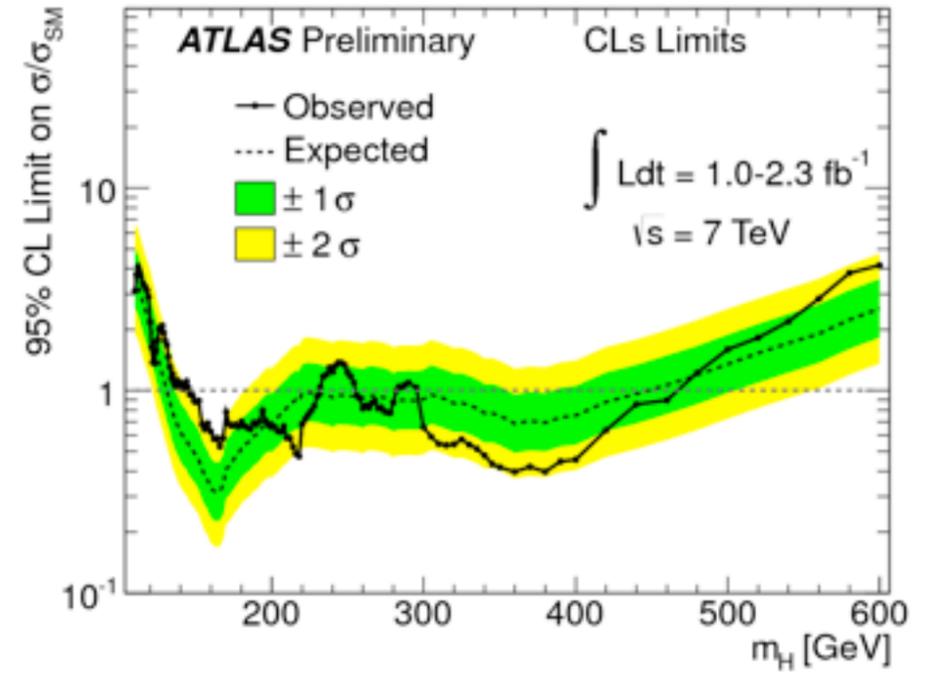
$\Gamma_H(H \rightarrow \nu\nu)$

$\text{Br}(h \rightarrow \nu\nu)$



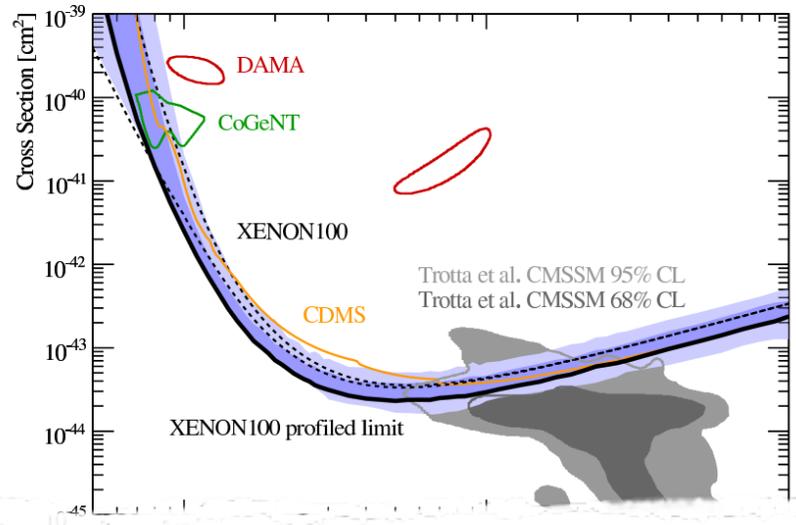
$\Gamma_H(H \rightarrow \nu\nu)$

$$\frac{\Gamma_H^{Inv}}{\sigma_{SI}^{S-p}} =$$



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# Invisible width of the Higgs



the decay

H ---

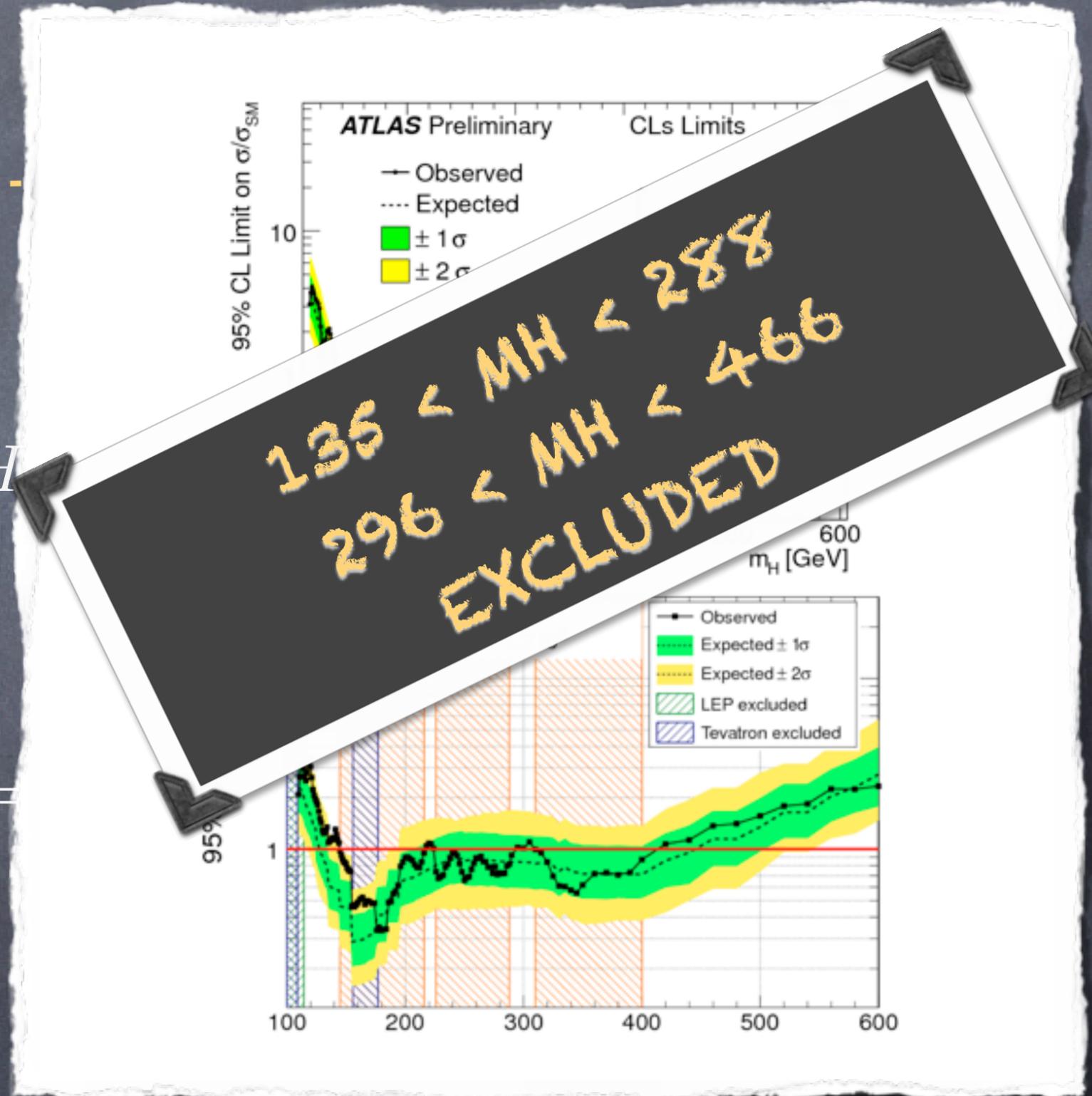
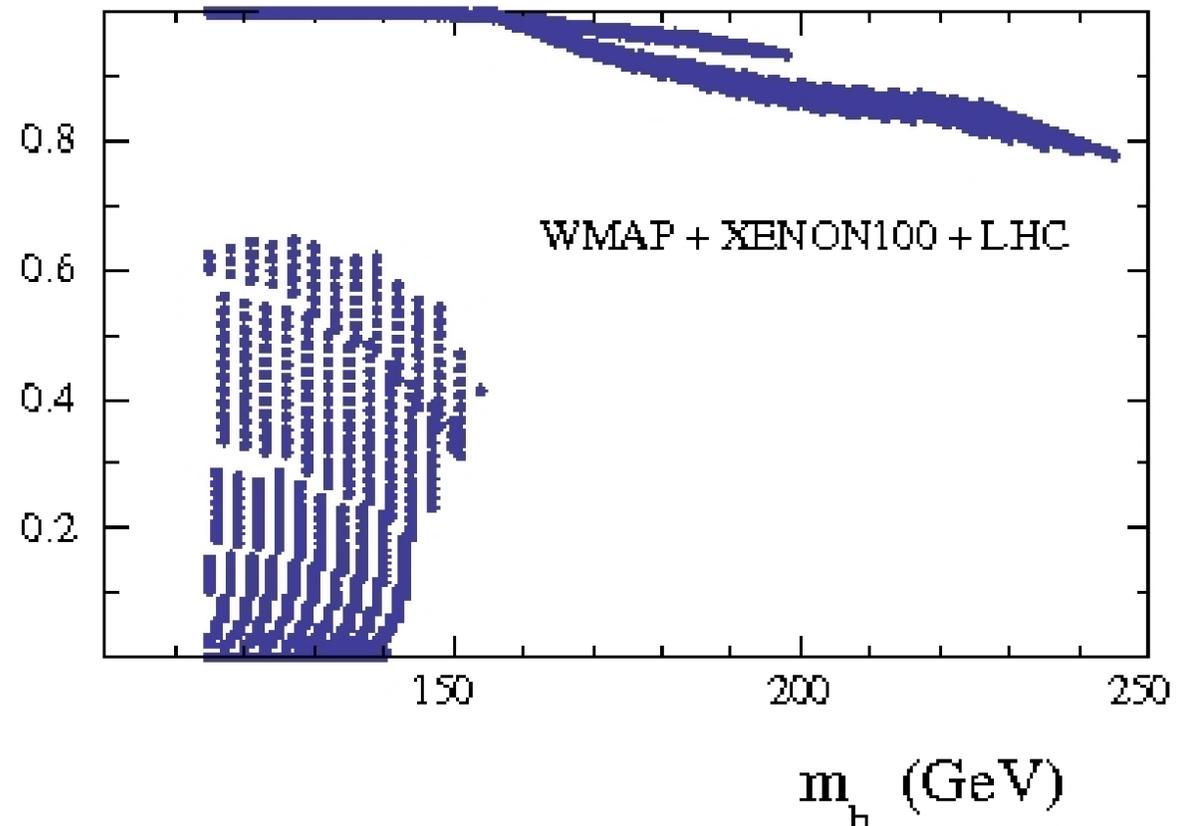
f ) 2

H

$\Gamma_H(H$

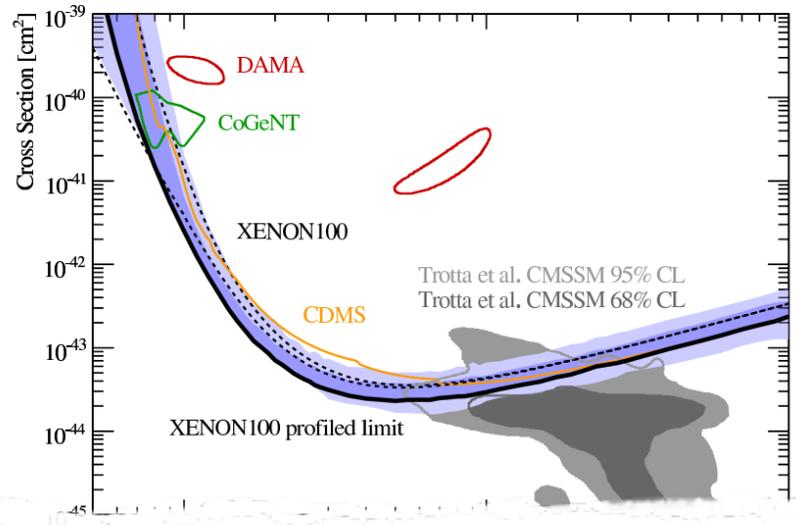
$$\frac{\Gamma_H^{Inv}}{\sigma_{SI}^{S-p}}$$

Br (h → SS)



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# Invisible width of the Higgs



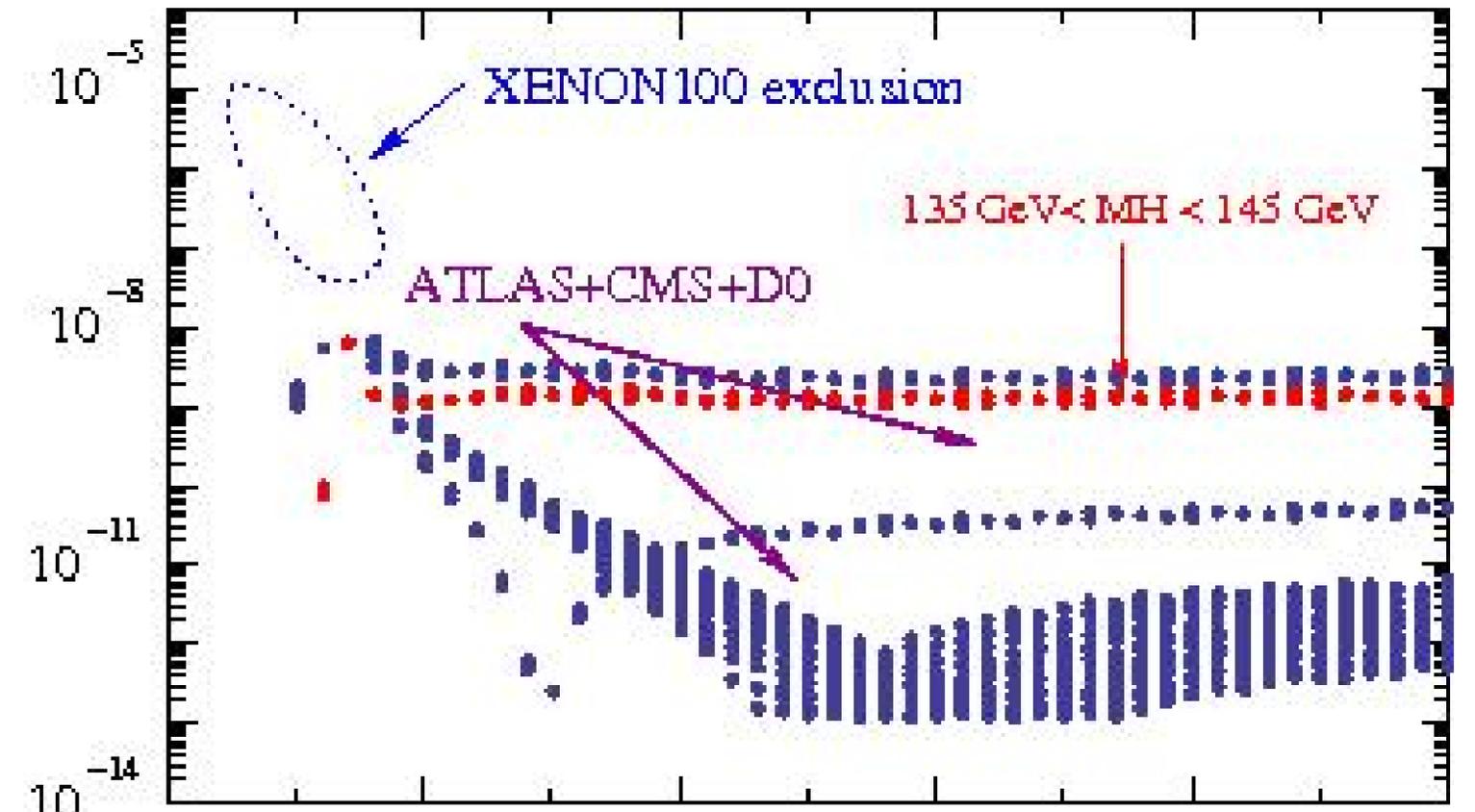
$\Gamma_H$  decay

$f_{SI}^2$

$\Gamma_H$

$$\frac{\Gamma_H^{Inv}}{\sigma_{SI}^{S-}}$$

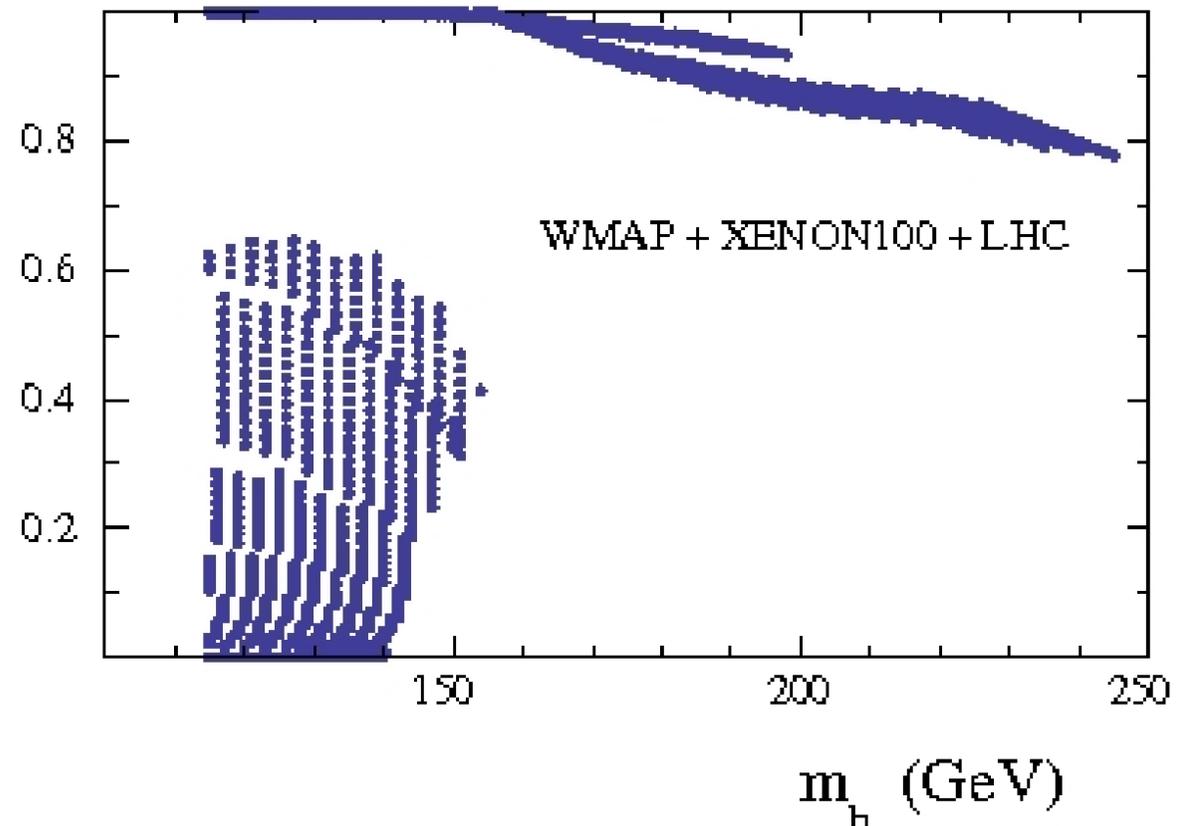
$\sigma_{SI}$



100 300 500  
m<sub>h</sub> (GeV)

100 200 300 400 500 600

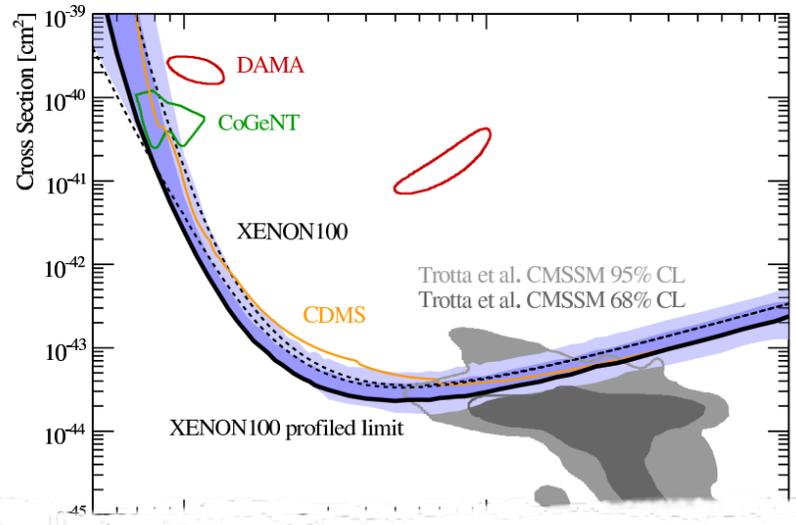
Br (h → SS)



$\Omega_{CDM}^2$

Y. Mambrini  
1108.0671

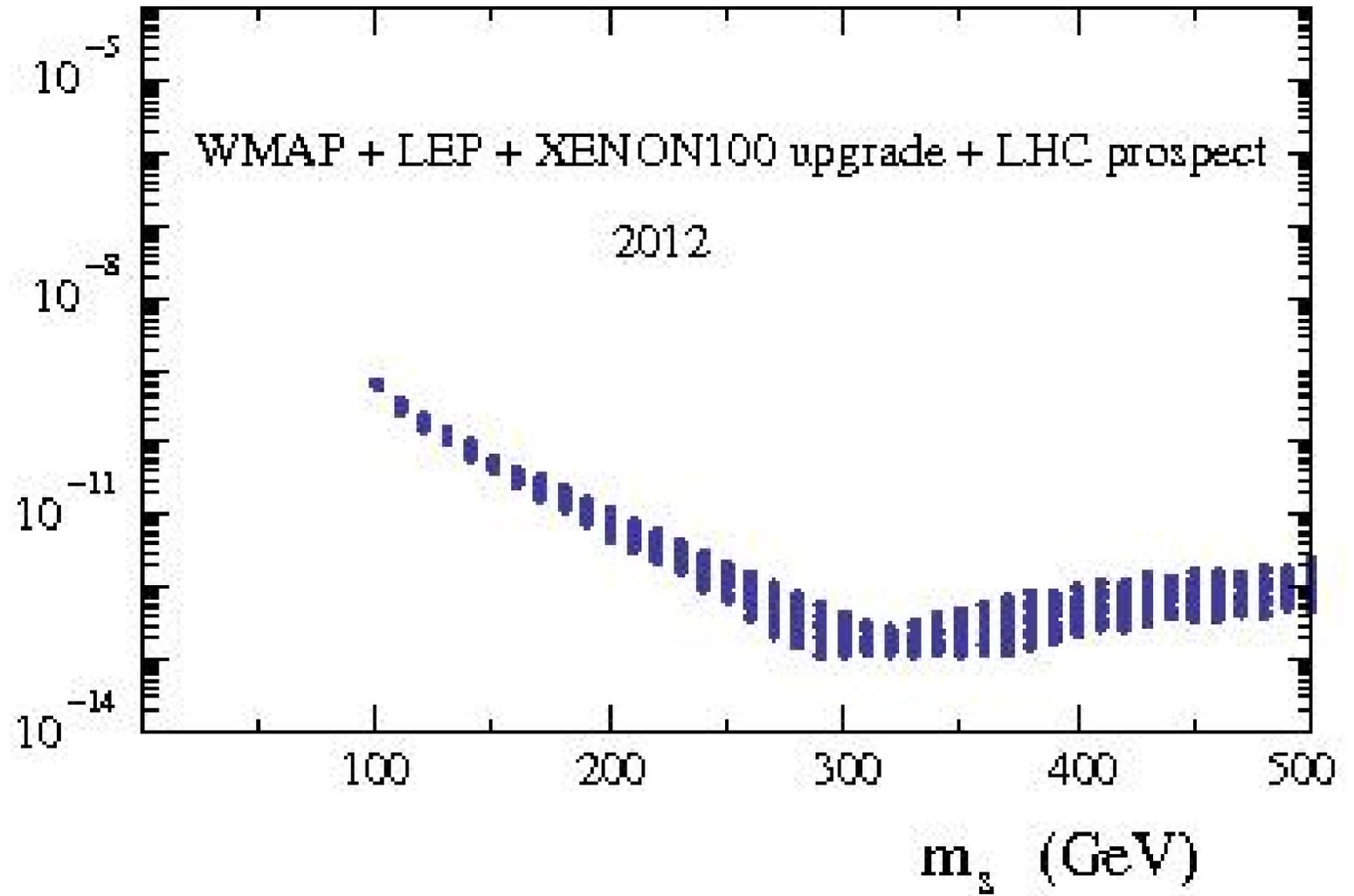
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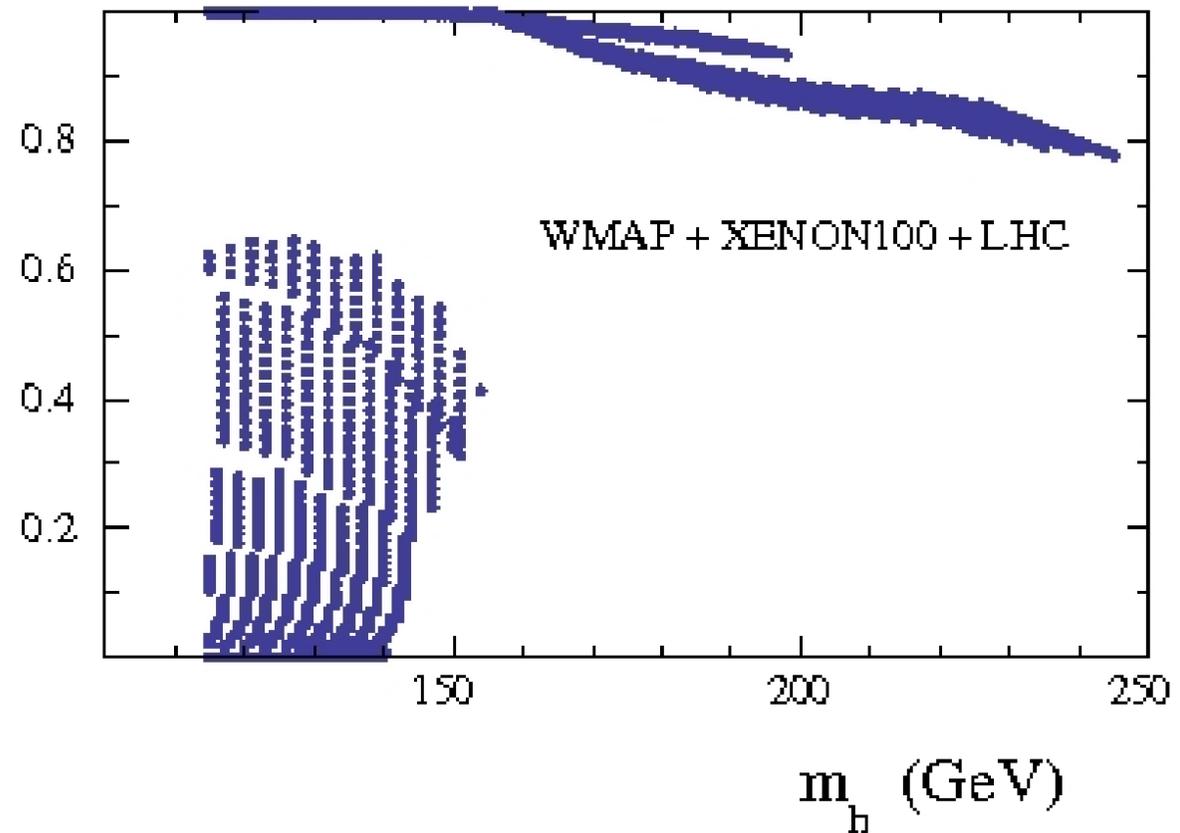
the decay

$f \sim 2$

$\sigma_{SI}$  (pb)



Br (h → SS)



# Vectorial Dark Matter?

O. Lebedev,  
H.M. Lee,  
Y. Mambrini  
1111.4482

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H.M. Lee,  
Y. Mambrini  
1111.4482

# Vectorial Dark Matter?

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{\lambda_{hV}}{4} H^\dagger H V^\mu V_\mu + \frac{\lambda}{4} (V^\mu V_\mu)^2 + \frac{1}{2} m^2 V^\mu V_\mu$$

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$$\Gamma_{H \rightarrow V_\mu V_\mu}^{inv} = \frac{\lambda_{hV}^2 M_W^2 m_H^3}{128\pi m_V^4} \left( 1 - 4 \frac{m_V^2}{m_H^2} + 12 \frac{m_V^4}{m_H^4} \right) \sqrt{1 - 4 \frac{m_V^2}{m_H^2}}$$

# Vectorial Dark Matter?

O. Lebedev,  
H.M. Lee,  
Y. Mambrini  
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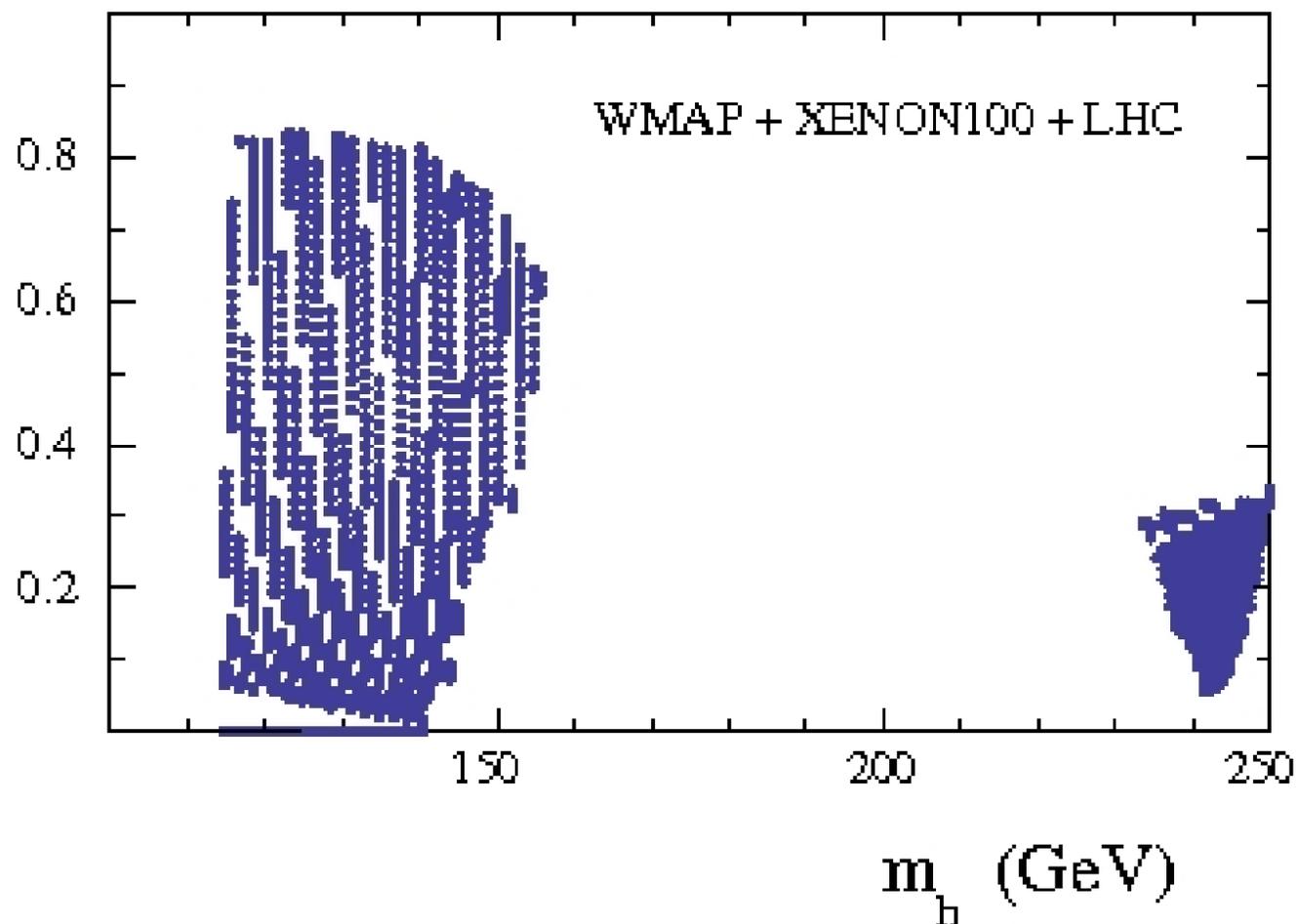
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# Vectorial Dark Matter?

O. Lebedev,  
H.M. Lee,  
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1111.4482

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Br (h → XX)



$$+ 12 \frac{m_V^4}{m_H^4} \sqrt{1 - 4 \frac{m_V^2}{m_H^2}}$$

$$m_S^2$$

A. Djouadi,  
O. Lebedev,  
Y. Mambrini,  
J. Quevillon  
1112.3299

# A 125 GeV Higgs?

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O. Lebedev,  
Y. Mambrini,  
J. Quevillon  
1112.3299

# A 125 GeV Higgs?

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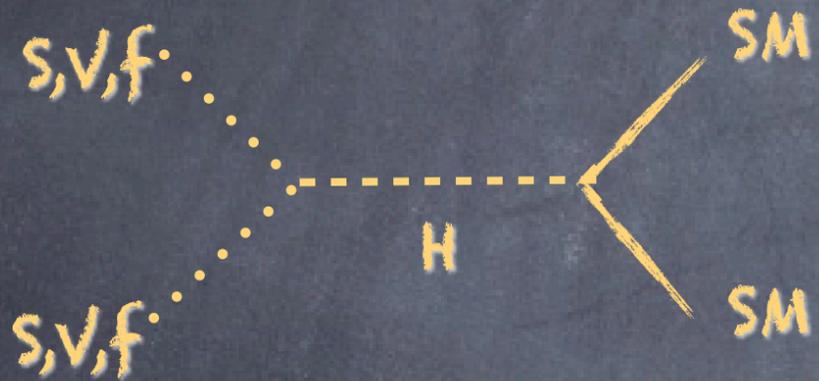
Idea : to combine ATLAS/CMS Higgs analysis in a model independant approach

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O. Lebedev,  
Y. Mambrini,  
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1112.3299

# A 125 GeV Higgs?

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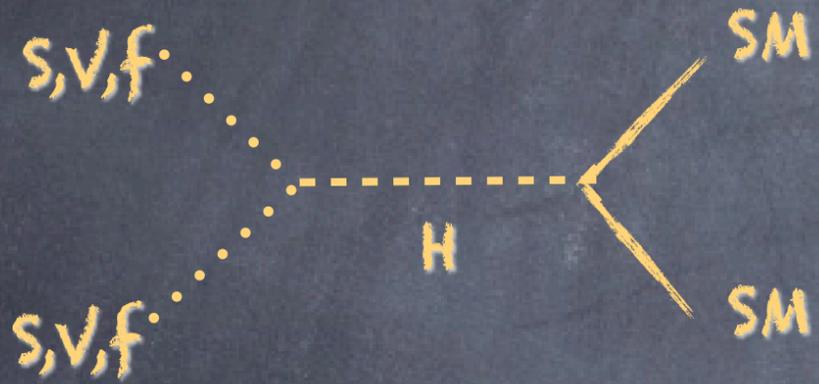
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# A 125 GeV Higgs?

Idea : to combine ATLAS/CMS Higgs analysis in a model independant approach



$$\mathcal{L}_S = \mathcal{L}_{SM} - \frac{1}{2}m_S^2 S^2 - \frac{1}{4}\lambda_S S^4 - \frac{1}{4}\lambda_{hSS} H^\dagger H S^2$$

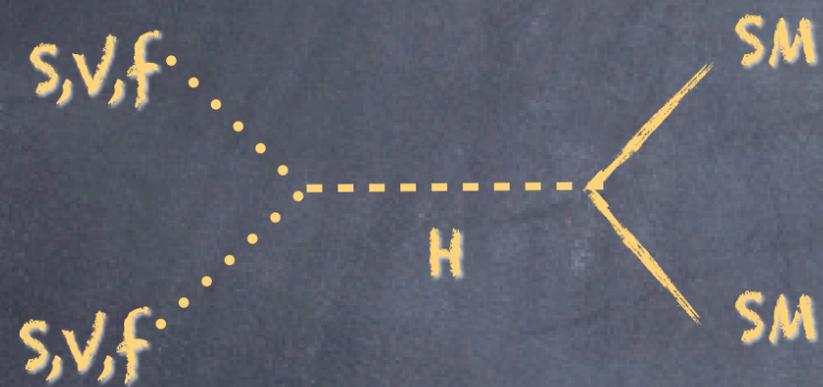
$$\mathcal{L}_V = \mathcal{L}_{SM} + \frac{1}{2}m_V^2 V_\mu V^\mu + \frac{1}{4}\lambda_V (V_\mu V^\mu)^2 + \frac{1}{4}\lambda_{hVV} H^\dagger H V_\mu V^\mu$$

$$\mathcal{L}_f = \mathcal{L}_{SM} - \frac{1}{2}m_f \bar{\chi}\chi - \frac{1}{4}\frac{\lambda_{hff}}{\Lambda} H^\dagger H \bar{\chi}\chi$$

# A 125 GeV Higgs?

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O. Lebedev,  
Y. Mambrini,  
J. Quevillon  
1112.3299

Idea : to combine ATLAS/CMS Higgs analysis in a model independant approach



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$$\mathcal{L}_V = \mathcal{L}_{SM} + \frac{1}{2}m_V^2 V_\mu V^\mu + \frac{1}{4}\lambda_V (V_\mu V^\mu)^2 + \frac{1}{4}\lambda_{hVV} H^\dagger H V_\mu V^\mu$$

$$\mathcal{L}_f = \mathcal{L}_{SM} - \frac{1}{2}m_f \bar{\chi}\chi - \frac{1}{4} \frac{\lambda_{hff}}{\Lambda} H^\dagger H \bar{\chi}\chi$$

$$\Gamma_{h \rightarrow SS}^{\text{inv}} = \frac{\lambda_{hSS}^2 v^2 \beta_S}{64\pi m_h}$$

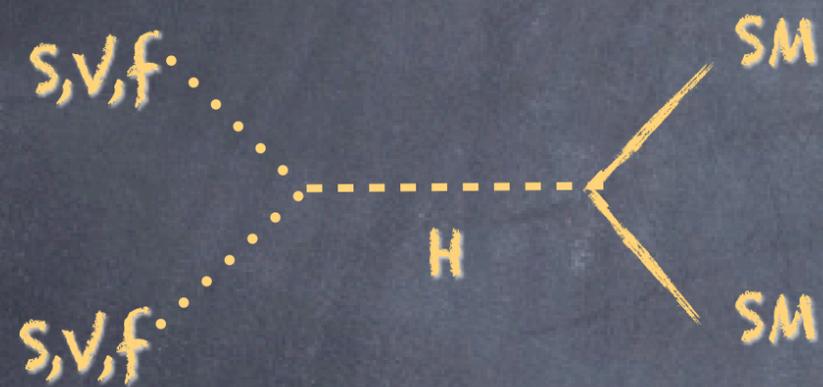
$$\Gamma_{h \rightarrow VV}^{\text{inv}} = \frac{\lambda_{hVV}^2 v^2 m_h^3 \beta_V}{256\pi M_V^4} \left( 1 - 4 \frac{M_V^2}{m_h^2} + 12 \frac{M_V^4}{m_h^4} \right)$$

$$\Gamma_{h \rightarrow \chi\chi}^{\text{inv}} = \frac{\lambda_{hff}^2 v^2 m_h \beta_f^3}{32\pi \Lambda^2}$$

# A 125 GeV Higgs?

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O. Lebedev,  
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1112.3299

Idea : to combine ATLAS/CMS Higgs analysis in a model independant approach



$$\mathcal{L}_S = \mathcal{L}_{SM} - \frac{1}{2}m_S^2 S^2 - \frac{1}{4}\lambda_S S^4 - \frac{1}{4}\lambda_{hSS} H^\dagger H S^2$$

$$\mathcal{L}_V = \mathcal{L}_{SM} + \frac{1}{2}m_V^2 V_\mu V^\mu + \frac{1}{4}\lambda_V (V_\mu V^\mu)^2 + \frac{1}{4}\lambda_{hVV} H^\dagger H V_\mu V^\mu$$

$$\mathcal{L}_f = \mathcal{L}_{SM} - \frac{1}{2}m_f \bar{\chi}\chi - \frac{1}{4} \frac{\lambda_{hff}}{\Lambda} H^\dagger H \bar{\chi}\chi$$

If  $m_h$  is fixed (125 GeV) for a given DM mass,  $\lambda_{hii}$  is determined by WMAP.

QS: is it compatible with direct detection constraints

( $\lambda_{hii} < \lambda_{hiiMAX}$ )

and a visible Higgs ( $\Gamma_{inv} < 10\%$ ),

$$\Gamma_{h \rightarrow SS}^{inv} = \frac{\lambda_{hSS}^2 v^2 \beta_S}{64\pi m_h}$$

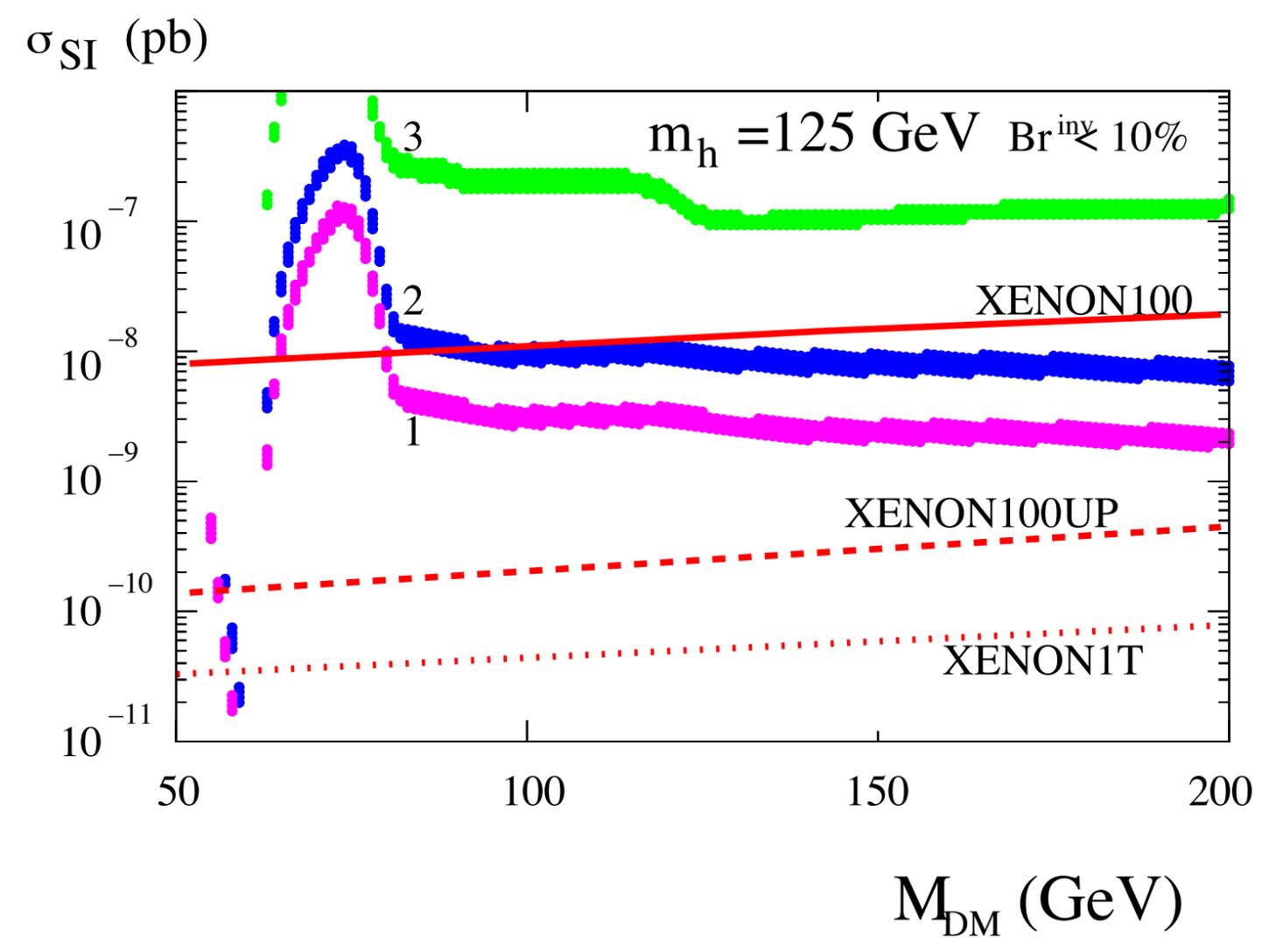
$$\Gamma_{h \rightarrow VV}^{inv} = \frac{\lambda_{hVV}^2 v^2 m_h^3 \beta_V}{256\pi M_V^4} \left( 1 - 4 \frac{M_V^2}{m_h^2} + 12 \frac{M_V^4}{m_h^4} \right)$$

$$\Gamma_{h \rightarrow \chi\chi}^{inv} = \frac{\lambda_{hff}^2 v^2 m_h \beta_f^3}{32\pi \Lambda^2}$$

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 1112.3299

# A 125 GeV Higgs?

Idea : to combine ATLAS/CMS Higgs analysis in a model independant approach



$$\mathcal{L}_{SM} - \frac{1}{2} m_S^2 S^2 - \frac{1}{4} \lambda_S S^4 - \frac{1}{4} \lambda_{hSS} H^\dagger H S^2$$

$$+ \frac{1}{2} g_V^2 V_\mu V^\mu + \frac{1}{4} \lambda_V (V_\mu V^\mu)^2 + \frac{1}{4} \lambda_{hVV} H^\dagger H V_\mu V^\mu$$

$$\mathcal{L}_f = \mathcal{L}_{SM} - \frac{1}{2} m_f \bar{\chi} \chi - \frac{1}{4} \frac{\lambda_{hff}}{\Lambda} H^\dagger H \bar{\chi} \chi$$

$$\Gamma_{h \rightarrow SS}^{inv} = \frac{\lambda_{hSS}^2 v^2 \beta_S}{64\pi m_h}$$

$$\Gamma_{h \rightarrow VV} = \frac{\lambda_{hVV}^2 v^2 m_h^3 \beta_V}{256\pi M_V^4} \left( 1 - 4 \frac{M_V^2}{m_h^2} + 12 \frac{M_V^4}{m_h^4} \right)$$

$$\Gamma_{h \rightarrow \chi\chi}^{inv} = \frac{\lambda_{hff}^2 v^2 m_h \beta_f^3}{32\pi \Lambda^2}$$

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- Several models could be excluded by the end of next year
- Complementarity with LHC is fundamental in any analysis
- Can be applied to any kind of SM extensions (SUSY, KK..)