

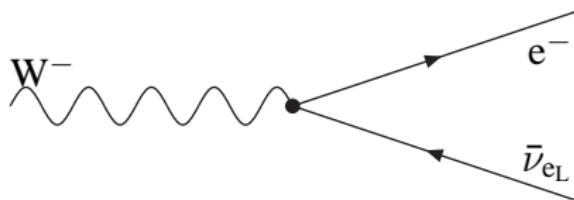
Particle Physics: The Standard Model

Dirk Zerwas

LAL
zerwas@lal.in2p3.fr

May 23, 2013

- Study properties of the W^\pm boson
- Study properties of the Z^0 boson
- Study properties of the t quark
- Theory+Experiment to “discover” the Higgs boson



- W^- polarization vector
- electron final state
- anti-neutrino final state
- EW vertex
- $\gamma^{0\dagger} = \gamma^0$ and $\gamma_\nu^\dagger = \gamma^0 \gamma_\nu \gamma^0$

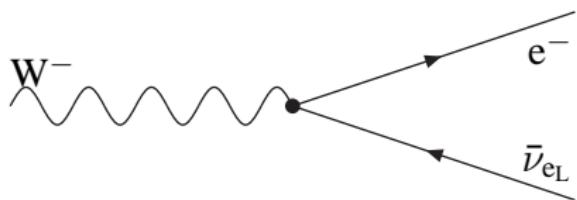
$$T_{fi} = \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma_\nu^\dagger i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^{0\dagger} u(\mathbf{k})$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma^0 \gamma_\nu \gamma^0 i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^0 u(\mathbf{k})$$



• W^- polarization vector

- electron final state
- anti-neutrino final state
- EW vertex
- $\gamma^{0\dagger} = \gamma^0$ and $\gamma_\nu^\dagger = \gamma^0 \gamma_\nu \gamma^0$

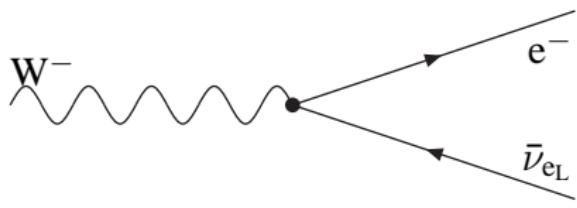
$$T_{fi} = \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma_\nu^\dagger i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^{0\dagger} u(\mathbf{k})$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma^0 \gamma_\nu \gamma^0 i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^0 u(\mathbf{k})$$



- W^- polarization vector
- electron final state
- anti-neutrino final state
- EW vertex
- $\gamma^{0\dagger} = \gamma^0$ and $\gamma_\nu^\dagger = \gamma^0 \gamma_\nu \gamma^0$

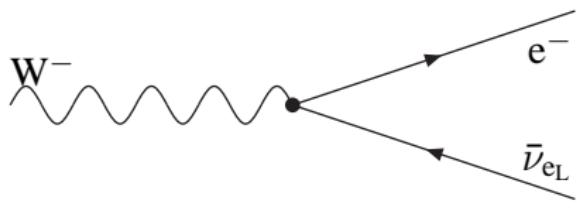
$$T_{fi} = \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma_\nu^\dagger i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^{0\dagger} u(\mathbf{k})$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma^0 \gamma_\nu \gamma^0 i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^0 u(\mathbf{k})$$



- W^- polarization vector
- electron final state
- anti-neutrino final state
- EW vertex
- $\gamma^{0\dagger} = \gamma^0$ and $\gamma_\nu^\dagger = \gamma^0 \gamma_\nu \gamma^0$

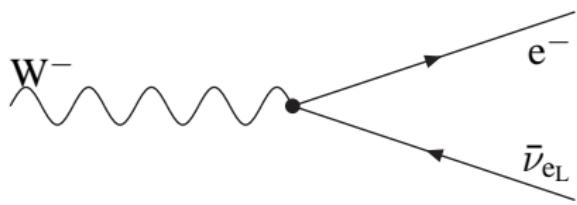
$$T_{fi} = \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma_\nu^\dagger i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^{0\dagger} u(\mathbf{k})$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma^0 \gamma_\nu \gamma^0 i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^0 u(\mathbf{k})$$



- W^- polarization vector
- electron final state
- anti-neutrino final state
- EW vertex
- $\gamma^{0\dagger} = \gamma^0$ and $\gamma_\nu^\dagger = \gamma^0 \gamma_\nu \gamma^0$

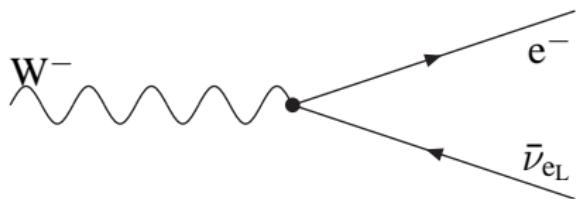
$$T_{fi} = \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma_\nu^\dagger i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^{0\dagger} u(\mathbf{k})$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma^0 \gamma_\nu \gamma^0 i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^0 u(\mathbf{k})$$



- W^- polarization vector
- electron final state
- anti-neutrino final state
- EW vertex
- $\gamma^{0\dagger} = \gamma^0$ and $\gamma_\nu^\dagger = \gamma^0 \gamma_\nu \gamma^0$

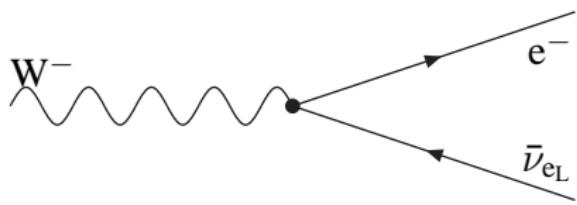
$$T_{fi} = \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma_\nu^\dagger i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^{0\dagger} u(\mathbf{k})$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma^0 \gamma_\nu \gamma^0 i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^0 u(\mathbf{k})$$



- W^- polarization vector
- electron final state
- anti-neutrino final state
- EW vertex
- $\gamma^{0\dagger} = \gamma^0$ and $\gamma_\nu^\dagger = \gamma^0 \gamma_\nu \gamma^0$

$$T_{fi} = \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma_\nu^\dagger i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^{0\dagger} u(\mathbf{k})$$

$$|\mathcal{M}|^2 = \sum' \epsilon^\mu \bar{u}(\mathbf{k}) \left(-i \frac{e}{\sqrt{2} \sin \theta_W} \gamma_\mu \frac{1-\gamma_5}{2} \right) v(\mathbf{k}')$$

$$\epsilon^{\nu*} v^\dagger(\mathbf{k}') \left(\frac{1-\gamma_5}{2} \right) \gamma^0 \gamma_\nu \gamma^0 i \frac{e}{\sqrt{2} \sin \theta_W} \gamma^0 u(\mathbf{k})$$

- common factors and \bar{v}
- initial state Spin average and polarisation vector relationship
- Spinor relationship
- $Tr(0\dots 3\gamma\gamma_5) = 0$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{e^2}{2\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu \frac{1-\gamma_5}{2} v(\mathbf{k}') v^\dagger(\mathbf{k}') \frac{1-\gamma_5}{2} \gamma^0 \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{8\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{3.8\sin^2\theta_W} \sum (-g^{\mu\nu}) \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{-e^2}{24\sin^2\theta_W} Tr(\not{k}\gamma_\mu (1 - \gamma_5)\not{k}'(1 + \gamma_5)\gamma^\mu) \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) - Tr(\not{k}\gamma_\mu \gamma_5 \not{k}'\gamma_5\gamma^\mu)] \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) + Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu)]
 \end{aligned}$$

● common factors and \bar{v}

- initial state Spin average and polarisation vector relationship
- Spinor relationship
- $Tr(0...3\gamma\gamma_5) = 0$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{e^2}{2\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu \frac{1-\gamma_5}{2} v(\mathbf{k}') v^\dagger(\mathbf{k}') \frac{1-\gamma_5}{2} \gamma^0 \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{8\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{3.8\sin^2\theta_W} \sum (-g^{\mu\nu}) \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{-e^2}{24\sin^2\theta_W} Tr(\not{k}\gamma_\mu (1 - \gamma_5)\not{k}'(1 + \gamma_5)\gamma^\mu) \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) - Tr(\not{k}\gamma_\mu \gamma_5 \not{k}'\gamma_5\gamma^\mu)] \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) + Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu)]
 \end{aligned}$$

- common factors and \bar{v}
- initial state Spin average and polarisation vector relationship
- Spinor relationship
- $Tr(0\dots 3\gamma\gamma_5) = 0$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{e^2}{2\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu \frac{1-\gamma_5}{2} v(\mathbf{k}') v^\dagger(\mathbf{k}') \frac{1-\gamma_5}{2} \gamma^0 \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{8\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{3.8\sin^2\theta_W} \sum (-g^{\mu\nu}) \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{-e^2}{24\sin^2\theta_W} Tr(\not{k}\gamma_\mu (1 - \gamma_5)\not{k}'(1 + \gamma_5)\gamma^\mu) \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) - Tr(\not{k}\gamma_\mu \gamma_5 \not{k}'\gamma_5\gamma^\mu)] \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) + Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu)]
 \end{aligned}$$

- common factors and \bar{v}
- initial state Spin average and polarisation vector relationship
- Spinor relationship
- $Tr(0...3\gamma\gamma_5) = 0$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{e^2}{2\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu \frac{1-\gamma_5}{2} v(\mathbf{k}') v^\dagger(\mathbf{k}') \frac{1-\gamma_5}{2} \gamma^0 \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{8\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{3.8\sin^2\theta_W} \sum (-g^{\mu\nu}) \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{-e^2}{24\sin^2\theta_W} Tr(\not{k}\gamma_\mu (1 - \gamma_5) \not{k}' (1 + \gamma_5) \gamma^\mu) \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}' \gamma^\mu) - Tr(\not{k}\gamma_\mu \gamma_5 \not{k}' \gamma_5 \gamma^\mu)] \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}' \gamma^\mu) + Tr(\not{k}\gamma_\mu \not{k}' \gamma^\mu)]
 \end{aligned}$$

- common factors and \bar{v}
- initial state Spin average and polarisation vector relationship
- Spinor relationship
- $Tr(0\dots 3\gamma\gamma_5) = 0$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{e^2}{2\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu \frac{1-\gamma_5}{2} v(\mathbf{k}') v^\dagger(\mathbf{k}') \frac{1-\gamma_5}{2} \gamma^0 \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{8\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{3.8\sin^2\theta_W} \sum (-g^{\mu\nu}) \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{-e^2}{24\sin^2\theta_W} Tr(\not{k}\gamma_\mu (1 - \gamma_5) \not{k}' (1 + \gamma_5) \gamma^\mu) \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}' \gamma^\mu) - Tr(\not{k}\gamma_\mu \gamma_5 \not{k}' \gamma_5 \gamma^\mu)] \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}' \gamma^\mu) + Tr(\not{k}\gamma_\mu \not{k}' \gamma^\mu)]
 \end{aligned}$$

- common factors and \bar{v}
- initial state Spin average and polarisation vector relationship
- Spinor relationship
- $Tr(0...3\gamma\gamma_5) = 0$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{e^2}{2\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu \frac{1-\gamma_5}{2} v(\mathbf{k}') v^\dagger(\mathbf{k}') \frac{1-\gamma_5}{2} \gamma^0 \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{8\sin^2\theta_W} \sum' \epsilon^\mu \epsilon^{\nu*} \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{e^2}{3.8\sin^2\theta_W} \sum (-g^{\mu\nu}) \bar{u}(\mathbf{k}) \gamma_\mu (1 - \gamma_5) v(\mathbf{k}') \bar{v}(\mathbf{k}') (1 + \gamma_5) \gamma_\nu u(\mathbf{k}) \\
 &= \frac{-e^2}{24\sin^2\theta_W} Tr(\not{k}\gamma_\mu (1 - \gamma_5)\not{k}'(1 + \gamma_5)\gamma^\mu) \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) - Tr(\not{k}\gamma_\mu \gamma_5 \not{k}'\gamma_5\gamma^\mu)] \\
 &= \frac{-e^2}{24\sin^2\theta_W} [Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu) + Tr(\not{k}\gamma_\mu \not{k}'\gamma^\mu)]
 \end{aligned}$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\mathbf{k}' \approx 2\mathbf{k}\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\mathbf{k} \gamma_\mu \mathbf{k}' \gamma^\mu)] \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\
 &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\
 &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2
 \end{aligned}$$

$$\begin{aligned}
 d\Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega \\
 \Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}^2} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi
 \end{aligned}$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\cdot\mathbf{k}' \approx 2\mathbf{k}\cdot\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\not{k} \gamma_\mu \not{k}' \gamma^\mu)] \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\
 &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\
 &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2
 \end{aligned}$$

$$\begin{aligned}
 d\Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega \\
 \Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}^2} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi
 \end{aligned}$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\mathbf{k}' \approx 2\mathbf{k}\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\mathbf{k} \gamma_\mu \mathbf{k}' \gamma^\mu)] \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\
 &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\
 &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2
 \end{aligned}$$

$$\begin{aligned}
 d\Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega \\
 \Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}^2} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi
 \end{aligned}$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\mathbf{k}' \approx 2\mathbf{k}\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\mathbf{k} \gamma_\mu \mathbf{k}' \gamma^\mu)] \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\
 &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\
 &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2
 \end{aligned}$$

$$\begin{aligned}
 d\Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega \\
 \Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi
 \end{aligned}$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\mathbf{k}' \approx 2\mathbf{k}\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\mathbf{k} \gamma_\mu \mathbf{k}' \gamma^\mu)] \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\
 &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\
 &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2
 \end{aligned}$$

$$\begin{aligned}
 d\Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega \\
 \Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}^2} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi
 \end{aligned}$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\mathbf{k}' \approx 2\mathbf{k}\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned} |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\mathbf{k} \gamma_\mu \mathbf{k}' \gamma^\mu)] \\ &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\ &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\ &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\ &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 \end{aligned}$$

$$\begin{aligned} d\Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega \\ \Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}^2} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi \end{aligned}$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\mathbf{k}' \approx 2\mathbf{k}\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\mathbf{k} \gamma_\mu \mathbf{k}' \gamma^\mu)] \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\
 &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\
 &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2
 \end{aligned}$$

$$d\Gamma_{e\bar{\nu}_e L} = \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega$$

$$\Gamma_{e\bar{\nu}_e L} = \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi$$

- $Tr(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma) = 4(g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma})$
- $m_{W^\pm}^2 = s = (\mathbf{k} + \mathbf{k}')^2 = \mathbf{k}^2 + \mathbf{k}'^2 + 2\mathbf{k}\mathbf{k}' \approx 2\mathbf{k}\mathbf{k}'$
- ME isotropical (as should be for unpolarized decay!)

$$\begin{aligned}
 |\mathcal{M}|^2 &= \frac{-e^2}{12 \sin^2 \theta_W} [Tr(\mathbf{k} \gamma_\mu \mathbf{k}' \gamma^\mu)] \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} k^\mu k'^\rho g^{\nu\sigma} (g_{\mu\nu}g_{\rho\sigma} + g_{\mu\sigma}g_{\nu\rho} - g_{\mu\rho}g_{\nu\sigma}) \\
 &= \frac{-e^2}{3 \sin^2 \theta_W} (\mathbf{k} \cdot \mathbf{k}' + \mathbf{k} \cdot \mathbf{k}' - \mathbf{k} \cdot \mathbf{k}' \cdot 4) \\
 &= \frac{2e^2}{3 \sin^2 \theta_W} \mathbf{k} \cdot \mathbf{k}' \\
 &= \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2
 \end{aligned}$$

$$\begin{aligned}
 d\Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{|\mathcal{M}|^2}{\sqrt{s}} d\Omega \\
 \Gamma_{e\bar{\nu}_e L} &= \frac{1}{64\pi^2} \frac{1}{m_{W^\pm}^2} \frac{e^2}{3 \sin^2 \theta_W} m_{W^\pm}^2 4\pi
 \end{aligned}$$

W^\pm partial widths

$$\Gamma_{e\bar{\nu}_e L} = \frac{1}{16\pi} \frac{e^2}{3\sin^2\theta_W} m_{W^\pm}$$

$$\frac{e^2}{\sin^2\theta_W} = 8 \frac{G}{\sqrt{2}} m_{W^\pm}^2$$

$$\Gamma_{e\bar{\nu}_e L} = G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3$$

$$\begin{aligned}\Gamma_{\bar{u}d} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 3 \cdot |\mathbf{V}_{ud}|^2 \\ &\quad (1 + \frac{\alpha_s(m_{W^\pm}^2)}{\pi})\end{aligned}$$

To include the radiative corrections means inclusive:
 $\Gamma_{\bar{u}d} = \Gamma_{\bar{u}d}(g)$

Branching ratios

$$\Gamma = \Gamma_{e\nu_e L} + \Gamma_{\mu\nu_\mu L} + \Gamma_{\tau\nu_\tau L}$$

$$\Gamma_{ud} + \Gamma_{cs}$$

$$\approx (3 + 2 \cdot 3) \Gamma_{e\nu}$$

$$\mathcal{B}(\ell\nu) \approx 30\%$$

- tb final state excluded by mass $180\text{GeV} > 80\text{GeV}$

- V_{ij} for $i \neq j$ excluded
 $V_{ij} \ll 1$

W^\pm partial widths

$$\Gamma_{e\bar{\nu}_e L} = \frac{1}{16\pi} \frac{e^2}{3\sin^2\theta_W} m_{W^\pm}$$

$$\frac{e^2}{\sin^2\theta_W} = \frac{8G}{\sqrt{2}} m_{W^\pm}^2$$

$$\Gamma_{e\bar{\nu}_e L} = G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3$$

$$\begin{aligned}\Gamma_{\bar{u}d} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \cdot 3 \cdot |\mathbf{V}_{ud}|^2 \\ &\quad (1 + \frac{\alpha_s(m_{W^\pm}^2)}{\pi})\end{aligned}$$

To include the radiative corrections means inclusive:
 $\Gamma_{\bar{u}d} = \Gamma_{\bar{u}d}(g)$

Branching ratios

$$\Gamma = \Gamma_{e\nu_e L} + \Gamma_{\mu\nu_\mu L} + \Gamma_{\tau\nu_\tau L}$$

$$\Gamma_{ud} + \Gamma_{cs}$$

$$\approx (3 + 2 \cdot 3) \Gamma_{e\nu}$$

$$\mathcal{B}(\ell\nu) \approx 30\%$$

- tb final state excluded by mass $180\text{GeV} > 80\text{GeV}$

- V_{ij} for $i \neq j$ excluded
 $V_{ij} \ll 1$

W^\pm partial widths

$$\begin{aligned}\Gamma_{e\bar{\nu}_e L} &= \frac{1}{16\pi} \frac{e^2}{3\sin^2\theta_W} m_{W^\pm} \\ \frac{e^2}{\sin^2\theta_W} &= 8 \frac{G}{\sqrt{2}} m_{W^\pm}^2 \\ \Gamma_{e\bar{\nu}_e L} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \\ \Gamma_{\bar{u}d} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \color{red}{3 \cdot |\mathcal{V}_{ud}|^2} \\ &\quad \color{red}{(1 + \frac{\alpha_s(m_{W^\pm}^2)}{\pi})}\end{aligned}$$

To include the radiative corrections means inclusive:
 $\Gamma_{\bar{u}d} = \Gamma_{\bar{u}d}(g)$

Branching ratios

$$\begin{aligned}\Gamma &= \Gamma_{e\nu_e L} + \Gamma_{\mu\nu_\mu L} + \Gamma_{\tau\nu_\tau L} \\ &\quad \Gamma_{ud} + \Gamma_{cs}\end{aligned}$$

$$\approx (3 + 2 \cdot 3) \Gamma_{e\nu}$$

$$\mathcal{B}(\ell\nu) \approx 30\%$$

- tb final state excluded by mass $180\text{GeV} > 80\text{GeV}$
- V_{ij} for $i \neq j$ excluded
 $V_{ij} \ll 1$

W^\pm partial widths

$$\begin{aligned}\Gamma_{e\bar{\nu}_e L} &= \frac{1}{16\pi} \frac{e^2}{3\sin^2\theta_W} m_{W^\pm} \\ \frac{e^2}{\sin^2\theta_W} &= 8 \frac{G}{\sqrt{2}} m_{W^\pm}^2 \\ \Gamma_{e\bar{\nu}_e L} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \\ \Gamma_{\bar{u}d} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \color{red}{3 \cdot |\mathcal{V}_{ud}|^2} \\ &\quad \color{red}{(1 + \frac{\alpha_s(m_{W^\pm}^2)}{\pi})}\end{aligned}$$

To include the radiative corrections means inclusive:
 $\Gamma_{\bar{u}d} = \Gamma_{\bar{u}d}(g)$

Branching ratios

$$\begin{aligned}\Gamma &= \Gamma_{e\nu_e L} + \Gamma_{\mu\nu_\mu L} + \Gamma_{\tau\nu_\tau L} \\ &\quad \Gamma_{ud} + \Gamma_{cs} \\ &\approx (3 + 2 \cdot 3) \Gamma_{e\nu} \\ \mathcal{B}(\ell\nu) &\approx 30\%\end{aligned}$$

- tb final state excluded by mass $180\text{GeV} \gg 80\text{GeV}$
- V_{ij} for $i \neq j$ excluded
 $V_{ij} \ll 1$

W^\pm partial widths

$$\begin{aligned}\Gamma_{e\bar{\nu}_e L} &= \frac{1}{16\pi} \frac{e^2}{3\sin^2\theta_W} m_{W^\pm} \\ \frac{e^2}{\sin^2\theta_W} &= 8 \frac{G}{\sqrt{2}} m_{W^\pm}^2 \\ \Gamma_{e\bar{\nu}_e L} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \\ \Gamma_{\bar{u}d} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \color{red}{3 \cdot |\mathcal{V}_{ud}|^2} \\ &\quad \color{red}{(1 + \frac{\alpha_s(m_{W^\pm}^2)}{\pi})}\end{aligned}$$

To include the radiative corrections means inclusive:
 $\Gamma_{\bar{u}d} = \Gamma_{\bar{u}d}(g)$

Branching ratios

$$\begin{aligned}\Gamma &= \Gamma_{e\nu_e L} + \Gamma_{\mu\nu_\mu L} + \Gamma_{\tau\nu_\tau L} \\ &\quad \Gamma_{ud} + \Gamma_{cs} \\ &\approx (3 + 2 \cdot 3) \Gamma_{e\nu} \\ \mathcal{B}(\ell\nu) &\approx 30\%\end{aligned}$$

- tb final state excluded by mass $180\text{GeV} \gg 80\text{GeV}$
- V_{ij} for $i \neq j$ excluded
 $V_{ij} \ll 1$

W^\pm partial widths

$$\begin{aligned}\Gamma_{e\bar{\nu}_e L} &= \frac{1}{16\pi} \frac{e^2}{3\sin^2\theta_W} m_{W^\pm} \\ \frac{e^2}{\sin^2\theta_W} &= 8 \frac{G}{\sqrt{2}} m_{W^\pm}^2 \\ \Gamma_{e\bar{\nu}_e L} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \\ \Gamma_{\bar{u}d} &= G \frac{1}{6\pi\sqrt{2}} m_{W^\pm}^3 \color{red}{3 \cdot |\mathcal{V}_{ud}|^2} \\ &\quad \color{red}{(1 + \frac{\alpha_s(m_{W^\pm}^2)}{\pi})}\end{aligned}$$

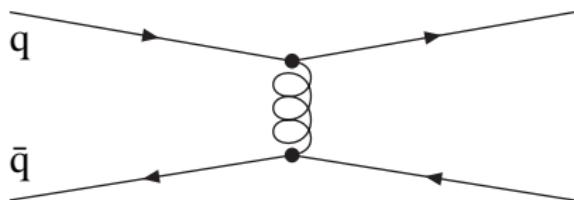
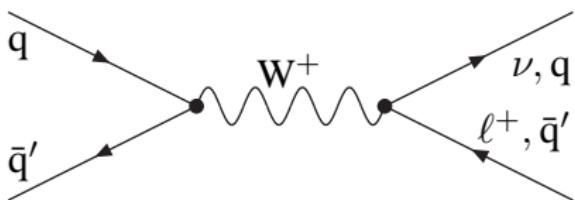
To include the radiative corrections means inclusive:
 $\Gamma_{\bar{u}d} = \Gamma_{\bar{u}d}(g)$

Branching ratios

$$\begin{aligned}\Gamma &= \Gamma_{e\nu_e L} + \Gamma_{\mu\nu_\mu L} + \Gamma_{\tau\nu_\tau L} \\ &\quad \Gamma_{ud} + \Gamma_{cs} \\ &\approx (3 + 2 \cdot 3) \Gamma_{e\nu} \\ \mathcal{B}(\ell\nu) &\approx 30\%\end{aligned}$$

- tb final state excluded by mass $180\text{GeV} \gg 80\text{GeV}$
- V_{ij} for $i \neq j$ excluded
 $V_{ij} \ll 1$

Hadron colliders



- leptons or jets
- each vertex G_F
- LHC: sea \bar{q}
- LHC NLO: gluon (splitting)
- $\sqrt{\hat{s}} = m_W^\pm$
- $\sqrt{x_1 \cdot x_2} \cdot \sqrt{s} = m_W^\pm$
- $\sqrt{s} \approx 3 \cdot 2m_W^\pm$ (gluon 50%)

- no leptons :)
- each vertex g_S
- jet-jet mass only difference to signal
- LHC: 10^4 lepton-jet rejection

The discovery machine

- SPPS transformed to SP \bar{P} S
- $\sqrt{s} \approx 600$ GeV
- stochastic cooling

TeVatron

- p \bar{p}
- $\sqrt{s} \leq 1.96$ TeV
- $\approx 10\text{fb}^{-1}$
- fb cross sections measureable

LHC

- pp
- $\sqrt{s} \leq 14$ TeV
- luminosity of anti-protons
- gluon PDFs dominate
- QCD cross sections increase more rapidly than the signal
- 2011: 7 TeV $\approx 5\text{fb}^{-1}$
- 2012: 8 TeV $\approx 25\text{fb}^{-1}$
- 2015: 14 TeV-X

The discovery machine

- SPPS transformed to SP \bar{P} S
- $\sqrt{s} \approx 600$ GeV
- stochastic cooling

TeVatron

- p \bar{p}
- $\sqrt{s} \leq 1.96$ TeV
- $\approx 10\text{fb}^{-1}$
- fb cross sections measureable

LHC

- pp
- $\sqrt{s} \leq 14$ TeV
- luminosity of anti-protons
- gluon PDFs dominate
- QCD cross sections increase more rapidly than the signal
- 2011: 7 TeV $\approx 5\text{fb}^{-1}$
- 2012: 8 TeV $\approx 25\text{fb}^{-1}$
- 2015: 14 TeV-X

The discovery machine

- SPPS transformed to SP \bar{P} S
- $\sqrt{s} \approx 600$ GeV
- stochastic cooling

TeVatron

- p \bar{p}
- $\sqrt{s} \leq 1.96$ TeV
- $\approx 10\text{fb}^{-1}$
- fb cross sections measureable

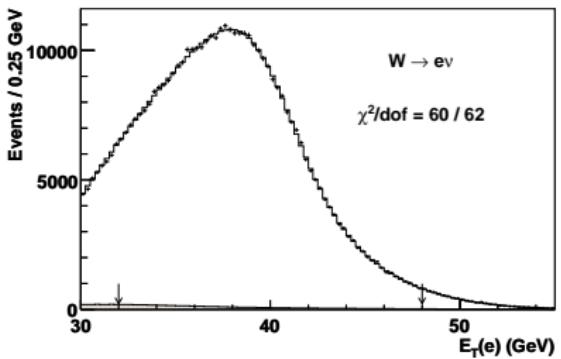
LHC

- pp
- $\sqrt{s} \leq 14$ TeV
- luminosity of anti-protons
- gluon PDFs dominate
- QCD cross sections increase more rapidly than the signal
- 2011: 7 TeV $\approx 5\text{fb}^{-1}$
- 2012: 8 TeV $\approx 25\text{fb}^{-1}$
- 2015: 14 TeV-X

Transverse Energy

$$E_T \leq \frac{m_W^\pm}{2}$$

- Problem Solving
- 2-body decay
- particles massless
- maximal energy half of the mother mass



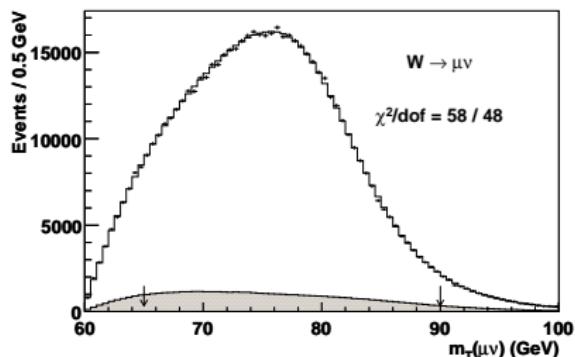
Experimental issues

- finite width
- final state radiation
- detector effects
- ISR

Transverse Mass

$$m_T = m(\cancel{E}_T, \mathbf{p}_\ell)$$

- \cancel{E}_T negative sum of activity
- ignore longitudinal component
- maximal mass is m_W^\pm
- less sensitive to NLO

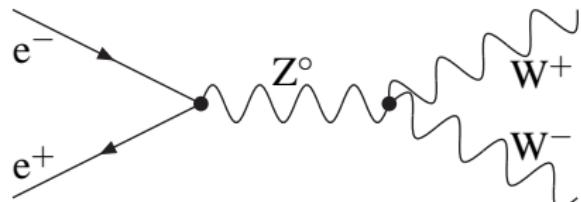
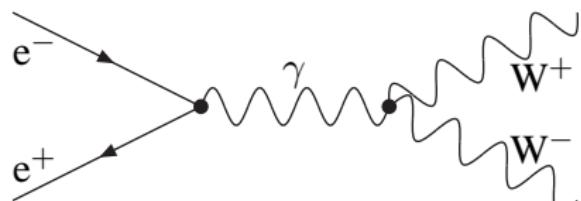
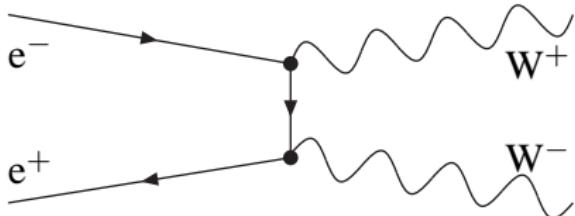


Experimental Results

CDF	80.413	\pm	0.048	GeV
D0	80.401	\pm	0.043	GeV

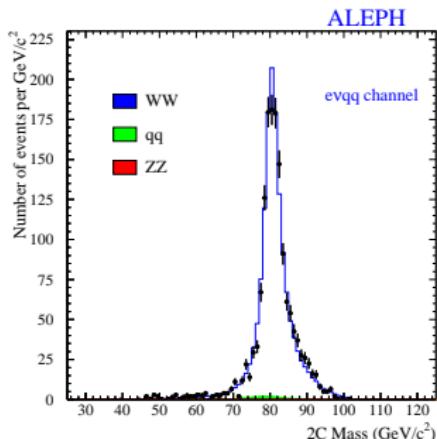
Lepton Colliders

- CM system is lab system
- $\sqrt{\hat{s}} = \sqrt{s}$ (modulo ISR)
- initial state charge zero
- pair production



Signatures

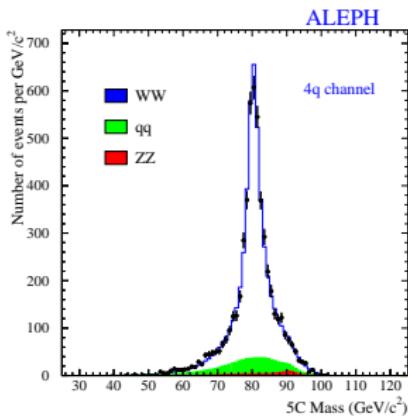
- $W^+W^- \rightarrow q\bar{q}q\bar{q}$ 50%
- $W^+W^- \rightarrow \ell\nu q\bar{q}$ 40%
- $W^+W^- \rightarrow \ell\nu\ell\nu$ 10%
- 90% useful with $\nu = E$



Semi-leptonic

- low background
- use separately leptonic decay and hadronic decay





Fully hadronic

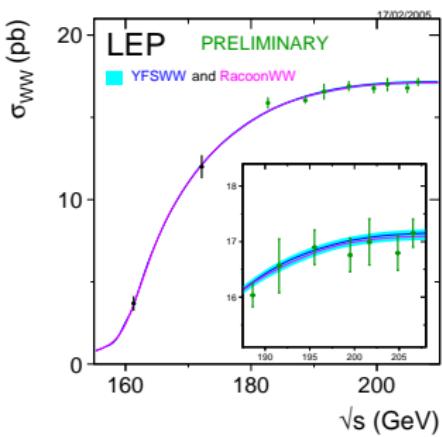
- QCD background
- jet reconstruction and calibration
- Color reconnection (jet algo)

Experimental Results

ALEPH	80.440	\pm	0.051 GeV
DELPHI	80.336	\pm	0.067 GeV
L3	80.270	\pm	0.055 GeV
OPAL	80.415	\pm	0.052 GeV
PDG	80.399	\pm	0.023 GeV

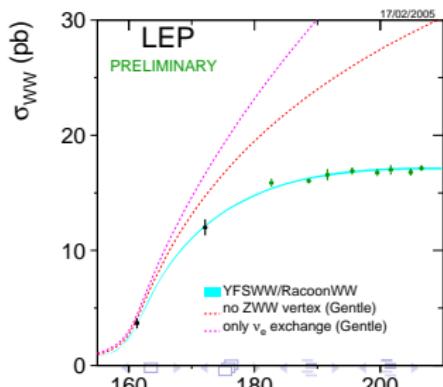
WW Cross section

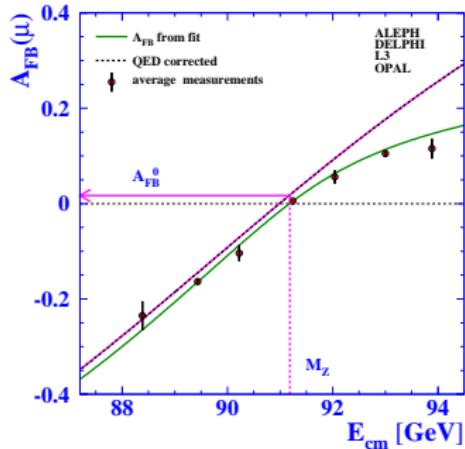
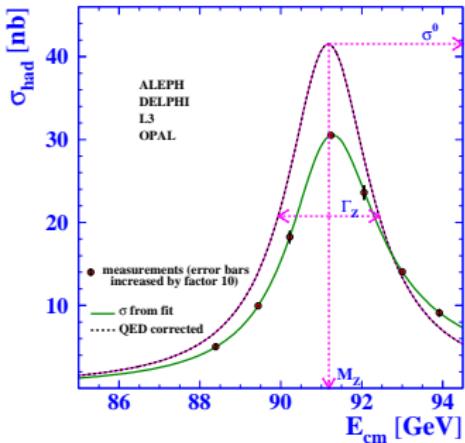
- threshold scan sensitive to W mass
- width washes out threshold



WW Interpretation

- t channel alone: insufficient
- adding $\gamma W^+ W^-$ still insufficient
- unitarity not proven



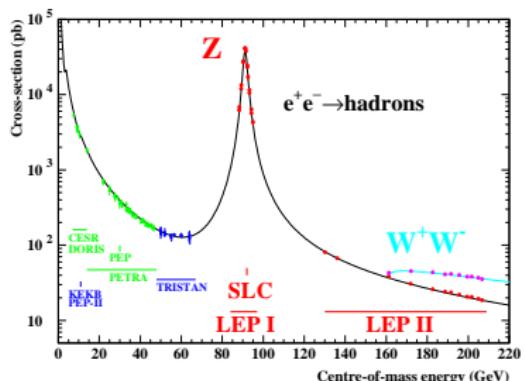
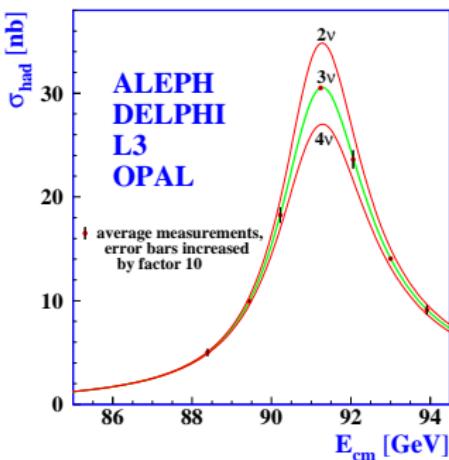


Experimental Issues

- measure beam energy precisely
- measure efficiency precisely

Predictions/Measurements

- Radiative corrections important
- Interference zero on peak



Results

- 3 generations
- 91.1875 ± 0.0021 GeV
- 2.4952 ± 0.0023 GeV

Results

Impressive agreement over

- decades in time
- different machines
- different energies

Forward/backward asymmetry

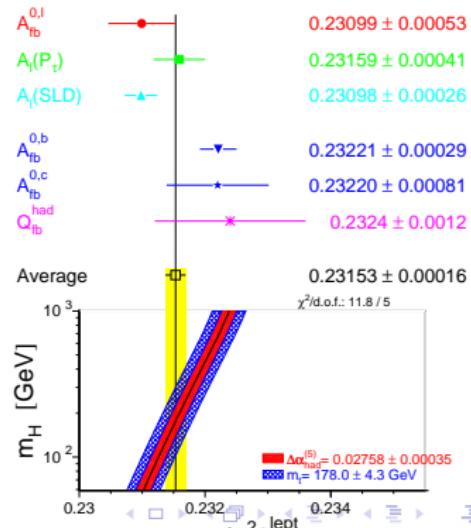
$$A_{FB}(f) = \frac{\int_0^1 \frac{d\sigma}{d \cos \theta} d \cos \theta - \int_{-1}^0 \frac{d\sigma}{d \cos \theta} d \cos \theta}{\int_0^1 \frac{d\sigma}{d \cos \theta} d \cos \theta + \int_{-1}^0 \frac{d\sigma}{d \cos \theta} d \cos \theta}$$

Measurements

- A_{FB} is a measurement of $\sin^2 \theta_W$
- leptons easy
- hadrons hard
- leptons with polarization

Open question

Do leptons behave differently than hadrons?

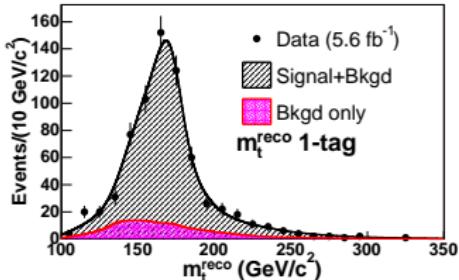
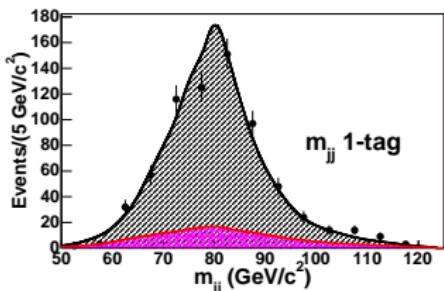


Experimental issues

- jet reconstruction
- jet calibration
- b tagging
- W mass constraint
- all final states used

Experimental result

PDG $172.9 \pm 0.6 \pm 0.9$ GeV



Perturbativity and Unitarity

- unitarity of WW scattering (see WW)
- width smaller than mass
- upper limit on Higgs boson mass

Triviality

RGE solution Quartic coupling (only Higgs sector for large λ):

$$\lambda(Q^2) = \lambda(v^2) \left[1 - \frac{3}{4\pi^2} \lambda(v^2) \log \frac{Q^2}{v^2} \right]^{-1}$$

- $Q \gg v \rightarrow \lambda = \infty$ Landau pole
- $Q = 0 \rightarrow \lambda = 0$ (trivial)
- ϕ^4 must be trivial if valid at all scales
- navigate between the poles

Perturbativity and Unitarity

- unitarity of WW scattering (see WW)
- width smaller than mass
- upper limit on Higgs boson mass

Triviality

RGE solution Quartic coupling (only Higgs sector for large λ):

$$\lambda(Q^2) = \lambda(v^2) \left[1 - \frac{3}{4\pi^2} \lambda(v^2) \log \frac{Q^2}{v^2} \right]^{-1}$$

- $Q \gg v \rightarrow \lambda = \infty$ Landau pole
- $Q = 0 \rightarrow \lambda = 0$ (trivial)
- ϕ^4 must be trivial if valid at all scales
- navigate between the poles

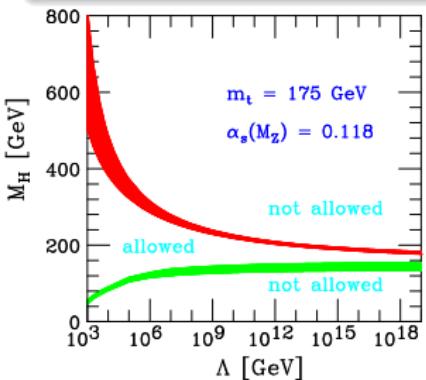


Stability

RGE solution including fermions (small λ):

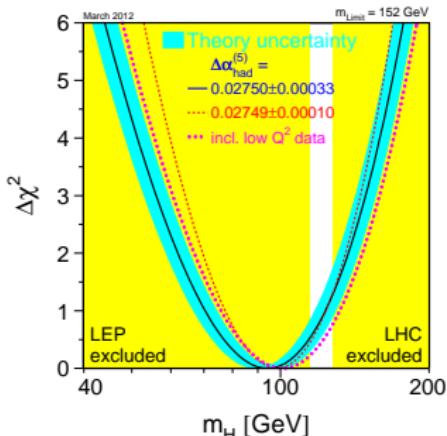
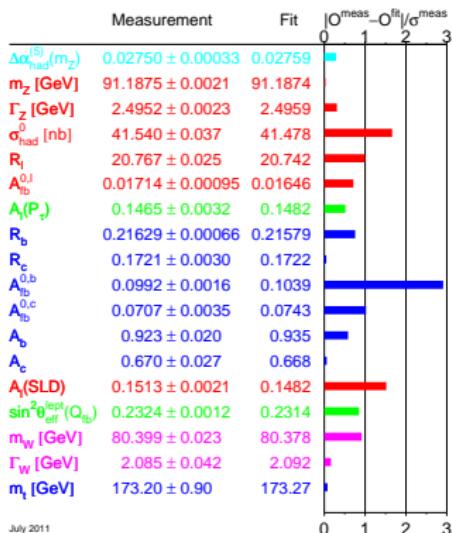
$$\lambda(Q^2) = \lambda(v^2) \frac{1}{16\pi^2} \left[+ \frac{3}{16} \left(-12 \frac{m_t^4}{v^4} + 2g_2^4 + (g_2^2 + g_1^2)^2 \right) \right] \log \frac{Q^2}{v^2}$$

- λ small \rightarrow top yukawa can turn $\lambda < 0$
- Higgs potential unbounded



Consequences

- Λ : cut-off scale for new physics
- $m_H \leq 700 \text{ GeV}$
 - window for SM valid up to GUT scale at 125GeV



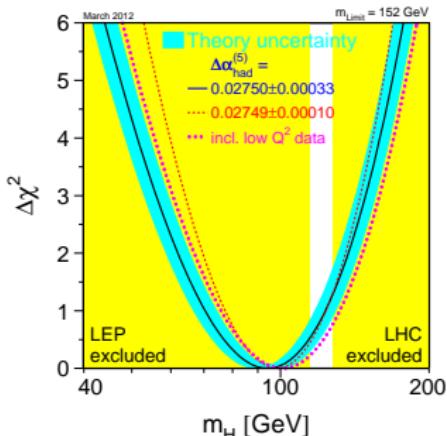
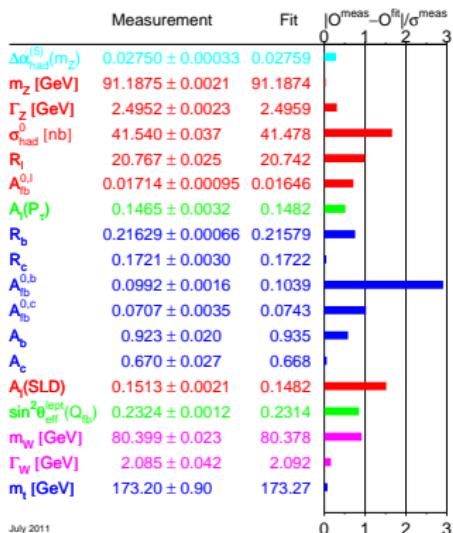
EW measurements

- internally compatible
- largest deviation:
asymmetry

Indirect Higgs Search

logarithmic dependence on
Higgs boson mass:

$$m_H = 94^{+29}_{-24} \text{ GeV}$$



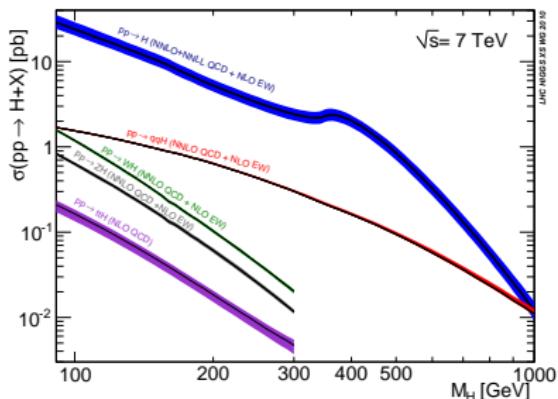
EW measurements

- internally compatible
- largest deviation:
asymmetry

Indirect Higgs Search

logarithmic dependence on
Higgs boson mass:

$$m_H = 94^{+29}_{-24} \text{ GeV}$$

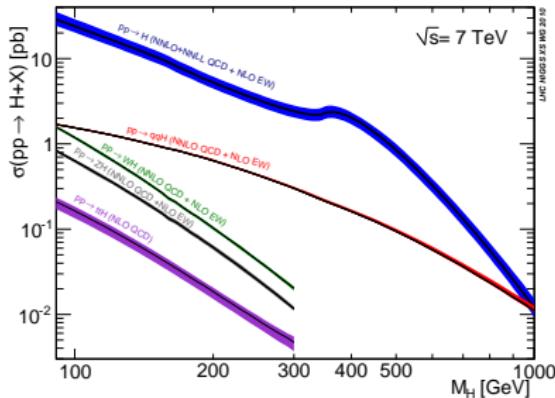


Higgs boson production

- gluon fusion (NLO > 2)
- VBF ($\mathcal{O}(20\%)$)
- associated production
- radiation off heavy quark

Experimental issues

- huge QCD background for jets
- mass information?
- mass resolution?
- total cross section
 $\sim \text{mb} = 10^{-3} \text{b}$
- signal $\sim \text{pb} = 10^{-12} \text{b}$
- pile-up

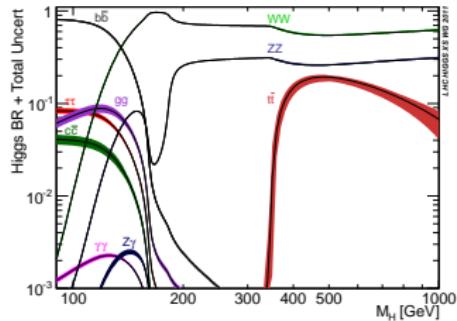
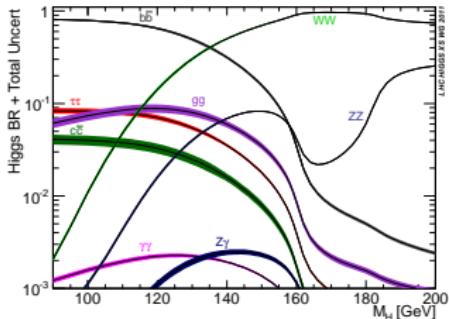


Higgs boson production

- gluon fusion (NLO > 2)
- VBF ($\mathcal{O}(20\%)$)
- associated production
- radiation off heavy quark

Experimental issues

- huge QCD background for jets
- mass information?
- mass resolution?
- total cross section
 $\sim \text{mb} = 10^{-3} \text{b}$
- signal $\sim \text{pb} = 10^{-12} \text{b}$
- pile-up

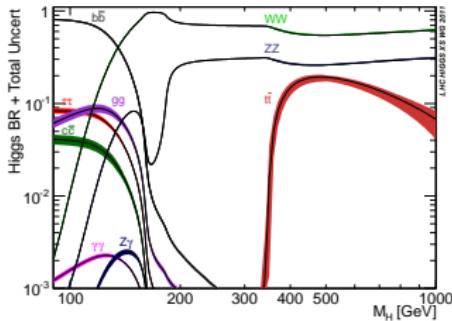
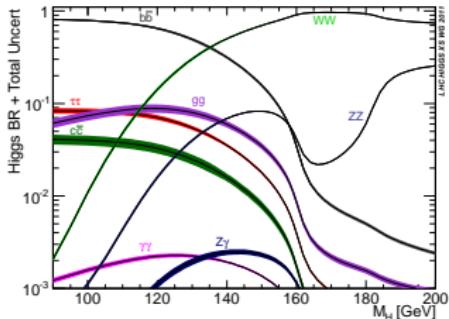


Low mass Higgs boson

- heaviest fermion dominates: b
- gauge boson pairs (one off-shell) kick in
- thresholds visible

High mass Higgs boson

- gauge bosons dominate
- top threshold 350 GeV
- \sqrt{s} dependence in gauge bosons



Low mass Higgs boson

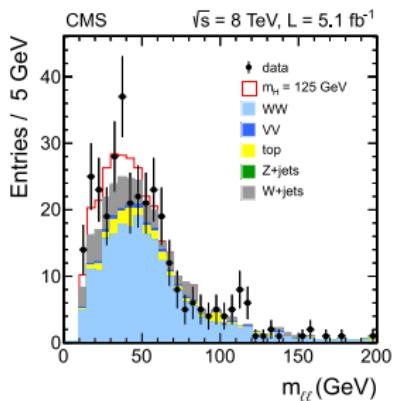
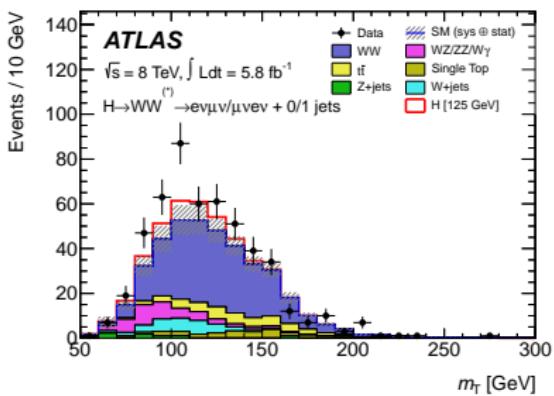
- heaviest fermion dominates: b
- gauge boson pairs (one off-shell) kick in
- thresholds visible

High mass Higgs boson

- gauge bosons dominate
- top threshold 350 GeV
- \sqrt{s} dependence in gauge bosons

Higgs to WW*

- WW EW production
- no mass peak

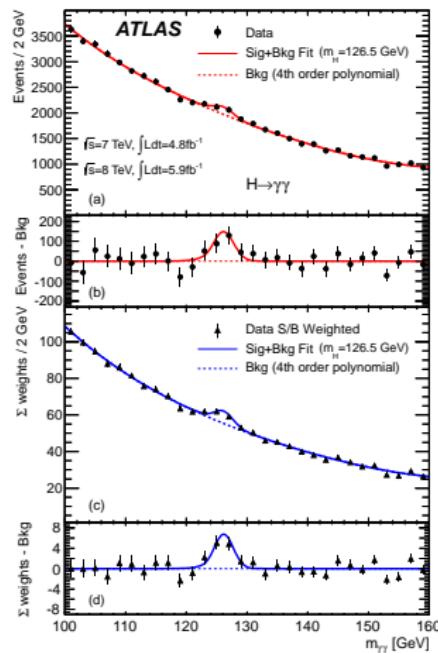


Using the spin correlation

- Higgs Spin-0
- W opposite
- charged leptons emitted in the same direction

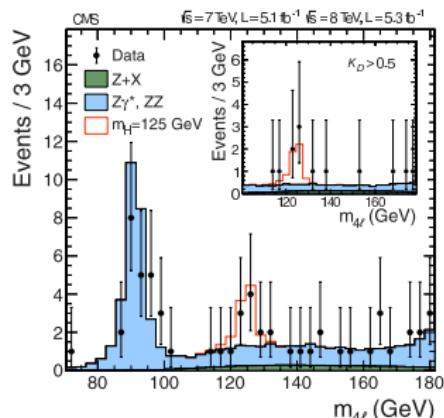
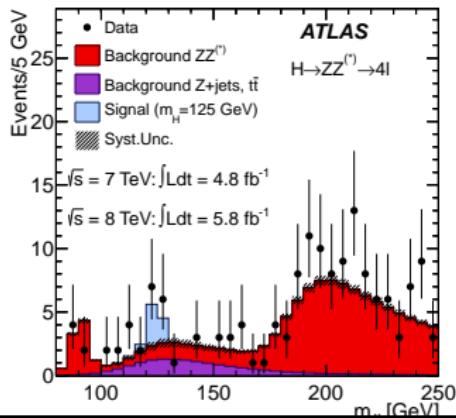
Higgs to $\gamma\gamma$

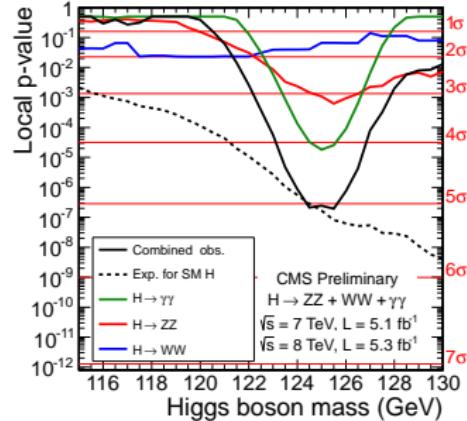
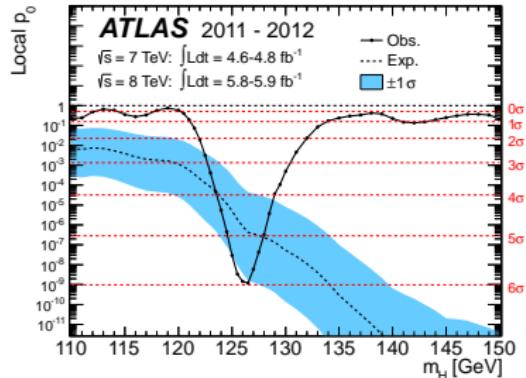
- Rare decay (via triangle) as γ is massless 10^{-3}
- SM background from continuum di-photon production (EM couplings)
- excellent mass resolution necessary



Higgs to ZZ^*

- ZZ EW production
- good mass resolution required
- e, μ : low \mathcal{B}





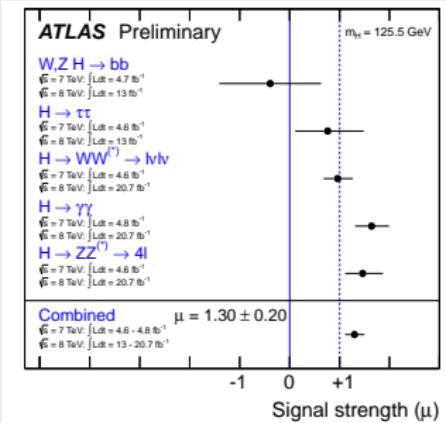
A new resonance

- discovery convention: 5σ
- independent discovery in two experiments

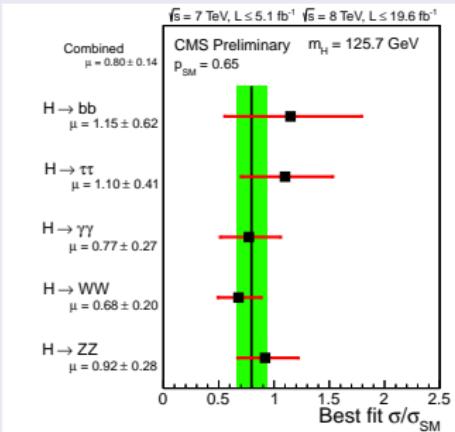
A success of particle physics

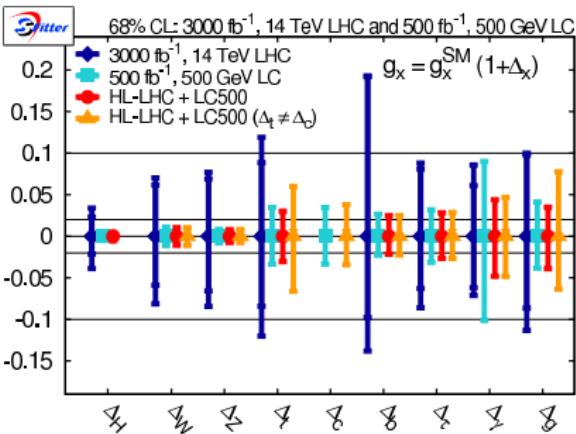
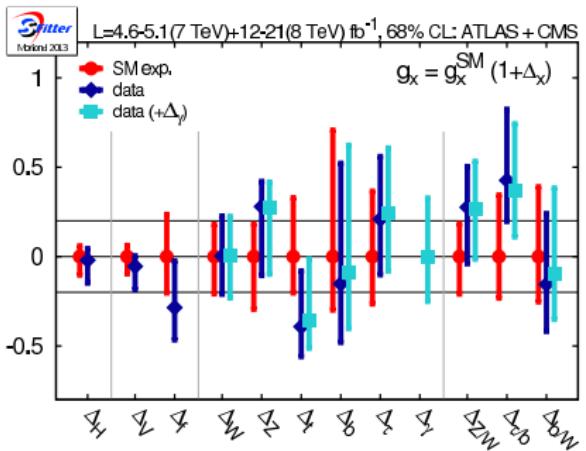
- discovery 50 years after the prediction

Signal Strengths ATLAS (Moriond 2013)



Signal Strengths CMS (Moriond 2013)



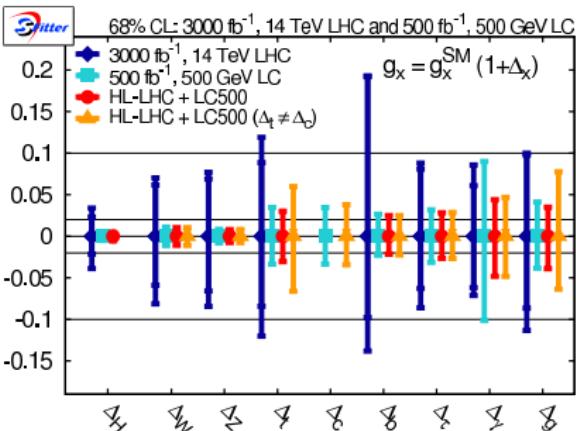
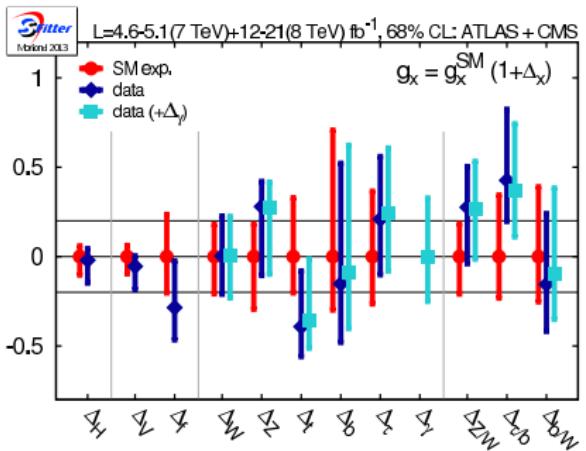


Higgs couplings

- global sensitivity 15%
- smallest errors for gauge boson couplings

Higgs couplings future

- luminosity 3000fb^{-1}
- global sensitivity 5%
- precision of the order 10-20%



Higgs couplings

- global sensitivity 15%
- smallest errors for gauge boson couplings

Higgs couplings future

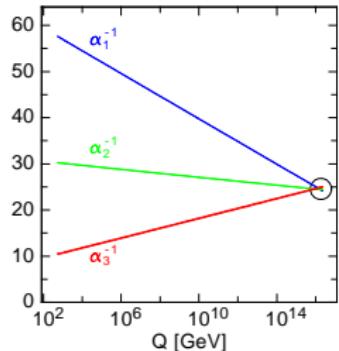
- luminosity 3000fb^{-1}
- global sensitivity 5%
- precision of the order 10-20%

Gauge couplings

- Planck scale (gravity) 10^{18}GeV
- unification of g_1, g_2, g_3 at 10^{16}GeV ?
- Standard Model: close miss

Supersymmetry

- fermionic degree of freedom has a bosonic counter part
- unification g_1, g_2, g_3 at 10^{16}GeV
- light $< 140\text{GeV}$ Higgs boson
- stabilizes the Higgs boson mass
- dark matter candidate



Experimental evidence

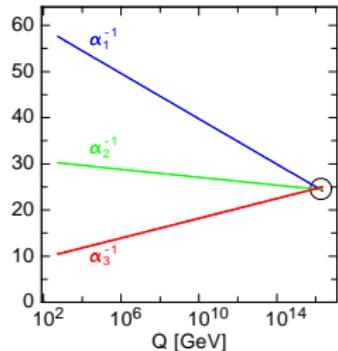
- half discovered?
- or nada?

Gauge couplings

- Planck scale (gravity) 10^{18}GeV
- unification of g_1, g_2, g_3 at 10^{16}GeV ?
- Standard Model: close miss

Supersymmetry

- fermionic degree of freedom has a bosonic counter part
- unification g_1, g_2, g_3 at 10^{16}GeV
- light $< 140\text{GeV}$ Higgs boson
- stabilizes the Higgs boson mass
- dark matter candidate



Experimental evidence

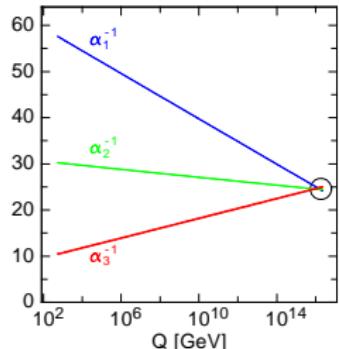
- half discovered?
- or nada?

Gauge couplings

- Planck scale (gravity) 10^{18}GeV
- unification of g_1, g_2, g_3 at 10^{16}GeV ?
- Standard Model: close miss

Supersymmetry

- fermionic degree of freedom has a bosonic counter part
- unification g_1, g_2, g_3 at 10^{16}GeV
- light $< 140\text{GeV}$ Higgs boson
- stabilizes the Higgs boson mass
- dark matter candidate



Experimental evidence

- half discovered?
- or nada?

Future

- The Standard Model is alive and kicking
- The Higgs boson be discovered:
 - What are the couplings of the Higgs boson?
 - What is the self-coupling of the Higgs boson?
 - Is it a Spin-0 particle?
- Neutrinos
 - Mixing in the leptonic sector
 - Is the Neutrino Majorana or DIRAC?
- Are there new physics beyond the Standard Model?
 - Supersymmetry
 - Extra Dimensions
 - ...
- What happens up to the Planck scale at 10^{18} GeV ?
 - Desert from EW to GUT scale?
 - Unification of the 3 forces at the GUT scale?

