

# Accurate predictions for Higgs pair production at the LHC



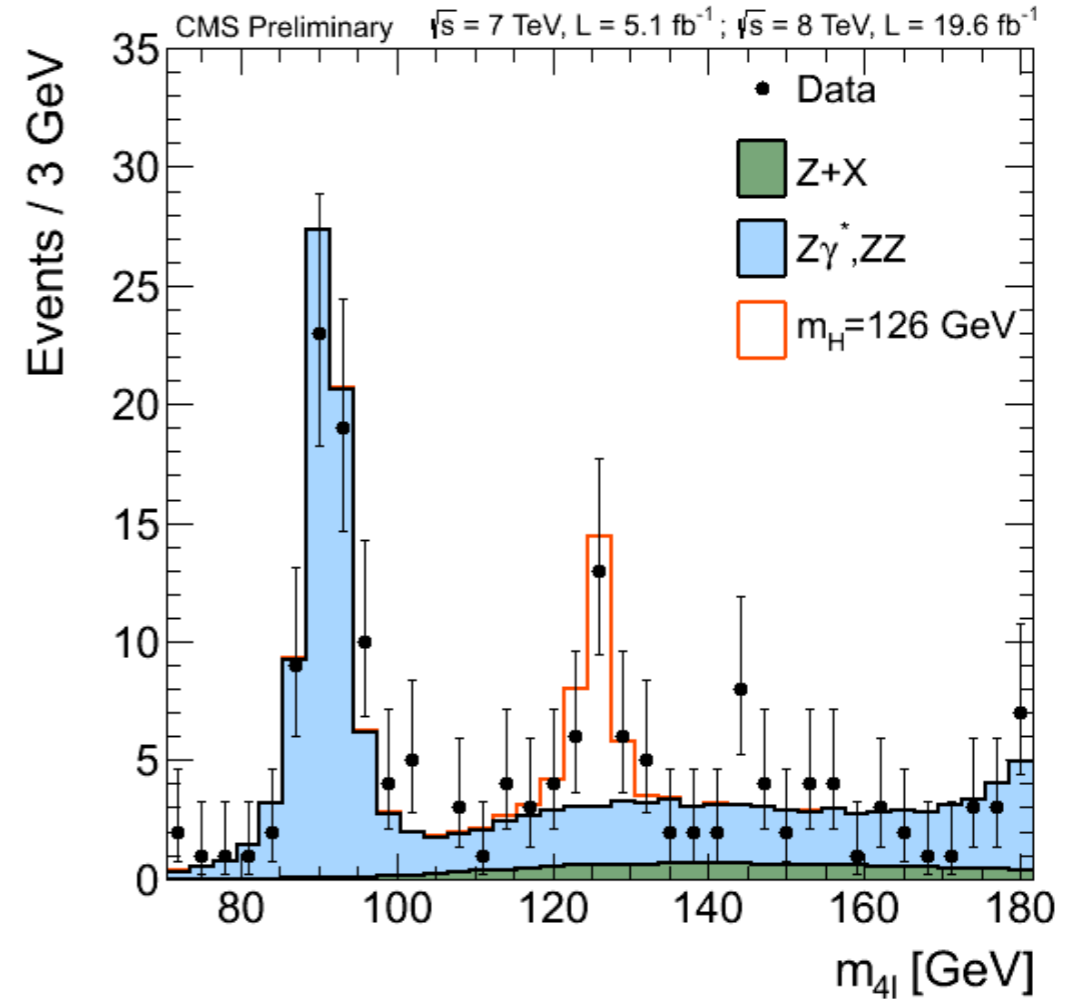
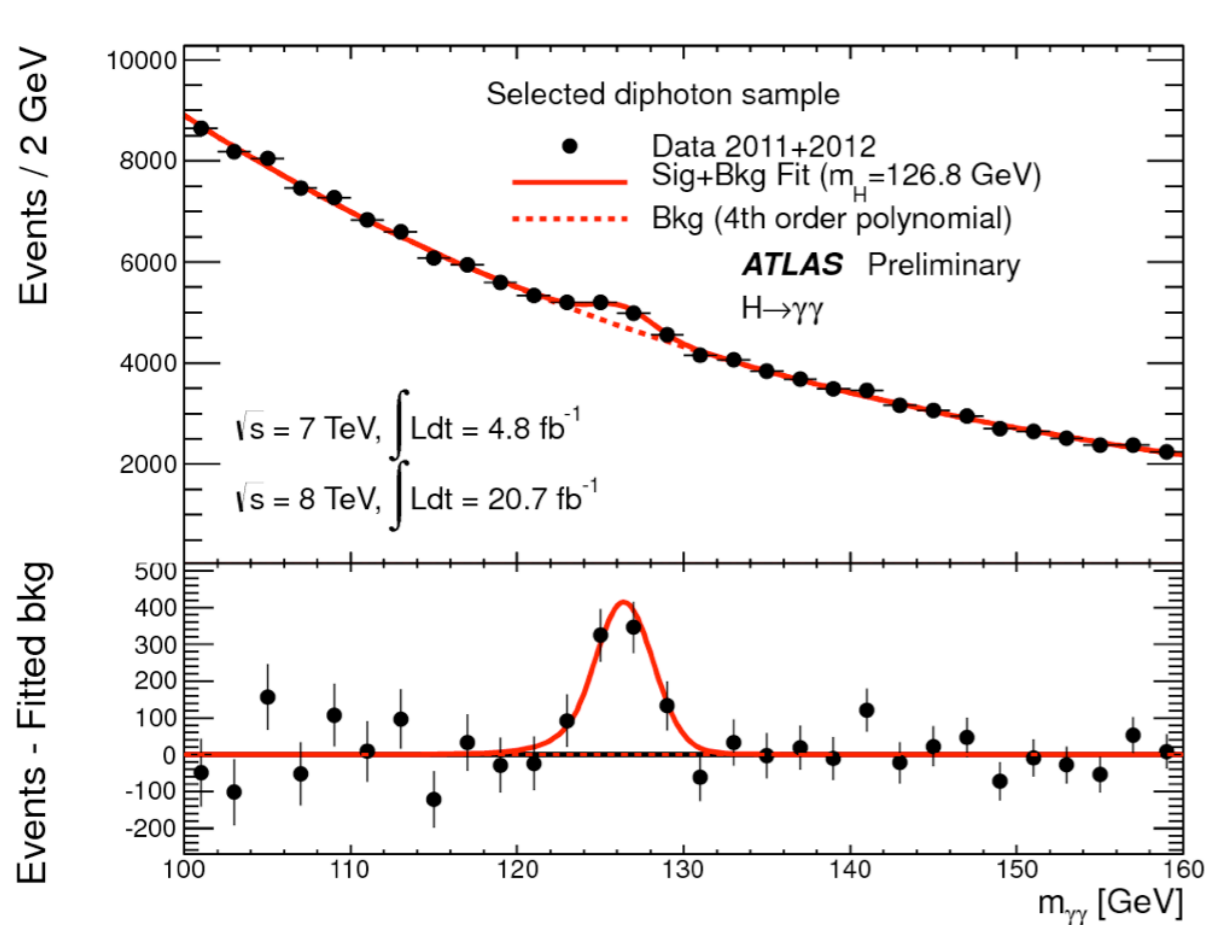
R. Frederix, S. Frixione, V. Hirschi, F. Maltoni, O. Mattelaer, P. Torrielli, E. Vryonidou, MZ  
arXiv: 1401.7340, PLB

Marco Zaro, LPTHE - UPMC Paris VI

*HiggsHunting 2014, Orsay*

*July 22, 2014*

# A new 125 GeV resonance has been found (long time ago...)



# Is it THE Higgs boson *as expected in the SM?*



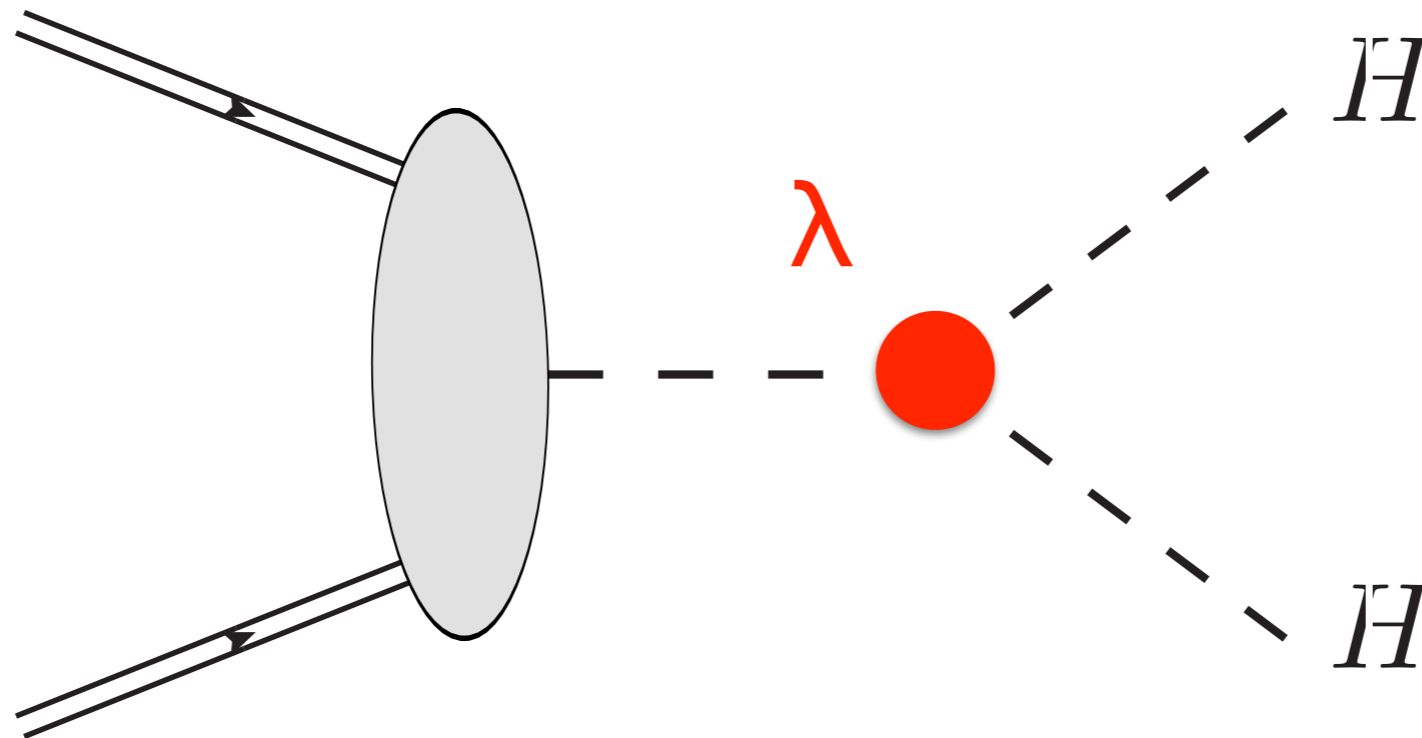
# Many properties have been measured...





... but one!

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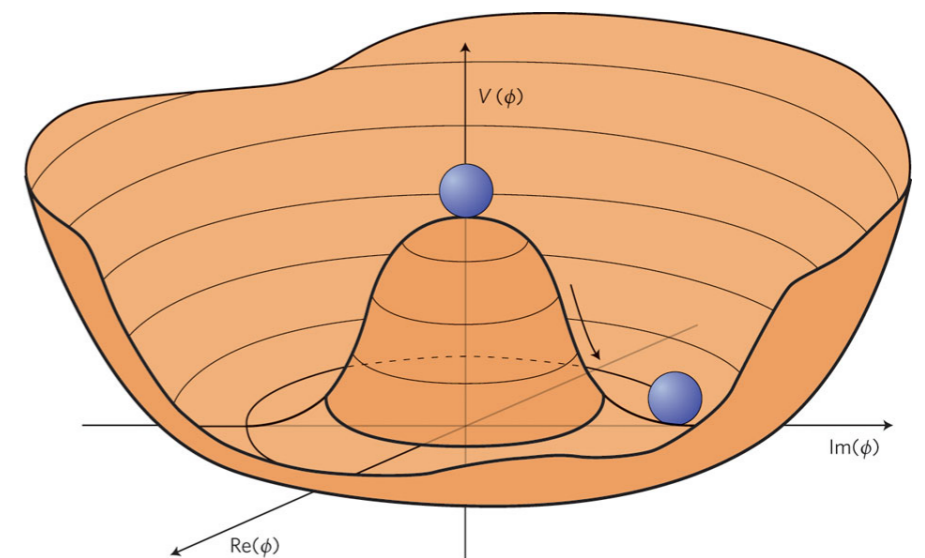




# The Higgs self-coupling

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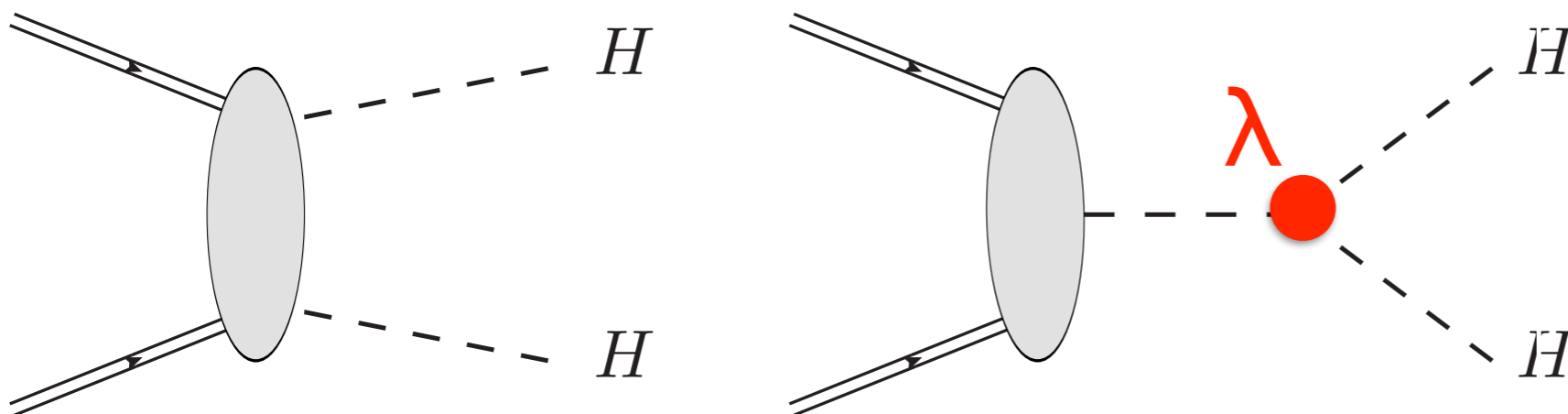
- The only parameter of the Higgs boson Lagrangian which cannot be measured in single Higgs production is its self-coupling  $\lambda$
- $\lambda$  drives the Higgs potential shape:  $V(\phi) = \mu^2\phi^2/2 + \lambda\phi^4/4$
- In the SM:  $M_H^2 = 2\lambda v^2 = -2\mu^2$



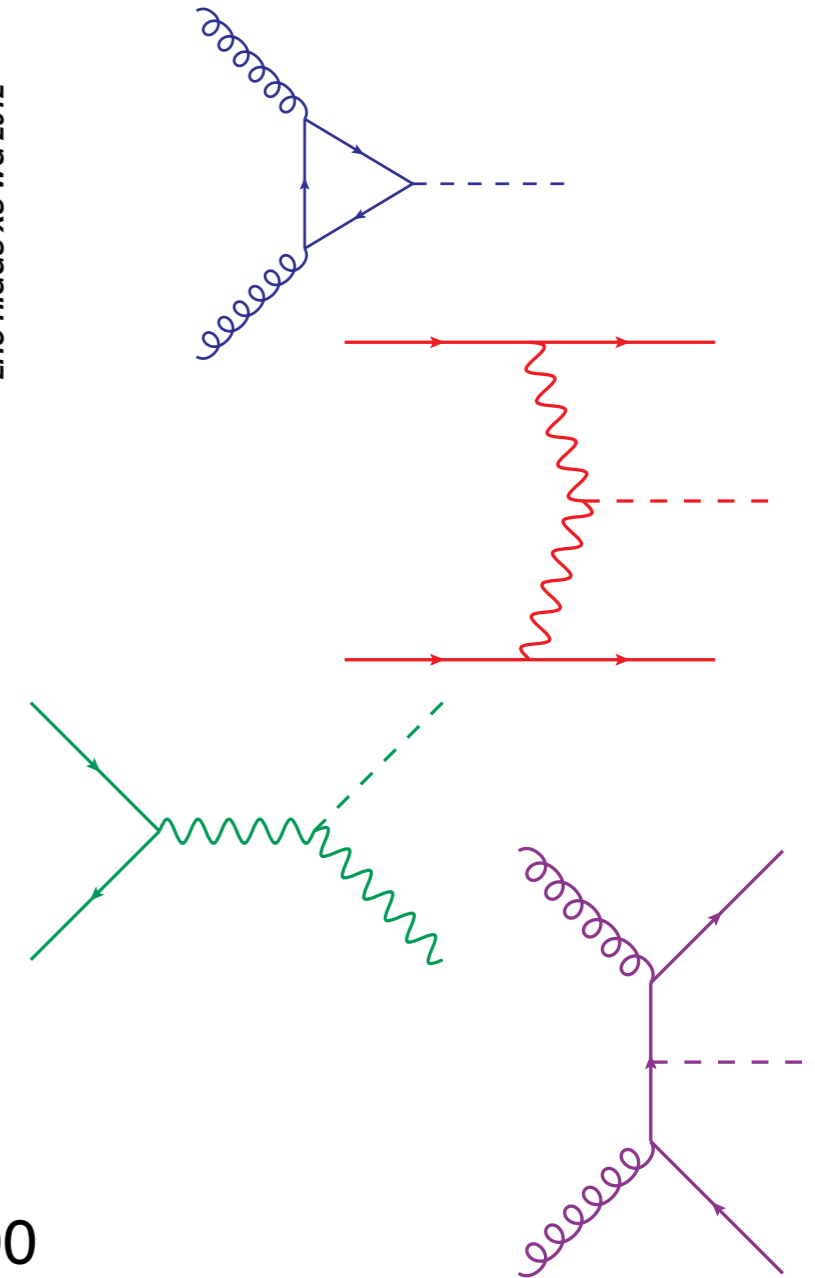
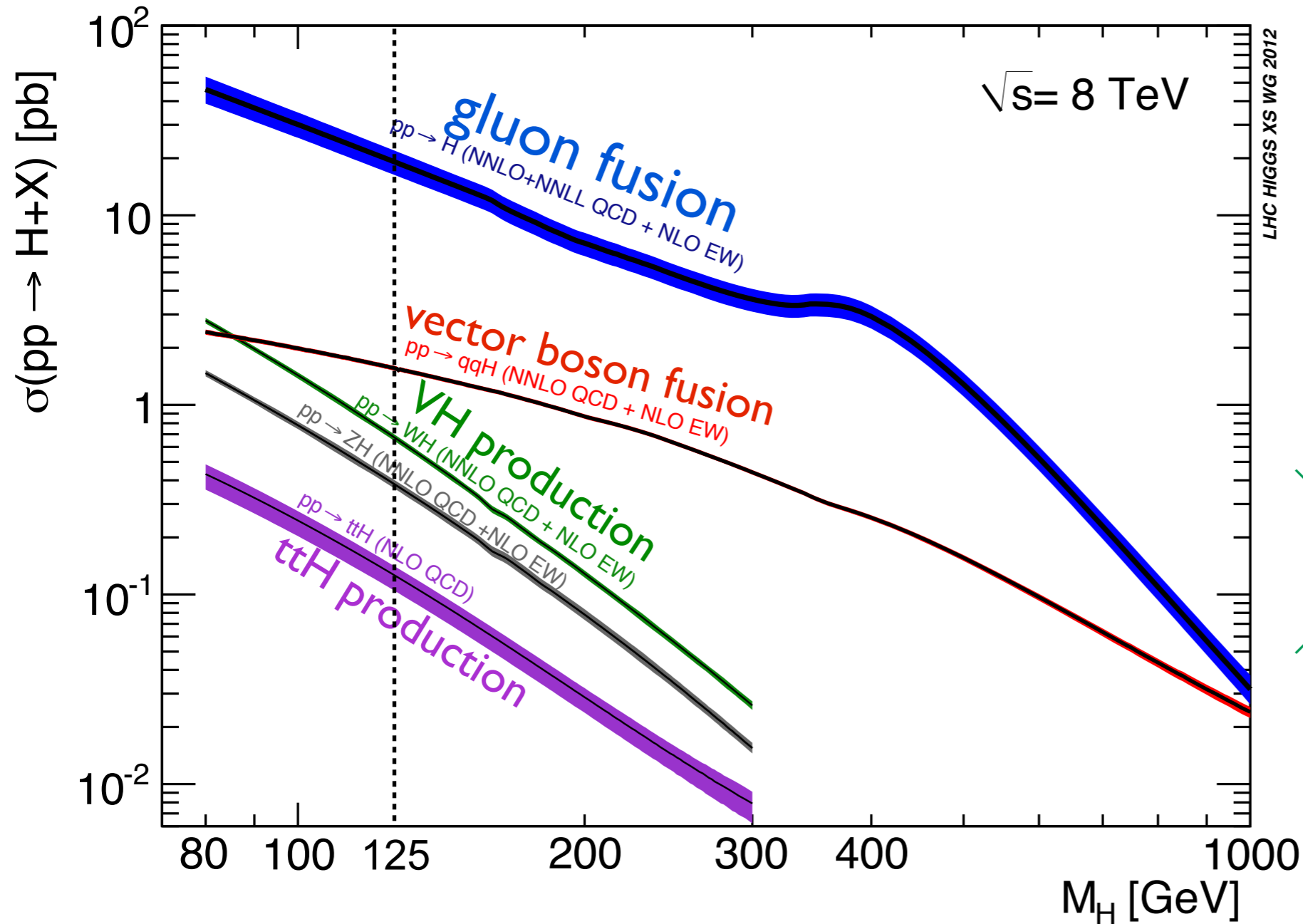


# The Higgs self-coupling

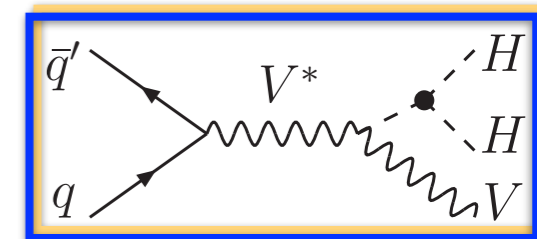
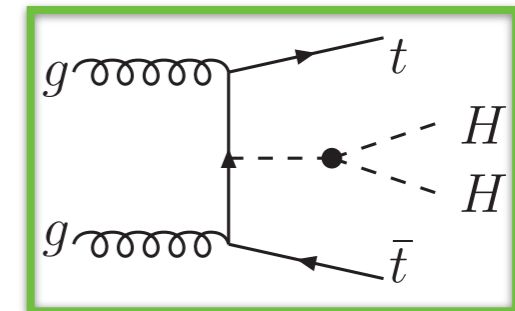
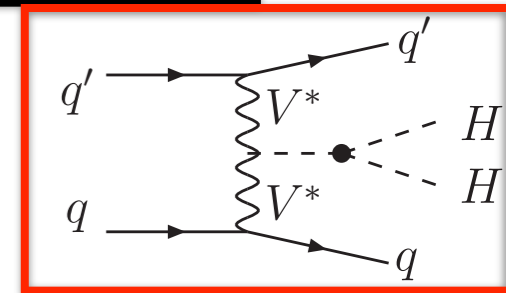
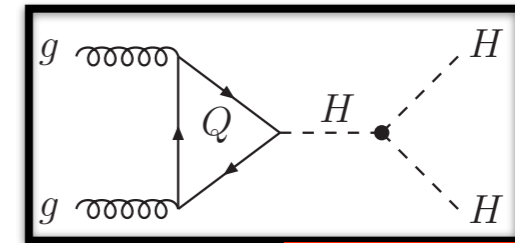
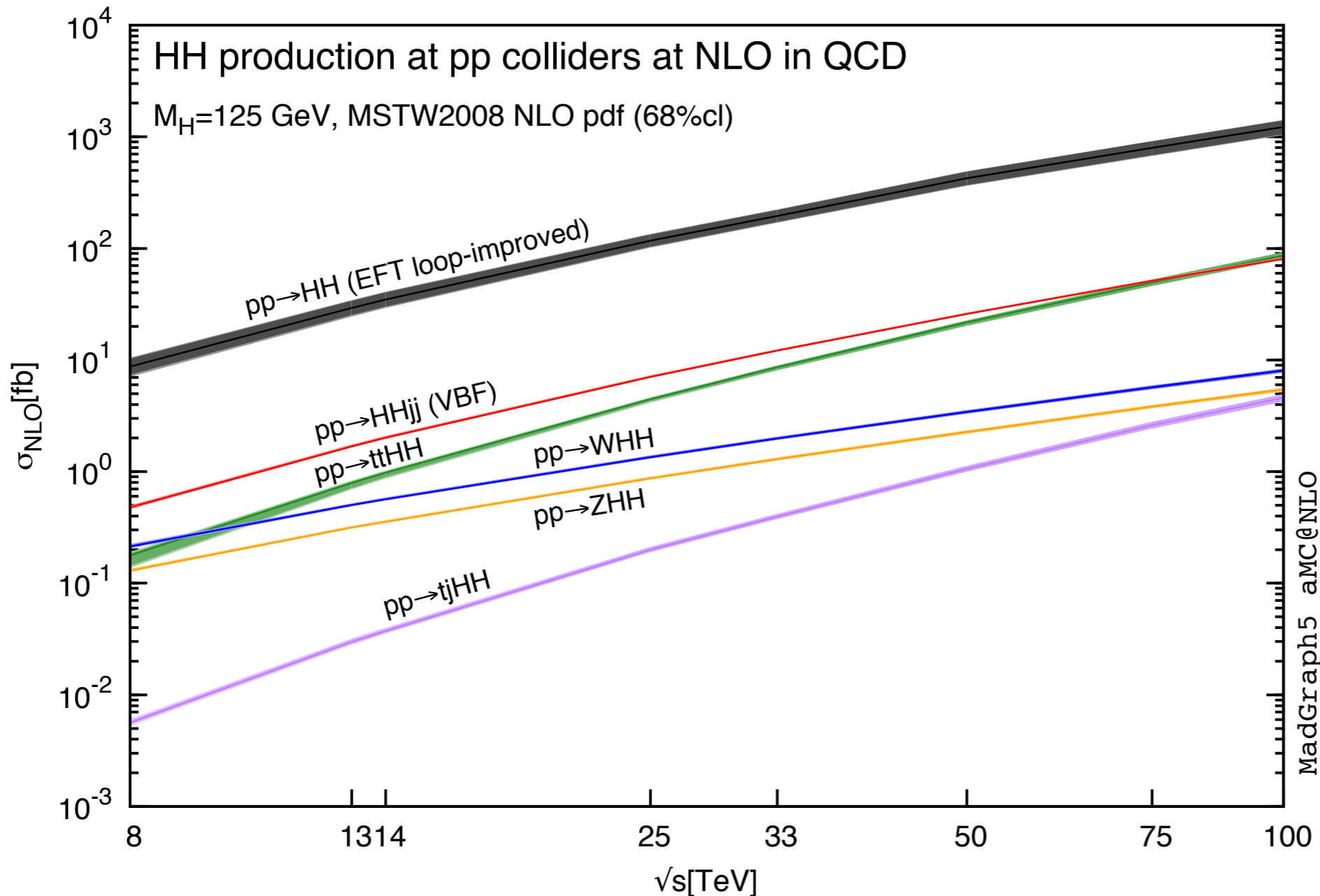
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  - In the SM:  $M_H^2 = 2\lambda v^2 = -2\mu^2$
- To measure lambda one has to look at double Higgs production



# Single Higgs production channels



# Double Higgs production channels



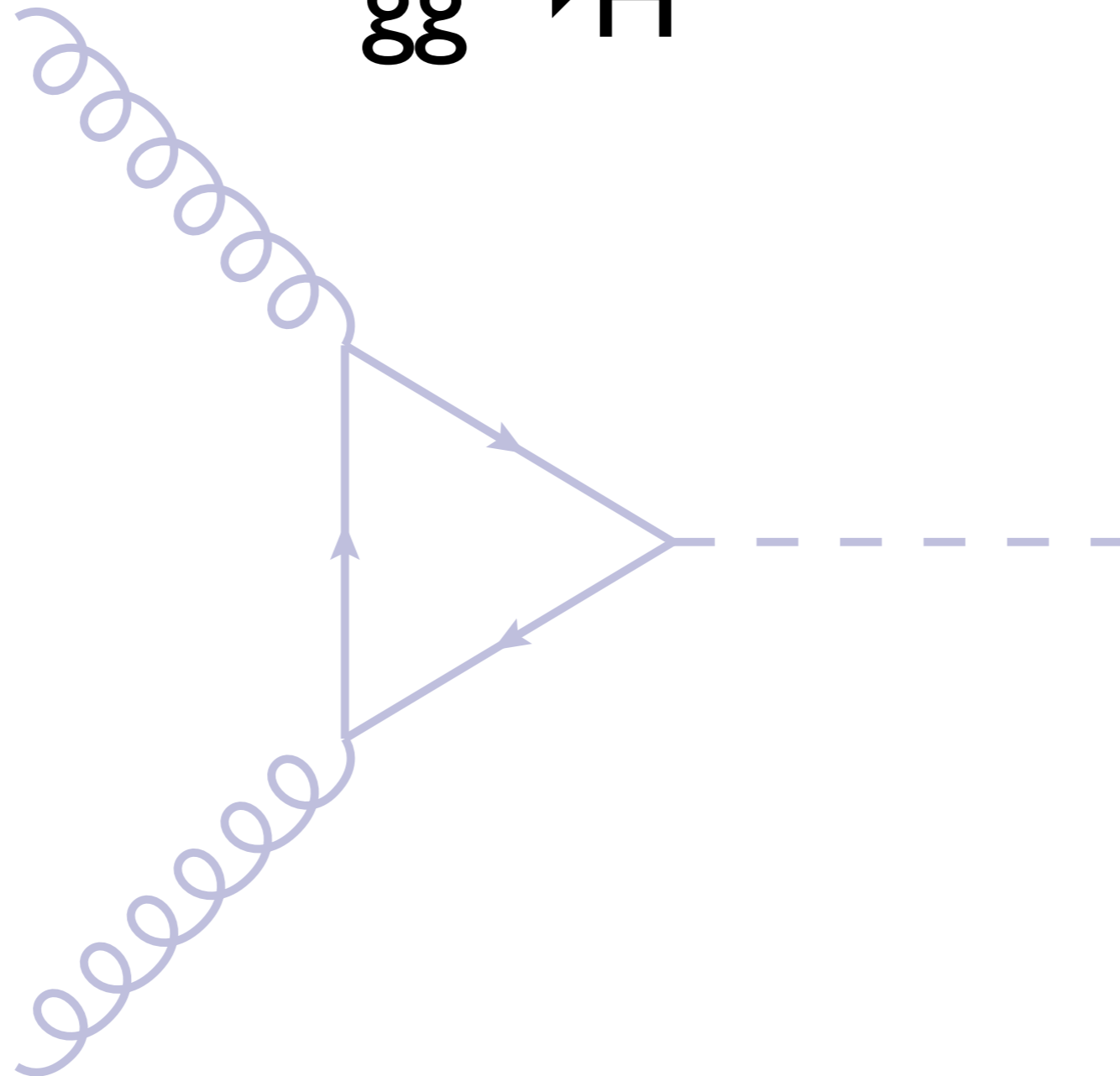
ggHH also known at NNLO (in the EFT) (de Florian, Mazzitelli, arXiv:1309.6594)

HH-VBF also known at NNLO (Liu-Sheng, Ren-You, Wen-Gan, Lei, Wei-Hua, Xiao-Zhou, arXiv:1401.7754)

VHH also known at NNLO (Baglio, Djouadi, Grober, Muhlleitner, Quevillon, Spira, arXiv:1212.5581)

# The tricky case:

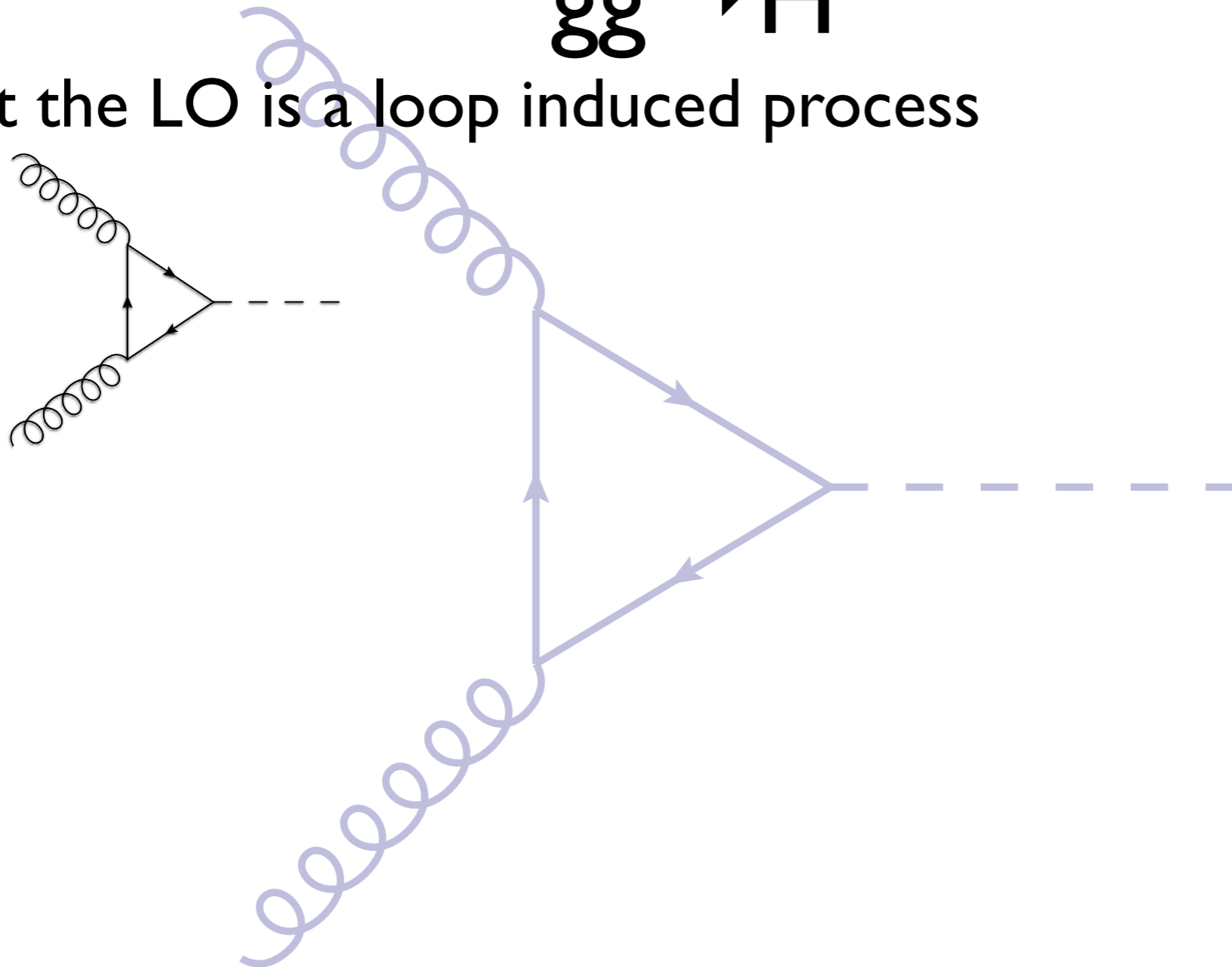
$$gg \rightarrow H$$



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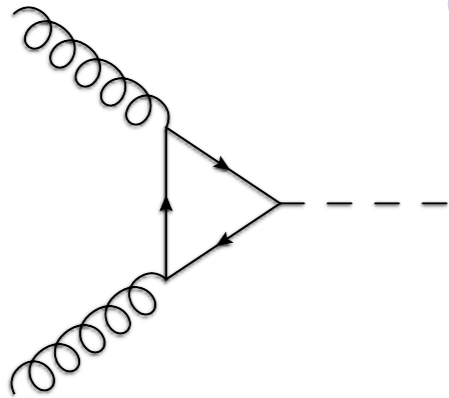
- $ggH$  at the LO is a loop induced process



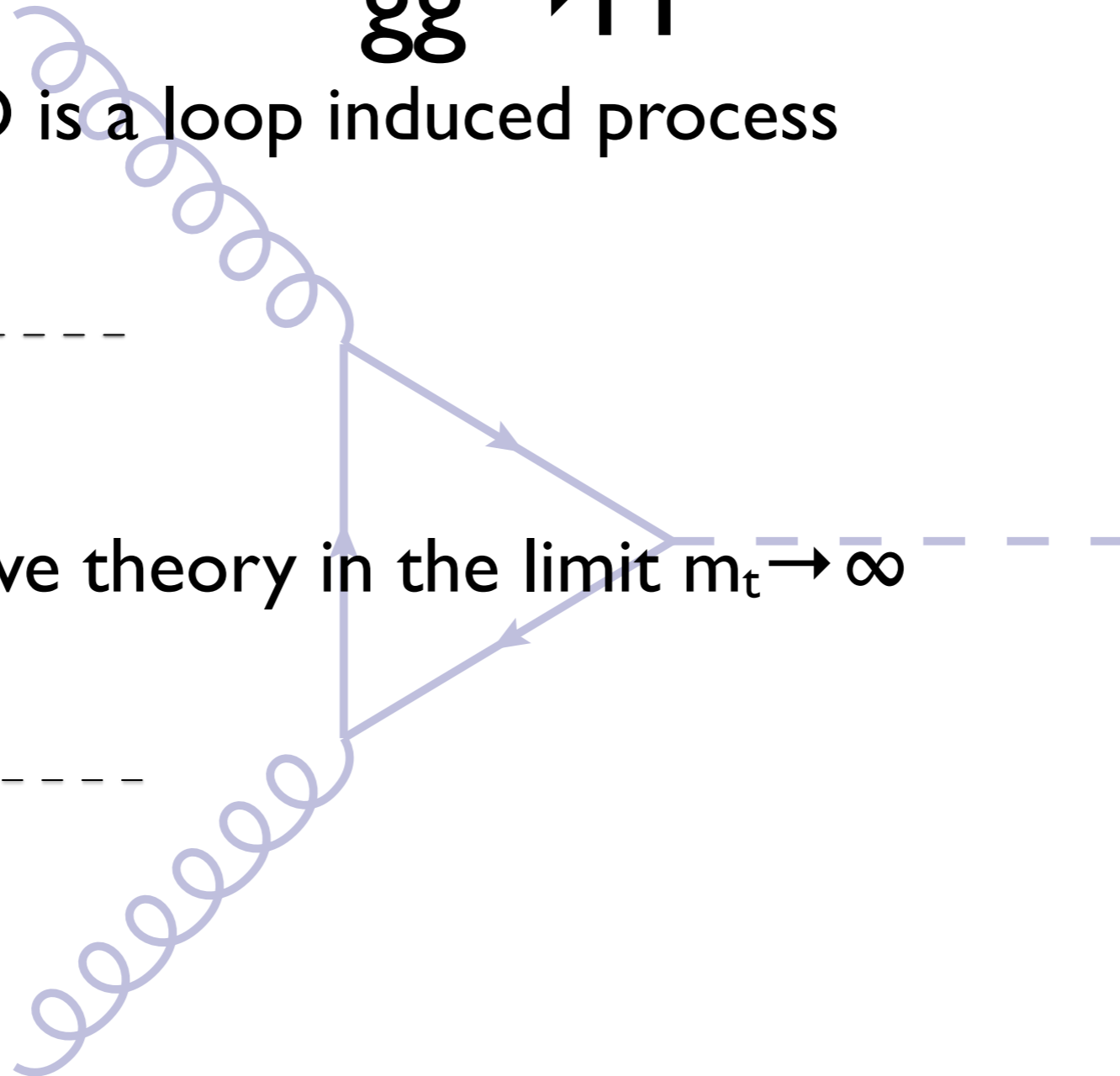
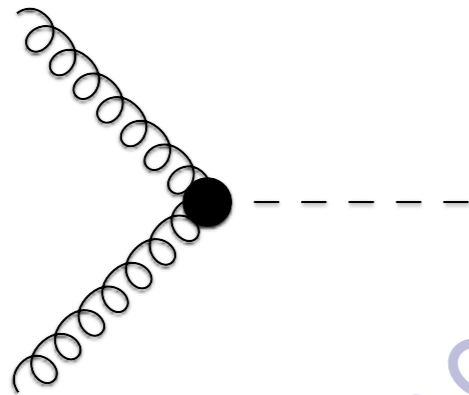
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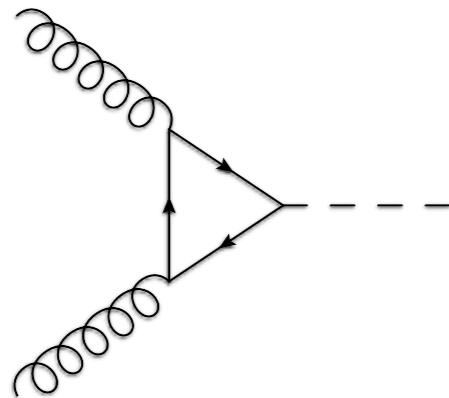
- Use an effective theory in the limit  $m_t \rightarrow \infty$



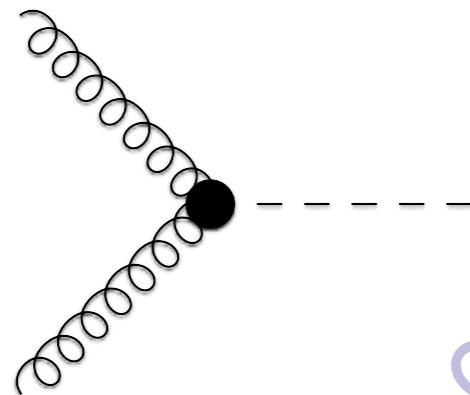
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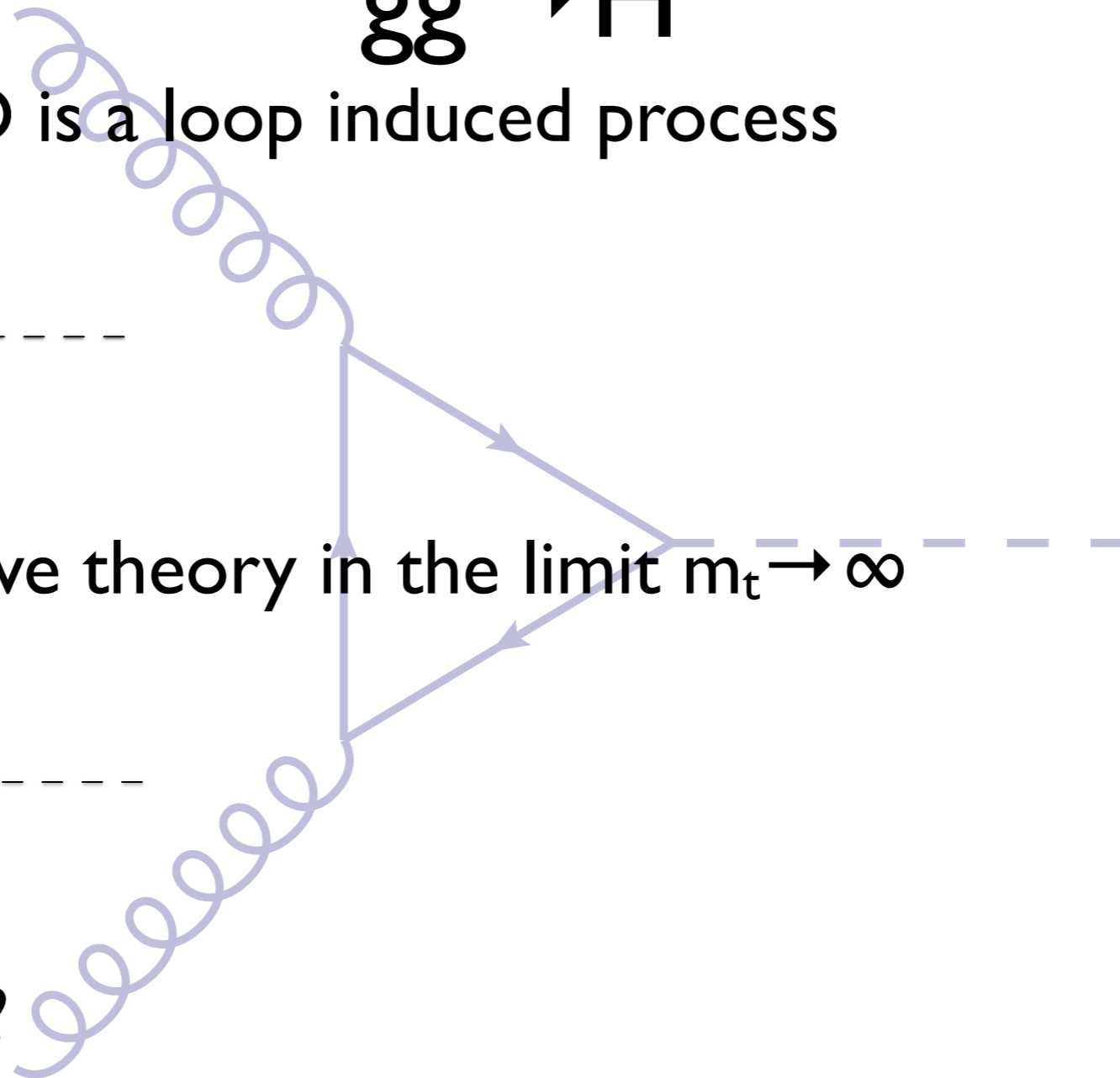
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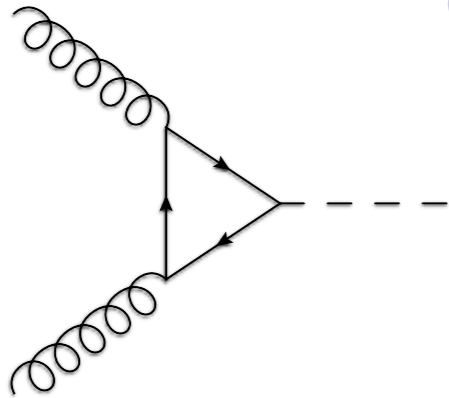
- Does it work?



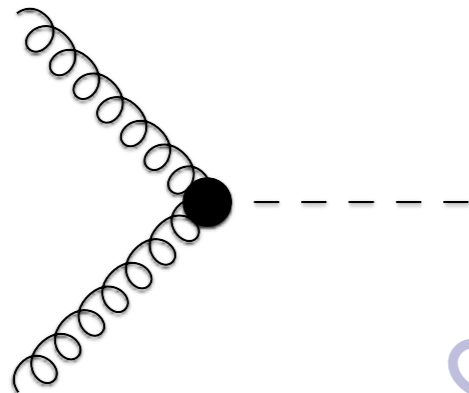
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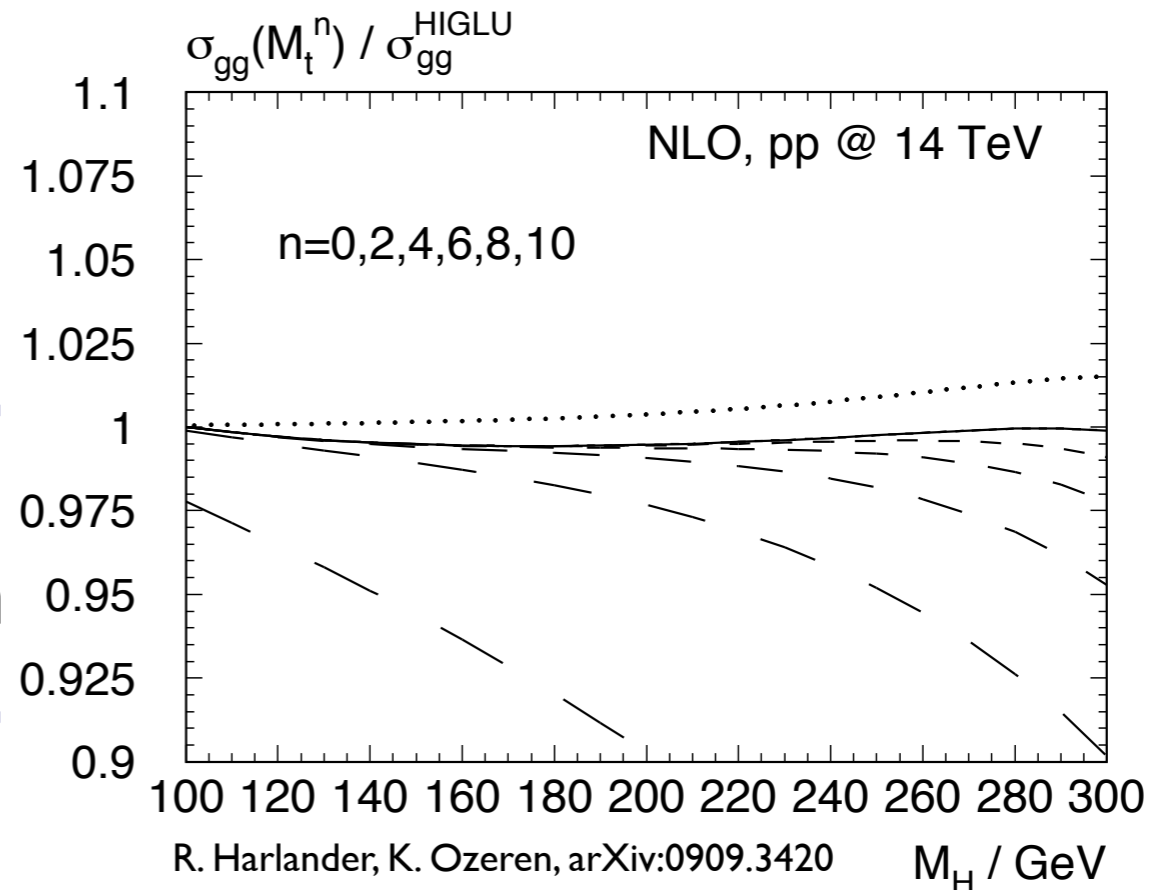


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- Does it work?

- For the cross-section it does quite well (if  $m_h < m_t$ )

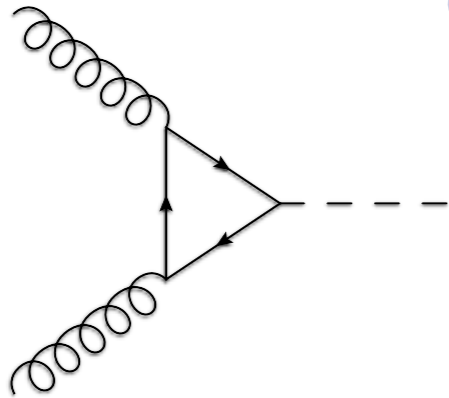




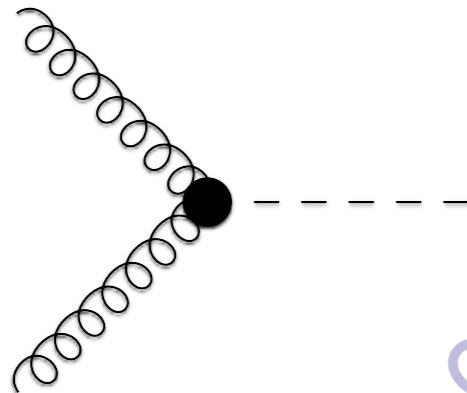
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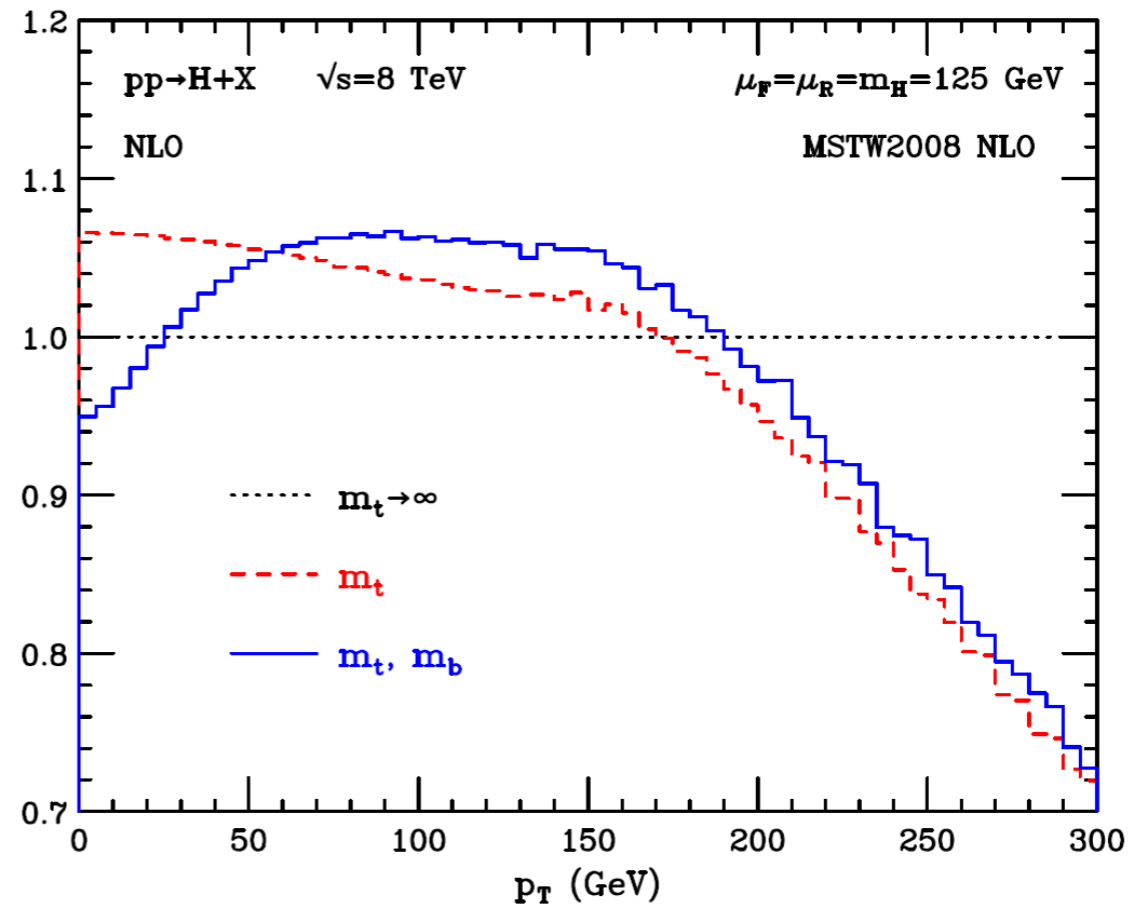


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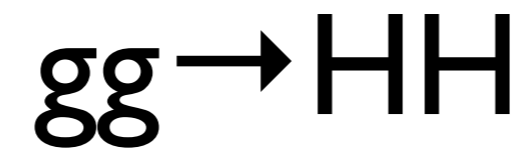
- Does it work?

- For the cross-section it does quite well (if  $m_h < m_t$ )
- Be careful for differential distributions (e.g Higgs  $p_T$ )





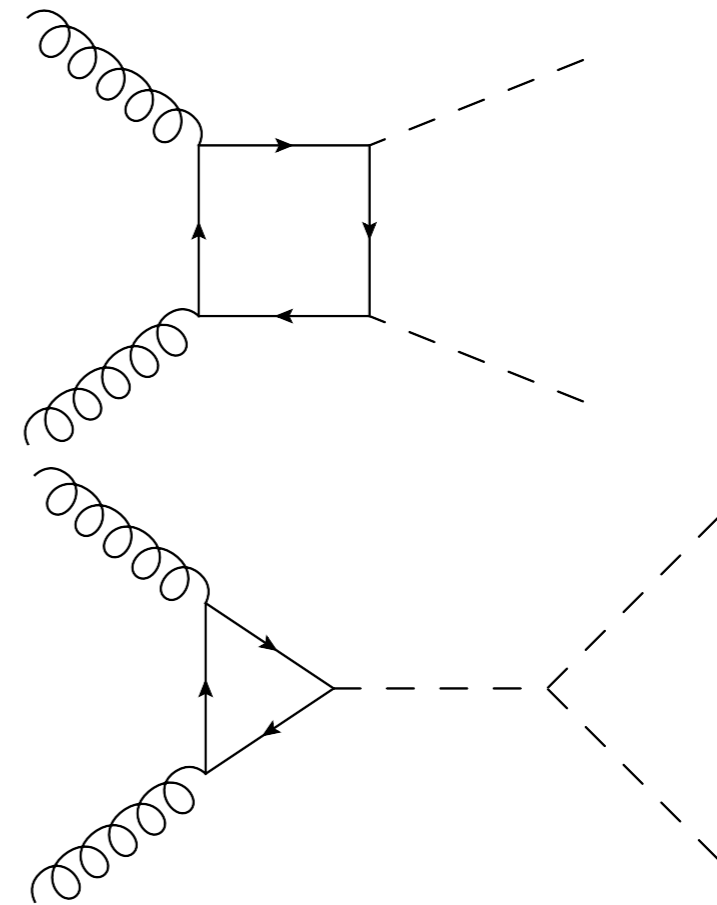
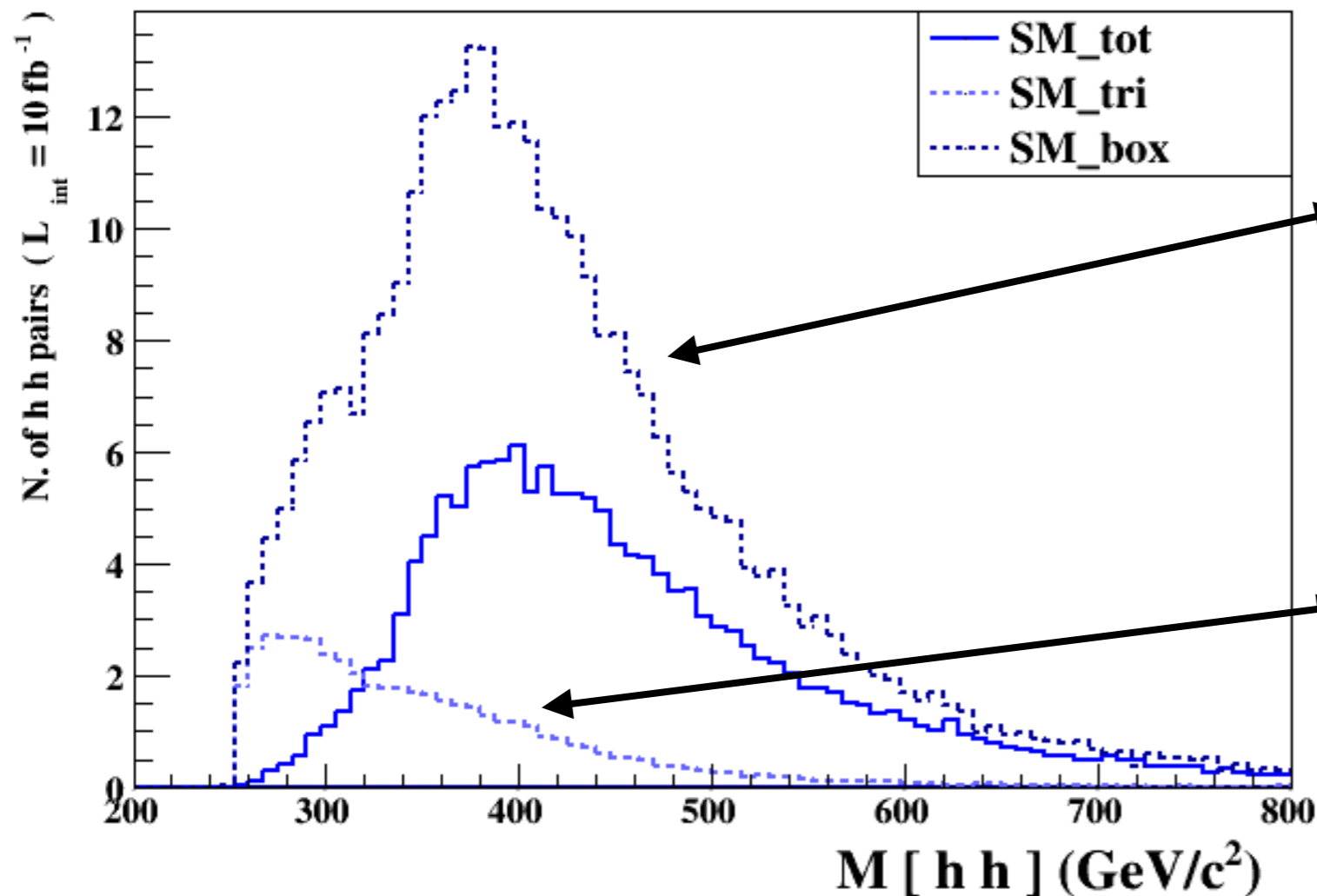
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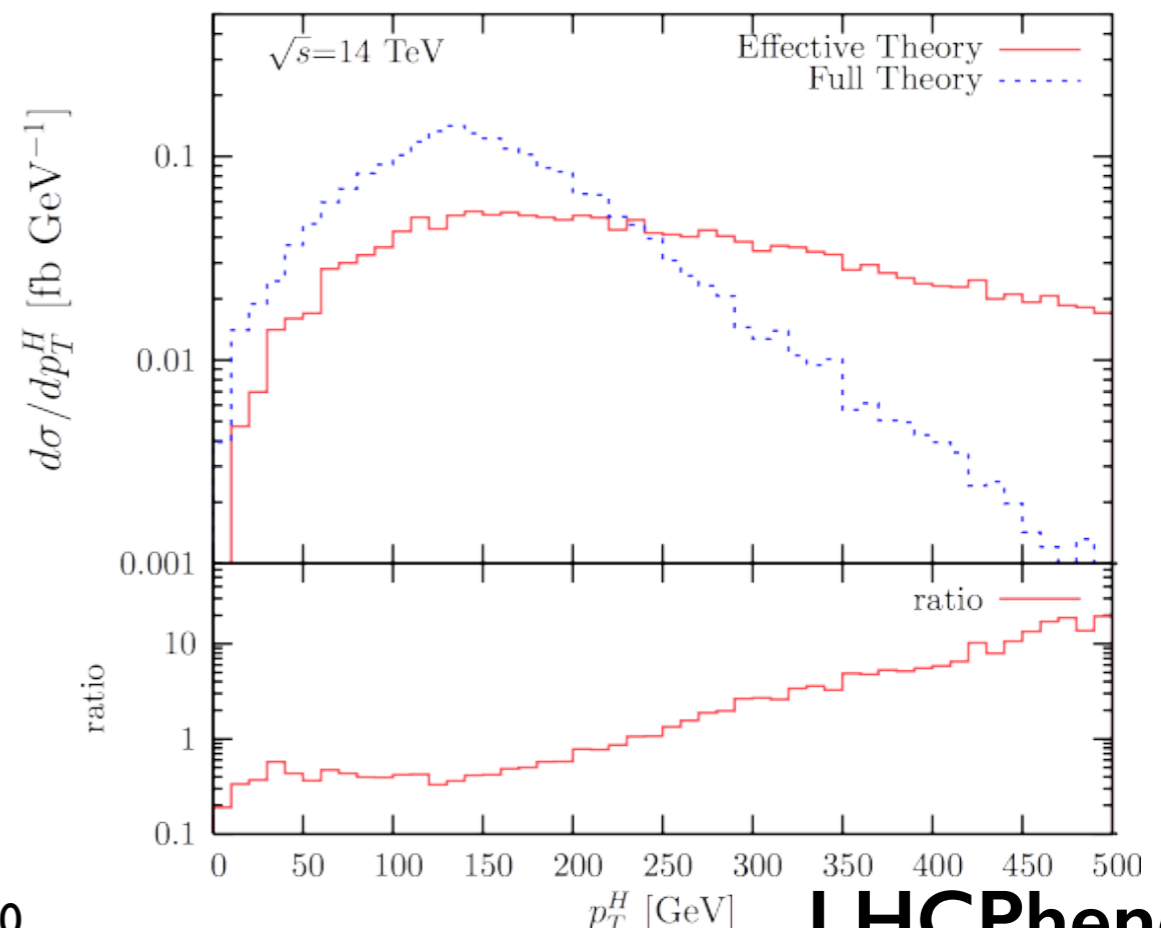
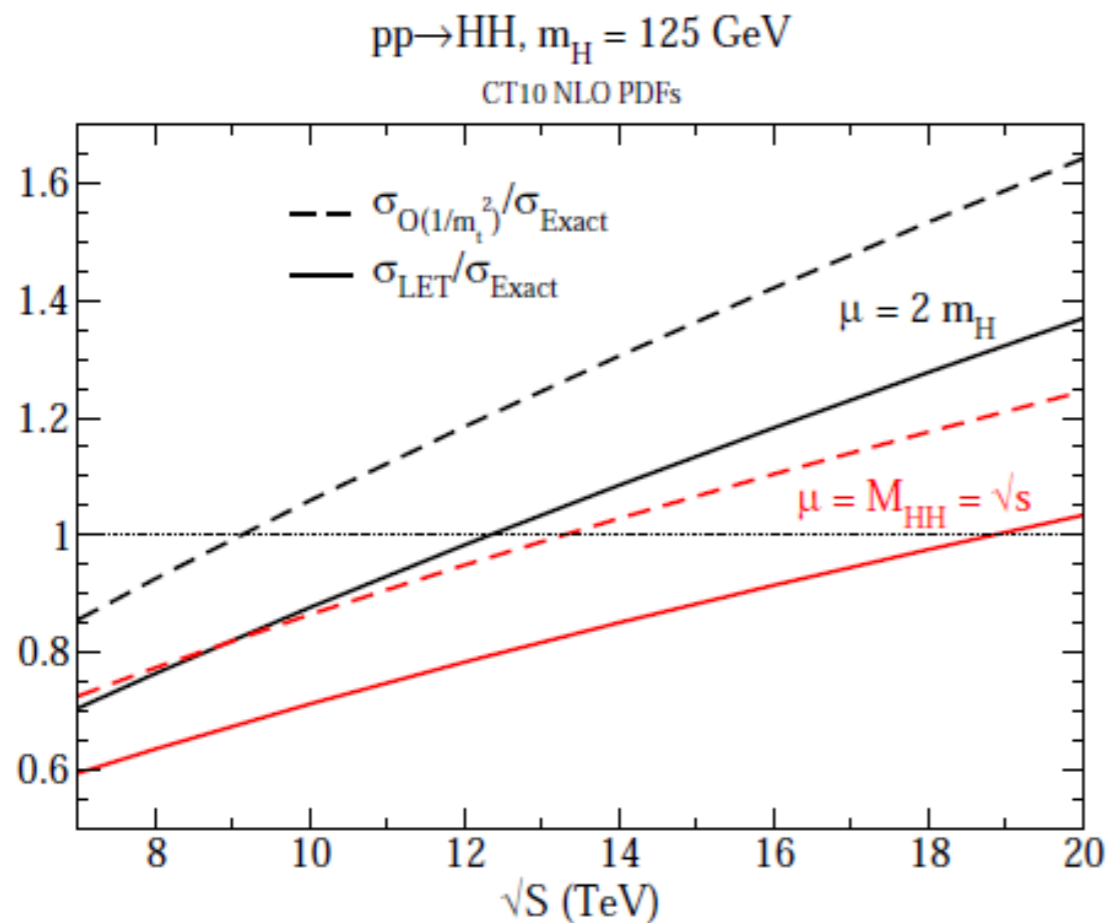
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- Unlike the single-Higgs case, EFT ( $m_t \rightarrow \infty$ ) does not work well



# The trickier case:

$$gg \rightarrow HH$$

- $gg \rightarrow HH$  at the LO is a loop induced process too!
  - Triangle and box diagrams interfere destructively
- Unlike the single-Higgs case, EFT ( $m_t \rightarrow \infty$ ) does not work well
  - Need to consistently take into account loop effects
    - **Include the exact one-loop matrix elements**



$gg \rightarrow HH @NLO$

with **MADGRAPH5\_AMC@NLO**

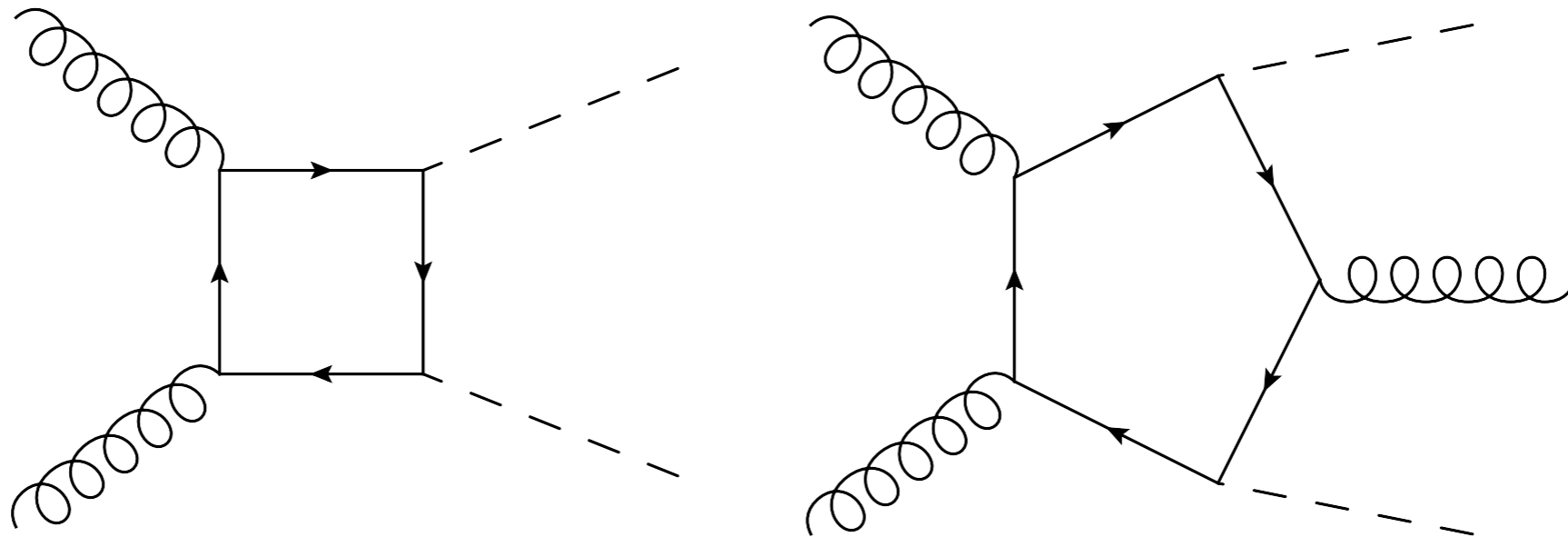
$$d\sigma_{NLO}^n = d\sigma_{LO}^n + d\sigma_V^n + \int d\Phi_1 d\sigma_R^{n+1}$$

# $gg \rightarrow HH @NLO$

with **MADGRAPH5\_AMC@NLO**

$$d\sigma_{NLO}^n = d\sigma_{LO}^n + d\sigma_V^n + \int d\Phi_1 d\sigma_R^{n+1}$$

- Include exact one-loop born and real emission ME

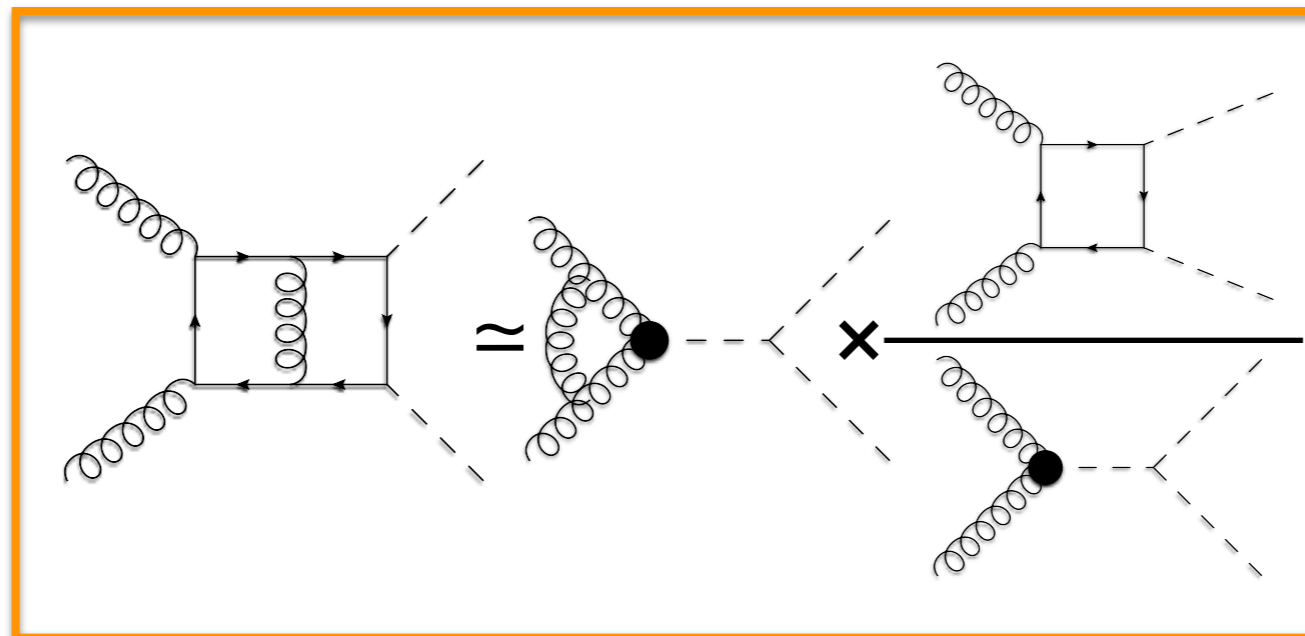


# gg → HH @NLO

with MADGRAPH5\_AMC@NLO

$$d\sigma_{NLO}^n = d\sigma_{LO}^n + d\sigma_V^n + \int d\Phi_1 d\sigma_R^{n+1}$$

- Include exact one-loop born and real emission ME
- Two-loop virtual ME is currently unknown
- Approximate with the born-rescaled EFT

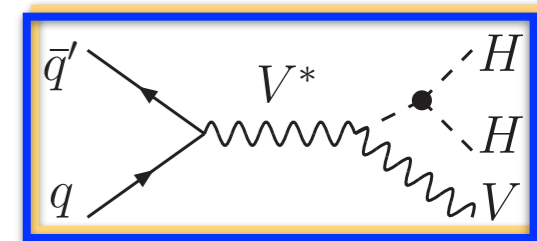
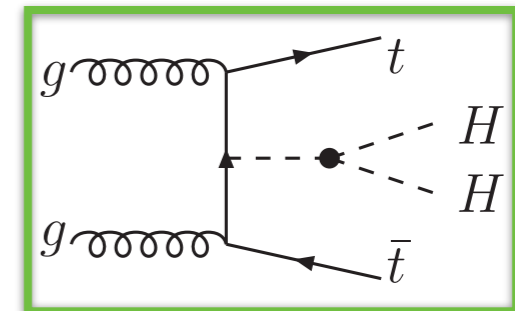
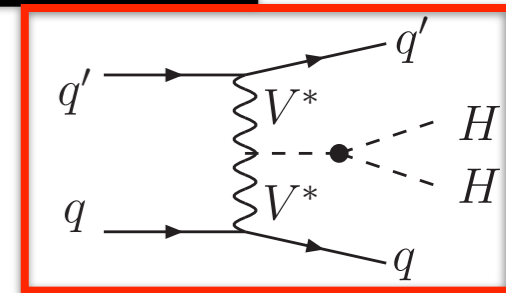
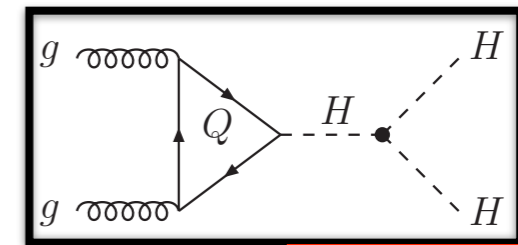
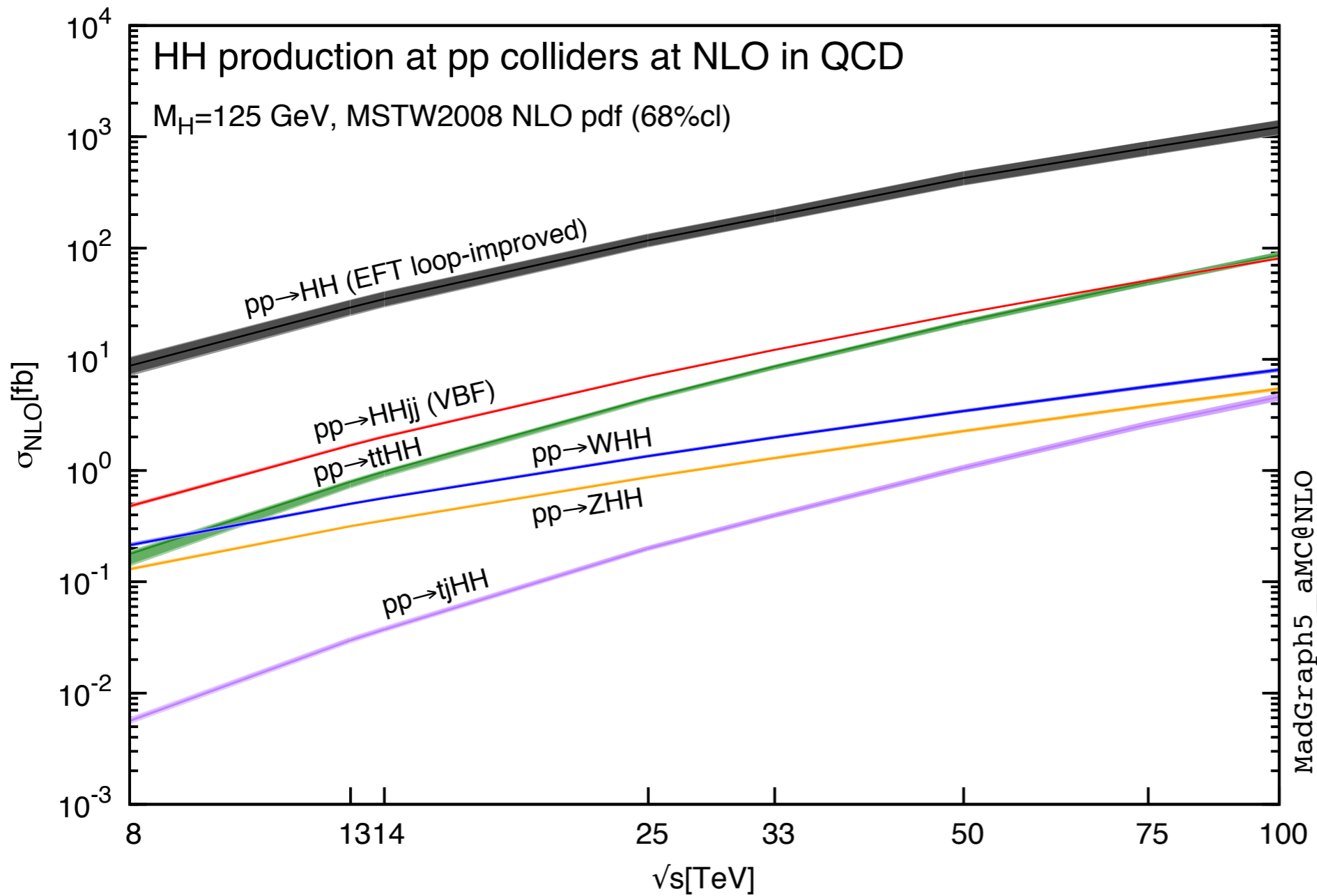




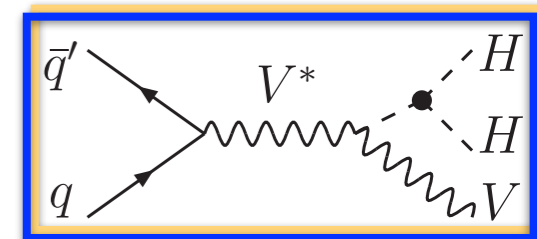
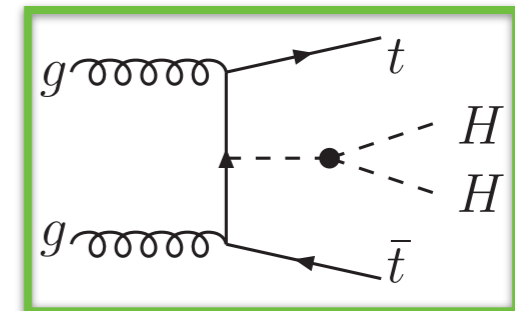
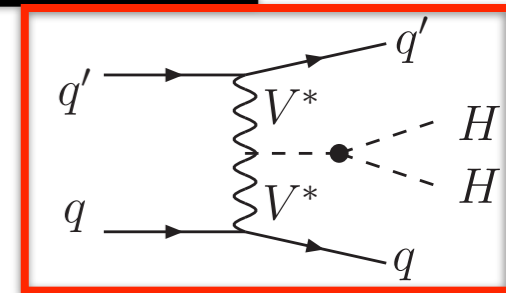
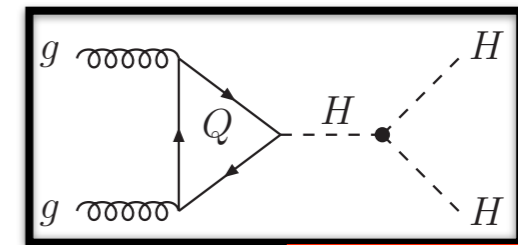
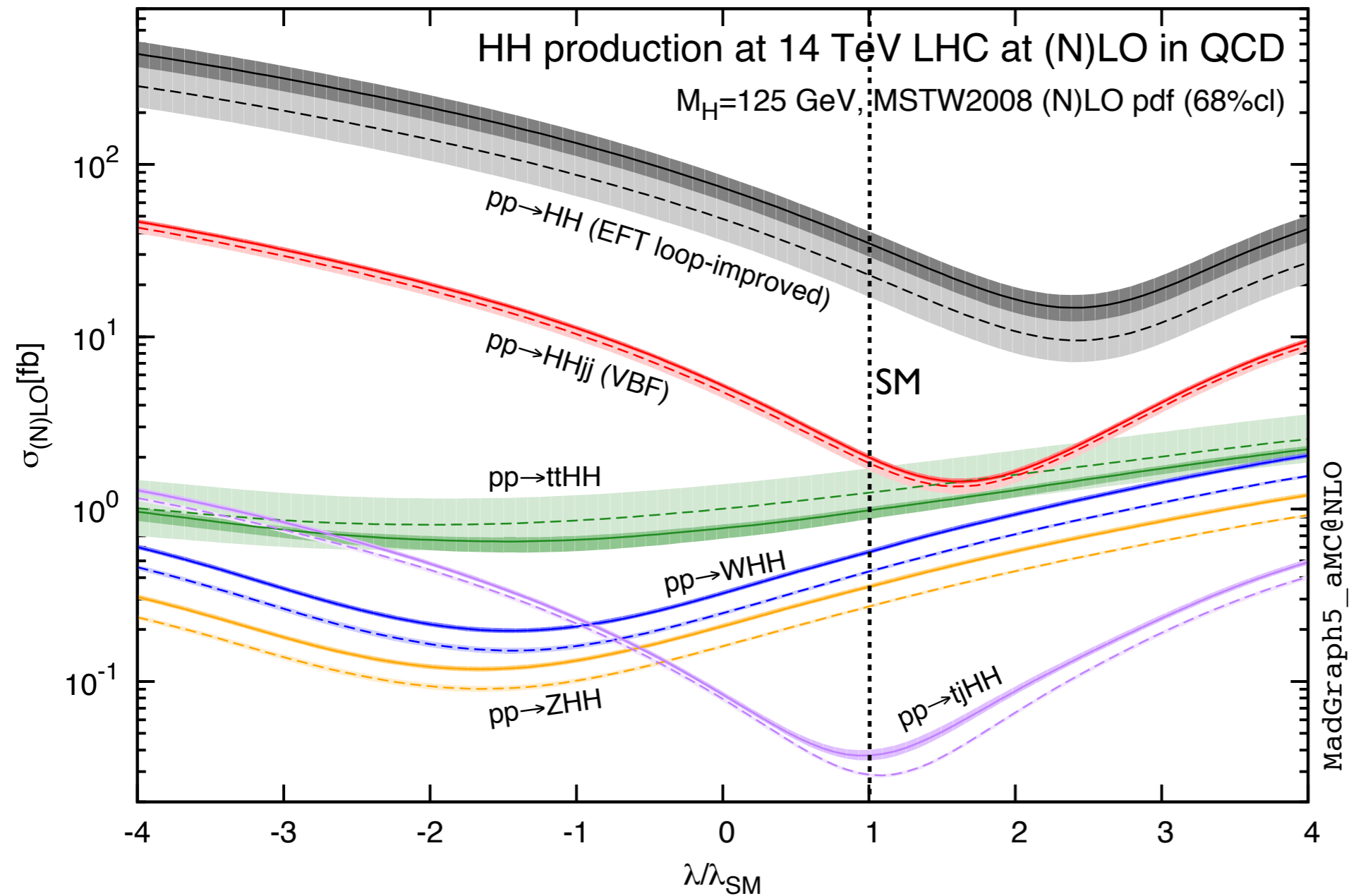


# Results

# Total cross-section

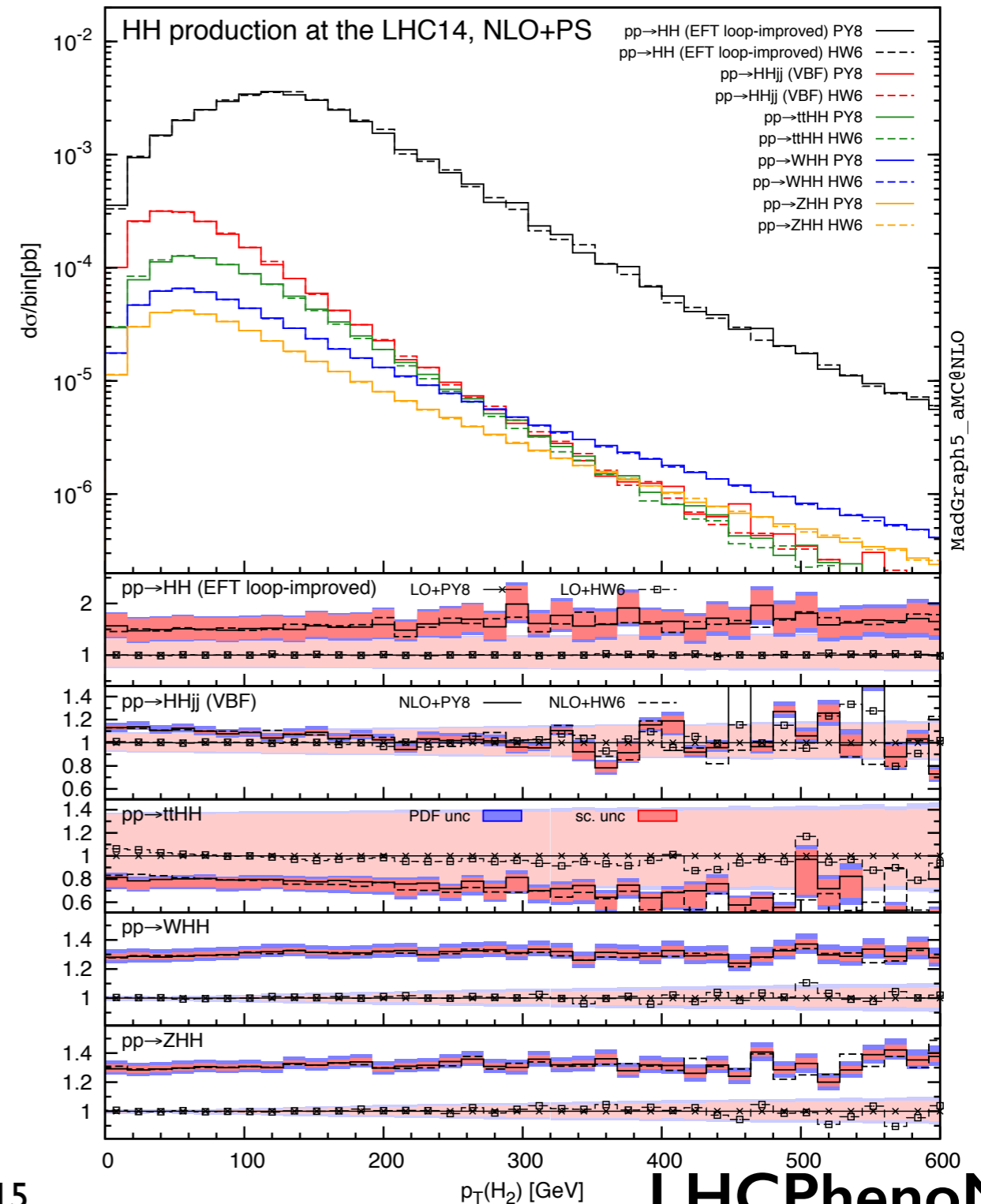
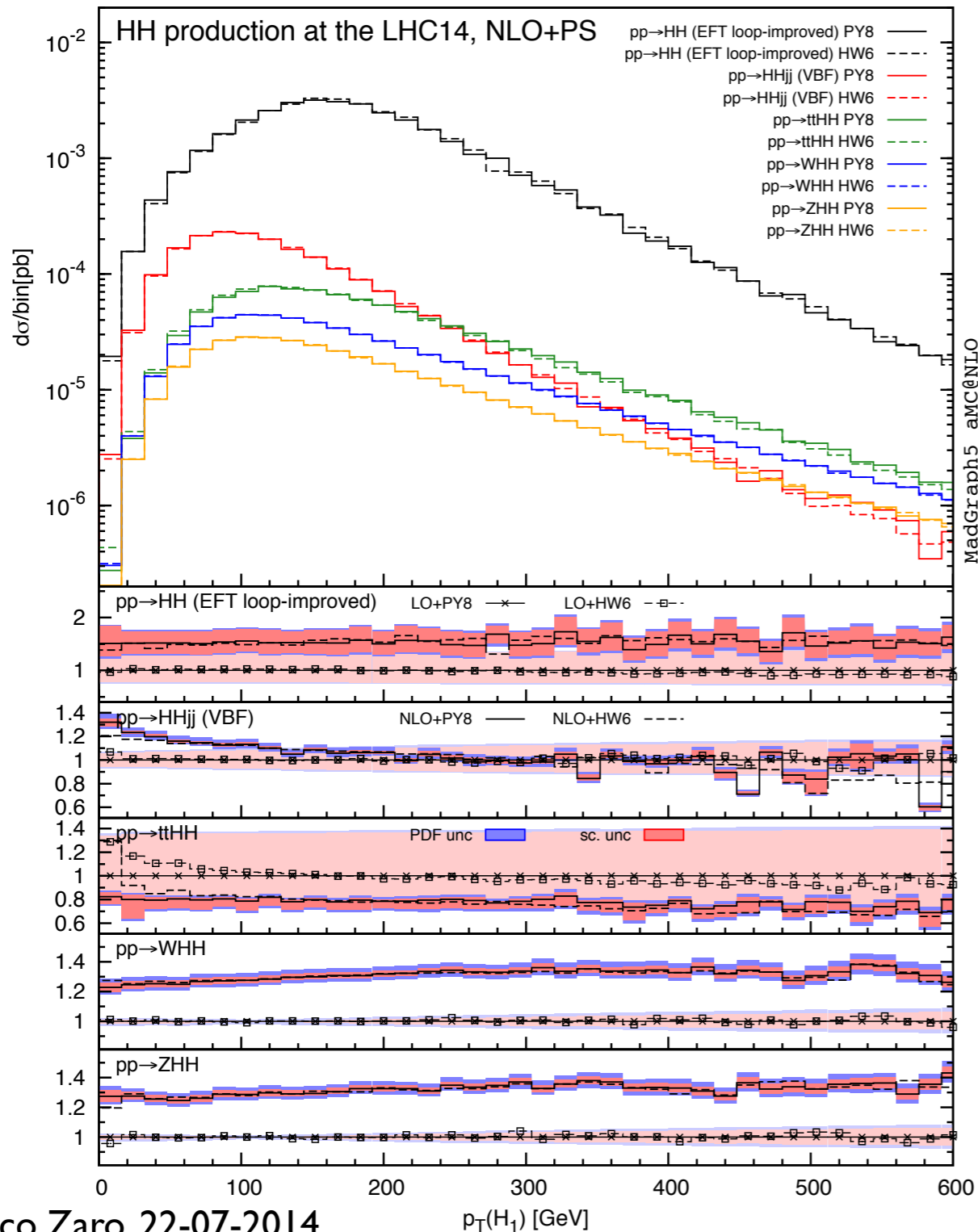


# Lambda dependence



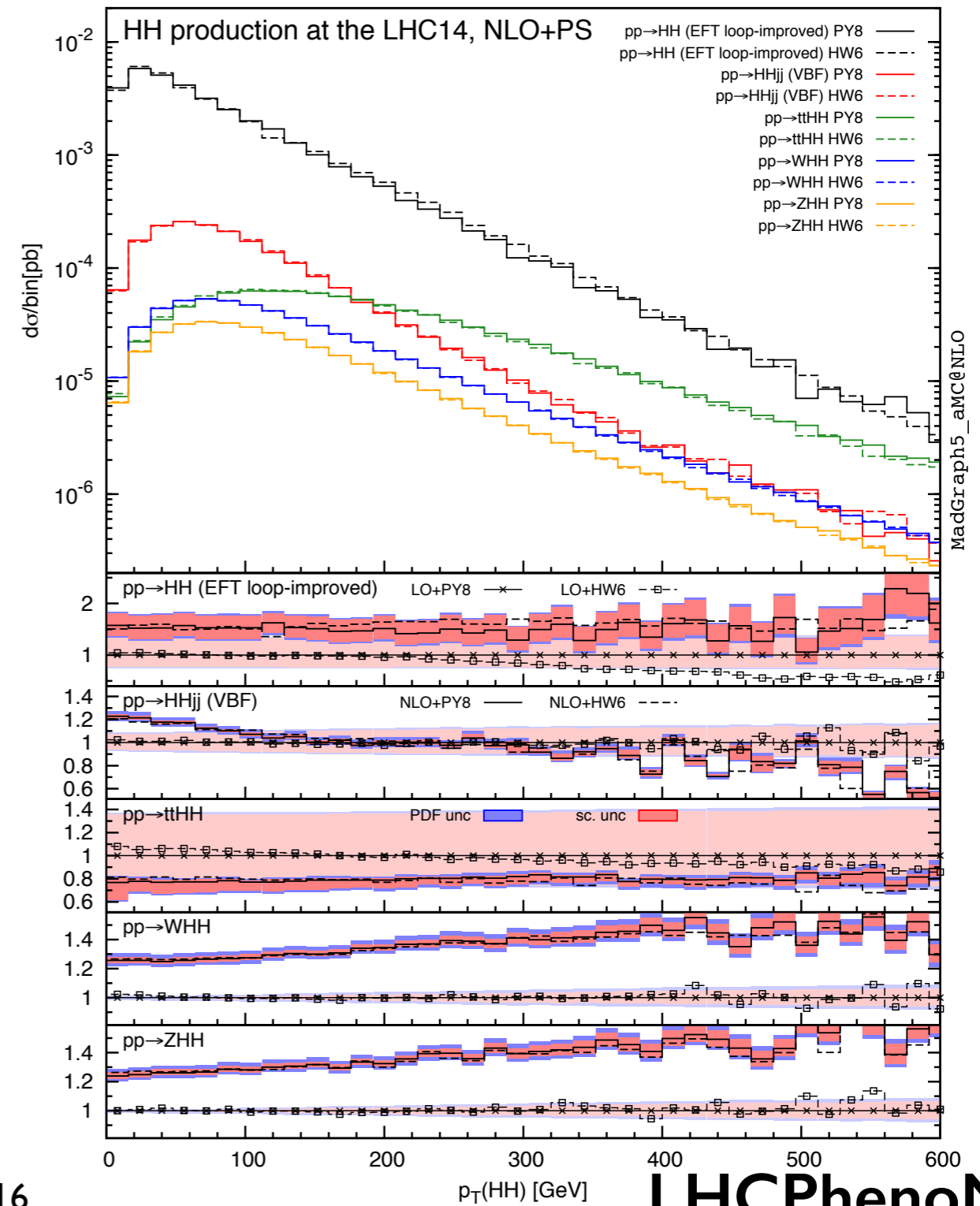
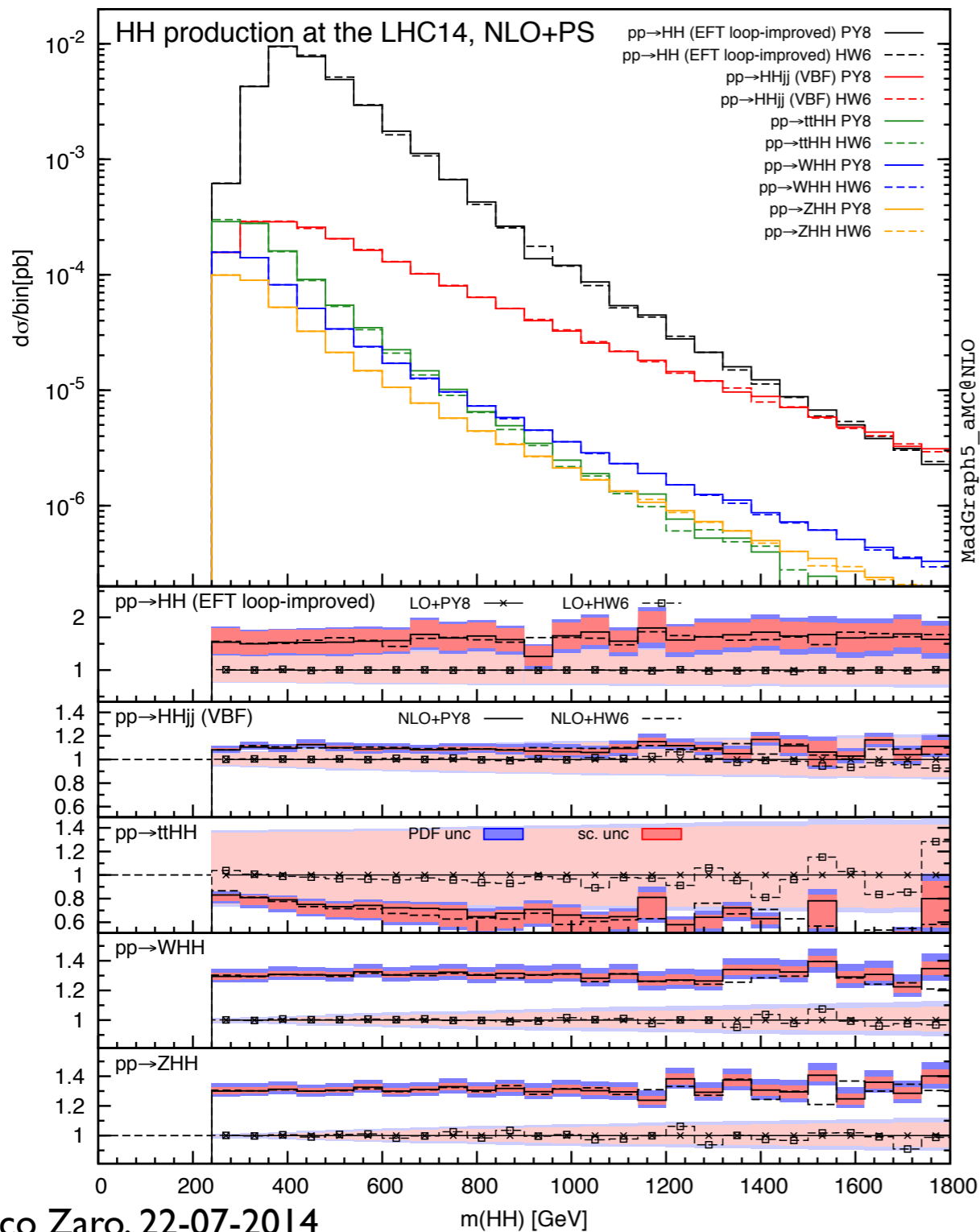
# HH differential observables

(for the first time at NLO + PS)



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# Conclusions

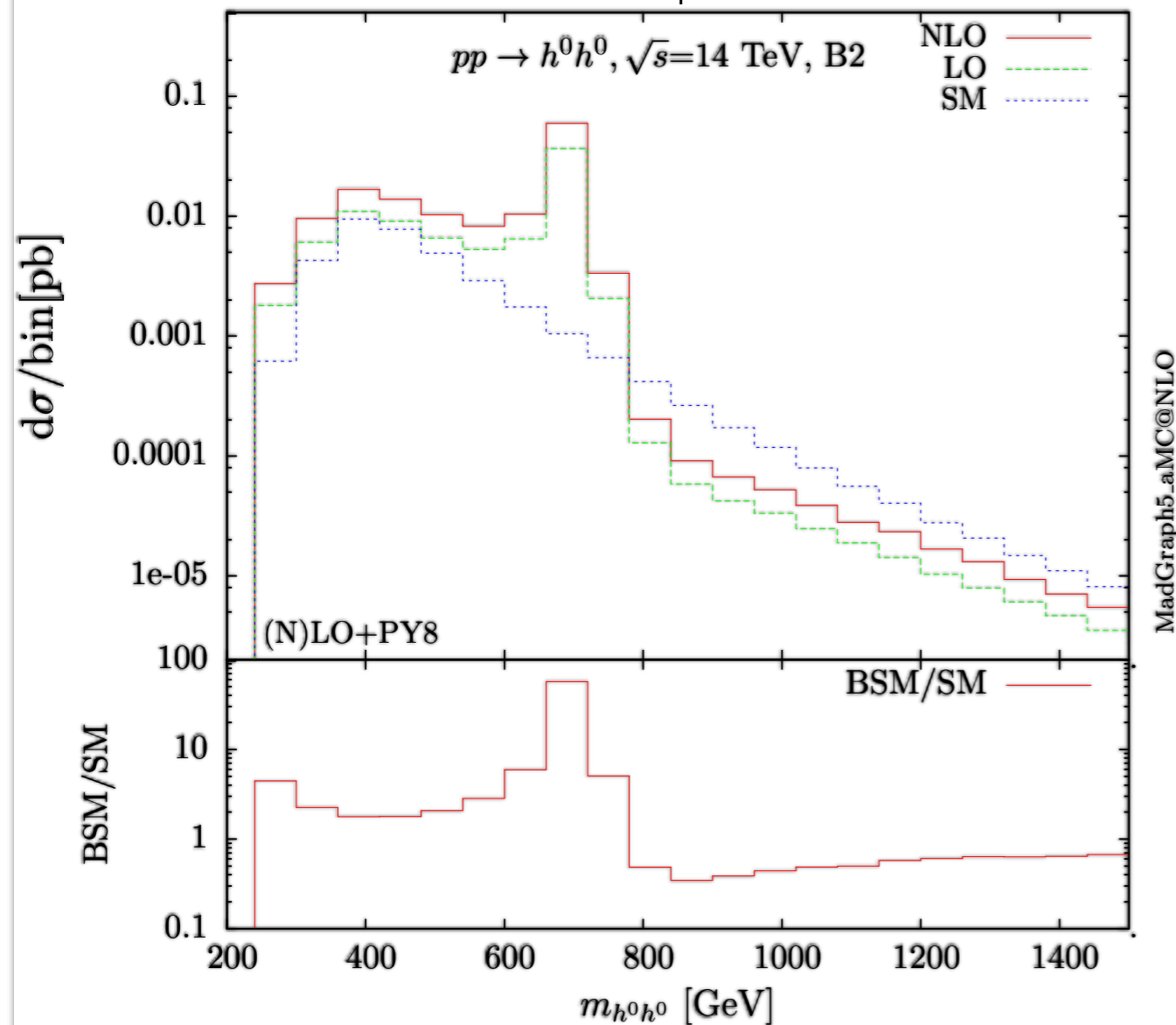
- Higgs pair production will be a key process to be looked at the next run of the LHC and at a FC
- It is the simplest class of processes which is sensitive to the Higgs boson self coupling  $\lambda$
- Accurate predictions for the Higgs pair production mechanisms are needed
- All production mechanisms can be computed at NLO accuracy within the `MADGRAPH5_AMC@NLO` framework,
  - All automated, but  $gg \rightarrow HH$
  - General approach to include loop-ME
- Fully differential predictions at NLO + PS are available for the first time for all production channels

# Conclusions

Hespel, Lopez-Val, Vryonidou arXiv:1407.0281

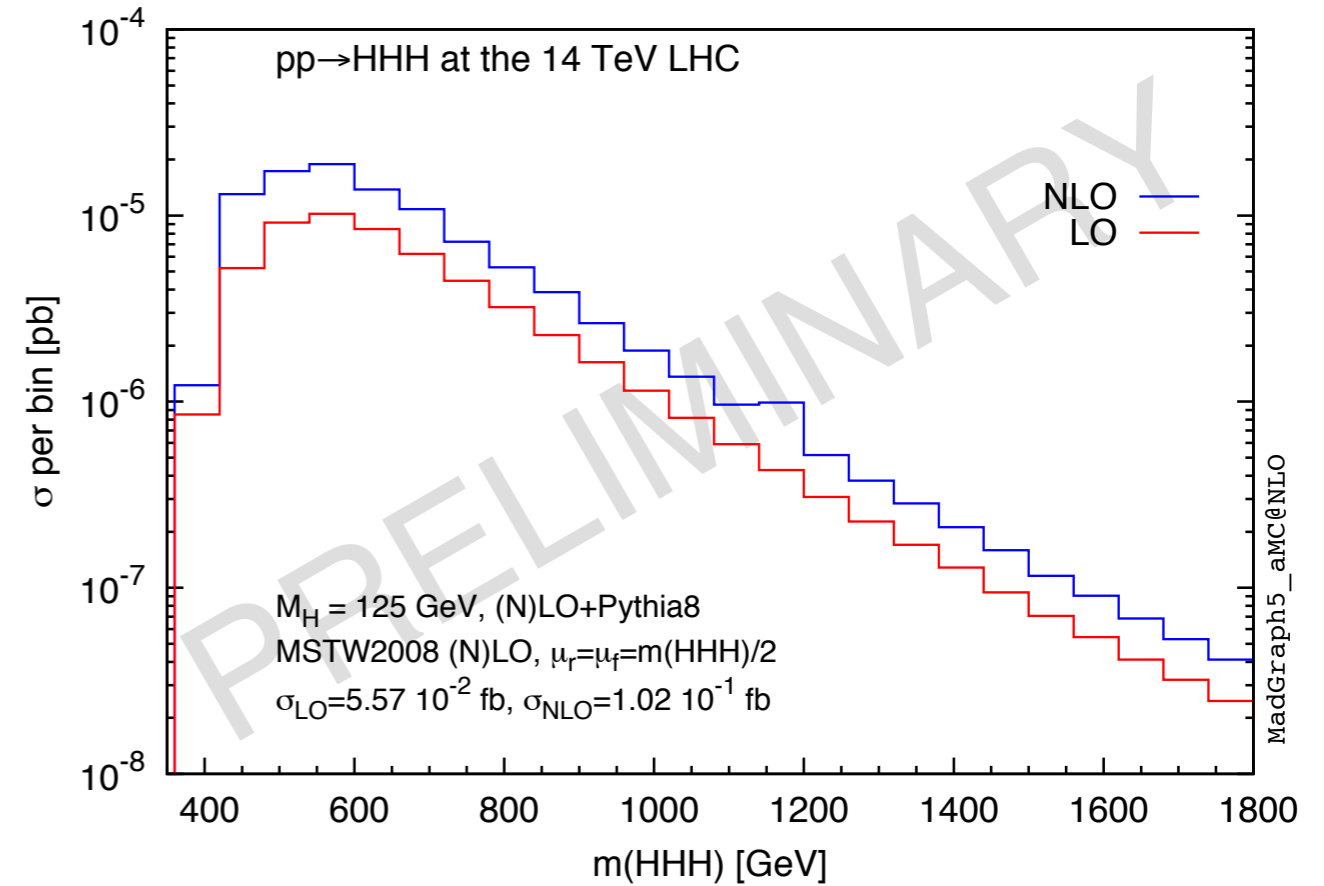
2HDM

$\tan \beta$	$\alpha/\pi$	$m_{H^0}$	$m_{A^0}$	$m_{H^\pm}$	$m_{12}^2$
1.50	-0.2162	700	701	670	180000



key process to be looked at the LHC

Maltoni, Vryonidou, MZ, in prep.



channels

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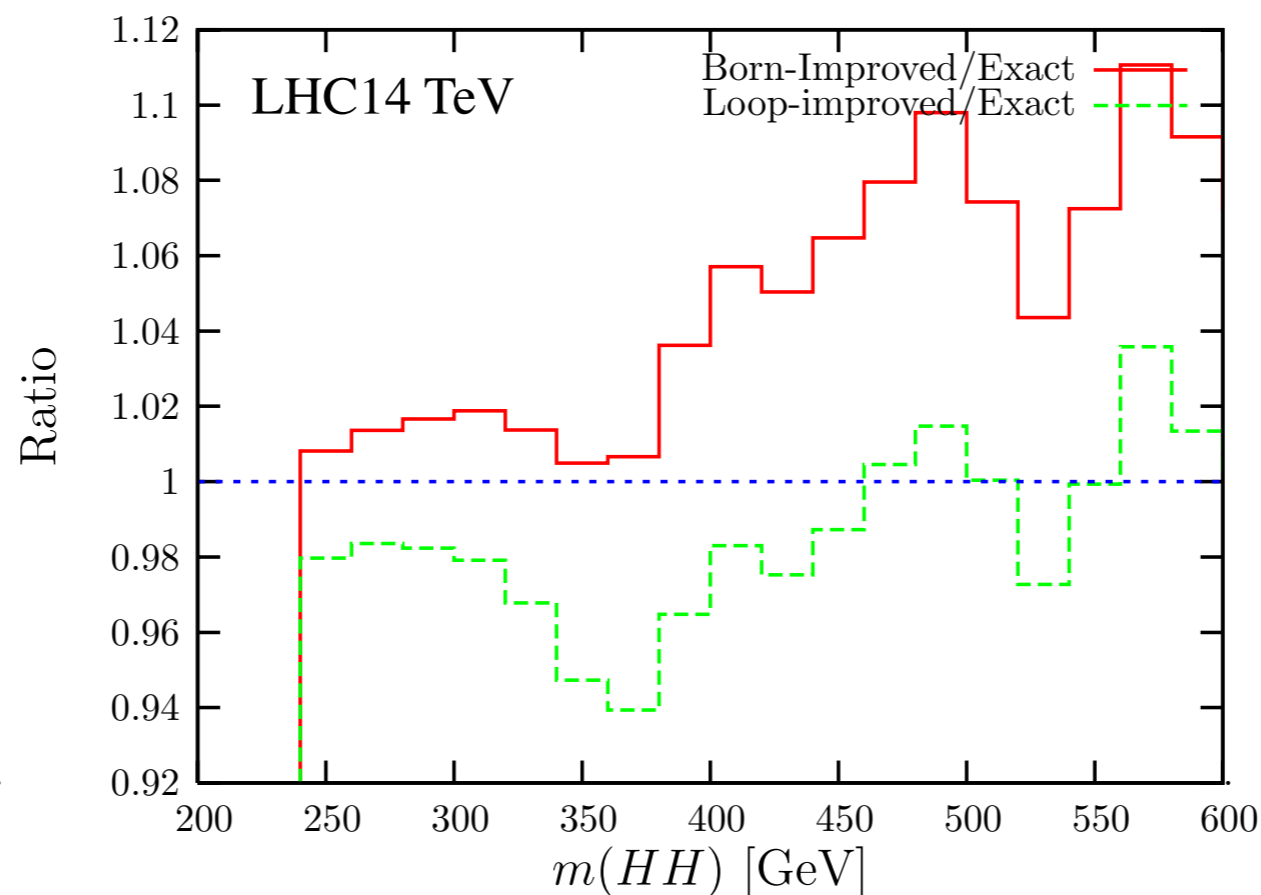




# Backup slides

# Different approximations for $gg \rightarrow HH$

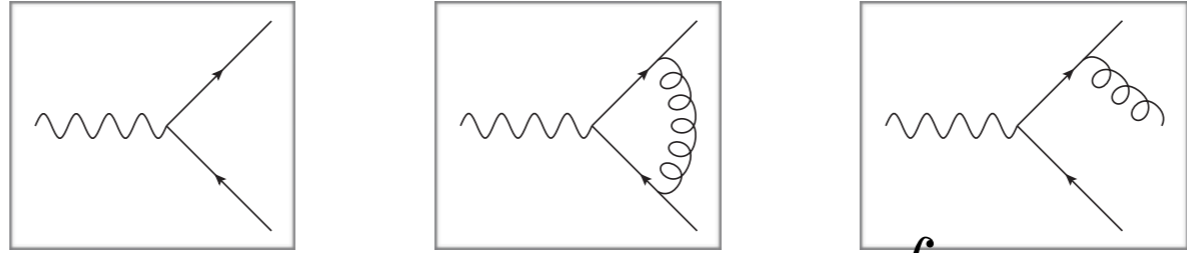
- Only consider terms  $\sim \lambda$
- Reweight everything with the Born ME (as HPAIR) or with the Born and real ME





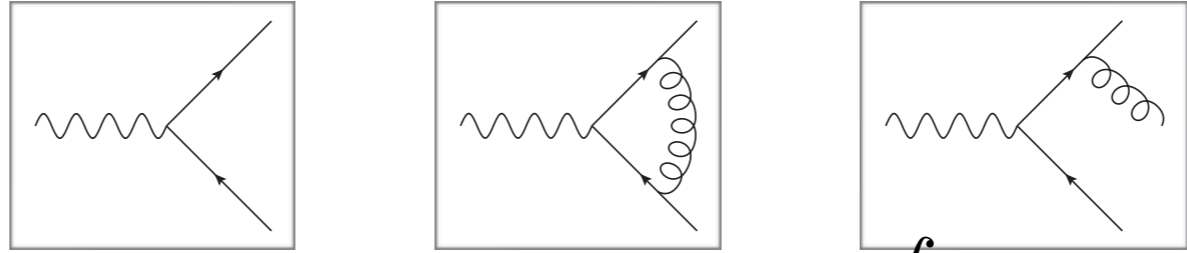
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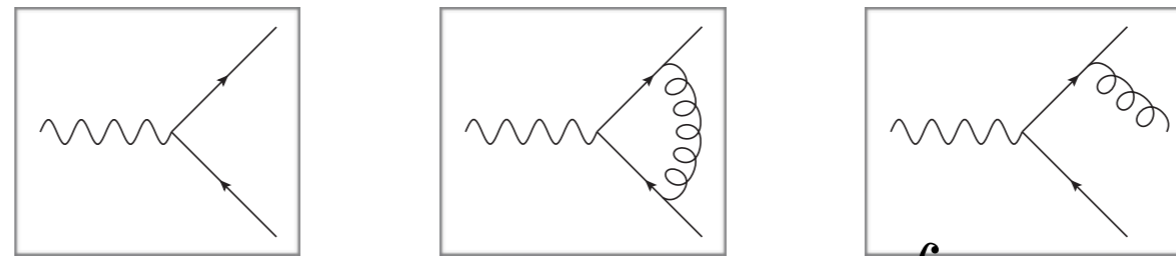
The equation is accompanied by three Feynman diagrams in boxes. The first diagram shows a wavy line (photon or gluon) splitting into two fermion lines. The second diagram shows a wavy line splitting into a fermion line and a loop of fermions. The third diagram shows a wavy line splitting into a fermion line and a gluon line, with a fermion line loop attached to the gluon line.

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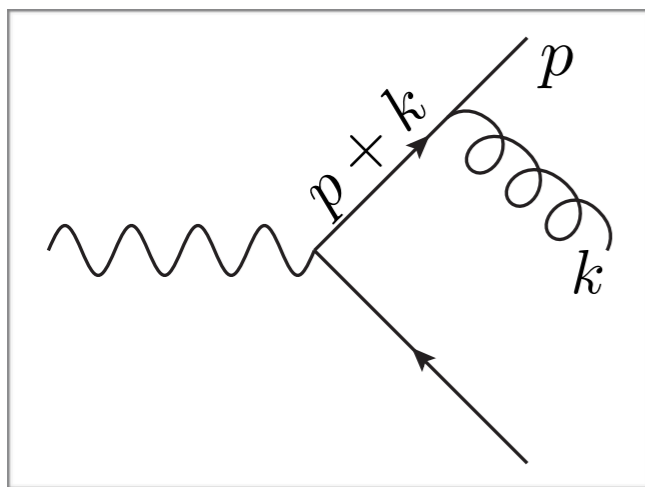
- Warning! Real emission ME is divergent!
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  - Need to cancel them before numerical integration (in  $D=4$ )

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- Structure of divergences is universal:

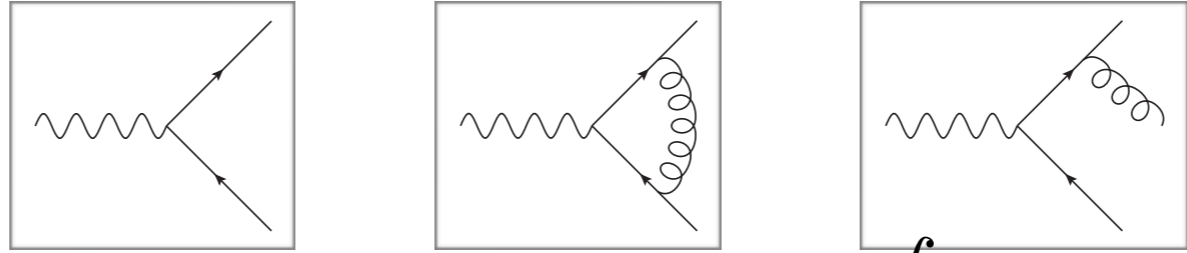


$$(p+k)^2 = 2E_p E_k (1 - \cos \theta_{pk})$$

$$\lim_{p//k} |M_{n+1}|^2 \simeq |M_n|^2 P^{AP}(z)$$

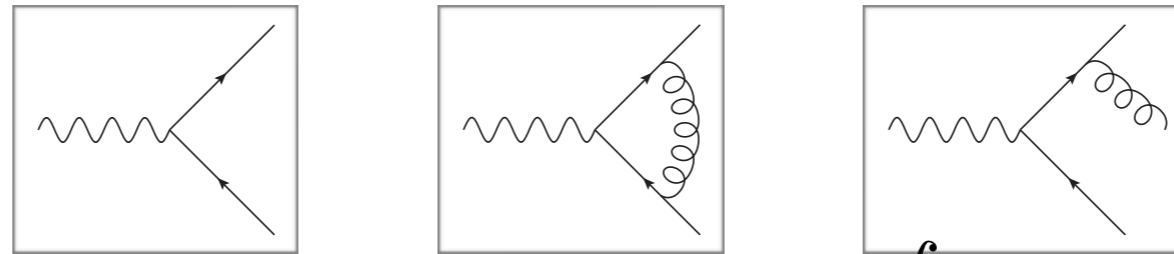
$$\lim_{k \rightarrow 0} |M_{n+1}|^2 \simeq \sum_{ij} |M_n^{ij}|^2 \frac{p_i p_j}{p_i k p_j k}$$

# NLO: how to?

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The equation is accompanied by three Feynman diagrams in boxes. The first diagram shows a wavy line (representing a photon or gluon) splitting into two straight lines (representing fermions). The second diagram shows a wavy line splitting into a fermion and a fermion loop. The third diagram shows a wavy line splitting into a fermion and a fermion with a gluon emission (represented by a curly line).

# NLO: how to?



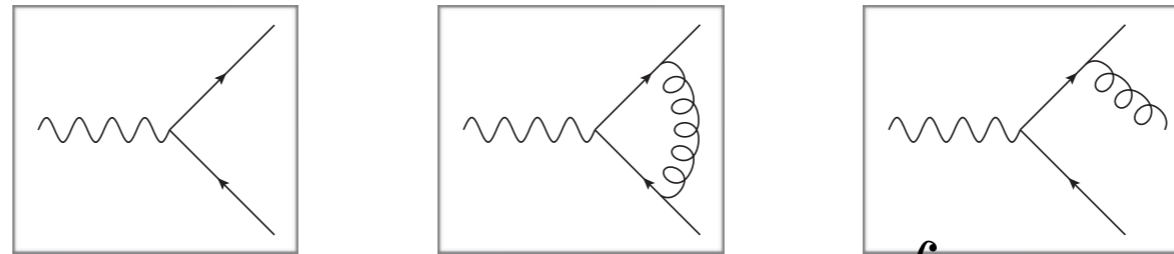
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- Add local counterterms in the singular regions and subtract its integrated finite part (poles will cancel against the virtuals)
- The  $n$  and  $n+1$  body integral now are finite in 4 dimension
  - Can be integrated numerically



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How to do this in an efficient way?

# The FKS subtraction

Frixione, Kunszt, Signer, arXiv:hep-ph/9512328

- Soft/collinear singularities arise in many PS regions
- Find parton pairs  $i, j$  that can give collinear singularities
- Split the phase space into regions with one collinear sing
  - Soft singularities are split into the collinear ones

$$|M|^2 = \sum_{ij} S_{ij} |M|^2 = \sum_{ij} |M|_{ij}^2 \quad \sum S_{ij} = 1$$

$$S_{ij} \rightarrow 1 \text{ if } k_i \cdot k_j \rightarrow 0 \quad S_{ij} \rightarrow 0 \text{ if } k_{m \neq i} \cdot k_{n \neq j} \rightarrow 0$$

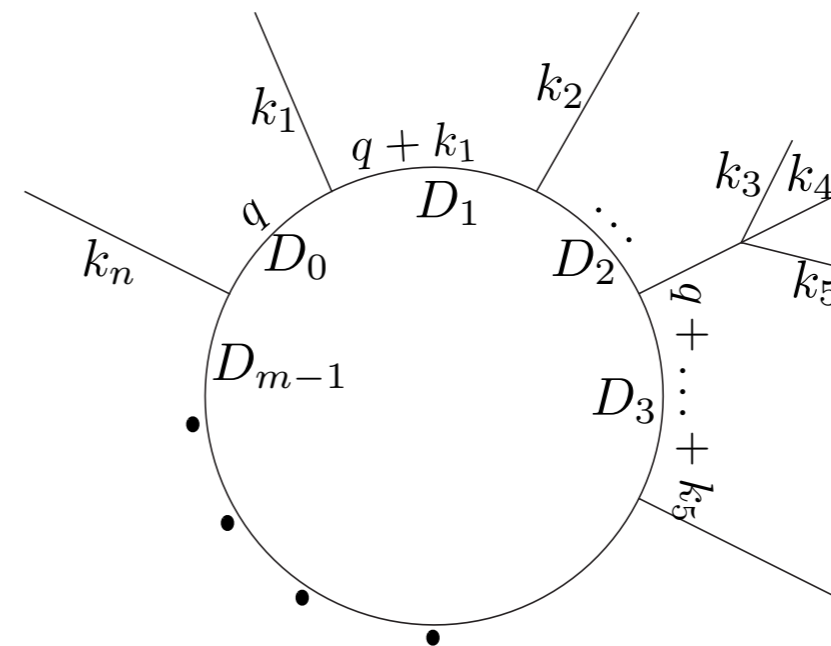
- Integrate them independently
  - Parallelize integration
  - Choose ad-hoc phase space parameterization
- Advantages:
  - # of contributions  $\sim n^2$
  - Exploit symmetries: 3 contributions for  $X \ Y \ > \ ng$

# Loops: the OPP Method

Ossola, Papadopoulos, Pittau, arXiv:hep-ph/0609007 & arXiv:0711.3596

- Passarino & Veltman reduction:
  - Write the amplitude at the integrand level as linear combination of 1-...-4-point scalar integrals

$$\begin{aligned}
 A(q) &= \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} d(i_0 i_1 i_2 i_3) D_0(i_0 i_1 i_2 i_3) \\
 &+ \sum_{i_0 < i_1 < i_2}^{m-1} c(i_0 i_1 i_2) C_0(i_0 i_1 i_2) \\
 &+ \sum_{i_0 < i_1}^{m-1} b(i_0 i_1) B_0(i_0 i_1) \\
 &+ \sum_{i_0}^{m-1} a(i_0) A_0(i_0) \\
 &+ R
 \end{aligned}$$



- Do this at the integrand level

# Loops: the OPP Method

Ossola, Papadopoulos, Pittau, arXiv:hep-ph/0609007 & arXiv:0711.3596

$$\begin{aligned}
 A(\bar{q}) &= \frac{N(q)}{\bar{D}_0 \bar{D}_1 \cdots \bar{D}_{m-1}} & N(q) &= \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3) \right] \prod_{i \neq i_0, i_1, i_2, i_3}^{m-1} D_i \\
 & & &+ \sum_{i_0 < i_1 < i_2}^{m-1} \left[ c(i_0 i_1 i_2) + \tilde{c}(q; i_0 i_1 i_2) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\
 & & &+ \sum_{i_0 < i_1}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i \\
 & & &+ \sum_{i_0}^{m-1} \left[ a(i_0) + \tilde{a}(q; i_0) \right] \prod_{i \neq i_0}^{m-1} D_i \\
 & & &+ \tilde{P}(q) \prod_i^{m-1} D_i.
 \end{aligned}$$

- Sample the numerator at complex values of the loop momenta in order to reconstruct the  $a, b, c, d$  coefficients and part of the rational terms (R1)
- Use CutTools: fed with the loop numerator outputs the coefficients of the scalar integrals and CC rational terms (R1)
- Add R2-rational terms/UV counterterms
  - Model dependent but process-independent

# Loop ME evaluation: MadLoop

Hirschi et al. arXiv:1103.0621

- Load the NLO UFO model
- Generate Feynman diagrams to evaluate the loop ME
- Add R2/UV renormalisation counter terms
- Interface to CutTools or to tensor reduction programs (in progress)
- Check PS point stability (and switch to QP if needed)
- Improved with the OpenLoops method Cascioli, Maierhofer, Pozzorini  
arXiv:1111.5206
- And much more (can be used as standalone or external OLP via the BLHA, handle loop-induced processes, ...)

# Matching in MC@NLO

- Use suitable counterterms to avoid double counting the emission from shower and ME, keeping the correct rate at order  $\alpha_s$ :

$$\frac{d\sigma_{MC@NLO}}{dO} = \underbrace{\left( \mathcal{B} + \mathcal{V} + \int d\Phi_1 MC \right) d\Phi_n I_{MC}^n(O)}_{\text{S-events}} + \underbrace{(\mathcal{R} - MC) d\Phi_n d\Phi_1 I_{MC}^{n+1}(O)}_{\text{H-events}}$$

- MC depends on the PSMC's Sudakov:

$$MC = \left| \frac{\partial (t^{MC}, z^{MC}, \phi)}{\partial \Phi_1} \right| \frac{1}{t^{MC}} \frac{\alpha_s}{2\pi} \frac{1}{2\pi} P(z^{MC}) \mathcal{B}$$

- Available for Herwig6, Pythia6 (virtuality-ordered), Herwig++, Pythia8 (in the new release)
- MC acts as local counterterm
- Some weights can be negative (unweighting up to sign)
  - Only affects statistics