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Off-shell effects and signal-background interference in  $H \rightarrow VV$  at the (I)LC



Orsay - 22 July 2014

University of Hamburg



Aim of this talk: Discuss LHC inspired effects for linear collider:

▷ 1. Off-shell contributions in  $H \rightarrow VV^{(*)}$

[1206.4803; Kauer Passarino:

Inadequacy of zero-width approximation  
for a light H boson signal]

[1305.2092, 1310.7011; Kauer:

Interference effects for  $H \rightarrow WW/ZZ \rightarrow l\bar{\nu}_l\bar{l}\nu_l$   
searches in gluon fusion at the LHC]

[1307.4935; Caola Melnikov: Constraining the Higgs boson width with  $ZZ$  production at the LHC]

Further elaboration: [1311.3589, 1312.1628; Campbell Ellis Williams], Talk by K. Ellis!

Application: CMS [CMS-PAS-HIG-14-002, 1405.3455], Talk by M. Dalchenko!

ATLAS [ATLAS-CONF-2014-042]

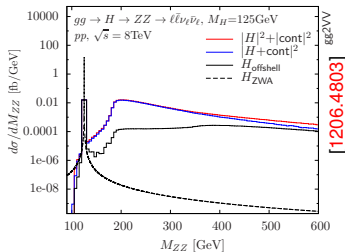
Further comments: [1310.1397, 1405.0285, 1405.1925, 1406.1757, 1406.6338]

▷ 2. Interferometry with background in  $H \rightarrow \gamma\gamma$

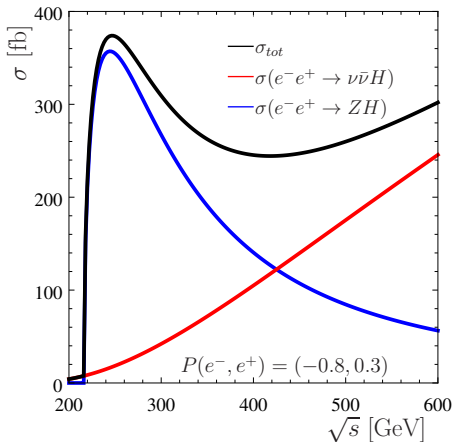
[1208.1533, 1303.3342; Martin: Shift in the  $H \rightarrow \gamma\gamma$  mass peak from interference with background]

Further elaboration: [1303.1397; de Florian et al., 1305.3854; Dixon Li]

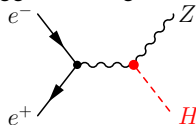
→ Can also be investigated at the (I)LC!



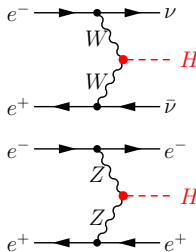
## Main production mechanisms of the SM Higgs at the (I)LC:



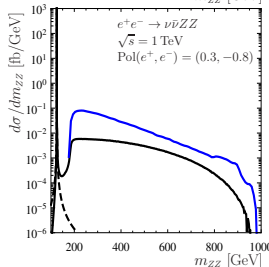
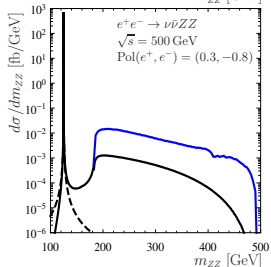
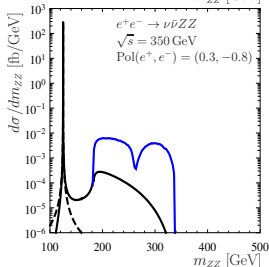
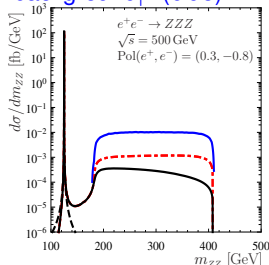
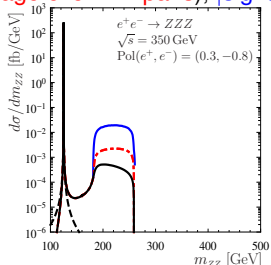
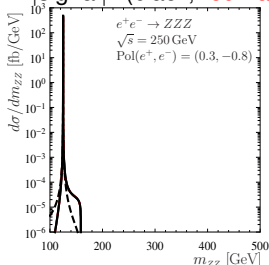
### Higgsstrahlung



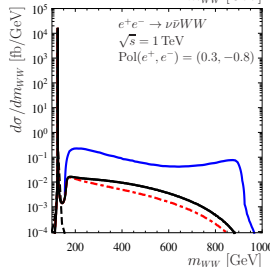
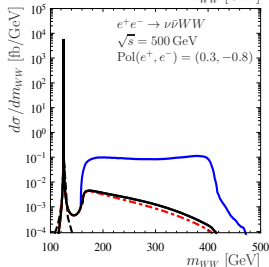
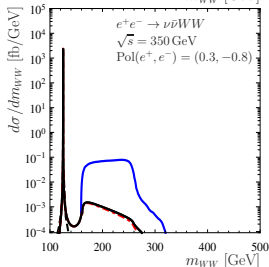
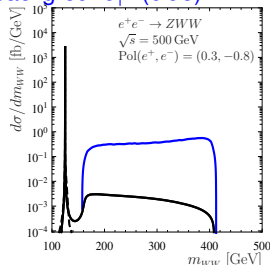
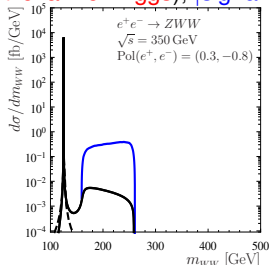
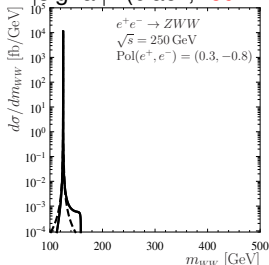
### Vector boson fusion



Quantification for  $H \rightarrow ZZ^{(*)}$  as function of  $\sqrt{s}$ :

 $|\text{signal}|^2$  (black, red - average over  $ZZ$  pairs),  $|\text{signal} + \text{background}|^2$  (blue)


Quantification for  $H \rightarrow WW^{(*)}$  as function of  $\sqrt{s}$ :

 $|\text{signal}|^2$  (black, red - with  $t$ -channel Higgs),  $|\text{signal} + \text{background}|^2$  (blue)


Relative contribution:  $\text{Pol}(e^+, e^-) = (0.3, -0.8)$

With  $\sigma_X(m_{VV}^d, m_{VV}^u) = \int_{m_{VV}^d}^{m_{VV}^u} dm_{VV} \left( \frac{d\sigma_X}{dm_{VV}} \right)$  we define

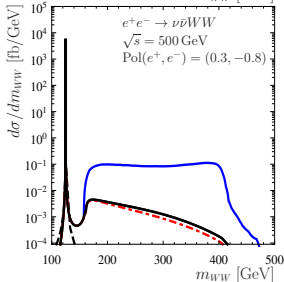
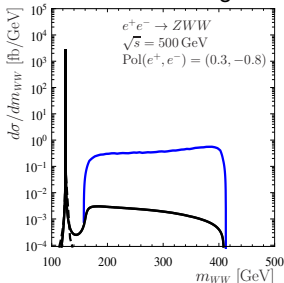
$$\Delta_{\text{off}}^{ZVV} = \frac{\sigma_{\text{off}}^{ZVV}(130\text{GeV}, \sqrt{s} - m_Z)}{\sigma_{\text{off}}^{ZVV}(0, \sqrt{s} - m_Z)} \quad \text{and} \quad \Delta_{\text{off}}^{\nu\bar{\nu}VV} = \frac{\sigma_{\text{off}}^{\nu\bar{\nu}VV}(130\text{GeV}, \sqrt{s})}{\sigma_{\text{off}}^{\nu\bar{\nu}VV}(0, \sqrt{s})}$$

$\sqrt{s}$	$\sigma_{\text{off}}^{ZZZ}$	$\Delta_{\text{off}}^{ZZZ}$	$\sigma_{\text{off}}^{\nu\bar{\nu}ZZ}$	$\Delta_{\text{off}}^{\nu\bar{\nu}ZZ}$
250 GeV	3.12(3.12) fb	0.03(0.03) %	0.490 fb	0.12 %
350 GeV	1.71(1.82) fb	1.82(7.77) %	1.91 fb	0.88 %
500 GeV	0.802(0.981) fb	7.20(24.1) %	4.78 fb	2.96 %
1 TeV	0.242(0.341) fb	30.9(50.9) %	15.0 fb	13.0 %
$\sqrt{s}$	$\sigma_{\text{off}}^{ZWW}$	$\Delta_{\text{off}}^{ZWW}$	$\sigma_{\text{off}}^{\nu\bar{\nu}WW}$	$\Delta_{\text{off}}^{\nu\bar{\nu}WW}$
250 GeV	76.3 fb	0.03 %	3.98(3.99) fb	0.13(0.12) %
350 GeV	41.4 fb	0.92 %	15.5(15.5) fb	0.49(0.43) %
500 GeV	18.6 fb	2.61 %	38.1(38.1) fb	1.21(0.96) %
1 TeV	4.58 fb	11.0 %	110.8(108.9) fb	4.45(2.78) %

Comments:

- ▷  $\Delta_{\text{off}}$  independent of the polarisation.
- ▷ For  $H \rightarrow ZZ \rightarrow 4l$  off-shell contributions accessible by  $m_{4l}$ .
- ↔ For  $H \rightarrow WW \rightarrow 2l2\nu$  not directly accessible! ↔ Coupling extraction!
- ▷ Important: On-shell XS strongly dependent on Higgs mass, off-shell not!

Comment on the background:


 Inclusive cross sections for  $m_{VV} > 130 \text{ GeV}$   
 for  $\text{Pol}(e^+, e^-) = (0.3, -0.8)$ :

$\sqrt{s}$	$\sigma_{\text{all}}^{ZZZ}$	$\Delta_{\text{SB}}^{ZZZ}$	$\sigma_{\text{all}}^{\nu\bar{\nu}ZZ}$	$\Delta_{\text{SB}}^{\nu\bar{\nu}ZZ}$
250 GeV	---	---	1.51 fb	0.04 %
350 GeV	1.19 fb	2.62(11.9) %	1.66 fb	1.01 %
500 GeV	2.06 fb	2.83(11.6) %	2.85 fb	4.96 %
1 TeV	1.71 fb	4.40(10.2) %	16.7 fb	11.6 %
$\sqrt{s}$	$\sigma_{\text{all}}^{ZWW}$	$\Delta_{\text{SB}}^{ZWW}$	$\sigma_{\text{all}}^{\nu\bar{\nu}WW}$	$\Delta_{\text{SB}}^{\nu\bar{\nu}WW}$
250 GeV	---	---	0.05 fb	9.87(9.87) %
350 GeV	29.2 fb	1.30 %	6.44 fb	1.18(1.03) %
500 GeV	91.8 fb	0.53 %	22.4 fb	2.05(1.63) %
1 TeV	136.7 fb	0.37 %	67.3 fb	7.31(4.49) %

 $\Delta_{\text{SB}} \leftrightarrow$  Signal/Background in off-shell region.

 Naturally: Very large interference term  
 guarantees unitarity in  $WW \rightarrow WW$ !

What can be done with the (basically  $m_H$  independent) off-shell contributions?

- ▷ They are needed for and allow to test unitarity in  $WW \rightarrow WW$ !
- ▷ They allow to test the influence of higher dimensional operators and thus can probe composite Higgs scenarios!

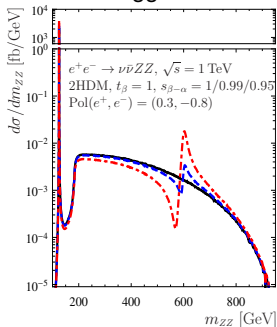
see e.g. [[hep-ph/0301097](#), Barger Han et al.]

Current study with e.g.  $e^+e^- \rightarrow \nu\bar{\nu} + 4\text{jets}$ : [[1309.7038](#), Contino Grojean et al.]

- ▷ They can test extended Higgs sectors!

- ▷ In the pure SM (without NLO effects) they allow to set a bound on  $\Gamma_H$ !

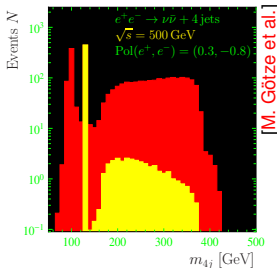
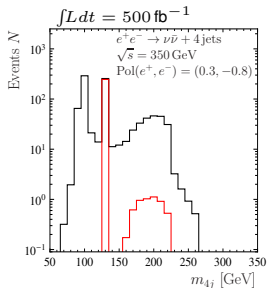
Note:  $\Gamma_H$  determination from  $Z$  recoil method is safe from off-shell effects for low  $\sqrt{s}$ !



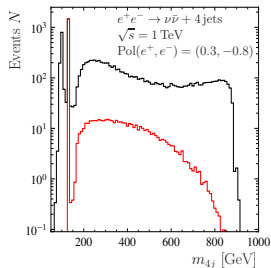


Bounding the Higgs width using e.g.  $e^+e^- \rightarrow \nu\bar{\nu} + 4\text{jets}$ :

MadGraph with  $\Delta_{R,j} > 0.4$ ,  $|y_j| < 5$ ,  $p_{T,j} > 20$  GeV,  $p_{T,4j} > 75$  GeV



[M. Götz et al.]

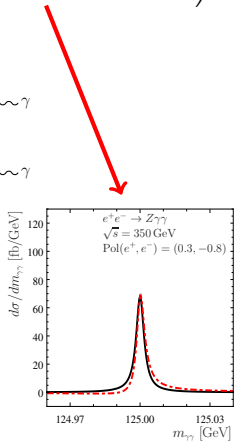
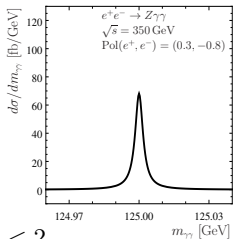
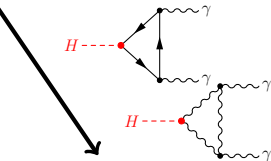
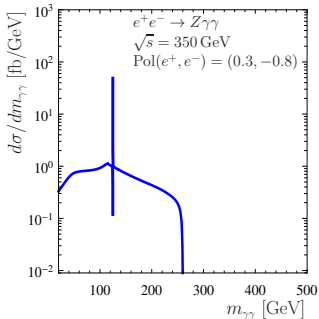
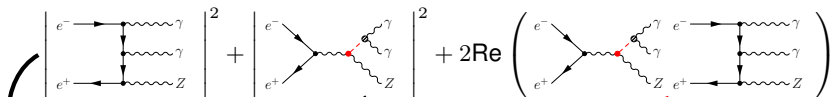


Rescaling couplings and the width (assuming pure SM!!!):

$$N(r) = N_0(1 + R_1\sqrt{r} + R_2r) + N_B \quad \text{with} \quad r = \Gamma_H/\Gamma_H^{SM}$$

$\sqrt{s}$	350 GeV	500 GeV	1 TeV
$N_0$ ( $\int L dt = 500 \text{ fb}^{-1}$ )	263	1775	8420
$R_1$	-0.017	-0.010	-0.098
$R_2$	0.026	0.019	0.048
Limit on $r$ ( $\int L dt = 500 \text{ fb}^{-1}$ )	4.1	2.5	2.3
Limit on $r$ ( $\int L dt = 1 \text{ ab}^{-1}$ )	3.2	2.1	2.0

Main limitation:  
 Negative interference!  
 In contrast to LHC:  
 Pure tree-level processes!

2. Interferometry with the background in  $H \rightarrow \gamma\gamma$ 

 Applied cuts:  $E_\gamma > 20 \text{ GeV}$ ,  $|\eta_\gamma| < 2$

Interferometry with the background in  $H \rightarrow \gamma\gamma$ 

$$\frac{d\sigma^{sig}}{dm_{\gamma\gamma}} = \frac{S}{(m_{\gamma\gamma}^2 - m_H^2)^2 + m_H^2 \Gamma_H^2} \rightarrow \sigma^{sig} = \frac{\pi S}{2m_H^2 \Gamma_H}$$

$$\frac{d\sigma^{int}}{dm_{\gamma\gamma}} = \frac{(m_{\gamma\gamma}^2 - m_H^2)R + m_H \Gamma_H I}{(m_{\gamma\gamma}^2 - m_H^2)^2 + m_H^2 \Gamma_H^2} \rightarrow \sigma^{int} = \frac{\pi I}{2m_H}$$

Relevant part:  $R$  induces shift of the peak without changing the incl. XS!

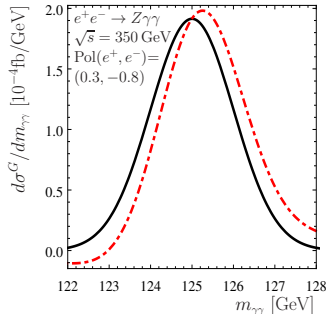
Smearing due to detector resolution:

Gaussian  $G$  with e.g.  $\hat{\sigma}^G = 1$  GeV

$$\frac{d\sigma^G}{dm_{\gamma\gamma}} = \int_0^\infty dm'_{\gamma\gamma} G(m_{\gamma\gamma} - m'_{\gamma\gamma}, \hat{\sigma}^G) \frac{d\sigma}{dm'_{\gamma\gamma}}$$

→ Visible shift  $\Delta m_H$  of the mass peak!

Depending on  $\hat{\sigma}^G$ ,  $E_\gamma$ ,  $\eta_\gamma$ ,  $\sqrt{s}$ ,  $\delta_\gamma$ , (Pol).

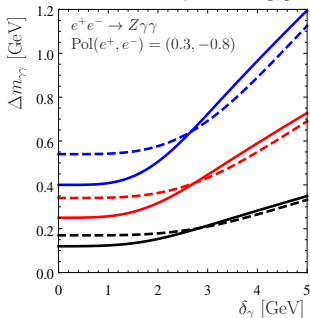


Mimic the method of peak extraction:

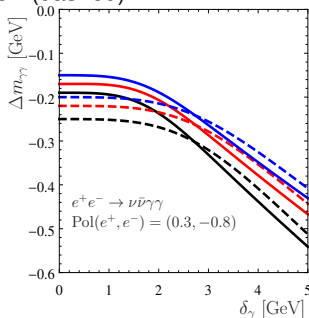
$$\langle m_{\gamma\gamma} \rangle_{\delta, X} = \frac{1}{N} \int_{m_p - \delta}^{m_p + \delta} dm_{\gamma\gamma} m_{\gamma\gamma} \frac{d\sigma_X^G}{dm_{\gamma\gamma}} \quad \rightarrow \Delta m_{\gamma\gamma} = \langle m_{\gamma\gamma} \rangle_{\delta_\gamma, S+I} - \langle m_{\gamma\gamma} \rangle_{\delta_\gamma, S}$$

Obtain  $\langle m_{\gamma\gamma} \rangle_{\delta_\gamma, S}$  from different cuts or other final states.

$\hat{\sigma}^G = 1 \text{ GeV}$  (solid),  $1.5 \text{ GeV}$  (dashed)



$\sqrt{s} = 250 \text{ GeV}, 350 \text{ GeV}, 500 \text{ GeV}$



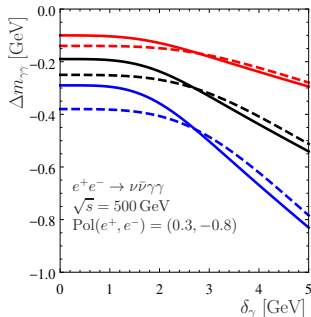
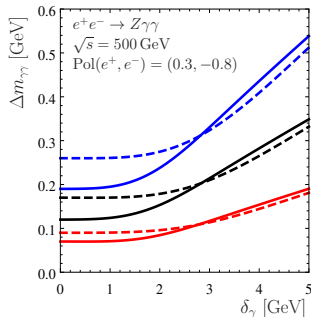
$\sqrt{s} = 350 \text{ GeV}, 500 \text{ GeV}, 1 \text{ TeV}$

## Higgs width dependence?

Perform similar rescaling of the couplings  $g_{HZZ}, g_{HWW}, g_{HAA}$  and the width  $\Gamma_H$  to keep  $\sigma_{ZWA}$  constant.

$$\hat{\sigma}^G = 1 \text{ GeV (solid), } 1.5 \text{ GeV (dashed)}$$

$$\Gamma_H = 1 \text{ MeV, } 4.07 \text{ MeV, } 15 \text{ MeV}$$



Further studies: Perform analysis with detector simulation?!

## Conclusions:

- ▶ Off-shell contributions in  $H \rightarrow VV^{(*)}$  are naturally large at a linear collider (except  $\sqrt{s}$  is below 300 GeV). Dependent on the assumptions, they can be used to test unitarity, higher dimensional operators, extended Higgs sectors or to set a bound on  $\Gamma_H$ .
- ▶ LC offers unique possibility to measure  $\Gamma_H$  through  $Z$  recoil method in  $e^+e^- \rightarrow ZH$  at 250 GeV, which is safe from off-shell contributions.
- ▶ Signal-background interference in  $H \rightarrow \gamma\gamma$  shifts the mass peak by a few 100 MeV! The shift also allows to access  $\Gamma_H$ .
- ▶ For all purposes a well determined Higgs mass is necessary.

Conclusions II: Precision machine LC now or never!

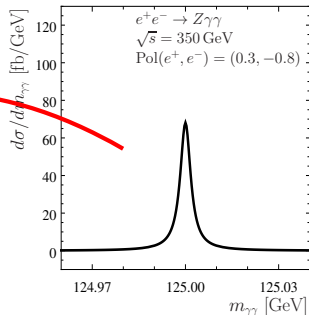
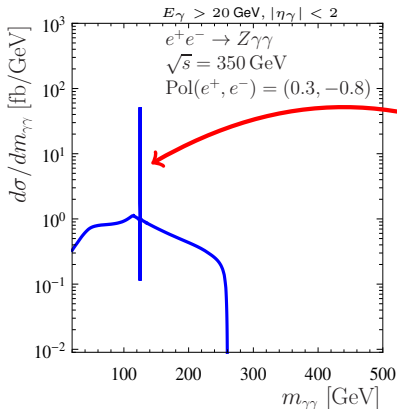
Thank you for your attention!

How to obtain information about the total Higgs width  $\Gamma_H$ ?

→ Measure the Breit-Wigner peak e.g. in  $H \rightarrow \gamma\gamma$ ?

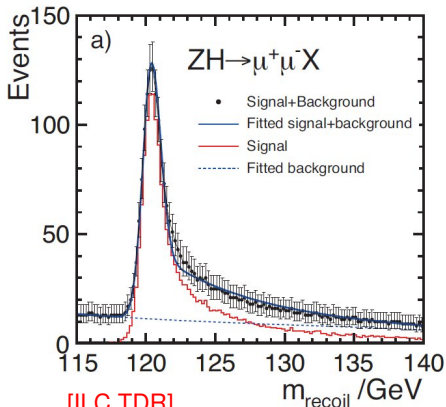
$$\frac{d\sigma_{\text{ZWA}}^{Z\gamma\gamma}}{dm_{\gamma\gamma}} = \sigma^{ZH}(m_H) \frac{2m_{\gamma\gamma}}{(m_{\gamma\gamma}^2 - m_H^2)^2 + m_H^2 \Gamma_H^2} \frac{m_H \Gamma_{H \rightarrow \gamma\gamma}(m_H)}{\pi}$$

Problem:  $m_H = 125 \text{ GeV} \leftrightarrow \Gamma_H = 4.07 \text{ MeV}$       →  $\sigma_{\text{ZWA}}^{Z\gamma\gamma} = \sigma^{ZH} \frac{\Gamma_{H \rightarrow \gamma\gamma}}{\Gamma_H}$



→ Detector resolution smears out the Breit-Wigner peak!

→ LC unique method: Higgs width  $\Gamma_H$  through the Z recoil at  $\sqrt{s} = 250$  GeV



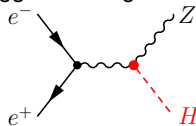
[ILC TDR]

$250 \text{ fb}^{-1} @ 250 \text{ GeV}$

$\Delta\sigma_P / \sigma_P = 2.5\%$

$\Delta m_H = 30 \text{ MeV}$

Higgsstrahlung



Observe:  $Z \rightarrow \mu^+ \mu^-$

Reconstruct:

$$\sigma_P = \sigma(e^+e^- \rightarrow HZ) \propto g_{HZZ}^2$$

(needs defined initial state)

Obtain absolute BR:

$$\text{BR}(H \rightarrow X) = (\sigma_P \text{BR}_X) / \sigma_P$$

Reconstruct (example):

$$\Gamma_H \propto \Gamma(H \rightarrow ZZ) / \text{BR}(H \rightarrow ZZ)$$

$$\propto g_{HZZ}^2 / \text{BR}(H \rightarrow Z)$$

Details: [1311.7155: Han, Liu, Sayre]

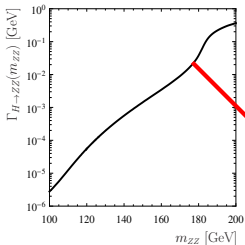


## Discussion of off-shell contributions $m_{ZZ} > 2m_Z$ in $H \rightarrow ZZ$

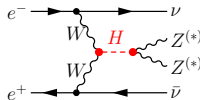
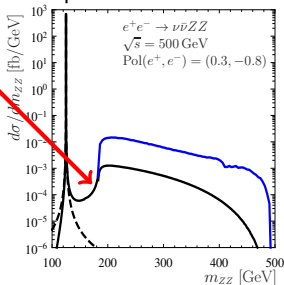
$$\left( \frac{d\sigma_{ZWA}^{\nu\bar{\nu}ZZ}}{dm_{ZZ}} \right) = \sigma^{\nu\bar{\nu}H}(m_H) \frac{2m_{ZZ}}{(m_{ZZ}^2 - m_H^2)^2 + (m_H\Gamma_H)^2} \frac{m_H\Gamma_{H\rightarrow ZZ}(m_H)}{\pi}$$

$$\left( \frac{d\sigma_{\text{off}}^{\nu\bar{\nu}ZZ}}{dm_{ZZ}} \right) = \sigma^{\nu\bar{\nu}H}(m_{ZZ}) \frac{2m_{ZZ}}{(m_{ZZ}^2 - m_H^2)^2 + (m_H\Gamma_H)^2} \frac{m_{ZZ}\Gamma_{H\rightarrow ZZ}(m_{ZZ})}{\pi}$$

Second equation describes the proper calculation of  $e^+e^- \rightarrow \nu\bar{\nu}ZZ$  at LO!



Consequences:



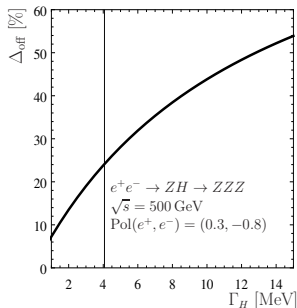
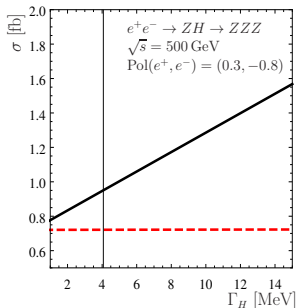
How can the width be determined from off-shell contributions?

$$\sigma_{\text{ZWA}}^{\text{ZZZ}} = \sigma^{\text{ZH}}(m_H) \frac{\Gamma_{H \rightarrow \text{ZZ}}(m_H)}{\Gamma_H} \propto \frac{g_{\text{HZZ}}^4}{\Gamma_H}$$

Rescaling  $g'_{\text{HZZ}} = \xi g_{\text{HZZ}}$ ,  $\Gamma'_H = \xi^4 \Gamma_H$  does not change  $\sigma_{\text{ZWA}}^{\text{ZZZ}}$ !

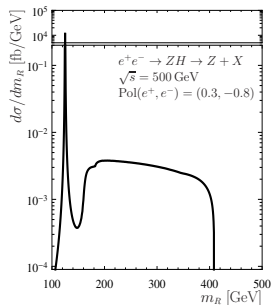
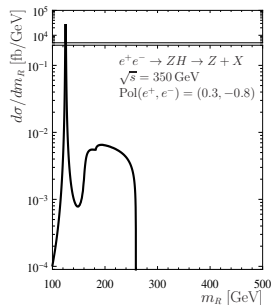
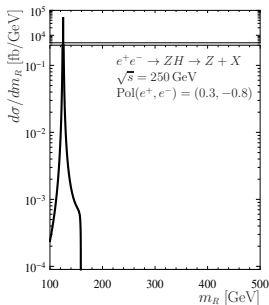
→ Vary  $\Gamma_H$  (in reasonable interval!) and leave  $\sigma_{\text{ZWA}}$  constant!

→ Off-shell contributions  $\propto g_{\text{HZZ}}^4 \rightarrow \Delta_{\text{off}}$  changes!.



Can the off-shell cont. be discriminated from the background?

### Off-shell contributions in the $Z$ recoil method:



### Recoil mass:

$$m_R^2 = s + \hat{m}_Z^2 - 2E_Z\sqrt{s}$$

$\sqrt{s}$	250 GeV	300 GeV	350 GeV	500 GeV	1 TeV
$\Delta_{\text{off}}$	0.02%	0.12%	0.30%	0.91%	1.84%