



Bounding the Higgs boson width from
off-shell production and decay to $ZZ \rightarrow 4\ell$
and $ZZ \rightarrow 2\ell 2\nu$

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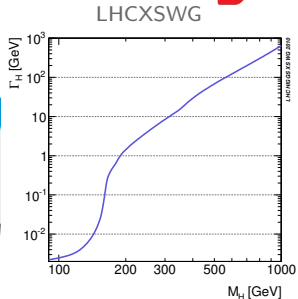


The Higgs boson width



The total decay width is a *fundamental* parameter

- Relates to the couplings to all massive particles
⇒ sensitive to the BSM physics
- Hard to measure: $\Gamma_H \simeq 4 \text{ MeV} @ m_H \simeq 126 \text{ GeV}$

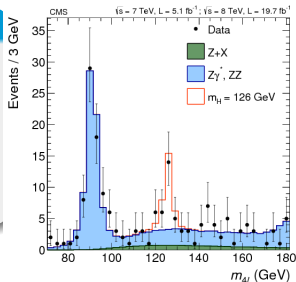


Direct constraints

- Current (excellent) experimental resolution is too large for direct measurements: $\sigma_m \sim \mathcal{O}(1) \text{ GeV}$
- Constraint from the resonance width in $H \rightarrow ZZ \rightarrow 4\ell$: $\Gamma_H < 3.4 \text{ GeV} @ 95\% \text{ CL}$
- Constraint from the resonance width in $H \rightarrow \gamma\gamma$: $\Gamma_H < 2.4 \text{ GeV} @ 95\% \text{ CL}$

(see S. Chhibra and C. Martin talks)

PRD 89 (2014) 092007





Width constraints from off-shell Higgs



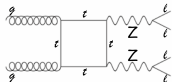
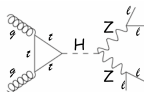
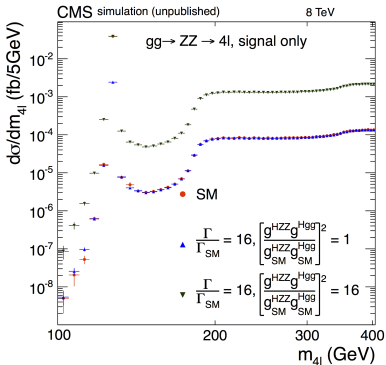
- Off-shell Higgs boson production is small but the BR to 2 real Z is large above $2m_Z$
- Under change of width, peak yield remains constant if we scale the couplings in appropriate way

$$\sigma_{pp \rightarrow H \rightarrow ZZ} \sim \frac{g_{Hgg}^2 g_{HZZ}^2}{\Gamma}$$

- Off-shell Higgs yield is proportional to this scale

$$\sigma_{pp \rightarrow H \rightarrow ZZ} \sim g_{Hgg}^2 g_{HZZ}^2$$

- Interference between the signal and background is sizeable and must be accounted for



(see D. de Florian talk)



Analysis overview



- Use $gg \rightarrow 4\ell$ (GG2VV and MCFM for $ZZ \rightarrow 4\ell$, GG2VV only for $ZZ \rightarrow 2\ell 2\nu$) and $VV \rightarrow 4\ell$, $2\ell 2\nu$ (Phantom) MC samples
- Build probability templates for signal, background and interference using $m_{4\ell}$ and a kinematic discriminant D_{gg} in case of $ZZ \rightarrow 4\ell$
$$P_{\text{total}} = \mu r P_{gg \rightarrow H \rightarrow 4\ell} + \sqrt{\mu r} P_{\text{interference}} + P_{gg \rightarrow 4\ell}$$

 $r = \Gamma / \Gamma_{\text{SM}}$, μ - signal strength
- Build similar probability templates for VBF
- Build similar probability templates for signal, background and interference using $m_{\ell\ell}^T$ in case of $ZZ \rightarrow 2\ell 2\nu$
- Perform a combined fit of the off-shell and on-shell regions



4 ℓ analysis

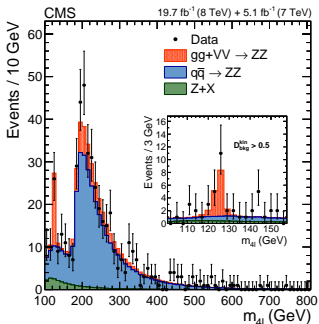


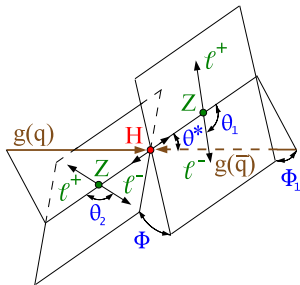
Selection settings and reconstruction are identical to PRD 89 (2014) 092007

- Two pairs of OS/SF leptons. Z_1 closest to the Z mass, Z_2 the remaining with highest sum of p_t
- $40 < m_{Z_1} < 120\text{GeV}$ $12 < m_{Z_2} < 120\text{GeV}$
- One lepton with $p_t > 20 \text{ GeV}/c$, another with $p_t > 10 \text{ GeV}/c$.
- $m_{4\ell} > 100 \text{ GeV}$, $m_{l+l-} > 4 \text{ GeV}$

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- Reducible background ("Z + X") estimated from data in control regions with Z_1 + at least 1 loose lepton.
- Irreducible bkg calculated from MC, $q\bar{q}$ annihilation from POWHEG
- Phenomenological model for the $q\bar{q} \rightarrow ZZ$ shape





- ggMELA discriminant was developed in the context of the PRD 89 (2014) 092007
- High performances for separating $gg \rightarrow ZZ$ from $q\bar{q} \rightarrow ZZ$ where $gg \rightarrow ZZ$ includes signal, continuum and their interference for any relative signal strength a .

Built from signal and background probabilities: $D_{gg,a} = \frac{P_{gg,a}}{P_{gg,a} + P_{q\bar{q},a}}$, where $P_{gg,a} = a \times P_{\text{sig}}^{gg} + \sqrt{a} \times P_{\text{int}}^{gg} + P_{\text{bkg}}^{gg}$ and $P_{q\bar{q},a} = P_{\text{bkg}}^{q\bar{q}}$

- Signal strength a must be chosen when building the discriminant.
- From preliminary studies we expected sensitivity for run 1 data to be around $10 \times \text{SM}$, so we chose $D_{gg,10}$.

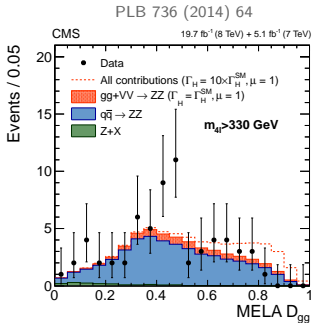
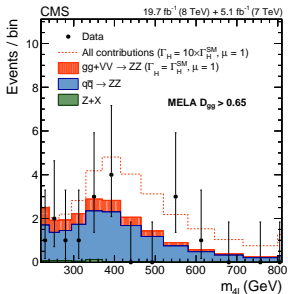
About 30% improvement when including it in the fit procedure



Distributions of selected events

$$H \rightarrow ZZ \rightarrow 4\ell$$

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signal-enriched region

■ $m_{4\ell} > 330 \text{ GeV}$

■ $D_{gg} > 0.65$

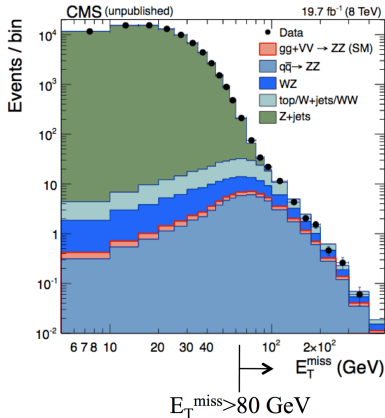
	4ℓ	$2\ell 2\nu$
(a) total gg ($\Gamma_H = \Gamma_H^{\text{SM}}$)	1.8 ± 0.3	9.6 ± 1.5
gg signal component ($\Gamma_H = \Gamma_H^{\text{SM}}$)	1.3 ± 0.2	4.7 ± 0.6
gg background component	2.3 ± 0.4	10.8 ± 1.7
(b) total gg ($\Gamma_H = 10 \times \Gamma_H^{\text{SM}}$)	9.9 ± 1.2	39.8 ± 5.2
(c) total VBF ($\Gamma_H = \Gamma_H^{\text{SM}}$)	0.23 ± 0.01	0.90 ± 0.05
VBF signal component ($\Gamma_H = \Gamma_H^{\text{SM}}$)	0.11 ± 0.01	0.32 ± 0.02
VBF background component	0.35 ± 0.02	1.22 ± 0.07
(d) total VBF ($\Gamma_H = 10 \times \Gamma_H^{\text{SM}}$)	0.77 ± 0.04	2.40 ± 0.14
(e) qq background	9.3 ± 0.7	47.6 ± 4.0
(f) other backgrounds	0.05 ± 0.02	35.1 ± 4.2
(a+c+e+f) total expected ($\Gamma_H = \Gamma_H^{\text{SM}}$)	11.4 ± 0.8	93.2 ± 6.0
(b+d+e+f) total expected ($\Gamma_H = 10 \times \Gamma_H^{\text{SM}}$)	20.1 ± 1.4	124.9 ± 7.8
observed	11	91



2l2ν analysis



- 6 times higher branching ratio w.r.t. 4l final state
- Use 4l results in the on-shell region
- High Z + jets background (fake E_T^{miss} from hadronic energy mismeasurement)
- Other backgrounds:
 - Irreducible: ZZ, WZ (from MC)
 - Non-resonant: top, WW
- Analysis variable: transverse mass



$$M_T^2 = \left[\sqrt{p_{T,\ell\ell}^2 + m_{\ell\ell}^2} + \sqrt{E_{miss,T}^2 + m_{\ell\ell}^2} \right]^2 - \left[\vec{p}_{T,\ell\ell} + \vec{E}_{miss,T} \right]^2$$

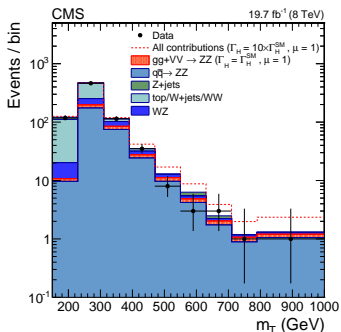


Distributions of selected events

$$H \rightarrow ZZ \rightarrow 2\ell 2\nu$$



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signal-enriched region

- $m_T > 350$ GeV
- $E_T^{miss} > 100$ GeV

		4ℓ	$2\ell 2\nu$
(a)	total gg ($\Gamma_H = \Gamma_H^{SM}$)	1.8 ± 0.3	9.6 ± 1.5
	gg signal component ($\Gamma_H = \Gamma_H^{SM}$)	1.3 ± 0.2	4.7 ± 0.6
	gg background component	2.3 ± 0.4	10.8 ± 1.7
(b)	total gg ($\Gamma_H = 10 \times \Gamma_H^{SM}$)	9.9 ± 1.2	39.8 ± 5.2
(c)	total VBF ($\Gamma_H = \Gamma_H^{SM}$)	0.23 ± 0.01	0.90 ± 0.05
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(d)	total VBF ($\Gamma_H = 10 \times \Gamma_H^{SM}$)	0.77 ± 0.04	2.40 ± 0.14
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	observed	11	91



Systematic uncertainties



- $gg \rightarrow ZZ$
 - Part of cross section uncertainties cancel in the ratio between off-shell and on-shell
 - Shape uncertainties obtained varying PDFs: CT10, MSTW and NNPDF
 - Correlated shape-yield uncertainties produced varying the scales and applying corresponding K-factor
 - $K_{bkg} = K_{sig} \times (1.0 \pm 0.1)$
- $q\bar{q} \rightarrow ZZ$
 - QCD scale: correlate shape and yield uncertainties
 - PDFs: constant 4%
 - $q\bar{q} \rightarrow ZZ$ EWK corrections (2-7%)
- Lepton efficiency for the trigger, reconstruction and selection
- Background estimation from data (mostly $2\ell 2\nu$, up to 25%)

All systematic uncertainties correlated between the on-shell and off-shell regions affect μ but not Γ_H in the combined measurement



Results



Joint unbinned likelihood fit

- Perform fit simultaneously on-shell (4ℓ) and off-shell ($4\ell + 2\ell 2\nu$)
- Fitted parameters: $\mu_F, \mu_V, \Gamma_H/\Gamma_H^{SM}$ ($\Gamma_H^{SM} = 4.15$ MeV)

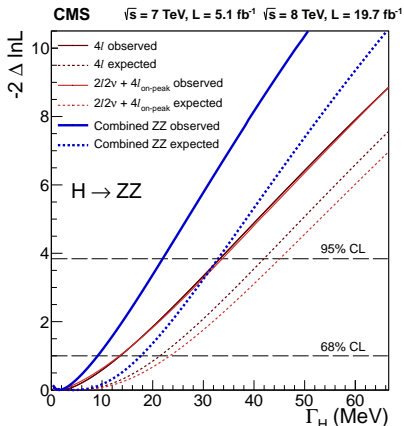
95% CL exclusion limit

- Observed: 22 MeV
- Expected: 33 MeV

Best fit Γ_H

- Observed: $1.8_{-1.8}^{+7.7}$ MeV
- Expected: $4.2_{-4.2}^{+13.5}$ MeV

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Conclusions



First experimental constraint on the Higgs width using $H^* \rightarrow ZZ$ events

- $\Gamma_H < 22 \text{ MeV}$ @95% CL
- More than 2 order of magnitude improvement w.r.t. direct on-peak measurement

Mild model dependency

- Assume essentially no new particles in the gluon fusion loop
- Assume no contribution from BSM physics in the background

Perspectives

- The expected tail/peak cross-sections ratio is ~ 2 times higher @13 TeV than @8 TeV
- The measurement may become limited by systematic uncertainties
- Improved determination of theory cross sections, in particular for the gg background is needed



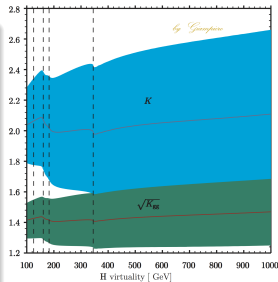
Cross sections and events generation



Gluon fusion, GG2VV and MCFM

- Latest versions of generators include signal, background and interference
- Assume $m_H = 125.6$ and $\Gamma_{SM} = 4.15$ MeV for generations as in PRD 89 (2014) 092007
- Set factorization and renormalization running scale of $m_{4\ell}/2$
- Apply $m_{4\ell}$ -dependent K-factor (NNLO/LO) for signal from G. Passarino (Eur. Phys. J. C74 (2014) 2866)
- Use $K_{background} = K_{signal}$ assigning a 10% uncertainty on this assumption following M. Bonvini et al. (Phys. Rev. D88 (2013) 034032)

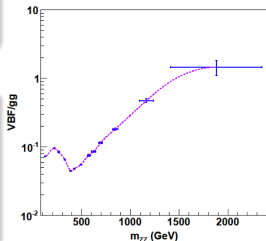
Eur. Phys. J. C74 (2014) 2866



Vector boson fusion, PHANTOM

- VBF production is 7% of the total at H(125.6) peak
- Fraction of VBF increases with $m_{4\ell}$
- Same settings, similar lineshape except second enhancement due to $t\bar{t}$ in the $gg \rightarrow H$

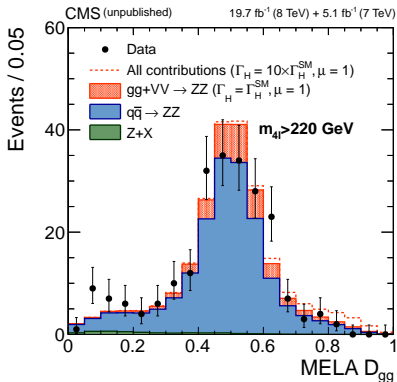
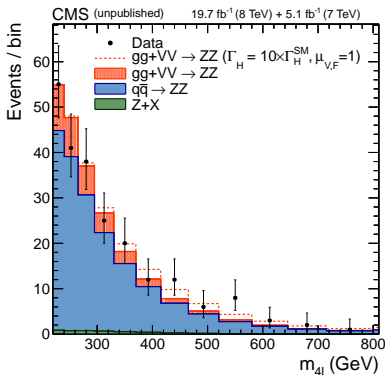
Comparison of gg and VBF $2e^+e^-4\mu$ Rates (Pure H)





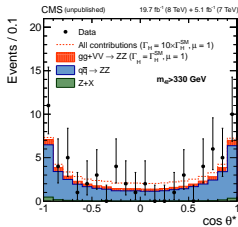
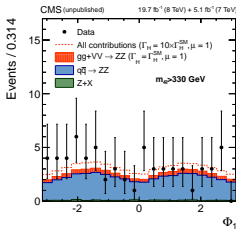
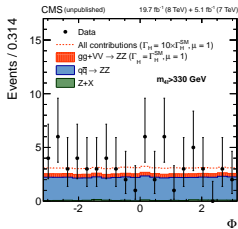
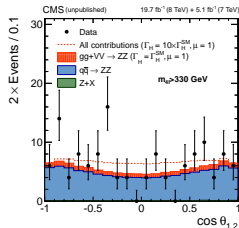
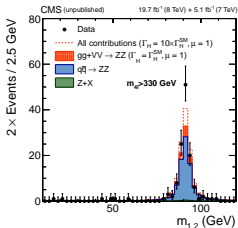
$H \rightarrow ZZ \rightarrow 4\ell$

1D distributions in the full analysis range





$H \rightarrow ZZ \rightarrow 4\ell$ ggMELA inputs, signal-enriched

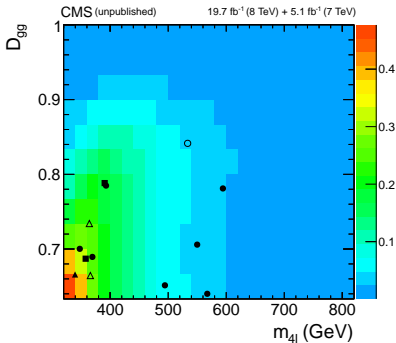




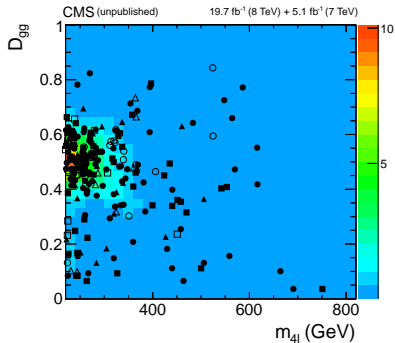
$H \rightarrow ZZ \rightarrow 4\ell$
2D templates



Signal-enriched



All analysis range



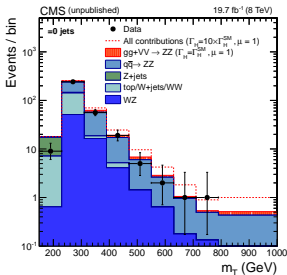


$$H \rightarrow ZZ \rightarrow 2\ell 2\nu$$

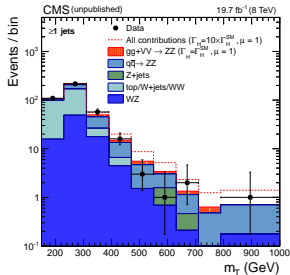
$m_{\ell\ell}^T$ and E_T^{miss}



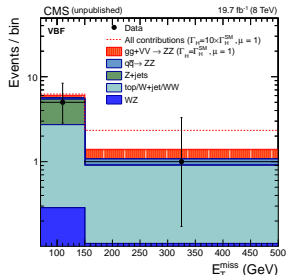
0 jets



≥ 1 jet



VBF-type



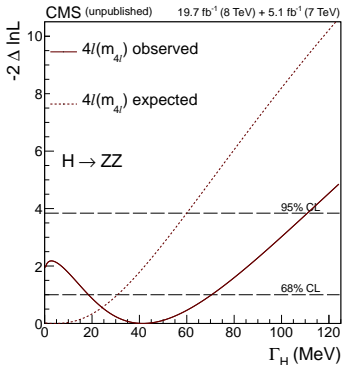


$H \rightarrow ZZ \rightarrow 4\ell$

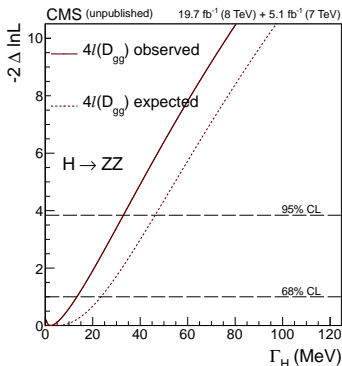
1D likelihood scans



1D $m_{4\ell}$



1D D_{gg}





$$H \rightarrow ZZ \rightarrow 2\ell 2\nu$$

1D likelihood scan, on-shell measurements from 4ℓ



1D $m_{\ell\ell}^T$

