



Higgs and cosmology

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Higgs boson discovery at the LHC

- Atlas $M_H = 125.36 \pm 0.41~{
 m GeV}$
- CMS $M_H = 125.03 \pm 0.29~\mathsf{GeV}$
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Therefore, we can describe the evolution of the Universe from the very early stages till the present days!

Higgs coupling to gravity

Higgs field in general must have non-minimal coupling to gravity:

$$S_G = \int d^4x \sqrt{-g} \Biggl\{ -rac{M_P^2}{2}R - rac{m{\xi}h^2}{2}R \Biggr\}$$

Jordan, Feynman, Brans, Dicke,...

Consider large Higgs fields $h > M_P / \sqrt{\xi}$, which may have existed in the early Universe

The Higgs field not only gives particles their masses $\propto h$, but also determines the gravity interaction strength:

 $M_P^{
m eff} = \sqrt{M_P^2 + \xi h^2} \propto h$

For $h > \frac{M_P}{\sqrt{\xi}}$ (classical) physics is the same $(M_W/M_P^{\text{eff}}$ does not depend on h)!

Potential in Einstein frame



 χ - canonically normalized scalar field in Einstein frame.

This form of the potential is universal for (Bezrukov, MS) $y_t(173.2) < y_t^{crit}$:

$$y_t^{ ext{crit}} = 0.9223 + 0.00118 \left(rac{lpha_s - 0.1184}{0.0007}
ight) + 0.00085 \left(rac{M_H - 125.03}{0.3}
ight) + 0.0023 \left(rac{\log \xi}{6.9}
ight)$$

 $y_t(173.2)$ - top Yukawa coupling in $\overline{\mathrm{MS}}$ - scheme at $\mu = 173.2~\mathrm{GeV},\, lpha_s(M_Z)$ - strong coupling

theoretical uncertainty: $\delta y_t/y_t \simeq 2 \times 10^{-4}$ equivalent to changing of M_H by ~ 70 MeV, or m_t by ~ 35 MeV Buttazzo et al

Numerically for $\xi = 1$, y_t^{crit} coincides with the metastability bound on the top Yukawa coupling



Complicated problem: - extraction of top Yukawa coupling from available data

- FNAL and LHC "Monte Carlo \simeq pole ± 1 GeV" top quark mass
- top quark pole mass is not well defined theoretically: hadronisation, renormalons
- unknown higher order perturbative effects: $\mathcal{O}(\alpha_s^4)$. Estimate of Kataev and Kim: $\delta y_t/y_t \simeq -750 (\alpha_s/\pi)^4 \simeq -0.0015$, corresponding to $\delta m_t \sim 300~{
 m MeV}$
- \square unknown non-perturbative QCD effects, $\delta m_t \simeq \Lambda_{QCD} \simeq 300$ MeV , $\delta y_t/y_t \simeq 0.0015$
- \checkmark Alekhin et al. Theoretically clean is the extraction of y_t from $t\bar{t}$ cross-section. However, the experimental errors in $p\bar{p} \rightarrow t\bar{t} + X$ are quite large, leading to $\delta m_t \simeq \pm 2.8$ GeV, $\delta y_t/y_t \simeq 0.015$

Precision measurements of m_H, y_t and α_s are needed! ILC, TLEP stage of FCC.

Comparison with experiments for $\xi = 1$:

If the Monte Carlo mass is identified with the pole mass and theoretical uncertainties in the pole mass are disregarded, then

 $y_t(173.2) = 0.937 - 1.6$ % above the critical value 0.922 : $1 - 3 \sigma$ away from the boundary, if systematic uncertainties are included



Tevatron - LHC combination : $M_t = 173.34 \pm 0.27 \pm 0.71$ GeV CMS Higgs mass value : $M_H = 125.03 \pm 0.3$ GeV $lpha_s = 0.1184 \pm 0.0007$ The SM vacuum may be absolutely stable, and potential for the Higgs field may be flat at large values of h

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Inflation, Big Bang - all in the framework of the Standard Model!

Role of the Higgs field in cosmology

- Can make the Universe flat, homogeneous and isotropic
- Can produce fluctuations leading to structure formation: clusters of galaxies, etc
- Can lead to Hot Big Bang
- Can play a crucial role in baryogenesis leading to charge asymmetric Universe
- Can play a crucial role in Dark Matter production

Cosmological inflation



- χ_{COBE} χ
- Makes the Universe flat, homogeneous and isotropic
- Produces fluctuations leading to structure formation: clusters of galaxies, etc

CMB parameters - spectrum and tensor modes, $\xi \gtrsim 1000$



Bezrukov, MS For y_t very close to $y_t^{\rm crit}$: critical Higgs inflation tensor-to-scalar ratio can be large, $\xi \sim 10$

Critical point





$$\lambda(z) = \lambda_0 + b \left(\log z\right)^2, \ \ z = rac{\mu}{qM_P}, \ \ M_P = 2.44 imes 10^{18} {
m GeV}$$

Numerically $\lambda_0 \ll 1$, $q \sim 1$, $b \simeq 2.3 \times 10^{-5}$.

Orsay, 22 July, 2014 - p. 13

Effective potential

$$U(\chi) \simeq rac{\lambda(z')}{4\xi^2} ar{\mu}^4 \;, \;\; z' = rac{ar{\mu}}{\kappa M_P}, \;\; ar{\mu}^2 = M_P^2 \left(1 - e^{-rac{2\chi}{\sqrt{6}M_P}}
ight)$$

The parameter μ that optimises the convergence of the perturbation theory is related to $\bar{\mu}$ as

$$\mu^2 = lpha^2 rac{y_t(\mu)^2}{2} rac{ar{\mu}^2}{\xi(\mu)} \,, \;\; lpha \simeq 0.6$$

Behaviour of effective potential for $\lambda_0 \simeq b/16$:



The inflationary indexes



 n_s

 n_s

r can be large! BICEP 2? see also Hamada, Kawai, Oda and Park



Critical Higgs inflation only works if both Higgs and top quark masses are close to their experimental values.



- All particles of the Standard Model are produced
- Coherent Higgs field disappears
- The Universe is heated up to $T \propto M_P / \xi \sim (3 15) \times 10^{13}$ GeV

For further discussion, we need to go beyond the Standard Model, which cannot explain matter-antimatter asymmetry of the Universe and dark matter. The Neutrino Minimal Standard Model - ν MSM will be used.

Three new particles - heavy neutral leptons - HNL - with masses from keV to GeV - explain in addition neutrino masses and oscillations

Heavy neutral leptons interact with the Higgs boson via Yukawa interactions - exactly in the same way other fermions do:

These interactions lead to

- active neutrino masses due to GeV scale see-saw
- creation of matter-antimatter asymmetry at temperatures $T \sim 100 \text{ GeV}$
- Ito dark matter production at $T \sim 100 \text{ MeV}$

Stage 3: Baryogenesis

- Nothing essentially interesting happens between $10^3 \text{ GeV} < T < 10^{13} \text{ GeV}$: all SM elementary particles are nearly in thermal equilibrium.
- Heavy neutral leptons $N_{2,3}$ are out of equilibrium. They are created in interaction with the Higgs boson $H \leftrightarrow N\nu, \ t\bar{t} \leftrightarrow N\nu, \text{ etc}$
- CP- violation in these reactions lead to lepton asymmetry of the Universe
- Electroweak baryon number violation due to SM sphalerons convert lepton asymmetry to baryon asymmetry of the Universe
- These processes freeze out at $T \simeq 140 \text{ GeV}$

Electroweak cross-over

No phase transition in the electroweak theory for Higgs masses larger than **73** GeV the Higgs field vacuum expectation value smoothly grows from small values up to **250** GeV. The crossover temperature

 $T^{crit} = 109.2 \pm 0.8 \; GeV$ $M_{H}^{crit} = 72.3 \pm 0.7 \; GeV$

Т

Stage 4: Dark matter production

Production temperature of Dark matter HNL via processes like $l\bar{l} \rightarrow \nu N_1$:

History of the Universe

Crucial experiments to confirm or to rule out this picture

Experiments, which will be done anyway

Unitarity of PMNS neutrino mixing matrix: $\theta_{13}, \theta_{23} - \pi/4, \text{ type of neutrino mass hierarchy, Dirac CP-violating phase}$

Absolute neutrino mass. The *ν*MSM prediction: $m_1 \leq 10^{-5}$ eV (from DM). Then $m_2 \simeq 5 \cdot 10^{-2}$ eV, $m_3 \simeq 9 \cdot 10^{-3}$ eV or $m_{2,3} \simeq 5 \cdot 10^{-2}$ eV.
(Double β decay, Bezrukov)
Normal hierarchy: 1.3 meV < $m_{\beta\beta}$ < 3.4 meV</p>
Inverted hierarchy: 13 meV < $m_{\beta\beta}$ < 50 meV</p>

- Crucial experimental test the LHC, precise determination of the Higgs mass, $\Delta M_H \simeq 200 \text{ MeV}$
- Crucial cosmological test precise measurements of cosmological parameters n_s, r

New dedicated experiments

High energy frontier

Construction of t-quark factory $-e^+e^-$ or $\mu^+\mu^-$ linear collider with energy $\simeq 200 \times 200$ GeV.

Precise measurement of top and Higgs masses, to elucidate the stability of the EW vacuum and possibility of Higgs inflation.

Search for N_1

X-ray telescopes similar to *Chandra* or *XMM-Newton* but with better energy resolution: narrow X-ray line from decay $N_e \rightarrow \nu \gamma$ One needs:

- Improvement of spectral resolution up to the natural line width $(\Delta E/E \sim 10^{-3}).$
- **FoV** $\sim 1^{\circ}$ (size of a dwarf galaxies).
- Wide energy scan, from $\mathcal{O}(100)$ eV to $\mathcal{O}(50)$ keV.

Detection of An Unidentified Emission Line in the Stacked X-ray spectrum of Galaxy Clusters. E. Bulbul, M. Markevitch, A. Foster, R. K. Smith, M. Loewenstein, S. W. Randall. e-Print: arXiv:1402.2301

An unidentified line in X-ray spectra of the Andromeda galaxy and Perseus galaxy cluster. A. Boyarsky, O. Ruchayskiy, D. lakubovskyi, J. Franse. e-Print: arXiv:1402.4119

Searches for HNL in space

- Has been previously searched with XMM-Newton, Chandra, Suzaku, INTEGRAL
- Spectral resolution is not enough (required $\Delta E/E \sim 10^{-3}$)
- Proposed/planned X-ray missions with sufficient spectral resolution:

Search for N_2 , N_3

Challenge: for baryon asymmetry generation the heavy neutral leptons must be very weakly coupled, to satisfy the Sakharov condition of out-of-equilibrium

Proposal to Search for Heavy Neutral Leptons at the SPS arXiv:1310.1762: general purpose beam dump facility for investigation of the hidden sector

W. Bonivento, A. Boyarsky, H. Dijkstra, U. Egede, M. Ferro-Luzzi, B. Goddard, A. Golutvin, D. Gorbunov, R. Jacobsson, J. Panman, M. Patel, O. Ruchayskiy, T. Ruf, N. Serra, M. Shaposhnikov, D. Treille

Fixed target SPS: SHIP

FCC-ee for 10^{12} Z

very preliminary

Conclusions

The Standard Model Higgs field can play an important role in cosmology:

- It can make the Universe flat, homogeneous and isotropic
- Quantum fluctuations of the Higgs field can lead to structure formation
- Coherent oscillations of the Higgs field can make the Hot Big Bang and produce all the matter in the Universe
- Real and virtual Higgs boson can play a crucial role in baryogenesis leading to charge asymmetric Universe
- Dark Matter production may come about as an effect of mixing between neutrinos and heavy neutral leptons, induced by the Higgs field
- A number of new experiments is needed to reveal the "secret" couplings of the Higgs boson