

CLOSE COLLISIONS DURING STOCHASTIC DEFLECTION OF HIGH-ENERGY CHARGED PARTICLES BY A BENT CRYSTAL

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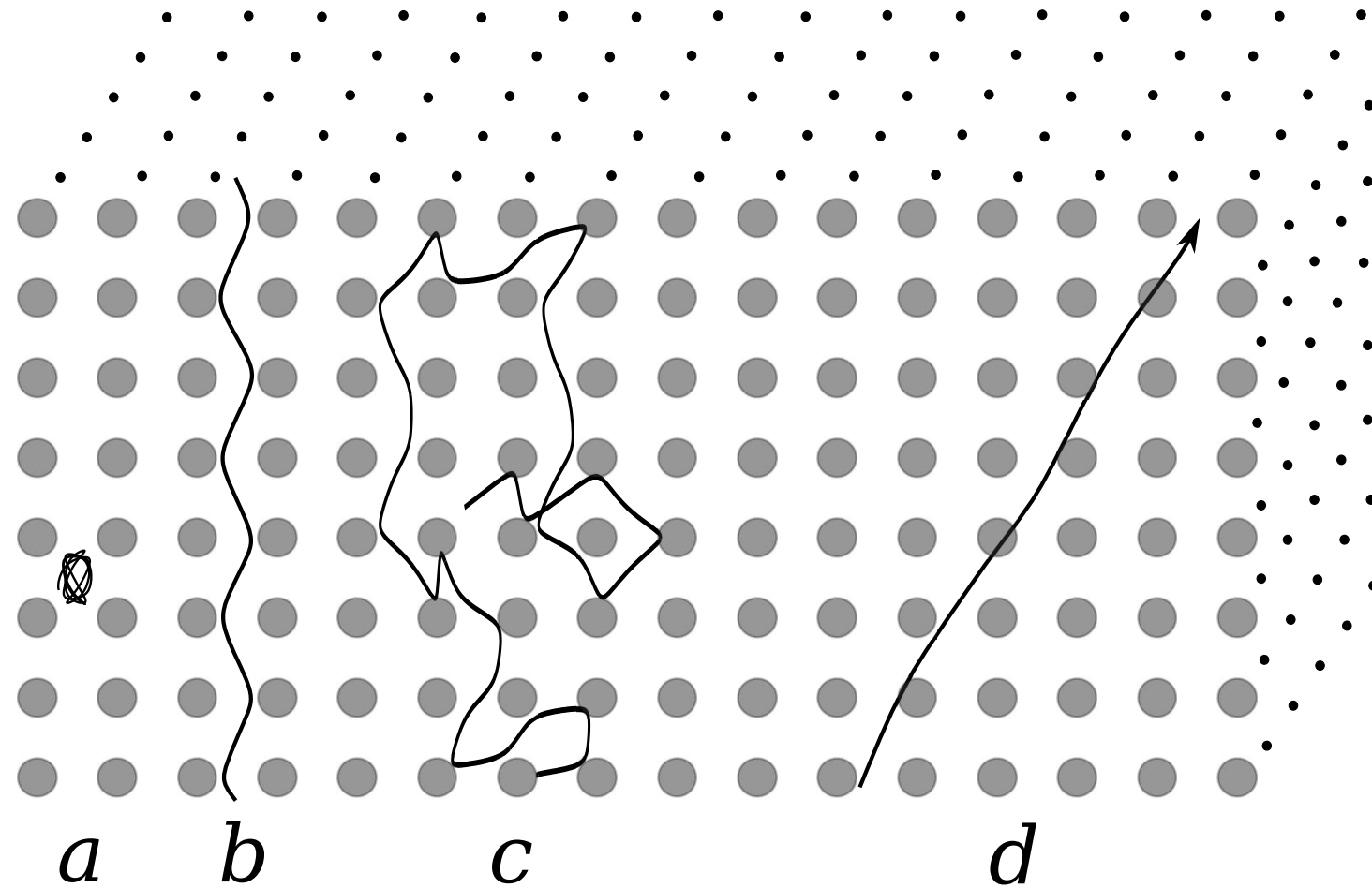
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French-Ukrainian workshop on instrumentation developers for high energy physics

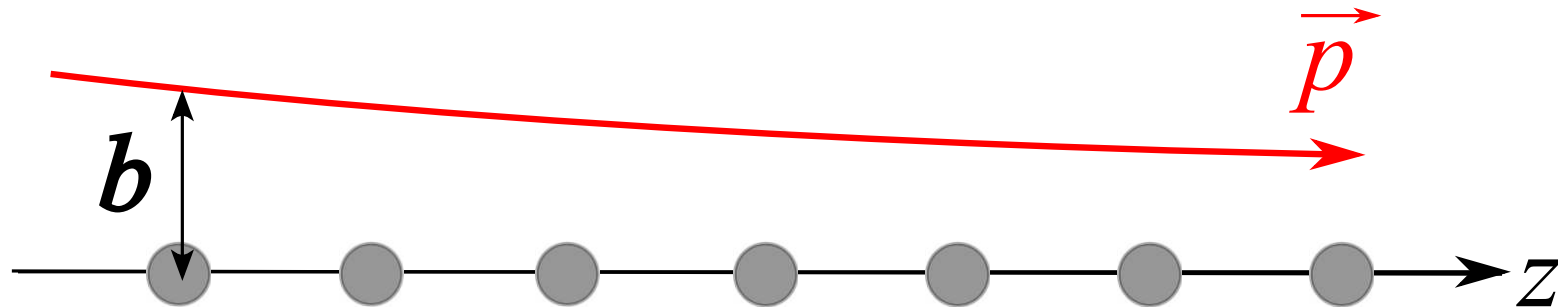
Orsay, 2014

The crystal lattice of Si is a face-centered cubic diamond-type lattice

Diamond, silicon, germanium, and gray tin have this type of crystal lattice.



- a)* axial channeling;
- b)* planar channeling;
- c)* stochastic scattering;
- d)* strongly above-barrier motion.



$$U(\vec{r}) = \sum_n u(\vec{r} - \vec{r}_n)$$

$$U_R(\vec{\rho}) = \frac{1}{L} \int_{-\infty}^{\infty} dz \sum_n u(\vec{r} - \vec{r}_n)$$

$$\ddot{\vec{\rho}} = - \frac{c^2 q}{E} \frac{\partial}{\partial \vec{\rho}} U_R(\vec{\rho})$$

Crystal potential

$$\text{if } \Phi_a(r) = \frac{Z|e|}{r} \exp(-r/R) \quad \text{then} \quad \Phi_s(\rho) = \frac{2Z|e|}{d} K_0(\rho/R),$$

so summation $\sum_n U_R(\vec{\rho} - \vec{\rho}_n)$ couldn't be done analytically.

However, if we use Doyle-Terner approximation

$$\Phi_a(r) = \frac{2\pi\hbar^2}{|e|m_e} \sum_{i=1}^4 \alpha_i \left(\frac{4\pi}{\beta_i + B} \right)^{3/2} \exp\left(-\frac{4\pi^2 r^2}{\beta_i + B} \right) \quad \text{then}$$

$$\Phi_s(\rho) = \frac{1}{d_a} \int_{-\infty}^{\infty} dz \Phi_a(\rho, z) = \frac{8\pi^2\hbar^2}{|e|m_e d_a} \sum_{i=1}^4 \frac{\alpha_i}{\beta_i + B} \exp\left(-\frac{4\pi^2 \rho^2}{\beta_i + B} \right)$$

and summation could be done analytically

Crystal potential

for Si $\langle 100 \rangle$ crystal axis

$$\begin{aligned} \langle \Phi_{\langle 100 \rangle}(\vec{\rho}) \rangle &= \sum_{n=-\infty}^{\infty} \Phi_s(\vec{\rho} - \vec{\rho}_n) = \\ &= \frac{2\pi\hbar^2}{|e|m_e d_a d_s^2} \sum_{i=1}^4 \alpha_i \theta_3 \left(-\frac{x}{d_s} \middle| \frac{i(\beta_i + B)}{4\pi d_s^2} \right) \theta_3 \left(-\frac{y}{d_s} \middle| \frac{i(\beta_i + B)}{4\pi d_s^2} \right), \end{aligned}$$

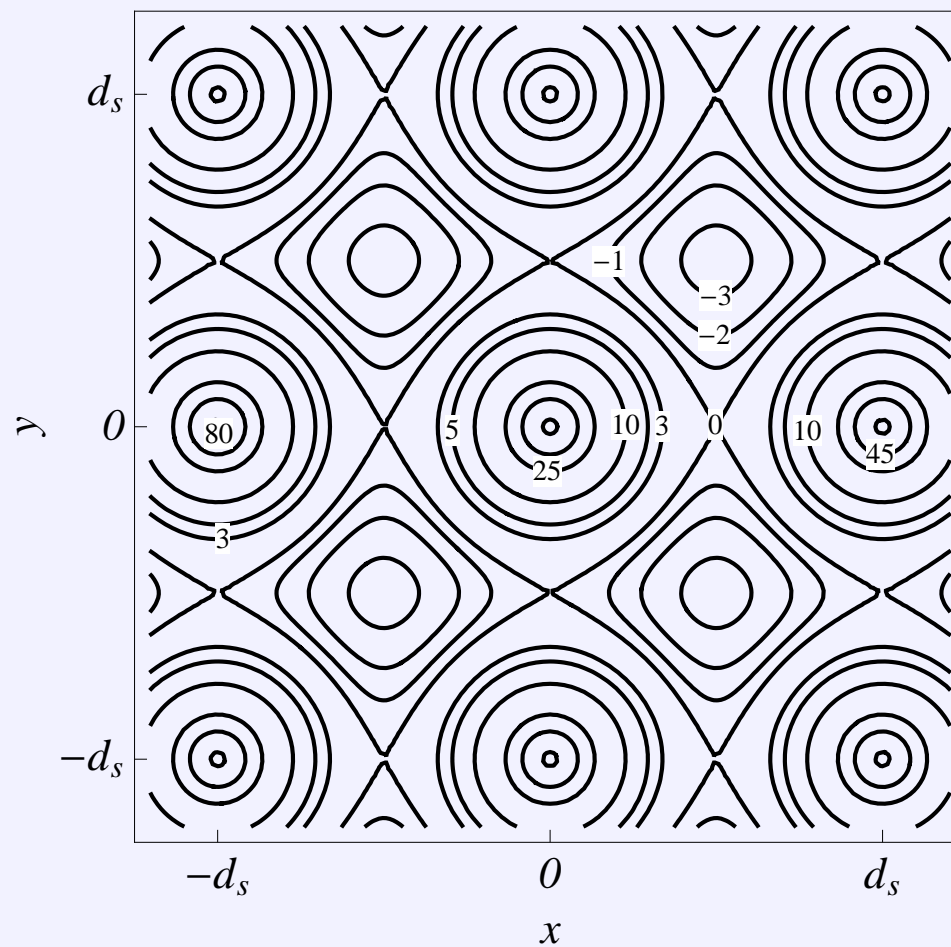
where

d_s — distance between neighboring crystal atomic strings,

$$\theta_3(v|w) = \sum_{n=-\infty}^{\infty} \exp(\pi i w n^2) \exp(2\pi i v n) \text{ — Jacobi theta function of the third type}$$

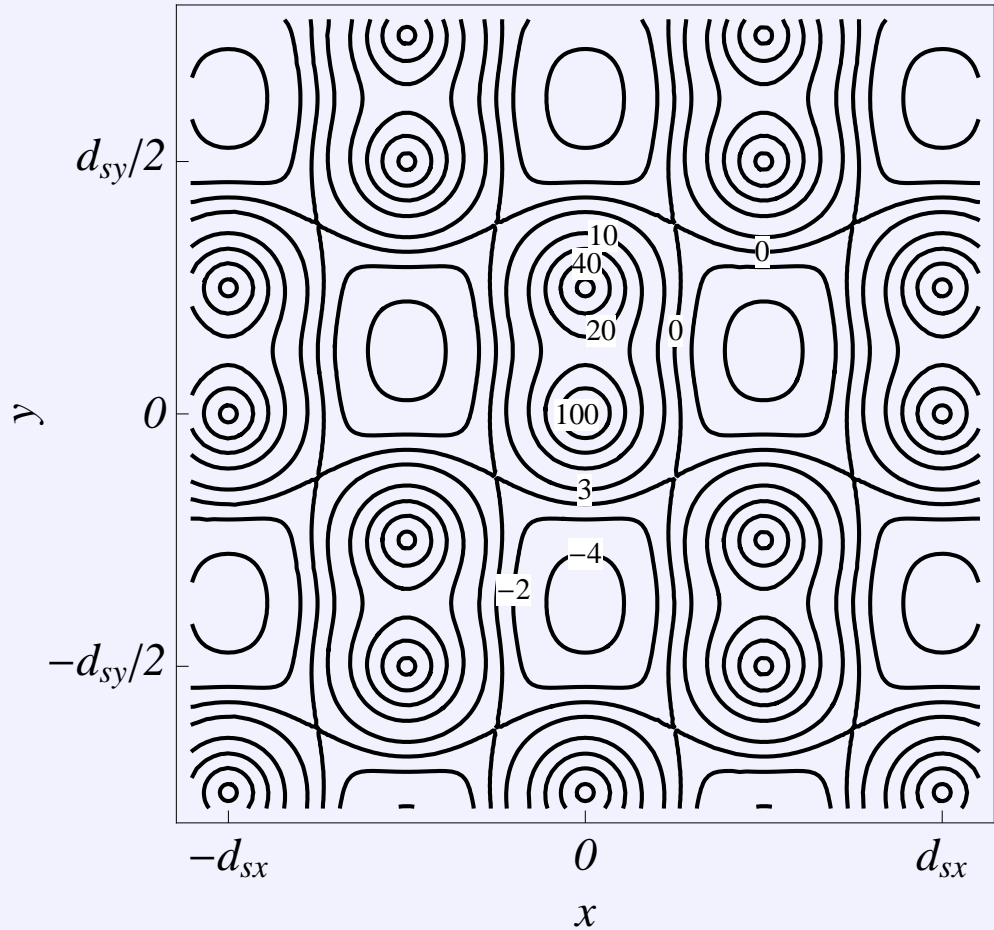
Crystal potential

Si $\langle 100 \rangle$

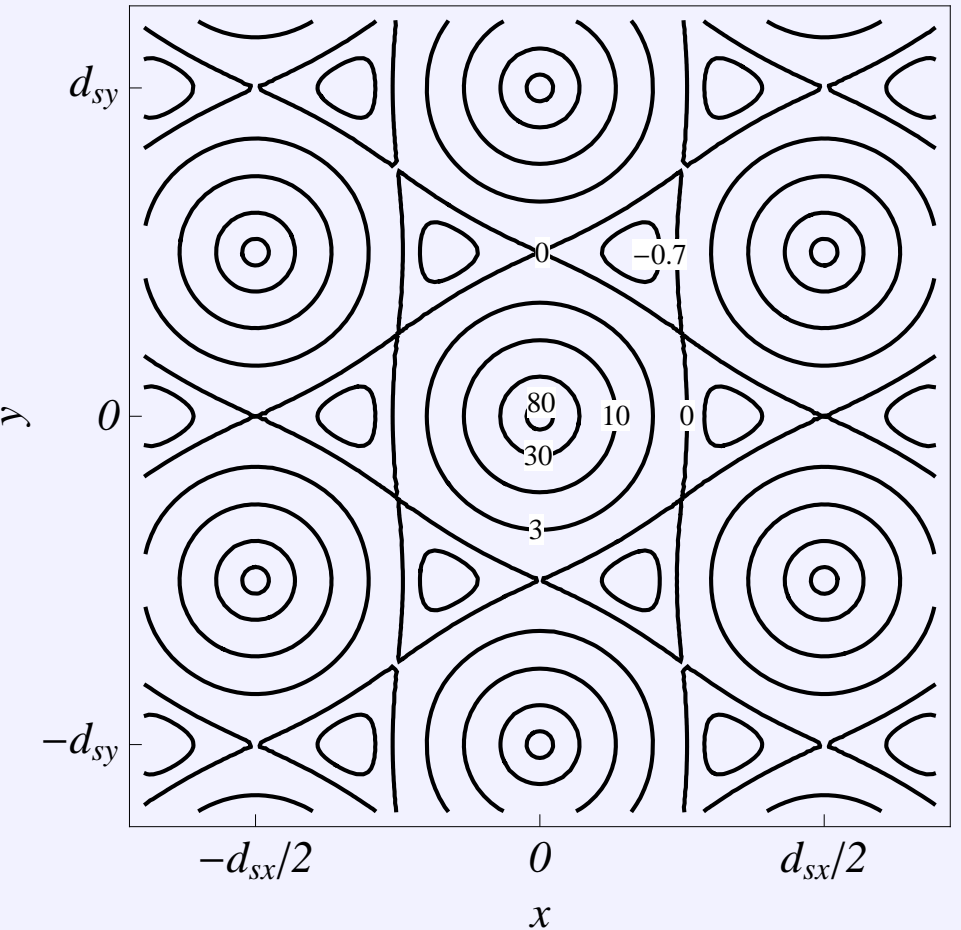


Crystal potential

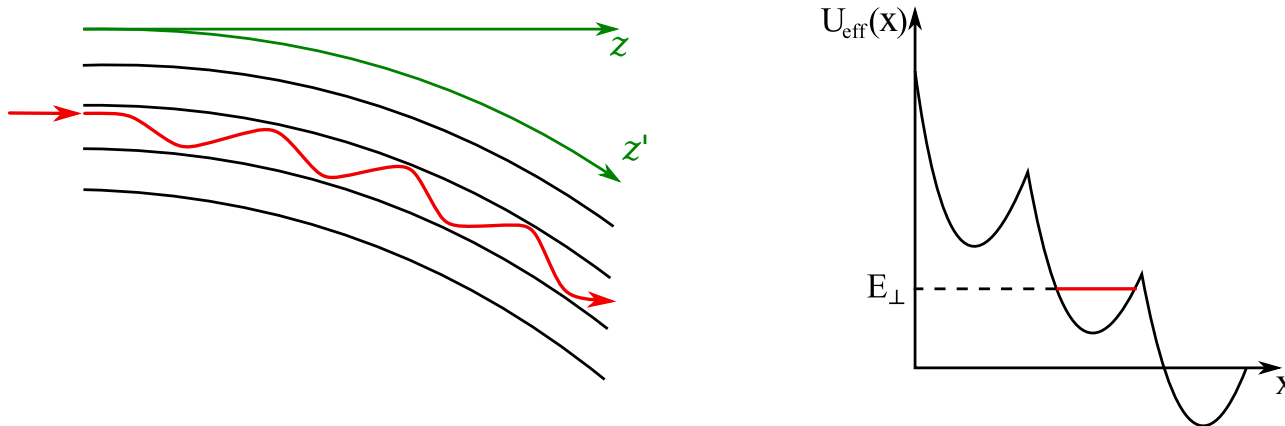
Si $\langle 110 \rangle$



Si $\langle 111 \rangle$



Planar channeling in a bent crystal



If the distance between particle and the center of crystal curvature is $\rho(t) = R + x(t)$, where R is the radius of crystal curvature, $x(t) \ll R$, then

$$\frac{d^2x}{dt^2} = -\frac{c^2}{E} \frac{\partial}{\partial x} U_{eff}(x),$$

where E is particle energy,

$$U_{eff}(x) = U(x) - x \frac{E}{R}.$$

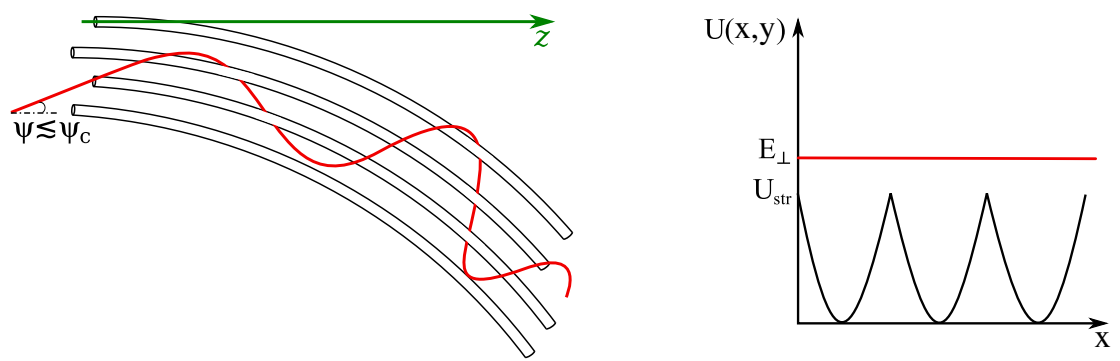
Radius of curvature R_c , at which local minima of function $U_{eff}(x)$ disappear is called the critical. If inter-planar potential is

$$U_p(x) = U_{max} \frac{x^2}{(d/2)^2},$$

where d is the distance between neighboring crystal atomic planes, then

$$R_c = d \frac{E}{4U_{max}}.$$

Stochastic deflection mechanism



Greenenko-Shul'ga criterion: $\frac{l_{\perp}}{R\psi_c} \frac{L}{R\psi_c} < 1$

R is crystal curvature radius;

$\psi_c = \sqrt{4Z|qe|/(pvd)}$ is critical angle of axial channeling;

$Z|e|$ – charge of the nucleus of an atom of the crystal;

q – particle charge;

v и p – particle velocity and momenta;

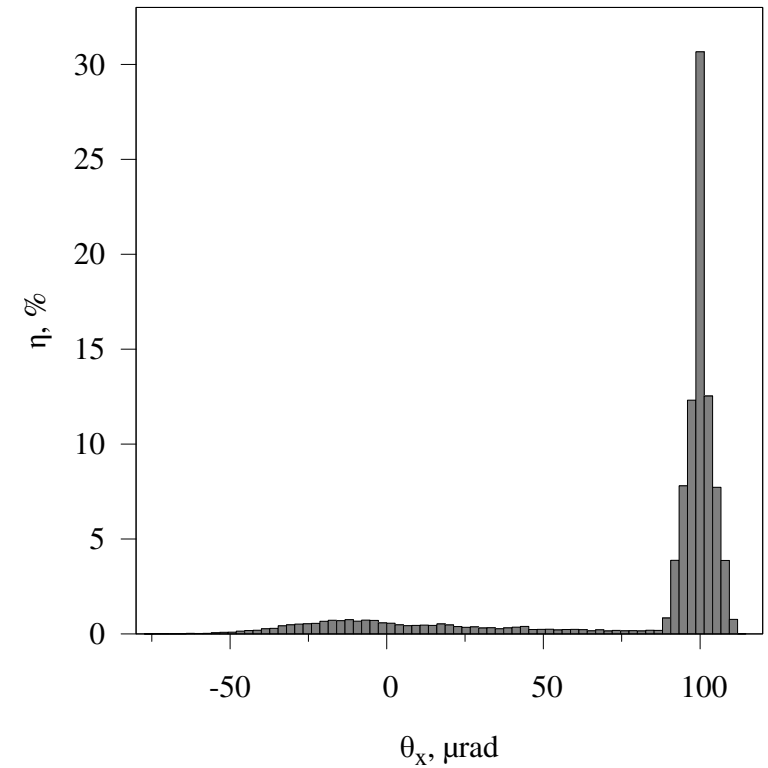
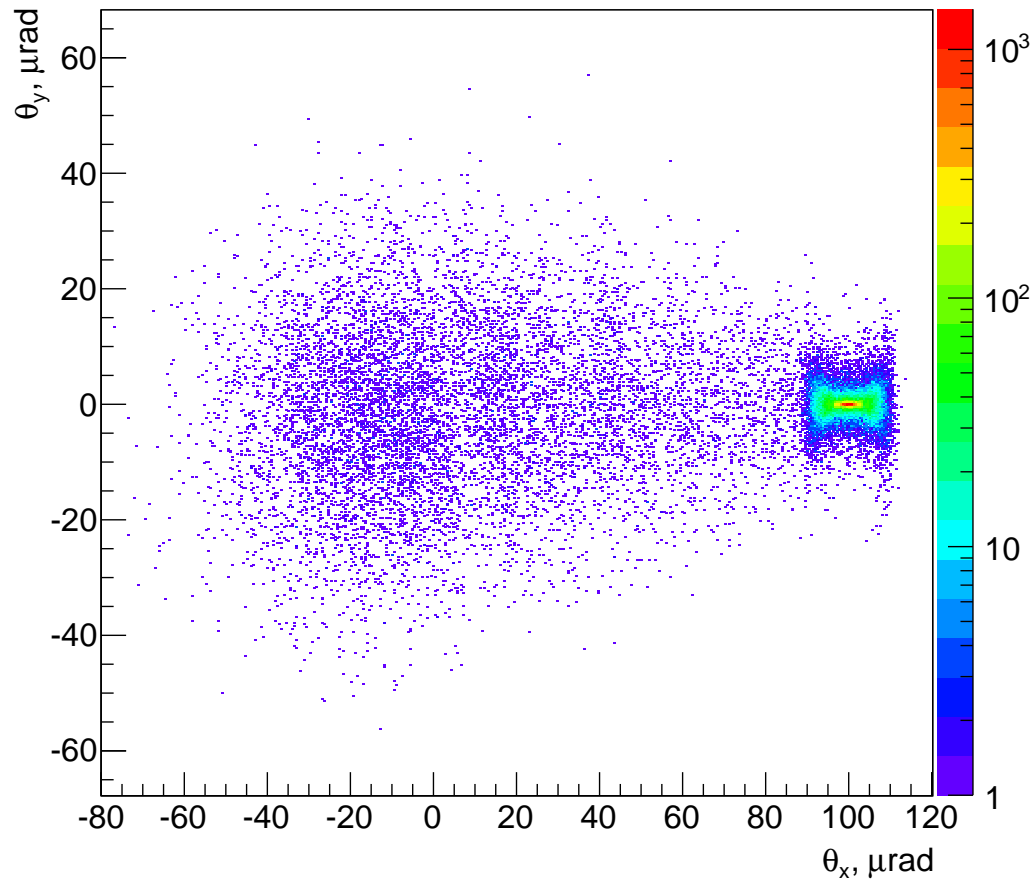
d – distance between neighboring atoms in the string;

l_{\perp} – the mean free path of the particle between successive collisions with strings of atoms in a crystal;

L – the thickness of the crystal.

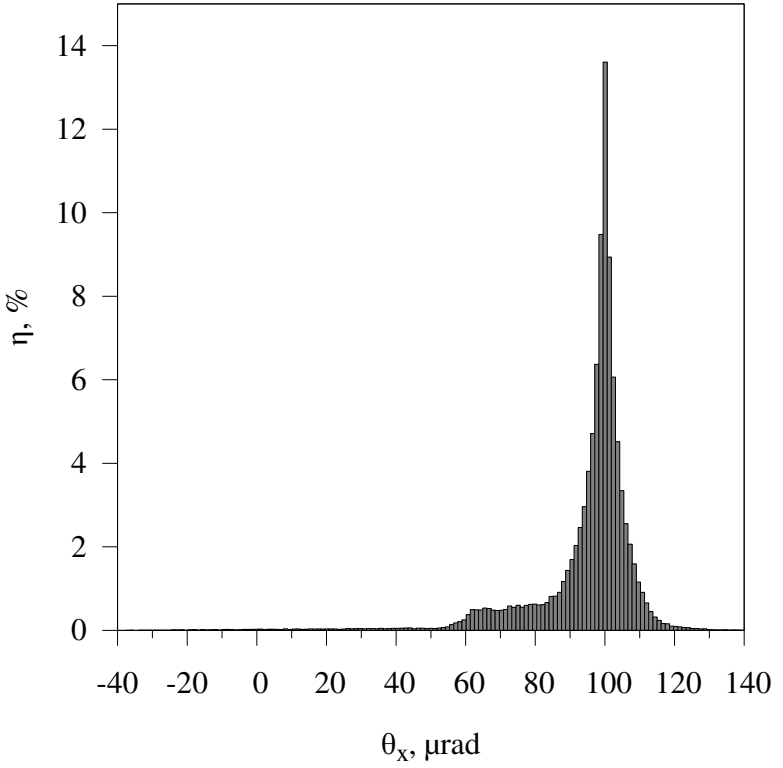
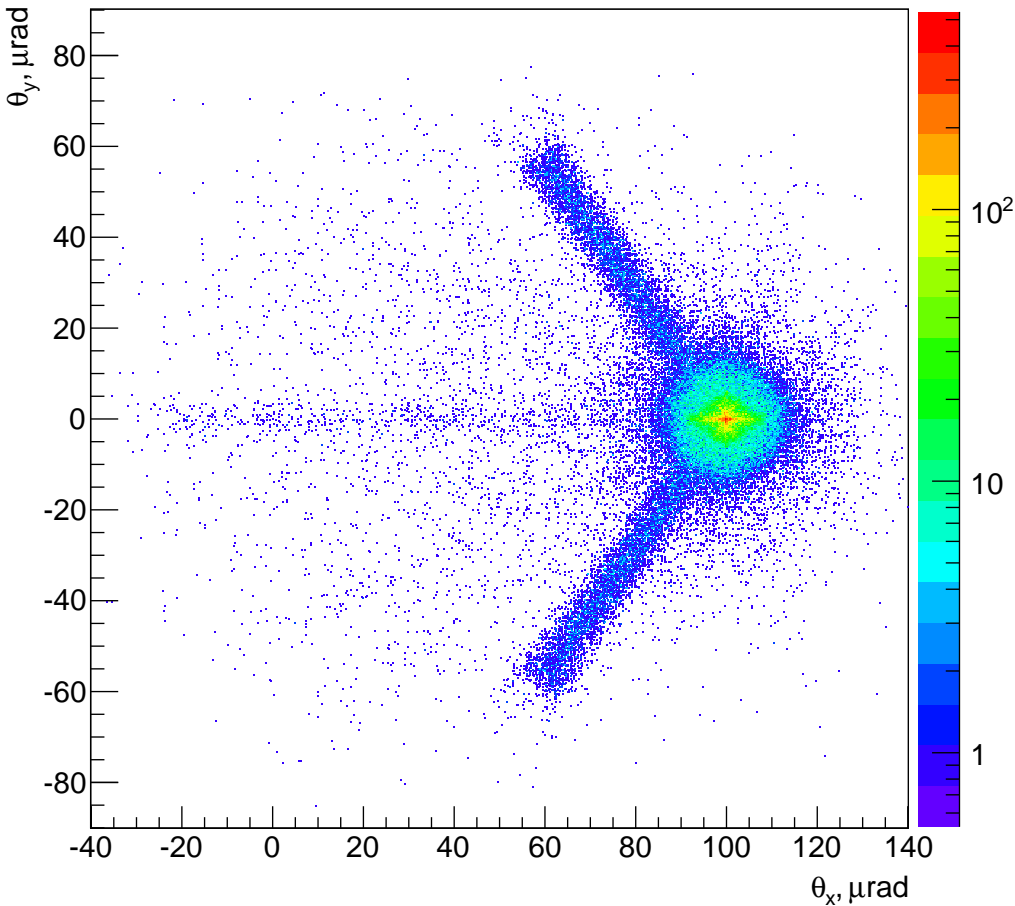
Angular distribution of 270 GeV/ c protons after passing through 5 mm Si crystal with radius of curvature 50 m.

$$\theta_y^{in} = 400 \mu\text{rad}$$



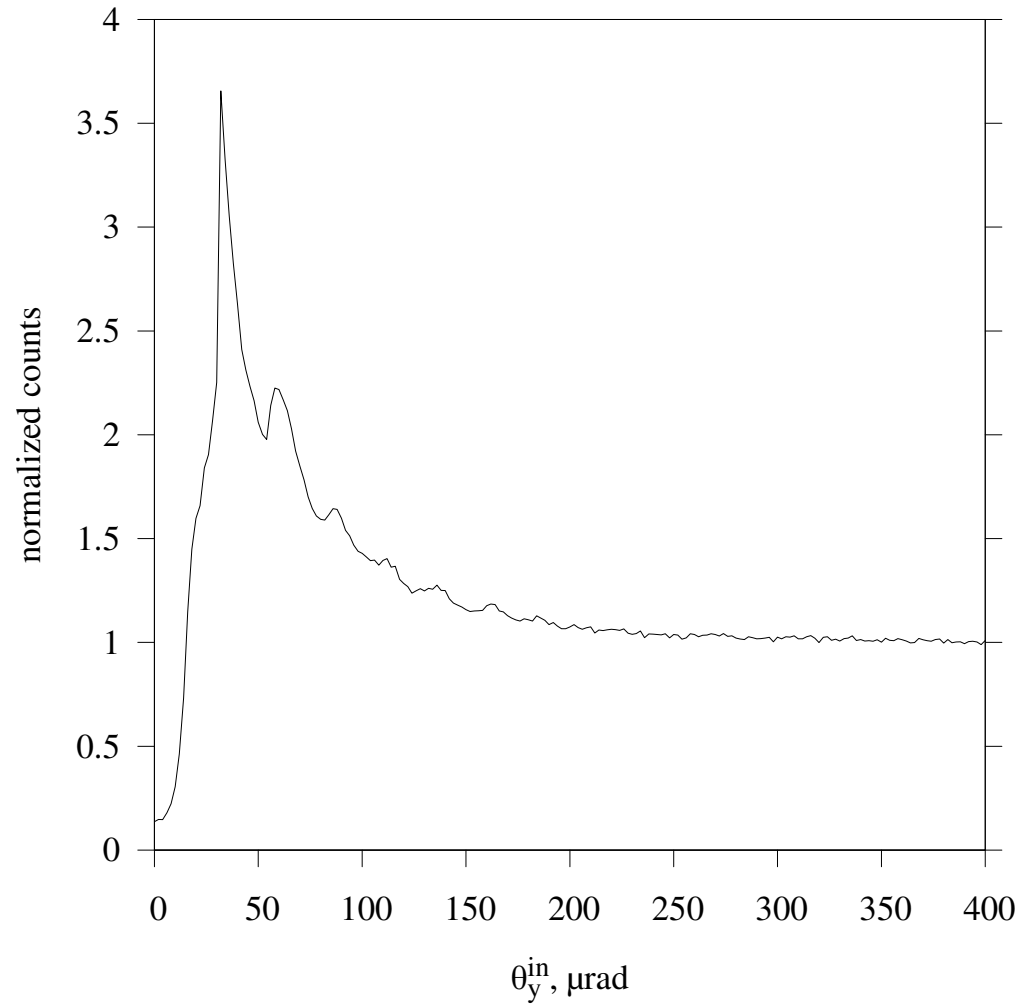
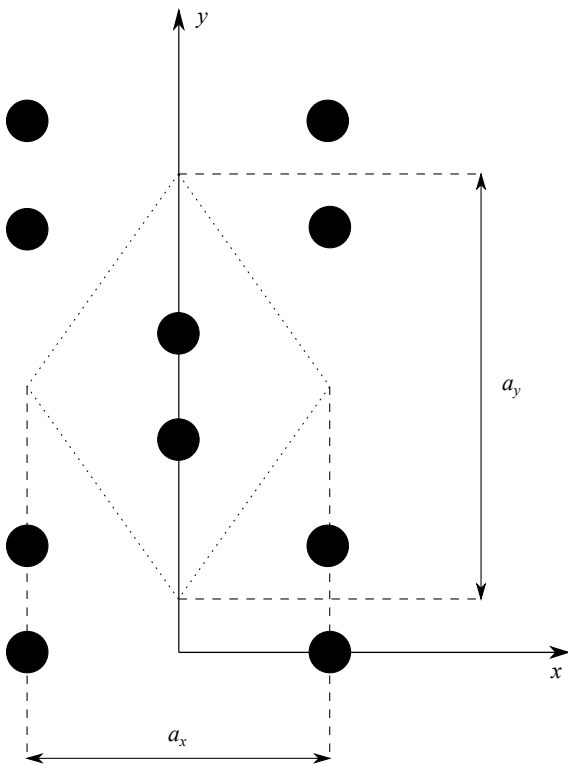
Angular distribution of 270 GeV/c protons after passing through 5 mm Si crystal with radius of curvature 50 m.

$$\theta_y^{in} = 0$$



Yu.A. Chesnokov, I.V. Kirillin, W. Scandale, N.F. Shul'ga, V.I. Truten'. *Physics Letters B* 731 (2014) 118–121

$\langle 110 \rangle$ axis of Si and normalized probability of close collisions of protons in a bent crystal.

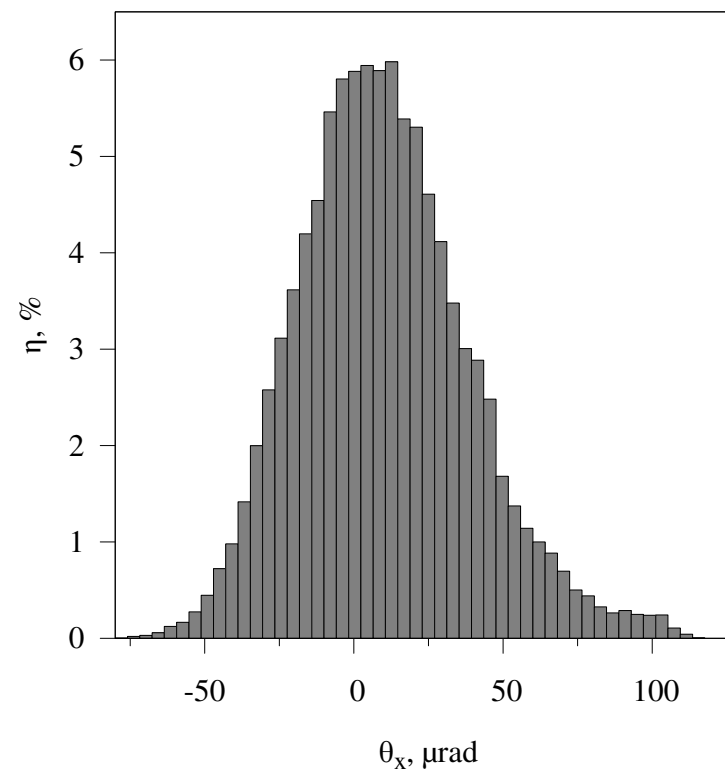
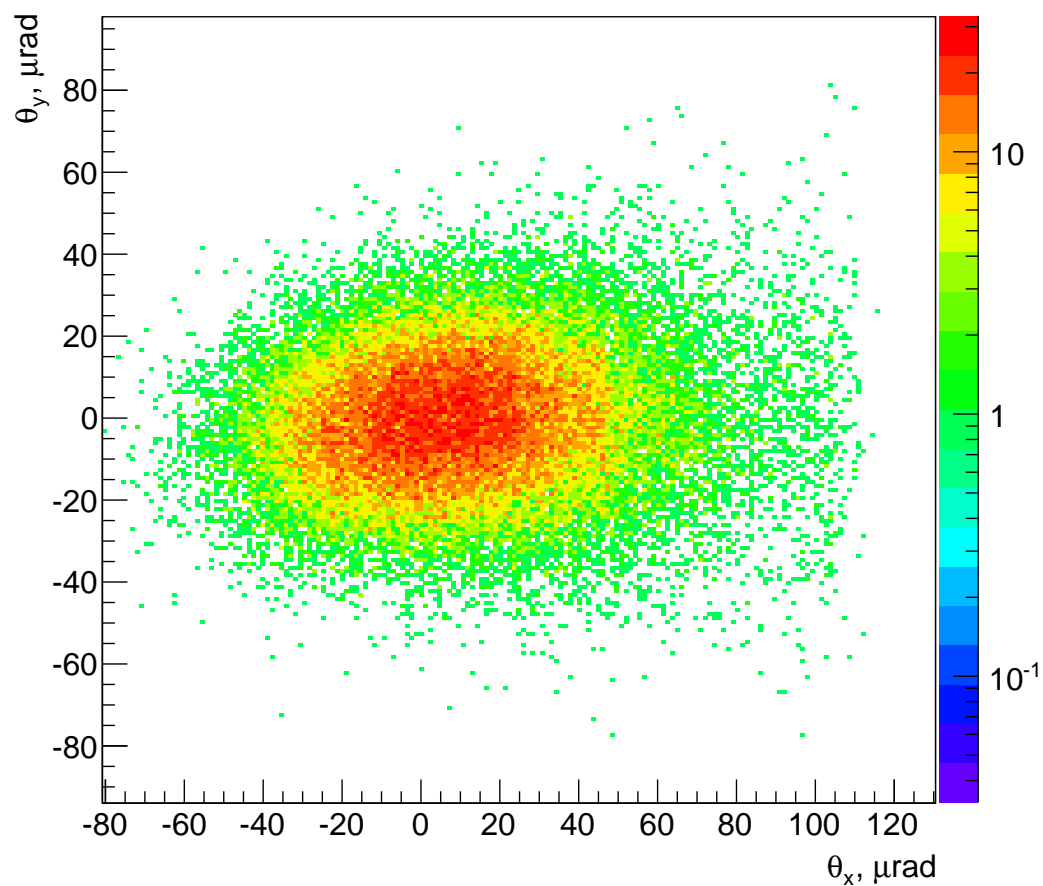


$$w_a = \frac{4\pi r_T^2}{a_x a_y} = 4\sqrt{2}\pi r_T^2 / a^2 \approx 3.39 * 10^{-3}$$

$$w_p = \frac{4r_T}{a_x} = 4\sqrt{2}r_T / a \approx 78.12 * 10^{-3}$$

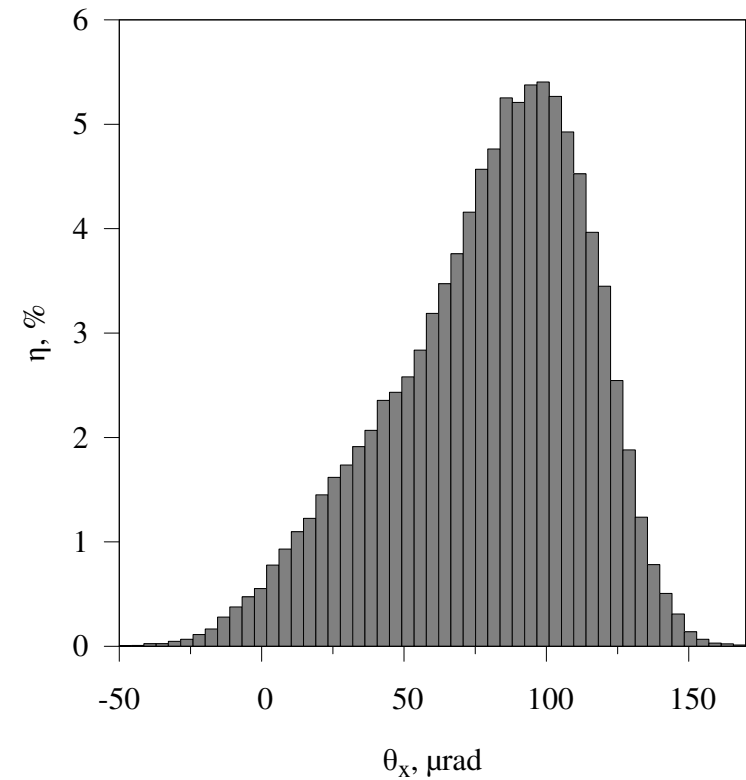
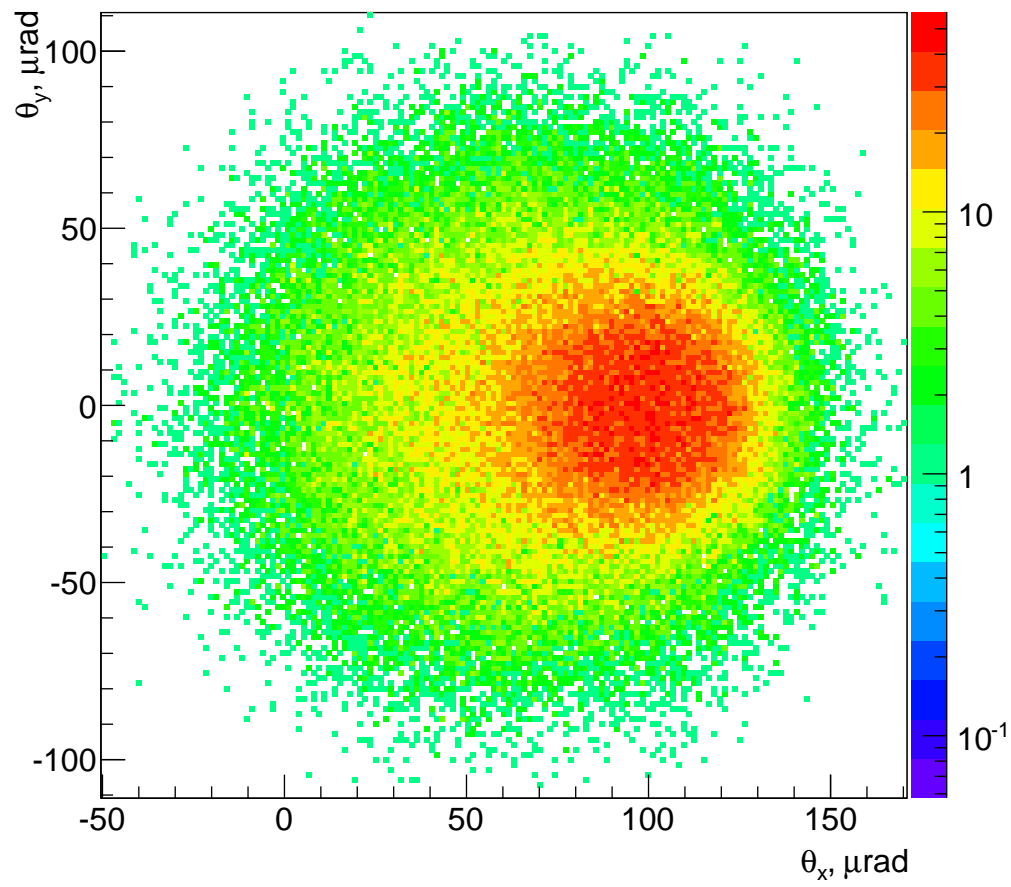
Angular distribution of 270 GeV/c π^- -mesons after passing through 5 mm Si crystal with radius of curvature 50 m.

$$\theta_y^{in} = 400 \mu\text{rad}$$

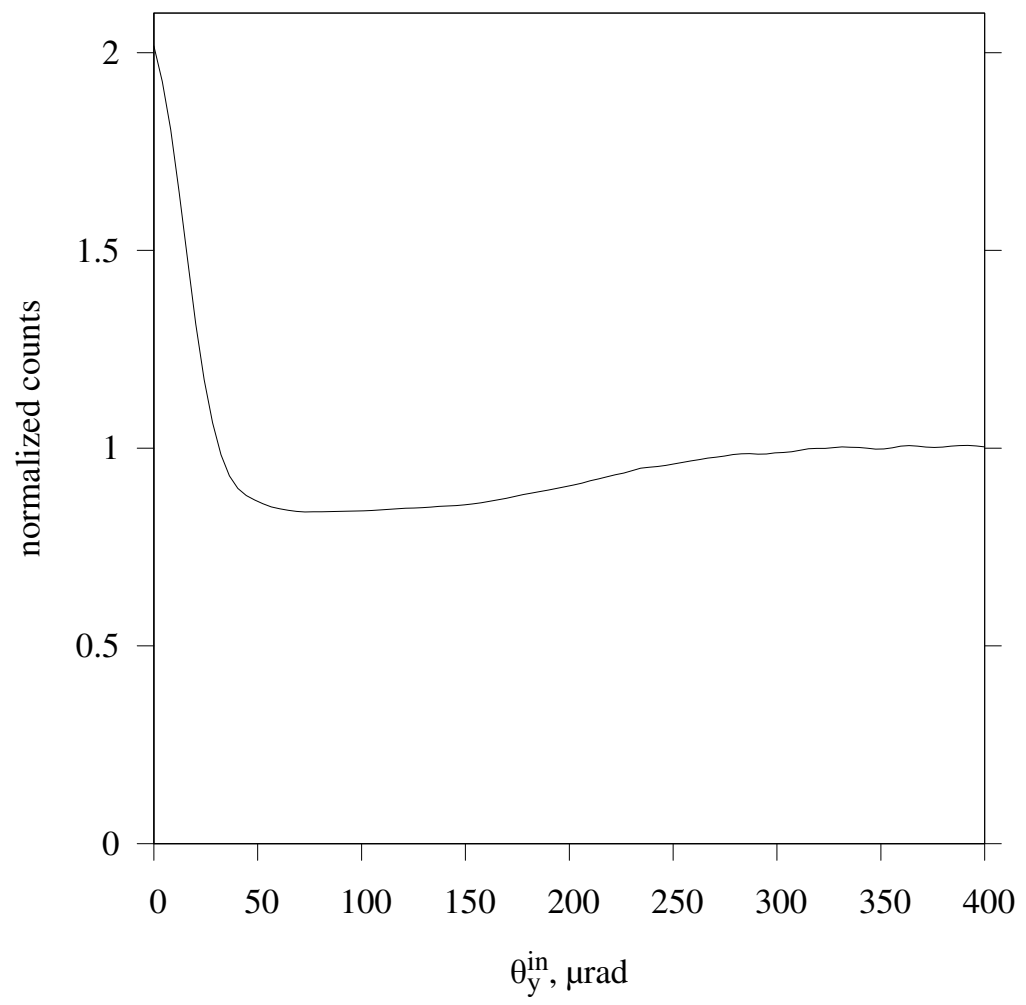
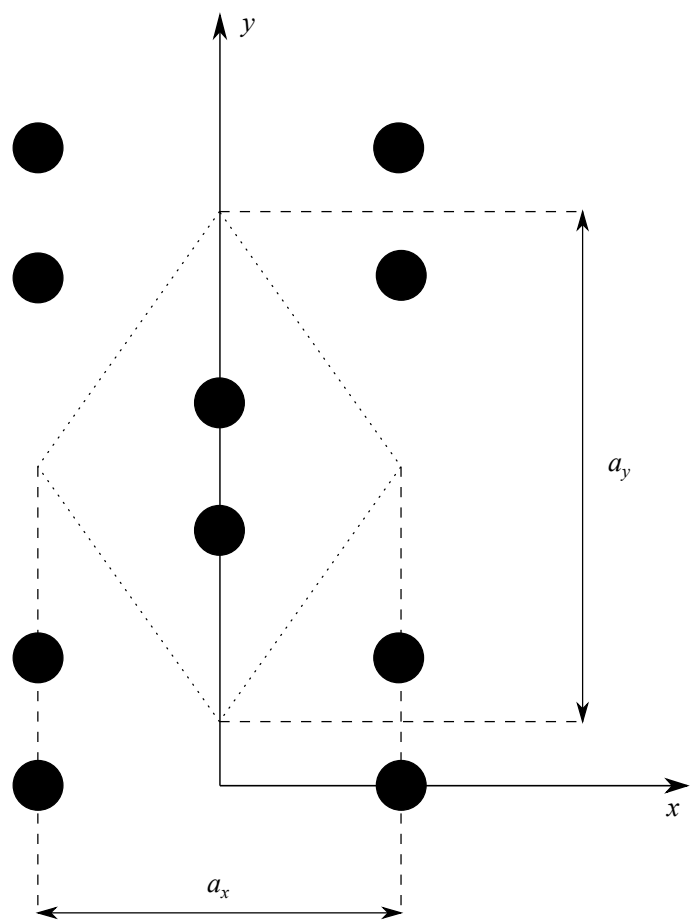


Angular distribution of 270 GeV/c π^- -mesons after passing through 5 mm Si crystal with radius of curvature 50 m.

$$\theta_y^{in} = 0$$



$\langle 110 \rangle$ axis of Si and normalized probability of close collisions of π^- -mesons in a bent crystal.



THANK YOU FOR ATTENTION!