

"This Year a New Era Has Commenced & You Can Say You Have Been Present" -- Goethe, the Cannonade of Valmy and Charm Dynamics

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Cannonade of Valmy: 1792 battle in Northern France that saved `New' France from having (that era's) Standard Model of governance imposed by the `Old' powers.

Tactically a draw, strategically a French victory

Goethe's statement to the Prussian soldiers at camp fire:

'From this place and from this day forth commences a new era in the world's history, and you can all say that you were present at its birth.'

But written up much later; i.e. Goethe -- not unheard of for a theorist -- bragged about a *post*-diction.

in 2007: Strong evidence has surfaced for D oscillations, which could become conclusive by the summer/fall.

A tactical draw in the struggle for gaps in the SM --

x_D & y_D while possibly generated by SM alone, could contain large contributions from NP --

yet a **strategic** victory in sight:

CP studies in the future will decide the issue

possibly paving the way for a **New SM** to emerge!

A historical analogy:

We had been talking about ~~CP~~ in B decays for years without much resonance - till B oscill. were **observed!**

- ☹ numerical size much smaller in D decays,
- ☹ no definitive predictions for ~~CP~~ from New Physics
- 😊 yet SM 'background' even tinier &
- 😊 experimentalists have become more experienced

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Prologue: New Physics Scenarios & Uniqueness of Charm

❖ New Physics in general induces FCNC

👉 their couplings could be substantially stronger for Up-type than for Down-type quarks

(actually happens in some models which 'brush the dirt of FCNC in the down-type sector under rug of the up-type sector')

❖ 2 conceivable scenarios

❑ specific New Physics effects observable/identifiable in charm decays only (much smaller SM 'background')

❑ New Physics effects observable in dynamics of down-type quarks -- B & K -- as well as up-type quarks

👉 still essential complementary info on New Physics!

up-type quarks: u c t

only up-type quark allowing full range of probes for New Phys.

☞ top quarks do not hadronize \implies no $T^0 - \bar{T}^0$ oscillations
hadronization while hard to force under theor. control
enhances observability of CP

☞ up quarks: no $\pi^0 - \pi^0$ oscillations possible
 CP asymmetries basically ruled out by CPT

basic contention:
charm transitions are a unique portal for obtaining a novel
access to flavour dynamics with the experimental
situation being a priori favourable (apart from absence of
Cabibbo suppression)!

I Inconclusiveness in Interpretation of D^0 Oscillations

(1.1) Basics

- 😊 fascinating quantum mechanical phenomenon
- 😐 ambiguous probe for New Physics (=NP)
- 😊 important ingredient for NP CP asymm. in D^0 decays

$$x_D = \frac{\Delta m_D}{\Gamma_D} \quad y_D = \frac{\Delta \Gamma_D}{2\Gamma_D}$$

general expectations

- $\Delta \Gamma$: on-shell contributions
↳ ~ insensitive to New Physics
 - Δm : virtual intermediate states
↳ sensitive to New Physics
- $x_D \sim O(\text{few \%})$ conceivable in models

👉 central theoretical issue:

duality at the charm scale?

- ↔ more averaging in x_D than in y_D
- ↳ duality better in x_D than in y_D

$D^0-\bar{D}^0$ oscillations 'slow' in the SM

How 'slow' is 'slow'?

$$x_D, y_D \sim \cancel{SU(3)_{Fl}} \times 2\sin^2 \theta_C < \text{few} \times 0.01$$

on-shell transitions

off-shell transitions

While the history of predicting x_D, y_D does not fill one of the glory pages of theoret. HEP, we are not completely off the mark either -- see for example:

hep-ph/9712475 (Lecture notes from 1997):

'CP Violation -- an Essential Mystery in Nature's Grand Design'
p.57f: "*It is often stated that the SM predicts ... $x_D, y_D \leq 3 \times 10^{-4}$
I myself am somewhat flabbergasted by the boldness of such predictions... I cannot see how anyone can make such a claim with the required confidence... [my estimate] $x_D, y_D|_{SM} \leq 10^{-2}$.*"

2 general comments:

(A) $x_D \ll y_D$ not a natural scenario!

If $D^0 \rightarrow f \rightarrow \bar{D}^0$ via an *on-shell* final state

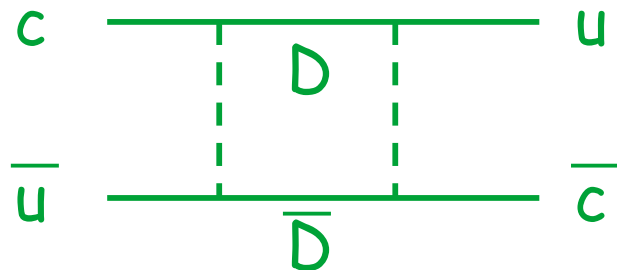
then $D^0 \rightarrow "f" \rightarrow \bar{D}^0$ via an *off-shell* final state

↔ dispersion relation connects Δm_D and $\Delta \Gamma_D$

(B)

GIM suppression $(m_s/m_c)^4$ of usual quark box diagram *un-typically severe!*

→ statement oscillations of mesons built from up-type quarks teach us about down-type quark dynamics



2 general comments:

(A) $x_D < y_D$ natural in SM, yet $x_D \ll y_D$ not!

If $D^0 \rightarrow f \rightarrow \bar{D}^0$ via an *on-shell* final state

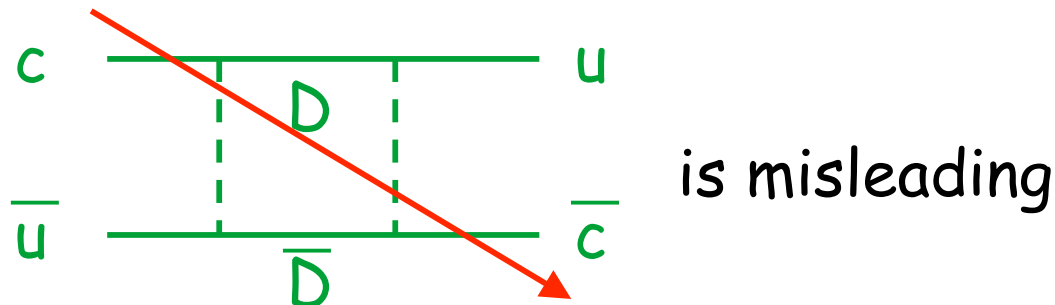
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(1.2) Theoretical Predictions

2 complement. approaches to evaluating Δm_D and $\Delta \Gamma_D$ in the SM:

`inclusive' vs. `exclusive'

↔ `inclusive':

quarks & gluons + nonperturb. contributions

OPE in powers of $1/m_c$, m_s , μ_{had} (quark condensates)

Uraltsev, IB, Nucl. Phys. B592 ('01)

$$m_s^2 m_{\text{had}}^4 / m_c^{(6)} \text{ (vs. } m_s^4 / m_c^{(4)})$$

power counting in $1/m_c$ can be quite iffy

leading contrib. **not**
given by partonic term

- $x_D(\text{SM})|_{\text{OPE}}, y_D(\text{SM})|_{\text{OPE}} \sim \mathcal{O}(10^{-3}) [x_D(\text{SM}) < y_D(\text{SM})]$
- unlikely uncertainties can be reduced
- violations of quark-hadron duality due to proximity of thresholds could enhance in particular y_D

❖ 'exclusive':

hadrons

$SU(3)_{FI}$ breaking from phase space for 2-, 3-, 4-body modes

A. Falk et al., Phys. Rev. D65 ('02)

$$y_D(SM) \sim 0.01 \quad \xrightarrow{\text{dispersion relation}} \quad 0.001 \leq |x_D(SM)| \leq 0.01$$

👉 my judgment: 2 questions

❑ most likely value in SM? $x_D(SM), y_D(SM) \sim \mathcal{O}(10^{-3})!$

❑ can one rule out 0.01? **No!**

(1.3) Data

(1.3.1) Lifetimes of $D^0 \rightarrow K^+K^-$, $\pi^+ \pi^-$ vs. $K^-\pi^+$

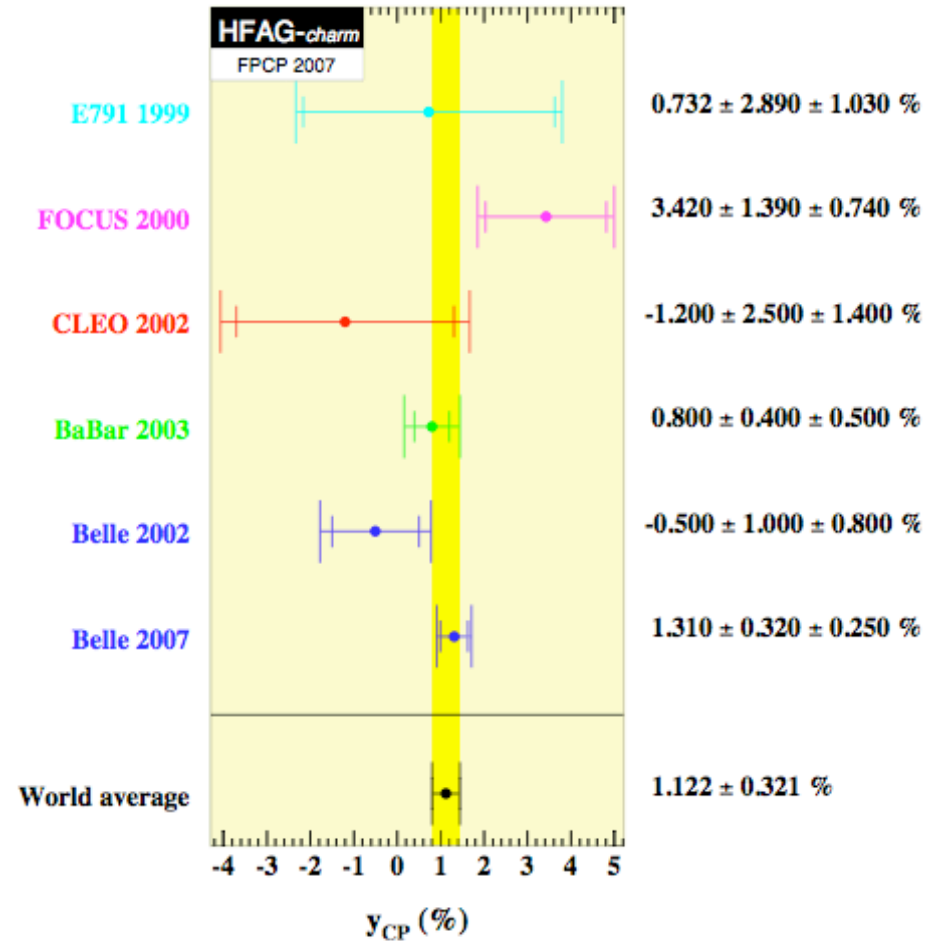
$$\gamma_{CP} = \bar{\tau}/\tau_{CP^+} - 1$$

with CP invariance $\gamma_{CP} = \gamma_D$

BELLE:

$$\gamma_{CP} = (1.31 \pm 0.30 \pm 0.15) \%$$

[$\neq 0$ with 3.2σ]



(1.3.2) $D^0(t) \rightarrow K^+\pi^-$

Rate($D^0(t) \rightarrow K^+\pi^-$) \propto

$$[|T_{DCS}|^2 + \gamma'_D (t\Gamma) T_{DCS} T_{osc} + R_M (t\Gamma)^2 |T_{osc}|^2]$$

$$\gamma'_D = -x_D \sin\delta + \gamma_D \cos\delta, \quad x'_D = x_D \cos\delta + \gamma_D \sin\delta$$

$$R_M = (x_D^2 + \gamma_D^2)/2$$

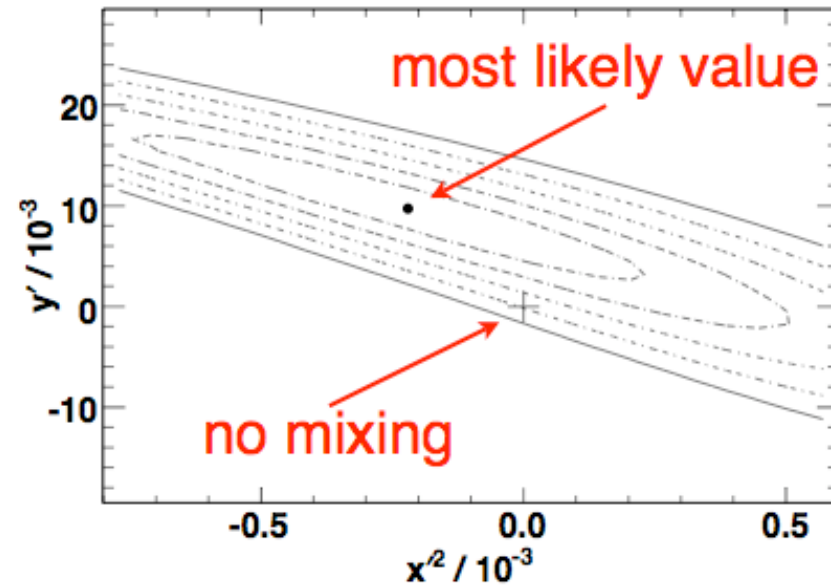
BABAR:

$$\gamma'_D = (0.97 \pm 0.44 \pm 0.31) \times 10^{-4}$$

$$(x'_D)^2 = (-2.2 \pm 3.0 \pm 2.1) \times 10^{-4}$$

[3.9 σ]

$$R_M = (-0.6 \pm 1.7) \times 10^{-4}$$



Decay-time distribution deviates from a pure exponential

$$\frac{\Gamma_{WS}(t)}{\Gamma_{RS}(t)} = R_D + y' \sqrt{R_D} (\Gamma t) + \frac{x'^2 + y'^2}{4} (\Gamma t)^2$$

interference

oscillation

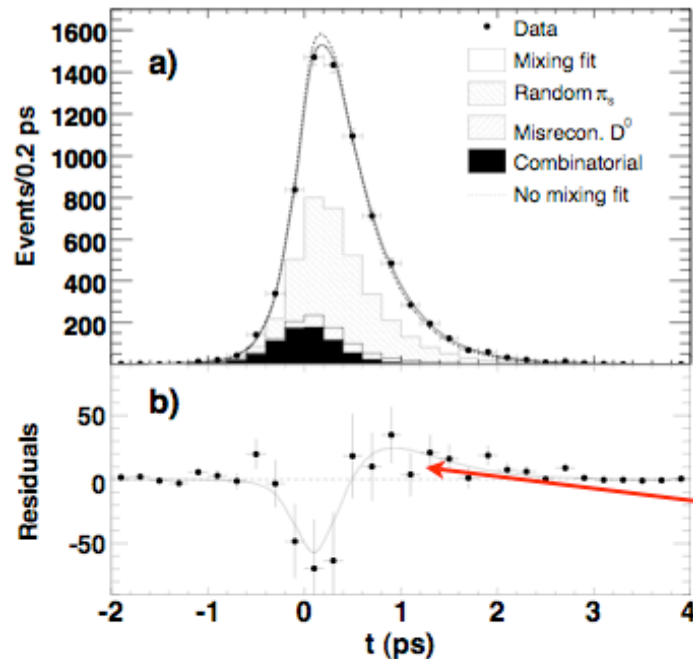
$$y' = y \cos \delta - x \sin \delta$$

$$x' = x \cos \delta + y \sin \delta$$

δ is an unknown strong-phase difference

BABAR observes
positive interference

$$(y' > 0)$$



$$(1.3.3) D^0(t) \rightarrow K_S \pi^+ \pi^-$$

Resonance	Amplitude	Phase (deg)	Fit fraction
$K^*(892)^-$	1.629 ± 0.005	134.3 ± 0.3	0.6227
$K_0^*(1430)^-$	2.12 ± 0.02	-0.9 ± 0.5	0.0724
$K_2^*(1430)^-$	0.87 ± 0.01	-47.3 ± 0.7	0.0133
$K^*(1410)^-$	0.65 ± 0.02	111 ± 2	0.0048
$K^*(1680)^-$	0.60 ± 0.05	147 ± 5	0.0002
$K^*(892)^+$	0.152 ± 0.003	-37.5 ± 1.1	0.0054
$K_0^*(1430)^+$	0.541 ± 0.013	91.8 ± 1.5	0.0047
$K_2^*(1430)^+$	0.276 ± 0.010	-106 ± 3	0.0013
$K^*(1410)^+$	0.333 ± 0.016	-102 ± 2	0.0013
$K^*(1680)^+$	0.73 ± 0.10	103 ± 6	0.0004
$\rho(770)$	1 (fixed)	0 (fixed)	0.2111
$\omega(782)$	0.0380 ± 0.0006	115.1 ± 0.9	0.0063
$f_0(980)$	0.380 ± 0.002	-147.1 ± 0.9	0.0452
$f_0(1370)$	1.46 ± 0.04	98.6 ± 1.4	0.0162
$f_2(1270)$	1.43 ± 0.02	-13.6 ± 1.1	0.0180
$\rho(1450)$	0.72 ± 0.02	40.9 ± 1.9	0.0024
σ_1	1.387 ± 0.018	-147 ± 1	0.0914
σ_2	0.267 ± 0.009	-157 ± 3	0.0088
NR	2.36 ± 0.05	155 ± 2	0.0615

← Cabibbo favored

← doubly Cabibbo suppressed

Belle

$D^0 \rightarrow K_S^0 \pi^+ \pi^-$ features

Doubly Cabibbo suppressed contributions are *enhanced* at high masses

$$\frac{A_{K^*(892)^+}}{A_{K^*(892)^-}} \approx 0.1 \quad \leftarrow \text{seen by CLEO}$$

$$\frac{A_{K_0^*(1430)^+}}{A_{K_0^*(1430)^-}} \approx 0.3$$

makes no sense to me --
Orsay group, check it!

$$\frac{A_{K_2^*(1430)^+}}{A_{K_2^*(1430)^-}} \approx 0.3$$

each corresponds to ~700 events;
comparable to BaBar's
 $D^0 \rightarrow K^+ \rho^- \rightarrow K^+ \pi^- \pi^0$
signal size

$$\frac{A_{K^*(1410)^+}}{A_{K^*(1410)^-}} \approx 0.5$$

$$\frac{A_{K^*(1680)^+}}{A_{K^*(1680)^-}} \approx 1.2$$

$D^0 \rightarrow K_S^0 \pi^+ \pi^-$ results

Possible CP violation not investigated; assuming CP invariance

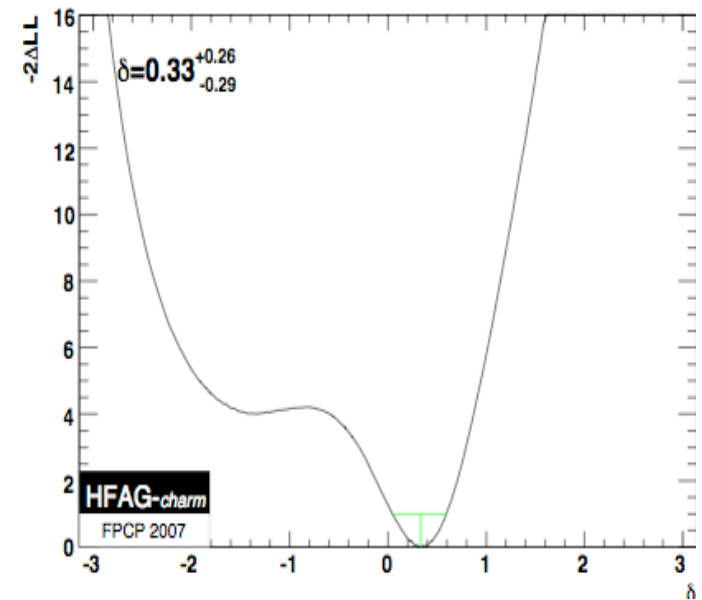
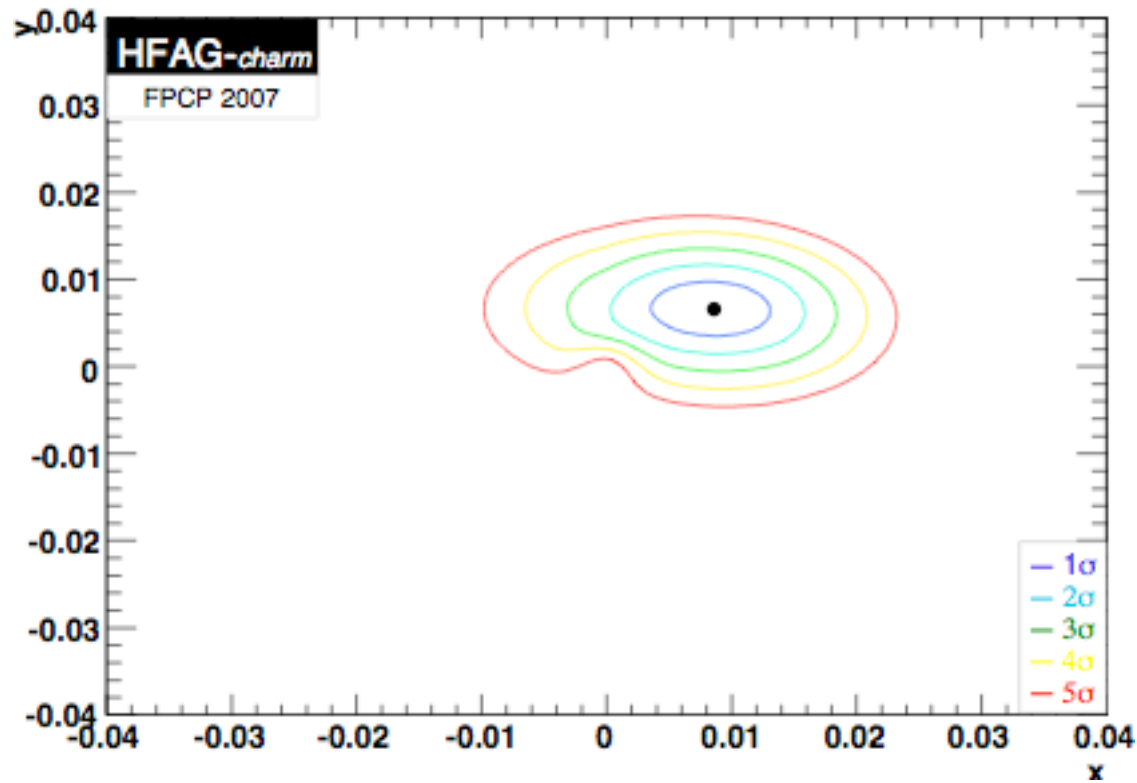
$$x = (0.80 \pm 0.29 \text{ (stat)} \begin{matrix} +0.09 \\ -0.07 \end{matrix} \text{ (syst)} \begin{matrix} +0.15 \\ -0.14 \end{matrix} \text{ (model)})\%$$

$$y = (0.33 \pm 0.24 \text{ (stat)} \begin{matrix} +0.07 \\ -0.12 \end{matrix} \text{ (syst)} \begin{matrix} +0.08 \\ -0.09 \end{matrix} \text{ (model)})\%$$

Reported systematics are much smaller than in other analyses.

$$\text{(Belle KK, pipi)} \quad y = (1.31 \pm 0.32 \text{ (stat)} \pm 0.25 \text{ (syst)})\%$$

$$\text{(BaBar Kpi)} \quad y' = (0.97 \pm 0.44 \text{ (stat)} \pm 0.31 \text{ (syst)})\%$$



in this exercise $(x_D, y_D) \neq (0,0)$ emerges with 5σ

$$x_D = (0.87^{+0.30}_{-0.34})\% , y_D = (0.66^{+0.21}_{-0.20})\% , \delta = 0.33^{+0.26}_{-0.29}$$

(1.4) Interpretation?

- 👉 $x_D > 1\% \gg y_D$ could be interpreted as manifestation of New physics -- yet such a scenario has basically been ruled out
- 👉 data suggest: x_D, y_D can be in range $\sim 0.5 - 1\%$
- 👉 could be due 'merely' to SM dynamics --
 - 👉 even then it would be a great discovery &
 - 👉 it should be measured accurately --
- 👉 must know (i) whether $(x_D, y_D) \neq 0$ & (ii) $x_D = ?$ vs. $y_D = ?$
irrespective of theory -- like for $\varepsilon'/\varepsilon_K$!
- 👉 yet might also contain large contributions from NP!

How to resolve this conundrum?

- theoretical breakthrough?
- CP violation!

II ~~CP~~ with & without D^0 Oscillations

☺ baryon # of Universe implies/requires NP in ~~CP~~ dynamics

☺ existence of three-level Cabibbo hierarchy

$$\text{SM rate } CF : CS : DCS \sim 1 : 1/20 : 1/400$$

☺ within SM:

☞ tiny weak phase in 1x Cabibbo supp. modes: $V(cs) = 1 \dots + i\lambda^4$

☞ no weak phase in Cab. favoured & 2 x Cab. supp. modes

(except for $D^\pm \rightarrow K_S h^\pm$)

☺ CP asymmetry linear in NP amplitude

☺ D^0 oscillations at an observable rate! | ←

☺ final state interactions large

☺ BR's for CP eigenstates large

☺ flavour tagging by $D^{\pm*} \rightarrow D\pi^\pm$

☺ many $H_c \rightarrow \geq 3 P, VV\dots$ with sizeable BR's

☞ CP observables also in final state distributions

(2.1) ~~CP~~ without D^0 Oscillations

direct ~~CP~~

(2.1.1) *time integrated partial widths*

final state interact. $\left\{ \begin{array}{l} \text{☹️ necessary evil} \\ \text{😊 cannot fake signal} \\ \text{😊 } \sim \text{ large in charm} \end{array} \right.$

😊 Cabibbo favour. (CF) modes: need New Physics (except *)

😊 2x Cabibbo supp. modes (DCS): need New Physics (except *)

exception *: $D^\pm \rightarrow K_{S[L]} \pi^\pm$

interference between $D^+ \rightarrow \underbrace{\bar{K}^0}_{CF} \pi^+$ and $D^+ \rightarrow \underbrace{K^0}_{DCS} \pi^+$

in KM only effect from ~~CP~~ in $K^0 - \bar{K}^0$: $A_S = [+]_S - [-]_S = -3.3 \times 10^{-3}$

exists model by G. D'Ambrosio ('01), which creates observable effect in DCS while not affecting oscillations.

LHCb specific: $D^\pm \rightarrow K^\pm \pi^+ \pi^-$

(2.1.2) Final state *distributions*: Dalitz plots, T-odd moments

Dalitz plots asymmetries

final state interact.

- ☹ will be there
- ☺ can *not* fake signal

considerable initial overhead -- yet will pay handsome dividends in the long run due to *overconstraints*

T-odd moments

final state interact.

- ☹ *not* necessary
- ☹ a nuisance: can fake signal
- ☺ can be disentangled

very promising -- most effective theoretical tools not developed yet for small asymmetries (except Dalitz plot)

Pilot study by Focus (CLEO-c?)

- ☺ `local' asymmetry likely to be larger than integrated one
- ☺ angular asymmetry can provide info on chirality of underlying effective operator!

An example for a T odd distribution

$$K_L \rightarrow \pi^+ \pi^- e^+ e^-$$

$$\text{BR} \sim 3 \times 10^{-7}$$

interference between
~~CP~~ E1 & CP M1 amplitude

ϕ = angle between $\pi^+ \pi^-$ & $e^+ e^-$ planes

forward-backward asymmetry in ϕ : $A = 14\%$ driven by $\varepsilon = 0.002$

-- i.e. trade BR for size of asymmetry!

$$D \rightarrow K \bar{K} \pi^+ \pi^-$$

ϕ = angle between $\pi^+ \pi^-$ & $K \bar{K}$ planes

$$d\Gamma/d\phi (D \rightarrow K \bar{K} \pi^+ \pi^-) = \Gamma_1 \cos^2 \phi + \Gamma_2 \sin^2 \phi + \Gamma_3 \cos \phi \sin \phi$$

$$d\Gamma/d\phi (\bar{D} \rightarrow K \bar{K} \pi^+ \pi^-) = \bar{\Gamma}_1 \cos^2 \phi + \bar{\Gamma}_2 \sin^2 \phi - \bar{\Gamma}_3 \cos \phi \sin \phi$$

• Γ_3 drops out after integrating over ϕ

→ Γ_1 vs. $\bar{\Gamma}_1$ & Γ_2 vs. $\bar{\Gamma}_2$: ~~CP~~ in partial widths

• T odd moments $\Gamma_3, \bar{\Gamma}_3 \neq 0$ can be faked by FSI

yet $\Gamma_3 \neq \bar{\Gamma}_3 \implies \text{CP!}$

even closer analogy

$$D \rightarrow K^+ K^- \mu^+ \mu^-$$

$$\text{BR} \sim O(10^{-6})$$

A few general remarks on CP in *final state distributions*

$D \rightarrow PPP$

A Catholic Scenario:

single path to heaven: asymmetries in the Dalitz plot

$D \rightarrow PPPP$

A Calvinist Scenario

many paths to heaven -- success reveals Heaven's blessing

$D \rightarrow K^+K^-\pi^+\pi^-$ ϕ = angle between $\pi^+\pi^-$ & $K K$ planes

① **Integrated** (over 2 quadrants) T odd moment

$$\langle A \rangle = 2\Gamma_3/\pi(\Gamma_1+\Gamma_2) \quad \text{vs.} \quad \langle A \rangle = 2\Gamma_3/\pi(\Gamma_1+\Gamma_2)$$

② **Differential** T odd moment

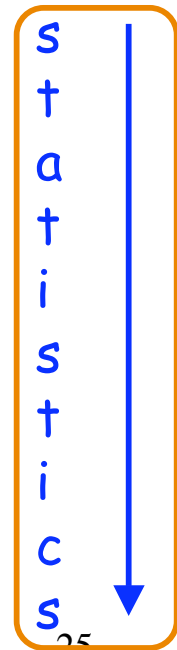
$$d\Gamma/d\phi(D \rightarrow K K \pi^+\pi^-) = \Gamma_1 \cos^2\phi + \Gamma_2 \sin^2\phi + \Gamma_3 \cos\phi \sin\phi$$

same dynamical info, yet valuable experim. check

③ Full amplitude analysis

☺ more dynamical info

☹ more model dependence (?)



(2.2) ~~CP~~ with D^0 Oscillations

All the previously given justifications for CP searches
plus

$$L(\Delta C=2) \neq 0$$

- provides a much wider stage for ~~CP~~ to surface
- allowing us to decide whether NP is involved.

Analogies with two other cases,
one from the past & one from the present:

K^0 & B_s oscillations

$\Delta S=2$:

Assume -- contrary to history -- that people had accepted the SM with 2 families when $\Delta M_K \neq 0$ was observed & knew about possibility of ~~CP~~.

They would have reasoned that LD dynamics could produce $\sim 1/3$ of ΔM_K via $K^0 \rightarrow \pi, \eta, \eta', \pi\pi, \dots \rightarrow \bar{K}^0$ and SD dynamics via the quark box diagram the rest.

This might have led to the proposal to search for $K_L \rightarrow \pi\pi$ to establish the presence of NP, namely the 3rd family (which is irrelevant for ΔM_K).

$\Delta B=2$ -- the topical example:

The observed value of $\Delta M(B_s)$ is fully consistent with SM expectations -- within sizable uncertainties. Yet a subdominant NP contribution to $\Delta M(B_s)$ could still provide the dominant source of time dependent ~~CP~~ in $B_s \rightarrow \psi\phi$!

oscillations can generate *time dependent* CP asymmetries

□ none seen so far down to the 1% ($1\%/ \tan^2 \theta_c$) level --

☞ they are $\sim (x_D \text{ or } y_D) (t/\tau_D) \sin \phi_{\text{weak}}$:

☞ with $x_D, y_D \leq 0.01$ a signal would not have been credible

☞ yet now it is getting interesting!

Scenario (A)

LD dynamics (involving barely 2 families) cannot generate ~~CP~~!

I.e., minimal scenario: *no* significant ~~CP~~ in $L(\Delta C=2)$,

direct ~~CP~~ only: (i) $|q|=|p|$,

whereas (ii) $|\mathcal{T}(D \rightarrow f)| \neq |\mathcal{T}(\bar{D} \rightarrow \bar{f})|$

(iii) $\text{Im}(q/p)\bar{\rho}(f) \neq 0$

□ CF: $K_S\pi^0, K_S\rho^0, \boxed{K_S\phi}$ $\text{Im}V(cs)V(ud) = \eta|V(cb)|^2 \sim 0.6 \times 10^{-3}$

□ DCS: $D^0 \rightarrow K^+\pi^-$ -- $\text{Im}V(cd)V(us) = 0$
yet NP models a la D'Ambrosio

□ CS: $D^0 \rightarrow K^+K^-, \pi^+\pi^-$ -- time depend. & indep. CP

Scenario (B)

NP contributes significantly to $L(\Delta C=2)$

→ expect significant source for ~~CP~~ in $L(\Delta C=2)$:

(i) $|q| \neq |p|$, (ii) $|\mathcal{T}(D \rightarrow f)| \neq |\mathcal{T}(\bar{D} \rightarrow \bar{f})|$, (iii) $\text{Im}(q/p)\bar{\rho}(f) \neq 0$

□ CF: $D^0 \rightarrow K_S \phi$ $A_{CP}(t) = (x_D \sin \phi_{NP} - y_D \varepsilon_{NP} \cos \phi_{NP})(t/\tau_D)$

$$L(\Delta C=2) \rightarrow \phi_{NP} \ \& \ \varepsilon_{NP} = 1 - |q/p|$$

□ CS: $D^0 \rightarrow K^+ K^-, \pi^+ \pi^-$ $A_{CP}(t) = (x_D \sin \phi'_{NP} - y_D \varepsilon_{NP} \cos \phi'_{NP})(t/\tau_D)$

$D^0 \rightarrow K^+ K^- \pi^+ \pi^-$ $\Gamma_3(t), \bar{\Gamma}_3(t)$ time dependence!

□ DCS: $D^0 \rightarrow K^+ \pi^-$ -- ditto (+NP models a la D'Ambrosio)

the SM amplitude suppressed by $\text{tg}^2 \theta_c$

The 'Dark Horse'

SL: $D^0 \rightarrow l^- \nu K^+$ vs. $D^0 \rightarrow l^+ \nu K^-$

$$R_M^{SL} = (1.7 \pm 3.9) \times 10^{-4}, R_M^{WA} = (2.1 \pm 1.1) \times 10^{-4}$$

$$\text{if } x_D = y_D = 7 \times 10^{-3}, \text{ then } R_M = 5 \times 10^{-5}$$

$$a_{SL} \sim \text{Min}[\Delta\Gamma/\Delta M, \Delta M/\Delta\Gamma] \sin\phi_{NP}, \quad \Delta\Gamma/\Delta M \sim O(1)$$

• $a_{SL} \sim 0.1$ conceivable (even few $\times 0.1$)

-- i.e. relatively few wrong-sign leptons, yet with a large asymmetry!

vs.

• $a_{SL}(K_L) = 3.3 \times 10^{-3}$ with $\Delta\Gamma/\Delta M \sim O(1)$ & $\sin\phi_{CKM,eff} \ll 1$

• $a_{SL}(B_d) \sim 4 \times 10^{-4}$ with $\Delta\Gamma/\Delta M \sim O(\text{few} \times 10^{-3})$

• $a_{SL}(B_s) \sim 2 \times 10^{-5}$ with $\Delta\Gamma/\Delta M \sim O(\text{few} \times 10^{-3})$
& $\sin\phi_{CKM,eff} \sim O(\text{few} \times 10^{-2})$

(2.3) Benchmarks

☞ Allowed **New Physics** scenarios could produce ~~CP~~ close to present **experim. bounds**, but **hardly higher!**

○ **time dependant CP asymmetries** in

☞ $D^0 \rightarrow K^+ K^-, \pi^+ \pi^-, K_S \phi$ down to $O(10^{-4})$

☞ $D^0 \rightarrow K^+ \pi^-$ down to $O(10^{-3})$

LHCb: $\sim 10^6$ $D^* \rightarrow D \pi \rightarrow [KK]_D \pi$ in 10^7 sec

○ **direct ~~CP~~** in partial widths of

☞ $D^\pm \rightarrow K_{S[L]} \pi^\pm$ down to $O(10^{-3})$

☞ in a host of **1xCS** channels down to $O(10^{-3})$

☞ in **2xCS** channels down to $O(10^{-2})$

○ **direct ~~CP~~** in the **final state distributions**:

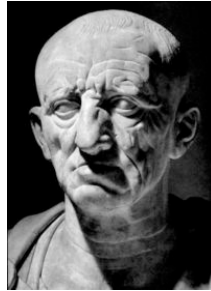
Dalitz plots, T-odd correlations etc. down to $O(10^{-3})$

IV Conclusions & Outlook

- a lot of work of **great importance** to be done
 - establish $(x_D, y_D) \neq 0$
 - determine $x_D = ?$ vs. $y_D = ?$
 - **go after CP** ← main message
 - in all of its possible manifestations
 - ☞ time dependent & independent,
 - ☞ partial widths, Dalitz plots, T odd moments ...
 - and on all Cabibbo levels
 - (i) $D^0 \rightarrow K_S \pi^+ \pi^- / K_S K^+ K^-$
 - (ii) $D^0 \rightarrow \pi^+ \pi^- / K^+ K^-$
 - (iii) $D^0 \rightarrow K^+ \pi^-$
 - down to the 0.001 (or even better) level
 - ☞ present **no-signal not** telling!
- can expect a positive learning curve for theorists -- yet do not count on miracles

The Big Picture

- detailed study of charm decays provides a novel & possibly unique window onto flavour dynamics
 - ➔ need the statistical muscle of the LHCb
 - interesting and worthy challenge for LHCb
 - $D^0 \rightarrow K^+K^-, \pi^+\pi^-, K^+\pi^-, K^+K^-\mu^+\mu^-$ good channels for LHCb
- 👉 yet need more statistics & channels!



"Ceterum Censeo Fabricam Super Saporis Esse Faciendam"
"Moreover I Advise a Super-Flavour Factory has to be Built"

example for a unique capability:

$$e^+ e^- \rightarrow \psi''(3770) \rightarrow D\bar{D} \rightarrow (K^+K^-)_D D_L \rightarrow K^+K^-\mu^+\mu^-$$

It is the task of the physicist to make the greatest use of a special gift from Nature

D^0 oscillations are such a gift

→ it is your duty -- & there is fame within your grasp!