# CMS Higgs Couplings and Spin/CP 



## Outline

- Couplings
- Signal strengths
- Coupling modifiers
- Vector bosons and fermions
- Generic modifier ratios
- New physics


## - Spin/CP

- Exotic spin
- Spin 0 anomalous couplings

Up to $5.1 \mathrm{fb}^{-1}$ ( 7 TeV ) and $19.7 \mathrm{fb}^{-1}$ (8 TeV)
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## Combined Higgs Channels

- Comprehensive set of production and decay modes targeted
- Over 250 mutually exclusive event categories


## Event category targets

$\checkmark$ Included in coupling combinations
$\checkmark$ Considered in certain interpretations


- $\mathrm{m}_{\mathrm{H}}=125.0 \mathrm{GeV}$ and narrow-width approximation assumed
- Off-shell measurements treated separately
- See David Sperka's talk yesterday on CMS diboson results


## Signal Strength (б/бзм)

- Best fit signal strength for production- and decay-tag pairs
- Tag by production and decay mode expected to dominate sensitivity in SM
- All signal contributions to tag pair scaled together
- p -value wrt SM $=0.84$
- Overall combination
$1.00 \pm 0.09$ (stat) ${ }_{-0.07}^{+0.08}$ (theo) $\pm 0.07$ (syst)
- Theory uncertainties: QCD scales PDFs, branching fractions, underlying event


## Coupling Modifiers ( $\mathrm{K}_{\mathrm{i}}$ )

- With additional assumption that signal arises from single particle with $\mathrm{JPC}=0^{++}$,

$$
(\sigma \mathcal{B})(x \rightarrow \mathrm{H} \rightarrow y y)=\frac{\sigma_{x} \Gamma_{y y}}{\Gamma_{\mathrm{tot}}}
$$

- $\Gamma_{\text {tot }}=\Sigma \Gamma_{\text {ii }}+\Gamma_{\text {BSM }}$, where $\Gamma_{\text {BSM }}=\Gamma_{\text {inv }}+\Gamma_{\text {undet }}$
- Introduce coupling modifiers ( $\mathrm{K}_{\mathrm{i}}$ ) to test for deviations from SM
- Production: Ki $^{2}=\sigma_{i} / \sigma_{i}{ }^{S M}$
- Decay: $\mathrm{ki}^{2}=\Gamma_{\mathrm{ij}} / \Gamma_{\mathrm{ij}} \mathrm{SM}$
- Total width: $\mathrm{KH}^{2}=\Gamma_{\text {tot }} / \Gamma_{\mathrm{sm}}$


## Couplings to Massive Vector Bosons and Fermions



## Test of custodial symmetry

- Likelihood scan of $\lambda_{w z}=K w / K z$ while profiling $k z$ and $K_{f}$
- Assume single $\mathrm{Kf}_{\mathrm{f}}$ and $\Gamma_{\text {BSm }}=0$
- Consistent with SM value of 1, resulting from protection against large radiative corrections


## Couplings to vector boson and fermions

- 68\% CL regions for kv and Kf
- Assume $\Gamma_{\mathrm{Bsm}}=0$
- Shows complementarity of combined channels
- Consistent with SM value of $(1,1)$



## Generic Modifier Ratio Model

- Most general model proposed by LHCXSWG (arXiv:1307.1347)
- Parameters are
- $\mathrm{Kgz}=\mathrm{Kg} \mathrm{Kz} / \mathrm{KH}$, where $\mathrm{KH}^{2}=\Gamma_{\mathrm{tot}} / \Gamma_{\mathrm{sm}}$ modifies the width
- Ratios of couplings $\lambda_{i j}=\mathrm{K}_{\mathrm{i}} / \mathrm{K}_{\mathrm{j}}$
- No assumption on scaling of total width.

- Most significant deviation is Parameter value driven by excess in ttH channels


## New Physics

## In loops

- ggH production and $\mathrm{H} \rightarrow \mathrm{\gamma} \mathrm{\gamma}$ decay are loopinduced at leading order
- Likelihood scan of $\mathrm{K}_{\mathrm{g}}$ and $\mathrm{k}_{\mathrm{y}}$ assuming SM tree-level couplings and $\Gamma_{\mathrm{BSM}}=0$
- Best fit $\left(\mathrm{Kg}_{\mathrm{g}}, \mathrm{K}_{\mathrm{y}}\right)=(1.14,0.89)$ is compatible with SM within $68 \%$ CL region



## Undetected and invisible decays

- Include $\mathrm{H} \rightarrow$ inv search results to constrain $\mathrm{BR}_{\text {inv }}=\Gamma_{\text {inv }} / \Gamma_{\text {tot }}$
- Uncombined: BRinv observed (expected) 95\% CL upper limit = 0.58 (0.44)
- Simultaneous fit for $\mathrm{BR}_{\text {inv }}$ and $\mathrm{BR}_{\text {undet }}=$ $\Gamma_{\text {undet }} / \Gamma_{\text {tot }}$ while profiling $\mathrm{K}_{\mathrm{v}}, \mathrm{K}_{\mathrm{g}}, \mathrm{Kv} \leq 1, \mathrm{~Kb}_{\mathrm{b}}, \mathrm{K}_{\mathrm{T}}, \mathrm{K}_{\mathrm{t}}$
- Very general!


## Exotic Spin Scenarios

- Spin-two
- with gravity-like minimal couplings excluded at 99.87\% CL in combination of $\mathrm{H} \rightarrow \mathrm{ZZ}, \mathrm{H} \rightarrow \mathrm{WW}$, and $\mathrm{H} \rightarrow \gamma \gamma$.
- Another ten models excluded at 99\% CL or higher.
- Any mixed-parity spin-one state is excluded at $>99.999 \% \mathrm{CL}$ in combination of $\mathrm{H} \rightarrow \mathrm{ZZ}$ and $\mathrm{H} \rightarrow \mathrm{WW}$
- Fraction of non-interfering exotic spin state in addition to $\mathrm{JP}^{\mathrm{P}}=\mathrm{O}^{+}$state also considered.




## Spin 0 Anomalous Couplings Phenomenology

- Generic $\mathrm{HV}_{1} \mathrm{~V}_{2}(\mathrm{~V}=\mathrm{W}, \mathrm{Z}, \mathrm{V}, \mathrm{g})$ scattering amplitude, expanded up to $\mathrm{q}^{2}$

$$
\begin{aligned}
& A(\mathrm{HVV}) \sim\left[a_{1}^{\mathrm{VV}}+\frac{\kappa_{1}^{\mathrm{VV}} q_{\mathrm{V} 1}^{2}+\kappa_{2}^{\mathrm{VV}} q_{\mathrm{V} 2}^{2}}{\left(\Lambda_{1}^{\mathrm{VV}}\right)^{2}}\right] \\
& \quad \text { tree level scalar (0+) leading momentum }
\end{aligned} m_{\mathrm{V} 1}^{2} \epsilon_{\mathrm{V} 1}^{*} \epsilon_{\mathrm{V} 2}^{*}+a_{2}^{\mathrm{VV}} f_{\mu \nu}^{*(1)} f^{*(2), \mu \nu}+a_{3}^{\mathrm{VV}} f_{\mu \nu}^{*(1)} \tilde{f}^{*(2), \mu \nu}
$$

## higher order scalar

| Interaction | Anomalous Coupling | Coupling Phase | Effective Fraction |
| :---: | :---: | :---: | :---: |
| HZZ | $\Lambda_{1}$ | $\phi_{\Lambda 1}$ | $f_{\Lambda 1}$ |
|  | $a_{2}$ | $\phi_{a 2}$ | $f_{a 2}$ |
|  | $a_{3}$ | $\phi_{a 3}$ | $f_{a 3}$ |
| HWW | $\Lambda_{1}^{\text {WW }}$ | $\phi_{\wedge 1}^{\mathrm{WW}}$ | $f_{\wedge 1}^{W W}$ |
|  | $a_{2}^{\text {WW }}$ | $\phi_{\text {a }}{ }^{W}{ }^{W}$ | $f_{\text {a }}{ }^{W} \mathrm{~W}$ |
|  | $a_{3}^{\text {WW }}$ | $\phi_{a 3}^{W W W}$ | $f_{a 3}^{W W}$ |
| $\mathrm{HZ} \gamma$ | $\Lambda_{1}^{\mathrm{Z} \gamma}$ | $\phi_{11}^{\mathrm{Z} \gamma}$ | $f_{\Lambda 1}^{Z \gamma}$ |
|  | $a_{2}^{Z \gamma}$ |  | $f_{a_{2}}^{Z \gamma}$ |
|  | $a_{3}^{Z_{\gamma}}$ | $\phi_{a 3}^{Z,}$ | $f_{a 3}^{Z \gamma}$ |
| $\mathrm{H} \gamma \gamma$ | $a_{2}^{\gamma \gamma}$ | $\phi_{02}^{\gamma \gamma}$ | $f_{0}^{\gamma \gamma}$ |
|  | $a_{3}^{\gamma \gamma}$ | $\phi_{a 3}^{\gamma 2 \gamma}$ | $f_{a 3}^{\gamma \gamma}$ |

Example phase and effective fraction:

$$
\phi_{a 2}=\arg \left(\frac{a_{2}}{a_{1}}\right) \quad f_{a 2}=\frac{\left|a_{2}\right|^{2} \sigma_{2}}{\left|a_{1}\right|^{2} \sigma_{1}+\left|a_{2}\right|^{2} \sigma_{2}+\left|a_{3}\right|^{2} \sigma_{3}+\tilde{\sigma}_{\Lambda_{1}} /\left(\Lambda_{1}\right)^{4}}
$$

where $\sigma_{i}$ is the cross section for $a_{i}=1$ and $a_{j \neq i}=0$
One non-zero anomalous coupling:
A. real, $\phi_{\mathrm{a}}=0, \pi$
B. complex, $\phi_{\text {ai }}$ unconstrained Two non-zero anomalous couplings:
C. real, $\phi_{\mathrm{a}, \mathrm{a}, \mathrm{j}}=0, \pi$
D. complex, $\phi_{\mathrm{ai}, \mathrm{aj}}$ unconstrained

## Simulated with JHUGen or POWHEG+JHUGen

## Observables



- Use 5 angles and 3 masses to describe $\mathrm{H} \rightarrow \mathrm{VV} \rightarrow 4$ l kinematics
- Matrix elements define event by event probabilities for observed kinematics (MELA)
- Construct kinematic discriminants from probabilities

$$
\text { e.g. } \mathcal{D}_{J^{p}}=\frac{\mathcal{P}_{\mathrm{SM}}}{\mathcal{P}_{\mathrm{SM}}+\mathcal{P}_{J^{p}}}
$$

- $\mathrm{H} \rightarrow \mathrm{WW} \rightarrow$ Ivlv contains reduced information due to v's
- Use $m_{\|}$and $m_{T}$ distinguish signal models




## HZZ: complex ai

## Example: $f_{\mathrm{a} 2}$


all consistent with SM



## HZZ: real $\mathrm{a}_{i}$

- Couplings assumed to be real, so $\phi_{\mathrm{ai}}=0$ or $\pi$ and $\cos \left(\phi_{\mathrm{ai}}\right)=1$ or -1





CMS Couplings and Spin/CP


all consistent with SM


## $\mathrm{H} \rightarrow \mathrm{ZZ}+\mathrm{H} \rightarrow \mathrm{WW}$ Combination

- A priori, no relationship between HZZ and HWW couplings
- Combine $\mathrm{H} \rightarrow \mathrm{ZZ}$ and $\mathrm{H} \rightarrow \mathrm{WW}$ after assuming a relationship

$$
r_{a i}=\frac{a_{i}^{\mathrm{WW}} / a_{1}^{\mathrm{WW}}}{a_{i} / a_{1}}, \text { or } R_{a i}=\frac{r_{a i}\left|r_{a i}\right|}{1+r_{a i}^{2}}
$$

- Custodial symmetry implies $a_{1}=a_{1} w w$



$$
a_{1}=a_{1} w w
$$



$$
\mathrm{r}_{\mathrm{ai}}=1, \text { or } \mathrm{R}_{\mathrm{a} 3}=0.5
$$

$f_{\wedge 1}$ and $f_{a 2}$ in backup all consistent with SM

- Additional term depending only on invariant mass of Higgs boson

$$
A(\mathrm{HVV}) \propto[a_{1} \underbrace{}_{+a_{2} f_{\mu \nu}^{*(1)} f^{*(2), \mu v}+a_{3} f_{\mu \nu}^{*(1)} \tilde{f}^{*(2), \mu v}, \frac{\left(q_{\mathrm{v} 1}+q_{\mathrm{V} 2}\right)^{2}}{\left(\Lambda_{Q}\right)^{2}}-e^{i \phi_{\Lambda 1}} \frac{\left(q_{\mathrm{V} 1}^{2}+q_{\mathrm{V} 2}^{2}\right)}{\left(\Lambda_{1}\right)^{2}} m_{\mathrm{V} 1}^{2} \epsilon_{\mathrm{V} 2}^{*} \epsilon^{*}}
$$

- Must be tested in off-shell region

$$
f_{\Lambda Q}=\frac{m_{\mathrm{H}}^{4} / \Lambda_{Q}^{4}}{\left|a_{1}\right|^{2}+m_{\mathrm{H}}^{4} / \Lambda_{\mathrm{Q}}^{4}}
$$

- Joint constraint on width and $\wedge_{Q}$ anomalous coupling





## Conclusions

- Comprehensive sets of Higgs measurements combined to test compatibility of couplings with SM
- Constraints placed on exotic spin states and spin-zero anomalous couplings,
- Including new results on f^Q
- All observations are consistent with the standard model scalar JPC= $0^{++}$





## Backup

## Scaling of couplings with mass

- Phenomenological parameterization relating masses to coupling modifiers with two parameters
$-\mathrm{K}_{\mathrm{f}}=\mathrm{v} \mathrm{mf}_{\mathrm{f}} / \mathrm{M}^{1+\varepsilon}$
$-k v=v m v^{2 \varepsilon} / M^{1+2 \varepsilon}$
- SM recovered for $(\mathrm{M}, \varepsilon)=(\mathrm{v}, 0)$, where $v=246 \mathrm{GeV}$
- Assume
- Coupling to massive SM particles only, one parameter per tree-level coupling
- SM loop structure




## $\mathrm{H} \rightarrow \mathrm{VV} \rightarrow 4 \mid$ Kinematics










## $\mathrm{H} \rightarrow \mathrm{ZZ}+\mathrm{H} \rightarrow$ WW Combination




Ben Kreis

## $\mathrm{HZ} \gamma$ and $\mathrm{H} \gamma \gamma$

- $\mathrm{H} \rightarrow \mathrm{VV} \rightarrow 4 \mathrm{I}$, where $\mathrm{VV}=\mathrm{Z} \gamma^{*}, \gamma^{*} \gamma^{*}$
- Currently, not competitive with direct cross section measurements from on-shell $\mathrm{H} \rightarrow \mathrm{Zy}$ or $\mathrm{H} \rightarrow \mathrm{\gamma} \mathrm{\gamma}$
- However, with sufficient luminosity, $f_{a 3} V_{r}$ and $f_{a 2} V_{r}$ can be measured separately in this channel. Also $f_{\wedge 1} z_{\gamma}$.




