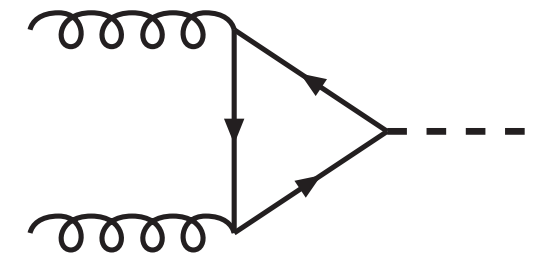
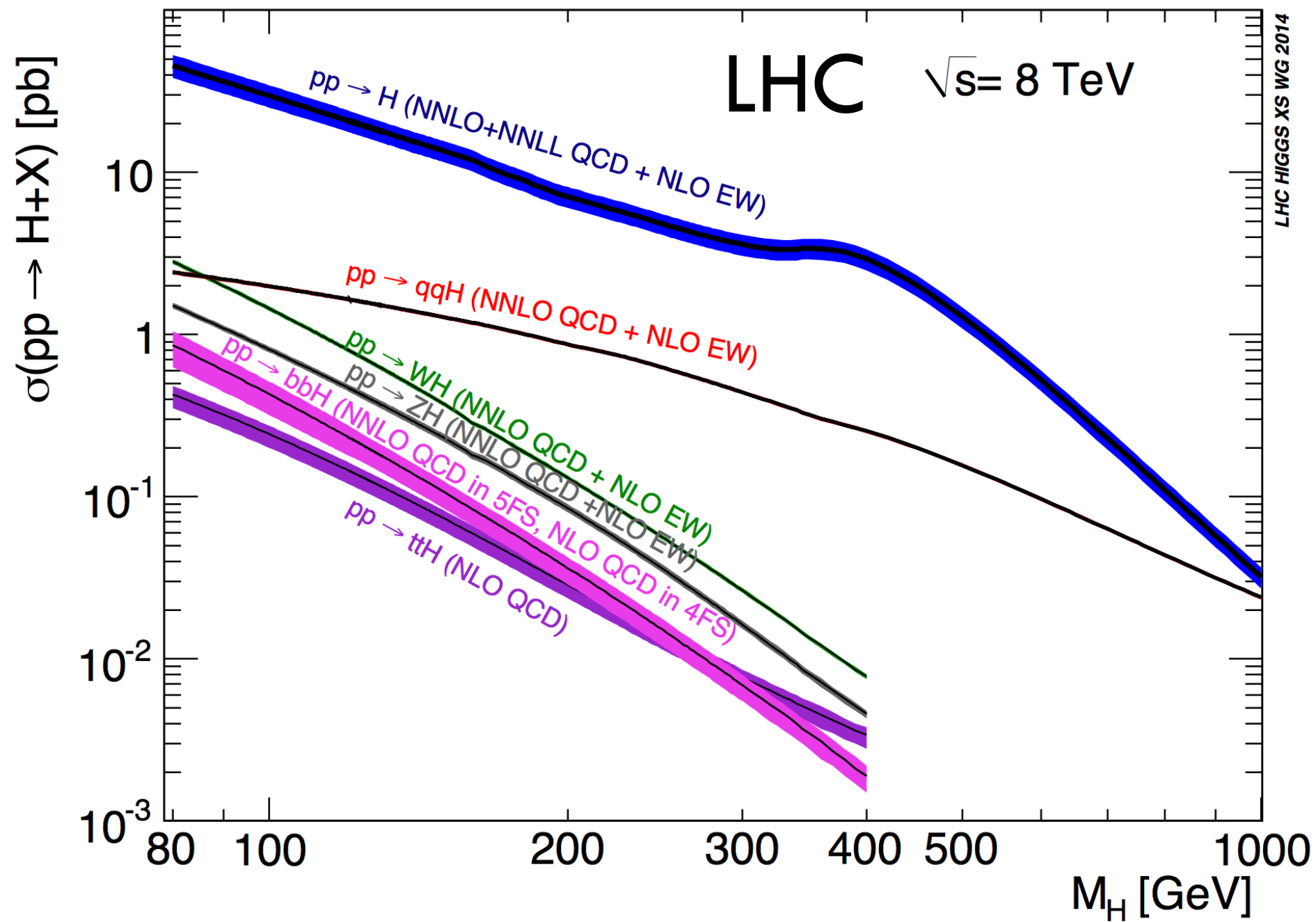


Theory Higgs production

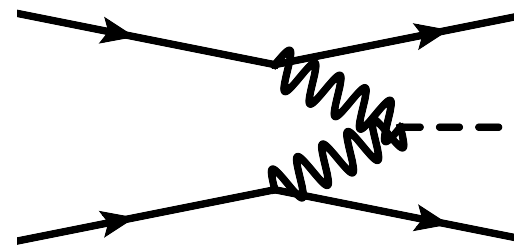
Claude Duhr
in collaboration with C. Anastasiou, F. Dulat, E. Furlan,
T. Gehrmann, F. Herzog, A. Lazopoulos, B. Mistlberger

Higgs Hunting 2015
LAL Orsay, 30/07/2015

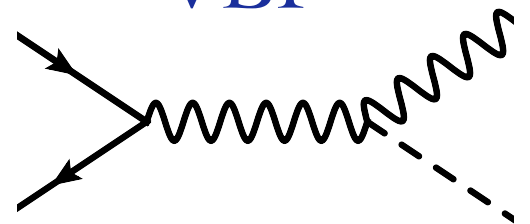
Higgs physics at the LHC



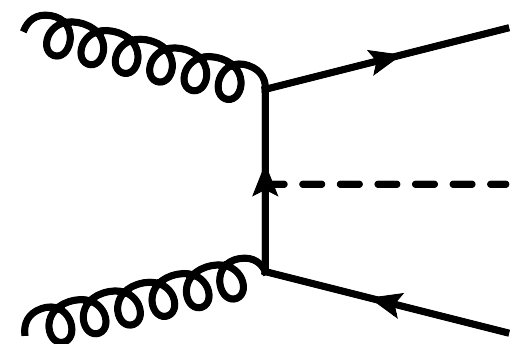
Gluon fusion



VBF

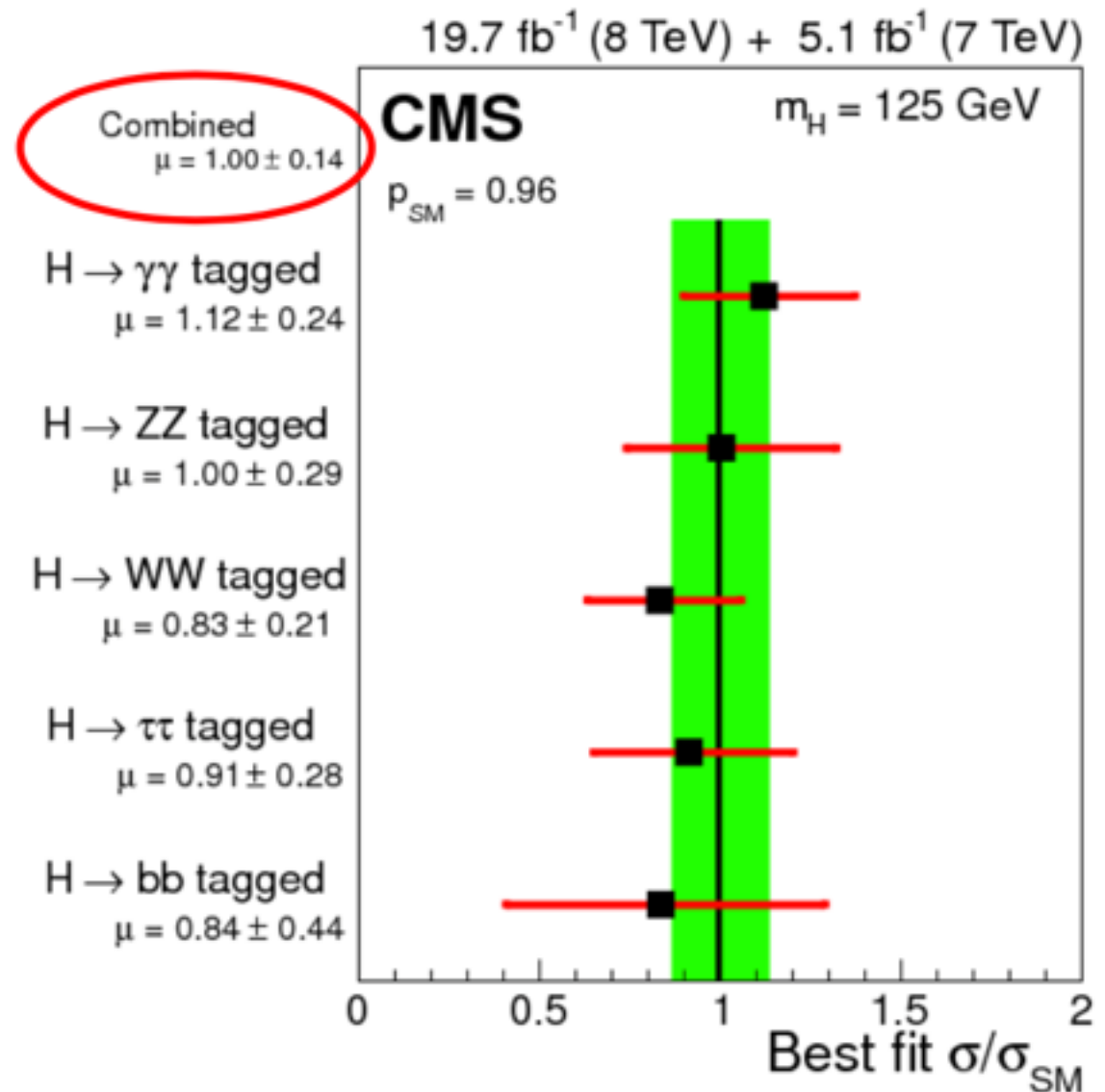
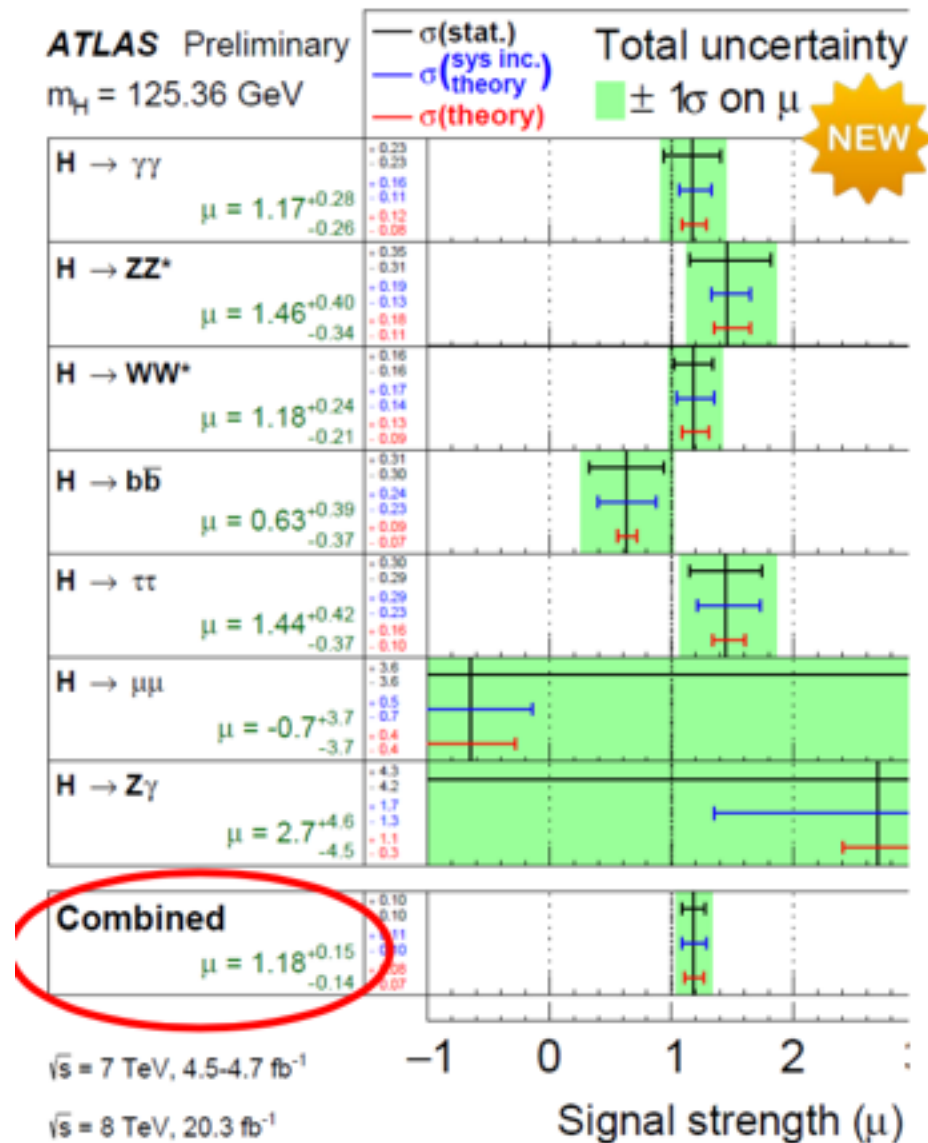


VH



TTH/bbH

Higgs physics at the LHC



$$\mu_{ATLAS} = 1.18^{+0.15}_{-0.14}$$

$$\text{stat.} = {}^{+0.10}_{-0.10}$$

$$\text{sys. (inc. theo.)} = {}^{+0.11}_{-0.10}$$

$$\text{theory} = {}^{+0.08}_{-0.07}$$

$$\mu_{CMS} = 1.00 \pm 0.14$$

[M. Dührssen @ Moriond EW 2015]

Higgs physics at the LHC

- **Outline:** important theoretical advances in predictions for Higgs production in the last few months:
 - ➔ Fully differential NNLO predictions for VBF.
 - ➔ Fully differential NNLO predictions for H + jet in gluon fusion.
 - ➔ Inclusive gluon fusion cross section at N³LO.
 - ➔ New PDF sets, with reduced uncertainty.
- These advances show the maturity of our tools to make precision computations in QCD!

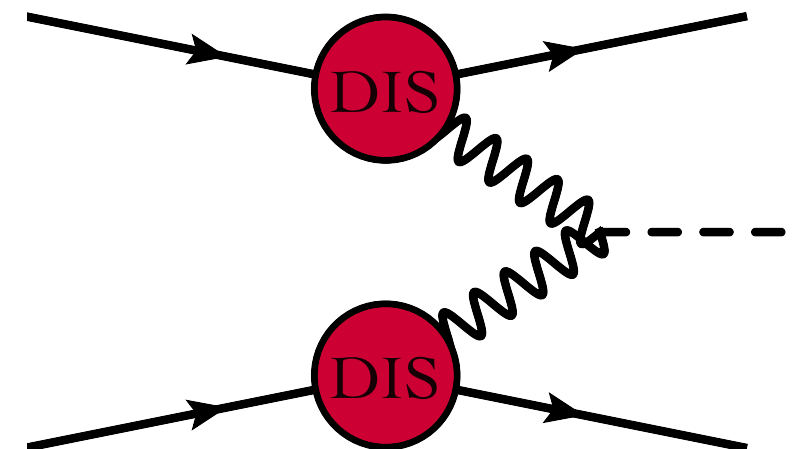
VBF @ NNLO

Vector boson fusion

- VBF is the 2nd largest production channel at the LHC.
 - ➔ Direct access to HVV coupling.
 - ➔ Non-zero H - p_T at leading order.
 - ➔ Radiation pattern allows one to disentangle ggH from VBF (VBF cuts).

Vector boson fusion

- VBF is the 2nd largest production channel at the LHC.
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 - ➔ Radiation pattern allows one to disentangle ggH from VBF (VBF cuts).
- LO process is 2-to-3: Two-loop corrections unknown!
- Form factor approach: no colour exchange between the two quark lines.
 - ➔ Exact at NLO.
 - ➔ $VBF = (DIS)^2$.
 - ➔ Used to compute inclusive VBF cross section at NNLO.



Vector boson fusion

- BMMZ: NNLO effects in inclusive cross section small ($\sim 1\%$).
- Inclusive cross section is not quite what we want:
 - ➔ No differential information.
 - ➔ Cannot impose VBF cuts.

Vector boson fusion

- BMMZ: NNLO effects in inclusive cross section small ($\sim 1\%$).
- Inclusive cross section is not quite what we want:
 - ➔ No differential information.
 - ➔ Cannot impose VBF cuts.
- Recently: first fully differential computation of VBF at NNLO in form factor approach. [Cacciari, Dreyer, Karlberg, Salam, Zanderighi]
 - ➔ Rather large NNLO corrections after VBF cuts ($\sim 5\text{-}6\%$)!

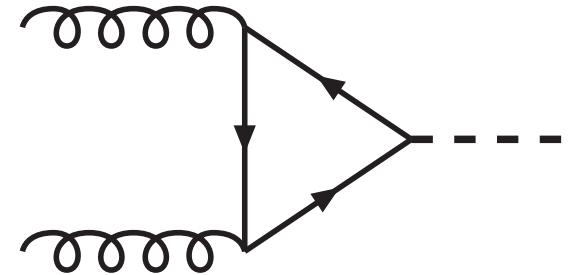
$p_{T,j_1}, p_{T,j_2} > 25 \text{ GeV}$		$\sigma^{(\text{no cuts})} [\text{pb}]$	$\sigma^{(\text{VBF cuts})} [\text{pb}]$
$ y_{j_1} , y_{j_2} < 4.5$		LO 4.032 $^{+0.057}_{-0.069}$	0.957 $^{+0.066}_{-0.059}$
$\Delta y_{j_1,j_2} > 4.5 \quad M_{j_1 j_2}^2 > (600 \text{ GeV})^2$		NLO 3.929 $^{+0.024}_{-0.023}$	0.876 $^{+0.008}_{-0.018}$
		NNLO 3.888 $^{+0.016}_{-0.012}$	0.826 $^{+0.013}_{-0.014}$
➔ See Dreyer's talk this afternoon!		~ 1	$\sim 5\text{-}6$

Gluon-fusion

The gluon fusion cross section

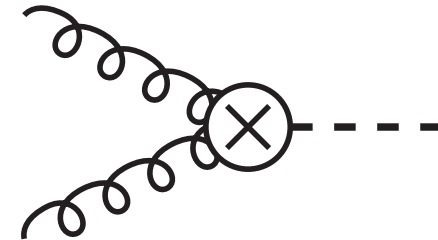
- The dominant Higgs production mechanism at the LHC is gluon fusion.

➔ Loop-induced process.



- For a light Higgs boson, the dimension five operator describing a tree-level coupling of the gluons to the Higgs boson

$$\mathcal{L} = \mathcal{L}_{QCD,5} - \frac{1}{4v} C_1 H G_{\mu\nu}^a G_a^{\mu\nu}$$



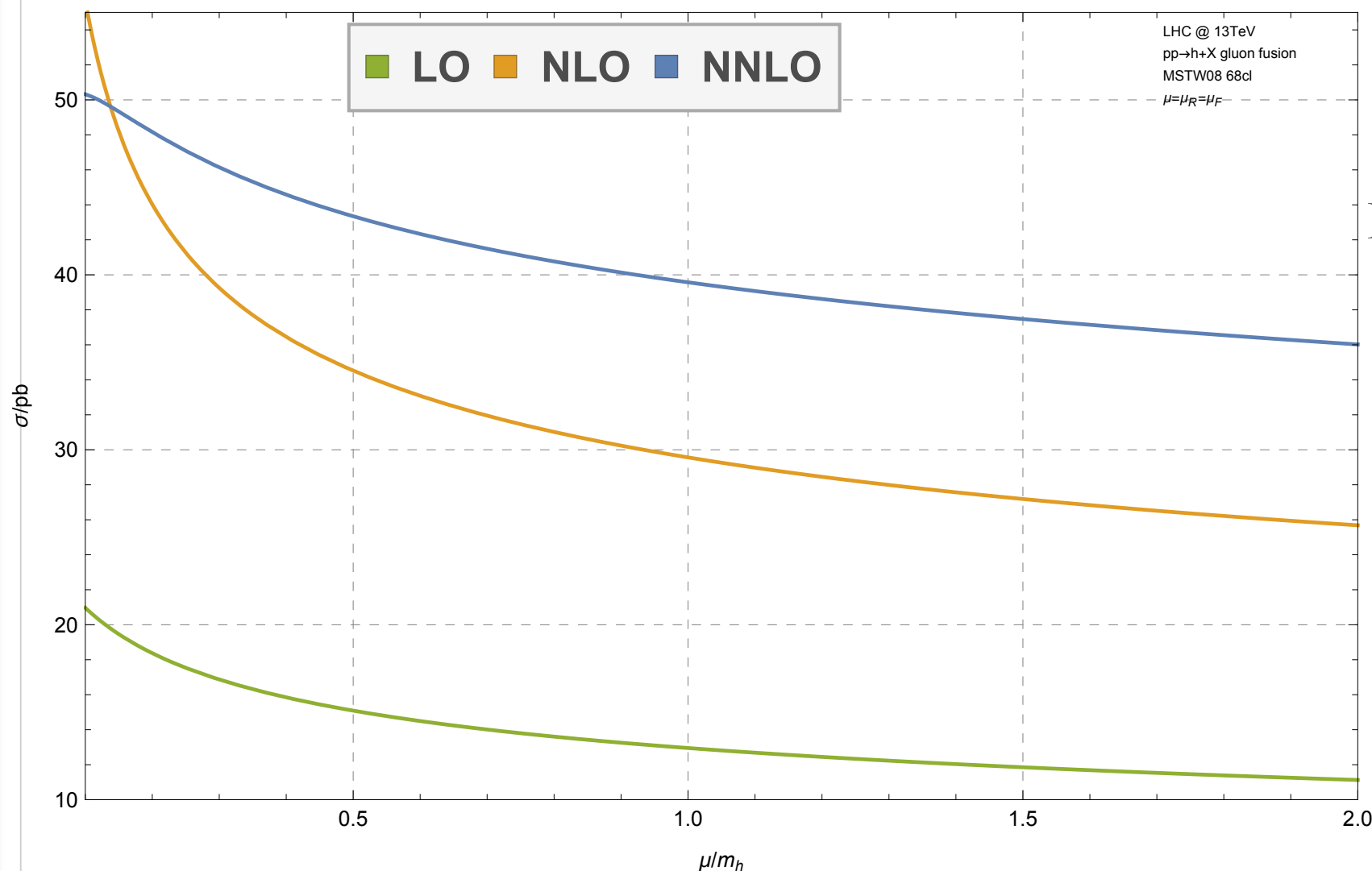
- Top-mass corrections known at NNLO.

[Harlander, Ozeren; Pak, Rogal, Steinhauser; Ball, Del Duca, Marzani, Forte, Vicini; Harlander, Mantler, Marzani, Ozeren]

- In the rest of the talk, I will only concentrate on the effective theory.

The gluon fusion cross section

- Known inclusively at NLO and NNLO, but plagued by large perturbative uncertainties.



[Dawson; Djouadi, Spira,
Zerwas; Harlander, Kilgore;
Anastasiou, Melnikov;
Ravindran, Smith, van Neerven]

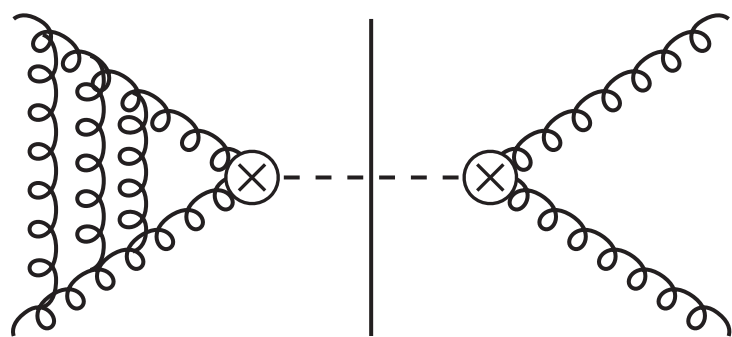
- Enormous progress over the last few months!
 - ➔ Both inclusively and differentially.

Gluon-fusion

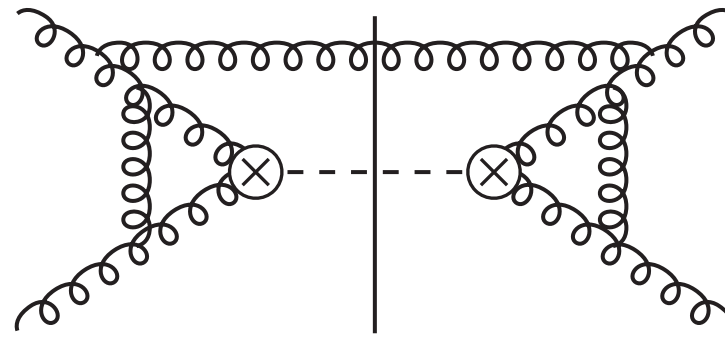
H @ N³LO

The gluon fusion cross section

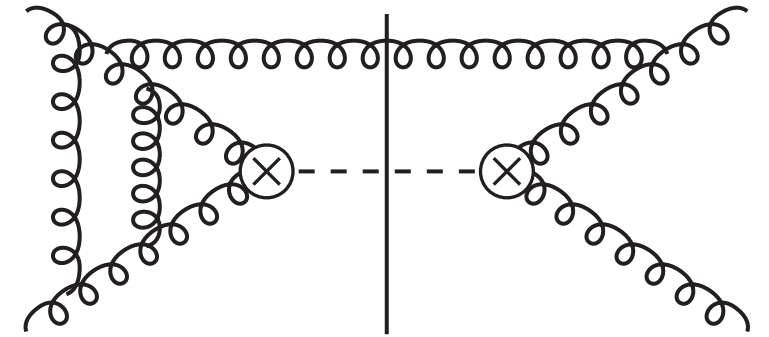
- At N³LO, there are five contributions:



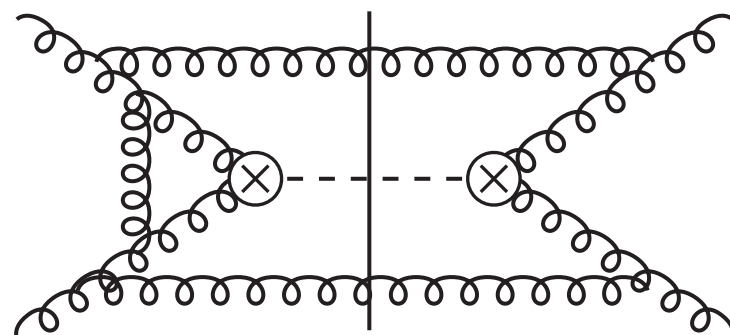
Triple virtual



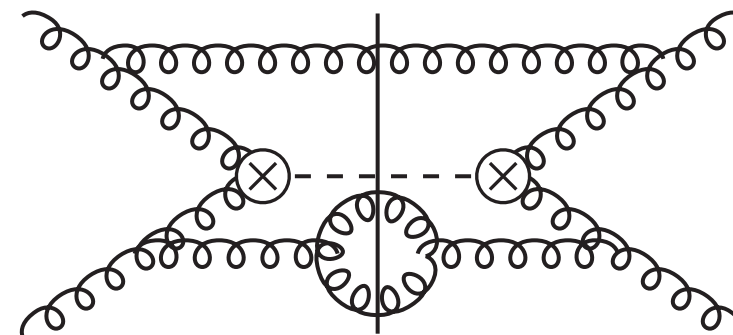
Real-virtual
squared



Double virtual
real



Double real
virtual



Triple real

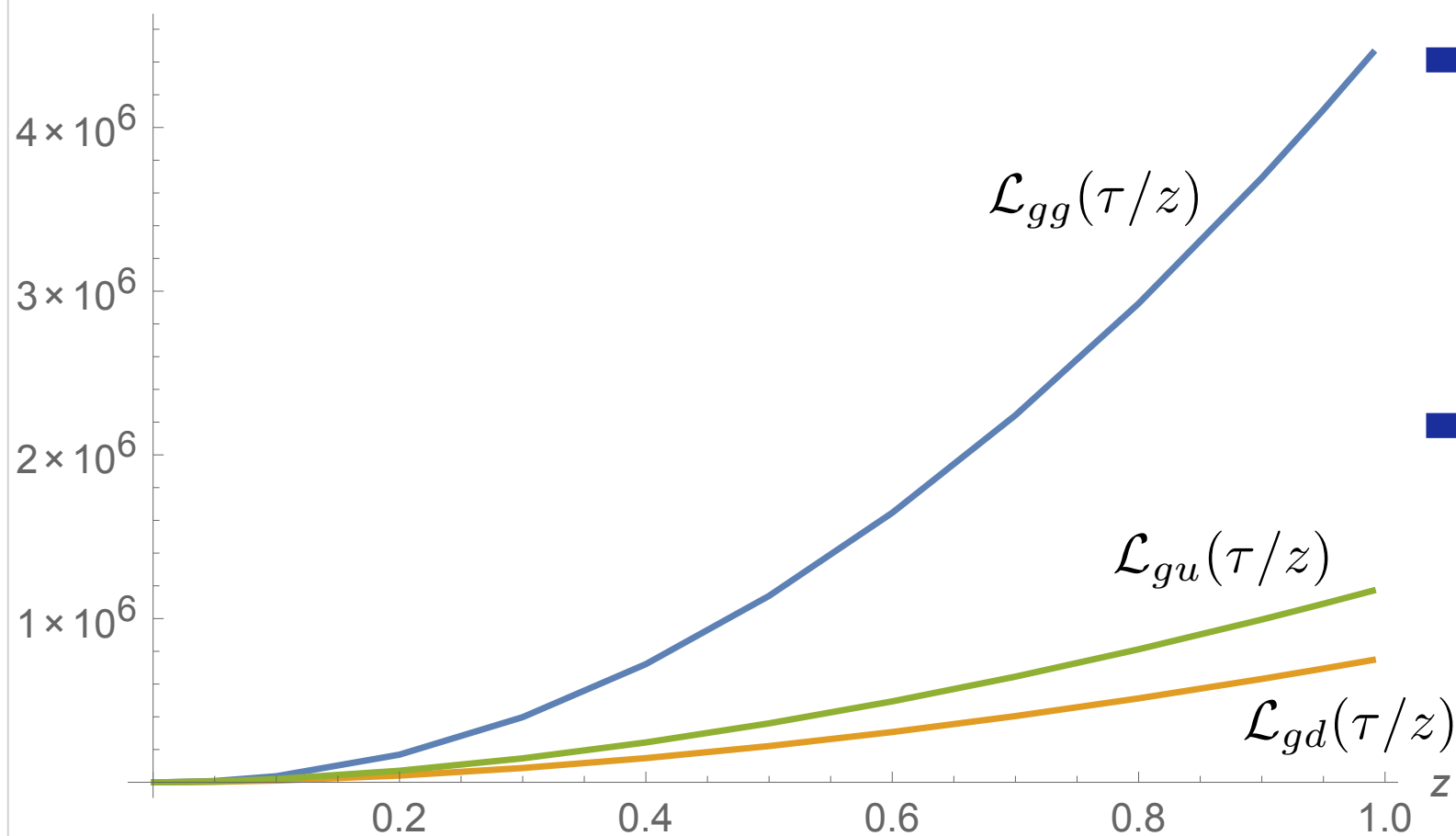
The gluon fusion cross section

- The gluon fusion cross section is given in perturbation theory by

$$\sigma = \tau \sum_{ij} \int_{\tau}^1 \frac{dz}{z} \mathcal{L}_{ij}(\tau/z) \frac{\hat{\sigma}_{ij}(z)}{z}$$

$$z = \frac{m_H^2}{\hat{s}}$$

$$\tau = \frac{m_H^2}{S} \simeq 10^{-4}$$



➔ Main contribution from region where $z \simeq 1$.

➔ Physically: production at threshold + emission of soft partons.

Systematics of the expansion

$$\frac{\hat{\sigma}_{ij}(z)}{z} = \hat{\sigma}^{SV} \delta_{ig} \delta_{jg} + \sum_{N=0}^{\infty} \hat{\sigma}_{ij}^{(N)} (1-z)^N$$

- Goal: Compute enough terms to establish convergence.

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- The first term is called the **soft-virtual** term and is distribution-valued:

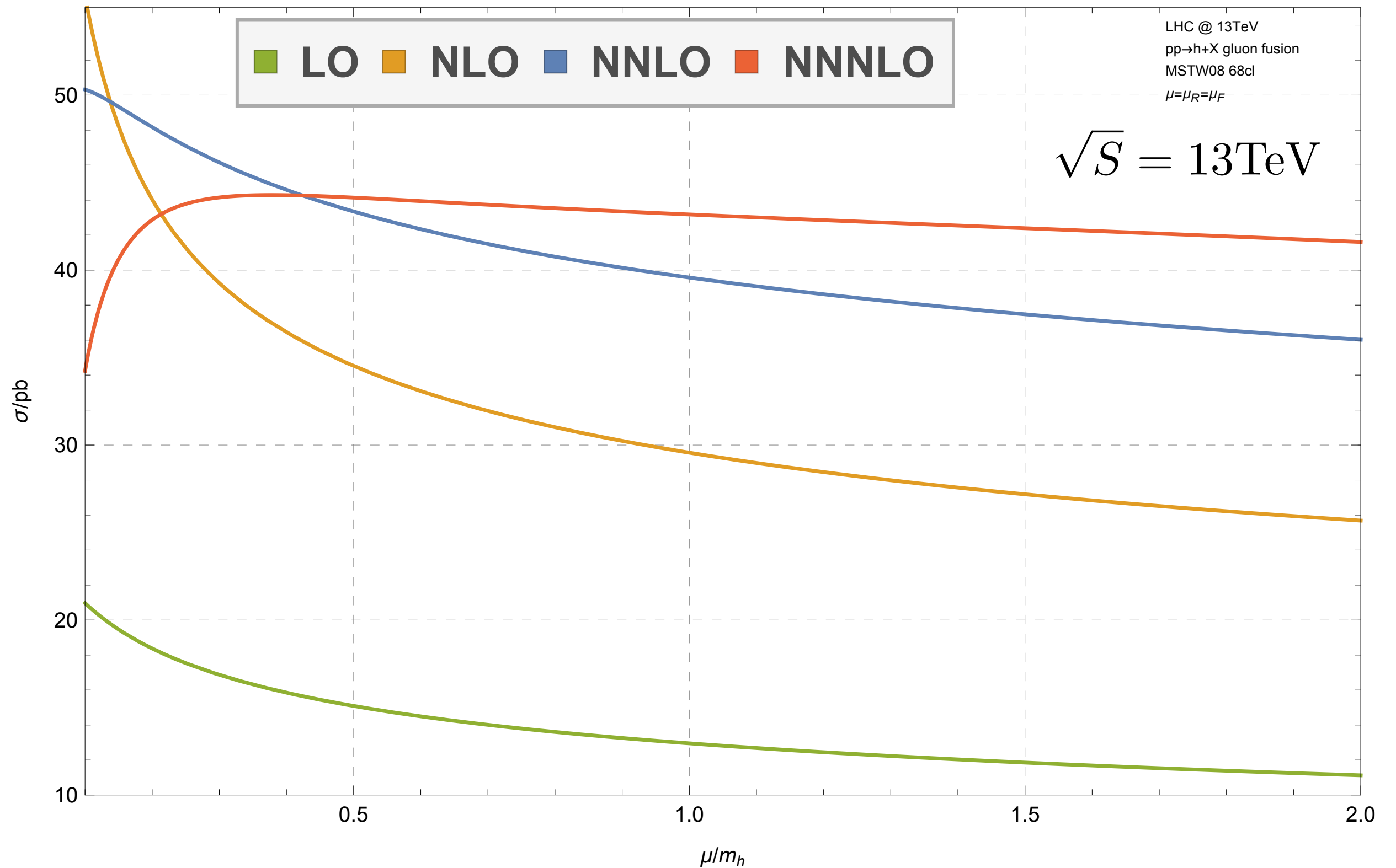
➔ **At N3LO:** $\hat{\sigma}^{SV} = a \delta(1-z) + \sum_{k=0}^5 b_k \left[\frac{\log^k(1-z)}{1-z} \right]_+$

Some numbers

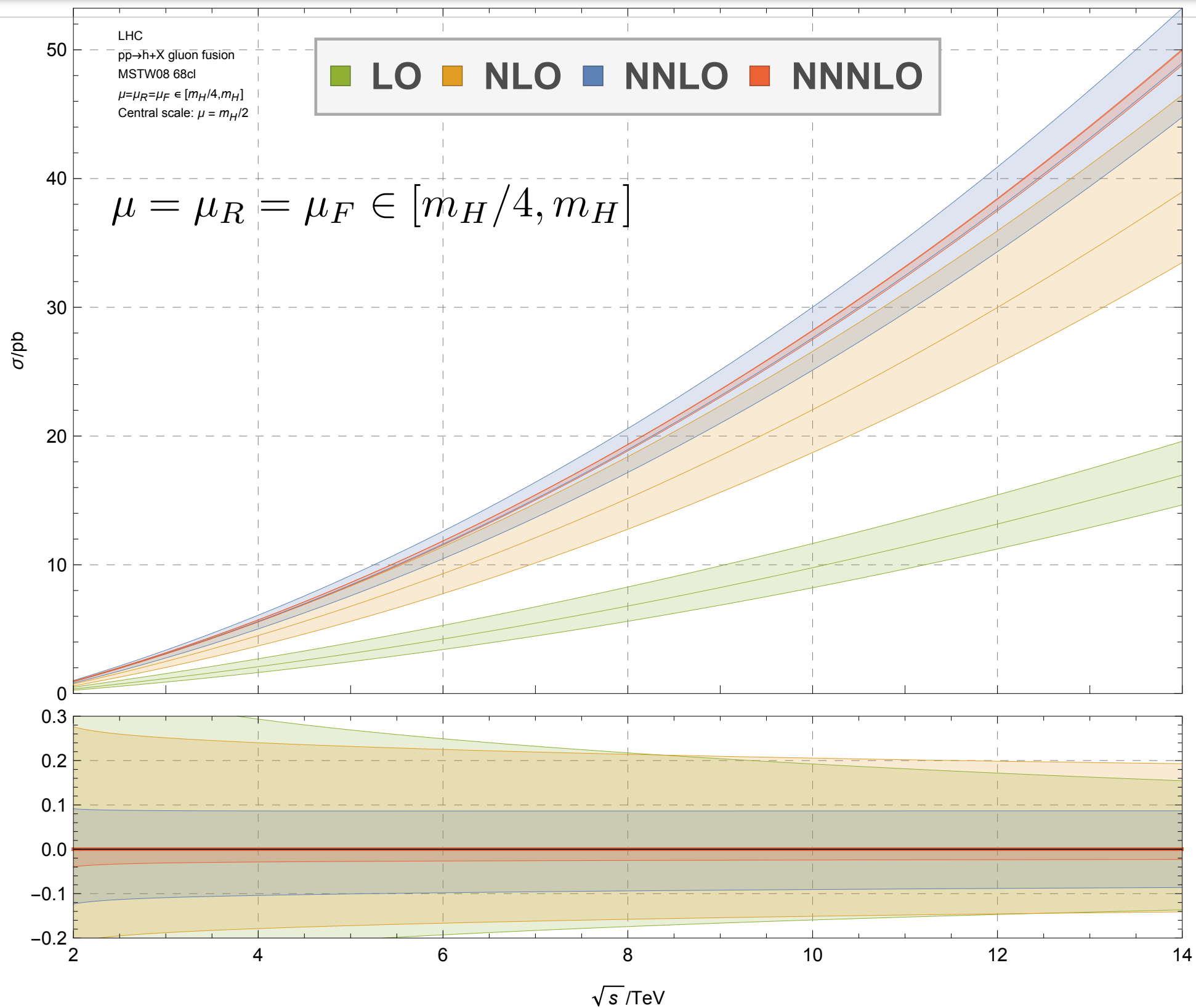
	NNLO	N3LO
# diagrams	~ 1.000	~ 100.000
# integrals	~ 50.000	517.531.178
# masters	27	1.028
# boundary conditions	5	78

- This brings you to the edge of what is technically possible at the moment.
 - ➔ A lot of cross talk with formal amplitudes community!

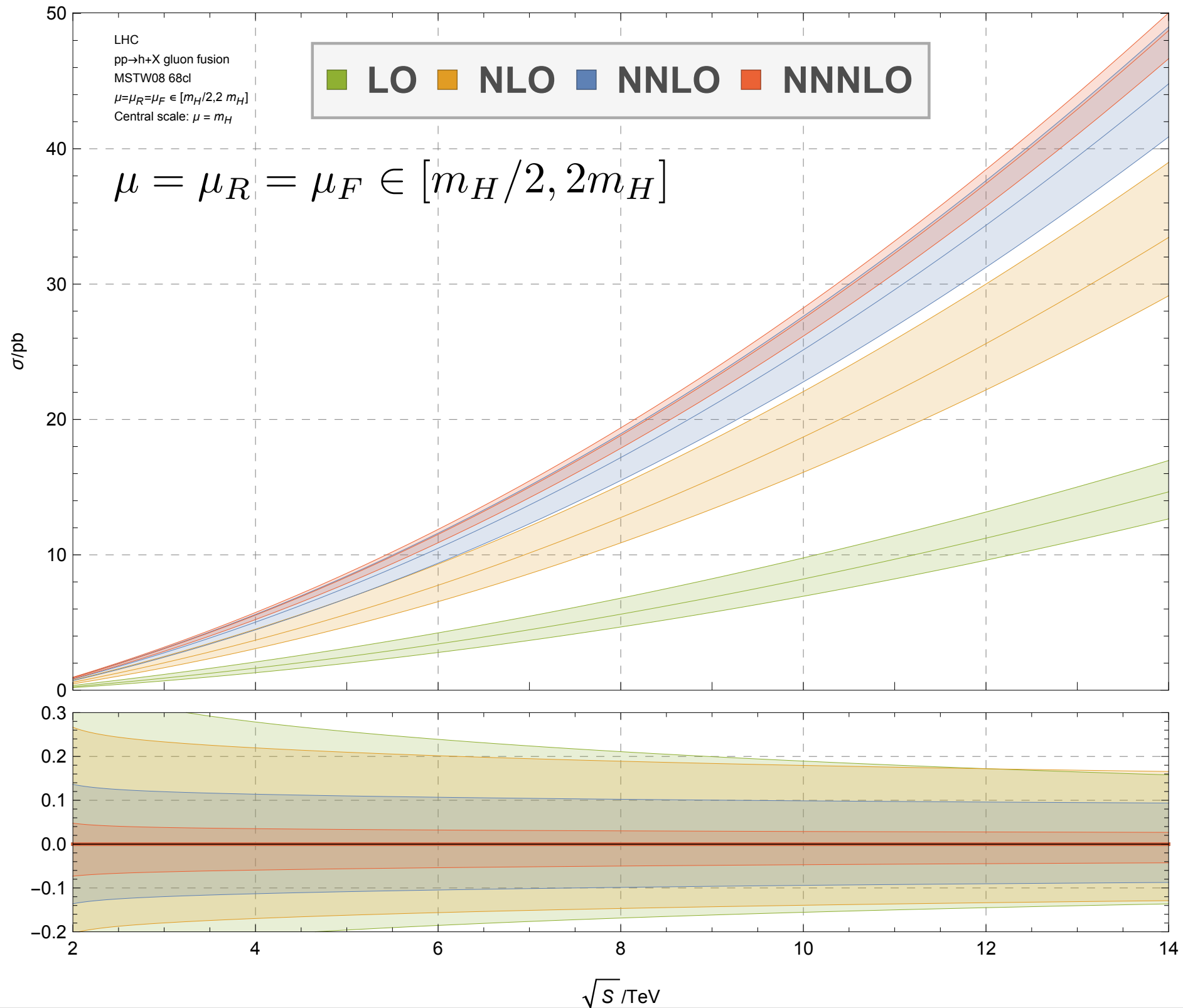
Scale variation



Energy variation



Energy variation



Uncertainties

- Remaining scale uncertainty at N³LO
 - ➔ We should think very carefully which other effects could be of the same size!
- Other sources of uncertainty:

Uncertainties

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 - ➔ $1/m_t$ corrections:

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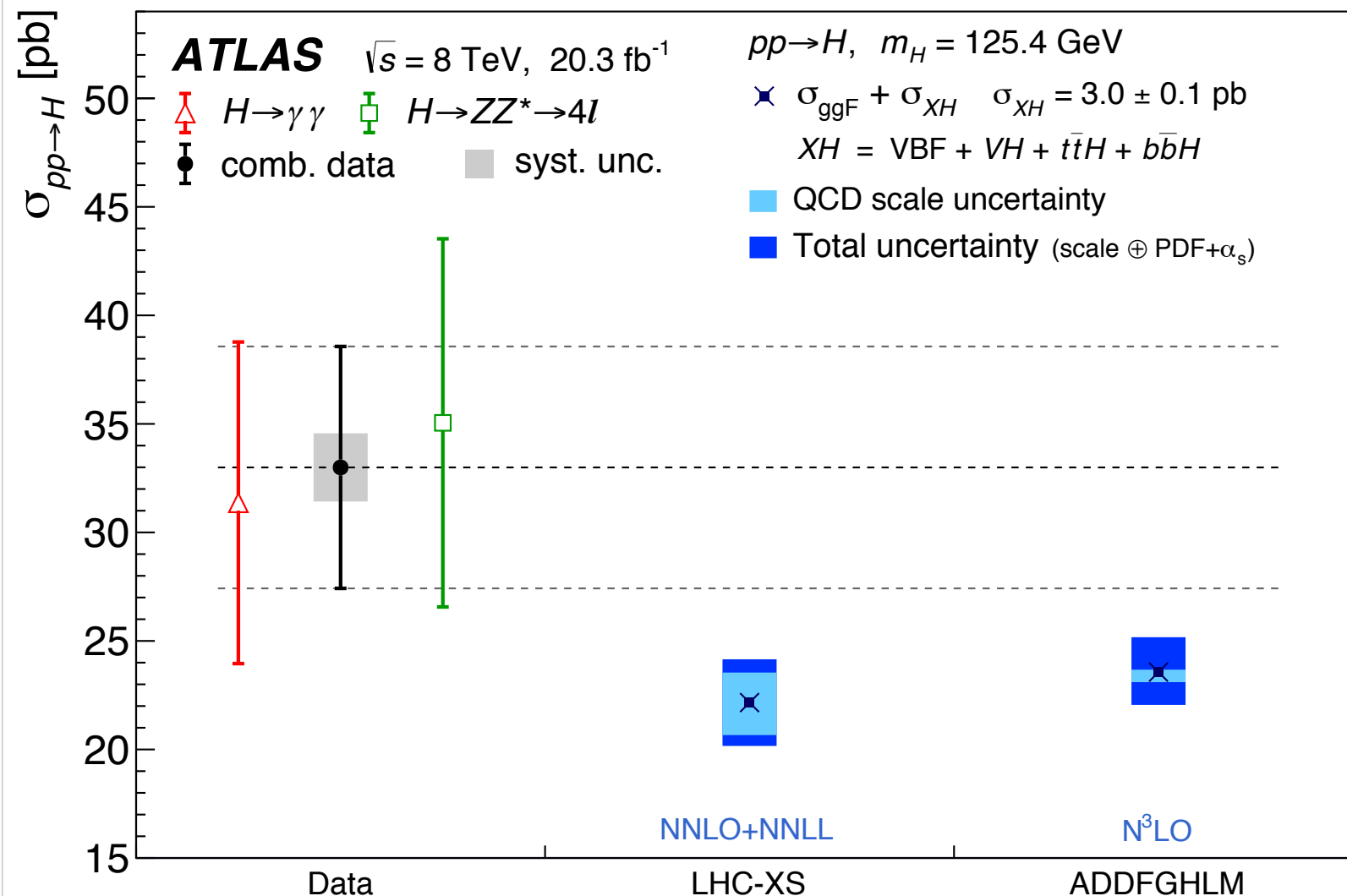
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 - ➔ PDF + α_S :

Scale vs. PDF uncertainty



	CT14	MMHT2014	NNPDF3.0	CT10
8 TeV	$18.66^{+2.1\%}_{-2.3\%}$	$18.65^{+1.4\%}_{-1.9\%}$	$18.77^{+1.8\%}_{-1.8\%}$	$18.37^{+1.7\%}_{-2.1\%}$
13 TeV	$42.68^{+2.0\%}_{-2.4\%}$	$42.70^{+1.3\%}_{-1.8\%}$	$42.97^{+1.9\%}_{-1.9\%}$	$42.20^{+1.9\%}_{-2.5\%}$

[CTEQ collaboration]

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 - ➔ PDF+ α_S : **$\sim 3\%$ with modern PDF sets**
 - ➔ NLO EW corrections:

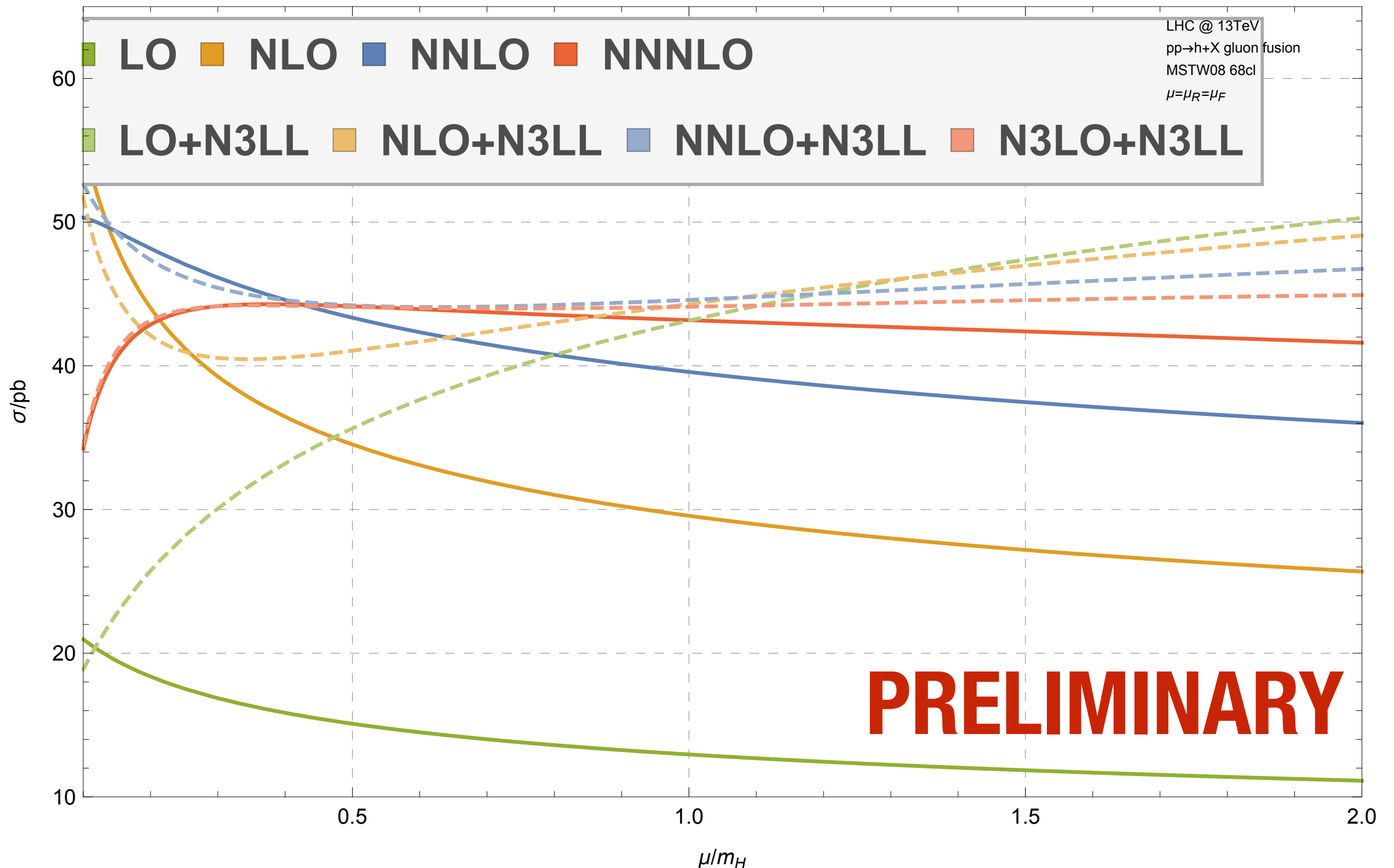
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[Djouadi, Gambino, Kniehl; Aglietti, Bonciani, Degrandi; Degrandi, Maltoni; Anastasiou, Boughezal, Petriello; Actis, Passarino, Sturm, Uccirati]
$$\sigma_0 (1 + \delta_{\text{QCD}} + \delta_{\text{EW}}) \quad \text{vs.} \quad \sigma_0 (1 + \delta_{\text{QCD}}) (1 + \delta_{\text{EW}})$$

Uncertainties

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 - ➔ Missing higher orders / threshold resummation

N3LL threshold resummation



Summary for H @ N3LO

- We are currently putting together all these effects (including different ‘flavours’ or threshold resummation).
- $\mu = m_H/2$ seems to be a good central scale choice.
 - ➔ Reduced scale uncertainty compared to $\mu = m_H$.
 - ➔ Series seems to converge.
 - ➔ Negligible impact of soft-gluon resummation.
 - ➔ Current recommendation of HXSWG: $\mu = m_H$.
- We are reaching the point where we should critically assess our method of estimating the uncertainty by scale variation!
 - ➔ A negligible scale variation does not mean that there are no more higher-order corrections!

Gluon-fusion

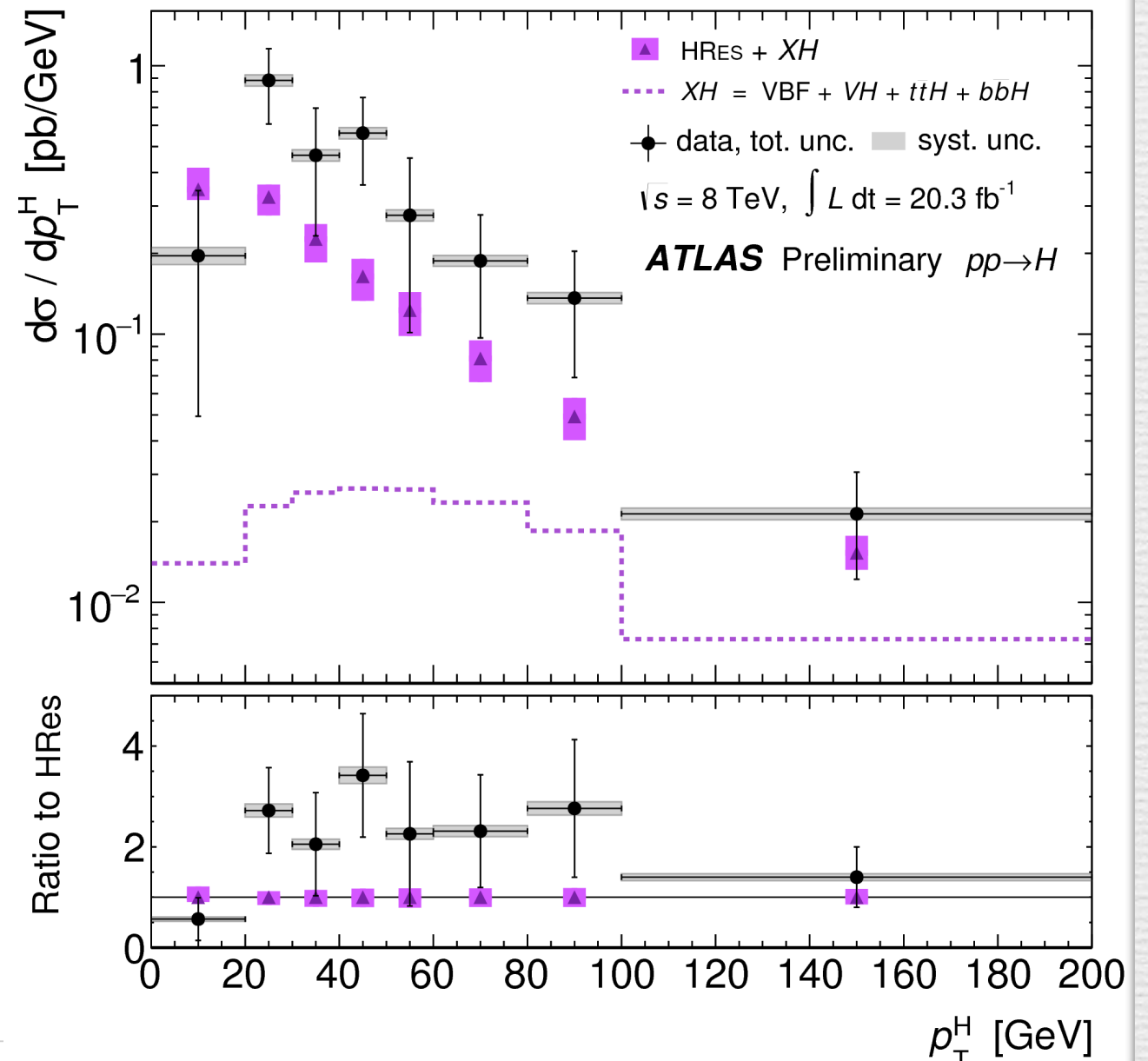
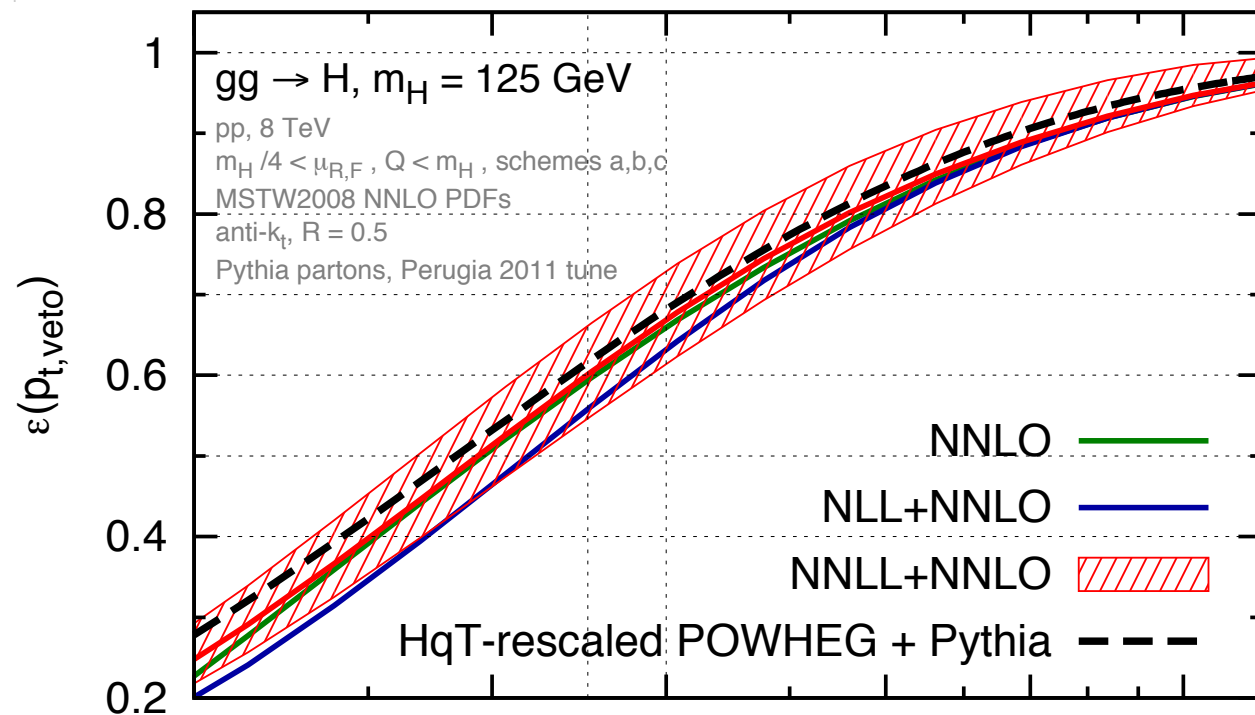
$H+j @ \text{NNLO}$

H+j @ NNLO

- Recently, H+j @ NNLO became available.

[Boughezal, Caola, Melnikov, Petriello, Schulze; Boughezal, Focke, Giele, Liu, Petriello; Chen, Gehrmann, Glover, Jaquier]

- Higgs-pT beyond NLO.
- Reduced uncertainties for jet-veto efficiencies.

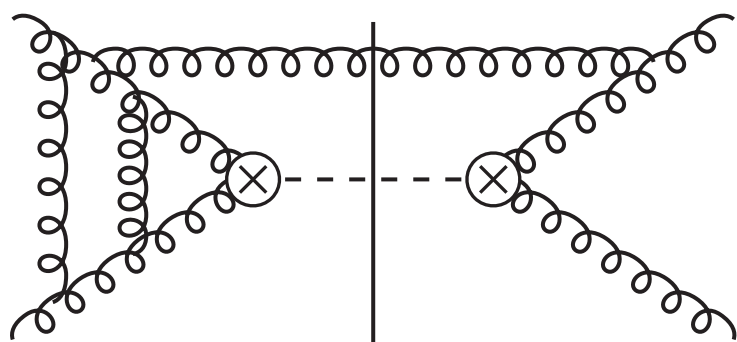


$H+j @ \text{NNLO}$

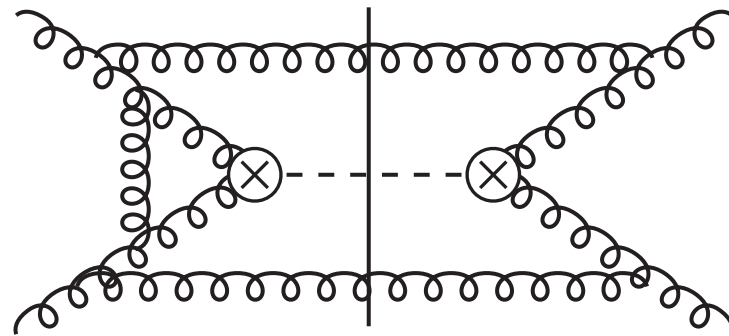
- Now: Predictions at NNLO accuracy for
 - ➔ (arbitrary) differential distributions
 - ➔ with (arbitrary) cuts on the final state (fiducial volume!)

$H+j$ @ NNLO

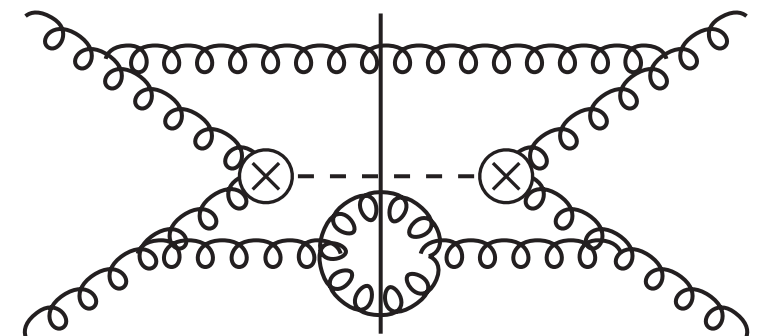
- Now: Predictions at NNLO accuracy for
 - ➔ (arbitrary) differential distributions
 - ➔ with (arbitrary) cuts on the final state (fiducial volume!)
- Only possible due to major advances in our understanding of how to cancel IR singularities at NNLO!



Double virtual



Real-virtual



Double-real

- ➔ Different contributions individually divergent.
- ➔ Divergences cancel in the sum.
- ➔ Different contributions live in different phase spaces.

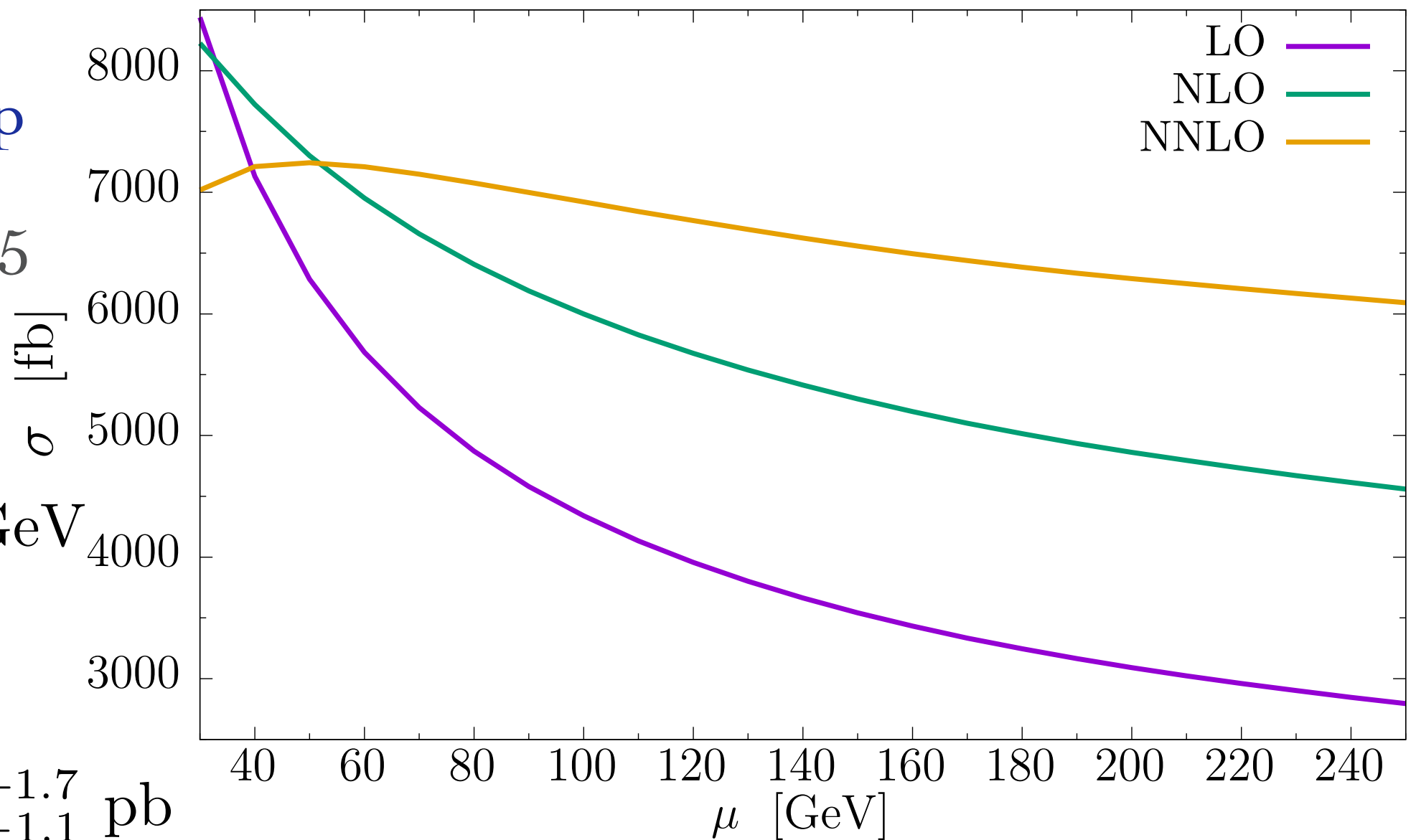
IR singularities

- **Basic Idea:** IR singularities of QCD amplitudes are known.
 - ➔ Use this to add and subtract counterterms that render all integrals finite.
 - ➔ Idea simple in principle, but very complicated in practise due to intricate nature of singularity structure at NNLO.
 - ➔ A lot of progress in the last few years!
- $H+j@NNLO$ was done using 3 different schemes to combine virtual and real corrections:
 - ➔ Antenna subtraction. [Kosower; Gehrmann, Gehrmann-de Ridder, Glover]
 - ➔ Stripper. [Czakon; Boughezal, Melnikov, Petriello; Czakon, Heymes]
 - ➔ N-jettiness subtraction. [Boughezal, Focke, Giele, Liu, Petriello; Gaunt, Stahlhofen, Tackmann, Walsh]

NNLO cross section for H+j

Sample setup

- anti-kt, R=0.5
- NNPF 2.3
- $p_{T,\text{cut}} > 30 \text{ GeV}$



$$\sigma_{\text{LO}} = 3.9^{+1.7}_{-1.1} \text{ pb}$$

$$\sigma_{\text{NLO}} = 5.6^{+1.3}_{-1.1} \text{ pb}$$

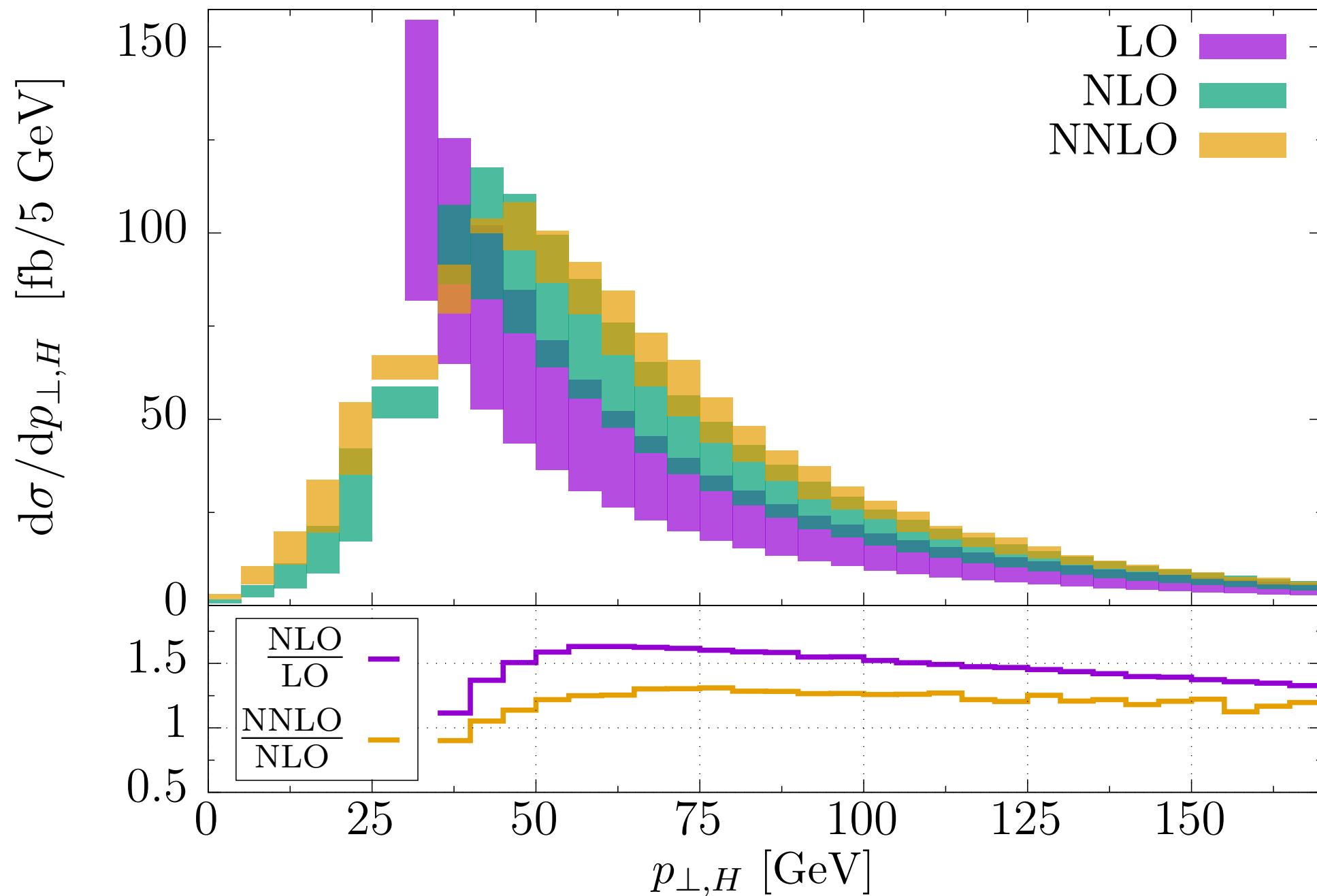
$$\sigma_{\text{NNLO}} = 6.7^{+0.5}_{-0.6} \text{ pb}$$

[Plot from Boughezal, Caola, Melnikov, Petriello, Schulze]

$$K_{\text{NNLO}}(\mu = m_H) \sim 20\%$$

$$K_{\text{NNLO}}(\mu = m_H/2) \sim 4\%$$

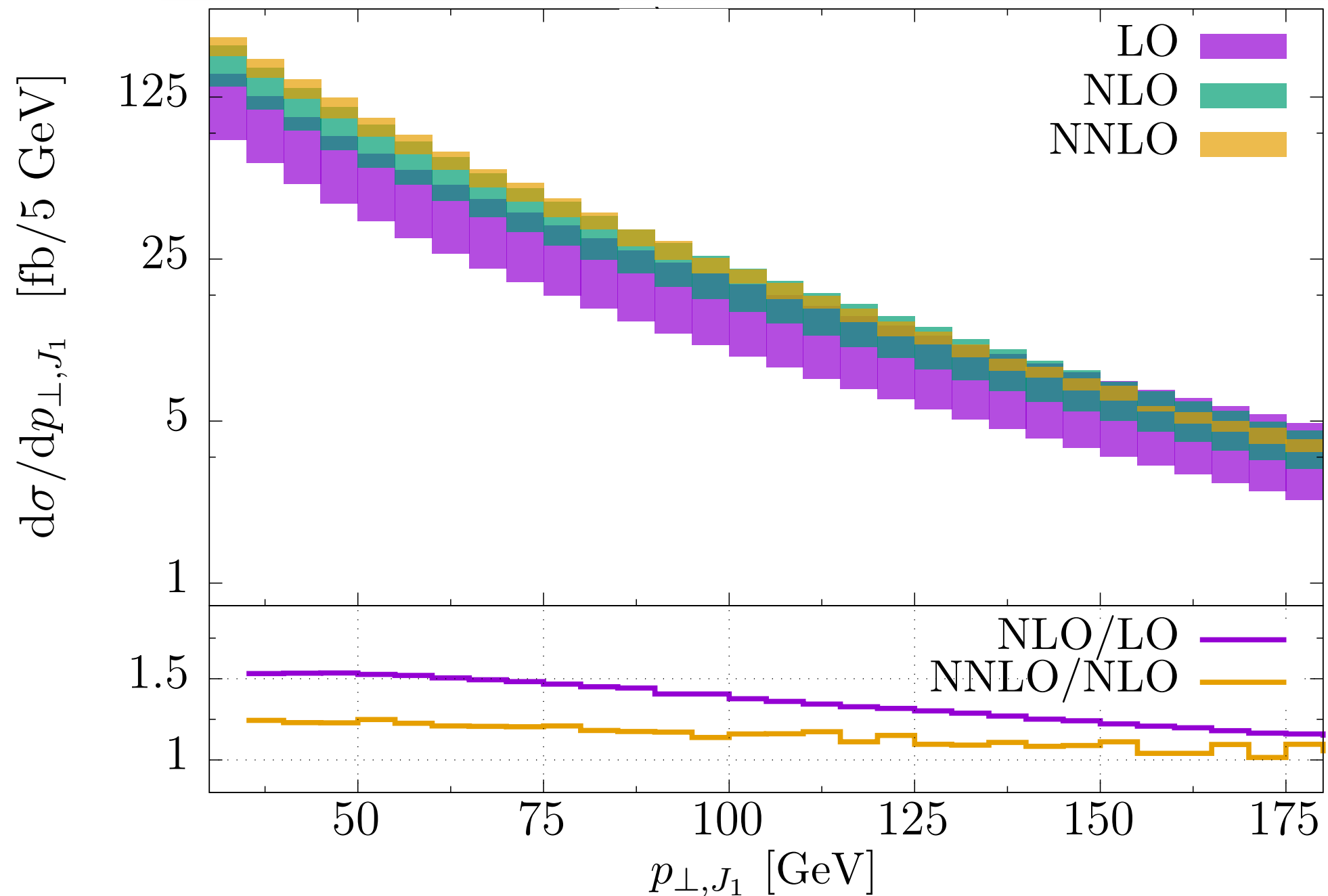
pT distributions



[Plot from Boughezal, Caola, Melnikov, Petriello, Schulze]

- Expect EFT to work within 2-3% up to $p_T \sim 150\text{GeV}$.

pT distributions



[Plot from Boughezal, Caola, Melnikov, Petriello, Schulze]

- Expect EFT to work within 2-3% up to $p_T \sim 150 \text{ GeV}$.

Jet veto efficiency

- $H + j$ @ NNLO gives very accurate predictions for the 1st jet bin.
- $H+j$ @ NNLO is at the same order in α_S as the inclusive cross section at N3LO.
 - ➔ Can combine the two and get very precise predictions for the 0 jet bin!

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- **Example:** Jet veto efficiency

$$\varepsilon(p_{T,\text{veto}}) = \frac{\Sigma_0(p_{T,\text{veto}}) + \Sigma_1(p_{T,\text{veto}}) + \Sigma_2(p_{T,\text{veto}}) + \Sigma_3(p_{T,\text{veto}}) + \dots}{\sigma_0 + \sigma_1 + \sigma_2 + \sigma_3 + \dots}$$

$$\Sigma_i(p_{T,\text{veto}}) = \sigma_i - \int_{p_{T,\text{veto}}}^{\infty} dp_T \frac{d\sigma_i}{dp_T}$$

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$$(H@N3LO) - (H+J@NNLO)$$

$$\varepsilon(p_{T,\text{veto}}) = \frac{\Sigma_0(p_{T,\text{veto}}) + \Sigma_1(p_{T,\text{veto}}) + \Sigma_2(p_{T,\text{veto}}) + \boxed{\Sigma_3(p_{T,\text{veto}})} + \dots}{\sigma_0 + \sigma_1 + \sigma_2 + \boxed{\sigma_3} + \dots}$$

$$H@N3LO$$

$$\Sigma_i(p_{T,\text{veto}}) = \sigma_i - \int_{p_{T,\text{veto}}}^{\infty} dp_T \frac{d\sigma_i}{dp_T}$$

Jet veto efficiency

0-jet bin

ord	$\sigma_{0\text{-jet}}^{\text{f.o.}} \text{ (JVE)}$	$\sigma_{0\text{-jet}}^{\text{f.o.}+\text{NNLL}} \text{ (JVE)}$	$\sigma_{0\text{-jet}}^{\text{f.o.}+\text{NNLL}} \text{ (scales)}$
NNLO	$26.2^{+4.0}_{-4.0} \text{ pb}$	$25.8^{+3.8}_{-3.8}$	$25.8^{+1.6}_{-1.6}$
N ³ LO	$27.2^{+2.7}_{-2.7} \text{ pb}$	$27.2^{+1.4}_{-1.4}$	$27.2^{+0.9}_{-0.9}$

≥ 1 -jet bin

ord	$\sigma_{\geq 1\text{-jet}}^{\text{f.o.}} \text{ (scales)}$	$\sigma_{\geq 1\text{-jet}}^{\text{f.o.}} \text{ (JVE)}$	$\sigma_{\geq 1\text{-jet}}^{\text{f.o.}+\text{NNLL}} \text{ (JVE)}$
NLO	$14.7^{+2.8}_{-2.8} \text{ pb}$	$14.7^{+3.4}_{-3.4}$	$15.1^{+2.7}_{-2.7}$
NNLO	$17.5^{+1.3}_{-1.3} \text{ pb}$	$17.5^{+2.6}_{-2.6}$	$17.5^{+1.1}_{-1.1}$

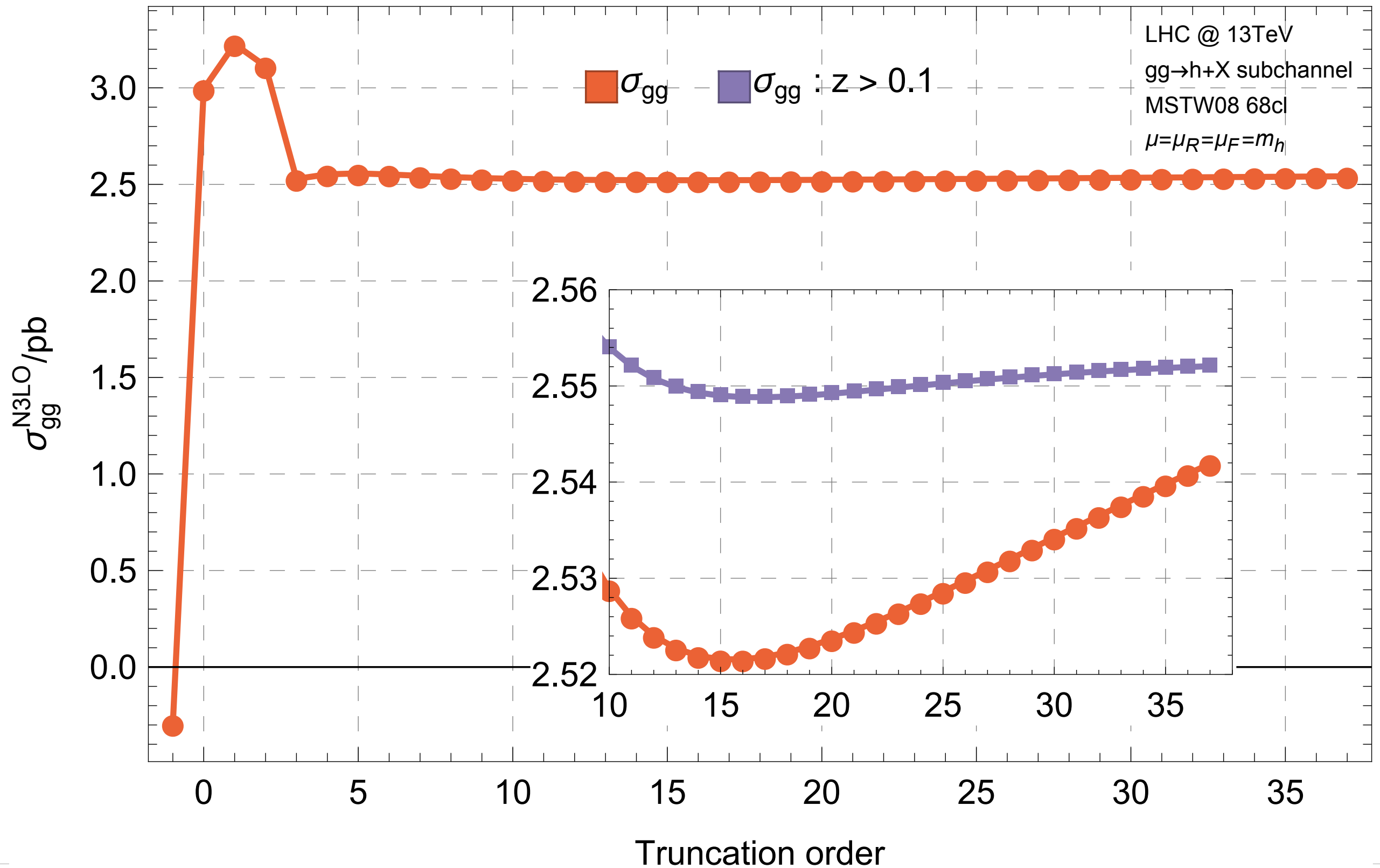
- Logs **completely under control**
(logR: see [Dasgupta, Dreyer, Salam, Soyez (2015)])
- No breakdown of f.o. perturbation theory for $p_T \sim 30 \text{ GeV}$
- Reliable error estimate from lower orders
- Logs **help in reducing uncertainties**
- **Significant decrease of pert. uncertainty**

Conclusion

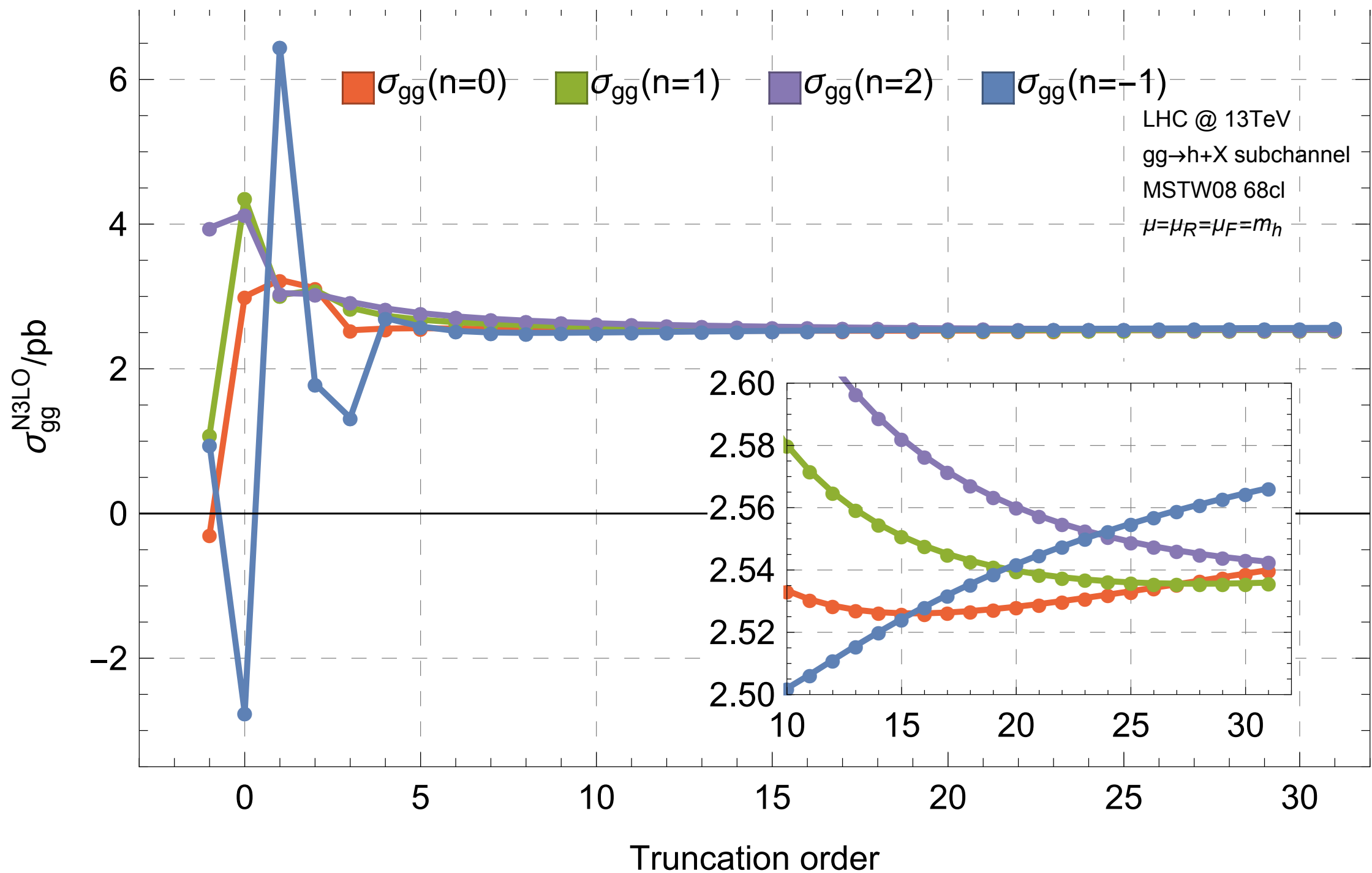
- In the last 6 months:
 - ➔ Fully differential VBF @ NNLO.
 - ➔ Fully differential $H+j$ @ NNLO.
 - ➔ Inclusive H @ N³LO.
- Drastic improvements of theoretical uncertainties!
 - ➔ Get theory predictions under control.
 - ➔ We are getting ready for precision Higgs physics!
- Our tools for QCD computations beyond NLO are getting more and more mature.

Backup slides

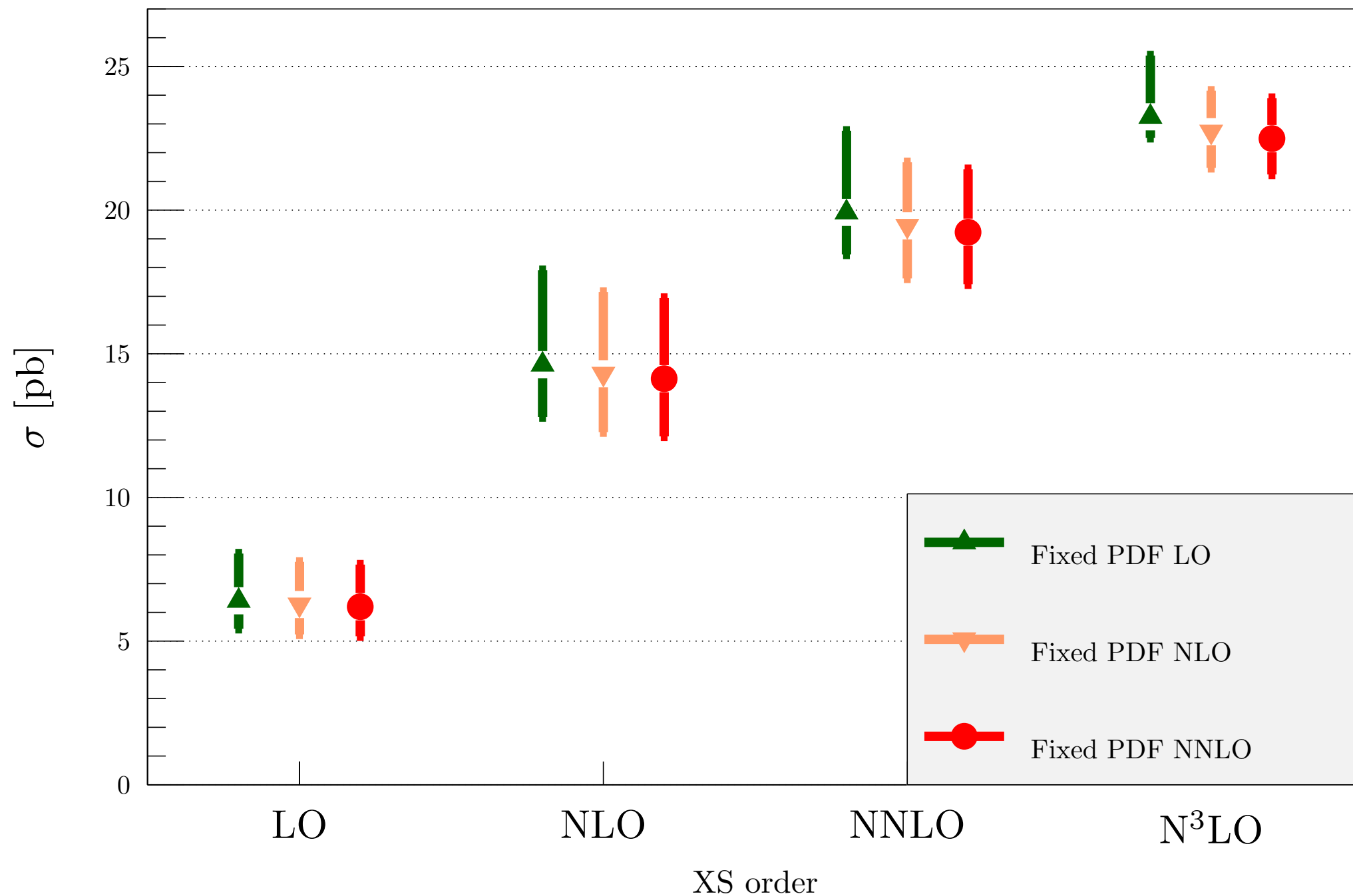
Threshold expansion



Threshold expansion

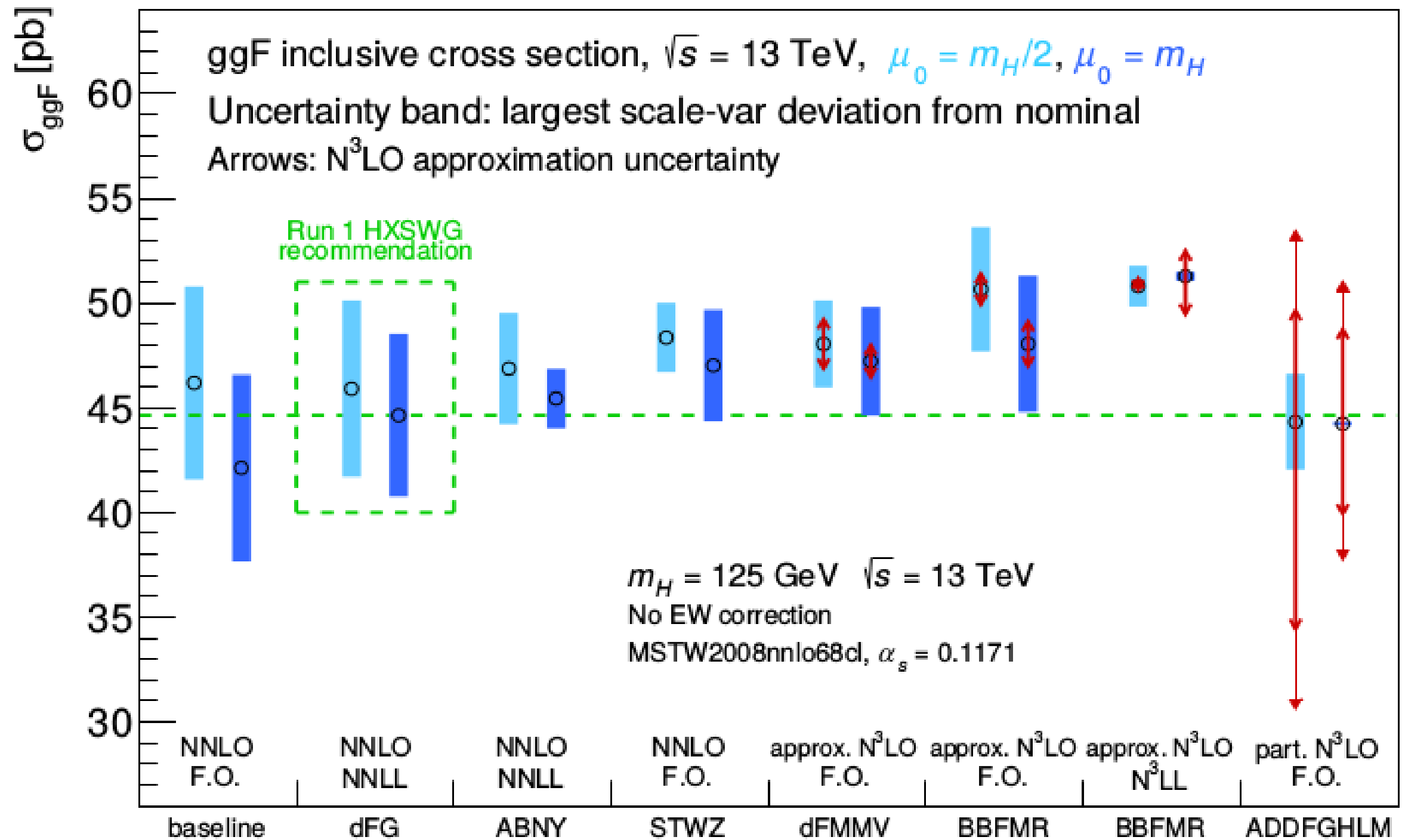


NNLO vs. N3LO PDFs?



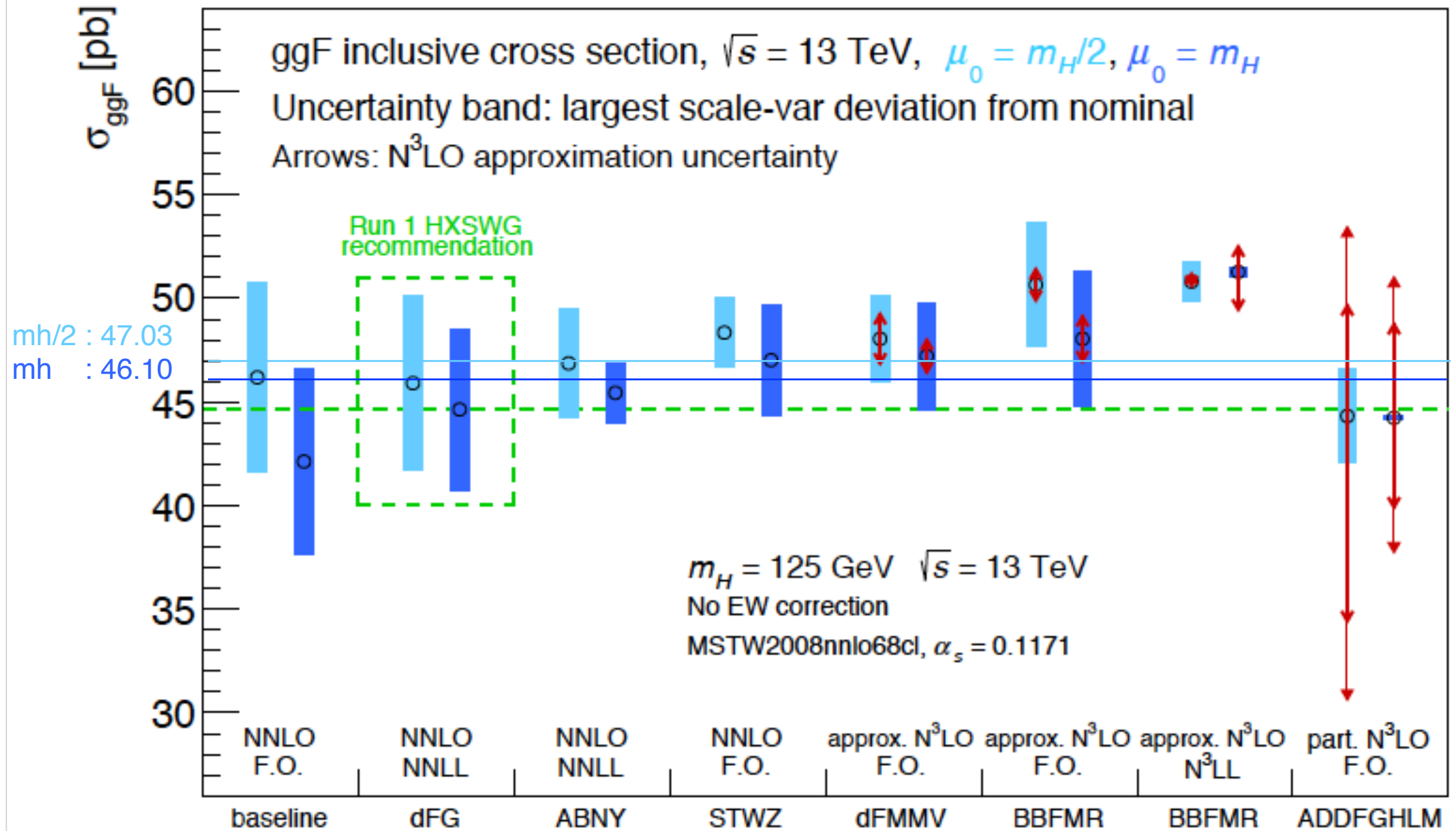
[Plot from Forte, Isgrò, Vita; N3LO is approximate]

Comparison to Approximate N3LO

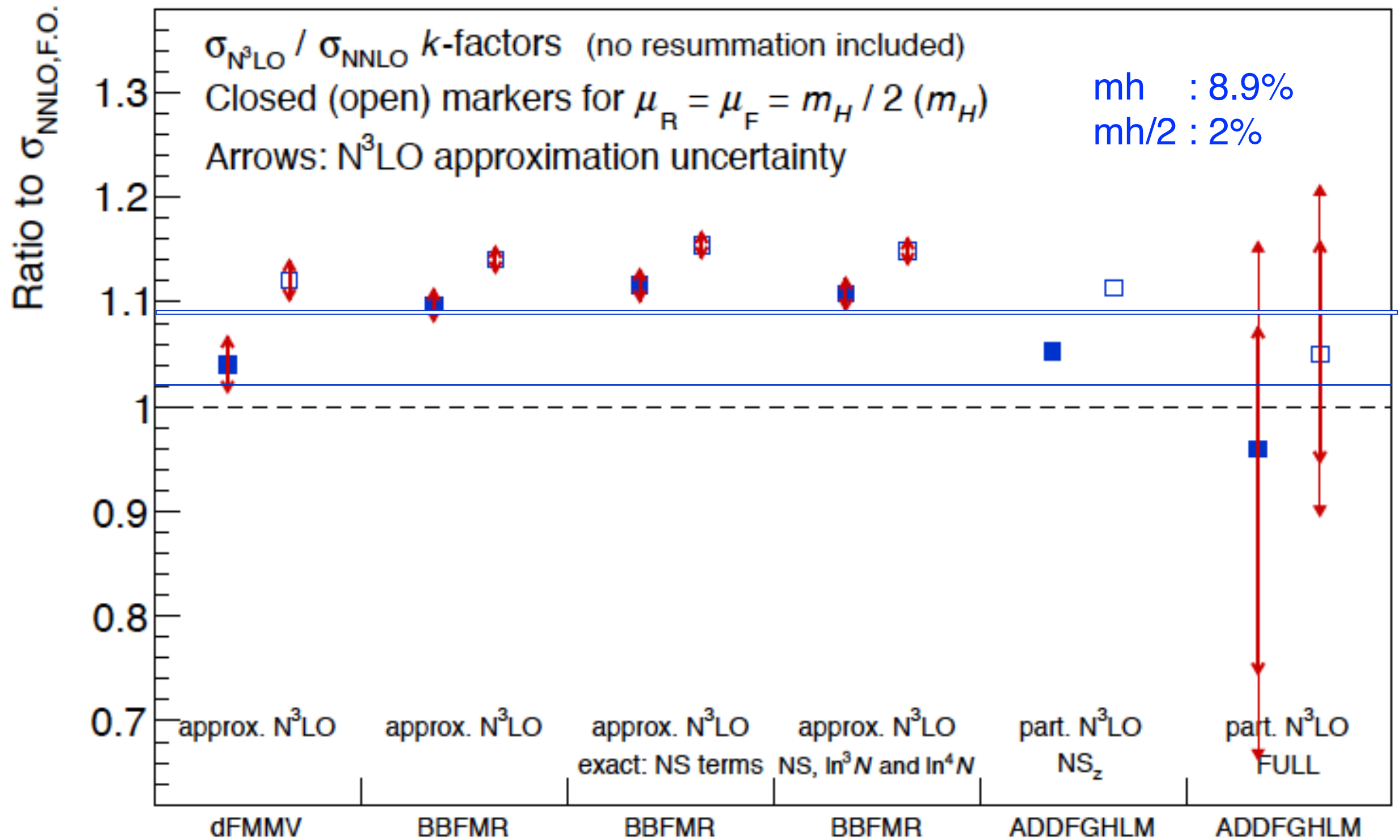


[Plot from HXSWG]

Comparison to Approximate N3LO



Comparison to Approximate N3LO



The soft-virtual contribution

$$\frac{\hat{\sigma}_{ij}(z)}{z} = \hat{\sigma}^{SV} \delta_{ig} \delta_{jg} + \sum_{N=0}^{\infty} \hat{\sigma}_{ij}^{(N)} (1-z)^N$$

$$a \delta(1-z) + \sum_{k=0}^5 b_k \left[\frac{\log^k(1-z)}{1-z} \right]_+$$

- Contributes to the gluon-channel only.
- Plus-distributions already known a decade ago.
 - ➔ Soft gluon emissions. [Moch, Vogt; Laenen, Magnea]
- delta-function contribution computed last year.
 - [Anastasiou, CD, Dulat, Furlan, Gehrmann, Herzog, Mistlberger; Li, von Manteuffel, Schabinger, Zhu]
 - ➔ Contains the complete three-loop corrections.
 - [Baikov, Chetyrkin, Smirnov², Steinhauser; Gehrmann, Glover, Huber, Ikizlerli, Studerus]

The regular contributions

$$\frac{\hat{\sigma}_{ij}(z)}{z} = \hat{\sigma}^{SV} \delta_{ig} \delta_{jg} + \sum_{N=0}^{\infty} \hat{\sigma}_{ij}^{(N)} (1-z)^N$$

$$\hat{\sigma}_{ij}^{(N)} = \sum_{k=0}^5 c_{ijk}^{(N)} \log^k(1-z)$$

- Describes subleading soft emissions.
- Single-emission contributions known exactly.
[Anastasiou, CD, Dulat, Herzog, Mistlberger; Kilgore; Gehrmann, Glover, Jaquier, Koukoutsakis, CD, Gehrmann, Jaquier; Dulat, Mistlberger]
- Double- and triple-emissions only known as an expansion around threshold.
[Anastasiou, CD, Dulat, Herzog, Mistlberger]
- ➔ Exact result for qq' channel was recently published.
[Anzai, Hasselhuhn, Hoff, Höschele, Kilgore, Steinhauser, Ueda]

Threshold resummation

- Soft gluon emissions exponentiate in Mellin space!

$$a \delta(1-z) + \sum_{k=0}^5 b_k \left[\frac{\log^k(1-z)}{1-z} \right]_+ \longrightarrow \tilde{a} + \sum_{k=1}^6 \tilde{b}_k \log^k N$$

$$\hat{\sigma}_{gg}^{resum} = g_0(\alpha_s) \exp \left[\frac{1}{\alpha_s^2} \sum_{k=1}^{\infty} \alpha_s^k g_k(\alpha_s \log N) \right] \quad [\text{Catani, Trentadue; Sterman}]$$

- Resummation functions g_i known up to N3LL (k=4).

[Moch, Vermaseren, Vogt; Bonvini, Marzani; Catani, Cieri, de Florian, Ferrara, Grazzini]

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- N3LL resummation needs 4-loop cusp anomalous dimension.

➔ Only known via Pade approximation, assuming Casimir scaling. [Moch, Vermaseren, Vogt]

➔ Casimir scaling assumption likely to fail at four loops.

➔ Numerical impact small!