

THEORY IMPROVEMENTS FOR THE HL-LHC

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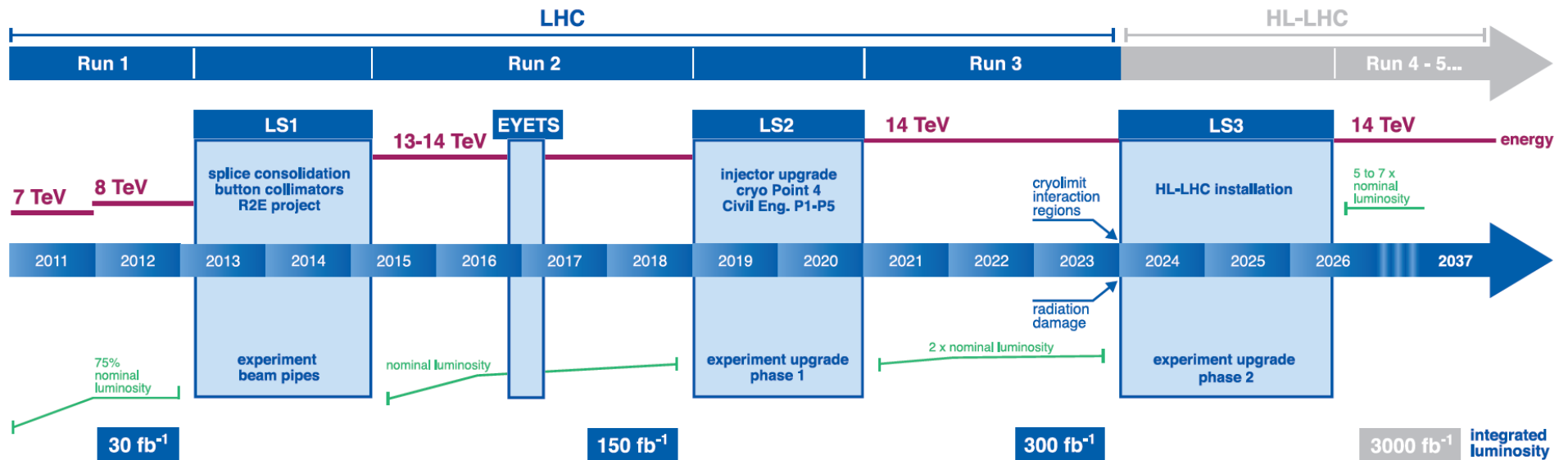


HIGGS HUNTING

ORSAY, AUGUST 1, 2015

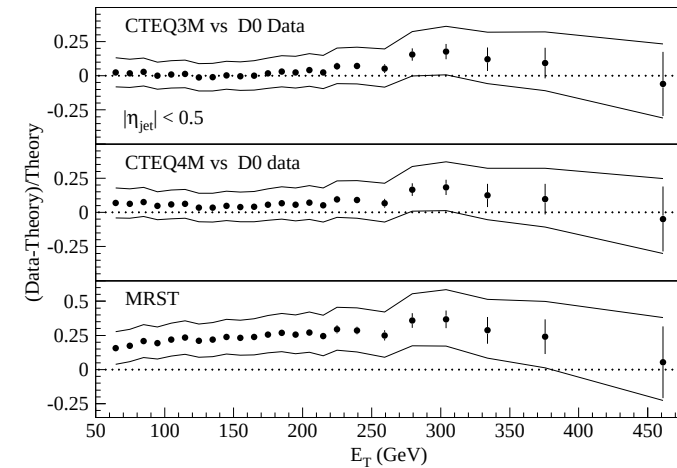
THE HL-LHC: 2025-2035

LHC / HL-LHC Plan



THE HL-LHC: 2025-2030

HOW WERE THINGS 15 YEARS AGO?



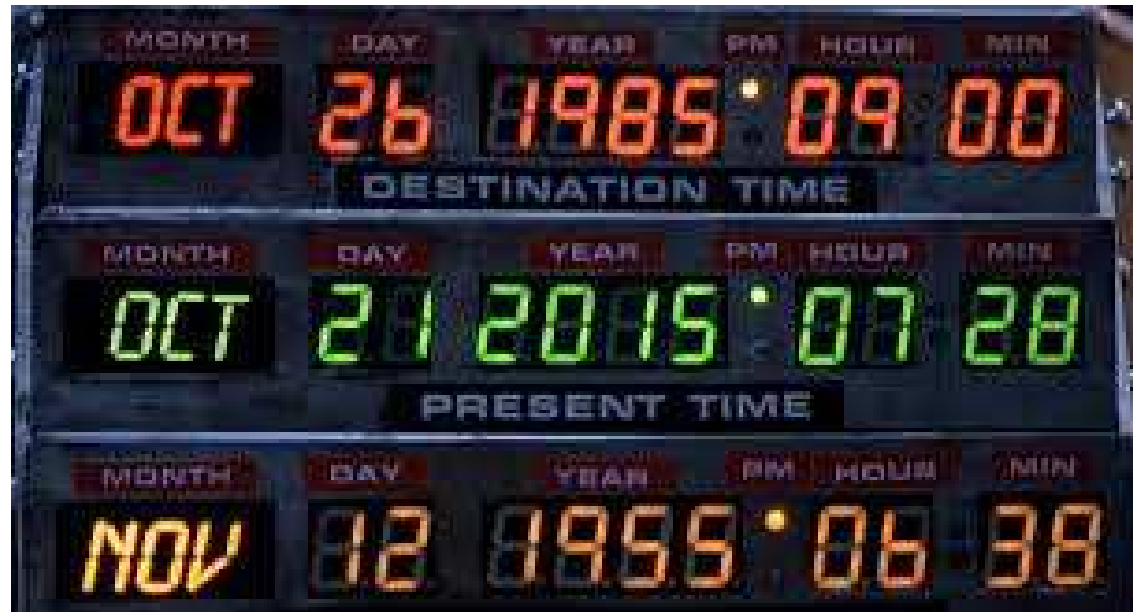
“The current agreement between theory and data is at the level of 30% over 8 orders of magnitude...”

(M. Mangano, EPS 1999)

(IT IS SUPPOSED TO BE A GOOD THING)

QCD IN THE YEAR 2000

- DGLAP SPLITTING FUNCTIONS KNOWN UP TO NLO
NNLO SPLITTING FUNCTIONS PUBLISHED IN 2002-2004
(Moch, Vermaseren, Vogt)
- ONLY INCLUSIVE TOTAL CROSS SECTIONS WITH COLORLESS FINAL STATES
(DRELL-YAN, HIGGS) KNOWN AT NNLO
RAPIDITY DISTRIBUTIONS COMPUTED TO NNLO IN 2003 (HIGGS),
2004 (DY) (Anastasiou, Melnikov, Petriello)
- PDF UNCERTAINTIES ESTIMATED BY COMPARING DIFFERENT PDF SETS
FIRST GLOBAL PDF SETS WITH UNCERTAINTIES PUBLISHED 2002
(CTEQ, MRST)
- EW CORRECTIONS MOSTLY AVAILABLE FOR LEPTON COLLIDER PROCESSES
EW CORRECTIONS TO HIGGS PRODUCTION 2004-2006
(Aglietti, Degrandi et al., Passarino et al., Dittmaier et al.)



QCD IN THE YEAR 2030?

- THE STRUCTURE OF THE QCD EXPANSION
 - RESUMMATION AND ITS USES
 - ANALYTIC STRUCTURE OF CROSS-SECTIONS
 - HIGHER ORDERS: UNCERTAINTY AND DIVERGENCE
- ELECTROWEAK CORRECTIONS
 - THE PHOTON PDF
 - MIXED QCD-EW CORRECTIONS
- PDFs (& α_s)
 - DATA: FROM LHC TO LHC-II (AND LHEC?)
 - PDFs BEYOND NNLO & THEORETICAL UNCERTAINTIES
 - MATCHING TO MONTE CARLO

HIGHER ORDERS AND RESUMMATION

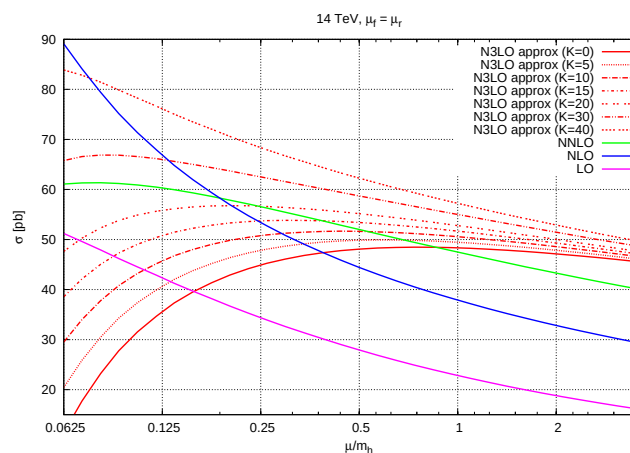


HIGHER ORDERS: WHAT DO WE KNOW? N³LO HIGGS IN GLUON FUSION

UNRESUMMED VS RESUMMED

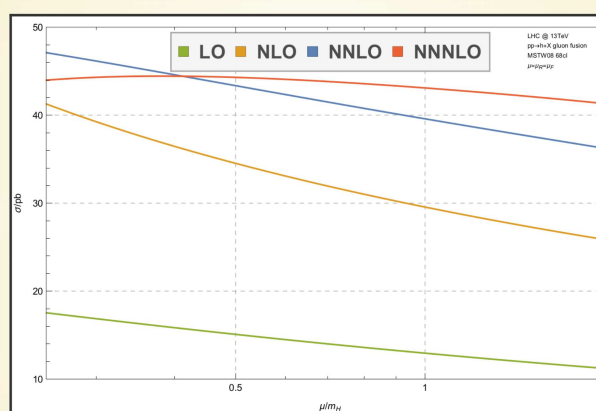
EXPANSION

SCALE DEP VS N³LO SIZE

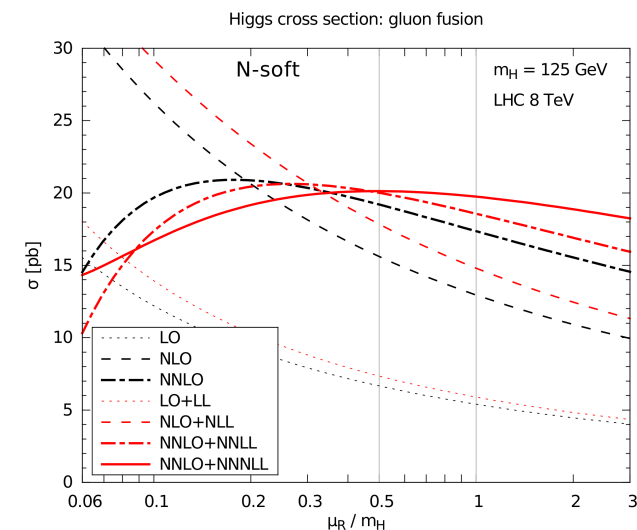


(Bühler, Lazopoulos, 2013)

EXACT N³LO



(Anastasiou et al, 2015)



(Bonvini, Marzani, 2014)

- THE PERTURBATIVE EXPANSION IS WELL-BEHAVED
- SOFT RESUMMATION ACCELERATES CONVERGENCE

Q: WHEN (AND WHY) IS SOFT RESUMMATION USEFUL?

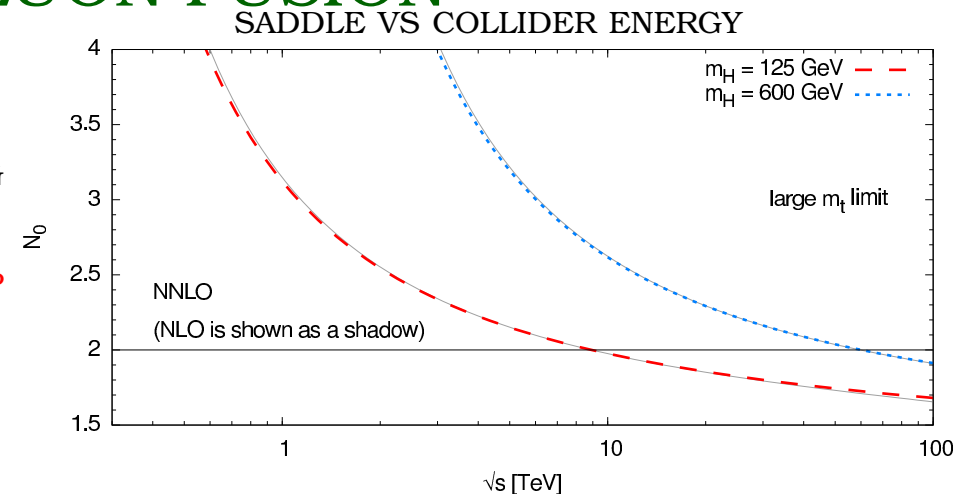
- PARTONIC CROSS SECTIONS $\hat{\sigma}(z, \alpha_s)$ ARE DISTRIBUTIONS, $0 \leq z = \frac{M_h^2}{s} \leq 1$, CONVOLUTED WITH LUMINOSITY TO GIVE HADRONIC CROSS-SECTION
- THEIR MELLIN TRANSFORMS $\hat{\sigma}(N, \alpha_s)$ ARE ORDINARY FUNCTIONS
- FOR GIVEN m_H^2 AND s ONLY ONE “SADDLE” N VALUE CONTRIBUTES;
POSITION OF SADDLE DETERMINED ESSENTIALLY BY PDF
 $\Rightarrow$ PDF TELLS YOU WHICH FRACTION OF HADRON MOMENTUM GOES INTO PARTONIC PROCESS

A: WHEN SADDLE N IS LARGE ENOUGH

- SOFT RESUMMATION INCLUDES TO ALL ORDERS $\alpha_s \ln(1 - z)$, $z = \frac{M_X^2}{s} \Leftrightarrow \alpha_s \ln N$
- SOFT $\Leftrightarrow z \rightarrow 1 \Leftrightarrow N \rightarrow \infty$.

HIGGS IN GLUON FUSION

- RESUMMATION CURRENTLY INCLUDED IN HXSWG RECOMMENATION: NNLO+NNLL
- BUT $N_{\text{saddle}} \sim 2$, $\Rightarrow$ RESUMMATION SUMS UP $\beta_0 \alpha_s \ln 2 \approx 0.04 \ll 1$
- WHAT IS “LARGE ENOUGH”??



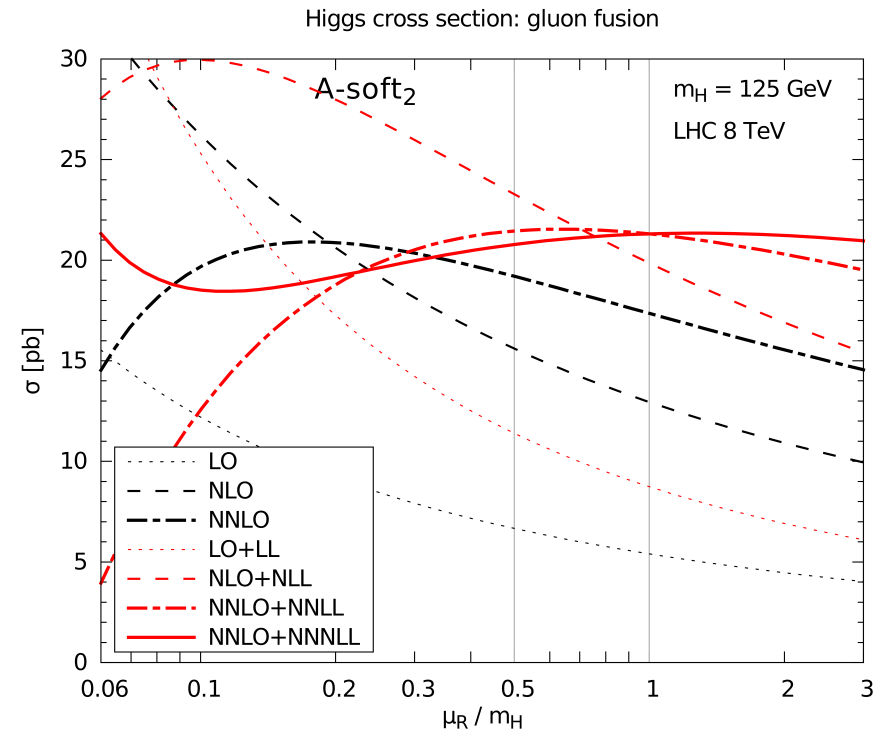
(Bonvini, SF, Ridolfi, 2012)

HIGGS IN GLUON FUSION: SHOULD RESUMMATION HELP? NO? YES!

$$\begin{aligned}
 M \left[\eta_{gg}^{(3)} \right] (N) &\simeq 36 \log^6 N & (\rightarrow 0.0013\%) \\
 &+ 170.679 \dots \log^5 N & (\rightarrow 0.0226\%) \\
 &+ 744.849 \dots \log^4 N & (\rightarrow 0.2570\%) \\
 &+ 1405.185 \dots \log^3 N & (\rightarrow 1.0707\%) \\
 &+ 2676.129 \dots \log^2 N & (\rightarrow 4.0200\%) \\
 &+ 1897.141 \dots \log N & (\rightarrow 5.1293\%) \\
 &+ 1783.692 \dots & (\rightarrow 8.0336\%) \\
 &+ 108 \frac{\log^5 N}{N} & (\rightarrow 0.0105\%) \\
 &+ 615.696 \dots \frac{\log^4 N}{N} & (\rightarrow 0.1418\%) \\
 &+ 2036.407 \dots \frac{\log^3 N}{N} & (\rightarrow 0.9718\%) \\
 &+ 3305.246 \dots \frac{\log^2 N}{N} & (\rightarrow 2.9487\%) \\
 &+ 3459.105 \dots \frac{\log N}{N} & (\rightarrow 5.2933\%) \\
 &+ 703.037 \dots \frac{1}{N} & (\rightarrow 1.7137\%).
 \end{aligned}$$

(Anastasiou et al, 2015)

- “LEADING” LOGS DO NOT PROVIDE THE DOMINANT CONTRIBUTION
- EXPANSION IN POWERS OF $1/N$ CONVERGES VERY SLOWLY



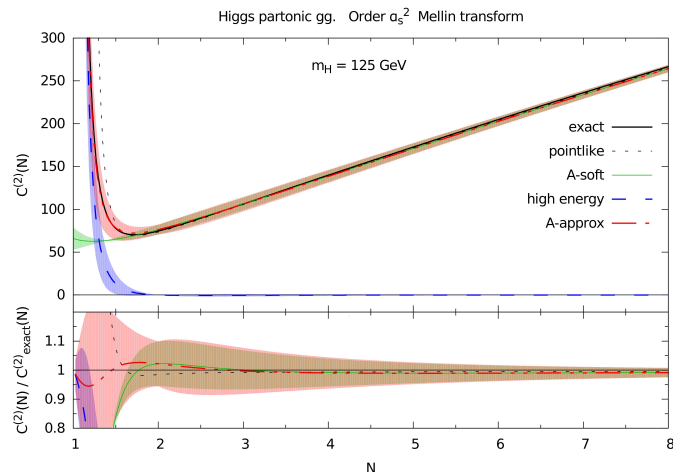
Bonvini and Marzani, 2014

- RESUMMED EXPANSION CONVERGES FASTER THAN UNRESUMMED ONE
- SCALE DEPENDENCE ALWAYS SMALLER AT RESUMMED LEVEL

ANALYTIC STRUCTURE IN N SPACE

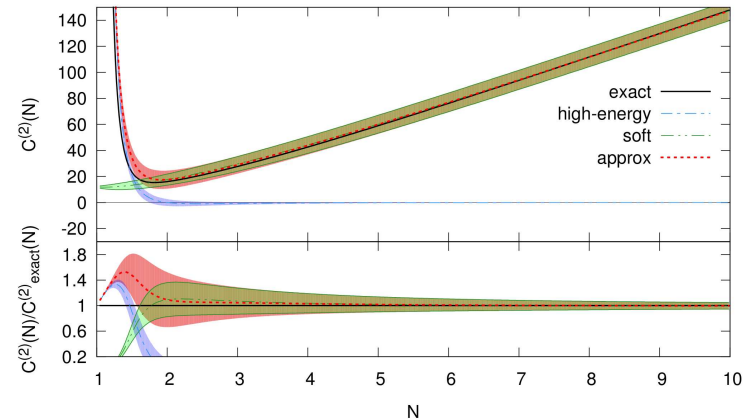
- MELLIN-SPACE PARTONIC CROSS SECTIONS ARE ANALYTIC FUNCTIONS OF N
- AS $N \rightarrow \infty$ THEY BEHAVE AS POWERS OF $\ln N$
- FINITE- N SINGULARITIES ARE ISOLATED POLES ON THE REAL AXIS FOR INTEGER $N \leq 1$
- RECONSTRUCT FROM SINGULARITIES, WITHOUT INTRODUCING SPURIOUS CUTS
EXAMPLE: LARGE- $N \ln N \Rightarrow \psi(N)$.

HIGGS IN GLUON FUSION



(Ball et al., 2013)

TOP PAIRS



(Muselli et al., 2015)

WHAT DO WE KNOW?

- **EIKONAL OR SOFT EXPANSION:** AS $N \rightarrow \infty$ (I.E. $z \rightarrow 1$)

$$\hat{\sigma}^{(3)}(N) = c_{06} \ln^6 N + c_{05} \ln^5 N + \dots + c_{02} \ln^2 N + c_{01} \ln N + c_{00} + \frac{1}{N} (c_{15} \ln^5 N + c_{14} \ln^4 N + \dots + c_{10}) + \frac{1}{N^2} (c_{24} \ln^4 N + \dots + c_{20}) + \dots$$

CURRENT KNOWLEDGE: EIKONAL, UP TO NNLL FOR MANY DIFFERENTIAL DISTRIBUTIONS

- **BFKL OR HIGH-ENERGY EXPANSION:** AS N DECREASES (I.E. $z \rightarrow 0$)

RIGHTMOST SINGULARITY AT $N = 1$:

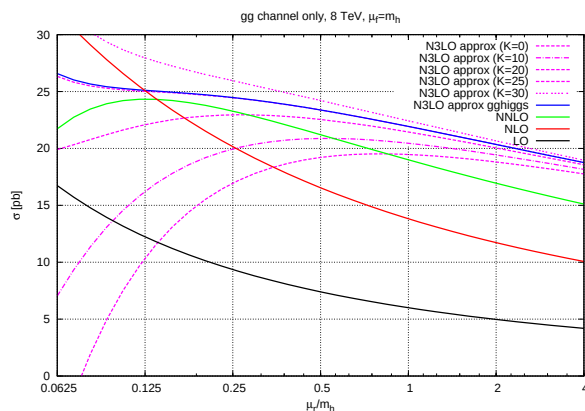
$$\hat{\sigma}^{(3)}(N) = \left(d_{13} \frac{1}{(N-1)^3} + d_{12} \frac{1}{(N-1)^2} + d_{11} \frac{1}{(N-1)} \right) + \left(d_{06} \frac{1}{N^6} + \dots + d_{01} \frac{1}{N} \right) + \dots$$

CURRENT KNOWLEDGE: LL, FOR SEVERAL TOTAL CROSS SECTIONS; ONLY GGHIGGS RAPDITY DISTRIBUTION

APPROXIMATE HIGHER ORDERS?

HIGGS GLUON FUSION: SCALE DEP

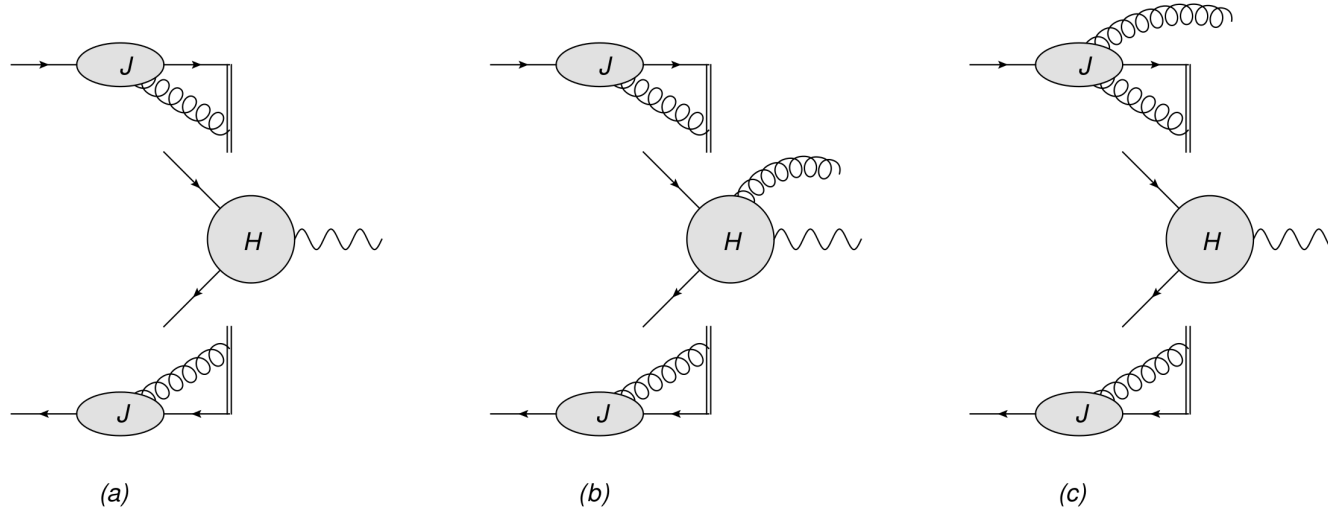
APPROX VS. GENERAL



(Bühler, Lazopoulos, 2013)

- AT NNLO METHOD WORKS WELL
- FOR GGHIGGS, FROM SCALE VARIATION RESULT SEEMS IN GOOD AGREEMENT WITH EXACT (AWAITING FOR FULL COMPARISON)
- N^4 LO COMPUTATIONS?
- **IMPROVING** BOTH THE SOFT AND THE BFKL **EXPANSION?**

IMPROVING THE SOFT EXPANSION BEYOND THE LEADING POWER



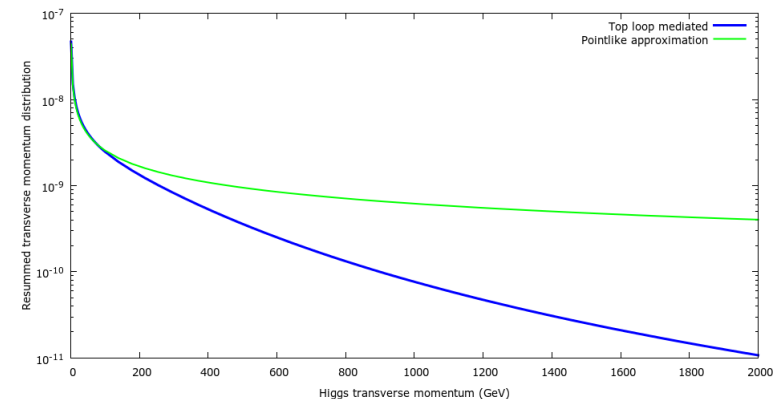
- **IMPRESSIVE PROGRESS** IN THE ORGANIZATION OF **RESUMMATION BEYOND THE EIKONAL LEVEL**: NEXT-TO-EIKONAL $\frac{\ln^k N}{N}$; NEXT-TO-NEXT-TO-EIKONAL $\frac{\ln^k N}{N^2}$ ETC (Laenen, Magnea, C.White et al, 2012-2105)
- **“METHOD OF REGIONS”**: SEPARATE INTEGRATION REGIONS OVER MOMENTA THROUGH APPROPRIATE SCALING
- **COMBINE WITH FACTORIZATION** TO CLASSIFY EMISSIONS FROM INTERNAL BLOBS, DRESSED WITH RADIATION FROM EXTERNAL LINES

IMPROVING HIGH ENERGY RESUMMATION BEYOND INCLUSIVE OBSERVABLES

- CURRENTLY KNOWN TO LL FOR TOTAL CROSS SECTIONS (HIGGS, DRELL-YAN, HEAVY QUARKS)
- RAPIDITY DISTRIBUTIONS KNOWN FOR HIGGS (Caola, SF, Marzani 2011)

- p_T DISTRIBUTIONS :
 - CAN STUDY SENSITIVITY TO m_t
 - BEHAVIOUR OF p_t RESUMMATION IN REGION $m_b < p_t \ll m_h$

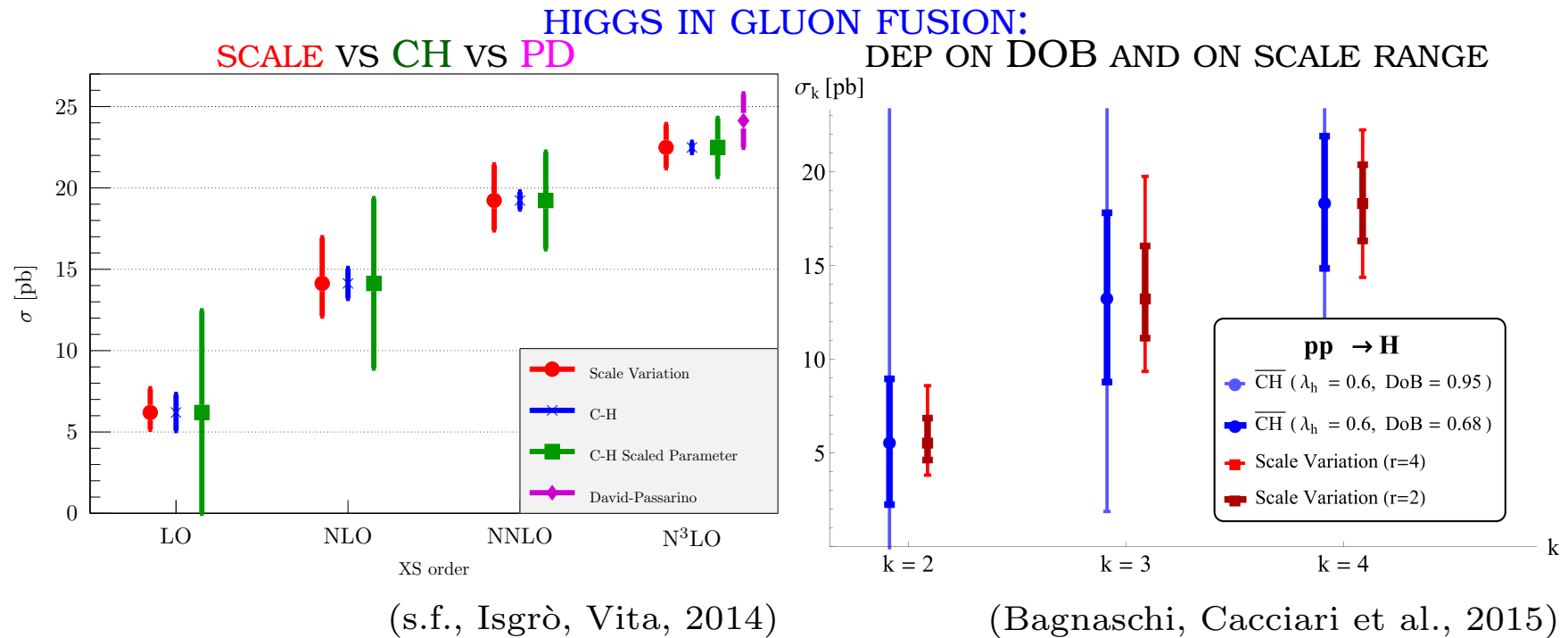
HIGGS IN GLUON FUSION: $LL \times p_t$
SPECTRUM
POINTLIKE VS. EXACT



FOR $p_t \gtrsim m_t$
EXACT DEVIATES FROM POINTLIKE
(Muselli, SF, Vita, in preparation)

THEORETICAL UNCERTAINTIES

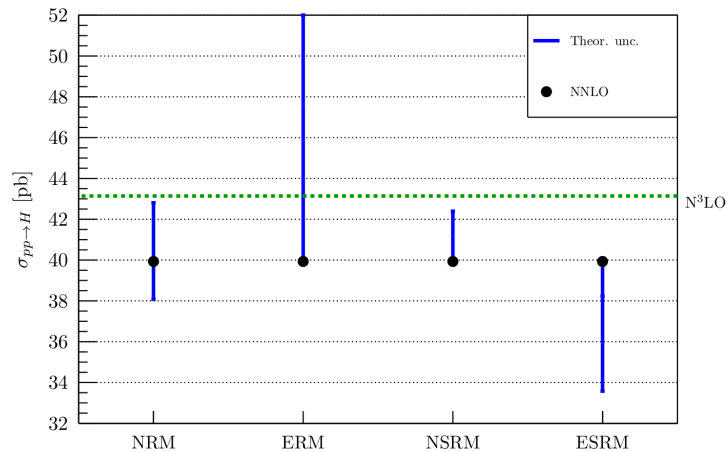
- **SCALE VARIATION** IS A **CRUDE** METHOD WHICH OFTEN FAILS
- CACCIARI-HOUDEAU: LOOK AT THE **BEHAVIOUR OF THE PERTURBATIVE EXPANSION**
- PASSARINO-DAVID: TRY TO **DEFINE SUM OF THE SERIES BASED ON SERIES ACCELERATION METHODS**
- **BETTER UNDERSTANDING AS MORE PROCESSES KNOWN**



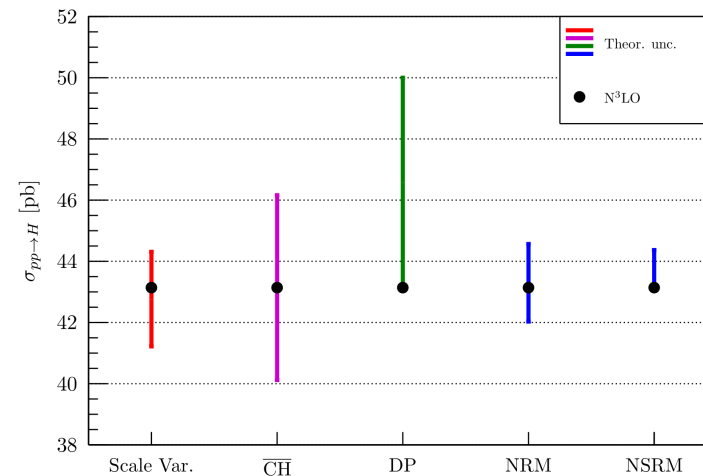
WHEN WILL IT START DIVERGING?

- THE QCD PERTURBATIVE EXPANSION IS ASYMPTOTIC:
DO WE HAVE TO WORRY ABOUT ITS DIVERGENCE?
- CAN WE USE THE ASYMPTOTIC SUM AS A
MEANS TO ESTIMATE THE THEORETICAL UNCERTAINTY?

HIGGS IN GLUON FUSION: NNLO UNCERTAINTY VS TRUE N³LO VARIOUS MODELS



NNNLO UNCERTAINTY



(s.f., Isgrò, in preparation)

- ORDER OF THE DIVERGENCE: LANDAU POLE $\Rightarrow n \sim 70$; RENORMALONS $\Rightarrow n \sim 70$
- ASSUME ASYMPTOTIC SUM GOVERNED BY RENORMALONS AFTER SUBTRACTING SOFT CONTRIBUTION (NSRM) $\Rightarrow$ REASONABLE RESULTS AT NNLO
- BECOMES MORE RELIABLE AS ORDER INCREASES

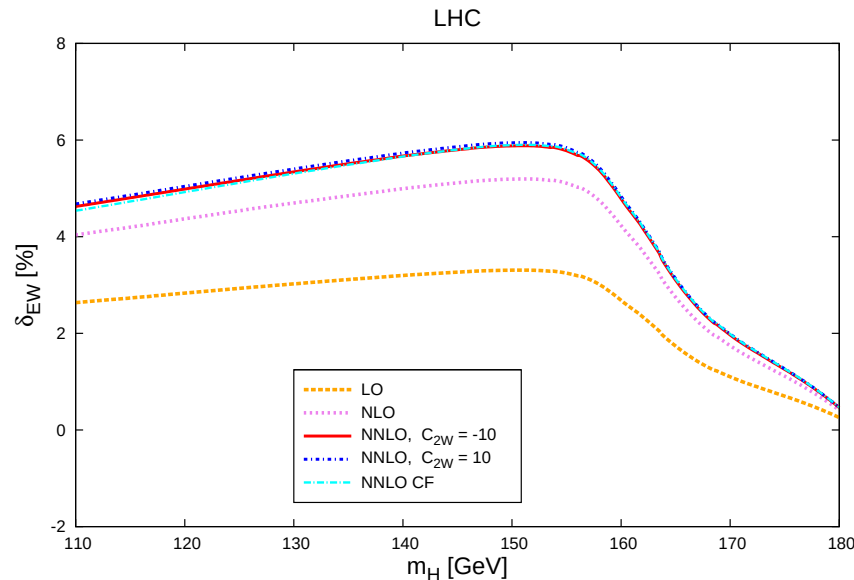
ELECTROWEAK CORRECTIONS



ELECTROWEAK CORRECTIONS

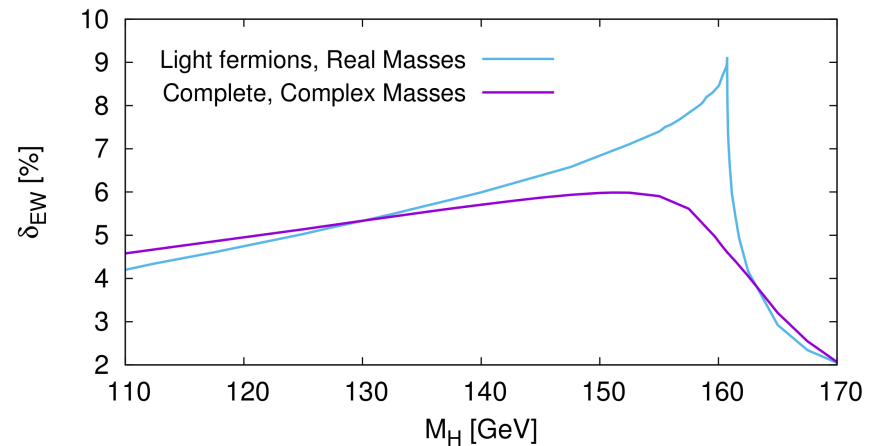
- EW CORRECTIONS HAVE A **SIZABLE IMPACT** ON THE GLUON FUSION XSECT 5%
- CURRENTLY **INCLUDED USING COMPLETE FACTORIZATION**:
IF ONLY APPLIED TO LO (PARTIAL FACTORIZATION), CHANGES UP TO 3%
- **WILSON COEFFICIENT COMPUTED TO $O(\lambda_{EW}\alpha_s)$ IN EFT**
(EXPANSION IN $\frac{m_H}{2m_t}, \frac{m_h}{m_W}$) (Anastasiou, Boughezal, Petriello 2008)
100% VIOLATION OF FACTORIZATION NOT ENOUGH TO PRODUCE SIGNIFICANT EFFECT
- FINITE m_t, m_W CORRECTIONS PROBABLY SMALL
- **CAN WE REALLY TRUST FACTORIZATION TO $O(\lambda_{EW}\alpha_s^2)$?**

DEPENDENCE OF FACTORIZATION



(Anastasiou, Boughezal, Petriello 2008)

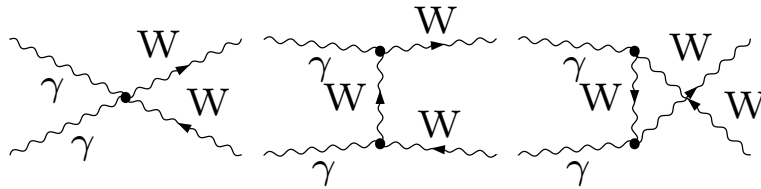
ACCURACY OF EFT



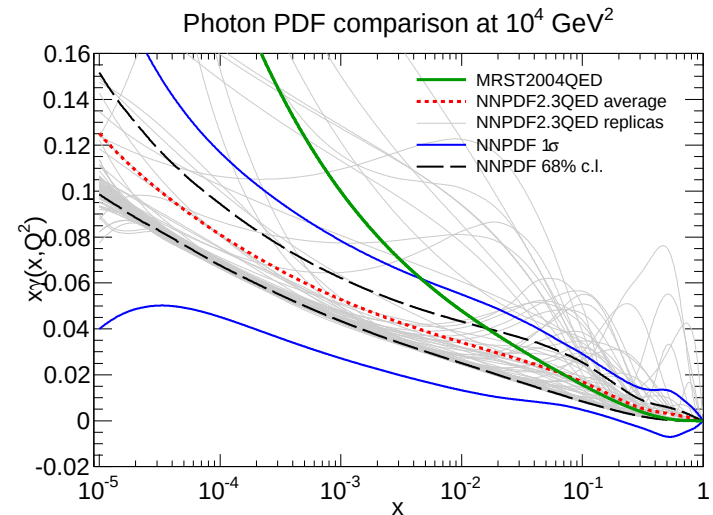
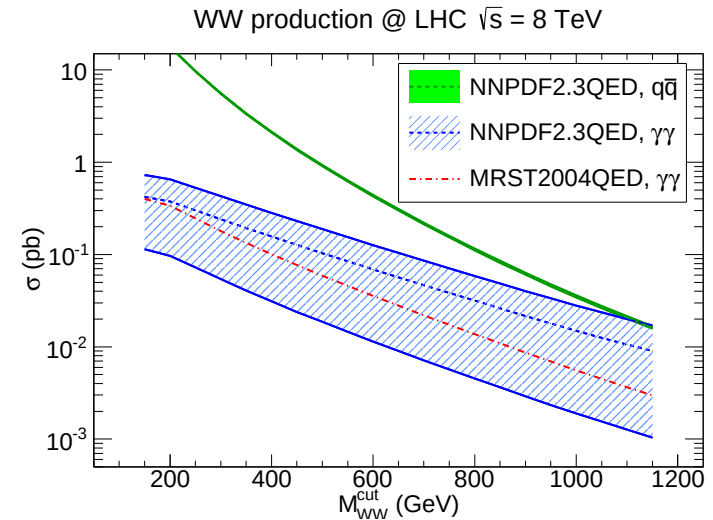
(Uccirati, Les Houches 2015)

ELECTROWEAK CORRECTIONS

QED PDFs



- PHOTON PDF CONTRIBUTION **CAN BE SIZABLE**:
EXAMPLE: W PAIR PRODUCTION
AT LARGE INVARIANT MASS
- CURRENTLY DETERMINED ONLY BY W
(ON- AND OFF-SHELL) PRODUCTION DATA,
AFFECTED BY VERY LARGE UNCERTAINTIES
- WILL BE **DETERMINED IN GLOBAL PDF FITS**
- **DIRECT PHOTON** PRODUCTION AT THE **LHeC**?



NNPDF2.3QED, Carrazza et al. 2013

PDFs (AND α_s)



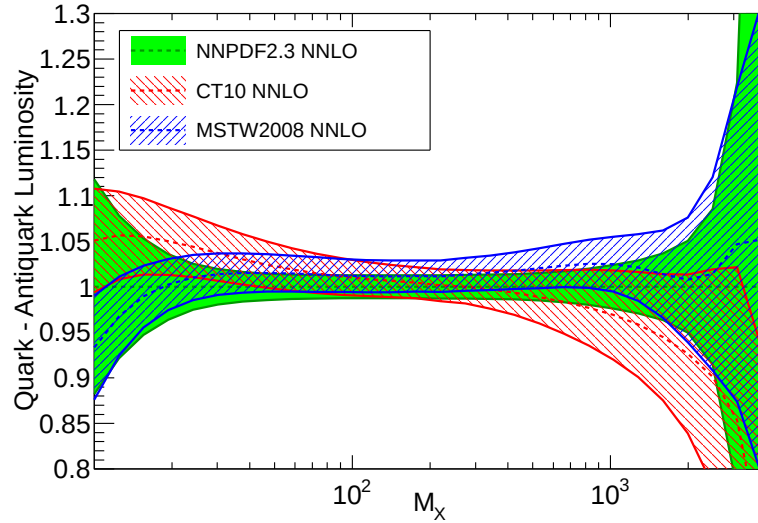
PDFs: RECENT PROGRESS

PARTON LUMINOSITIES

QUARK-QUARK

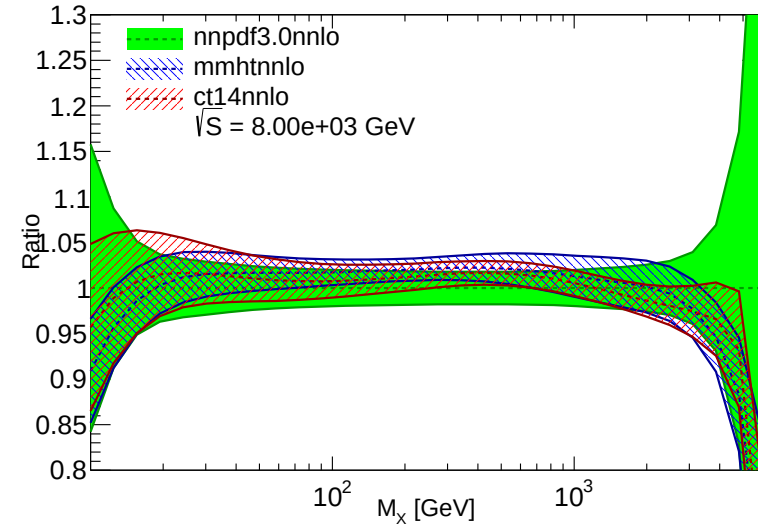
2012

LHC 8 TeV - Ratio to NNPDF2.3 NNLO - $\alpha_s = 0.118$



2015

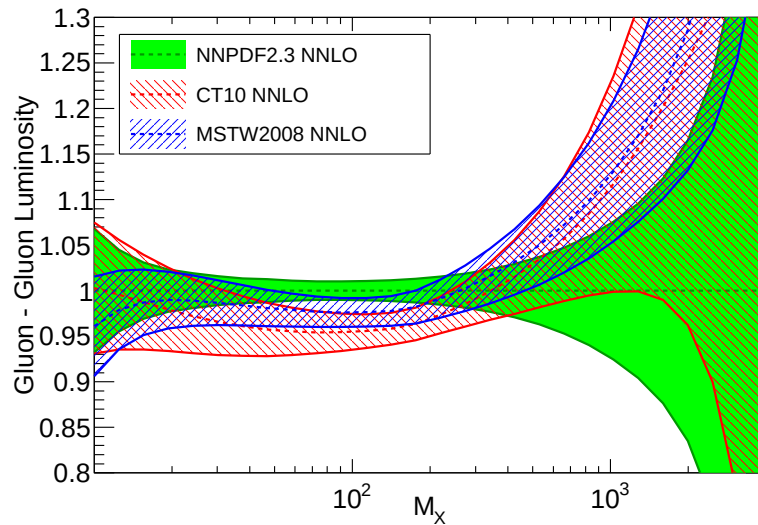
Quark-Quark, luminosity



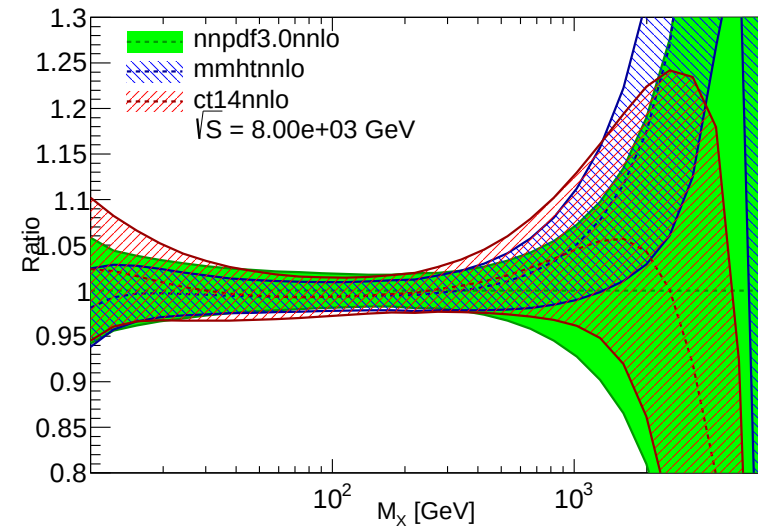
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GLUON-GLUON

LHC 8 TeV - Ratio to NNPDF2.3 NNLO - $\alpha_s = 0.118$



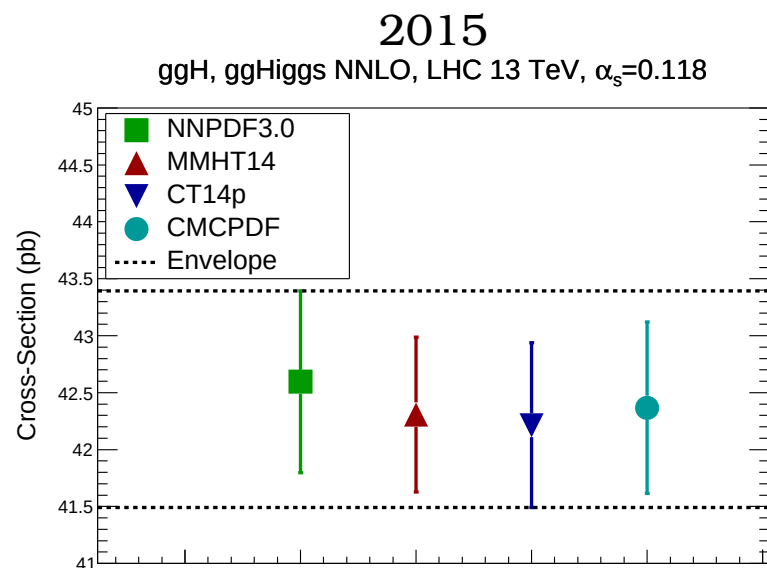
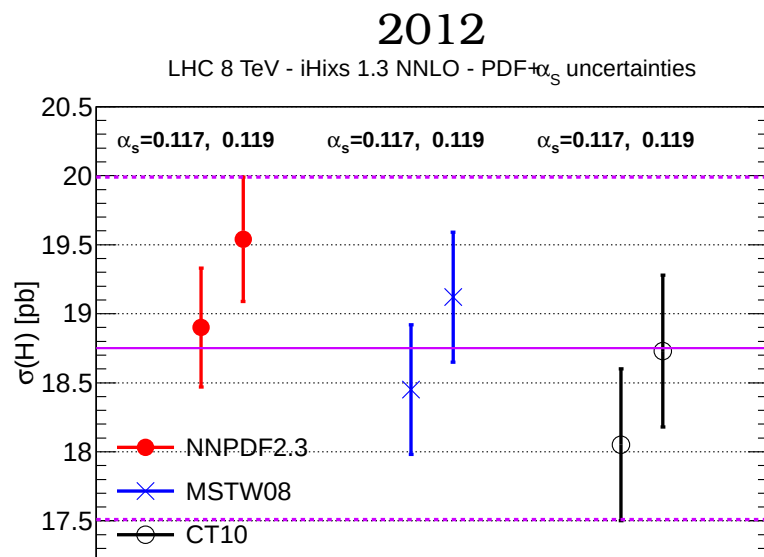
Gluon-Gluon, luminosity



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PDFs: RECENT PROGRESS

HIGGS IN GLUON FUSION

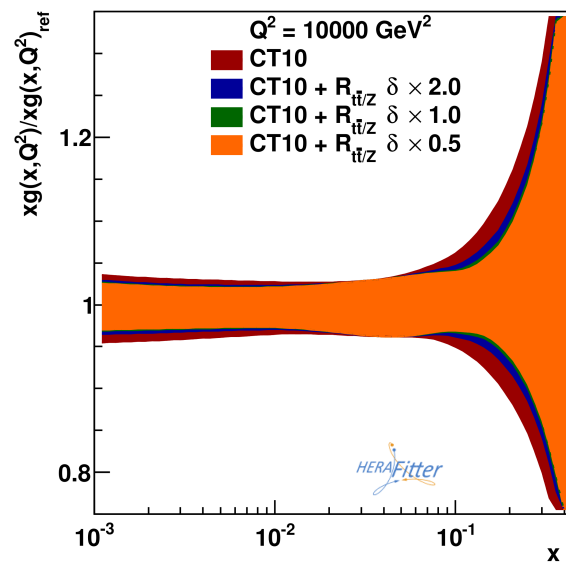
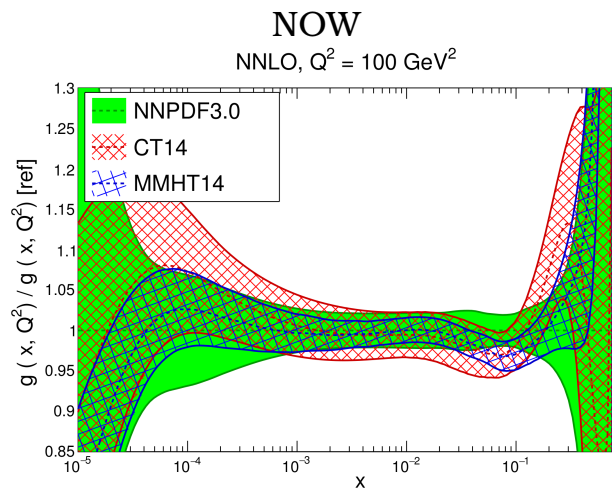


- **MAJOR UPGRADES** FROM ALL GLOBAL FITTING GROUPS:
NNPDF2.3→3.0 (10/2015); MSTW08 → MMHT14 (12/2015); CT10 ⇒ CT14 (06/2015)
- **METHODOLOGICAL IMPROVEMENTS:** CLOSURE TESTS (NNPDF); EXTENDED PARAMETRIZATIONS (CT, MMHT)
- LHC-I DATA INCLUDED
- PDF **UNCERTAINTY ON HIGGS PRODUCTION DOWN** TO ABOUT 2%
ENVELOPE NO LONGER NECESSARY

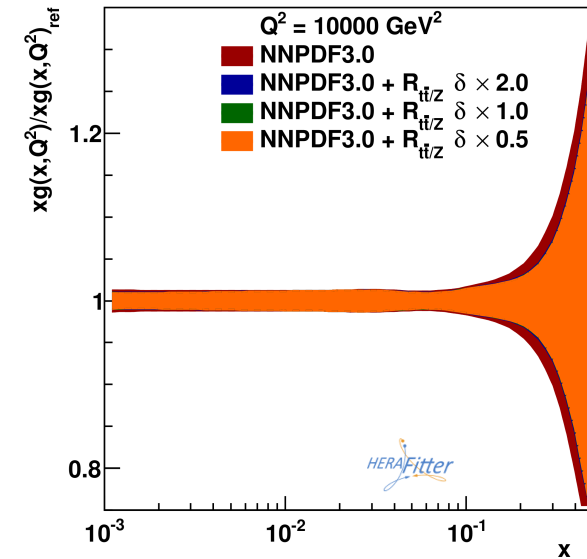
PDFS AT LHC RUN II

- DATA AT HIGHER CM ENERGY & INFO ON CORRELATION TO LOW ENERGY
→ EXTENDED KINEMATIC COVERAGE & REDUCED SYSTEMATICS
- POSSIBLY REDUCED STAT. UNCERTAINTIES
- EXPECT REDUCTION IN MODEL DEPENDENCE
- MODERATE REDUCTION IN UNCERTAINTY
- FURTHER IMPROVEMENTS DUE TO THEORY
EXAMPLE: GLUON FROM NNLO Z p_T DISTRIBUTION

THE GLUON CTEQ AFTER RUN II



NNPDF AFTER RUN II

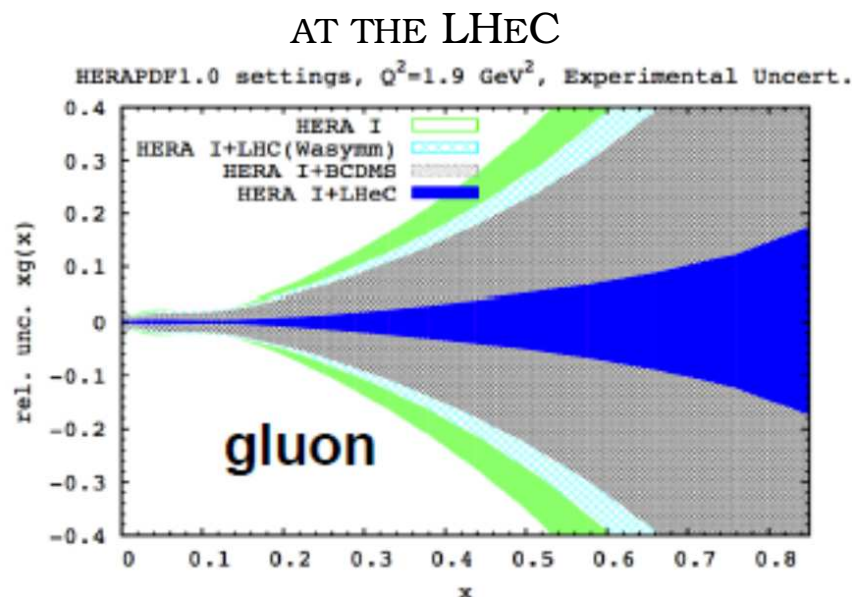
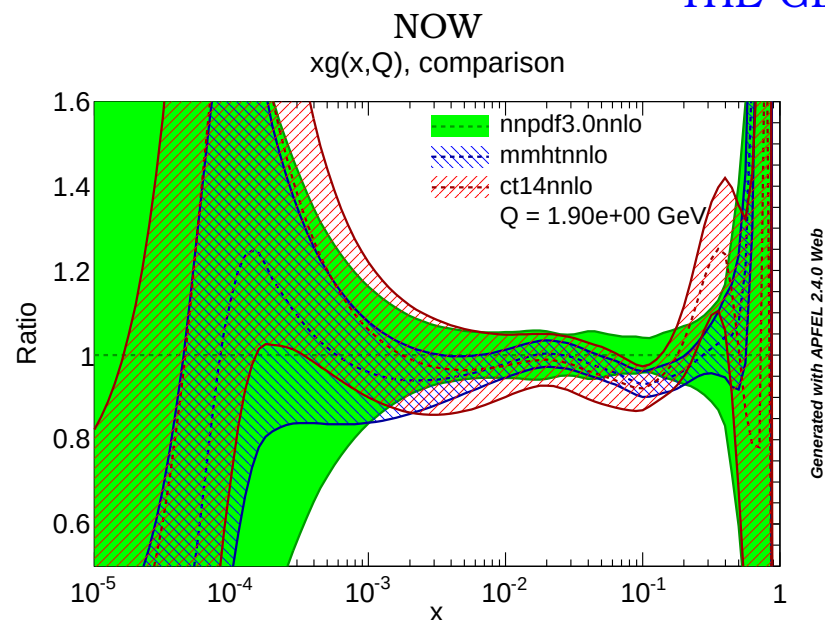


(PDF4LHC: 1507.00556)

PDFS AT THE LHeC?

- ESPECIALLY IF CORRELATION INFO AVAILABLE, **LHC RUN II** CAN LIKELY
 - **REMOVE SYSTEMATIC BIAS & DISCREPANCIES**
 - **BRING DOWN PDF UNCERTAINTIES TO $\sim 4\%$ LEVEL**
- **VERY DIFFICULT TO REDUCE UNCERTAINTIES FURTHER** AT A HADRON COLLIDER
- **HIGH-ENERGY DEEP-INELASTIC DATA** COULD LEAD TO PDFS AT 1% LEVEL

THE GLUON AGAIN



(A. Cooper-Sarkar & Voica Radescu, 2015)

THE VALUE OF α_s

HIGGS IN GLUON FUSION $O(\alpha_s^2)$, LARGE NLO CORRNS: $\Delta\sigma \sim 2.5\Delta\alpha_s$

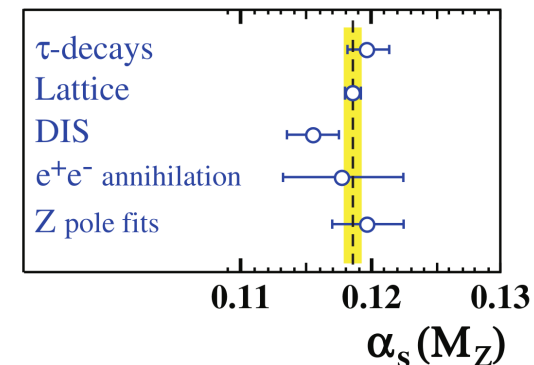
PDG VALUE (AUGUST 2014): $\alpha_s(M_Z) = 0.1185 \pm 0.0006$

- LATTICE UNCERTAINTY CURRENTLY ESTIMATED BY FLAG (arXiv:1310.8555) TO BE TWICE THE PDG VALUE (± 0.0012)
(IF PDG WERE TO ADOPT THIS, COMBINED UNCERTAINTY LIKELY TO DOUBLE)
- IT IS AN AN AVERAGE OF AVERAGES
- SOME SUB-AVERAGES (E.G. DIS) INCLUDE MUTUALLY INCONSISTENT VALUES
- SOME SUB-AVERAGES (E.G. τ OR JETS) INCLUDE DETERMINATIONS WHICH DIFFER FROM EACH OTHER BY EVEN FOUR-FIVE σ
- AVERAGING THE TWO MOST RELIABLE VALUES (GLOBAL EW FIT & τ , BOTH N³LO, NO DEP. ON HADRON STRUCTURE) GIVES

$$\alpha_s = 0.1196 \pm 0.0010$$

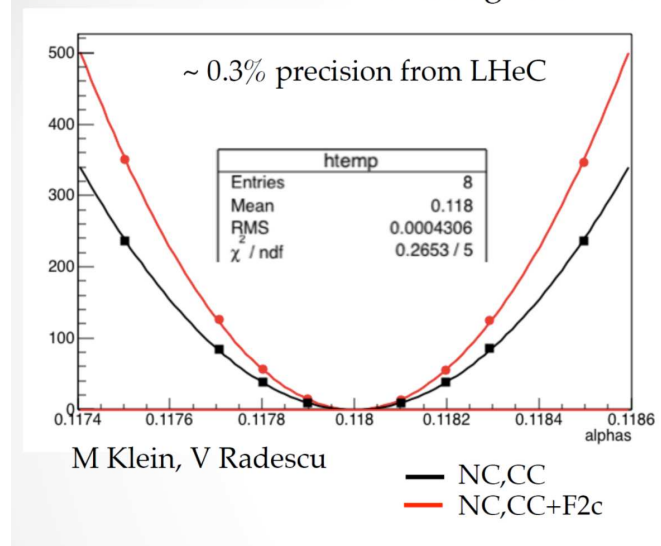
- LITTLE PROGRESS FOR MANY YEARS: PDG 1998-2006 $\Delta\alpha_s(M_Z) = 0.002$; PDG 2010-2014 $\Delta\alpha_s(M_Z) = 0.0006 \div 0.0008$ (CHANGE OF AUTHOR)
- UNLIKELY TO CHANGE - SHOULD DETERMINE α_s FROM HIGGS IN GLUON FUSION?
- COULD BE DETERMINED ACCURATELY AT THE LHeC

α_s DETERMINATIONS IN PDG



α_s AT THE LHeC

combined fit to PDFs+ α_s using LHeC data



THEORETICAL UNCERTAINTIES ON PDFs:

- PDFs ARE DETERMINED BY COMPARING TO DATA THEORY AT SOME FINITE ORDER
- AFFECTED BY THEORETICAL UNCERTAINTY JUST LIKE HARD CROSS-SECTIONS
- NOT INCLUDED IN CURRENT “PDF UNCERTAINTY”
(ACCOUNTS ONLY DATA & METHODOLOGY)

CAN WE ESTIMATE THEM?

- SCALE VARIATION DIFFICULT:
CORRELATED BETWEEN PROCESSES? HOW DOES IT CORRELATE WITH PROCESSES
IN WHICH PDFs ARE USED?
- AT NLO: WE KNOW THE SHIFT TO NNLO
- AT NNLO: LOOK AT THE BEHAVIOUR OF THE PERTURBATIVE EXPANSION
(CACCIARI-HOUDEAU)

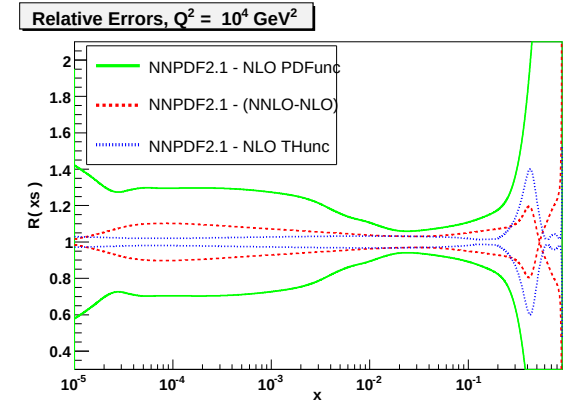
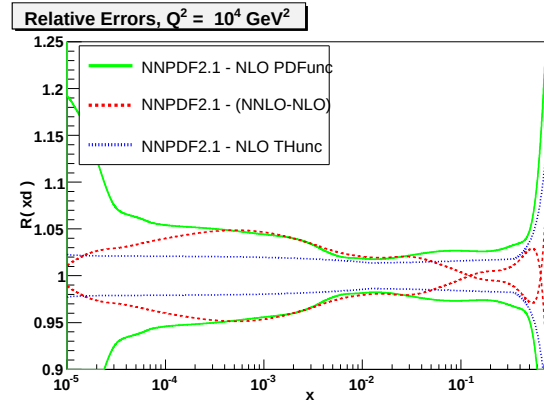
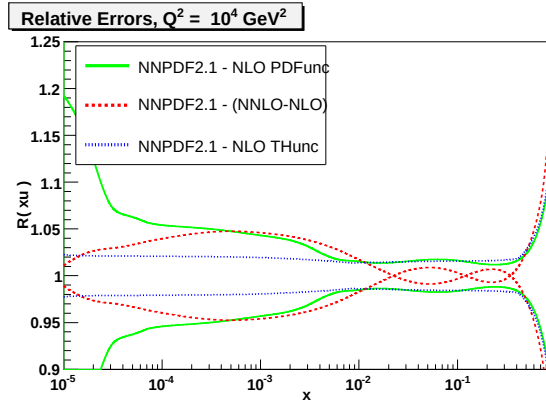
THEORETICAL UNCERTAINTIES

NLO PDF UNC. VS NLO-NNLO SHIFT VS NLO CACCIARI-HOUDEAU (NNPDF2.1)

UP

DOWN

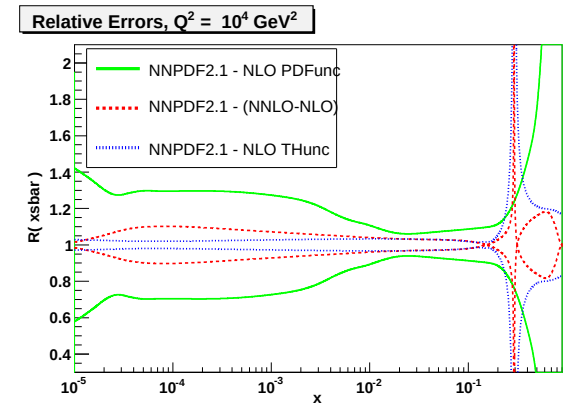
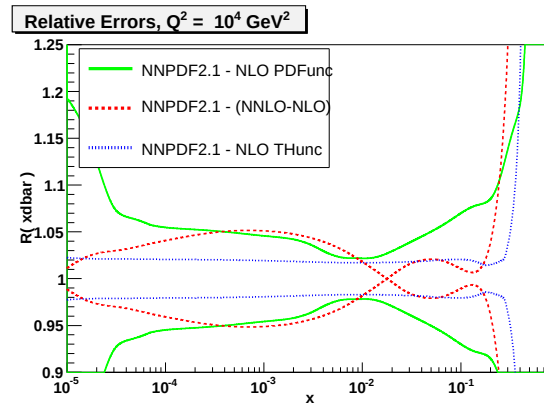
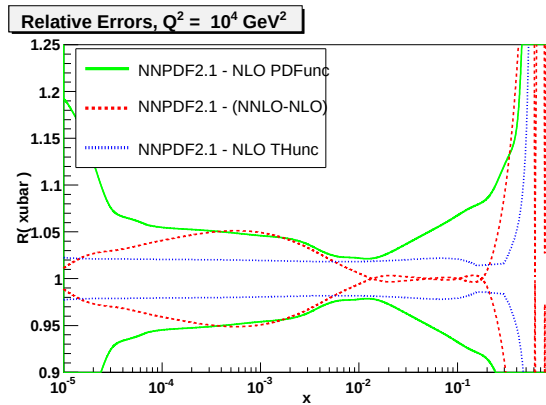
STRANGE



ANTIUP

ANTIDOWN

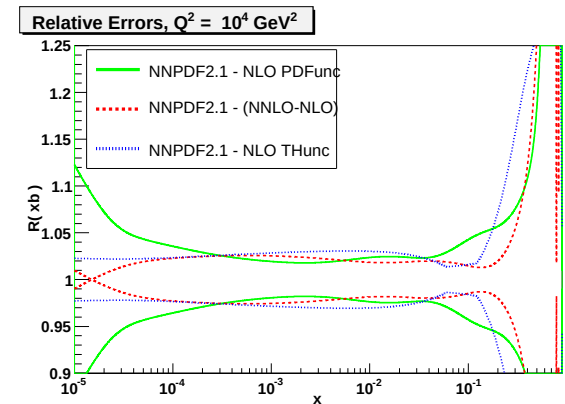
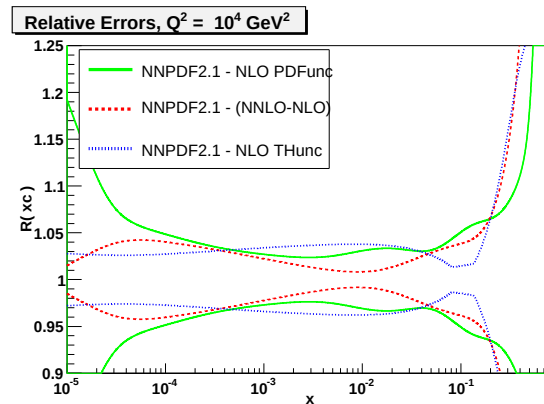
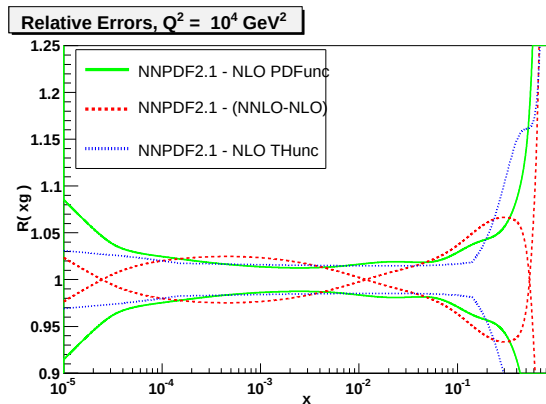
ANTISTRANGE



GLUON

CHARM

BOTTOM



THEORETICAL UNCERTAINTY ON NLO PDF IS ORDER 5% $\Rightarrow$ COMPARABLE TO PDF UNCERTAINTY

THEORETICAL UNCERTAINTIES

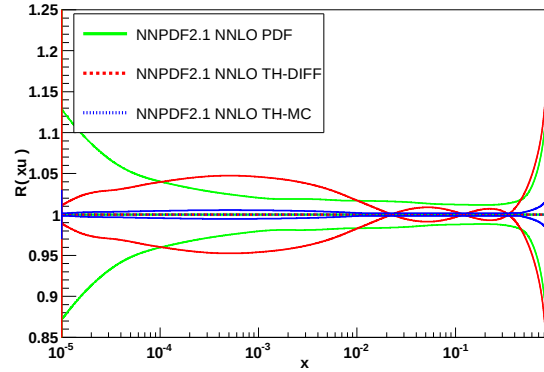
NNLO PDF UNC. VS NLO-NNLO SHIFT VS NNLO CACCIARI-HOUDEAU (NNPDF2.1)

UP

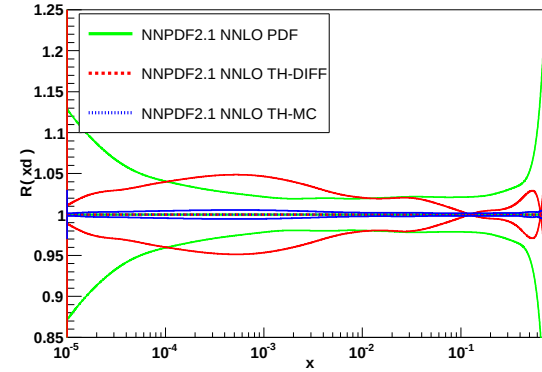
DOWN

STRANGE

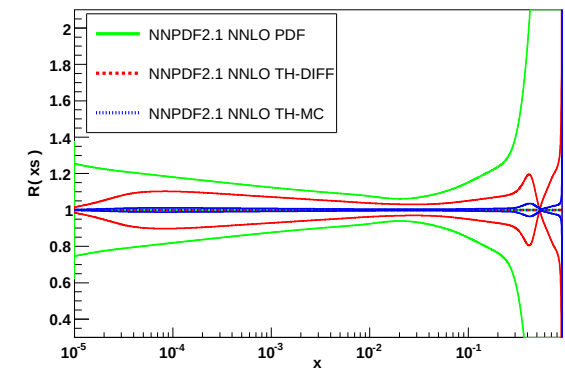
Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$



Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$



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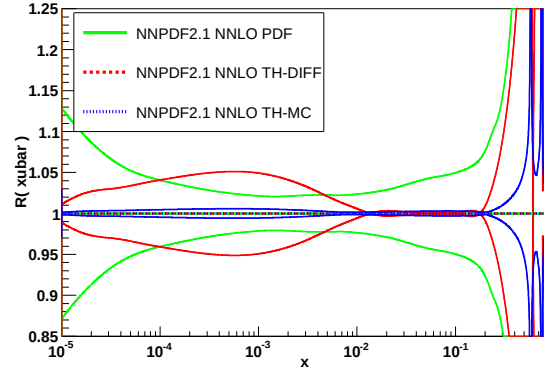


ANTIUP

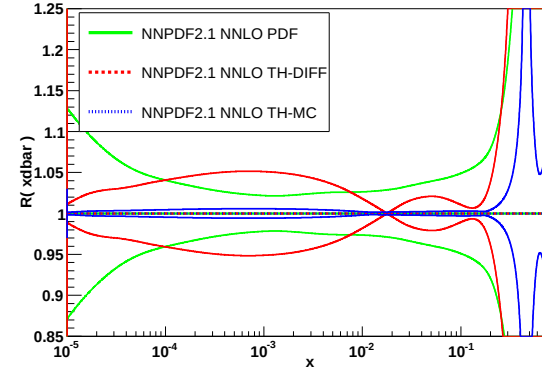
ANTIDOWN

ANTISTRANGE

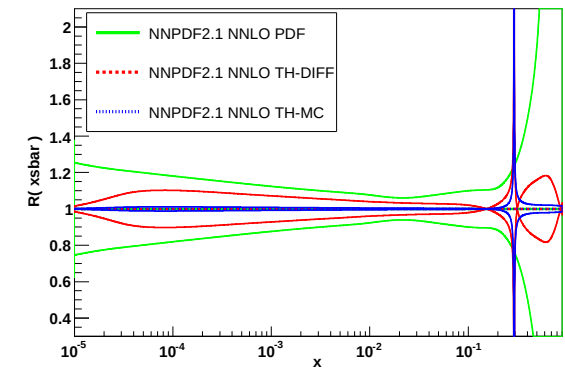
Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$



Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$



Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$

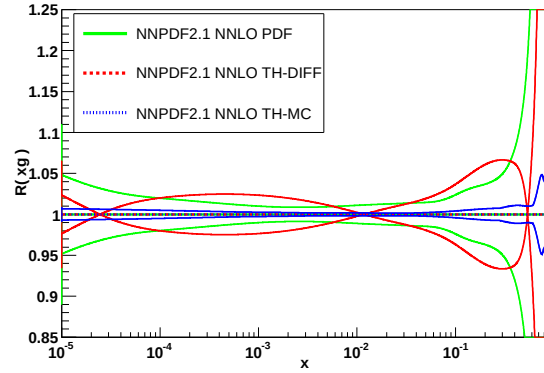


GLUON

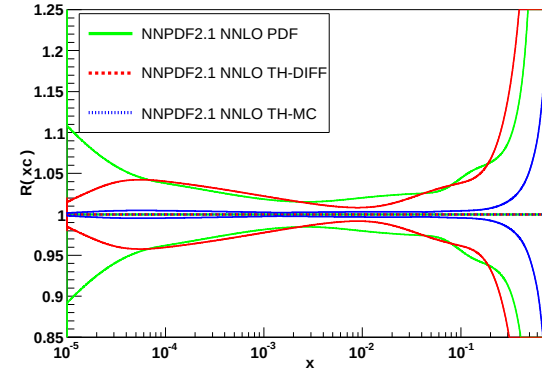
CHARM

BOTTOM

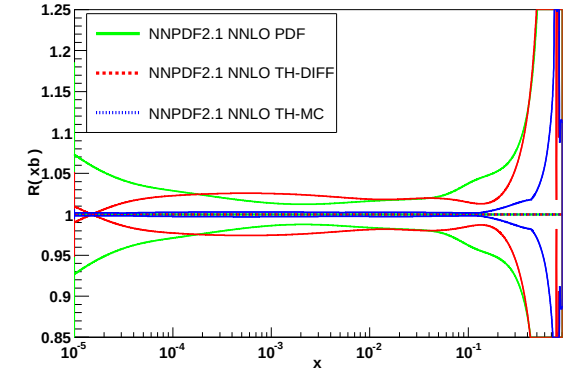
Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$



Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$



Relative Errors, $Q^2 = 10^4 \text{ GeV}^2$

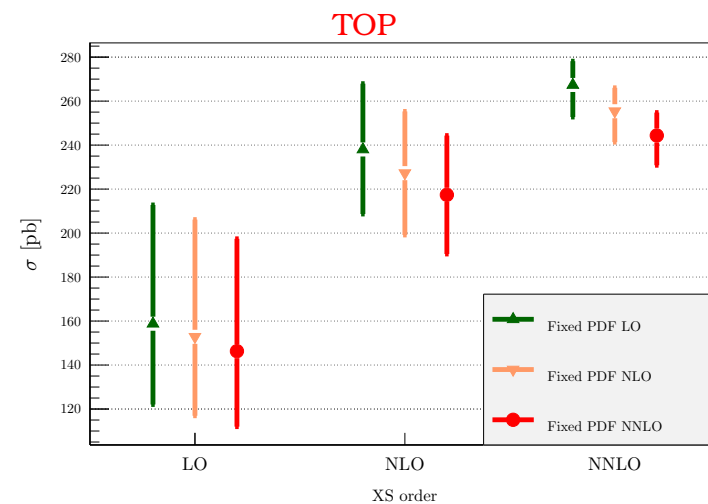
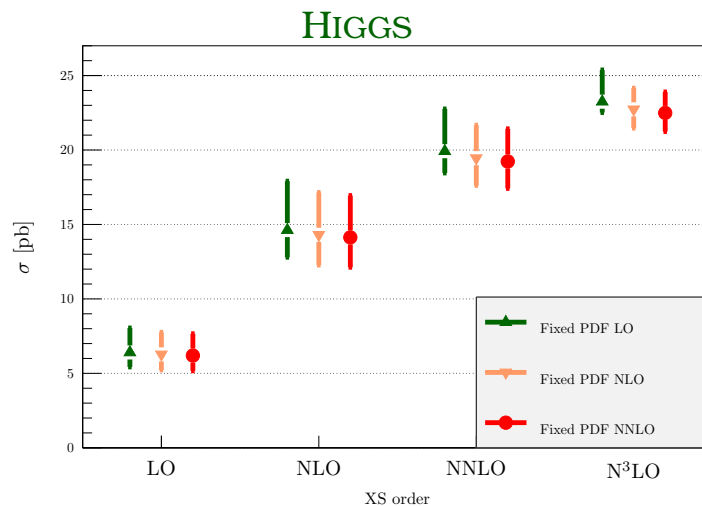


THEORETICAL UNCERTAINTY ON NNLO PDF IS ORDER 1% $\Rightarrow$ SMALLER THAN PDF UNCERTAINTY

N^3 LO PDFs:

- **NEEDED** AT THE 1% ACCURACY LEVEL
- **IMPACT OF N^3 LO DEPENDS ON PROCESS:**
 - **HIGGS GLUON FUSION:** PERTURBATIVE DEP. OF PDF NEGLIGIBLE IN COMPARISON TO MATRIX ELEMENT $\Rightarrow N^3$ LO **NOT NEEDED**
 - **TOP:** PERTURBATIVE DEP. OF PDF SMALLER, BUT NOT NEGLIGIBLE IN COMPARISON TO MATRIX ELEMENT, **ANTICORRELATED** TO IT $\Rightarrow N^3$ LO **NECESSARY**

SCALE UNCERTAINTY & DEP. ON PERTURBATIVE ORDER

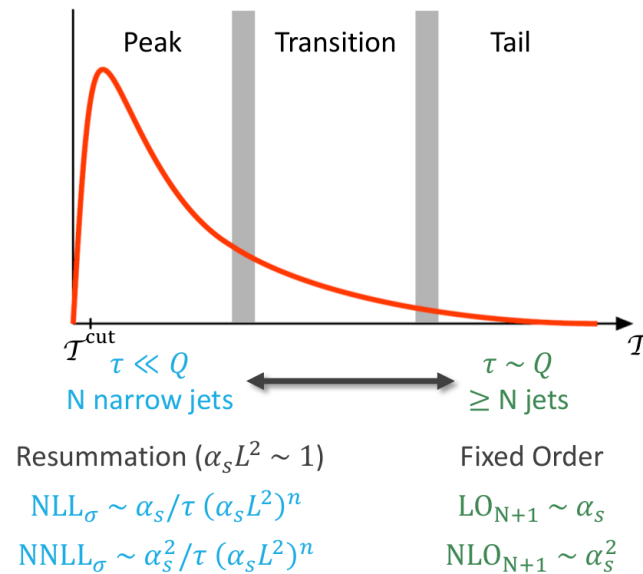


(s.f., Isgrò, Vita, 2014)

WHEN WILL WE HAVE THEM?

- N^3 LO **DIS** COEFFICIENT FUNCTIONS KNOWN
- **BOTTLENECK:** N^3 LO **ANOMALOUS DIMENSIONS**
- **ANOMALOUS DIMENSIONS:** LO: 1974; NLO: 1981; NNLO: 2004; N^3 LO: 2030?

MONTE CARLO PDFs



(Alioli et al., 2012)

- THEORETICAL CALCULATIONS GETTING CLOSER TO THE FINAL STATE:
 - FULLY DIFFERENTIAL
 - PRODUCTION COMBINED WITH DECAY
 - FIDUCIAL OBSERVABLES
- REQUIRES MATCHING OF FIXED-ORDER, RESUMMATION, PARTON SHOWERING (POSSIBLY ALSO HADRONIZATION)
 - $\Rightarrow$ (AT LEAST) THREE METHODS ALREADY AVAILABLE: GENEVA (SCET-BASED); NNLOPS (POWHEG-BASED); UN²LOPS (SHERPA-BASED)
- EVENTUALLY, THIS WILL LEAD TO MONTE CARLO PDFs $\Rightarrow$ SAME TOOLS USED IN PDF DETERMINATION AS IN PHYSICAL PREDICTION

WHAT SHOULD WE AIM FOR 2030?

TYPICAL ACCURACY BELOW 5%, WITH RELIABLE UNCERTAINTY ESTIMATION:

- NNLO AS A DEFAULT FOR QCD CALCULATIONS
- SOME INCLUSIVE N^3 LO OBSERVABLES
(MOSTLY FOR UNCERTAINTY ESTIMATION)
- CONSIDERABLE PARTIAL INFORMATION ON HIGHER ORDERS FROM RESUMMATION (EIKONAL, NEXT-TO-EIKONAL, HIGH ENERGY)
- MIXED NLO QCD-EW CORRECTIONS AS A DEFAULT, SOMETIMES NNLO
- CONSISTENCY BETWEEN GLOBAL PDF FITS
- N^3 LO PDFs (MOSTLY FOR UNCERTAINTY ESTIMATION)

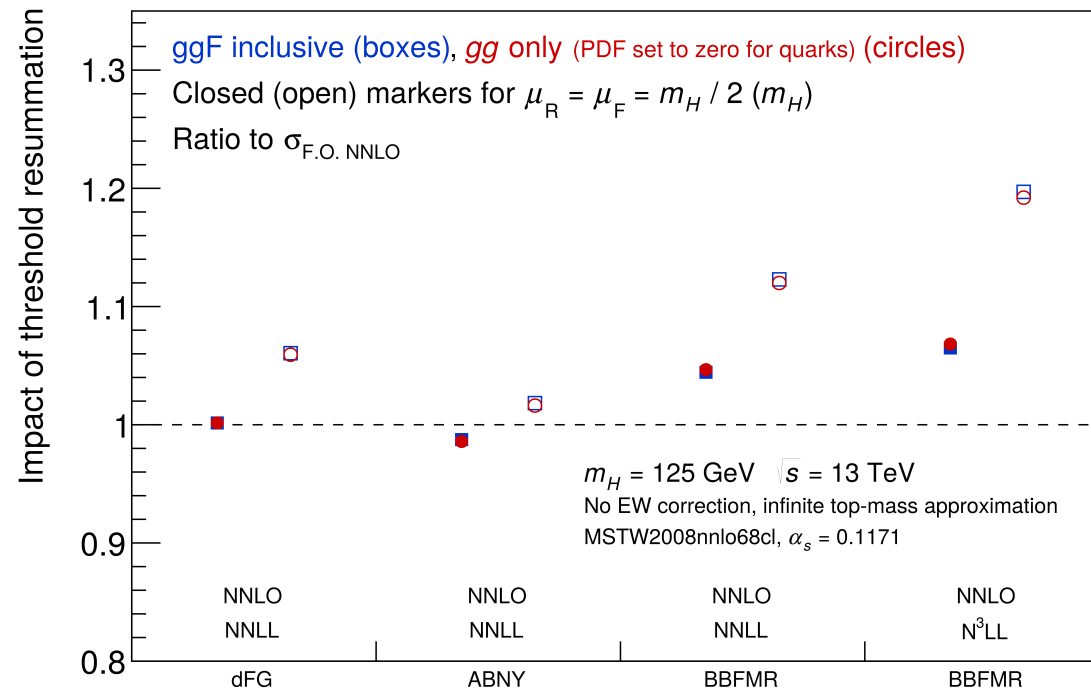
WHEN UNCERTAINTIES ARE SMALL, CAN WE ESTIMATE THEM RELIABLY?

EXTRAS

HIGGS IN GLUON FUSION: RESUMMATION AMBIGUITIES

- LARGE SUBLEADING TERMS $\Rightarrow$ LARGE AMBIGUITIES
- SCET RESUMMATION SMALLER, DUE TO $1/N$ TERMS (Bonvini, SF, Ridolfi, Rottoli)
- RESUMMATION EFFECTIVELY AMOUNTS TO APPROXIMATE HIGHER ORDERS

RESUMMED/UNRESUMMED RATIO (NO CONST. EXPONENTIATION)



note k factor computed wr to NNLO at respective scale

- De Florian, Grazzini (dFG) (HXSWG REFERENCE) RESUM $\ln N$
- Becher, Neubert et al. (ABNY): SCET (z SPACE) APPROACH TO NNLL (REALLY N³LL*)
- Ball et al. (BBFMR): RESUM $\ln N$, DIFFERS FROM dFG BECAUSE OF 'ANALYTIC' RESUMMATION (correct small- N singularities when expanded to finite order)
- BBFMR: N³LL ALSO AVAILABLE

GLOBAL FITS: PROGRESS

- **MMHT** (SUCCESSOR OF MSTW08) DECEMBER 2014
- **CT14** JUNE 2015
- FOR ALL, **PROGRESS IN METHODOLOGY & DATASET**

METHODOLOGY

	NNPDF3.0	MMHT14	CT14
NO. OF FITTED PDFs	7	7	6
PARAMETRIZATION	NEURAL NETS	$x^a(1-x)^b \times$ CHEBYSHEV	$x^a(1-x)^b \times$ BERNSTEIN
FREE PARAMETERS	259	37	30-35
UNCERTAINTIES	REPLICAS	HESSIAN	HESSIAN
TOLERANCE	NONE	DYNAMICAL	DYNAMICAL
CLOSURE TEST	✓	✗	✗
REWEIGHTING	REPLICAS	EIGENVECTORS	EIGENVECTORS

- **MMHT, CT10** LARGER # OF PARMS., ORTHOGONAL POLYNOMIALS
- **NNPDF** CLOSURE TEST

DATASET

	NNPDF3.0	MMHT14	CT14(PREL)
SLAC P,D DIS	✓	✓	✗
BCDMS P,D DIS	✓	✓	✓
NMC P,D DIS	✓	✓	✓
E665 P,D DIS	✗	✓	✗
CDHSW NU-DIS	✗	✗	✓
CCFR NU-DIS	✗	✓	✓
CHORUS NU-DIS	✓	✓	✗
CCFR DIMUON	✗	✓	✓
NUTeV DIMUON	✓	✓	✓
HERA I NC,CC	✓	✓	✓
HERA I CHARM	✓	✓	✓
H1,ZEUS JETS	✗	✓	✗
H1 HERA II	✓	✗	✗
ZEUS HERA II	✓	✗	✗
E605 & E866 FT DY	✓	✓	✓
CDF & D0 W ASYM	✗	✓	✓
CDF & D0 Z RAP	✓	✓	✓
CDF RUN-II JETS	✓	✓	✓
D0 RUN-II JETS	✗	✓	✓
ATLAS HIGH-MASS DY	✓	✓	✓
CMS 2D DY	✓	✓	✗
ATLAS W,Z RAP	✓	✓	✓
ATLAS W PT	✓	✓	✗
CMS W ASY	✓	✓	✓
CMS W +C	✓	✗	✗
LHCb W,Z RAP	✓	✓	✓
ATLAS JETS	✓	✓	✓
CMS JETS	✓	✓	✓
TTBAR TOT XSEC	✓	✓	✗
TOTAL NLO	4276	2996	3248
TOTAL NNLO	4078	2663	3045

OUTLOOK: NNPDF3.1

NEW DATA (MINIMAL SET)

- HERA-II COMBINED DATA arxiv:1506.06042
- DO W ASYMMETRY arXiv:1412.2862
- ATLAS LOW-MASS DY arXiv:1404.1212
- ATLAS PROMPT PHOTON arXiv:1311.1440
- ATLAS $W + c$ arXiv:1402.6263
- ATLAS $Z p_t$ arXiv:1406.3660
- ATLAS $t\bar{t}$ RAPIDITY arXiv:1407.0371
- ATLAS INCLUSIVE JETS 7 TeV 5 FB^{-1} arXiv:1410.8857
- CMS DOUBLE-DIFFERENTIAL 8 TeV arXiv:1412.1115
- CMS $Z p_t$ arXiv:1504.03511
- LHCb $W \rightarrow \mu\nu$ RAP. arxiv:1408.4354

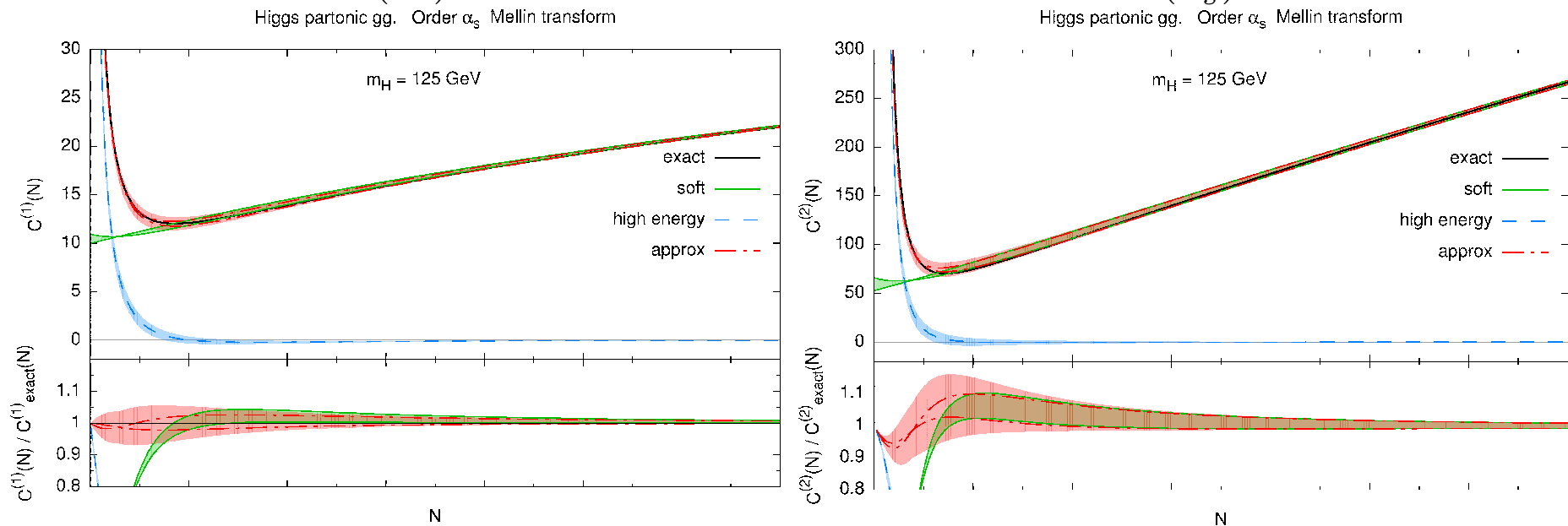
COMBINING INFORMATION

$$C_{\text{APPROX}}^{(k)}(N) = C_{\text{SOFT}}^{(k)}(N) + C_{\text{H.E.}}^{(k)}(N)$$

NOTHING ELSE TO DO: SINGULARITIES TAKE CARE OF THEMSELVES

NLO AND NNLO

EXACT+SOFT+HIGH ENERGY+APPROX.
 $O(\alpha_s)$ $O(\alpha_s^2)$



NOTE FOR $N \sim 1.2$ WHERE SMALL DISCRAPANCY WITH NNLO IS OBSERVED, “EXACT” IS BASED ON AN UNRELIABLE APPROX.

THE PIECES OF THE PUZZLE

https://twiki.cern.ch/twiki/bin/view/LHCPhysics/CERNYellowReportPageAt8TeV#gluon_gluon_Fusion_Process

$$m_h = 125 \text{ GeV}; \quad \sigma = 19.27 \text{ pb}_{-7.8}^{+7.2} (\text{SCALE}) \quad {}_{-6.9}^{+7.5} (\text{PDF}+\alpha_s)$$

$$\mu_r = \mu_F = m_H$$

(NOTE: IF $m_H = 125.5 \rightarrow \sigma = 19.12 \text{ pb}$)

(De Florian, Grazzini, 2012)

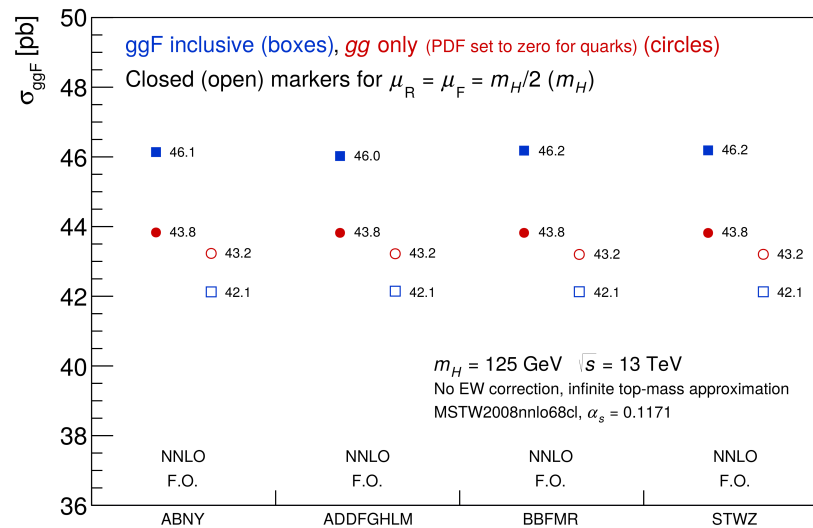
- NNLO+NNLL
- TOP MASS DEP, AND BOTTOM AND CHARM CONTRIBUTIONS AT NLO+NLL
DECREASE $\sim 1.5\%$ DUE TO TOP+BOTTOM (MOSTLY BOTTOM)+
FURTHER $\sim 1\%$ DUE TO CHARM
- NLO ELECTROWEAK CORRECTIONS (COMPLETE FACTORIZATION)
INCREASE $\sim 5\%$ (WOULD CHANGE BY A FEW PERCENT WITH
PARTIAL FACTORIZATION)
- COMPLEX-POLE (OFFP)
FINITE-WIDTH BELOW $\sim 1\%$ FOR LIGHT HIGGS
- NO MIXED QCD-EW AND REAL EW RADIATION
BELOW $\sim 1\%$

BENCHMARKING

THE GROUPS INVOLVED:

- **ABNY** Ahrens, Becher, Neubert, Yang → **NNLO+NNLL RESUMMED (SCET)**
- **STWZ** Stewart, Tackmann, Walsh, Zuberi → **NNLO WITH EXP. CONSTANTS**
- **DFMMV** de Florian, Mazzitelli, Moch, Vogt → **APPROXIMATE N³LO**
- **BBFMR** Ball, Bonvini, SF, Marzani, Ridolfi → **APPROXIMATE N³LO+N³LL**
- **ADDFGLHM** Anastasiou, Duhr, Dulat, Furlan, Gehrmann, Lazopoulos, Herzog, Mistlberger → **PARTIAL N³LO**

THE NNLO CROSS-SECTION

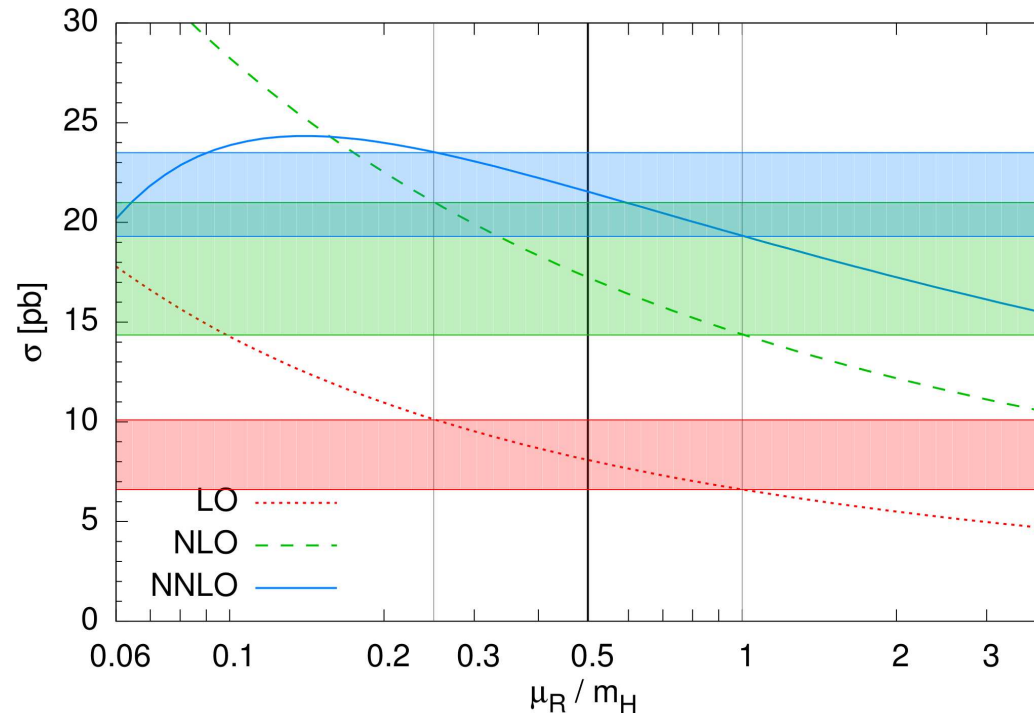


COMMON SETTINGS FOR PDFs, EW, HW, ETC ADOPTED;
NOT MEANT TO BE REALISTIC, FOCUS ON SIZE OF QCD CORRECTIONS

BEYOND NNLO

LO, NLO AND NNLO

Higgs cross section $m_H = 125 \text{ GeV}$ @ LHC 8 TeV



- THE PERTURBATIVE SERIES FOR $gg \rightarrow H$ **CONVERGES SLOWLY**
- **SCALE VARIATION UNDERESTIMATES** SHIFT TO NEXT ORDER
- FACTORIZATION SCALE DEPENDENCE NEGLIGIBLE, ONLY RENORMALIZATION SCALE DEPENDENCE SIGNIFICANT

QUESTIONS

- **DOES RESUMMATION IMPROVE** THE FIXED-ORDER RESULT? How **AMBIGUOUS** IS IT?
- CAN WE **CONSTRUCT A REASONABLE APPROXIMATION** TO $N^3\text{LO}$ BASED ON **CURRENT KNOWLEDGE**?

RESUMMATION

$$\sigma(N, m_H^2) \equiv \int_0^1 d\tau \tau^{N-2} \sigma(\tau, m_H^2); \quad z = \frac{M_h^2}{s} \left(\frac{\ln^{k-1}(1-x)}{1-x} \right)_+ \Leftrightarrow \ln^k N + O(\ln^{k-1} N) + O\left(\frac{1}{N}\right)$$

$$\sigma_{\text{res}}(N, \alpha_s) = \sigma_0 g_0(\alpha_s) \exp \left[\frac{1}{\alpha_s} g_1(\alpha_s \ln N) + g_2(\alpha_s \ln N) + \alpha_s g_3(\alpha_s \ln N) + \dots \right];$$

$$g_0(\alpha_s) = 1 + \alpha_s g_{0,1} + \alpha_s^2 g_{0,2} + \mathcal{O}(\alpha_s^3);$$

N^kLL RESUMMATION: ALL g_i WITH $i < k + 1$, ALL g_{0j} WITH $j < k$
TO ORDER α_s^n , $\ln^m N$ PREDICTED WITH $2(n - k) \leq m \leq 2n$

in SCET literature, this is called N^kLL', at N^kLL one less power of log: $2(n - k) + 1 \leq m \leq 2n$

AMBIGUITIES

- SHOULD g_0 SIT IN THE EXPONENT, ENTIRELY OR IN PART?

NOTE $g_{0,1}$ NATURALLY PART OF g_{i+2}

⇒ DIFFERENCE SUBLEADING!

- HOW DOES ONE FIX SUBLEADING TERMS?

$\ln N$, $\ln N + 1$, $\psi_0(N)$ (MELLIN OF $\left(\frac{1}{1-x} \right)_+$) BEHAVE IN THE SAME WAY AS $N \rightarrow \infty$, UP TO $O(1)$ &

$O(1/N)$ TERMS

⇒ DIFFERENCE SUBLEADING!

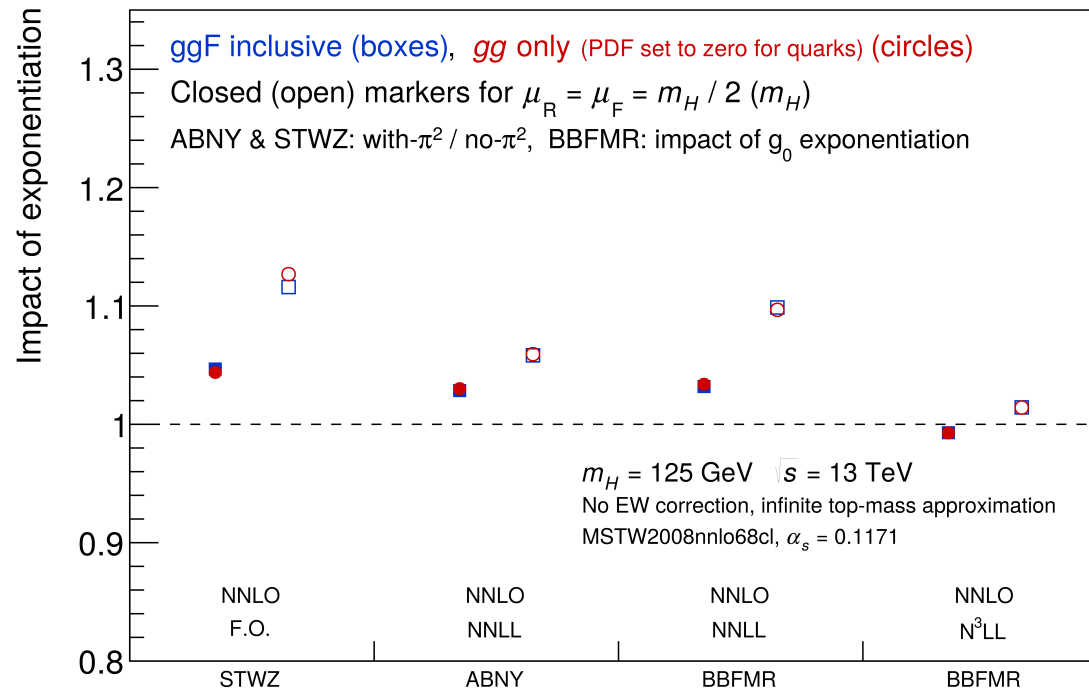
NOTE IN SCET APPROACH, RESUMMATION NATURALLY IN z SPACE ⇒ SUBLEADING

DIFFERENCES

CONSTANT EXPONENTIATION

- IN SCET APPROACH, π^2 EXPONENTIATION W. IMAGINARY SCALE:
STWZ: ADDED TO FIXED NNLO; ABNY: ADDED TO NNLO+NNLL
- IN dQCD RESUMMATION DIS-SCHEME CONSTANT EXPONENTIATION
BBFMR: BOTH NNLO+NNLL AND NNLO+N³LL AVAILABLE

EXPONENTIATED/UNEXPONENTIATED RATIO (NO RESUMMATION)



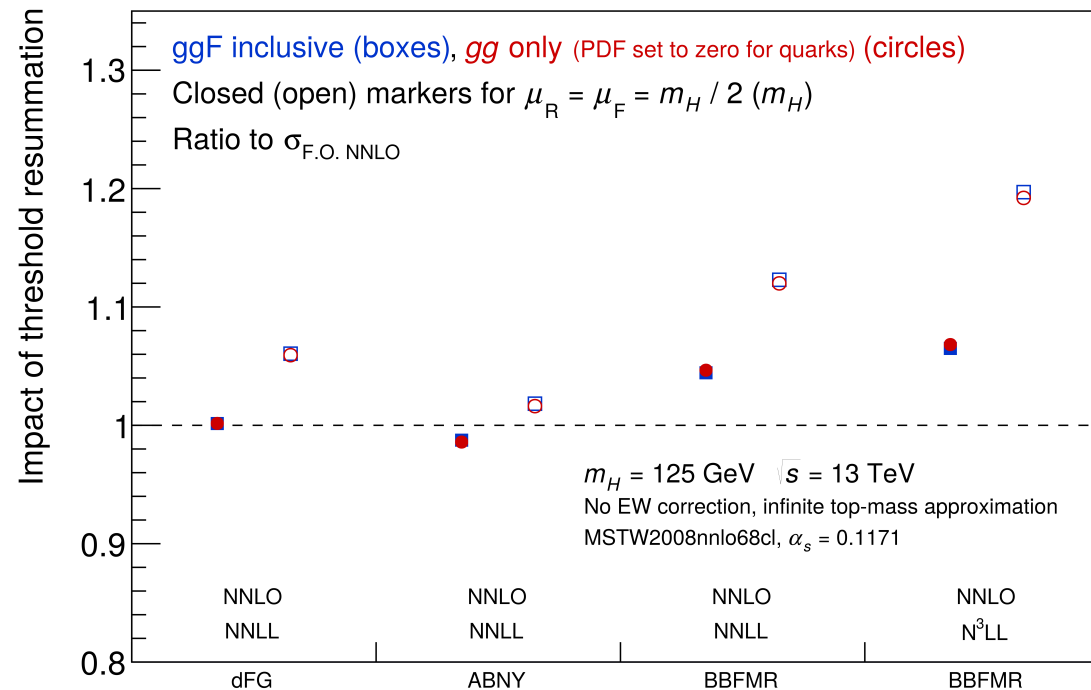
note k factor computed wr to NNLO at respective scale

- IMPACT SMALLER FOR $\mu_R = m_H/2$ THAN FOR $\mu_R = m_H$
- IMPACT SMALLER AT RESUMMED THAN AT FIXED ORDER
- IMPACT SMALLER AT N³LL THAN AT NNLL

CHOICE OF LOGS

- DFG: CURRENT REFERENCE $\mu_R = m_H$
- ABNY: SCET (z SPACE) APPROACH TO NNLL (REALLY N^3LL^* , OR SCET- N^3LL)
- BBFMR: DIFFERS FROM DFG BECAUSE OF 'ANALYTIC' RESUMMATION
(correct small- N singularities when expanded to finite order)
- BBFMR: N^3LL ALSO AVAILABLE

RESUMMED/UNRESUMMED RATIO (NO CONST. EXPONENTIATION)



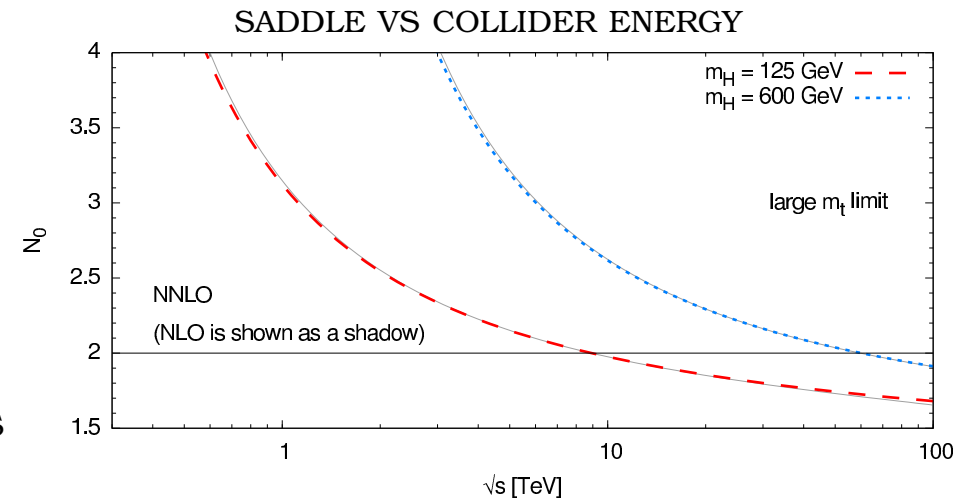
note k factor computed wr to NNLO at respective scale

- IMPACT **SMALLER** FOR $\mu_R = m_H / 2$ THAN FOR $\mu_R = m_H$
- z -SPACE: **LITTLE** EFFECT, ANALYTIC: **LARGER** EFFECT
- IMPACT **LARGER** AT N^3LL THAN AT NNLL

FROM RESUMMATION TO N³LO

TECHNICAL INTERLUDE

- PARTONIC CROSS SECTIONS $\hat{\sigma}(z, \alpha_s)$ ARE DISTRIBUTIONS, $0 \leq z = \frac{M_h^2}{s} \leq 1$
- THEIR MELLIN TRANSFORMS $\hat{\sigma}(N, \alpha_s)$ ARE ORDINARY FUNCTIONS
- FOR GIVEN m_H^2 AND s ONLY ONE “SADDLE” N VALUE CONTRIBUTES
- BONUS: HAD. XSCT PRODUCT OF PARTONIC TIMES LUMI: $\sigma(N, m_H^2) = \hat{\sigma}(m_H^2, \alpha_s) \mathcal{L}(N)$



Q: WHY ARE THERE LARGE RESUMMATION AMBIGUITIES?

A: BECAUSE $\beta_0 \alpha_s \ln 2 \approx 0.04 \ll 1$: RESUMMATION IS PERTURBATIVE!

IN ALL FLAVOURS OF RESUMMATION DISCUSSED ABOVE

BETWEEN 2/3 AND 3/4 OF THE NNLL OR N³LL RESUMMATION COMES FROM N³LO,

ALL REMAINING INFINITE ORDERS CONTRIBUTING $\sim 1 \div 2\%$

RESUMMATION IS AMBIGUOUS BECAUSE N³LO IS APPROXIMATE!

N³LO: WHAT DO WE KNOW

THE SOFT EXPANSION

ABOUT $N \rightarrow \infty$ OR $z = 1$

$$\begin{aligned} \hat{\sigma}^{(3)}(N) &= c_{06} \ln^6 N + c_{05} \ln^5 N + \dots + c_{02} \ln^2 N + c_{01} \ln N + c_{00} + && \text{SOFT (S)} \\ &+ \frac{1}{N} (c_{15} \ln^5 N + c_{14} \ln^4 N + \dots + c_{10}) + && \text{NEXT-TO-SOFT (NS)} \\ &+ \frac{1}{N^2} (c_{25} \ln^5 N + c_{24} \ln^4 N + \dots + c_{20}) + \dots && \text{NNS} \end{aligned}$$

$$\begin{aligned} \hat{\sigma}^{(3)}(z) &= c'_{06} \left(\frac{\ln^5(1-z)}{1-z} \right)_+ + \dots + c'_{01} \left(\frac{1}{1-z} \right)_+ + c'_{00} \delta(1-z) + && \text{SOFT (S)} \\ &+ (c'_{15} \ln^5(1-z) + c'_{14} \ln^4(1-z) + \dots + c'_{10}) + && \text{NEXT-TO-SOFT (NS)} \\ &+ (1-z) (c'_{25} \ln^5(1-z) + c'_{24} \ln^4(1-z) + \dots + c'_{20}) + \dots && \text{NNS} \end{aligned}$$

N SPACE c COEFFICIENTS DETERMINE THE z SPACE c' AND CONVERSELY:

- N -SPACE N^k SOFT $\Leftrightarrow z$ -SPACE N^k SOFT
- N -SPACE N^k LL $\Leftrightarrow z$ -SPACE N^k LL
- MELLIN OF z SPACE N^k SOFT EQUALS N -SPACE N^k SOFT PLUS N^{k+1} TERMS
- MELLIN OF z SPACE N^k LL EQUALS N -SPACE N^k LL SOFT PLUS N^{k+1} LL TERMS

N³LO: WHAT DO WE KNOW

THE SOFT EXPANSION ABOUT $N \rightarrow \infty$ OR $z = 1$

$$\begin{aligned}\hat{\sigma}^{(3)}(N) &= c_{06} \ln^6 N + c_{05} \ln^5 N + \dots + c_{02} \ln^2 N + c_{01} \ln N + c_{00} + && \text{SOFT (S)} \\ &+ \frac{1}{N} (c_{15} \ln^5 N + c_{14} \ln^4 N + \dots + c_{10}) + && \text{NEXT-TO-SOFT (NS)} \\ &+ \frac{1}{N^2} (c_{25} \ln^5 N + c_{24} \ln^4 N + \dots + c_{20}) + \dots && \text{NNS}\end{aligned}$$

$$\begin{aligned}\hat{\sigma}^{(3)}(z) &= c'_{06} \left(\frac{\ln^5(1-z)}{1-z} \right)_+ + \dots + c'_{01} \left(\frac{1}{1-z} \right)_+ + c'_{00} \delta(1-z) + && \text{SOFT (S)} \\ &+ (c'_{15} \ln^5(1-z) + c'_{14} \ln^4(1-z) + \dots + c'_{10}) + && \text{NEXT-TO-SOFT (NS)} \\ &+ (1-z) (c'_{25} \ln^5(1-z) + c'_{24} \ln^4(1-z) + \dots + c'_{20}) + \dots && \text{NNS}\end{aligned}$$

- c_{06} TO c_{02} DETERMINED BY NNLL RESUMMATION
- c_{01} COMPUTED (Moch, Vogt, 2005; Laenen, Magnea, 2005)
- COMPLETE z -SPACE SOFT $\Rightarrow c'_{00}$ (Anastasiou et al. 2014)
- $c_{15}, c_{24}, \dots, c_{51}$ KNOWN FROM LL RESUMMATION (Catani, deFlorian, Grazzini, 2001)
- c_{14} COMPUTED BASED ON CONJECTURE (deFlorian, Mazzitelli, Moch, Vogt, 2014)
- COMPLETE z -SPACE NS $\Rightarrow$ ALL c'_{1i} (Anastasiou et al. 2015) $\rightarrow$ CONJECTURE VALIDATED
- COMPLETE z -SPACE $\ln^4(1-z)$ & $\ln^3(1-z) \Rightarrow$ ALL c'_{i4}, c'_{i3} (Anastasiou et al. 2015)
(FORMALLY LEADING LOGS)

N³LO: WHAT DO WE KNOW

THE SOFT EXPANSION ABOUT $N \rightarrow \infty$ OR $z = 1$

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N³LO: WHAT DO WE KNOW

THE HIGH-ENERGY EXPANSION

SMALL N POLES OR ABOUT $z = 0$

rightmost singularity at $N = 1$:

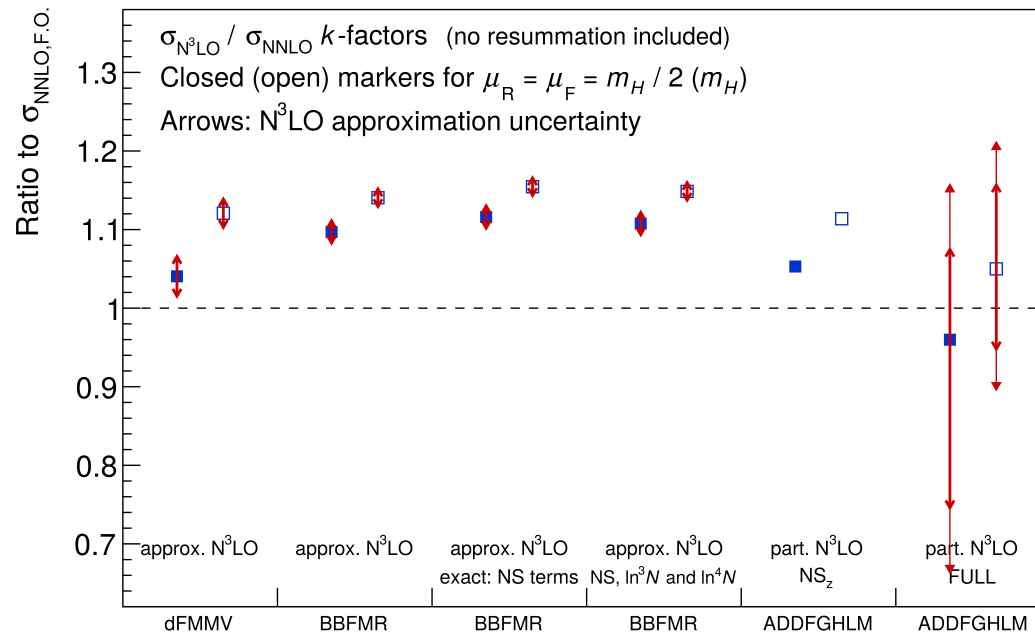
$$\begin{aligned}
 \hat{\sigma}^{(3)}(N) &= \left(d_{13} \frac{1}{(N-1)^3} + d_{12} \frac{1}{(N-1)^2} + d_{11} \frac{1}{(N-1)} \right) + && \text{LEADING LOG } x \text{ (LL}x\text{)} \\
 &+ \left(d_{06} \frac{1}{N^6} + \cdots + d_{01} \frac{1}{N} \right) + && \text{NLL}x \\
 &+ \left(d_{-16} \frac{1}{(N+1)^6} + \cdots + d_{-11} \frac{1}{N+1} \right) + + \dots && \text{NNLL}x \\
 \hat{\sigma}^{(3)}(z) &= \left(d_{13} \frac{\ln^2 z}{z} + d_{12} \frac{\ln z}{z} + d_{11} \frac{1}{z} \right) + && \text{LEADING LOG } x \text{ (LL}x\text{)} \\
 &+ \left(d_{06} \ln^6 z + \cdots + d_{01} \ln z \right) + && \text{NLL}x \\
 &+ z \left(d_{-16} \ln^6 z + \cdots + d_{-11} \ln z \right) + \dots && \text{NNLL}x
 \end{aligned}$$

- AS ABOVE, z -SPACE $\Leftrightarrow$ N -SPACE, ORDER BY ORDER (LOG & POWER), UP TO SUBLEADING TERMS
- ONLY d_{13} KNOWN EXACTLY FROM BFKL+SMALL x RESUMMATION
- CLASS OF DOMINANT (?) CONTRIBUTIONS TO d_{12} , d_{11} KNOWN FROM NL BFKL
- ALL SUBLEADING SINGS UNKNOWN

APPROXIMATE N³LO

- **DFMMV**: N -SPACE APPROXIMATE NS; NOW NS KNOWN EXACTLY: SMALL EXPECTED CHANGE
- **BBFMR**: ALL KNOWN SINGULARITIES (SOFT & HIGH ENERGY):
NS AND BEYOND APPROXIMATELY DETERMINED
→ IMPROVED WITH EXACT NS → IMPROVED WITH EXACT $\ln^4 N$, $\ln^3 N$ (SMALL CHANGES)
- **ADDFGHLM**: ALL KNOWN z SPACE TERMS (FULL) → PURE z SPACE NS ALSO SHOWN

N³LO/NNLO k -FACTOR



note k factor computed wr to NNLO at respective scale

UNCERTAINTIES (ARROWS)

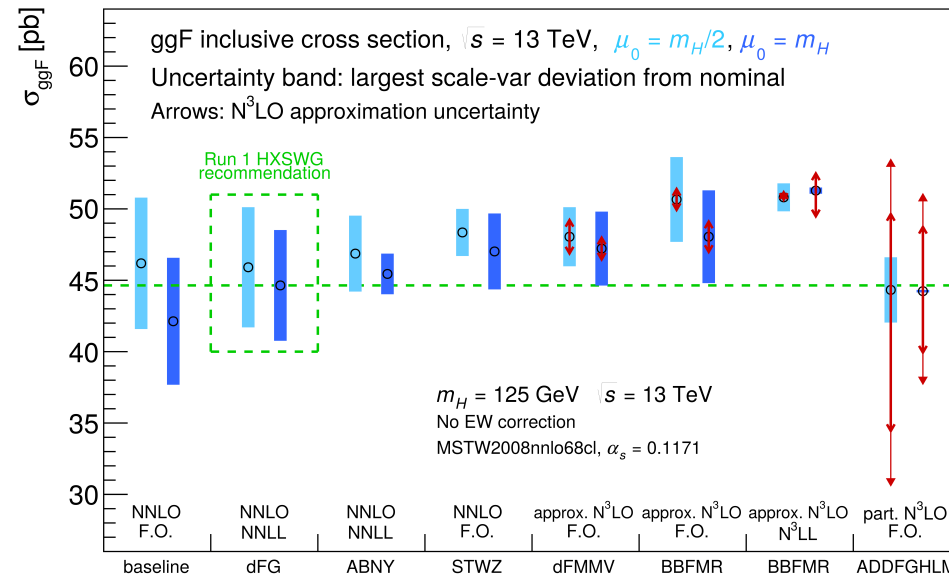
- **DFMMV**: S-NS BAND
- **BBFMR**: ESTIMATE OF SUBLEADING SINGULARITIES VALIDATED BY KNOWN ORDERS
- **ADDFGHLM**: ENVELOPE OF SCALE VAR. OF NS-FULL DIFFERENCE AT m_H , $\times 1.5$ OR 1

SIZE OF UNCERTAINTY HOTLY DEBATED!!

WHAT IS THE CROSS-SECTION?

- **STWZ**: DIFFERS FROM NNLO BY CONST. EXPONENTIATION
- **ABNY**: DIFFERS FROM NNLO+NNLL (REF) BY CONST. EXPONENTIATION & DIFFERENT RESUMMATION
- **DFMMV**: APPROXIMATE $N^3\text{LO}$ BASED ON APPROXIMATE N -SPACE NS; ARROWS: S-NS DIFFERENCE
- **BBFMR**: APPROXIMATE $N^3\text{LO}$ BASED ON ANALYTIC APPROX (INCLUDING NS+BFKL); ARROWS: ESTIMATED MISSING SINGS+ FACT SCALE VARIATION FROM MISSING QUARK CHANNELS; $\rightarrow N^3\text{LO}+N^3\text{LL}$ BEST PREDICTION
- **ADDFGHLM**: ALL KNOWN z SPACE TERMS; ARROWS: ENVELOPE OF SCALE VARIATION OF NS-FULL DIFFERENCE

THE GLUON FUSION CROSS-SECTION



all numbers given with benchmark settings (K -factor to dFG very stable)
 band is seven-point fact. and ren. scale variation