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# Tianlai and the Foreground Removal Problem

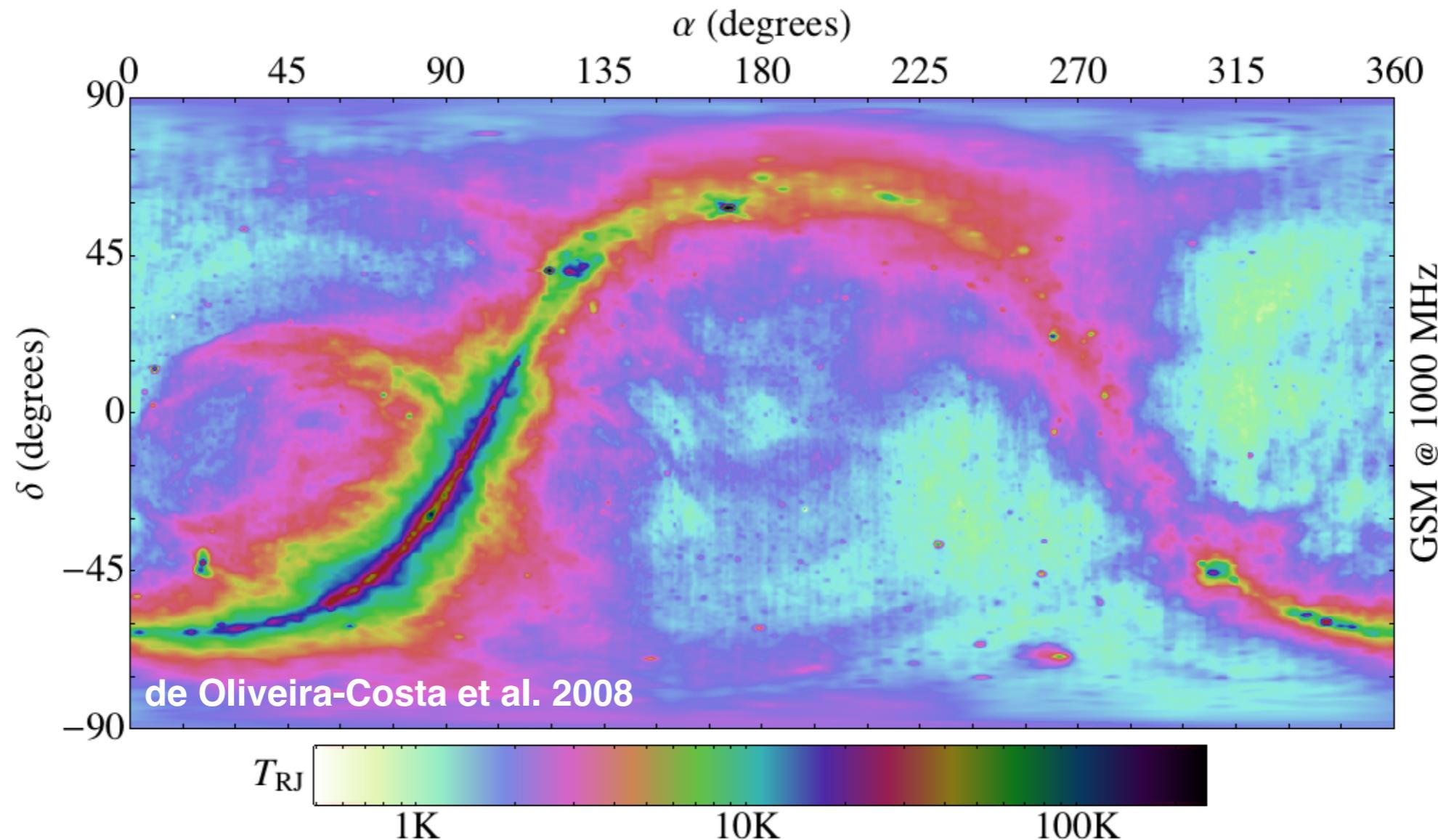
Albert Stebbins

2015 Tianlai 21cm Workshop

Balikun, Xinjiang, China

9 September 2015

# Foreground removal efficacy remains a significant issue



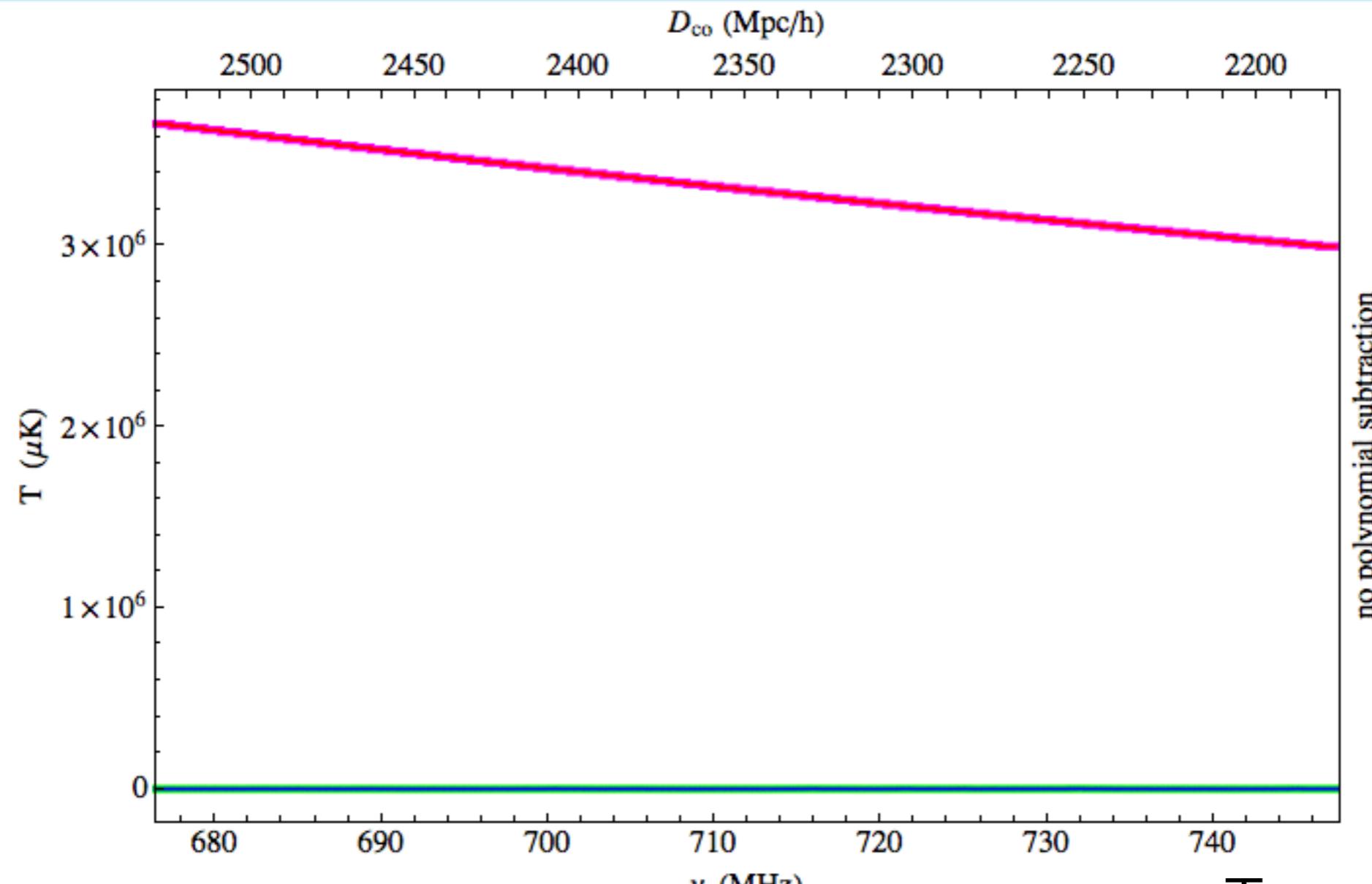
Even for dark radio sky  $\sim 1K$  foreground is  $\sim 10^4$  larger than  $\sim 100\mu K$  signal  
Foregrounds are expected to be smooth in frequency  
... but are they?

# Non-Smooth Spectrum Foregrounds

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- While it is true that optically thin free-free or synchrotron emission is smooth in the optically thin limit for any electron energy distribution - yet it need not be so when self-absorption is present.
- There is evidence for synchrotron self-absorption in gigahertz peaked sources (GPS).
- Faraday rotation linearly polarized light can cause the linearly polarization to have oscillatory behavior which can leak into the inferred intensity.
- Peter Timbie presented evidence yesterday for this effect in the GBT data.
- It is unlikely that these could have spectral features as sharp as those expected in 21cm spectrum it can contaminate the low  $k$  modes which are important for measuring quantities like  $f_{NL}$ .

# Ideally Smooth foreground subtraction should work well



$$T_{\text{fit}} = \sum_p a_p \ln[v]^p$$

- Many algorithms proposed to take advantage of this
- See excellent overview yesterday by Yichao Li
- Would like to get noise down to  $\sim 50\mu\text{K}$  level to see real non-linear structures

# Mode Mixing

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- An interferometer with a finite number of elements will only “see” a finite number of “beams” on the sky.
- The Hilbert space of all linear combinations of beams we call the **space of beams**.
- This space of beams generally depends frequency. This frequency dependence of the Hilbert spaces is called **mode mixing** because it mixes frequency dependence and angle dependence.
- If we could remove mode mixing one can directly measure the angular average spectrum with no assumption about angular structure.
- For this filtering smooth spectrum foregrounds is better when the amount of mode mixing is minimized.

# Beam Projection and Purity

- Given a metric,  $\cdot$ , on the space of beam define the **beam projection operator**:

$$\mathfrak{B}[v] \equiv \sum_{i,j} \mathbf{B}_i[v] \cdot (\mathbf{B}_i[v] \cdot \mathbf{B}_j[v])^{-1} \cdot \mathbf{B}_j[v]$$

where  $\mathbf{B}_i[v]$  are the frequency dependent beam in the Stokes parameter space.

- $\mathfrak{B}[v]$  has  $n_{\text{beam}}$  (number of beams) unit eigenvalues and the rest zero.
- Define the **purity operator** by

$$\mathcal{P} \equiv \int dv W[v] \mathfrak{B}[v] = \sum_a p_a \ \boldsymbol{p}_a \otimes \boldsymbol{p}_a$$

where  $W[v]$  is a  $v$  weight function (or **purity band**) such that:  $\int dv W[v] = 1$

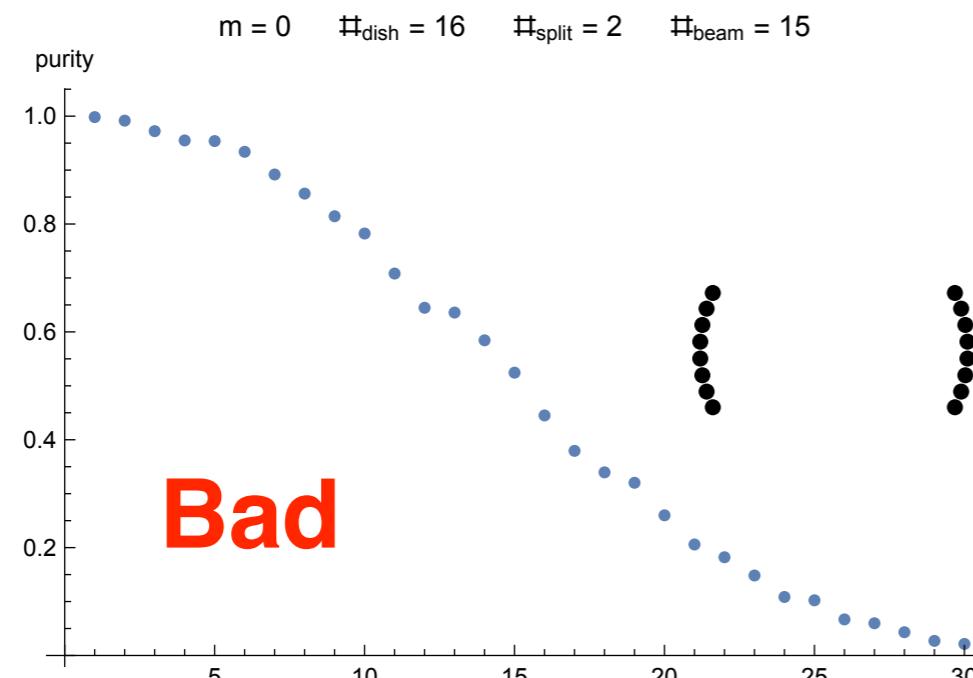
the  $\boldsymbol{p}_a$  (eigenvectors) are **purity eigenbeams**:  $\boldsymbol{p}_i \cdot \boldsymbol{p}_j = \delta_{ij}$

the  $p_a$  (eigenvalues) are **purities**:  $0 \leq p_a \leq 1$  and  $\sum_a p_a = 1$

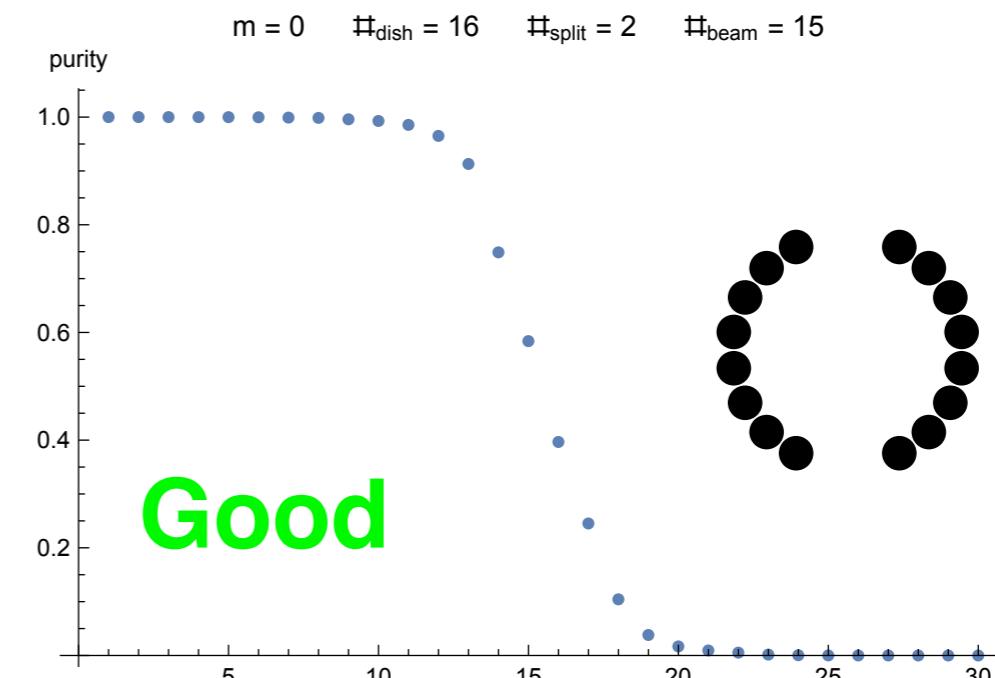
- The  $\boldsymbol{p}_a$  with the largest purity have the least mode mixing!
- The  $\boldsymbol{p}_a$  with  $p_a \ll 1$  have large amounts of mode mixing and high pass filtering is less effective at removing smooth but highly anisotropic foregrounds.

# Purity and Telescope Design

- There are at most  $n_{beam}$  very pure ( $p_a \approx 1$ ) modes  $\mathbf{p}_a$ .
- N.B.  $p_a \rightarrow 1$  in the limit of zero bandwidth:  $W[v] \rightarrow \delta[v - v_0]$
- A **high purity interferometer** is one which for a given bandwidth has close to  $n_{beam}$  very pure modes. They are useful for understanding the underlying spectra of the emission.



**Bad**



**Good**

- Define the **purity number**  $= -\ln[1-p_a]$  which is large for very pure modes
- Dense arrays with large overlap,  $\mathbf{B}_i[v] \cdot \mathbf{B}_j[v]$ , do better than sparse arrays (see Reza's talk).
- One never does worse by adding an additional element to an existing array (but  $\neq \$$ )

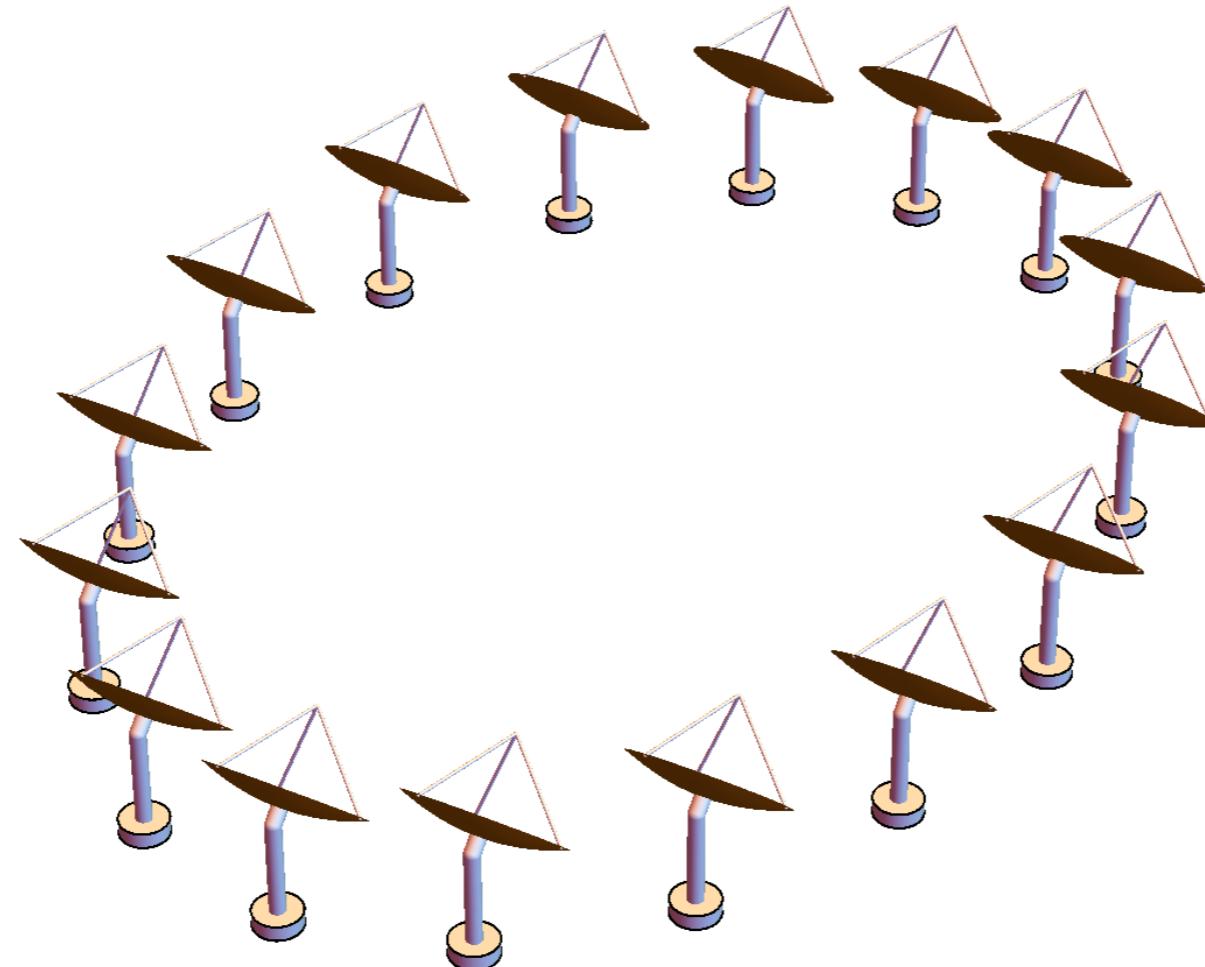
# Purity and Telescope Design and Tianlai



# Polarscope

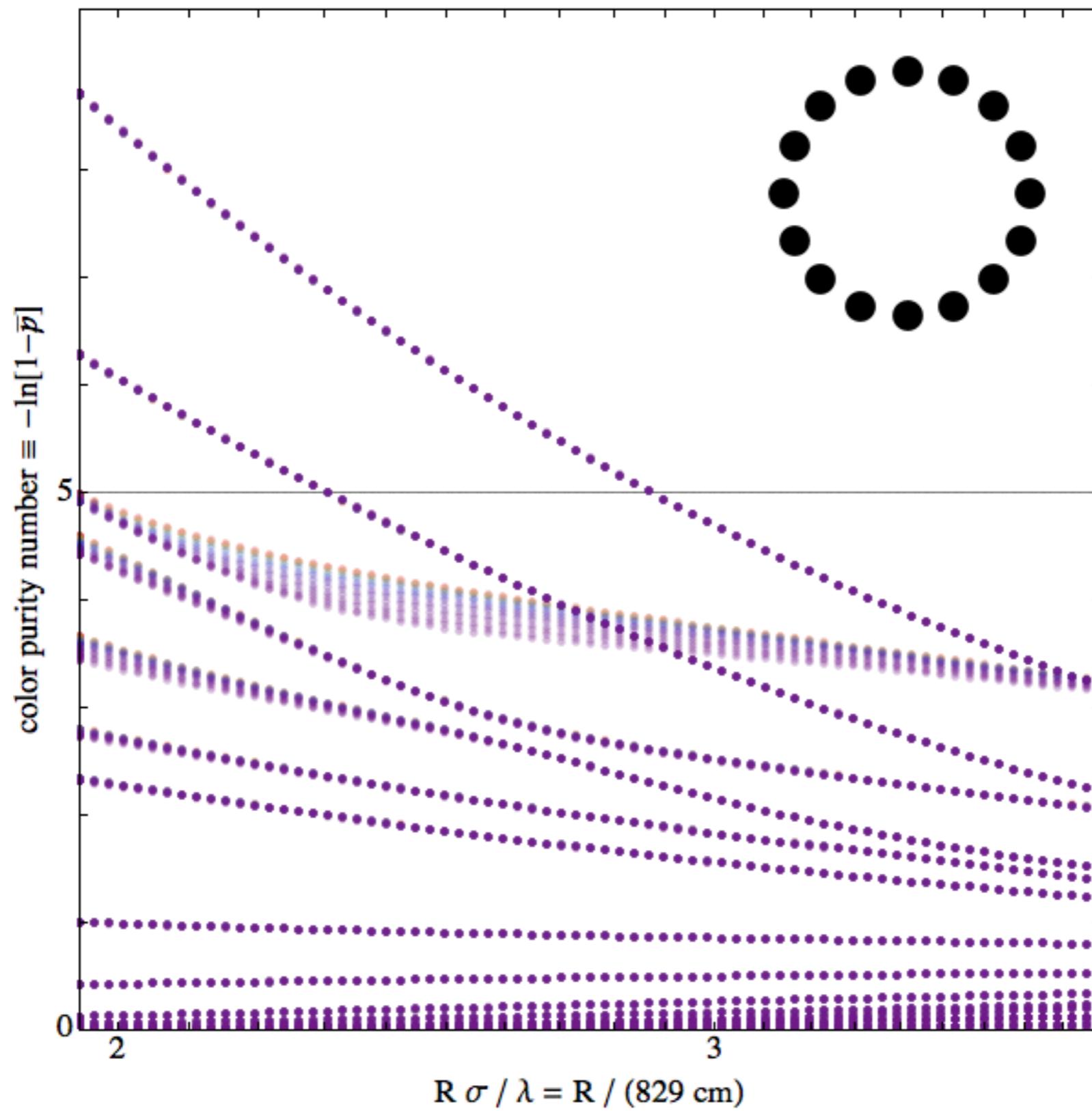
A **polarscope** is a transit interferometer consisting of a number antenna each pointing directly at a Celestial Pole, North or South.

Since it always points at same spot it integrates to low noise very rapidly

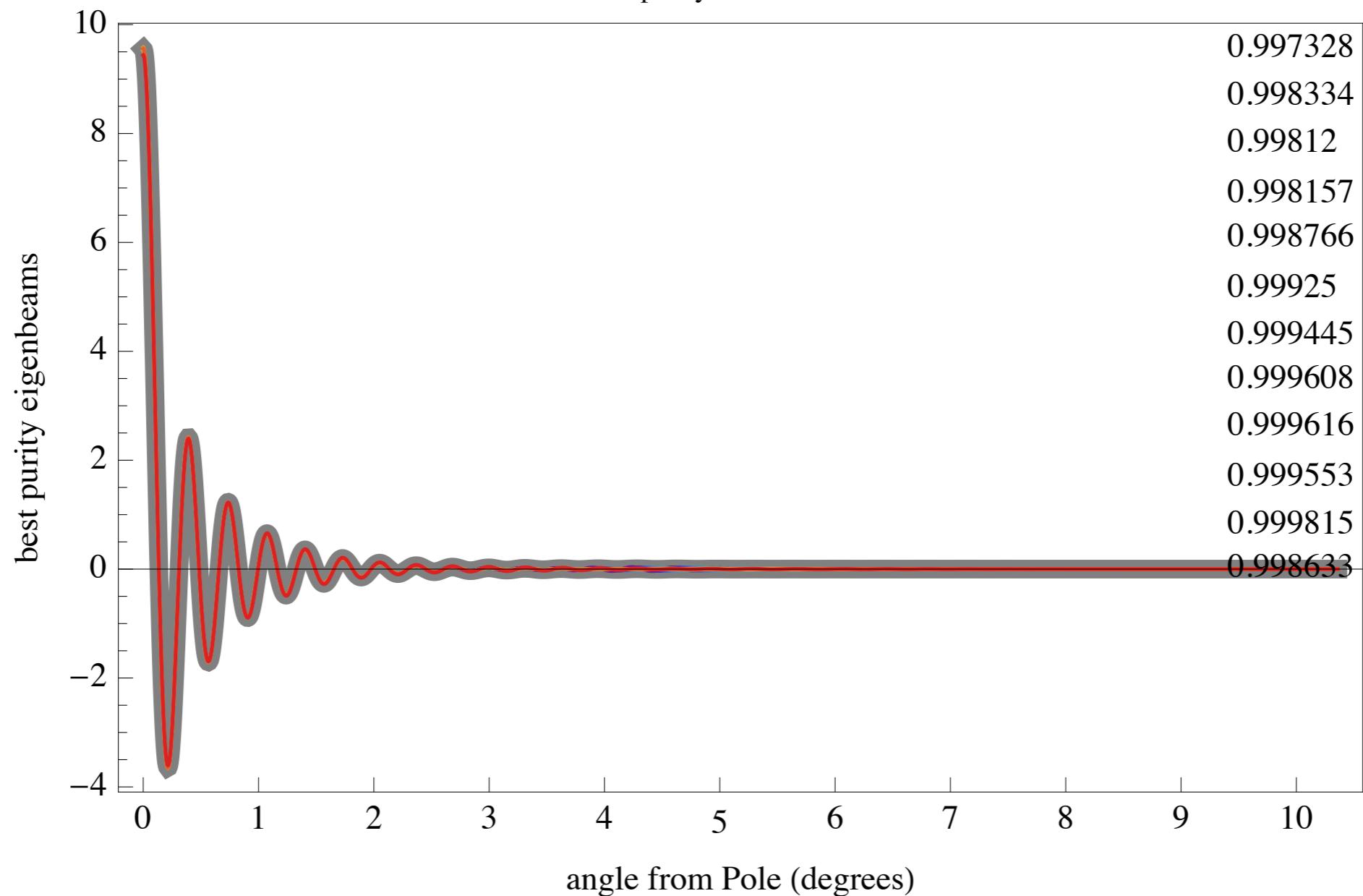


DISCLAIMER: Celestial sources near poles move slowly so a polarscope has very little handle on day timescale transients, *e.g.* ground pickup. N.B. Signals repeat every half day not every day.

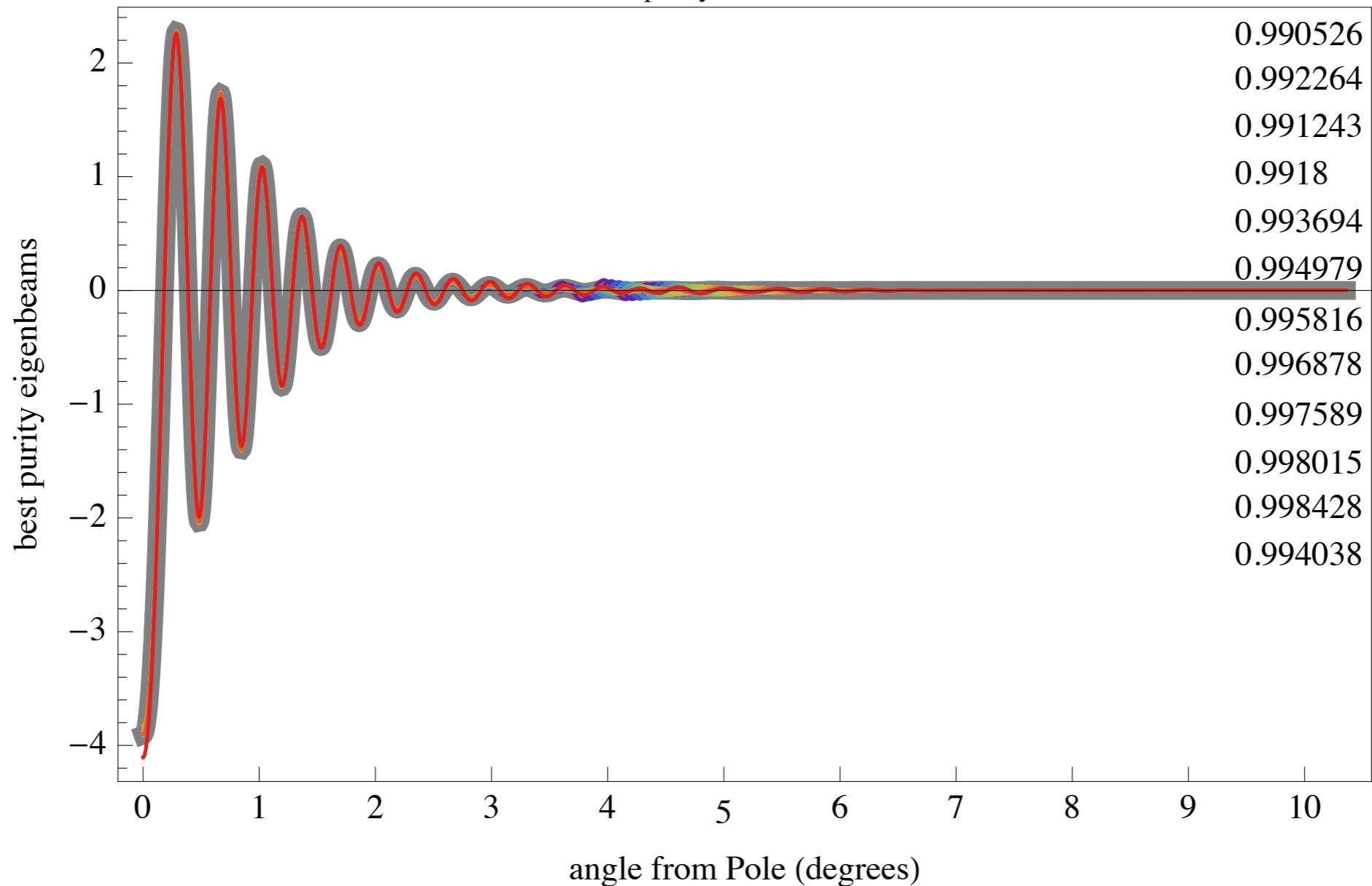
$\#_{\text{dish}} = 16$     $\#_{\text{split}} = 0$     $\nu \in [700, 800] \text{ MHz}$    spaced 630 cm



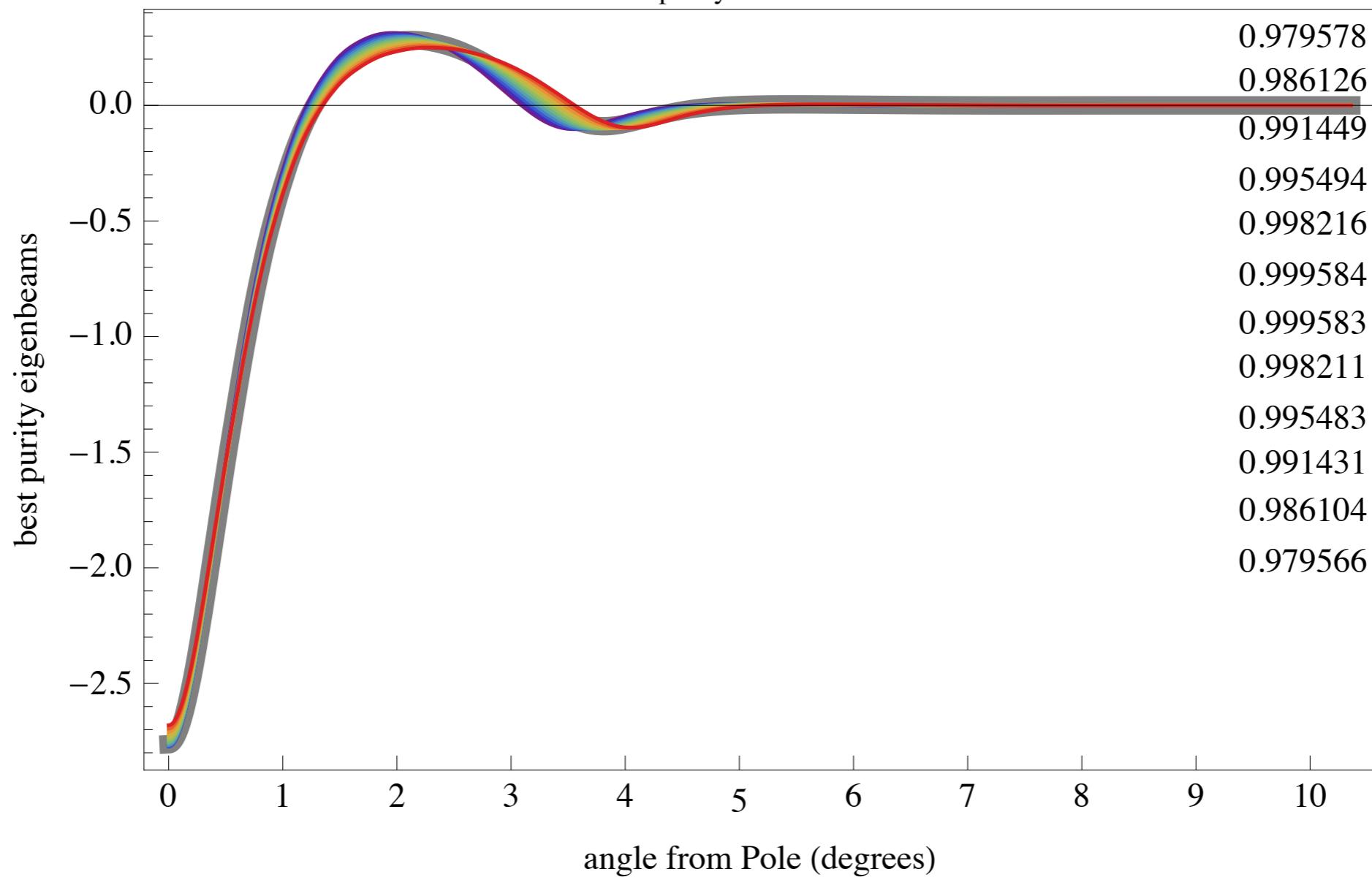
$m = 0$   $\#_{\text{beams}} = 15$   $i_{\text{purity}} = 1$  mean purity = 0.998885



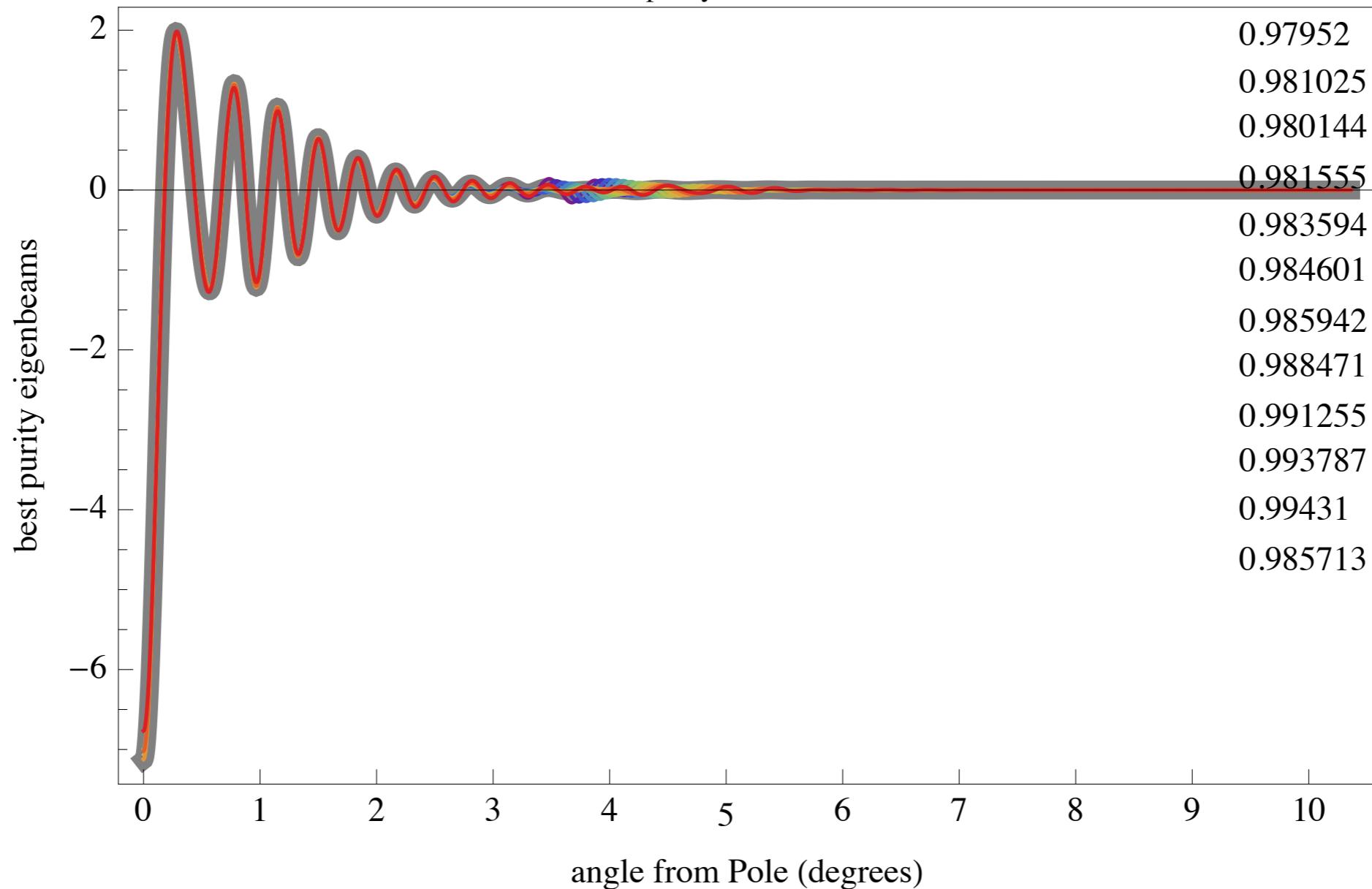
$m = 0$     $\#\text{beams} = 15$     $i_{\text{purity}} = 2$    mean purity = 0.994606



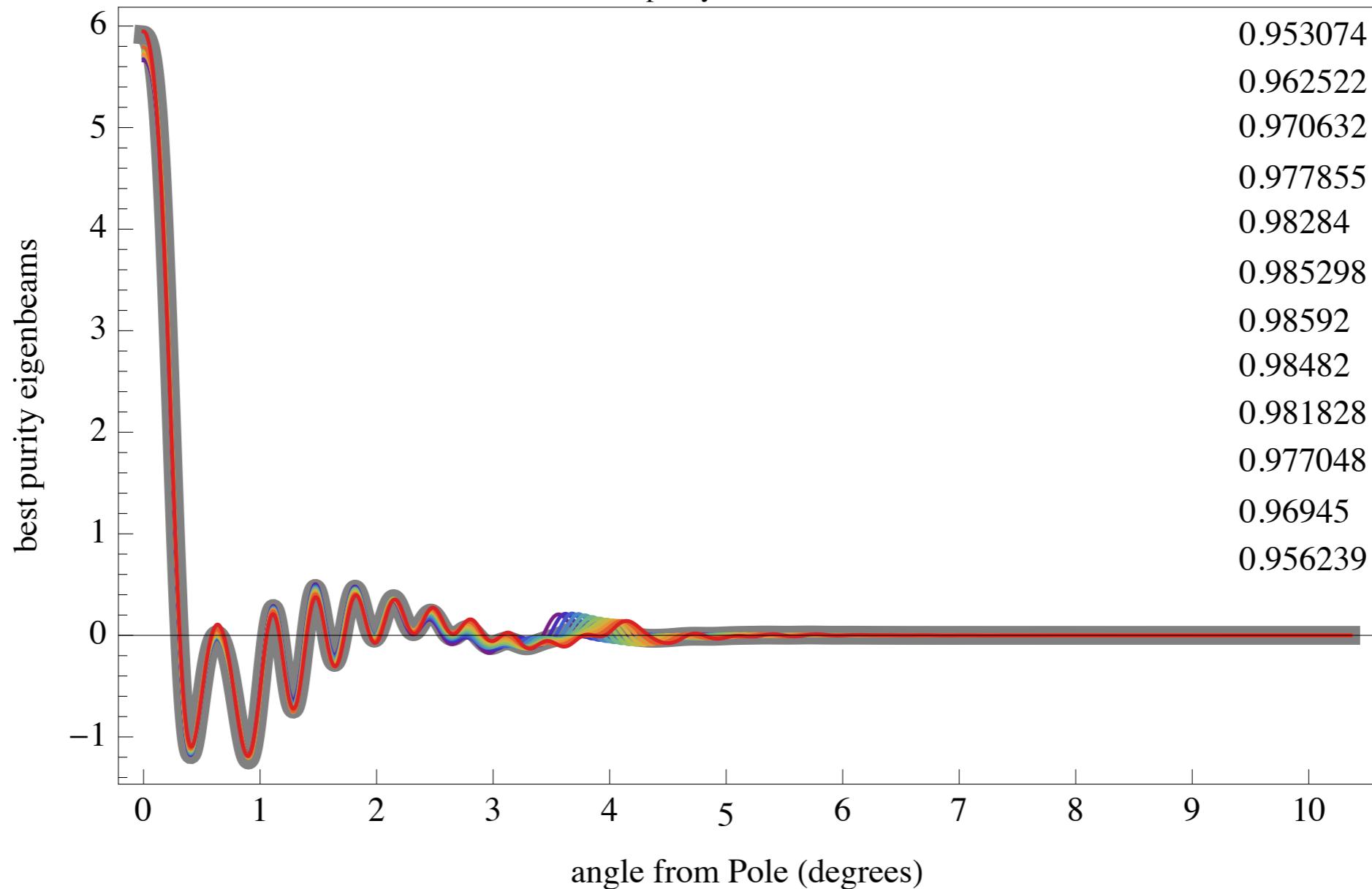
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 3$  mean purity = 0.991735



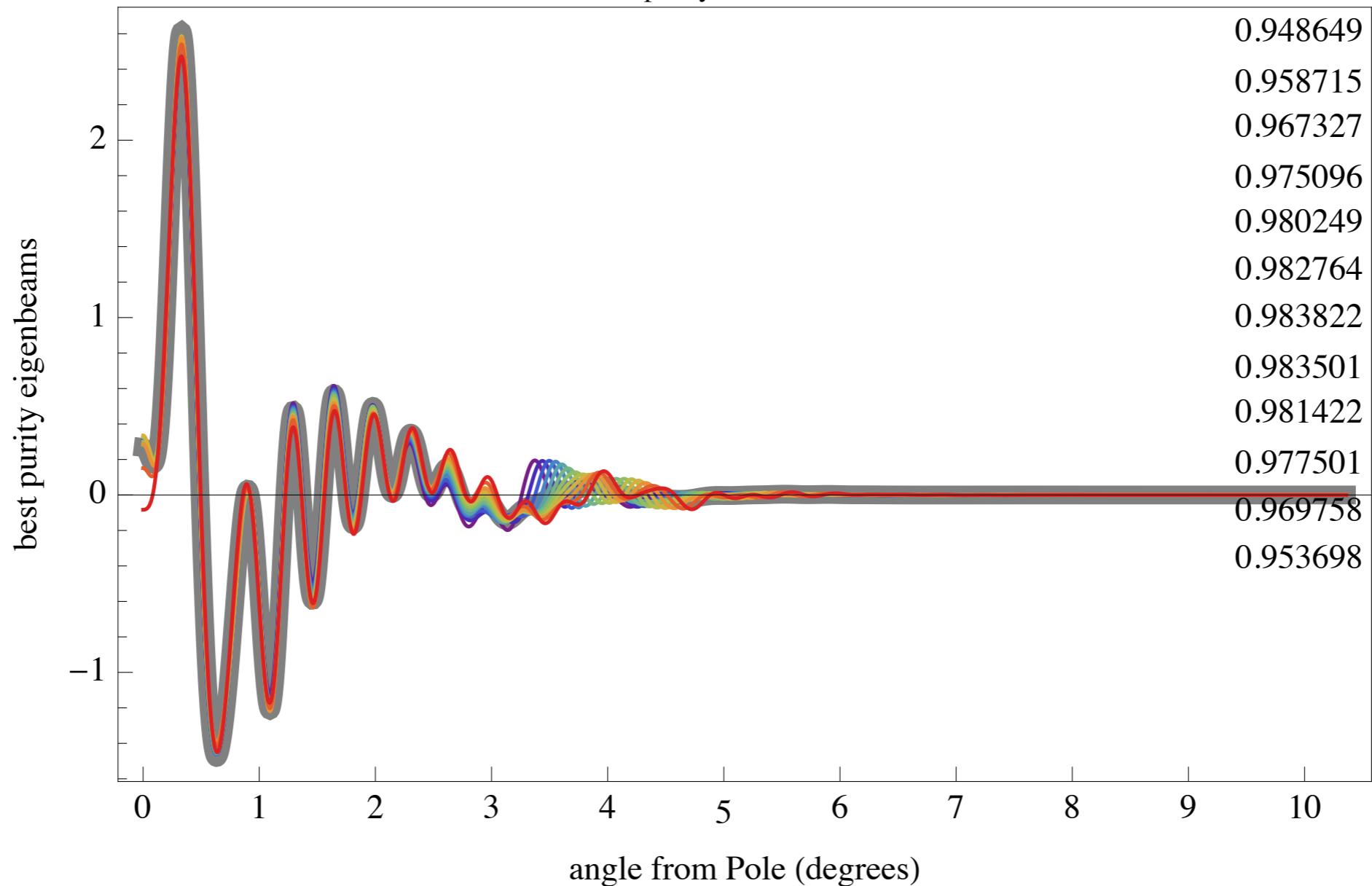
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 4$  mean purity = 0.985826



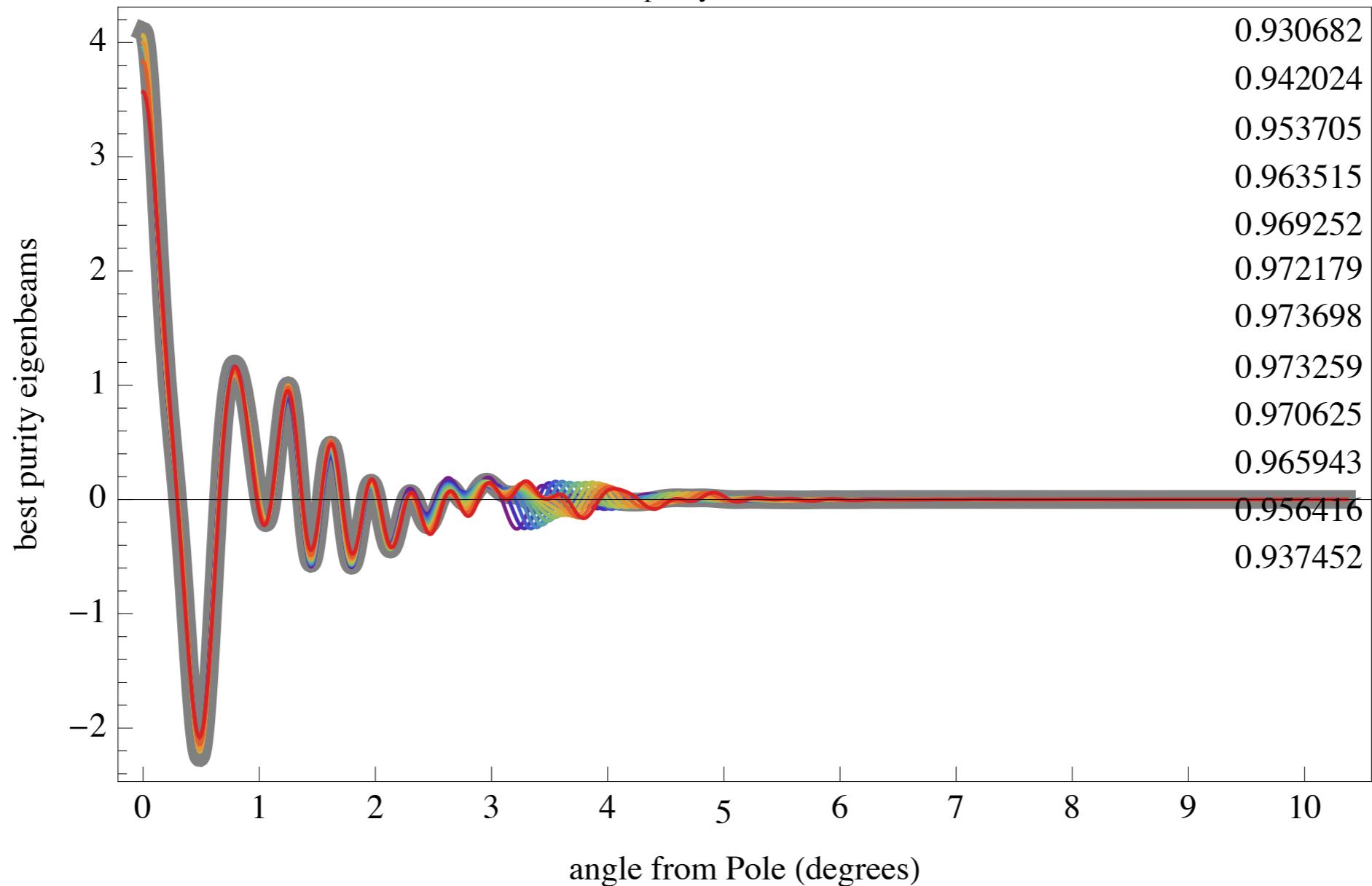
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 5$  mean purity = 0.973961



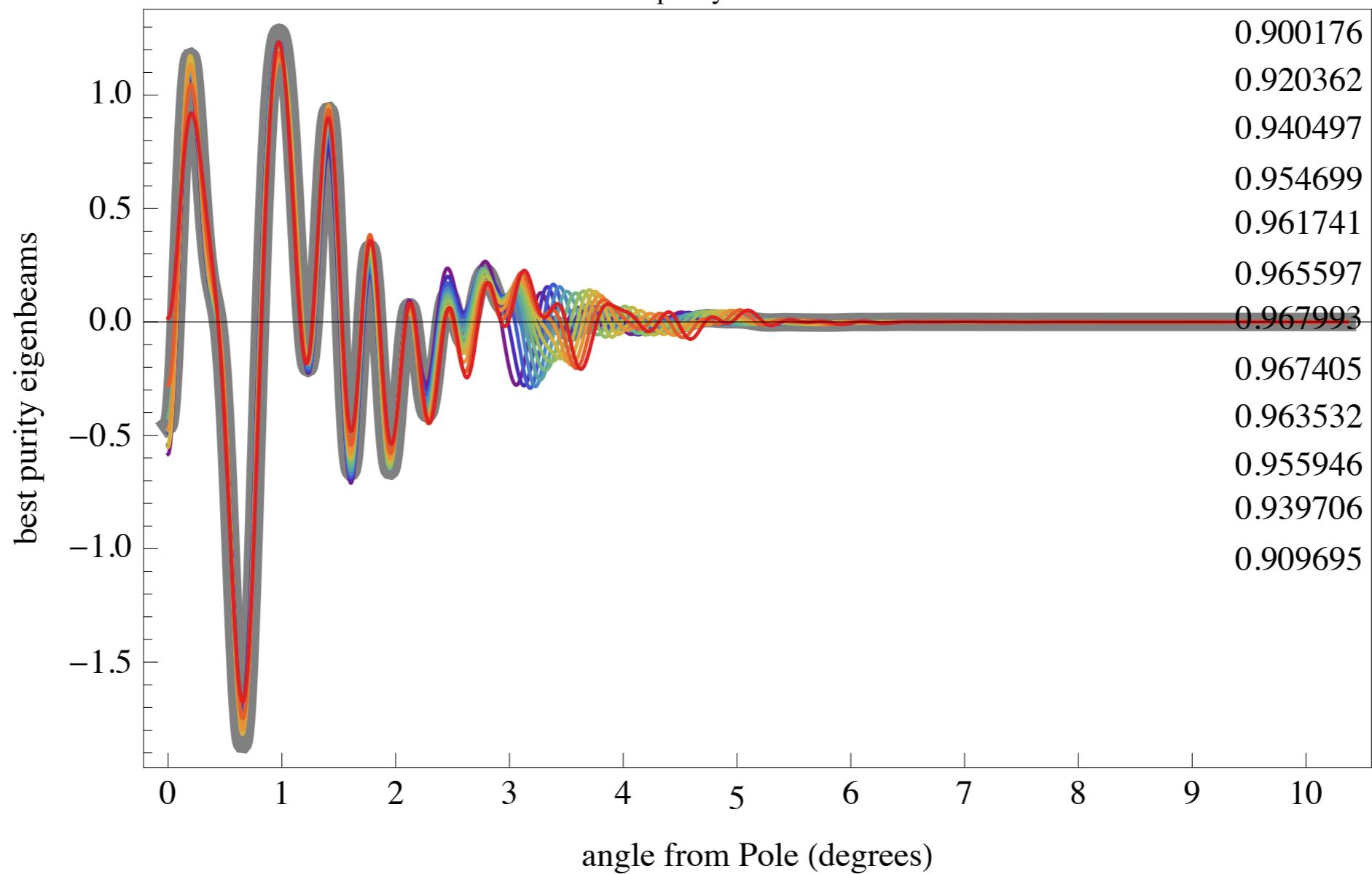
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 6$  mean purity = 0.971875



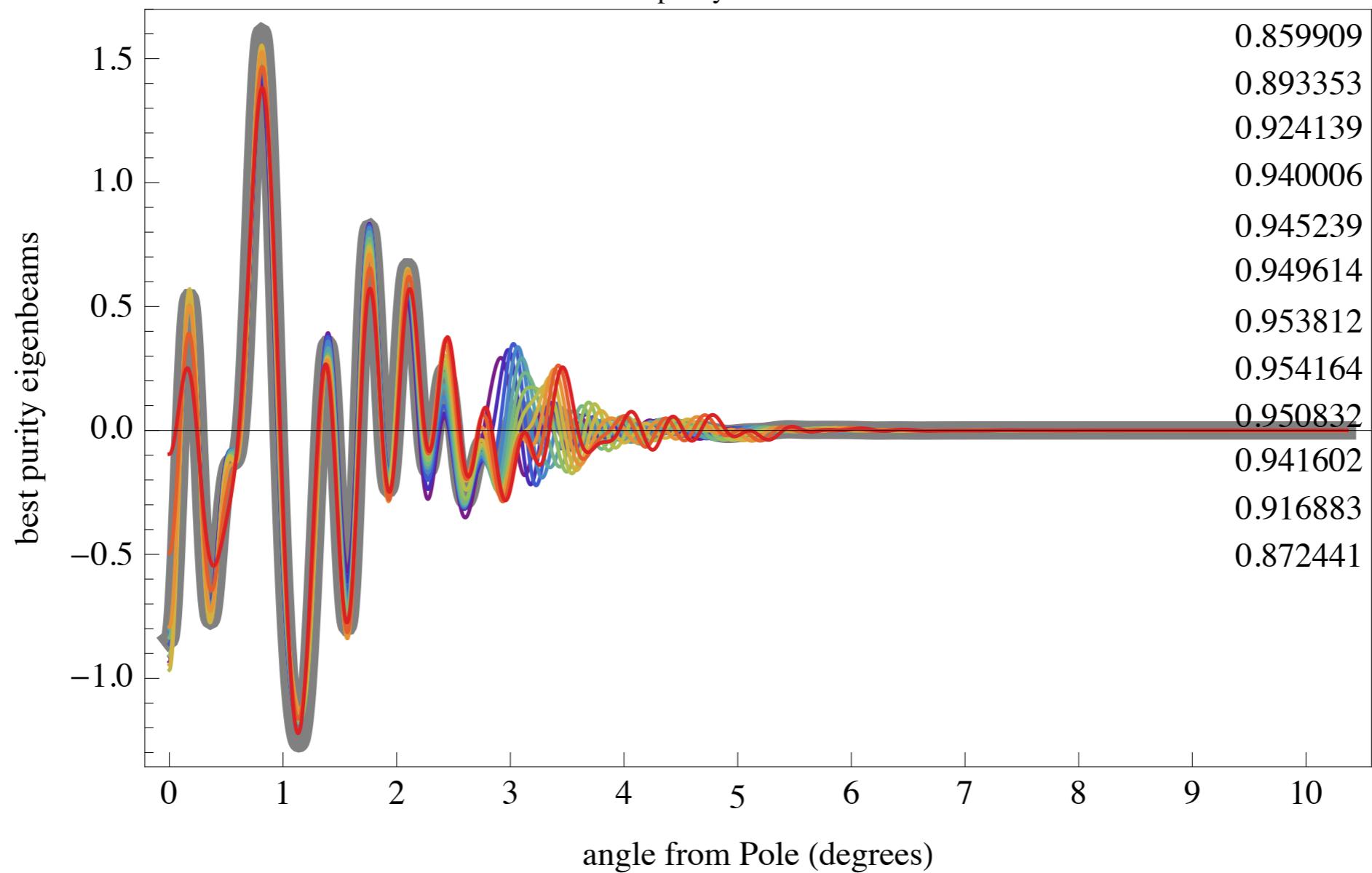
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 7$  mean purity = 0.959062



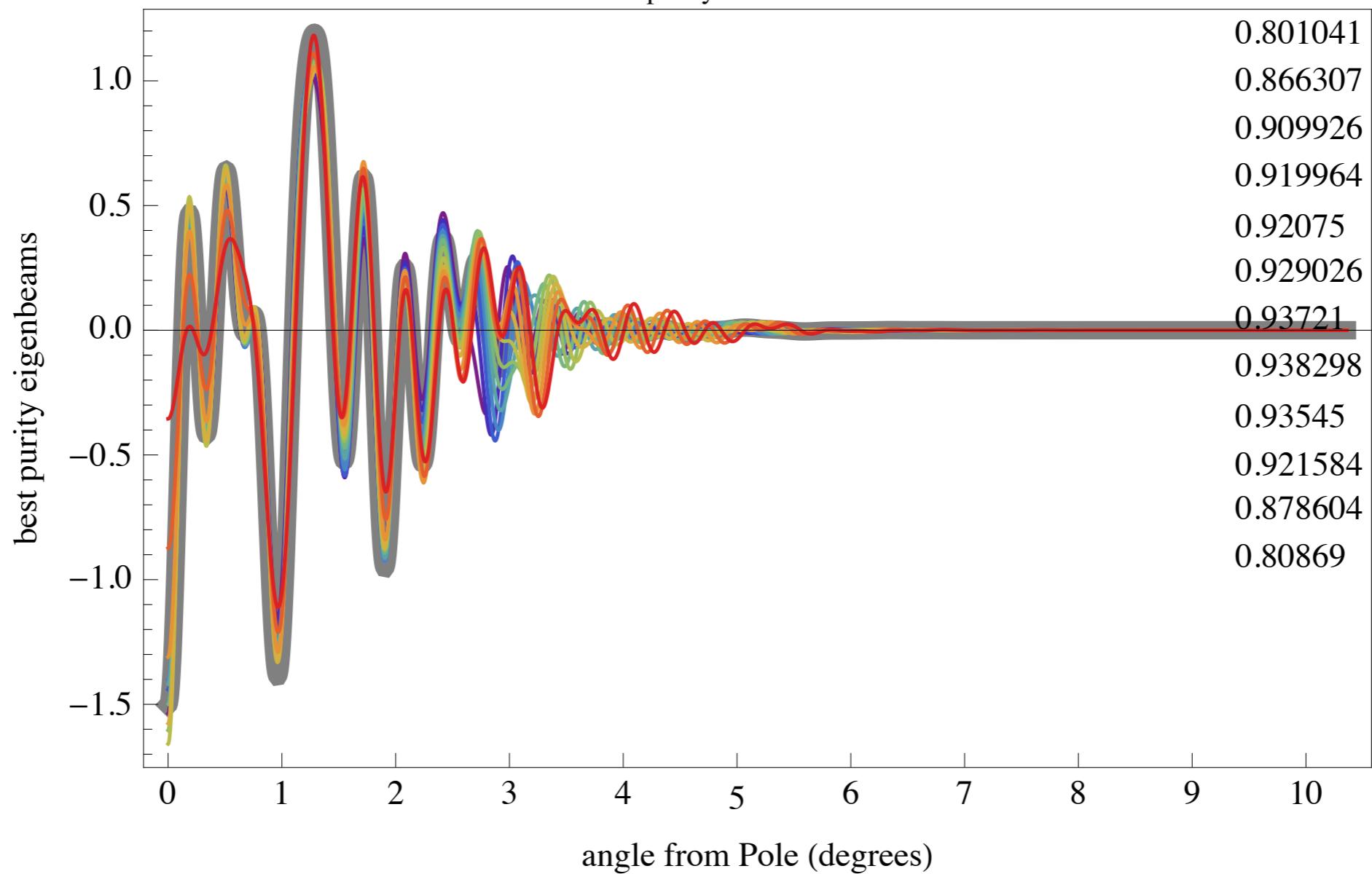
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 8$  mean purity = 0.945612



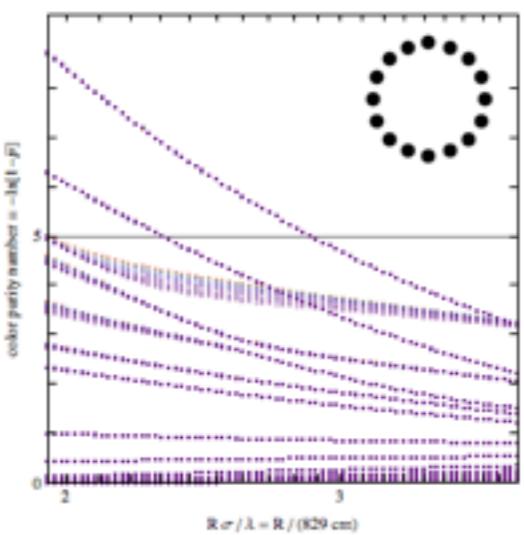
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 9$  mean purity = 0.925166



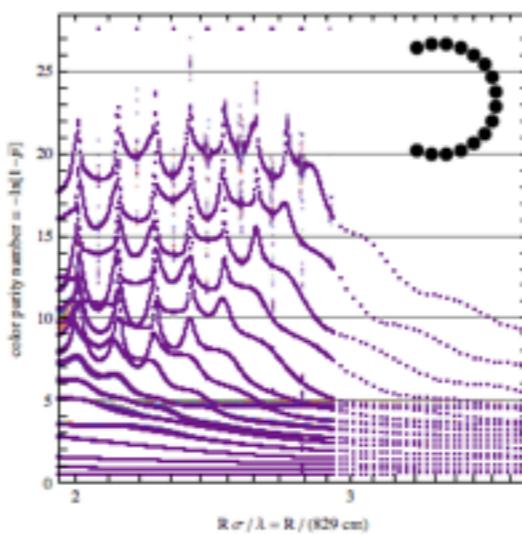
$m = 0$   $\#_{\text{beams}} = 15$   $i_{\text{purity}} = 10$  mean purity = 0.897237



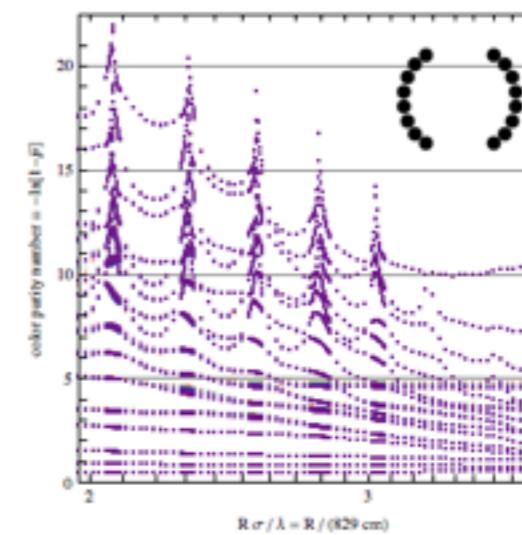
$\Pi_{\text{disk}} = 16$   $\Pi_{\text{split}} = 0$   $\nu \in [700, 800]$  MHz spaced 630 cm



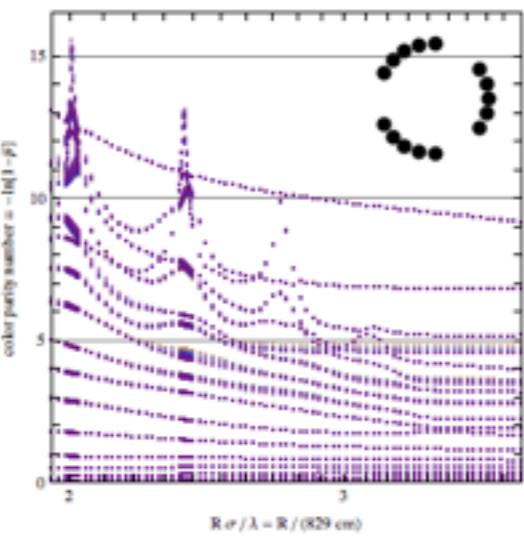
$\Pi_{\text{disk}} = 16$   $\Pi_{\text{split}} = 1$   $\nu \in [700, 800]$  MHz spaced 630 cm



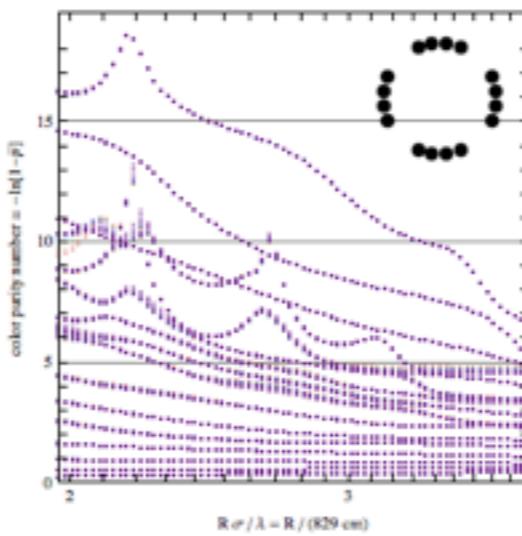
$\Pi_{\text{disk}} = 16$   $\Pi_{\text{split}} = 2$   $\nu \in [700, 800]$  MHz spaced 630 cm



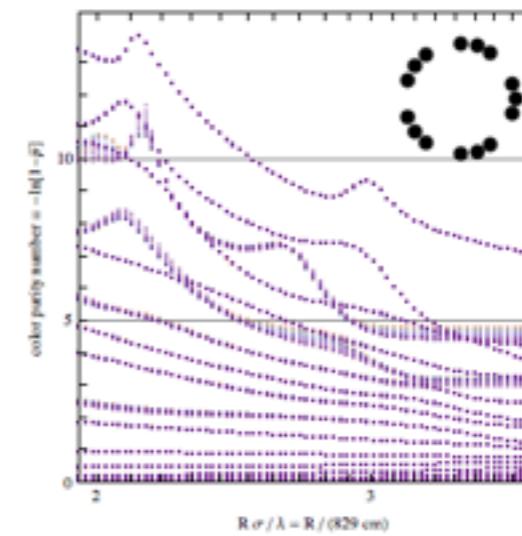
$\Pi_{\text{disk}} = 15$   $\Pi_{\text{split}} = 3$   $\nu \in [700, 800]$  MHz spaced 630 cm



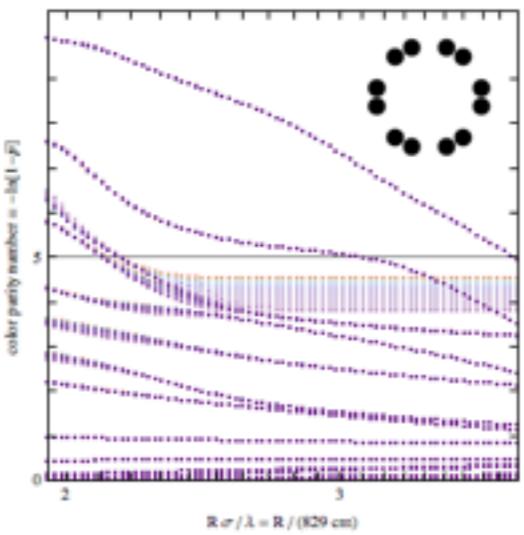
$\Pi_{\text{disk}} = 16$   $\Pi_{\text{split}} = 4$   $\nu \in [700, 800]$  MHz spaced 630 cm



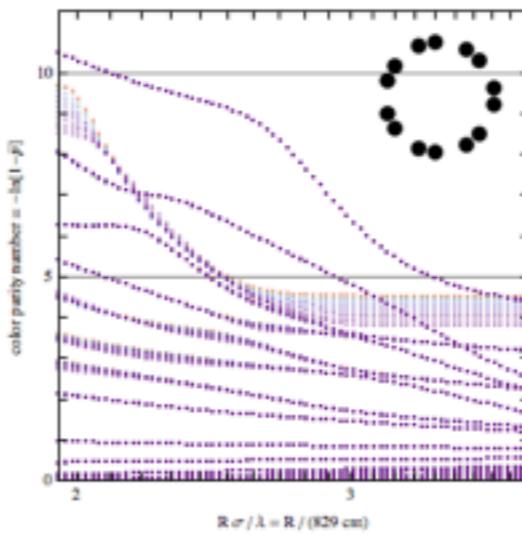
$\Pi_{\text{disk}} = 15$   $\Pi_{\text{split}} = 5$   $\nu \in [700, 800]$  MHz spaced 630 cm



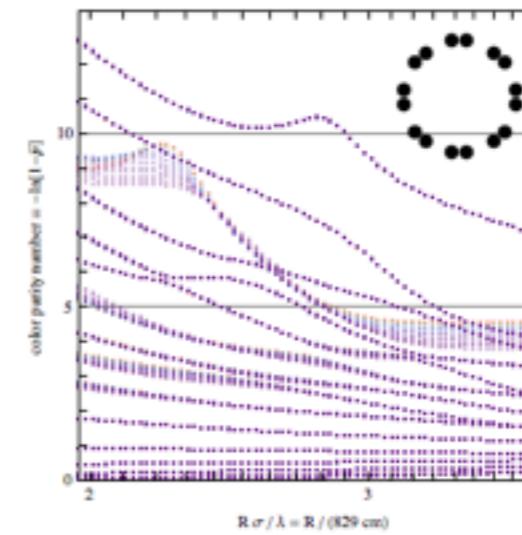
$\Pi_{\text{disk}} = 12$   $\Pi_{\text{split}} = 6$   $\nu \in [700, 800]$  MHz spaced 630 cm



$\Pi_{\text{disk}} = 14$   $\Pi_{\text{split}} = 7$   $\nu \in [700, 800]$  MHz spaced 630 cm

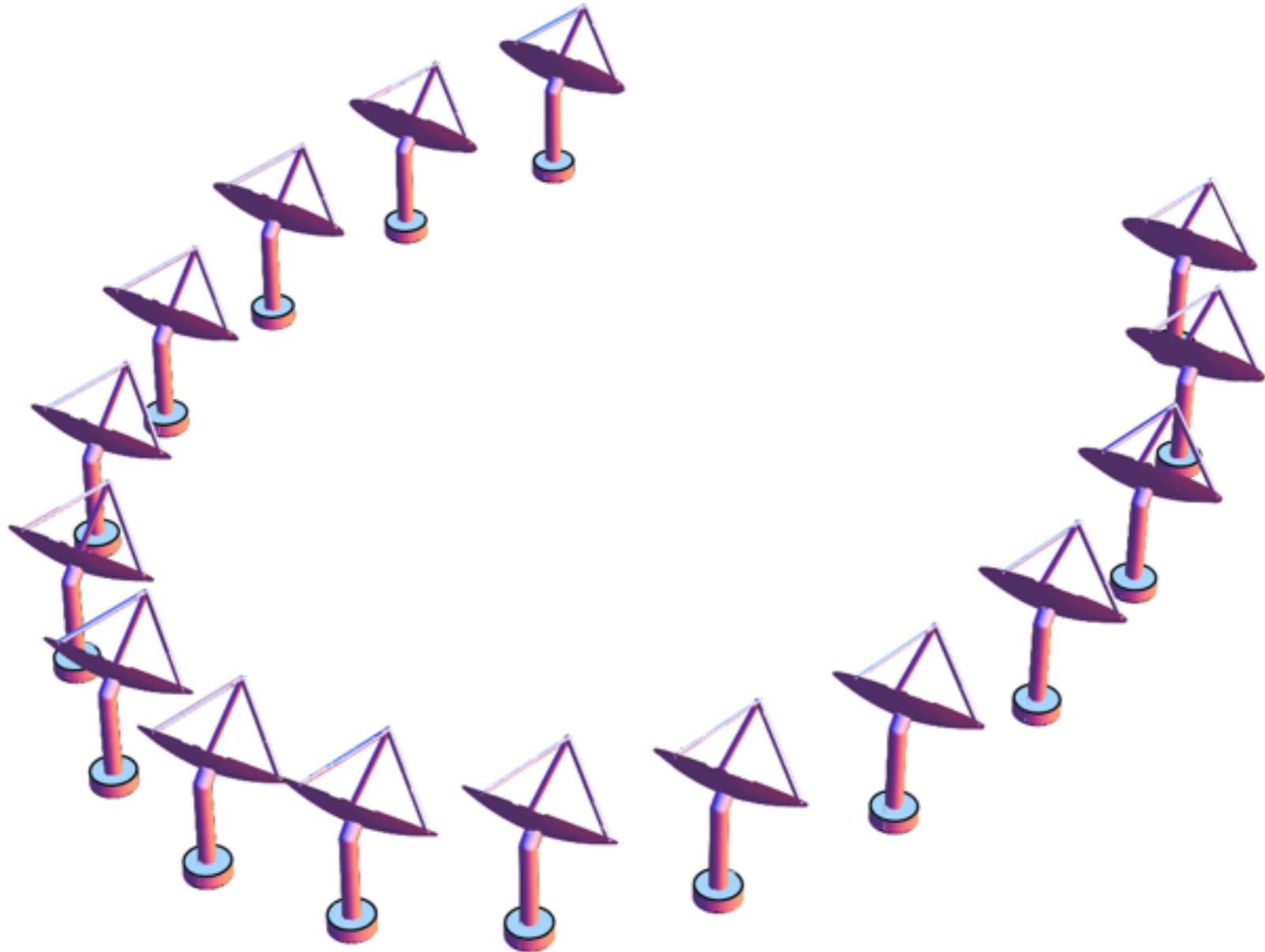


$\Pi_{\text{disk}} = 16$   $\Pi_{\text{split}} = 8$   $\nu \in [700, 800]$  MHz spaced 630 cm

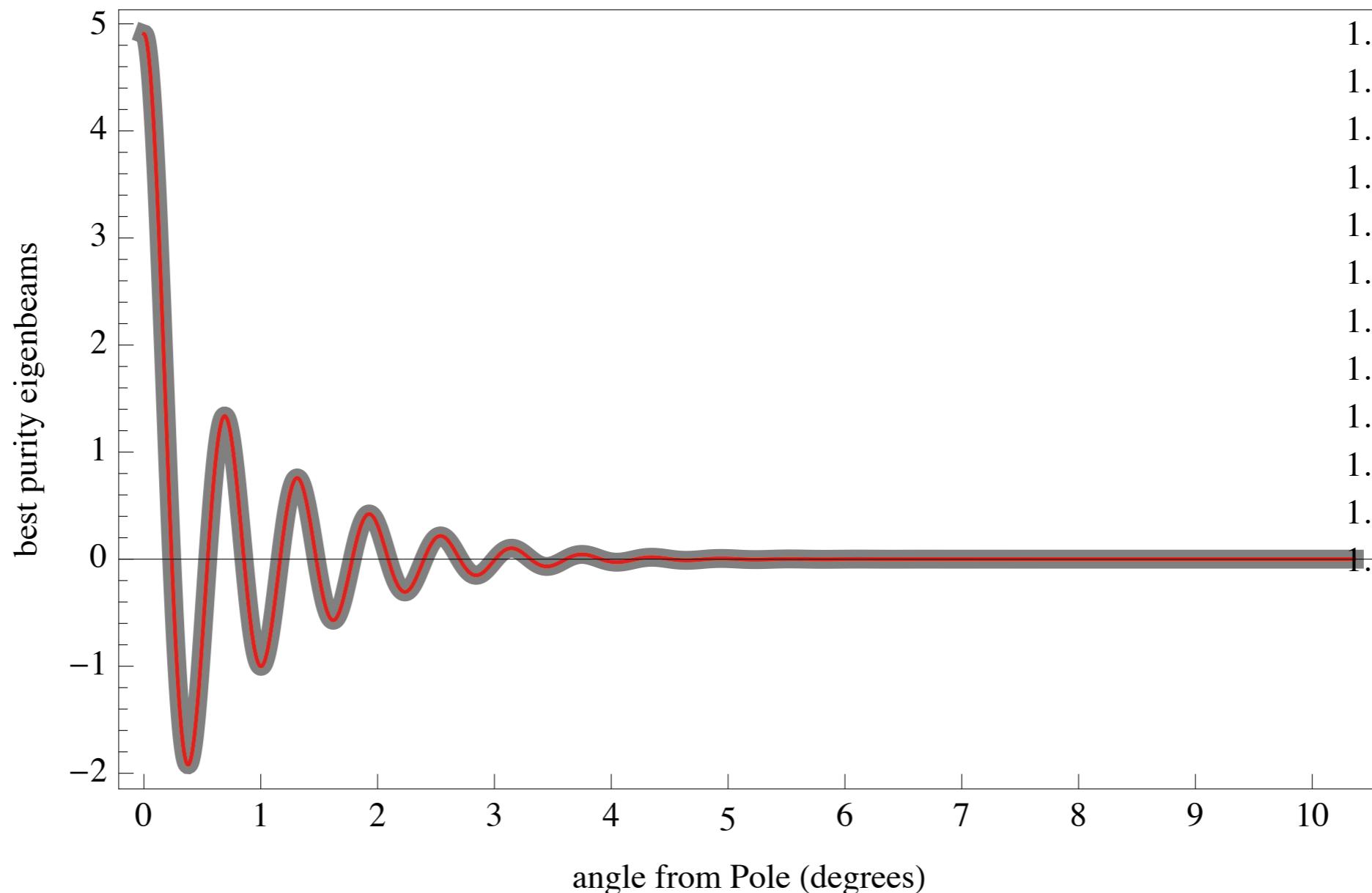


# A Very Pure Polarscope

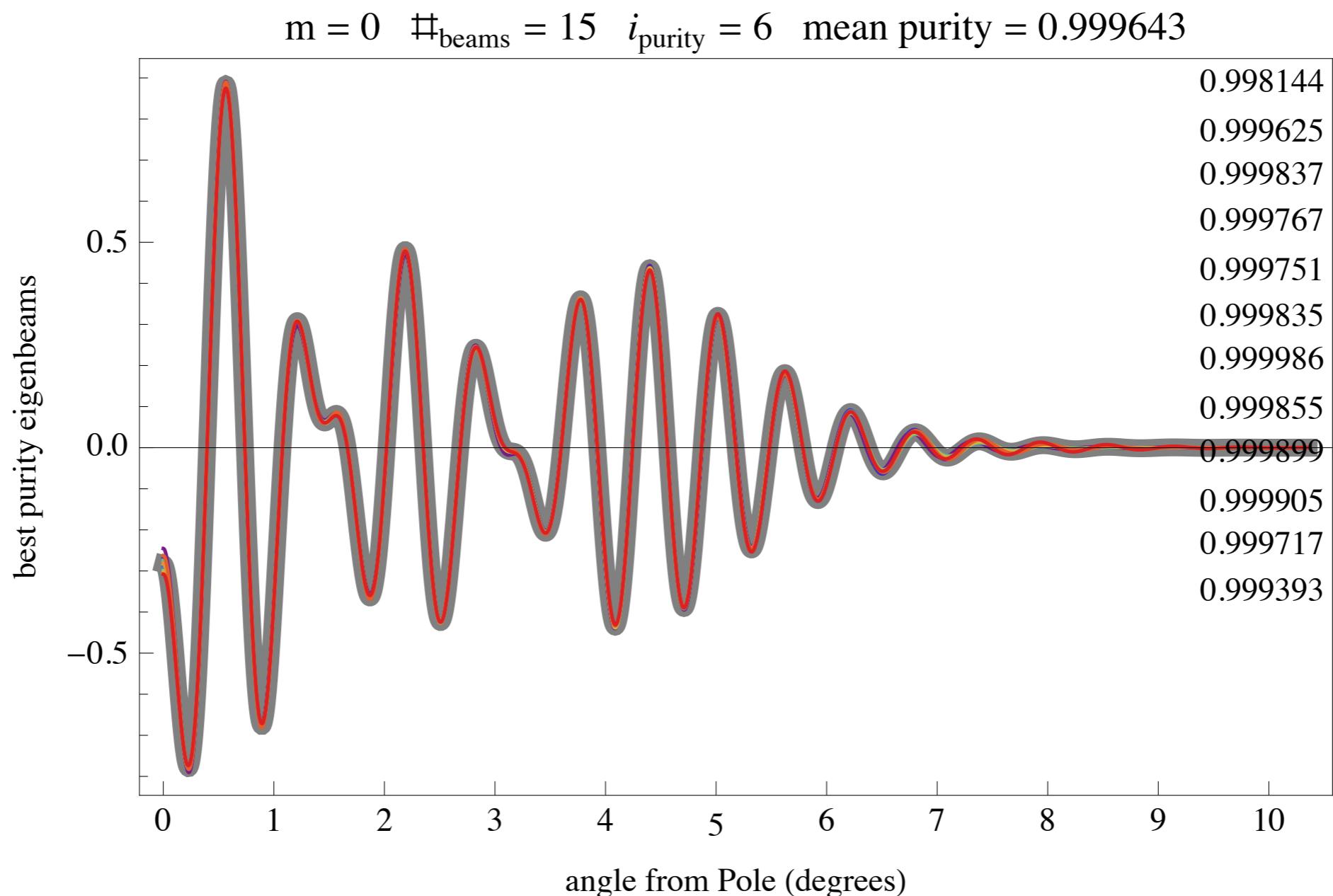
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$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 1$  mean purity = 1.

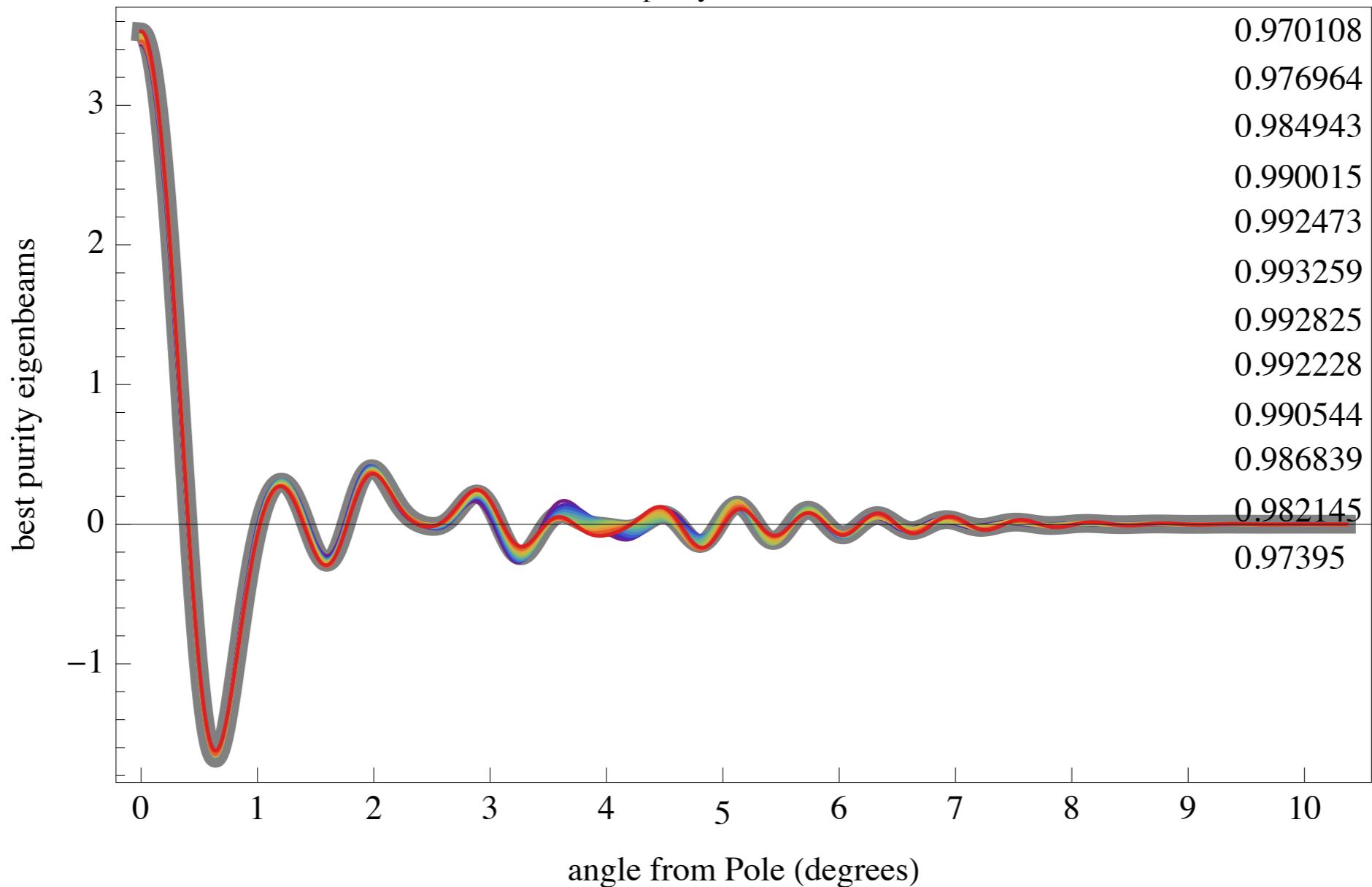


# Skip to 6th purity eigenmode

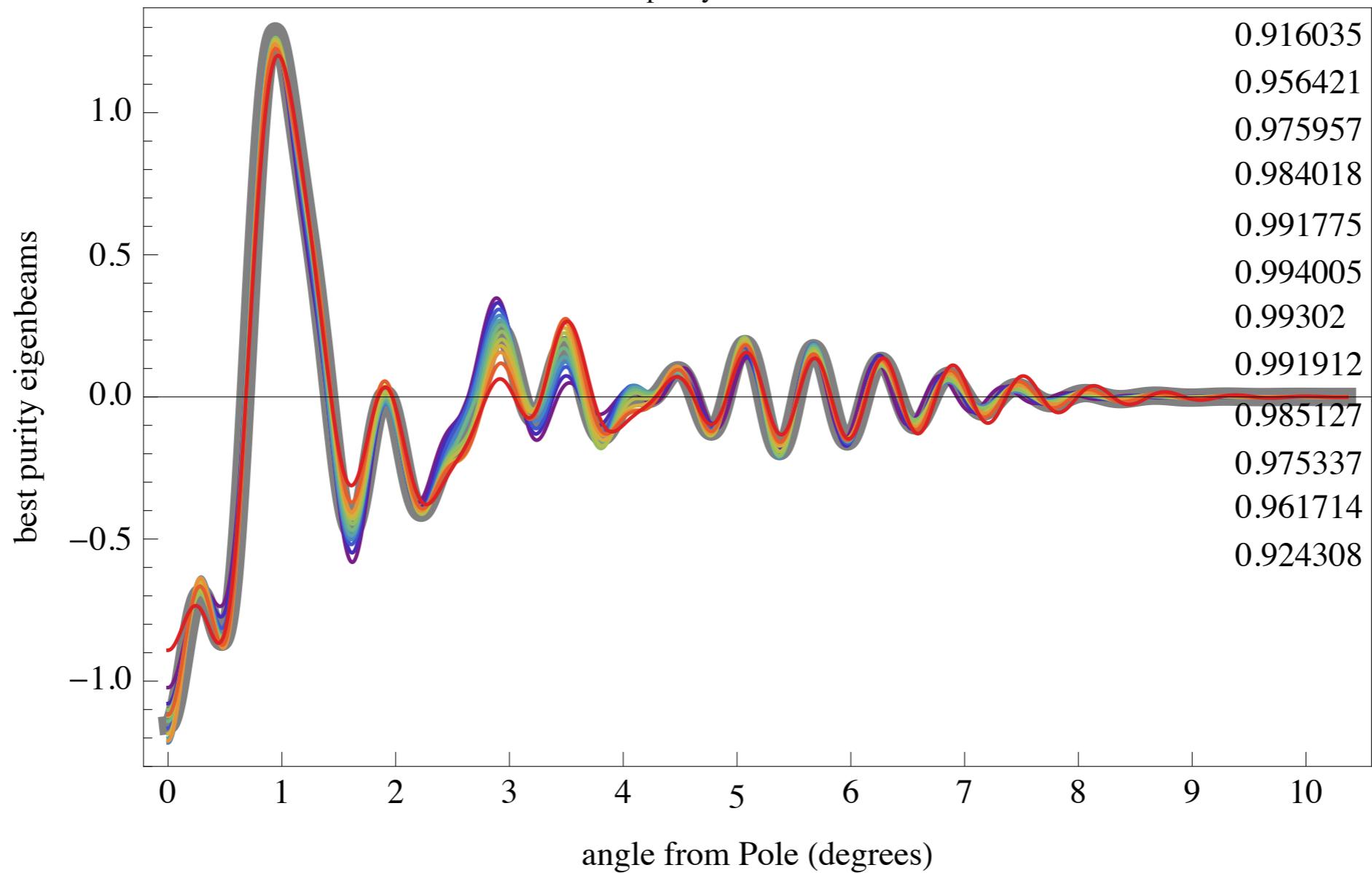


# Skip to 9th purity eigenmode

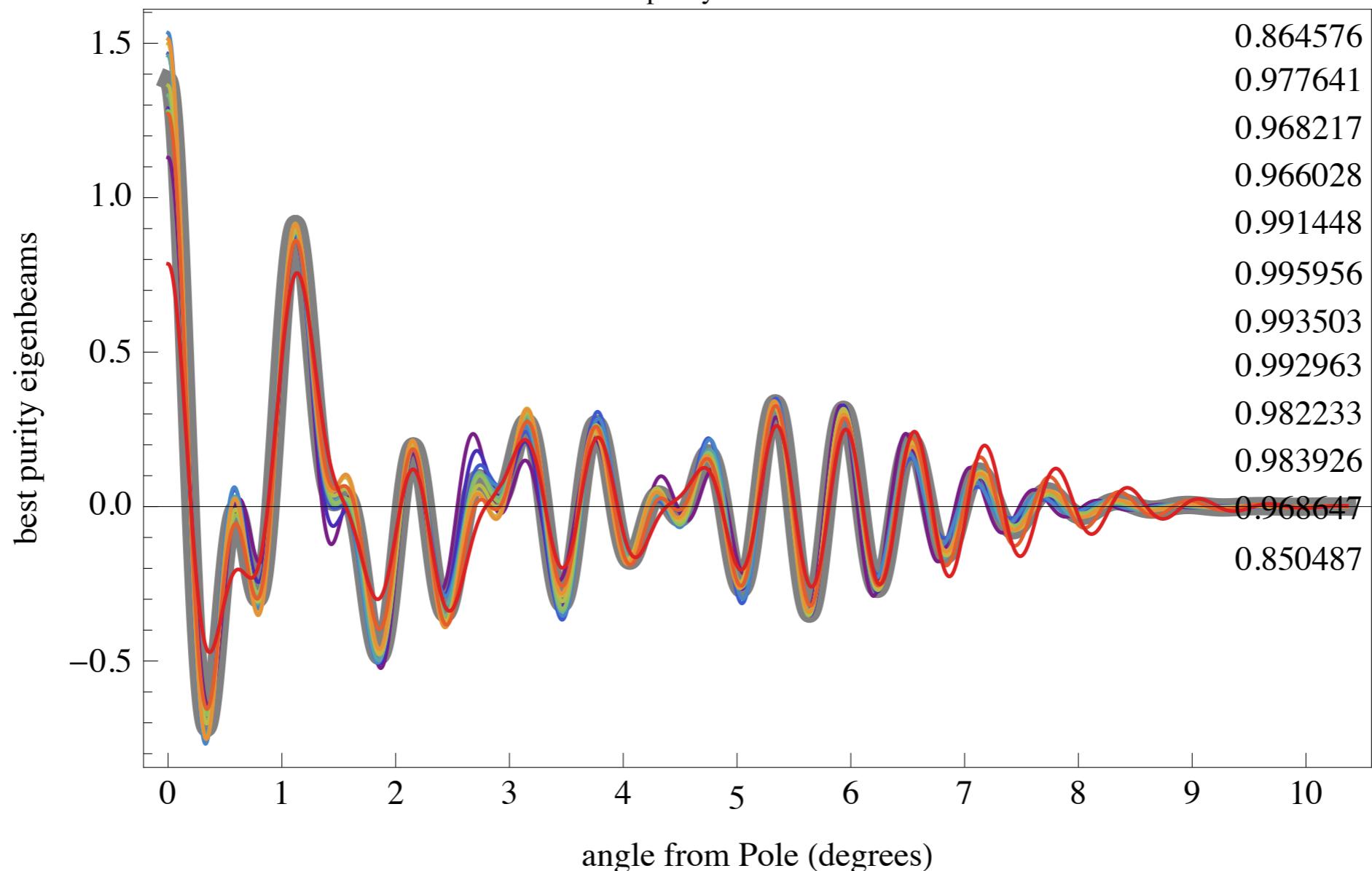
$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 9$  mean purity = 0.985524



$m = 0$   $\#\text{beams} = 15$   $i_{\text{purity}} = 10$  mean purity = 0.970802



$m = 0$   $\#_{\text{beams}} = 15$   $i_{\text{purity}} = 11$  mean purity = 0.961302



# Suggestion for Next Tianlai-16 Configuration

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The Tianlai-16 dish array allows us the opportunity to

- by pointing toward the NCP (a polarscope or close to one) to integrate down rapidly to low noise levels
- experiment with array configuration to try to demonstrate a high purity interferometer
- in so doing rapidly test the efficacy of smooth spectrum foreground subtraction
- look for more concrete evidence of non-smooth foreground components

These results will have important implications for Tianlai cylinder and 21cm intensity mapping more generally.

# Unfortunately Other Surveys Have Avoided NCP

