

Higgs EFT and kinematic distributions

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based on 1602.05202
(with Johann Brehmer and Tilman Plehn)

Higgs Hunting in Paris, September 1, 2016

Higgs effective field theory

- New physics at $\Lambda \gg E_{\text{LHC}} \sim v?$ [Grzadkowski et al 1008.4884; ...]

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \underbrace{\sum_i^{59} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)}}_{\text{e.g. } \mathcal{O}_{WW} = -\frac{g^2}{4} (\phi^\dagger \phi) W_{\mu\nu}^a W^{\mu\nu a}, \mathcal{O}_W = \frac{ig}{2} (D^\mu \phi)^\dagger \sigma^k (D^\nu \phi) W_{\mu\nu}^k} + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$

- Framework for indirect searches at the electroweak scale
- reproducible and (mostly) model independent

EFT **cannot** describe LHC Higgs physics

- LHC accuracy $\sim 10\%$ translates to new physics reach of:

$$\left| \frac{\sigma \times \text{BR}}{(\sigma \times \text{BR})_{\text{SM}}} - 1 \right| = \frac{g^2 m_h^2}{\Lambda^2} \gtrsim 10\% \quad \Leftrightarrow \quad \Lambda < \frac{g m_h}{\sqrt{10\%}} \approx 400 \text{ GeV}$$

- scenarios with $\Lambda \gg E$ not measurable at the LHC
- D8 not sufficiently suppressed

Johann Brehmer, Ayres Freitas, David Lopez-Val, Tilman Plehn
[1510.03443]

EFT can describe LHC Higgs physics

→ answer some remaining questions

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D6 can describe LHC Higgs physics

→ answer some remaining questions

Outline

D6 description

- To square or not to square dimension-6 amplitudes?
- Vector triplet model
- Higgs-strahlung and WBF
- Which observable to study for WBF?

D6 description - to square or not to square?

$$|\mathcal{M}_{4+6}|^2 = |\mathcal{M}_4|^2 + 2 \operatorname{Re} \mathcal{M}_4^* \mathcal{M}_6 + |\mathcal{M}_6|^2$$

D6 description - to square or not to square?

$$|\mathcal{M}_{4+6}|^2 = |\mathcal{M}_4|^2 + 2 \operatorname{Re} \mathcal{M}_4^* \mathcal{M}_6 + |\mathcal{M}_6|^2$$

Preferable to include $D6^2$ when neglecting $D8$?



Study for vector triplet model

NOT a consistent EFT → **practical** question

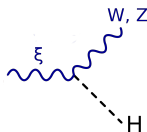
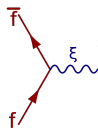
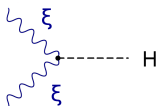
Vector triplet model

Full model

$$\begin{aligned}
 \mathcal{L} \supset & -\frac{1}{4} V_{\mu\nu}^a V^{\mu\nu a} + \frac{M_V^2}{2} V_\mu^a V^{\mu a} \\
 & + i \frac{g_V}{2} c_H V_\mu^a \left[\phi^\dagger \sigma^a \overleftrightarrow{D}^\mu \phi \right] \\
 & + \frac{g_w^2}{2g_V} V_\mu^a \sum_{\text{fermions}} c_F \bar{F}_L \gamma^\mu \sigma^a F_L \\
 & + g_V^2 c_{VVHH} V_\mu^a V^{\mu a} (\phi^\dagger \phi)
 \end{aligned}$$

D6 approximation

$$\begin{aligned}
 \mathcal{O}_{WW} &= -\frac{g^2}{4} (\phi^\dagger \phi) W_{\mu\nu}^a W^{\mu\nu a} \\
 \mathcal{O}_W &= \frac{ig}{2} (D^\mu \phi)^\dagger \sigma^k (D^\nu \phi) W_{\mu\nu}^k \\
 \Lambda &= m_\xi
 \end{aligned}$$

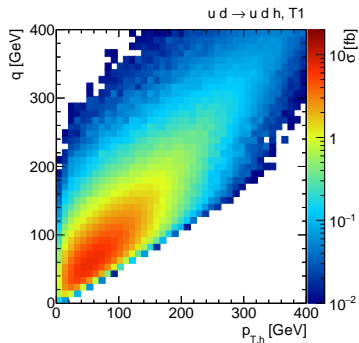
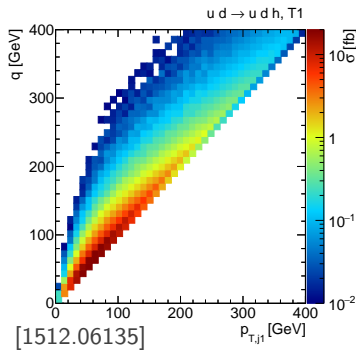
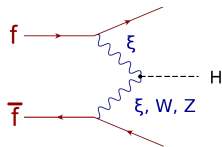


new ξ resonance - modification of Higgs couplings

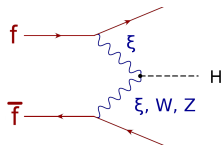
[1510.03443; 1211.2229; 1406.7320; 1506.03631]

WBF - momentum transfer

- study parton-level process $ud \rightarrow u'd'h$
- momentum transfer q



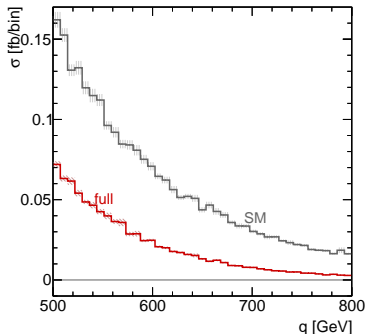
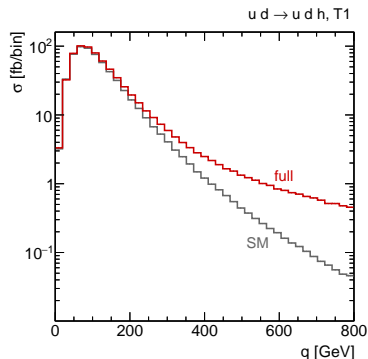
WBF - parton level



- study momentum transfer q (p_{T,j_1})
- two benchmarks with $m_\xi = 1200$ GeV

constructive

destructive

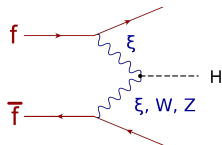


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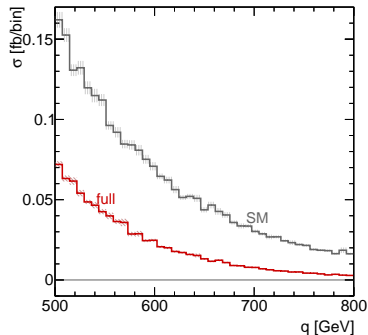
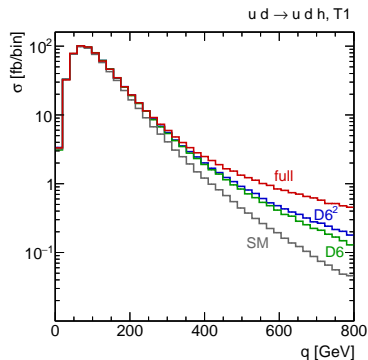
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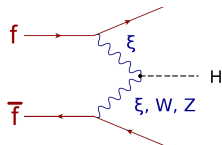


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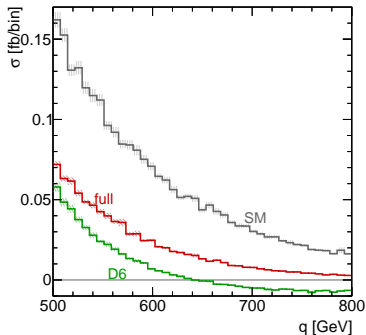
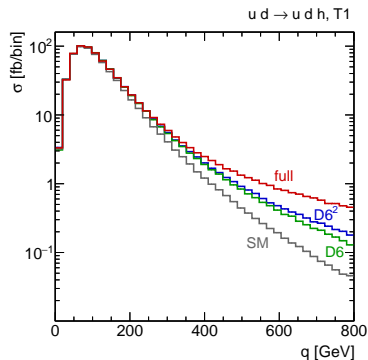
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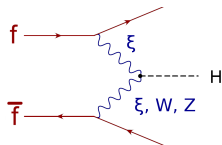


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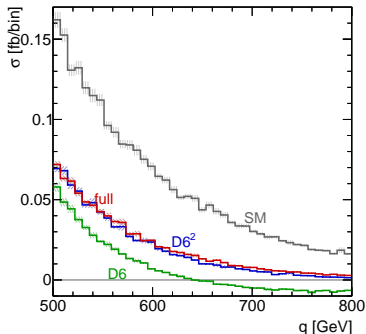
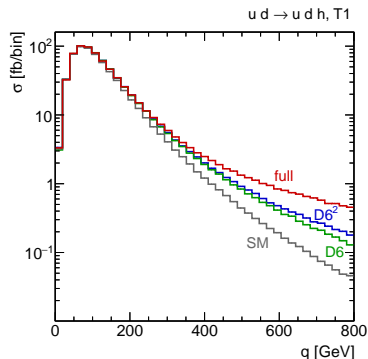
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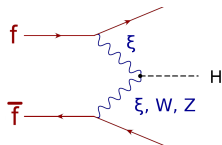
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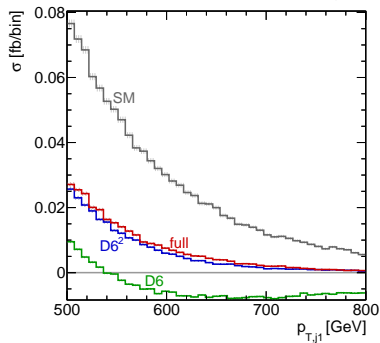
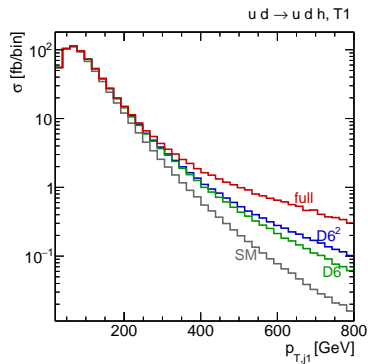
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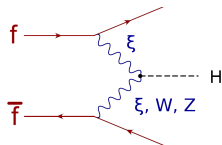
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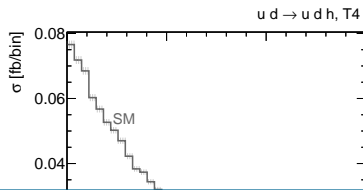
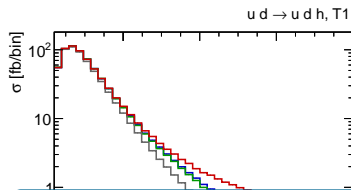
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Reality check

- Valid for larger parameter range?
- Valid for full, hadron level process?

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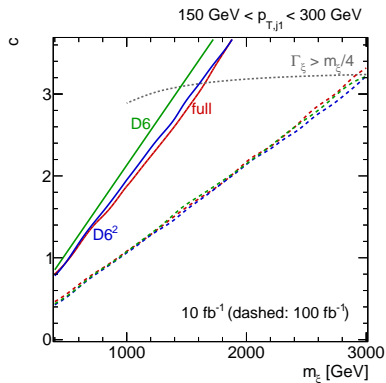
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WBF - Comparison of expected exclusion limits

- choose universal coupling rescaling c

$$g_V = 1, \quad c_H = c, \quad c_F = \frac{g_V^2}{2g^2} c, \quad c_{HHVV} = c^2$$

- m_ξ mass of the new heavy vector

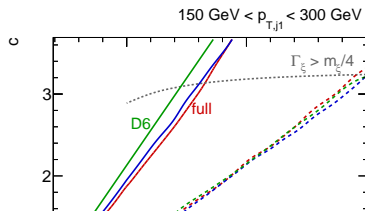


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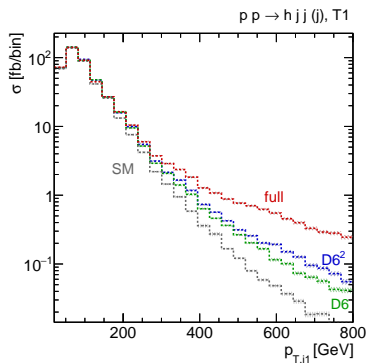
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WBF - getting realistic

- hadron level analysis $pp \rightarrow h jj (+j)$ using PYTHIA6 and FastJet
- apply WBF cuts

$$p_{T,j} > 20 \text{ GeV}, \quad m_{jj} > 500 \text{ GeV}, \quad \Delta\eta_{jj} > 3.6$$

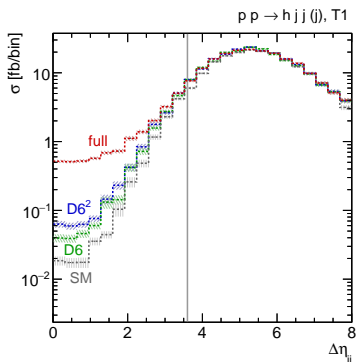
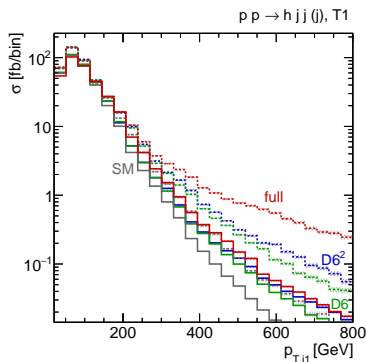


dotted: WBF diagrams only, without $\Delta\eta_{jj}$ cut

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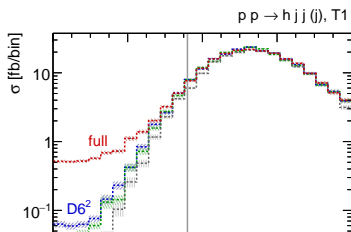
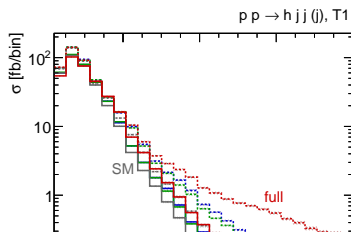


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Conclusions

- Including $D6^2$ terms improves agreement with full model and avoids negative cross sections
- Leading tagging jet p_T highly correlated with momentum transfer q for WBF
- Results survive in a realistic environment

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Thank you for your attention!
Any questions?

D6 operators

HISZ basis	
$\mathcal{O}_{\phi 1} = (D_\mu \phi)^\dagger (\phi \phi^\dagger) (D^\mu \phi)$	$\mathcal{O}_{\phi 2} = \frac{1}{2} \partial^\mu (\phi^\dagger \phi) \partial_\mu (\phi^\dagger \phi)$
$\mathcal{O}_{\phi 3} = \frac{1}{3} (\phi^\dagger \phi)^3$	
$\mathcal{O}_{GG} = (\phi^\dagger \phi) G_{\mu\nu}^A G^{\mu\nu A}$	$\mathcal{O}_{BW} = -\frac{g g'}{4} (\phi^\dagger \sigma^k \phi) B_{\mu\nu} W^{\mu\nu k}$
$\mathcal{O}_{BB} = -\frac{g'^2}{4} (\phi^\dagger \phi) B_{\mu\nu} B^{\mu\nu}$	$\mathcal{O}_{WW} = -\frac{g^2}{4} (\phi^\dagger \phi) W_{\mu\nu}^k W^{\mu\nu k}$
$\mathcal{O}_B = \frac{ig}{2} (D^\mu \phi^\dagger) (D^\nu \phi) B_{\mu\nu}$	$\mathcal{O}_W = \frac{ig}{2} (D^\mu \phi^\dagger) \sigma^k (D^\nu \phi) W_{\mu\nu}^k$

Table: Bosonic CP-conserving Higgs operators in the HISZ basis.

Wilson coefficients

$$f_{\phi 2} = \frac{3}{4} (-2 c_F g^2 + c_H g_V^2) ,$$

$$f_{\phi 3} = -3\lambda (-2 c_F g^2 + c_H g_V^2) ,$$

$$f_{f\phi} = -\frac{1}{4} y_f c_H (-2 c_F g^2 + c_H g_V^2) ,$$

$$f_{WW} = c_F c_H$$

$$f_{BW} = c_F c_H \equiv f_{WW}$$

$$f_W = -2 c_F c_H .$$

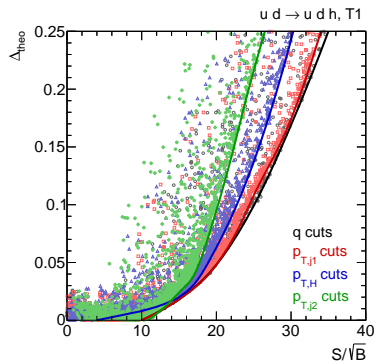
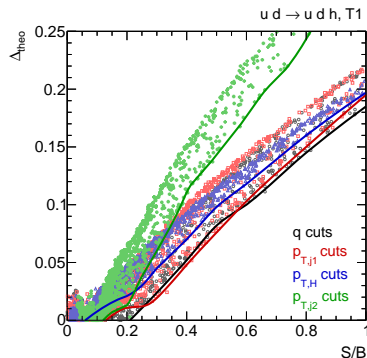
WBF - Which observable to study?

Compare deviations from the full model

$$\Delta_{\text{theo}}(x_{\min,\max}) = \left| \frac{\sigma_{\text{D6}} - \sigma_{\text{full}}}{\sigma_{\text{full}}} \right|, \quad x \in \{q, p_{T,j_1}, p_{T,j_2}, p_{T,h}\}$$

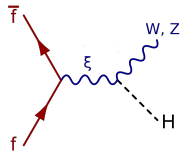
to statistics-driven and systematics-driven significances

$$\frac{S}{B}(x_{\min,\max}) = \left| \frac{\sigma_{\text{full}} - \sigma_{\text{SM}}}{\sigma_{\text{SM}}} \right| \quad \text{and} \quad \frac{S}{\sqrt{B}}(x_{\min,\max}) = \sqrt{L} \left| \frac{\sigma_{\text{full}} - \sigma_{\text{SM}}}{\sqrt{\sigma_{\text{SM}}}} \right|$$



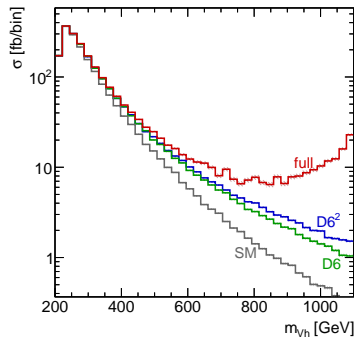
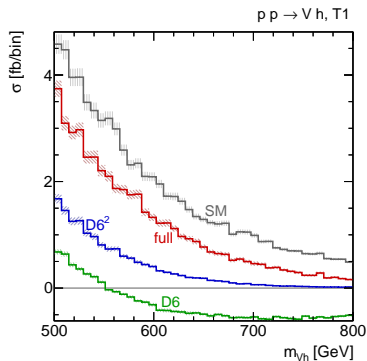
Higgs-strahlung

- study distribution of m_{Vh} ($p_{T,V}$)
- two benchmarks with $m_\xi = 1200$ GeV



destructive

constructive

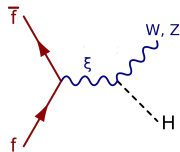


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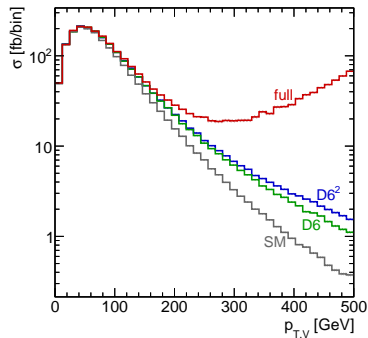
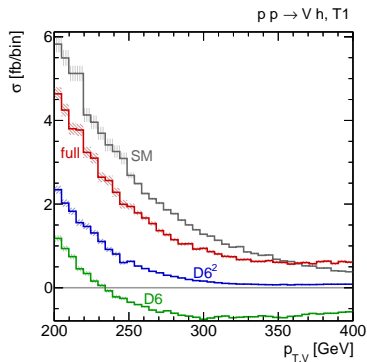
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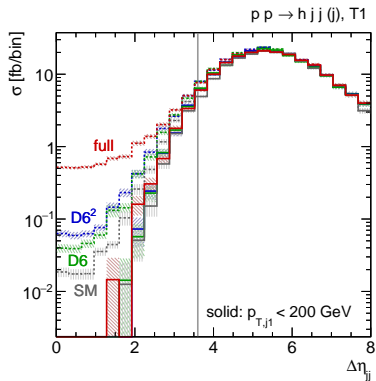
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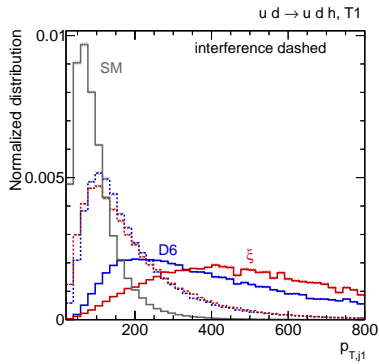
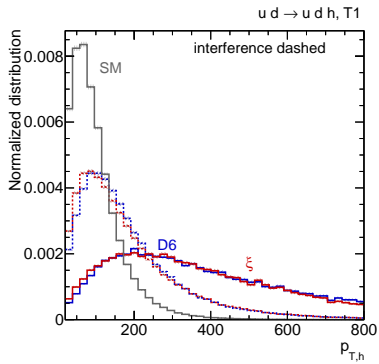


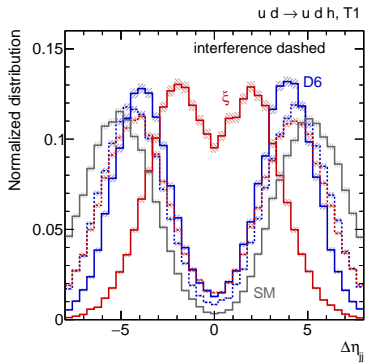
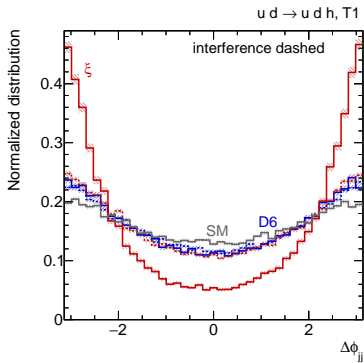
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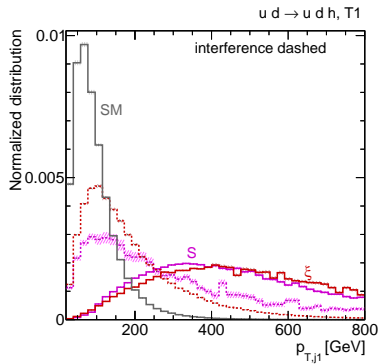
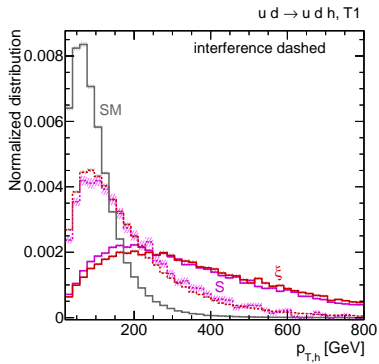
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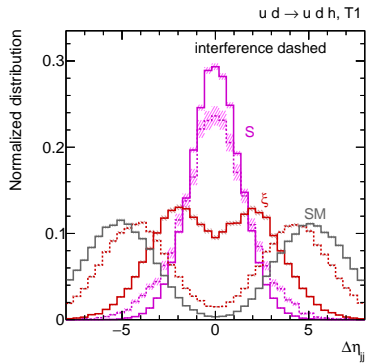
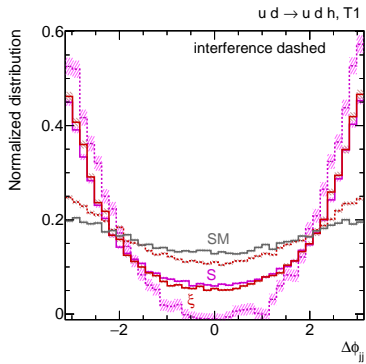
Only WBF diagrams, $\Delta\eta_{jj}$











Scalar splitting function

$$|\mathcal{M}(q \rightarrow q'S)|^2 = g_F^2 \frac{x^2 m_q^2}{1-x} + g_F^2 \frac{p_T^2}{1-x} + \mathcal{O}\left(\frac{m_q^2 p_T^2}{E^2}, \frac{m_q^4}{E^2}, \frac{p_T^4}{E^2}\right)$$

$$\sigma(qX \rightarrow q'Y) = \int dx dp_T F_S(x, p_T) \sigma(SX \rightarrow Y)$$

with the splitting function

$$F_S(x, p_T) = \frac{g_F^2}{16\pi^2} x \frac{p_T^3}{(m_S^2(1-x) + p_T^2)^2}$$

$$F_T(x, p_T) = \frac{g^2}{16\pi^2} \frac{1 + (1-x)^2}{x} \frac{p_T^3}{(m_W^2(1-x) + p_T^2)^2}$$

$$F_L(x, p_T) = \frac{g^2}{16\pi^2} \frac{(1-x)^2}{x} \frac{2m_W^2 p_T}{(m_W^2(1-x) + p_T^2)^2}$$

[9712400; S. Dawson (1985); G. L. Kane, W. W. Repko, W. B. Rolnick (1684); 0706.0536; 1202.1904]