

# "1/2 vs. 3/2" puzzle in $\overline{B} \rightarrow X_c l \bar{\nu}$

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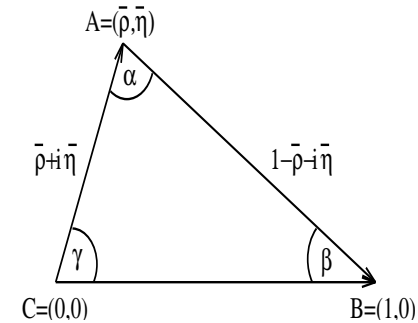
[I. Bigi, B. B., A. Le Yaouanc, L. Oliver, A. Oyanguren,  
O. Pène, J.C. Raynal, P. Roudeau, arXiv:0708.1621]

# Motivations

$$\underbrace{\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix}}_{\text{weak int.}} = V_{\text{CKM}} \underbrace{\begin{pmatrix} d \\ s \\ b \end{pmatrix}}_{\text{strong int.}} \quad V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}.$$

Without constraint:  $\delta V_{ij} < 5\%$ ,  $\delta V_{ij} > 5\%$ ,  $\delta V_{cb} \sim 1.5\%$

$$|\epsilon_K| = \bar{\eta} A^2 \hat{B}_K [1.11(5) A^2 (1 - \bar{\rho}) + 0.31(5)], \quad V_{cb} \sim \lambda^2 A$$



$$\delta \epsilon_K < 1\%, \quad \delta \hat{B}_K \sim 10\%, \quad \delta \bar{\eta}(V_{cb}) \sim 6\%$$

$$|V_{cb}|(\bar{B} \rightarrow D^* l \bar{\nu}) = (37.7 \pm 0.3 \pm 1.2 \pm_{1.4}^{1.2}) \times 10^{-3} \text{ [BABAR, '07]}$$

$$|V_{cb}|(\text{incl.}) = (41.7 \pm 0.7) \times 10^{-3} \text{ [PDG, '06]}$$

It is relevant to better figure out the **QCD nonperturbative dynamics** which enters in **all** processes involving **bounded quarks**  $\implies$  their SM contribution can be more easily distinguished from the contribution coming from a new physics.

What is the composition of the hadronic final state  $X_c$  in  $\bar{B} \rightarrow X_c l \bar{\nu}$ ?

$$\text{BR}(\bar{B}_d \rightarrow X_c l^- \bar{\nu}) = (10.33 \pm 0.28)\%$$

$$\text{BR}(\bar{B}_u \rightarrow X_c l^- \bar{\nu}) = (10.99 \pm 0.28)\%$$

		Mass (MeV)	Width (MeV)	$J^P$	$j_l^P$
$S: D^{(*)}$	$D^\pm$	$1869 \pm 0.5$	-	$0^-$	$\frac{1}{2}^-$
	$D^{*\pm}$	$2010 \pm 0.2$	$96 \pm 25$	$1^-$	$\frac{1}{2}^-$
$P: D^{**}$	$D_0^*$	$2352 \pm 50$	$261 \pm 50$	$0^+$	$\frac{1}{2}^+$
	$D_1^*$	$2427 \pm 26 \pm 25$	$384_{-75}^{+107} \pm 74$	$1^+$	$\frac{1}{2}^+$
	$D_1$	$2422.3 \pm 1.3$	$20.4 \pm 1.7$	$1^+$	$\frac{3}{2}^+$
	$D_2^*$	$2461.1 \pm 1.6$	$43 \pm 4$	$2^+$	$\frac{3}{2}^+$

$D^{**} \rightarrow D^{(*)} \pi$  is the main decay channel: parity and orbital momentum conservations  
 $\implies$  the decay occurs with the pion in a  $S$  wave or in a  $D$  wave

$D_{0,1}^* \rightarrow D^{(*)} \pi$ :  $S$  wave     $D_2^* \rightarrow D^{(*)} \pi$ :  $D$  wave

$D_1 \rightarrow D^* \pi$ :  $S$  and  $D$  wave are *a priori* allowed; however the  $S$  wave is forbidden by HQS

# Corroborated features

- Theory: – OPE and HQE  $\implies$  Bjorken, Uraltsev, Voloshin and moments sum rules  
 – Quark models that are **covariant** in the  $m_Q \rightarrow \infty$  limit  
 example: models *à la* Bakamijan-Thomas  
 – Lattice QCD

Experiment: B factories, LEP, Tevatron

States	% of $\Gamma(\bar{B} \rightarrow X_c l \bar{\nu})$
$D, D^*$	75 %
$D(3/2)$	$\sim 10$ %

[BABAR, '07]

[HFAG, '07]

[ALEPH, '97]

[DELPHI, '06]

[D0, '05]

[V. Morénas *et al*, '97] BT models

$D, D^*$  and  $D(3/2)$  do not saturate the total width;  $\sim 15$  % is composed of an unknown part  $D_X$ .

$$B^* - B \text{ splitting: } \mu_G^2(1 \text{ GeV}) = 0.35(3) \text{ GeV}^2$$

[O. Buchmüller, H, Flücher, '05]

$$\mu_\pi^2(\mu) > \mu_G^2(\mu)$$

[Belle, '06]

[BABAR, 07]

$$\mu_\pi^2(1 \text{ GeV})|_{\text{ref}} = 0.45 \text{ GeV}^2$$

[I. Bigi *et al*, '95] OPE

OPE treatment is successful for subclasses of inclusive transitions.

Generalisation of the IW function  $\xi(w)$

$$\Gamma(\bar{B} \rightarrow D_{1/2[3/2]}^{(n)} l \bar{\nu}) \propto |\tau_{1/2[3/2]}^{(n)}(w_n)|^2$$

$$\sum_n \left[ \tau_{3/2}^{(n)}(1) \right]^2 - \sum_n \left[ \tau_{1/2}^{(n)}(1) \right]^2 = \frac{1}{4}$$

$$\tau_{3/2}^0(1) > \tau_{1/2}^0(1)$$

$$\tau_{1/2}^0(1) \in [0.20, 0.40], \tau_{3/2}^0(1) \in [0.55, 0.70]$$

Suppression of  $D(1/2)$  with respect to  $D(3/2)$  due to kinematics

Factorisation in the Class I  $\bar{B} \rightarrow D^{**} \pi$ :  
from an analysis by Belle it is expected  
that  $\tau_{3/2}^0 > \tau_{1/2}^0$  as well

[V. Morénas *et al*, '97] BT models

[A. K. Leibovich *et al*, '98]

[D. Ebert *et al*, '98] Relativistic model

[N. Uraltsev, '01] Uraltsev sum rule

[D. Bećirević *et al*, '05] Lattice

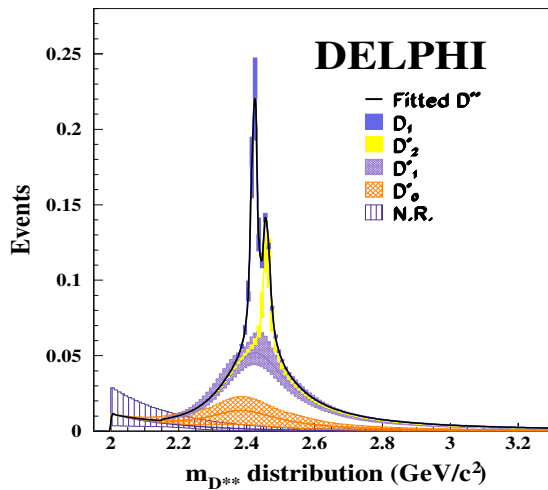
[Belle, '04]

$D(3/2)$  is expected to dominate  $D(1/2)$  in  $\bar{B} \rightarrow X_c l \bar{\nu}$ .

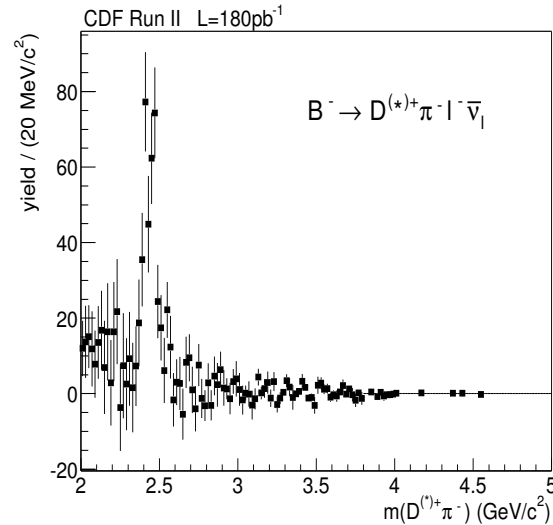
# Issues

DELPHI found a larger component of broad states than of the narrow states. Interpretation as  $D_0^*$  and  $D_1^*$ ??  $\implies$  Clear conflict with theory, '1/2' vs. '3/2' puzzle [V. Morénas et al, '01], [N. Uraltsev, '04]

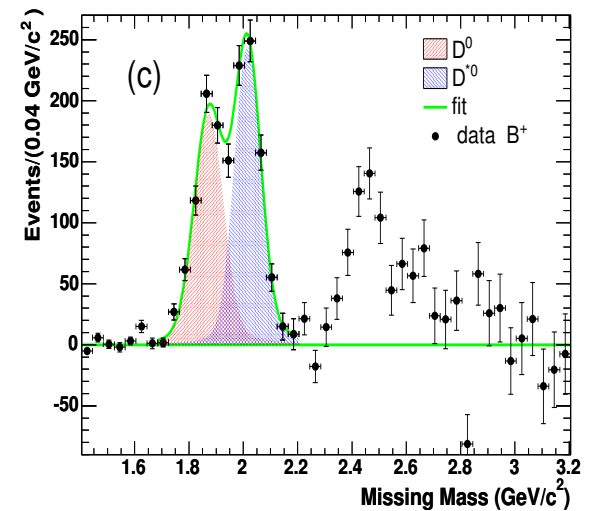
[DELPHI, '06]



[CDF, '05]



[BABAR: '06]



Up to now the experimental verdict about  $\bar{B} \rightarrow [D/D^* \pi]_{\text{broad}} l \bar{\nu}$  is not clear.

No obvious theoretical candidates for those broad states if the mass distribution is centered below 2.5  $\text{GeV}$ .

An important check of the theory is  $\langle M(D_X) \rangle$ : depending on BR ( $\bar{B} \rightarrow D^* l \bar{\nu}$ ) it varies from 2.4 and 2.7.

The extension of BT models to finite quark masses just started: predictions concerning the relative weight of  $\tau_{1/2}^0$  and  $\tau_{3/2}^0$  could change by including those corrections.

Some "exotic" possibilities similar to the nucleons Roper resonance could be investigated.

The study of the spectrum of radial and orbital excitations of the  $D$  meson on the lattice must be pursued.

[A. Green *et al*, '03]

[J. Foley *et al*, '07]

Nice results concerning  $\bar{B} \rightarrow D/D^* l \bar{\nu}$  are already available.

[G. M. de Divitiis *et al*, '07]

[J. Laiho, '07]

[S. Simula, '07]

The extension to  $D^{**}$  seems to be the next step, beyond the exploratory study performed before, in order to conclude about the relative weight of  $\tau_{1/2}^0$  and  $\tau_{3/2}^0$ .

[D. Bećirević *et al*, '05]

# Outlook

- The composition of the final state  $X_c$  in  $\bar{B} \rightarrow X_c l \bar{\nu}$  has received some attention since 10 years.
- Theoretically, it is expected that the states  $D, D^*$  and the 4 P wave states  $D^{**}$  do not saturate the total width. Moreover, covariant quark models and sum rules extracted from the OPE in the  $m_Q \rightarrow \infty$  limit lead to  $[\Gamma(\bar{B} \rightarrow D(\frac{1}{2}) l \bar{\nu}) < \Gamma(\bar{B} \rightarrow D(\frac{3}{2}) l \bar{\nu})]^{\text{TH}}$ .
- Experimentally, it was found at LEP that the total width is saturated by  $D, D^*, D^{**}$  and the measured branching ratios read  $[\Gamma(\bar{B} \rightarrow D(\frac{1}{2}) l \bar{\nu}) > \Gamma(\bar{B} \rightarrow D(\frac{3}{2}) l \bar{\nu})]^{\text{EXP}}$ .
- However there are strong theoretical assumptions that the broad states observed in the  $\bar{B} \rightarrow D^{**} l \bar{\nu}$  mass distribution are not the P wave states.
- An important experimental effort is demanded, in particular to have a better knowledge of the quantum numbers of those broad states.
- The answer will have an impact on the theoretical control over QCD nonperturbative dynamics of the heavy-light systems.
- On the theoretical side, taking account of  $1/m_Q$  corrections is crucial, either in the analytical treatment of QCD (OPE, quark models) or in its numerical one (lattice).