Anomalous enhancement of $B \to K \eta'$ branching ratio

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Contents

- Puzzle of $B \to K\eta'$ branching ratio
- Theoretical model predictions
- ullet Estimate of the density matrix for η'
- Numerical results
- Conclusions

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Puzzle of large branching ratio of $B \to K \eta'$

 \square CLEO '97: observation of large branching ratio for $B \to K\eta'$

$$Br(B^+ \to K^+ \eta') = (70.2 \pm 2.5) \times 10^{-6}$$

c.f. similar penguin dominant processes:

$$Br(B^+ \to K^+ \pi^0) = (12.9 \pm 0.6) \times 10^{-6}$$

 $Br(B^+ \to K^+ \eta) = (2.7 \pm 0.3) \times 10^{-6}$

 \square SU(3) relations: ratios of $B \to K[0^{-+}]$ branching ratios Lipkin (PLB '97), Dighe, Gronau, Rosner (PRL '98)

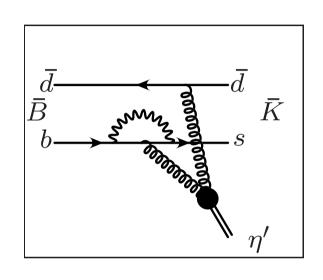
$$Br(K\eta'): Br(K\eta): Br(K\pi^0) = 3:0:1$$

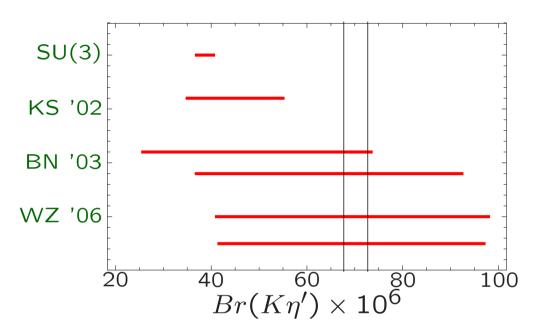
where $\theta=-19.5^\circ$ so that $\eta=(u\bar{u}+d\bar{d}-s\bar{s})/\sqrt{3},\ \eta'=(u\bar{u}+d\bar{d}+2s\bar{s})/\sqrt{6}.$

Theoretical predictions: status

☐ Inclusion of an U(1) breaking: anomaly diagram

Atwood, Soni (PLB '97), E.K, Ahmady, Sugamoto (PRD '98)



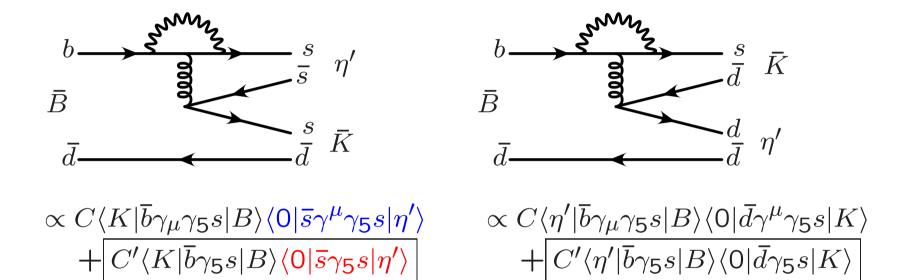


E.K. and Sanda (PLB '02), Beneke, Neubert (NPB '03), Williamson, Zupan (PRD '06)

rightharpoonup Large uncertainties could be reduced by the other observables, e.g. $\pi\eta', K^*\eta'$, etc.

see also Escribano, Matias, Virto ('07)

$B \to K \eta'$ and density matrix for η'



- rightharpoons Density matrix $\langle 0|\bar{s}\gamma_5 s|\eta'\rangle$: this talk
- $rac{rac}{rac}$ Decay constant $\langle 0|\bar{s}\gamma_{\mu}\gamma_{5}s|\eta'\rangle$: two mixing angle scheme Kaiser, Leutwyler ('98), Feldman, Kroll, Stech (PRD '98)
- $rac{raction}{raction}$ Form factor $\langle \eta' | \bar{s} \gamma_{\mu} \gamma_5 b | B \rangle$: QCD sum rule, lattice, experimental extraction.

Ball, Jones (JHEP '07), Pham ('07), Charng, Kurimoto, Li (PRD '06)

Anomaly contribution to the density matrix

 \square Example of the density matrix for K

$$\begin{array}{rcl} \langle 0|\bar{d}\gamma_{\mu}\gamma_{5}s|K\rangle &=& if_{K}p_{\mu} \\ & \downarrow & \partial^{\mu} \text{ and Eq. of motion} \\ \\ \langle 0|\bar{d}\gamma_{5}s|K\rangle &=& \frac{m_{K}^{2}}{m_{s}+m_{d}}f_{K} \end{array}$$

 \square Density matrix of $\eta - \eta'$:

Derivative of the singlet current leads to the anomaly term. Thus the usual technique like the K case does not work.

$$\partial^{\mu}(\bar{s}\gamma_{\mu}\gamma_{5}s) = 2im_{s}\bar{s}\gamma_{5}s + \frac{\alpha_{s}}{4\pi}G^{a}_{\mu\nu}\tilde{G}^{\mu\nu}_{a}$$

Estimate of density matrix

☐ FKS scheme:

see also Beneke, Neubert (NPB '03)

$$\begin{pmatrix} \eta \\ \eta' \end{pmatrix} = \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} (u\bar{u} + d\bar{d})/\sqrt{2} \\ s\bar{s} \end{pmatrix} : \quad \phi = \theta - \theta_I + \frac{\pi}{2}$$

$$\begin{vmatrix} \langle 0|\overline{s}\gamma_{\mu}\gamma_{5}s|\eta^{(\prime)}\rangle = if_{\eta(\prime)}^{s}p_{\mu} \\ f_{\eta}^{q} = f_{q}\cos\phi, \quad f_{\eta}^{s} = -f_{s}\sin\phi \\ f_{\eta'}^{q} = f_{q}\sin\phi, \quad f_{\eta'}^{s} = f_{s}\cos\phi \end{vmatrix} \Rightarrow \begin{vmatrix} 2m_{s}\langle 0|\overline{s}\gamma_{5}s|\eta^{(\prime)}\rangle = ih_{\eta(\prime)}^{s} \\ h_{\eta}^{q} = h_{q}\cos\phi, \quad h_{\eta}^{s} = -h_{s}\sin\phi \\ h_{\eta'}^{q} = h_{q}\sin\phi, \quad h_{\eta'}^{s} = h_{s}\cos\phi \end{vmatrix}$$

 \square Effective Lagrangian at large N_c (this talk)

E.K. and Gerard (PRL '07)

$$\mathcal{L} = \frac{f^2}{8} \langle \partial_{\mu} U \partial^{\mu} U^{\dagger} \rangle + \frac{m_0^2}{4N_c} \frac{f^2}{8} \langle \ln U - \ln U^{\dagger} \rangle^2 + \frac{f^2}{8} r \langle m U^{\dagger} + U m \rangle$$

$$+ \frac{f^2}{8} \left[-\frac{r}{\Lambda^2} \langle m \partial^2 U^{\dagger} \rangle + \frac{r^2}{2\Lambda_1^2} \langle m U^{\dagger} m U^{\dagger} \rangle + \frac{r}{2\Lambda_2^2} \langle m U^{\dagger} \partial_{\mu} U \partial^{\mu} U^{\dagger} \rangle \right] + h.c.$$

Parameters are fitted using $\eta - \eta'$ masses, J/ψ radiative decays and π, K masses and decay constants.

Numerical result

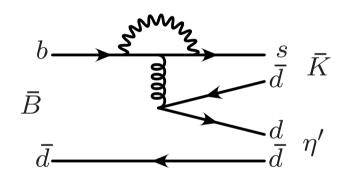
$$\zeta \equiv \frac{\langle 0|\bar{s}\gamma_5 s|\eta\rangle}{\langle 0|\bar{d}\gamma_5 s|K\rangle}/\sin\phi, \quad \zeta' \equiv \frac{\langle 0|\bar{s}\gamma_5 s|\eta'\rangle}{\langle 0|\bar{d}\gamma_5 s|K\rangle}/\cos\phi$$

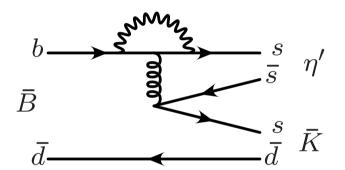
	our result	SU(3)	AG	BN
ζ	1.29 ± 0.19	1	1.38	1.34
ζ'	1.72 ± 0.26	1	1.12	1.07

AG: Ali, Greub (PRD '98), BN: Beneke, Neubert (NPB '03)

Impact on the SU(3) relation

$$\frac{A(B \to K\eta^{(\prime)})}{A(B \to K\pi^{0})}\Big|_{C'} = \frac{\langle K|\bar{b}\gamma_{5}s|B\rangle}{\langle \pi|\bar{b}\gamma_{5}d|B\rangle} \left| \frac{\langle \eta^{(\prime)}|\bar{b}\gamma_{5}d|B\rangle}{\langle K|\bar{b}\gamma_{5}s|B\rangle} + \underbrace{\frac{\langle 0|\bar{s}\gamma_{5}s|\eta^{(\prime)}\rangle}{\langle 0|\bar{d}\gamma_{5}s|K\rangle}}_{\propto \zeta^{(\prime)}} \right|.$$



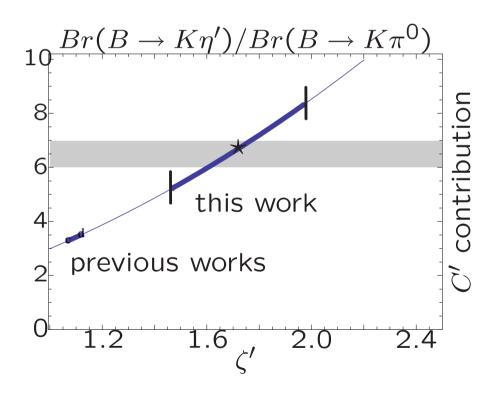


$$\frac{A(B \to K\eta)}{A(B \to K\pi^0)} \simeq -\left[\cos\phi - \zeta\sqrt{2}\sin\phi\right]^{\theta = -19.5^{\circ}} - \sqrt{\frac{2}{3}}[1 - \zeta]$$

$$\frac{A(B \to K\eta')}{A(B \to K\pi^0)} \simeq -\left[\sin\phi + \zeta'\sqrt{2}\cos\phi\right]^{\theta = -19.5^{\circ}} - \sqrt{\frac{1}{3}}[1 + 2\zeta']$$

At SU(3) limit, $\zeta^{(\prime)} = 1$.

Impact on $B \to K\eta'$



Conclusions

□ We emphasised that the puzzle of the large branching ratio for $B \to K \eta'$ has not been fully understood yet.

 \square We computed the density matrix for $\eta-\eta'$ using the effective Lagrangian at large N_c limit.

□ We found over 50 % of increase in the density matrix for η' comparing to the previous works, which would enhance the existing theoretical predictions for the $B \to K\eta'$ branching ratio.