

Progress on analytical expression of K_{l_3} form factors at two loop order

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In collaboration with

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Definitions of the form factors

J. Bijnens and P. Talavera hep-ph/0303103, Nucl. Phys. B669 (2003) 341-362

In the chiral conventions

$$U \doteq \exp\left(\frac{i\sqrt{2}}{F_0}\Phi\right) \quad \Phi \doteq \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & \pi^+ & K^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & K^0 \\ K^- & K^0 & -\frac{2\eta}{\sqrt{6}} \end{pmatrix}$$

We consider here the form factors defined as

$$\langle K^+(p) | \bar{u}\gamma_\mu s | \pi^0(q) \rangle \doteq \frac{1}{\sqrt{2}} \left[(p_\mu + q_\mu) f_+^{K\pi}(t) + (p_\mu - q_\mu) f_-^{K\pi}(t) \right] \quad (1)$$

and

$$f_0^{K\pi}(t) = f_+^{K\pi}(t) + \frac{t}{m_K^2 - m_\pi^2} f_-^{K\pi}(t)$$

with $t \doteq (q - p)^2$.

Specially, we are considering the limit

$$\lim_{t \rightarrow 0} f_0^{K\pi}(t) \doteq f_0^{K\pi} = f_+^{K\pi}(0) \quad (2)$$

Those coefficients are strongly connected to V_{us} .

J. Bijnens' web page <http://www.thep.lu.se/bijnens/chpt.html>

The chiral expansion of $f^{K\pi}$ is given by

$$f = \underbrace{f^{(2)}}_{\doteq 1} + f^{(4)} + f^{(6)} \quad (3)$$

From now, we are taking the two loops expression for $f^{K\pi}$ provided by J. Bijnens.

It involves scalar two loop D-dimensional integrals

$$V \doteq \int \frac{d^D k_1}{(2\pi)^D} \frac{d^D k_2}{(2\pi)^D} \frac{\text{Num.}}{[k_1^2 - m_1^2] [(k_1 - q)^2 - m_2^2] [(k_1 + k_2 - p)^2 - m_4^2]} \quad (4)$$

at the end 396 scalars integrals!

Necessity of a method to obtain a complete analytical expression.

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Necessity of a method to obtain a complete analytical expression.

To reduce the number of integrals to calculate we propose the following algorithm :

1. Use the Laporta's Algorithm to reduce the number of integral to a minimal set of Master Integrals.
2. Use the inverse multi-dimensional Converse Mapping theorem to evaluate the analytical expressions of the unknown Master Integrals.
3. Give the analytical expression of the form factors.

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Laporta's Algorithm

S.Laporta, *Int. J. Mod. Phys. A* 15 (2000) 5087

T. Gehrmann and E. Remiddi, *Nucl. Phys. B* 580 (2000) 485

R. Bonciani, P. Mastrolia and E. Remiddi, *Nucl. Phys. B* 661 (2003) 289

Every two loops amplitudes obey to the following form

$$\mathcal{I} = \int \frac{d^D k_1}{(2\pi)^D} \frac{d^D k_2}{(2\pi)^D} \frac{S_1^{n_1} \cdots S_N^{n_N}}{D_1^{\ell_1} \cdots D_L^{\ell_L}}$$

for S scalar products and D denominators.

1. Use the Integrate by parts relation (Stokes' theorem)

$$\int \frac{d^D k_1}{(2\pi)^D} \frac{d^D k_2}{(2\pi)^D} \frac{\partial}{\partial k_j^\mu} \left[v^\mu \frac{S_1^{n_1} \cdots S_N^{n_N}}{D_1^{\ell_1} \cdots D_L^{\ell_L}} \right] = 0$$

for $v = k_1, k_2, p, q$.

2. Use Lorentz' invariance and discrete symmetries

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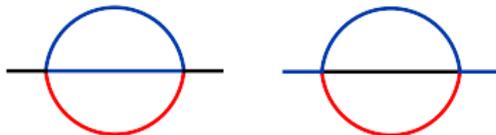
2. Use Lorentz' invariance and discrete symmetries

After the 2 points of the Laporta's Algorithm we obtain a linear system where every amplitudes are linked together.

Finally, using a recursive method and over constraining the generated system, every integrals can be deduced from a small set of **Master's integrals**.

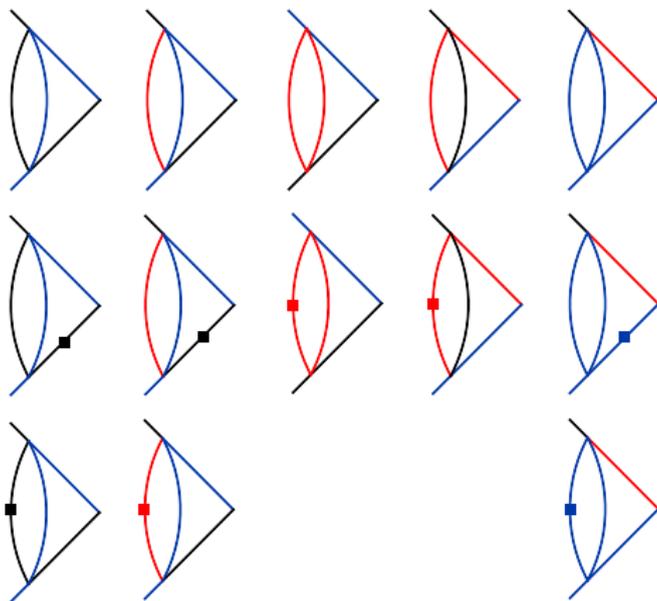
Here, after applying this algorithm, the only analytically unknown topologies of Master Integrals are

- Two points functions :



where m_π , m_K and m_η .

- Three points functions :



m_π , m_K and m_η .

One dimensional Mellin's Transform and *Converse Mapping theorem*

The Mellin's transform of a function f and its inverse transform are defined as

$$\mathcal{M}[f(x)](s) \doteq \int_0^{\infty} dx x^{s-1} f(x) \quad \longleftrightarrow \quad f(x) = \int_{c-i\infty}^{c+i\infty} \frac{ds}{2i\pi} x^{-s} \mathcal{M}[f(x)](s)$$

If and only if

$$c \doteq \operatorname{Re} s \in]\alpha, \beta[\quad \text{written} \quad \langle \alpha, \beta \rangle \quad \text{Fundamental strip}$$

It corresponds to the behaviours

$$f(x) \underset{x \rightarrow 0^+}{=} \mathcal{O}(x^{-\alpha}) \quad \& \quad f(x) \underset{x \rightarrow +\infty}{=} \mathcal{O}(x^{-\beta})$$

Examples :

f	\longleftrightarrow	$\mathcal{M}[f]$	
e^{-x}	\longleftrightarrow	$\Gamma(s)$	$\langle 0, \infty \rangle$
$(1+x)^{-\nu}$	\longleftrightarrow	$\frac{\Gamma(\nu-s)\Gamma(s)}{\Gamma(\nu)}$	$\langle 0, \operatorname{Re} \nu \rangle$
$\ln(1+x)$	\longleftrightarrow	$\frac{\pi}{s \sin \pi s}$	$\langle -1, 0 \rangle$

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Converse Mapping theorem

Flajolet et al. (1994)

Friot, Greynat and de Rafael (2005)

Idea : The singularities in the complex Mellin's plan manage completely the asymptotic behaviour of the associated function

Exemple :

$$\Gamma(s) \asymp \sum_{p=0}^{\infty} \frac{(-1)^p}{p!} \frac{1}{s+p} \quad \longleftrightarrow \quad e^{-x} \underset{x \rightarrow 0}{\sim} \sum_{p=0}^{\infty} \frac{(-1)^p}{p!} x^p$$

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$$\mathcal{M}[f(x)]_{\text{right}}(s) \asymp \sum_{p > \beta, n} c_{p,n} \frac{1}{(s-p)^n} \quad \leftrightarrow \quad f(x) \underset{x \rightarrow +\infty}{\sim} - \sum_{p > \beta, n} c_{n,p} x^{-p} \frac{(-1)^{n-1}}{(n-1)!} \ln^{n-1} x$$

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Multi-dimensional Mellin's Transform and Grothendieck's Residues theory

We define the n -dimensional Mellin's transform of function f as

$$\mathcal{M}[f](s_1, \dots, s_n) \doteq \int_0^\infty dx_1 \cdots \int_0^\infty dx_n x_1^{s_1-1} \cdots x_n^{s_n-1} f(x_1, \dots, x_n)$$

and its inverse transformation

$$f(x_1, \dots, x_n) \doteq \int_{c_1+i\mathbb{R}} \frac{ds_1}{2i\pi} \cdots \int_{c_n+i\mathbb{R}} \frac{ds_n}{2i\pi} x_1^{-s_1} \cdots x_n^{-s_n} \mathcal{M}[f](s_1, \dots, s_n)$$

This inversion formula is of course valid in the **fundamental polyhedra** defined as all the constraints on $\mathbf{c} \doteq {}^T(c_1, \dots, c_n)$ where the Mellin's transform is completely analytical.

If we want to extend the *Converse Mapping* Theorem to the multi-dimensional case we need to introduce "briefly" the **Grothendieck's Residue theory**.

A few words on Grothendieck's Residues theory

P. Griffiths, J.Harris, *Principles of Algebraic Geometry*, Wiley NYC 1978

A.K. Tsikh et al., hep-th 9609215

N.B. : From now, all vectors in n-dimension are written as $\mathbf{s} = {}^T(s_1, \dots, s_n)$

One way to see the residues in multi-dimensional complex analysis is to consider the quantity (for any h completely analytic)

$$\text{Res.} \left[\frac{h(\mathbf{s})}{f_1(\mathbf{s}) \cdots f_n(\mathbf{s})} \right]_{\mathbf{0}} = \oint_0 \frac{h(\mathbf{s})}{f_1(\mathbf{s}) \cdots f_n(\mathbf{s})} \frac{ds_1}{2i\pi} \wedge \cdots \wedge \frac{ds_n}{2i\pi} \doteq \oint_0 \omega$$

All the curves, the **divisors**, in the $2n$ -dimension complex space given by $j \in [1, n]$

$$D_j \doteq \{ \mathbf{s} \in \mathbb{C}^n, f_j(\mathbf{s}) = 0 \}$$

have intersections points in this space. They provide the calculation of the residue in a summation over

$$\bigcap_{j \in [1, n]} D_j$$

via a sequential Cauchy's theorem.

Multi-dimensional Converse Mapping Theorem

J.-Ph. Aguilar, D. Greynat and E. de Rafael, Work in progress (2008)

Idea : If you combine the calculation of the Grothendieck's residues and the *multi-dimensional Jordan's lemma* you can define sectors in complex plans where the x_j are bigger or smaller than 1 and their relative position and permit to generate the complete asymptotic behaviour in each variables

In the case of ratios of Euler's second function : the Γ function, the *multi-dimensional Converse Mapping theorem* for

$$f(\mathbf{s}) = \int_{\gamma+i\mathbb{R}^n} x_1^{-s_1} \cdots x_n^{-s_n} \frac{\prod_{j=1}^{j=p} \Gamma(\mathbf{a}_j \cdot \mathbf{s} + b_j)}{\prod_{k=1}^{k=q} \Gamma(\mathbf{c}_k \cdot \mathbf{s} + d_k)} \frac{ds_1}{2i\pi} \wedge \cdots \wedge \frac{ds_n}{2i\pi}$$

the divisors are

$$D_j^\ell = \{ \mathbf{s} \in \mathbb{C}^n, \mathbf{a}_j \cdot \mathbf{s} + b_j = -\ell, \ell \in \mathbb{N} \}$$

J.-Ph. Aguilar, D. Greynat and E. de Rafael, Work in progress (2008)

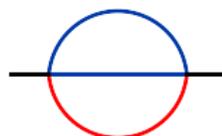
The multi-dimensional Jordan's lemma provides a sub-set \mathcal{J} of index j to permit the convergence of the series coming from the calculation of the Grothendieck's residue theorem. (we give here the theorem only in the case of simple non-degenerate poles)

$$\begin{aligned}
 f(\mathbf{s}) &= \sum_{j \in \mathcal{J}} \text{Res.} \left[\frac{\prod_{k=1}^{k=q} \Gamma(\mathbf{c}_k \cdot \mathbf{s} + d_k)^{-1}}{\prod_{j=1}^{j=p} \Gamma(\mathbf{a}_j \cdot \mathbf{s} + b_j)^{-1}} \right]_{\mathbf{s} \in \mathcal{D}_j} \\
 &= \sum_{j \in \mathcal{J}} \frac{(-1)^{|\ell|}}{\ell! \det(\mathbf{a}_j)} \frac{\prod_{j \neq \mathcal{J}} \Gamma(\mathbf{a}_j \cdot \mathbf{s}_j^\ell + b_j)}{\prod_{k=1} \Gamma(\mathbf{c}_k \cdot \mathbf{s}_j^\ell + d_k)} x_1^{-(s_j^\ell)_1} \dots x_n^{-(s_j^\ell)_n}
 \end{aligned}$$

Hopefully more clear on the following example...

An example of the calculation on the Master Integrals

We consider the sunrise type integral (m_π , m_K and m_η)



1. Feynman's parametrization $D = 4 - \epsilon$

$$\begin{aligned}
 & H(m_\eta^2, m_K^2, m_K^2, m_\pi^2) \\
 &= -\frac{\pi^D}{(2\pi)^{2D}} \Gamma(\epsilon - 1) \iint_{[0,1]^2} dx dy (1-x)^{\frac{3}{2}\epsilon-2} [1-Y+Yx]^{\frac{3}{2}\epsilon-3} \\
 &\quad \times (1-Y)^{1-\epsilon} x^{1-\epsilon} (1-x)^{1-\epsilon} \left(m_\pi^2\right)^{1-\epsilon} \\
 &\quad \times \left| 1 - \frac{1-Y+xY}{(1-x)(1-Y)} \rho_2 - \frac{1-Y+xY}{x(1-Y)} \rho_1 \right|^{1-\epsilon}
 \end{aligned}$$

where $Y = 1 - y(1 - y)$, $\rho_1 = m_K^2/m_\pi^2$ and $\rho_2 = m_\eta^2/m_\pi^2$.

2. Inverse Mellin's representation

Using the general functions inverse Mellin's representation ($c \in \mathbb{C}, \nu > 0$)

$$|1-x|^{-\nu} \text{sign}(1-x) = \int_{c+i\mathbb{R}} \frac{ds}{2i\pi} x^{-s} \Gamma(1-\nu) \left[\frac{\Gamma(s)}{\Gamma(s-\nu+1)} - \frac{\Gamma(\nu-s)}{\Gamma(1-s)} \right],$$

And using the polar coordinates we obtain the double inverse Mellin's representation of H :

$$\begin{aligned} H(m_\eta^2, m_K^2, m_K^2, m_\pi^2) &= -\frac{\Gamma(\epsilon-1)\Gamma(1-\epsilon)}{(4\pi)^D} \left(\frac{m_\pi^2}{2}\right)^{1-\epsilon} \sqrt{\pi} \\ &\times \int_{c+i\mathbb{R}^2} \frac{ds_1 \wedge ds_2}{(2i\pi)^2} (4\rho_1)^{-s_1} \rho_2^{-s_2} M(s_1, s_2) \\ &\times \left[h(s_1, s_2) - \frac{\rho_1}{4} h(s_1+1, s_2) - \frac{\rho_2}{4} h(s_1, s_2+1) \right], \end{aligned}$$

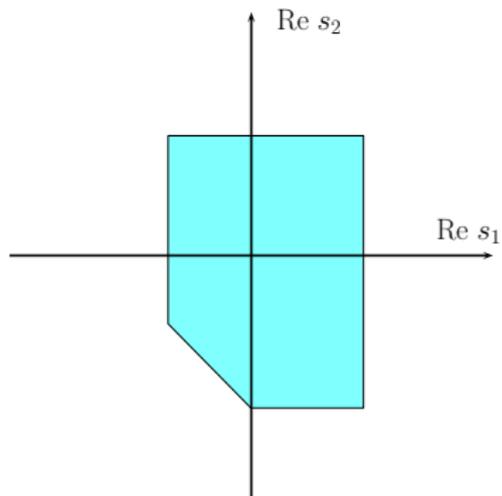
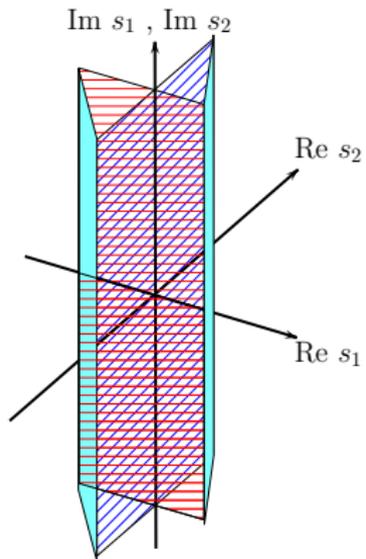
with

$$M(s_1, s_2) = \frac{\Gamma(s_1)\Gamma(s_2)}{\Gamma(s_1+s_2)} \left[\frac{\Gamma(s_1+s_2)}{\Gamma(s_1+s_2-\epsilon+1)} - \frac{\Gamma(\epsilon-s_1-s_2)}{\Gamma(1-s_1-s_2)} \right]$$

and

$$h(s_1, s_2) = \frac{\Gamma(2-\epsilon+s_1)\Gamma(1-\frac{\epsilon}{2}+s_2)\Gamma(1-\frac{\epsilon}{2}+s_1)}{\Gamma(3-\frac{3}{2}\epsilon+s_1+s_2)\Gamma(\frac{3}{2}-\frac{\epsilon}{2}+s_1)}$$

3. Fundamental polyhedra and pertinent sector



We obtain here 6 different 2-forms :

$$\omega_1 \doteq \frac{ds_1 \wedge ds_2}{(2i\pi)^2} (4\rho_1)^{-s_1} \rho_2^{-s_2} \Gamma \left[\begin{matrix} s_1, s_2, 2-\epsilon+s_1, 1-\frac{\epsilon}{2}+s_2, 1-\frac{\epsilon}{2}+s_1 \\ 1-\epsilon+s_1+s_2, 3-\frac{3}{2}\epsilon+s_1+s_2, \frac{3}{2}-\frac{\epsilon}{2}+s_1 \end{matrix} \right]$$

$$\omega_2 \doteq \frac{ds_1 \wedge ds_2}{(2i\pi)^2} (4\rho_1)^{-s_1} \rho_2^{-s_2} \Gamma \left[\begin{matrix} s_1, s_2, 3-\epsilon+s_1, 1-\frac{\epsilon}{2}+s_2, 2-\frac{\epsilon}{2}+s_1 \\ 1-\epsilon+s_1+s_2, 4-\frac{3}{2}\epsilon+s_1+s_2, \frac{5}{2}-\frac{\epsilon}{2}+s_1 \end{matrix} \right]$$

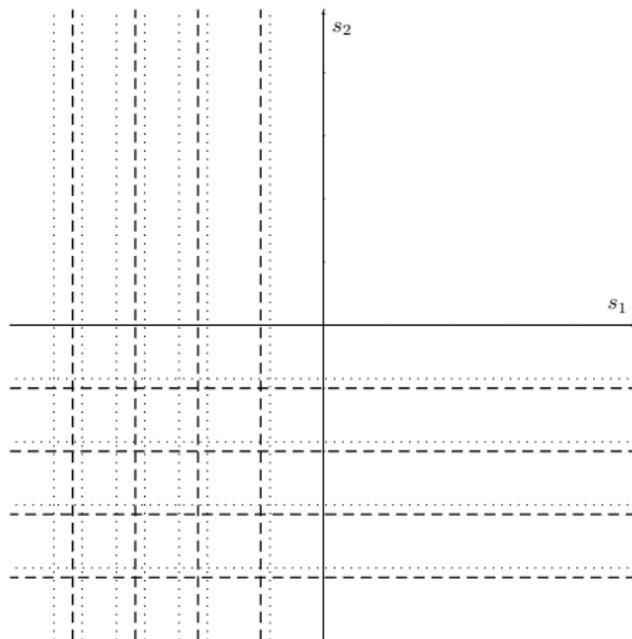
$$\omega_3 \doteq \frac{ds_1 \wedge ds_2}{(2i\pi)^2} (4\rho_1)^{-s_1} \rho_2^{-s_2} \Gamma \left[\begin{matrix} s_1, s_2, 2-\epsilon+s_1, 2-\frac{\epsilon}{2}+s_2, 1-\frac{\epsilon}{2}+s_1 \\ 1-\epsilon-s_1-s_2, 4-\frac{3}{2}\epsilon+s_1+s_2, \frac{3}{2}-\frac{\epsilon}{2}+s_1 \end{matrix} \right]$$

$$\omega_4 \doteq -\frac{ds_1 \wedge ds_2}{(2i\pi)^2} (4\rho_1)^{-s_1} \rho_2^{-s_2} \Gamma \left[\begin{matrix} s_1, s_2, \epsilon-s_1-s_2, 2-\epsilon+s_1, 1-\frac{\epsilon}{2}+s_2, 1-\frac{\epsilon}{2}+s_1 \\ s_1+s_2, 1-s_1-s_2, 3-\frac{3}{2}\epsilon+s_1+s_2, \frac{3}{2}-\frac{\epsilon}{2}+s_1 \end{matrix} \right]$$

$$\omega_5 \doteq -\frac{ds_1 \wedge ds_2}{(2i\pi)^2} (4\rho_1)^{-s_1} \rho_2^{-s_2} \Gamma \left[\begin{matrix} s_1, s_2, \epsilon-s_1-s_2, 3-\epsilon+s_1, 1-\frac{\epsilon}{2}+s_2, 2-\frac{\epsilon}{2}+s_1 \\ s_1+s_2, 1-s_1-s_2, 4-\frac{3}{2}\epsilon+s_1+s_2, \frac{5}{2}-\frac{\epsilon}{2}+s_1 \end{matrix} \right]$$

$$\omega_6 \doteq -\frac{ds_1 \wedge ds_2}{(2i\pi)^2} (4\rho_1)^{-s_1} \rho_2^{-s_2} \Gamma \left[\begin{matrix} s_1, s_2, \epsilon-s_1-s_2, 2-\epsilon+s_1, 2-\frac{\epsilon}{2}+s_2, 1-\frac{\epsilon}{2}+s_1 \\ s_1+s_2, 1-s_1-s_2, 4-\frac{3}{2}\epsilon+s_1+s_2, \frac{3}{2}-\frac{\epsilon}{2}+s_1 \end{matrix} \right]$$

The divisors are for example for ω_1 the following lines



Multi-dimensional Converse Mapping theorem implies to sum over intersections in the fourth quadrant.

Obtaining then the following representation

$$\int \omega_1 = \sum_{n=0}^{\infty} \sum_{k=0}^{\infty} \frac{(4\rho_1)^n \rho_2^k}{n! k!} \left(\Gamma \left[\begin{matrix} 2-\epsilon-n, 1-\frac{\epsilon}{2}-k, 1-\frac{\epsilon}{2}-n \\ -n-k-\epsilon+2, 3-\frac{3}{2}\epsilon-n-k, \frac{3}{2}-\frac{\epsilon}{2}-n \end{matrix} \right] \right. \\ \left. + \rho_2^{1-\frac{\epsilon}{2}} \Gamma \left[\begin{matrix} -k-1+\frac{\epsilon}{2}, 2-\epsilon-n, 1-\frac{\epsilon}{2}-n \\ \frac{3}{2}-\frac{\epsilon}{2}-n, 2-\epsilon-n-k, 1-\frac{\epsilon}{2}-n-k \end{matrix} \right] + 4\rho_1 \Gamma \left[\begin{matrix} 1-\frac{\epsilon}{2}-k, 1-\frac{\epsilon}{2}-n, -1+\frac{\epsilon}{2}-n \\ \frac{1}{2}-n, 2-\epsilon-k-n, 1-\frac{\epsilon}{2}-k-n \end{matrix} \right] \right)$$

We do the same process for all 6 2-forms ω_j ... To obtain finally the following behaviour after an ϵ -expansion

$$\bar{H}(m_\eta^2, m_K^2, m_\pi^2, m_\pi^2) \sim \frac{1}{\epsilon^2} \left[-\frac{1}{8} (2m_K^2 + m_\eta^2) \right] + \dots$$

in agreement with literature

Of course the epsilon-expansion and the cut of the infinite series are not obligatory and in this sense, we have an analytic expansion of the Master Integrals .

M. Caffo et al., Nuovo Cimmento Vol. III, A, N 4 (1998)

CONCLUSIONS

A work in progress... closed to the end...