

# Bounds in $\pi\pi$ scattering from dispersion relations

Constraining  $\bar{l}_1$  &  $\bar{l}_2$  and  $L_1, L_2$  &  $L_3$

Vicent Mateu

[mateu@ific.uv.es](mailto:mateu@ific.uv.es)

IFIC-Universitat de Valencia CSIC

Euroflavour07, Paris



UNIVERSITAT DE VALÈNCIA



CSIC

FlaviA  
net

# Outline

- 1 Motivations
- 2  $SU(2)$ 
  - $\pi\pi$  scattering
  - Fixed  $t$  dispersion relations & positivity conditions
  - Bounds on chiral LECs and the Linear Sigma Model
  - Equivalence with Pennington & Portoles
- 3  $SU(3)$ 
  - $SU(3)_V$  limit
  - Symmetry breaking
  - Results
- 4 Conclusions

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.
- 2 **Model independent** approach  $\implies$  dispersion relations.

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.
- 2 **Model independent** approach  $\implies$  dispersion relations.  
Dispersion relations and unitarity  $\implies$  positivity conditions.

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.
- 2 **Model independent** approach  $\implies$  dispersion relations.  
Dispersion relations and unitarity  $\implies$  positivity conditions.  
Positivity conditions  $\implies$  **bounds** for  $\bar{l}_1$  and  $\bar{l}_2$  (and  $L_{1,2,3}$ ).

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.
- 2 **Model independent** approach  $\implies$  dispersion relations.  
Dispersion relations and unitarity  $\implies$  positivity conditions.  
Positivity conditions  $\implies$  **bounds** for  $\bar{l}_1$  and  $\bar{l}_2$  (and  $L_{1,2,3}$ ).
- 3 Other (**less restrictive**) methods are contained in ours.  
[Distler et al '06 hep-ph/0604255 \[1\]](#)  
[Pennington et al '95 hep-ph/9409426 \[2\]](#)

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.
- 2 **Model independent** approach  $\implies$  dispersion relations.  
 Dispersion relations and unitarity  $\implies$  positivity conditions.  
 Positivity conditions  $\implies$  **bounds** for  $\bar{l}_1$  and  $\bar{l}_2$  (and  $L_{1,2,3}$ ).
- 3 Other (**less restrictive**) methods are contained in ours.  
[Distler et al '06 hep-ph/0604255 \[1\]](#)  
[Pennington et al '95 hep-ph/9409426 \[2\]](#)
- 4 **Linear Sigma Model**  $\bar{l}_1$  and  $\bar{l}_2$  in contradiction with bounds

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.
- 2 **Model independent** approach  $\implies$  dispersion relations.  
 Dispersion relations and unitarity  $\implies$  positivity conditions.  
 Positivity conditions  $\implies$  **bounds** for  $\bar{l}_1$  and  $\bar{l}_2$  (and  $L_{1,2,3}$ ).
- 3 Other (**less restrictive**) methods are contained in ours.  
[Distler et al '06 hep-ph/0604255 \[1\]](#)  
[Pennington et al '95 hep-ph/9409426 \[2\]](#)
- 4 **Linear Sigma Model**  $\bar{l}_1$  and  $\bar{l}_2$  in contradiction with bounds
- 5 3-flavour  $\chi$ PT  $\implies$  generalization for  $SU(3)_V$  breaking.

# Motivations

- 1 Beyond tree-level  $\chi$ PT has undetermined LECs. Estimates (resonance saturation) and fitted values have big errors.
- 2 **Model independent** approach  $\implies$  dispersion relations. Dispersion relations and unitarity  $\implies$  positivity conditions. Positivity conditions  $\implies$  **bounds** for  $\bar{l}_1$  and  $\bar{l}_2$  (and  $L_{1,2,3}$ ).
- 3 Other (**less restrictive**) methods are contained in ours.  
 Distler et al '06 hep-ph/0604255 [1]  
 Pennington et al '95 hep-ph/9409426 [2]
- 4 **Linear Sigma Model**  $\bar{l}_1$  and  $\bar{l}_2$  in contradiction with bounds
- 5 3-flavour  $\chi$ PT  $\implies$  generalization for  $SU(3)_V$  breaking.
- 6  $SU(2) \rightarrow$  A.Manohar & V.M. in preparation [3]  
 $SU(3) \rightarrow$  V.M. work in preparation [4]

# Outline

- 1 Motivations
- 2  **$SU(2)$** 
  - **$\pi\pi$  scattering**
  - Fixed  $t$  dispersion relations & positivity conditions
  - Bounds on chiral LECs and the Linear Sigma Model
  - Equivalence with Pennington & Portoles
- 3  $SU(3)$ 
  - $SU(3)_V$  limit
  - Symmetry breaking
  - Results
- 4 Conclusions

## $\pi\pi$ scattering : generalities

Symmetry constrains

$SU(2)_V \Rightarrow$  only **three** independent amplitudes  $l = 0, 1, 2$ .

## $\pi\pi$ scattering : generalities

### Symmetry constrains

$SU(2)_V \Rightarrow$  only **three** independent amplitudes  $l = 0, 1, 2$ .

Crossing  $\Rightarrow$  only **one** independent function

# $\pi\pi$ scattering : generalities

## Symmetry constrains

$SU(2)_V \Rightarrow$  only **three** independent amplitudes  $l = 0, 1, 2$ .

Crossing  $\Rightarrow$  only **one** independent function

## Chew–Mandelstam representation

$$T(ab \rightarrow cd) = A(s, t) \delta^{ab} \delta^{cd} + A(t, s) \delta^{ac} \delta^{bd} + A(u, t) \delta^{ad} \delta^{bc}$$

with  $A(x, y) = A(x, 4m_\pi^2 - x - y)$ .

# $\pi\pi$ scattering : generalities

## Symmetry constrains

$SU(2)_V \Rightarrow$  only **three** independent amplitudes  $l = 0, 1, 2$ .

Crossing  $\Rightarrow$  only **one** independent function

## Chew–Mandelstam representation

$$T(ab \rightarrow cd) = A(s, t) \delta^{ab} \delta^{cd} + A(t, s) \delta^{ac} \delta^{bd} + A(u, t) \delta^{ad} \delta^{bc}$$

with  $A(x, y) = A(x, 4m_\pi^2 - x - y)$ .

Isospin amplitudes

$$T^3(s, t) = 3A(s, t) + A(t, s) + A(u, s), \quad T^{1,2}(s, t) = A(t, s) \pm A(u, s)$$

# $\pi\pi$ scattering : generalities

## Symmetry constrains

$SU(2)_V \Rightarrow$  only **three** independent amplitudes  $l = 0, 1, 2$ .

Crossing  $\Rightarrow$  only **one** independent function

## Chew–Mandelstam representation

$$T(ab \rightarrow cd) = A(s, t) \delta^{ab} \delta^{cd} + A(t, s) \delta^{ac} \delta^{bd} + A(u, t) \delta^{ad} \delta^{bc}$$

with  $A(x, y) = A(x, 4m_\pi^2 - x - y)$ .

Isospin amplitudes

$$T^3(s, t) = 3A(s, t) + A(t, s) + A(u, s), \quad T^{1,2}(s, t) = A(t, s) \pm A(u, s)$$

Then we can write

$$T^I(s, t) = C_u^{II'} T^{I'}(u, t), \quad C_u^{II'} C_u^{I'J} = \delta_{IJ}, \quad C_u = \frac{1}{6} \begin{pmatrix} 2 & -6 & 10 \\ -2 & 3 & 5 \\ 2 & 3 & 1 \end{pmatrix},$$

# $\pi\pi$ scattering : generalities

## Symmetry constrains

$SU(2)_V \Rightarrow$  only **three** independent amplitudes  $l = 0, 1, 2$ .

Crossing  $\Rightarrow$  only **one** independent function

## Chew–Mandelstam representation

$$T(ab \rightarrow cd) = A(s, t) \delta^{ab} \delta^{cd} + A(t, s) \delta^{ac} \delta^{bd} + A(u, t) \delta^{ad} \delta^{bc}$$

with  $A(x, y) = A(x, 4m_\pi^2 - x - y)$ .

Isospin amplitudes

$$T^3(s, t) = 3A(s, t) + A(t, s) + A(u, s), \quad T^{1,2}(s, t) = A(t, s) \pm A(u, s)$$

Then we can write

$$T^l(s, t) = C_u^{ll'} T^{l'}(u, t), \quad C_u^{ll'} C_u^{l'J} = \delta_{lJ}, \quad C_u = \frac{1}{6} \begin{pmatrix} 2 & -6 & 10 \\ -2 & 3 & 5 \\ 2 & 3 & 1 \end{pmatrix},$$

$$T^l(s, t) = C_t^{ll'} T^{l'}(t, s), \quad C_t^{ll'} C_t^{l'J} = \delta_{lJ}, \quad C_t = \frac{1}{6} \begin{pmatrix} 2 & 6 & 10 \\ 2 & 3 & -5 \\ 2 & -3 & 1 \end{pmatrix}$$

## $\pi\pi$ scattering : analyticity

### s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$

## $\pi\pi$ scattering : analyticity

### s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$
- ◆ **Unitarity**  $\Rightarrow T(s \leq 4m_\pi^2) \in \mathbb{R} \rightarrow$  **single-valued**

## $\pi\pi$ scattering : analyticity

### s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$
- ◆ **Unitarity**  $\Rightarrow T(s \leq 4m_\pi^2) \in \mathbb{R} \rightarrow$  **single-valued**  
for  $s > 4m_\pi^2$   $T(s + i\epsilon) - T(s - i\epsilon) = 2i\text{Im} T(s + i\epsilon) \neq 0 \rightarrow$  **multivalued**

## $\pi\pi$ scattering : analyticity

### s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$
- ◆ **Unitarity**  $\Rightarrow T(s \leq 4m_\pi^2) \in \mathbb{R} \rightarrow$  **single-valued**  
for  $s > 4m_\pi^2$   $T(s+i\epsilon) - T(s-i\epsilon) = 2i\text{Im} T(s+i\epsilon) \neq 0 \rightarrow$  **multivalued**
- ◆ **Branch cut** for  $s > 4m_\pi^2$  ! Rest of singularities and branch points lay on it.

# $\pi\pi$ scattering : analyticity

## s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$
- ◆ **Unitarity**  $\Rightarrow T(s \leq 4m_\pi^2) \in \mathbb{R} \rightarrow$  **single-valued**  
 for  $s > 4m_\pi^2$   $T(s + i\epsilon) - T(s - i\epsilon) = 2i \text{Im} T(s + i\epsilon) \neq 0 \rightarrow$  **multivalued**
- ◆ **Branch cut** for  $s > 4m_\pi^2$  ! Rest of singularities and branch points lay on it.
- ◆ Remaining branch cuts obtained by crossing.

## $\pi\pi$ scattering : analyticity

### s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$
  - ◆ **Unitarity**  $\Rightarrow T(s \leq 4m_\pi^2) \in \mathbb{R} \rightarrow$  **single-valued**  
 for  $s > 4m_\pi^2$   $T(s + i\epsilon) - T(s - i\epsilon) = 2i \text{Im} T(s + i\epsilon) \neq 0 \rightarrow$  **multivalued**
  - ◆ **Branch cut** for  $s > 4m_\pi^2$  ! Rest of singularities and branch points lay on it.
  - ◆ Remaining branch cuts obtained by crossing.
- **Analytic region**  $\Rightarrow s, t, u \leq 4m_\pi^2 \rightarrow$  Dispersion relations

## $\pi\pi$ scattering : analyticity

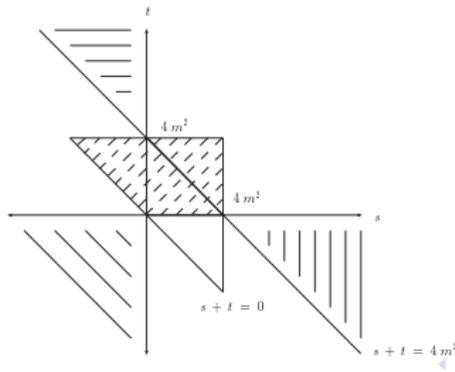
### s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$
- ◆ **Unitarity**  $\Rightarrow T(s \leq 4m_\pi^2) \in \mathbb{R} \rightarrow$  **single-valued**  
 for  $s > 4m_\pi^2$   $T(s + i\epsilon) - T(s - i\epsilon) = 2i \text{Im} T(s + i\epsilon) \neq 0 \rightarrow$  **multivalued**
- ◆ **Branch cut** for  $s > 4m_\pi^2$  ! Rest of singularities and branch points lay on it.
- ◆ Remaining branch cuts obtained by crossing.
  - **Analytic region**  $\Rightarrow s, t, u \leq 4m_\pi^2 \rightarrow$  Dispersion relations
  - Mandelstam triangle  $\Rightarrow 0 \leq s, t, u \leq 4m_\pi^2$  **wrongly** assumed in [1]

# $\pi\pi$ scattering : analyticity

## s channel

- ◆ No other state lighter than the pion. No lighter intermediate state  $m_\rho > 2m_\pi$
  - ◆ **Unitarity**  $\Rightarrow T(s \leq 4m_\pi^2) \in \mathbb{R} \rightarrow$  **single-valued**  
 for  $s > 4m_\pi^2$   $T(s + i\epsilon) - T(s - i\epsilon) = 2i \text{Im} T(s + i\epsilon) \neq 0 \rightarrow$  **multivalued**
  - ◆ **Branch cut** for  $s > 4m_\pi^2$  ! Rest of singularities and branch points lay on it.
  - ◆ Remaining branch cuts obtained by crossing.
- **Analytic region**  $\Rightarrow s, t, u \leq 4m_\pi^2 \rightarrow$  Dispersion relations
  - Mandelstam triangle  $\Rightarrow 0 \leq s, t, u \leq 4m_\pi^2$  **wrongly** assumed in [1]



# Outline

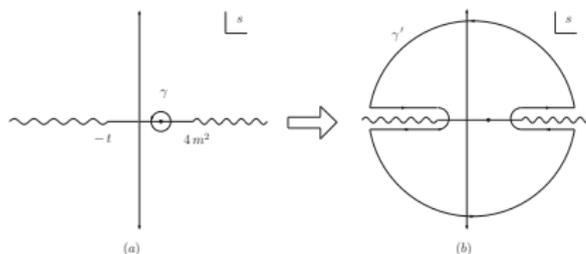
- 1 Motivations
- 2  **$SU(2)$** 
  - $\pi\pi$  scattering
  - **Fixed  $t$  dispersion relations & positivity conditions**
  - Bounds on chiral LECs and the Linear Sigma Model
  - Equivalence with Pennington & Portoles
- 3  $SU(3)$ 
  - $SU(3)_V$  limit
  - Symmetry breaking
  - Results
- 4 Conclusions

## Fixed t dispersion relations

For  $t \leq 4m_\pi^2$  and  $s \notin$  branch cut  $\rightarrow T^I(s, t) = \frac{1}{2\pi i} \oint_\gamma dx \frac{T^I(x, t)}{x-s}$

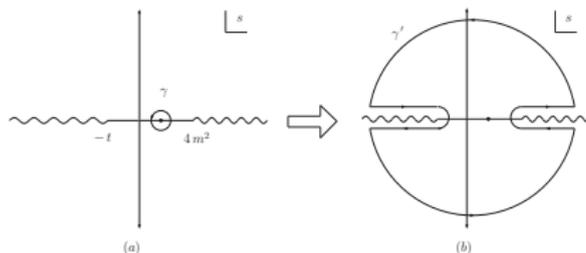
# Fixed $t$ dispersion relations

For  $t \leq 4m_\pi^2$  and  $s \notin$  branch cut  $\rightarrow T^I(s, t) = \frac{1}{2\pi i} \oint_{\gamma'} dx \frac{T^I(x, t)}{x-s}$



# Fixed $t$ dispersion relations

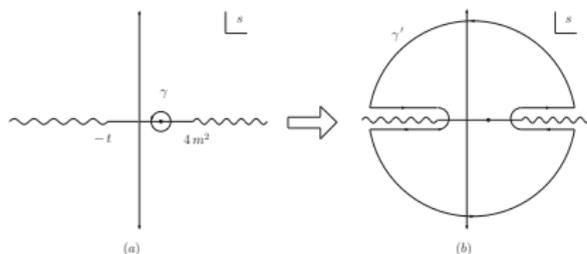
For  $t \leq 4m_\pi^2$  and  $s \notin$  branch cut  $\rightarrow T^I(s, t) = \frac{1}{2\pi i} \oint_\gamma dx \frac{T^I(x, t)}{x-s}$



if  $\oint_\gamma dx \frac{T^I(x, t)}{x-s} \neq 0$  “subtract”  $\frac{d^n}{ds^n} T^I(s, t) = \frac{n!}{2\pi i} \oint_\gamma dx \frac{T^I(x, t)}{(x-s)^{n+1}}$

# Fixed t dispersion relations

For  $t \leq 4m_\pi^2$  and  $s \notin$  branch cut  $\rightarrow T^I(s, t) = \frac{1}{2\pi i} \oint_\gamma dx \frac{T^I(x, t)}{x-s}$



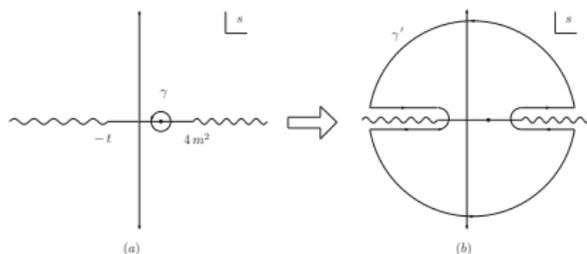
if  $\oint_\circ dx \frac{T^I(x, t)}{x-s} \neq 0$  “subtract”  $\frac{d^n}{ds^n} T^I(s, t) = \frac{n!}{2\pi i} \oint_\gamma dx \frac{T^I(x, t)}{(x-s)^{n+1}}$

Froissart bound  $\rightarrow n = 2$

$$\frac{d^2}{ds^2} T^I(s, t) = \frac{2}{\pi} \int_{4m_\pi^2}^{\infty} dx \left[ \frac{\delta''}{(x-s)^3} + \frac{C_u''}{(x-u)^3} \right] \text{Im } T^I(x + i\epsilon, t)$$

# Fixed t dispersion relations

For  $t \leq 4m_\pi^2$  and  $s \notin$  branch cut  $\rightarrow T^I(s, t) = \frac{1}{2\pi i} \oint_\gamma dx \frac{T^I(x, t)}{x-s}$



if  $\oint_\circ dx \frac{T^I(x, t)}{x-s} \neq 0$  “subtract”  $\frac{d^n}{ds^n} T^I(s, t) = \frac{n!}{2\pi i} \oint_\gamma dx \frac{T^I(x, t)}{(x-s)^{n+1}}$

Froissart bound  $\rightarrow n = 2$

$$\frac{d^2}{ds^2} T^I(s, t) = \frac{2}{\pi} \int_{4m_\pi^2}^{\infty} dx \left[ \frac{\delta^{II'}}{(x-s)^3} + \frac{C_u^{II'}}{(x-u)^3} \right] \text{Im } T^I(x + i\epsilon, t)$$

For  $s + t \geq 0$   $s \leq 4m_\pi^2$  both denominators  $\geq 0$  in the integral path

## Positivity conditions

- Partial wave expansion  $T'(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f'_{\ell}(s) P_{\ell}\left(1 + \frac{2t}{s - 4m_{\pi}^2}\right)$

## Positivity conditions

- Partial wave expansion  $T'(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f'_{\ell}(s) P_{\ell}\left(1 + \frac{2t}{s-4m_{\pi}^2}\right)$
- Optical theorem  $\Rightarrow \text{Im } f'_{\ell}(s) = s\beta(s)\sigma'_{\ell}(s)\theta(s - m_{\pi}^2) \geq 0$

## Positivity conditions

- **Partial wave expansion**  $T'(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f'_{\ell}(s) P_{\ell}\left(1 + \frac{2t}{s-4m_{\pi}^2}\right)$
- **Optical theorem**  $\Rightarrow \text{Im } f'_{\ell}(s) = s\beta(s)\sigma'_{\ell}(s)\theta(s - m_{\pi}^2) \geq 0$
- $P_{\ell}(z) \geq 1$  for  $z \geq 1$  for all  $\ell$

## Positivity conditions

- **Partial wave expansion**  $T'(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f'_{\ell}(s) P_{\ell}\left(1 + \frac{2t}{s - 4m_{\pi}^2}\right)$
- **Optical theorem**  $\Rightarrow \text{Im } f'_{\ell}(s) = s \beta(s) \sigma'_{\ell}(s) \theta(s - m_{\pi}^2) \geq 0$
- $P_{\ell}(z) \geq 1$  for  $z \geq 1$  for all  $\ell \Rightarrow$  If  $t \geq 0$  &  $s \geq 4m_{\pi}^2$  then  $z \geq 1$

## Positivity conditions

- **Partial wave expansion**  $T^l(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}^l(s) P_{\ell} \left( 1 + \frac{2t}{s-4m_{\pi}^2} \right)$
- **Optical theorem**  $\Rightarrow \text{Im } f_{\ell}^l(s) = s \beta(s) \sigma_{\ell}^l(s) \theta(s - m_{\pi}^2) \geq 0$
- $P_{\ell}(z) \geq 1$  for  $z \geq 1$  for all  $\ell \Rightarrow$  If  $t \geq 0$  &  $s \geq 4m_{\pi}^2$  then  $z \geq 1$   
 $\text{Im } T^l(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) s \beta(s) \sigma_{\ell}^l(s) P_{\ell} \left( 1 + \frac{2t}{s-4m_{\pi}^2} \right) \geq 0$

## Positivity conditions

- **Partial wave expansion**  $T^I(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}^I(s) P_{\ell} \left( 1 + \frac{2t}{s-4m_{\pi}^2} \right)$
- **Optical theorem**  $\Rightarrow \text{Im } f_{\ell}^I(s) = s \beta(s) \sigma_{\ell}^I(s) \theta(s - m_{\pi}^2) \geq 0$
- $P_{\ell}(z) \geq 1$  for  $z \geq 1$  for all  $\ell \Rightarrow$  If  $t \geq 0$  &  $s \geq 4m_{\pi}^2$  then  $z \geq 1$   
 $\text{Im } T^I(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) s \beta(s) \sigma_{\ell}^I(s) P_{\ell} \left( 1 + \frac{2t}{s-4m_{\pi}^2} \right) \geq 0$
- For certain  $\sum a_I T^I$  with  $a_I \geq 0 \rightarrow \sum a_I C_U^{IJ} T_J = \sum_K b_K T_K$  with  $b_K \geq 0$   
 They correspond to **physical processes with equal initial and final state.**

## Positivity conditions

- **Partial wave expansion**  $T^I(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) f_{\ell}^I(s) P_{\ell} \left( 1 + \frac{2t}{s-4m_{\pi}^2} \right)$
- **Optical theorem**  $\Rightarrow \text{Im } f_{\ell}^I(s) = s \beta(s) \sigma_{\ell}^I(s) \theta(s - m_{\pi}^2) \geq 0$
- $P_{\ell}(z) \geq 1$  for  $z \geq 1$  for all  $\ell \Rightarrow$  If  $t \geq 0$  &  $s \geq 4m_{\pi}^2$  then  $z \geq 1$   
 $\text{Im } T^I(s, t) = \sum_{\ell=0}^{\infty} (2\ell + 1) s \beta(s) \sigma_{\ell}^I(s) P_{\ell} \left( 1 + \frac{2t}{s-4m_{\pi}^2} \right) \geq 0$
- For certain  $\sum a_I T^I$  with  $a_I \geq 0 \rightarrow \sum a_I C_U^{IJ} T_J = \sum_K b_K T_K$  with  $b_K \geq 0$   
 They correspond to **physical processes with equal initial and final state.**

**Positivity conditions :**

Inside the region  $\mathcal{A} \equiv \{s \leq 4m_{\pi}^2, 0 \leq t \leq 4m^2 \text{ \& } s + t \geq 0\}$

$$\frac{d^2}{ds^2} T(\pi^0 \pi^0 \rightarrow \pi^0 \pi^0) [(s, t) \in \mathcal{A}] \geq 0, \quad \frac{d^2}{ds^2} T(\pi^+ \pi^+ \rightarrow \pi^+ \pi^+) [(s, t) \in \mathcal{A}] \geq 0,$$

$$\frac{d^2}{ds^2} T(\pi^+ \pi^0 \rightarrow \pi^+ \pi^0) [(s, t) \in \mathcal{A}] \geq 0,$$

# Outline

- 1 Motivations
- 2 **SU(2)**
  - $\pi\pi$  scattering
  - Fixed t dispersion relations & positivity conditions
  - **Bounds on chiral LECs and the Linear Sigma Model**
  - Equivalence with Pennington & Portoles
- 3 SU(3)
  - $SU(3)_V$  limit
  - Symmetry breaking
  - Results
- 4 Conclusions

## Bounds on chiral LECs

In the region  $\mathcal{A}$  we can apply  $\chi$ PT at  $\mathcal{O}(p^4)$  to obtain  $[\frac{d^2}{ds^2}\mathcal{O}(p^2) = 0]$

$$\sum_{i=1}^2 \alpha_{ji} \bar{l}_i - f_j[(\mathbf{s}, t) \in \mathcal{A}] \geq 0 \quad \implies \quad \sum_{i=1}^2 \alpha_{ji} \bar{l}_i \geq f_j[(\mathbf{s}, t) \in \mathcal{A}]_{\max}$$

## Bounds on chiral LECs

In the region  $\mathcal{A}$  we can apply  $\chi$ PT at  $\mathcal{O}(p^4)$  to obtain  $[\frac{d^2}{ds^2}\mathcal{O}(p^2) = 0]$

$$\sum_{i=1}^2 \alpha_{ji} \bar{l}_i - f_j[(s, t) \in \mathcal{A}] \geq 0 \implies \sum_{i=1}^2 \alpha_{ji} \bar{l}_i \geq f_j[(s, t) \in \mathcal{A}]_{\max}$$

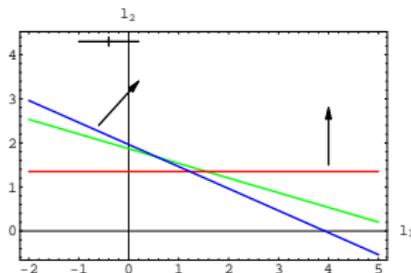
Process	LECs combination	Bound	Experimental value
$\pi^0\pi^0 \rightarrow \pi^0\pi^0$	$\bar{l}_1 + 2\bar{l}_2$ [1,2]	$\geq \frac{157}{40} = 3.925$	$8.2 \pm 0.6$
$\pi^+\pi^0 \rightarrow \pi^+\pi^0$	$\bar{l}_2$ [1,2]	$\geq \frac{27}{20} = 1.350$	$4.3 \pm 0.1$
$\pi^+\pi^+ \rightarrow \pi^+\pi^+$	$\bar{l}_1 + 3\bar{l}_2$ [3]	$\geq 5.604$	$12.5 \pm 0.7$

# Bounds on chiral LECs

In the region  $\mathcal{A}$  we can apply  $\chi$ PT at  $\mathcal{O}(p^4)$  to obtain  $[\frac{d^2}{ds^2}\mathcal{O}(p^2) = 0]$

$$\sum_{i=1}^2 \alpha_{ji} \bar{l}_i - f_j[(s, t) \in \mathcal{A}] \geq 0 \implies \sum_{i=1}^2 \alpha_{ji} \bar{l}_i \geq f_j[(s, t) \in \mathcal{A}]_{\max}$$

Process	LECs combination	Bound	Experimental value
$\pi^0\pi^0 \rightarrow \pi^0\pi^0$	$\bar{l}_1 + 2\bar{l}_2$ [1,2]	$\geq \frac{157}{40} = 3.925$	$8.2 \pm 0.6$
$\pi^+\pi^0 \rightarrow \pi^+\pi^0$	$\bar{l}_2$ [1,2]	$\geq \frac{27}{20} = 1.350$	$4.3 \pm 0.1$
$\pi^+\pi^+ \rightarrow \pi^+\pi^+$	$\bar{l}_1 + 3\bar{l}_2$ [3]	$\geq 5.604$	$12.5 \pm 0.7$



# Is the Linear Sigma Model consistent?

◆ Functional integration of  $\sigma$  particle  $\implies \bar{l}_1$  and  $\bar{l}_2$  in LSM : (at one-loop)

$$\bar{l}_1 = \frac{24\pi^2}{g} + \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{35}{6}, \quad \bar{l}_2 = \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{11}{6} \quad \text{Gasser et al '84}$$

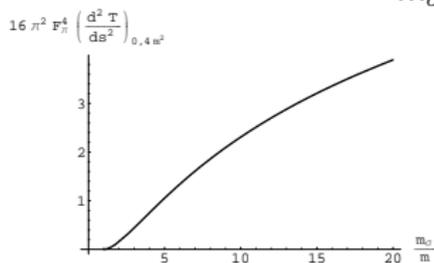
# Is the Linear Sigma Model consistent?

- ◆ **Functional integration** of  $\sigma$  particle  $\implies \bar{l}_1$  and  $\bar{l}_2$  in LSM : (at one-loop)
- $\bar{l}_1 = \frac{24\pi^2}{g} + \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{35}{6}$ ,  $\bar{l}_2 = \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{11}{6}$  Gasser et al '84
- ◆ Plugging into the second bound  $\log\left(\frac{m_\sigma}{m_\pi}\right) \geq \frac{191}{60}$  **violated for  $m_\sigma \lesssim 24 m_\pi$  !!!**



# Is the Linear Sigma Model consistent?

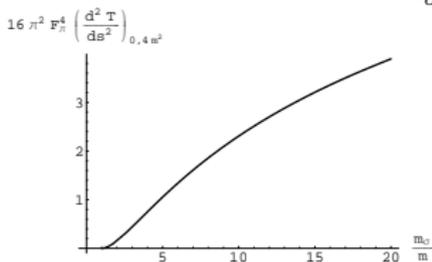
- ◆ **Functional integration** of  $\sigma$  particle  $\implies \bar{l}_1$  and  $\bar{l}_2$  in LSM : (at one-loop)
- $\bar{l}_1 = \frac{24\pi^2}{g} + \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{35}{6}$ ,  $\bar{l}_2 = \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{11}{6}$  Gasser et al '84
- ◆ Plugging into the second bound  $\log\left(\frac{m_\sigma}{m_\pi}\right) \geq \frac{191}{60}$  **violated for  $m_\sigma \lesssim 24 m_\pi$  !!!**
- ◆ But LSM consistent if  $m_\sigma \geq \sqrt{3} m_\pi$ . Apply directly positivity conditions.
- ◆ Integration of  $\sigma$  tantamount to  $\frac{1}{m_\sigma^2}$  expansion. To all orders we have :



It is **consistent** even for  $m_\sigma < m_\pi$

# Is the Linear Sigma Model consistent?

- ◆ **Functional integration** of  $\sigma$  particle  $\implies \bar{l}_1$  and  $\bar{l}_2$  in LSM : (at one-loop)  
 $\bar{l}_1 = \frac{24\pi^2}{g} + \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{35}{6}$ ,  $\bar{l}_2 = \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{11}{6}$  Gasser et al '84
- ◆ Plugging into the second bound  $\log\left(\frac{m_\sigma}{m_\pi}\right) \geq \frac{191}{60}$  **violated for  $m_\sigma \lesssim 24 m_\pi$  !!!**
- ◆ But LSM consistent if  $m_\sigma \geq \sqrt{3} m_\pi$ . Apply directly positivity conditions.
- ◆ Integration of  $\sigma$  tantamount to  $\frac{1}{m_\sigma^2}$  expansion. To all orders we have :

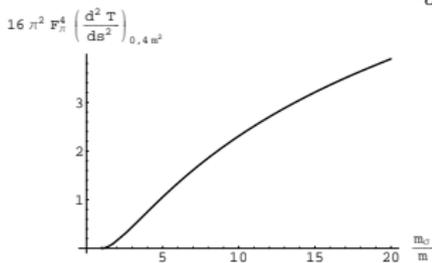


It is **consistent** even for  $m_\sigma < m_\pi$

↪ But the **non-Linear Sigma Model** is **inconsistent** for  $m_\sigma < 24m_\pi$

# Is the Linear Sigma Model consistent?

- ◆ **Functional integration** of  $\sigma$  particle  $\implies \bar{l}_1$  and  $\bar{l}_2$  in LSM : (at one-loop)  
 $\bar{l}_1 = \frac{24\pi^2}{g} + \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{35}{6}$ ,  $\bar{l}_2 = \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{11}{6}$  Gasser et al '84
- ◆ Plugging into the second bound  $\log\left(\frac{m_\sigma}{m_\pi}\right) \geq \frac{191}{60}$  **violated for  $m_\sigma \lesssim 24 m_\pi$  !!!**
- ◆ But LSM consistent if  $m_\sigma \geq \sqrt{3} m_\pi$ . Apply directly positivity conditions.
- ◆ Integration of  $\sigma$  tantamount to  $\frac{1}{m_\sigma^2}$  expansion. To all orders we have :



It is **consistent** even for  $m_\sigma < m_\pi$

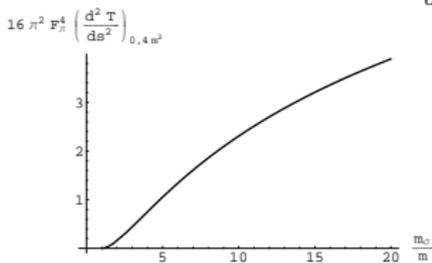
$\leadsto$  But the **non-Linear Sigma Model** is **inconsistent** for  $m_\sigma < 24m_\pi$

$\leadsto$  **Caveat!** For those who integrate  $\rho$  (as I do) for estimating chiral LECs

$m_\rho/m_\pi \ll 25$  (!).

# Is the Linear Sigma Model consistent?

- ◆ **Functional integration** of  $\sigma$  particle  $\implies \bar{l}_1$  and  $\bar{l}_2$  in LSM : (at one-loop)  
 $\bar{l}_1 = \frac{24\pi^2}{g} + \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{35}{6}$ ,  $\bar{l}_2 = \log\left(\frac{m_\sigma}{m_\pi}\right) - \frac{11}{6}$  Gasser et al '84
- ◆ Plugging into the second bound  $\log\left(\frac{m_\sigma}{m_\pi}\right) \geq \frac{191}{60}$  **violated for  $m_\sigma \lesssim 24 m_\pi$  !!!**
- ◆ But LSM consistent if  $m_\sigma \geq \sqrt{3} m_\pi$ . Apply directly positivity conditions.
- ◆ Integration of  $\sigma$  tantamount to  $\frac{1}{m_\sigma^2}$  expansion. To all orders we have :



It is **consistent** even for  $m_\sigma < m_\pi$

- ↪ But the **non-Linear Sigma Model** is **inconsistent** for  $m_\sigma < 24 m_\pi$
- ↪ **Caveat!** For those who integrate  $\rho$  (as I do) for estimating chiral LECs  $m_\rho/m_\pi \ll 25$  (!). But in this case LECs are generated at tree level.

# Outline

- 1 Motivations
- 2 **SU(2)**
  - $\pi\pi$  scattering
  - Fixed t dispersion relations & positivity conditions
  - Bounds on chiral LECs and the Linear Sigma Model
  - **Equivalence with Pennington & Portoles**
- 3 SU(3)
  - $SU(3)_V$  limit
  - Symmetry breaking
  - Results
- 4 Conclusions

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^l \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^l(s)}{(\frac{s}{4} - m^2)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^I \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^I(s)}{\left(\frac{s}{4} - m^2\right)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

P & P [2] quote  $a_2^0 + 2a_2^2 \geq 0, \quad a_2^0 - a_2^2 \geq 0 \quad (1)$

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^l \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^l(s)}{(\frac{s}{4} - m^2)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

P & P [2] quote  $a_2^0 + 2a_2^2 \geq 0, \quad a_2^0 - a_2^2 \geq 0 \quad (1)$

But in fact 
$$a_\ell^l = \frac{4^\ell \ell!}{(2\ell + 1)} C_t^{ll'} \left. \frac{d^\ell F^{ll'}(s, 4m^2)}{d s^\ell} \right|_{s=0} \quad [3]$$

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^l \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^l(s)}{\left(\frac{s}{4} - m^2\right)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

P & P [2] quote  $a_2^0 + 2a_2^2 \geq 0$ ,  $a_2^0 - a_2^2 \geq 0$  (1)

But in fact  $a_\ell^l = \frac{4^\ell \ell!}{(2\ell + 1)} C_t^{ll'} \left. \frac{d^\ell F^{ll'}(s, 4m^2)}{d s^\ell} \right|_{s=0}$  [3]

Since  $(s = 0, t = 4m^2) \in \mathcal{A}$  appropriate linear combinations  $\rightarrow$  (1)

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^l \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^l(s)}{\left(\frac{s}{4} - m^2\right)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

P & P [2] quote  $a_2^0 + 2a_2^2 \geq 0, \quad a_2^0 - a_2^2 \geq 0 \quad (1)$

But in fact  $a_\ell^l = \frac{4^\ell \ell!}{(2\ell + 1)} C_t^{ll'} \left. \frac{d^\ell F^{ll'}(s, 4m^2)}{ds^\ell} \right|_{s=0} \quad [3]$

Since  $(s = 0, t = 4m^2) \in \mathcal{A}$  appropriate linear combinations  $\rightarrow (1)$

- 1 They **only consider a single point** in  $\mathcal{A} \rightarrow$  less general, no necessarily the most stringent. No way to find third bound.

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^l \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^l(s)}{\left(\frac{s}{4} - m^2\right)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

P & P [2] quote  $a_2^0 + 2a_2^2 \geq 0$ ,  $a_2^0 - a_2^2 \geq 0$  (1)

But in fact  $a_\ell^l = \frac{4^\ell \ell!}{(2\ell + 1)} C_t^{ll'} \left. \frac{d^\ell F^{ll'}(s, 4m^2)}{ds^\ell} \right|_{s=0}$  [3]

Since  $(s = 0, t = 4m^2) \in \mathcal{A}$  appropriate linear combinations  $\rightarrow$  (1)

- 1 They **only consider a single point** in  $\mathcal{A} \rightarrow$  less general, no necessarily the most stringent. No way to find third bound.
- 2 We remind that in Ref. [1] **only the Mandelstam** triangle was taken into account. No way to find third bound.

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^j \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^j(s)}{\left(\frac{s}{4} - m^2\right)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

P & P [2] quote  $a_2^0 + 2a_2^2 \geq 0$ ,  $a_2^0 - a_2^2 \geq 0$  (1)

But in fact  $a_\ell^j = \frac{4^\ell \ell!}{(2\ell + 1)} C_t^{j\ell} \left. \frac{d^\ell F^j(s, 4m^2)}{ds^\ell} \right|_{s=0}$  [3]

Since  $(s = 0, t = 4m^2) \in \mathcal{A}$  appropriate linear combinations  $\rightarrow$  (1)

- 1 They **only consider a single point** in  $\mathcal{A} \rightarrow$  less general, no necessarily the most stringent. No way to find third bound.
- 2 We remind that in Ref. [1] **only the Mandelstam** triangle was taken into account. No way to find third bound.
- 3  $m_K/m_\pi \sim 3.5 \ll 24$  **Integrate K** in  $SU(3) \rightarrow SU(2)$ ? Yes...  
 bad results for  $\bar{l}_{1,2}$  but within bounds

# Equivalence with Pennington & Portoles[2]

## Scattering lengths

$$a_\ell^l \equiv \lim_{s \rightarrow 4m^2} \frac{f_\ell^l(s)}{\left(\frac{s}{4} - m^2\right)^\ell} \quad \text{and Bose symmetry implies} \quad a_{2k}^1 \equiv 0$$

P & P [2] quote  $a_2^0 + 2a_2^2 \geq 0$ ,  $a_2^0 - a_2^2 \geq 0$  (1)

But in fact  $a_\ell^l = \frac{4^\ell \ell!}{(2\ell + 1)} C_t^{ll'} \left. \frac{d^\ell F^{ll'}(s, 4m^2)}{ds^\ell} \right|_{s=0}$  [3]

Since  $(s = 0, t = 4m^2) \in \mathcal{A}$  appropriate linear combinations  $\rightarrow$  (1)

- 1 They **only consider a single point** in  $\mathcal{A} \rightarrow$  less general, no necessarily the most stringent. No way to find third bound.
- 2 We remind that in Ref. [1] **only the Mandelstam** triangle was taken into account. No way to find third bound.
- $m_K/m_\pi \sim 3.5 \ll 24$  **Integrate K** in  $SU(3) \rightarrow SU(2)$ ? Yes...  
 bad results for  $\bar{l}_{1,2}$  but within bounds
- **SU(3)  $\chi$ PT** ? Consistent with axiomatic principles?

# Outline

- 1 Motivations
- 2  $SU(2)$ 
  - $\pi\pi$  scattering
  - Fixed  $t$  dispersion relations & positivity conditions
  - Bounds on chiral LECs and the Linear Sigma Model
  - Equivalence with Pennington & Portoles
- 3  $SU(3)$ 
  - $SU(3)_V$  limit
  - Symmetry breaking
  - Results
- 4 Conclusions

# SU(3)<sub>V</sub> limit (LECs independent of $m_q$ )

## Octet-to-octet scattering

Missmatch **Clebsch-Gordan** ↔ **tensor analysis**

$$8 \otimes 8 = 27 \oplus 10 \oplus 10^* \oplus 8_1 \oplus 8_2 \oplus 1$$

# SU(3)<sub>V</sub> limit (LECs independent of $m_q$ )

## Octet-to-octet scattering

Missmatch **Clebsch-Gordan** ↔ **tensor analysis**

$$8 \otimes 8 = 27 \oplus 10 \oplus 10^* \oplus 8_1 \oplus 8_2 \oplus 1$$

$$T(ab \rightarrow cd) = A_1(s, t, u) \delta^{ab} \delta^{cd} + A_2(s, t, u) \delta^{ac} \delta^{bd} + A_3(s, t, u) \delta^{ad} \delta^{bc} \\ + B_1(s, t, u) d^{abe} d^{cde} + B_2(s, t, u) d^{ace} d^{bde}$$

# SU(3)<sub>V</sub> limit (LECs independent of $m_q$ )

## Octet-to-octet scattering

Missmatch **Clebsch-Gordan** ↔ **tensor analysis**

$$8 \otimes 8 = 27 \oplus 10 \oplus 10^* \oplus 8_1 \oplus 8_2 \oplus 1$$

$$T(ab \rightarrow cd) = A_1(s, t, u) \delta^{ab} \delta^{cd} + A_2(s, t, u) \delta^{ac} \delta^{bd} + A_3(s, t, u) \delta^{ad} \delta^{bc} \\ + B_1(s, t, u) d^{abe} d^{cde} + B_2(s, t, u) d^{ace} d^{bde}$$

Crossing symmetry  $\Rightarrow T_{10}(s, t) = T_{10^*}(s, t)$

# SU(3)<sub>V</sub> limit (LECs independent of $m_q$ )

## Octet-to-octet scattering

Missmatch **Clebsch-Gordan** ↔ **tensor analysis**

$$8 \otimes 8 = 27 \oplus 10 \oplus 10^* \oplus 8_1 \oplus 8_2 \oplus 1$$

$$T(ab \rightarrow cd) = A_1(s, t, u) \delta^{ab} \delta^{cd} + A_2(s, t, u) \delta^{ac} \delta^{bd} + A_3(s, t, u) \delta^{ad} \delta^{bc} \\ + B_1(s, t, u) d^{abe} d^{cde} + B_2(s, t, u) d^{ace} d^{bde}$$

Crossing symmetry  $\Rightarrow T_{10}(s, t) = T_{10^*}(s, t)$

Analogously to SU(2)  $I, J = 1, 8_1, 8_2, 10, 27$  (no isospin!)

$$\frac{d^2}{ds^2} T^I(s, t) = \frac{2}{\pi} \int_{4m^2}^{\infty} dx \left[ \frac{\delta^{II'}}{(x-s)^3} + \frac{C_u^{II'}}{(x-u)^3} \right] \text{Im } T^{I'}(x + i\epsilon, t)$$

# SU(3)<sub>V</sub> limit (LECs independent of $m_q$ )

## Octet-to-octet scattering

Missmatch **Clebsch-Gordan** ↔ **tensor analysis**

$$8 \otimes 8 = 27 \oplus 10 \oplus 10^* \oplus 8_1 \oplus 8_2 \oplus 1$$

$$T(ab \rightarrow cd) = A_1(s, t, u) \delta^{ab} \delta^{cd} + A_2(s, t, u) \delta^{ac} \delta^{bd} + A_3(s, t, u) \delta^{ad} \delta^{bc} \\ + B_1(s, t, u) d^{abe} d^{cde} + B_2(s, t, u) d^{ace} d^{bde}$$

Crossing symmetry  $\Rightarrow T_{10}(s, t) = T_{10^*}(s, t)$

Analogously to SU(2)  $I, J = 1, 8_1, 8_2, 10, 27$  (no isospin!)

$$\frac{d^2}{ds^2} T^I(s, t) = \frac{2}{\pi} \int_{4m^2}^{\infty} dx \left[ \frac{\delta^{II'}}{(x-s)^3} + \frac{C_U^{II'}}{(x-u)^3} \right] \text{Im } T^{I'}(x + i\epsilon, t)$$

$$\frac{d^2}{ds^2} T(\pi^+\pi^+ \rightarrow \pi^+\pi^+) [(s, t) \in \mathcal{A}] \geq 0, \quad \frac{d^2}{ds^2} T(\pi^0\pi^0 \rightarrow \pi^0\pi^0) [(s, t) \in \mathcal{A}] \geq 0,$$

$$\frac{d^2}{ds^2} T(\pi^+\pi^0 \rightarrow \pi^+\pi^0) [(s, t) \in \mathcal{A}] \geq 0, \quad \frac{d^2}{ds^2} T(\pi\eta \rightarrow \pi\eta) [(s, t) \in \mathcal{A}] \geq 0,$$

$$\frac{d^2}{ds^2} T(K\eta \rightarrow K\eta) [(s, t) \in \mathcal{A}] \geq 0, \quad \frac{d^2}{ds^2} T(K\pi^+ \rightarrow K\pi^+) [(s, t) \in \mathcal{A}] \geq 0,$$

## $SU(3)_V$ limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$

# SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^r(m^2) + \alpha_{2i} L_2^r(m^2) + \alpha_{3i} L_3^r \geq f_i[(\mathbf{s}, t) \in \mathcal{A}]|_{\max}$

## SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^r(m^2) + \alpha_{2i} L_2^r(m^2) + \alpha_{3i} L_3^r \geq f_i[(\mathbf{s}, \mathbf{t}) \in \mathcal{A}]|_{\max}$
- Experimental values for  $L_{1,2}(m_\rho) \Rightarrow$  run down to  $\mu = m$

## SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^r(m^2) + \alpha_{2i} L_2^r(m^2) + \alpha_{3i} L_3^r \geq f_i[(s, t) \in \mathcal{A}]|_{\max}$
- Experimental values for  $L_{1,2}(m_\rho) \Rightarrow$  run down to  $\mu = m$
- Which physical mass corresponds to  $m$ ?  $m_\pi$ ?  $m_K$ ? We take both.

## SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^r(m^2) + \alpha_{2i} L_2^r(m^2) + \alpha_{3i} L_3^r \geq f_i[(\mathbf{s}, t) \in \mathcal{A}]|_{\max}$
- Experimental values for  $L_{1,2}(m_\rho) \Rightarrow$  run down to  $\mu = m$
- Which physical mass corresponds to  $m$ ?  $m_\pi$ ?  $m_K$ ? We take both.
- $m = m_\pi$  not very stringent :-(. .

## SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^i(m^2) + \alpha_{2i} L_2^i(m^2) + \alpha_{3i} L_3^i \geq f_i[(s, t) \in \mathcal{A}]|_{\max}$
- Experimental values for  $L_{1,2}(m_\rho) \Rightarrow$  run down to  $\mu = m$
- Which physical mass corresponds to  $m$ ?  $m_\pi$ ?  $m_K$ ? We take both.
- $m = m_\pi$  not very stringent :-{ .
- $m = m_K$  severely violates bounds !

## SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^r(m^2) + \alpha_{2i} L_2^r(m^2) + \alpha_{3i} L_3^r \geq f_i[(s, t) \in \mathcal{A}]|_{\max}$
- Experimental values for  $L_{1,2}(m_\rho) \Rightarrow$  run down to  $\mu = m$
- Which physical mass corresponds to  $m$ ?  $m_\pi$ ?  $m_K$ ? We take both.
- $m = m_\pi$  not very stringent :- ( .
- $m = m_K$  severely violates bounds !

By including SU(3)<sub>V</sub> symmetry breaking :

- 1 The ambiguity disappears.

## SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^r(m^2) + \alpha_{2i} L_2^r(m^2) + \alpha_{3i} L_3^r \geq f_i[(s, t) \in \mathcal{A}]|_{\max}$
- Experimental values for  $L_{1,2}(m_\rho) \Rightarrow$  run down to  $\mu = m$
- Which physical mass corresponds to  $m$ ?  $m_\pi$ ?  $m_K$ ? We take both.
- $m = m_\pi$  not very stringent :- ( .
- $m = m_K$  severely violates bounds !

By including SU(3)<sub>V</sub> symmetry breaking :

- 1 The ambiguity disappears.
- 2 Bounds tighten

## SU(3)<sub>V</sub> limit (II)

Remarks :

- In  $\chi$ PT we adopt  $\mu = m_\pi = m_K \equiv m$
- $\alpha_{1i} L_1^i(m^2) + \alpha_{2i} L_2^i(m^2) + \alpha_{3i} L_3^i \geq f_i[(s, t) \in \mathcal{A}]|_{\max}$
- Experimental values for  $L_{1,2}(m_\rho) \Rightarrow$  run down to  $\mu = m$
- Which physical mass corresponds to  $m$ ?  $m_\pi$ ?  $m_K$ ? We take both.
- $m = m_\pi$  not very stringent :- ( .
- $m = m_K$  severely violates bounds !

By including SU(3)<sub>V</sub> symmetry breaking :

- 1 The ambiguity disappears.
- 2 Bounds tighten

but...need to reconsider the positivity conditions

# Outline

- 1 Motivations
- 2  $SU(2)$ 
  - $\pi\pi$  scattering
  - Fixed  $t$  dispersion relations & positivity conditions
  - Bounds on chiral LECs and the Linear Sigma Model
  - Equivalence with Pennington & Portoles
- 3  $SU(3)$ 
  - $SU(3)_V$  limit
  - **Symmetry breaking**
  - Results
- 4 Conclusions

# Symmetry breaking

## Analytic region

Consider the process  $a + b \rightarrow a + b$  ( $\text{Im } f_\ell \geq 0$ )  $m_a = M$ ,  $m_b = m$

# Symmetry breaking

## Analytic region

Consider the process  $a + b \rightarrow a + b$  ( $\text{Im } f_\ell \geq 0$ )  $m_a = M$ ,  $m_b = m$

If  $a + b \rightarrow c + d$ ,  $a + \bar{b} \rightarrow e + f$  and  $a + \bar{a} \rightarrow g + h$  exist analytic in

$$s \leq (m_c + m_d)^2, t \leq (m_e + m_f)^2, s + t \geq 2(m^2 + M^2) - (m_g + m_h)^2$$

# Symmetry breaking

## Analytic region

Consider the process  $a + b \rightarrow a + b$  ( $\text{Im } f_\ell \geq 0$ )  $m_a = M$ ,  $m_b = m$

If  $a + b \rightarrow c + d$ ,  $a + \bar{b} \rightarrow e + f$  and  $a + \bar{a} \rightarrow g + h$  exist analytic in

$$s \leq (m_c + m_d)^2, t \leq (m_e + m_f)^2, s + t \geq 2(m^2 + M^2) - (m_g + m_h)^2$$

## Dispersion relation

$$\frac{d^2}{ds^2} T(s, t) = \frac{2}{\pi} \int_{(m_c+m_d)^2}^{\infty} dx \frac{\text{Im } T(x + i\epsilon, t)}{(x - s)^3} + \frac{2}{\pi} \int_{(m_g+m_h)^2}^{\infty} dx \frac{\text{Im } T_u(x + i\epsilon, t)}{(x - u)^3}$$

# Symmetry breaking

## Analytic region

Consider the process  $a + b \rightarrow a + b$  ( $\text{Im } f_\ell \geq 0$ )  $m_a = M$ ,  $m_b = m$

If  $a + b \rightarrow c + d$ ,  $a + \bar{b} \rightarrow e + f$  and  $a + \bar{a} \rightarrow g + h$  exist analytic in

$$s \leq (m_c + m_d)^2, t \leq (m_e + m_f)^2, s + t \geq 2(m^2 + M^2) - (m_g + m_h)^2$$

## Dispersion relation

$$\frac{d^2}{ds^2} T(s, t) = \frac{2}{\pi} \int_{(m_c+m_d)^2}^{\infty} dx \frac{\text{Im } T(x + i\epsilon, t)}{(x - s)^3} + \frac{2}{\pi} \int_{(m_g+m_h)^2}^{\infty} dx \frac{\text{Im } T_u(x + i\epsilon, t)}{(x - u)^3}$$

Denominators positive for  $s \leq (m_c + m_d)^2$ ,  $s + t \geq 2(m^2 + M^2) - (m_g + m_h)^2$

$$P_\ell \left[ 1 + \frac{st}{(s + m^2 - M^2)^2 - 4m^2 s} \right] \geq 0 \text{ in the two integrals}$$

# Symmetry breaking

## Analytic region

Consider the process  $a + b \rightarrow a + b$  ( $\text{Im } f_\ell \geq 0$ )  $m_a = M$ ,  $m_b = m$

If  $a + b \rightarrow c + d$ ,  $a + \bar{b} \rightarrow e + f$  and  $a + \bar{a} \rightarrow g + h$  exist analytic in

$$s \leq (m_c + m_d)^2, t \leq (m_e + m_f)^2, s + t \geq 2(m^2 + M^2) - (m_g + m_h)^2$$

## Dispersion relation

$$\frac{d^2}{ds^2} T(s, t) = \frac{2}{\pi} \int_{(m_c + m_d)^2}^{\infty} dx \frac{\text{Im } T(x + i\epsilon, t)}{(x - s)^3} + \frac{2}{\pi} \int_{(m_g + m_h)^2}^{\infty} dx \frac{\text{Im } T_u(x + i\epsilon, t)}{(x - u)^3}$$

Denominators positive for  $s \leq (m_c + m_d)^2$ ,  $s + t \geq 2(m^2 + M^2) - (m_g + m_h)^2$

$$P_\ell \left[ 1 + \frac{st}{(s + m^2 - M^2)^2 - 4m^2 s} \right] \geq 0 \text{ in the two integrals}$$

$$\frac{st}{(s + m^2 - M^2)^2 - 4m^2 s} \geq 0 \quad \text{for } s \geq (m_c + m_d)^2 [(m_g + m_h)^2]$$

## Symmetry breaking (II)

- $t$  must be positive (if  $s \rightarrow \infty$  back to symmetric case)

## Symmetry breaking (II)

- $t$  must be positive (if  $s \rightarrow \infty$  back to symmetric case)
- $P_\ell \geq 0$  only for  $(M - m)^2 \geq s \geq (M + m)^2$

## Symmetry breaking (II)

- $t$  must be positive (if  $s \rightarrow \infty$  back to symmetric case)
- $P_\ell \geq 0$  only for  $(M - m)^2 \geq s \geq (M + m)^2$

### Theorem

Positivity conditions hold for processes of the type  $a + b \rightarrow a + b$  such that the lightest pair of particles that can arise off the scattering  $a + b$  is precisely  $a + b$ , and analogously for  $a + \bar{b}$ .

## Symmetry breaking (II)

- $t$  must be positive (if  $s \rightarrow \infty$  back to symmetric case)
- $P_\ell \geq 0$  only for  $(M - m)^2 \geq s \geq (M + m)^2$

### Theorem

Positivity conditions hold for processes of the type  $a + b \rightarrow a + b$  such that the lightest pair of particles that can arise off the scattering  $a + b$  is precisely  $a + b$ , and analogously for  $a + \bar{b}$ .

$$\begin{aligned} \frac{d^2}{ds^2} T(\pi^+\pi^+ \rightarrow \pi^+\pi^+) [(s, t) \in \mathcal{A}] &\geq 0, & \frac{d^2}{ds^2} T(\pi^0\pi^0 \rightarrow \pi^0\pi^0) [(s, t) \in \mathcal{A}] &\geq 0, \\ \frac{d^2}{ds^2} T(\pi^+\pi^0 \rightarrow \pi^+\pi^0) [(s, t) \in \mathcal{A}] &\geq 0, & \frac{d^2}{ds^2} T(\pi\eta \rightarrow \pi\eta) [(s, t) \in \mathcal{A}] &\geq 0, \\ \frac{d^2}{ds^2} T(K\pi^+ \rightarrow K\pi^+) [(s, t) \in \mathcal{A}] &\geq 0, & & \end{aligned}$$

# Outline

- 1 Motivations
- 2  $SU(2)$ 
  - $\pi\pi$  scattering
  - Fixed  $t$  dispersion relations & positivity conditions
  - Bounds on chiral LECs and the Linear Sigma Model
  - Equivalence with Pennington & Portoles
- 3  $SU(3)$ 
  - $SU(3)_V$  limit
  - Symmetry breaking
  - **Results**
- 4 Conclusions

# Results

	$10^3 \times L_1^r(\mu)$	$10^3 \times L_2^r(\mu)$	$10^3 \times L_3$
$\mu = m_\rho$	$0.43 \pm 0.12$	$0.43 \pm 0.12$	$-2.35 \pm 0.37$
$\mu = m_K$	$0.69 \pm 0.12$	$1.26 \pm 0.12$	$-2.35 \pm 0.37$
$\mu = m_\pi$	$2.78 \pm 0.12$	$1.26 \pm 0.12$	$-2.35 \pm 0.37$

# Results

	$10^3 \times L_1^r(\mu)$	$10^3 \times L_2^r(\mu)$	$10^3 \times L_3$
$\mu = m_\rho$	$0.43 \pm 0.12$	$0.43 \pm 0.12$	$-2.35 \pm 0.37$
$\mu = m_K$	$0.69 \pm 0.12$	$1.26 \pm 0.12$	$-2.35 \pm 0.37$
$\mu = m_\pi$	$2.78 \pm 0.12$	$1.26 \pm 0.12$	$-2.35 \pm 0.37$

Process	$10^3 \alpha_i L^i(\mu)$	$\mu = m_\pi$	$\mu = m_K$	$m_\pi = m_K$	$m_\pi \neq m_K$
$\pi^0 \pi^0$	$2L_1^r(\mu) + 2L_2^r(\mu) + L_3$	$6.20 \pm 0.5$	$1.6 \pm 0.5$	$\geq 2.27$	$\geq 2.28$
$\pi^+ \pi^0$	$L_2^r(\mu)$	$2.81 \pm 0.12$	$1.26 \pm 0.12$	$\geq 0.75$	$\geq 0.95$
$\pi^+ \pi^+$	$2L_1^r(\mu) + 3L_2^r(\mu) + L_3$	$9.0 \pm 0.6$	$2.8 \pm 0.6$	$\geq 3.32$	$\geq 3.91$
$K \eta$	$12L_2^r(\mu) + L_3$	$31.4 \pm 1.5$	$12.8 \pm 1.5$	$\geq 8.6$	-
$\pi \eta$	$3L_2^r(\mu) + L_3$	$6.1 \pm 0.5$	$1.4 \pm 0.5$	$\geq 2.51$	$\geq 6.00$
$K^+ \pi^+$	$4L_2^r(\mu) + L_3$	$8.9 \pm 0.6$	$2.7 \pm 0.6$	$\geq 3.50$	$\geq -5.55$

# Results

	$10^3 \times L_1^r(\mu)$	$10^3 \times L_2^r(\mu)$	$10^3 \times L_3$
$\mu = m_\rho$	$0.43 \pm 0.12$	$0.43 \pm 0.12$	$-2.35 \pm 0.37$
$\mu = m_K$	$0.69 \pm 0.12$	$1.26 \pm 0.12$	$-2.35 \pm 0.37$
$\mu = m_\pi$	$2.78 \pm 0.12$	$1.26 \pm 0.12$	$-2.35 \pm 0.37$

Process	$10^3 \alpha_i L^i(\mu)$	$\mu = m_\pi$	$\mu = m_K$	$m_\pi = m_K$	$m_\pi \neq m_K$
$\pi^0 \pi^0$	$2L_1^r(\mu) + 2L_2^r(\mu) + L_3$	$6.20 \pm 0.5$	$1.6 \pm 0.5$	$\geq 2.27$	$\geq 2.28$
$\pi^+ \pi^0$	$L_2^r(\mu)$	$2.81 \pm 0.12$	$1.26 \pm 0.12$	$\geq 0.75$	$\geq 0.95$
$\pi^+ \pi^+$	$2L_1^r(\mu) + 3L_2^r(\mu) + L_3$	$9.0 \pm 0.6$	$2.8 \pm 0.6$	$\geq 3.32$	$\geq 3.91$
$K \eta$	$12L_2^r(\mu) + L_3$	$31.4 \pm 1.5$	$12.8 \pm 1.5$	$\geq 8.6$	-
$\pi \eta$	$3L_2^r(\mu) + L_3$	$6.1 \pm 0.5$	$1.4 \pm 0.5$	$\geq 2.51$	$\geq 6.00$
$K^+ \pi^+$	$4L_2^r(\mu) + L_3$	$8.9 \pm 0.6$	$2.7 \pm 0.6$	$\geq 3.50$	$\geq -5.55$

The present accuracy of the experimental determinations for  $L_1$ ,  $L_2$  and  $L_3$  is not enough to discern whether SU(3)  $\chi$ PT at  $\mathcal{O}(p^4)$  satisfies the axiomatic principles

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$
  - **Caveat** for resonance saturation

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$
  - **Caveat** for resonance saturation
  - Generalizes scattering lengths method

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$
  - **Caveat** for resonance saturation
  - Generalizes scattering lengths method
- 3 SU(3)

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$
  - **Caveat** for resonance saturation
  - Generalizes scattering lengths method
- 3 SU(3)
  - Bounds on  $L_1$ ,  $L_2$  and  $L_3$

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$
  - **Caveat** for resonance saturation
  - Generalizes scattering lengths method
- 3 SU(3)
  - Bounds on  $L_1$ ,  $L_2$  and  $L_3$
  - **Discards**  $m_u = m_d \equiv m_s$

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$
  - **Caveat** for resonance saturation
  - Generalizes scattering lengths method
- 3 SU(3)
  - Bounds on  $L_1$ ,  $L_2$  and  $L_3$
  - **Discards**  $m_u = m_d \equiv m_s$
  - Consistent with  $m_s = m_d \equiv m_u$

# Conclusions

- 1 EFT & axiomatic principles  $\Rightarrow$  very interesting results
  - Bounds on LECs [as SU(2) and SU(3)  $\chi$ PT]
  - Only efficient for  $\mathcal{O}(p^4)$  LECs
  - Compare order of magnitude of LECs vs chiral logs
  - Permits testing the reliability of EFTs
- 2 SU(2)
  - Bounds on  $l_1$  and  $l_2$
  - **Discard** non-linear sigma Model for  $m_\sigma < 24m_\pi$
  - **Caveat** for resonance saturation
  - Generalizes scattering lengths method
- 3 SU(3)
  - Bounds on  $L_1$ ,  $L_2$  and  $L_3$
  - **Discards**  $m_u = m_d \equiv m_s$
  - Consistent with  $m_s = m_d \equiv m_u$
  - **Cannot discard/confirm** physical situation