

Hadronic Decays of the **TAU** Lepton within **RESONANCE CHIRAL THEORY ($R_\chi T$):**

$$\tau^- \rightarrow (2K \pi)^- \nu_\tau$$

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SUMMARY:

- **Hadronic decays** of the τ lepton
- Tools : χ PT, $R\chi$ T, Large N_c (inspired)
- Previous work: $\tau^- \rightarrow (\pi \pi \pi)^- \nu_\tau$ (Gómez Dumm, Pich, Portolés '04)
- $\tau^- \rightarrow (2K \pi)^- \nu_\tau$ (Gómez Dumm, Pich, Portolés, R. '07 to appear)
- Outlook : $\tau^- \rightarrow (h_1 h_2 h_3)^- \nu_\tau$

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- Outlook : $\tau^- \rightarrow (h_1 h_2 h_3)^- \nu_\tau$ $h_1 h_2 h_3 = 3K, 2\pi K, 2\pi \eta, K\pi\eta$

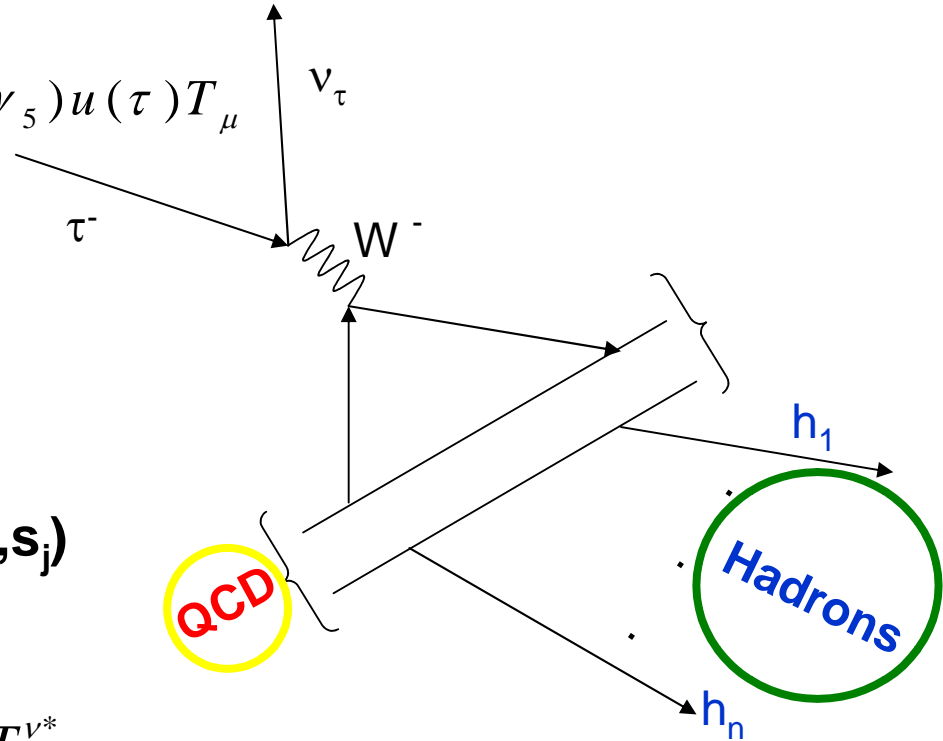
HADRONIC DECAYS OF THE τ LEPTON

$$M = \frac{G_F}{\sqrt{2}} V_{CKM} \bar{u}(v_\tau) \gamma^\mu (1 - \gamma_5) u(\tau) T_\mu$$

$$T_\mu = \langle \text{Hadrons} | (V-A)_\mu e^{iS_{\text{QCD}}} | 0 \rangle =$$

$$= \sum_i (\text{Lorentz Structure})^i F_i(Q^2, s_j)$$

$$d\Gamma = \frac{G_F^2}{4M_\tau^2} |V_{CKM}|^2 d\Phi^{(n+1)} L_{\mu\nu} T^\mu T^{\nu*}$$

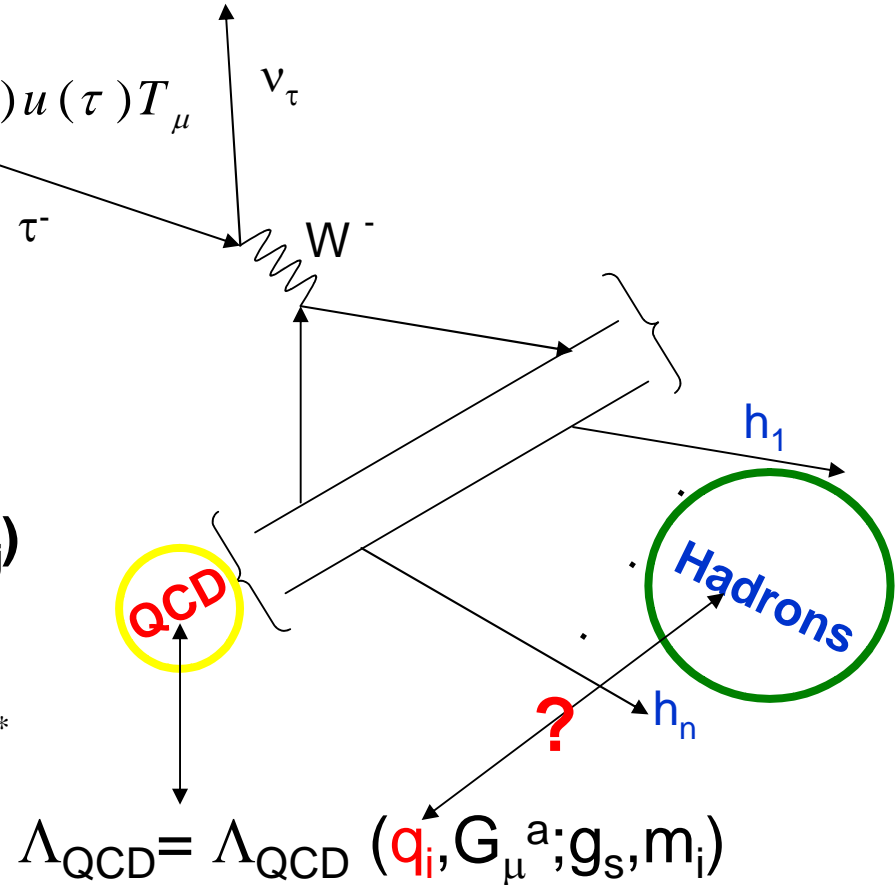


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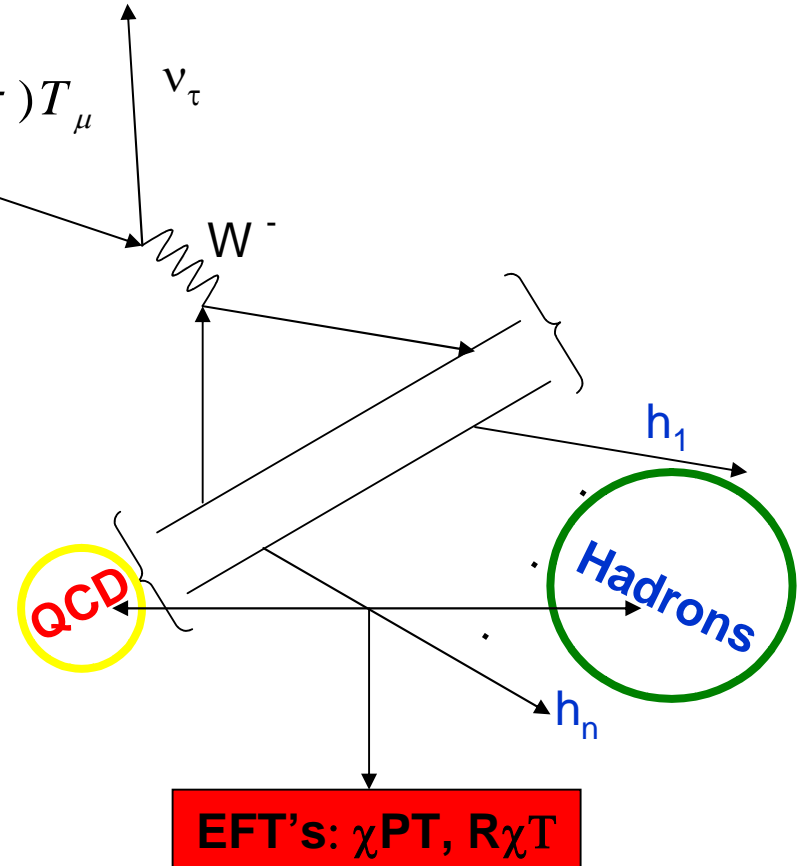


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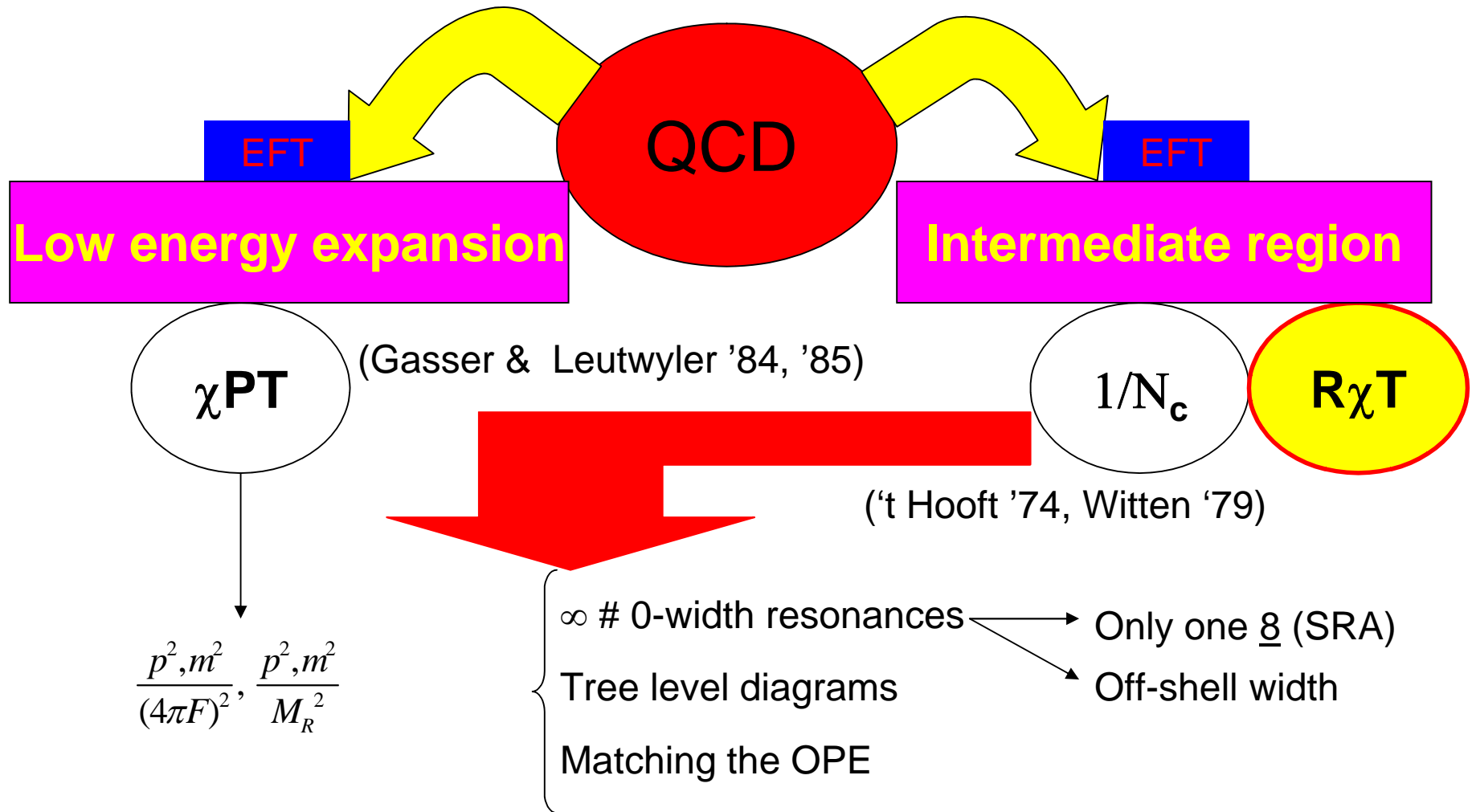
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TOOLS: EFFECTIVE FIELD THEORIES



TOOLS : χ PT

(Gasser & Leutwyler '84, '85)

χ PT
↓

$$\frac{p^2, m^2}{(4\pi F)^2}, \frac{p^2, m^2}{(M_R)^2}$$

(Gasser,
Leutwyler
'84, '85)

$$\varphi(x) = \begin{pmatrix} \frac{\pi^0}{\sqrt{2}} + \frac{\eta_8}{\sqrt{6}} & \pi^+ & K^+ \\ \pi^- & -\frac{\pi^0}{\sqrt{2}} + \frac{\eta_8}{\sqrt{6}} & K^0 \\ K^- & \bar{K}^0 & -\frac{2\eta_8}{\sqrt{6}} \end{pmatrix} SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_V$$

$$u(x) = \exp\left(\frac{i\varphi(x)}{\sqrt{2}F}\right), \quad u_\mu = i\left[u^\dagger(\partial_\mu - ir_\mu)u - u(\partial_\mu - il_\mu)u^\dagger\right]$$

$$\chi = 2B_0(s + ip), \quad \chi_\pm = u^\dagger \chi u^\dagger \pm u \chi u$$

$$f_\pm^{\mu\nu} = u F_L^{\mu\nu} u^\dagger \pm u^\dagger F_R^{\mu\nu} u$$

$$X \rightarrow h(g, \varphi) X h(g, \varphi)^\dagger$$

$$L_\chi^{(2)} = \frac{F^2}{4} \langle u_\mu u^\mu + \chi_+ \rangle$$

$$L_\chi^{(4)} = L_1 \langle u_\mu u^\mu \rangle^2 + \dots + L_4 \langle u_\mu u^\mu \rangle \langle \chi_+ \rangle + \dots + L_7 \langle \chi_- \rangle^2 + \dots - iL_9 \langle f_+^{\mu\nu} u_\mu u_\nu \rangle + \dots$$

$$L_{\chi, WZW}^{(4)} \text{ in the odd-intrinsic parity sector}$$

TOOLS : R χ T

(Ecker, Gasser, Pich, De Rafael '89)

(Ecker, Gasser, Leutwyler, Pich, De Rafael '89)

$$L^{(P_I=+)}_{R\chi T} = L^{(2)}_{\chi} + L^{kin}_{V,A} + L_V + L_A + L_{VAP};$$

$$L_V = \frac{F_V}{2\sqrt{2}} \langle V_{\mu\nu} f_+^{\mu\nu} \rangle + \frac{iG_V}{\sqrt{2}} \langle V_{\mu\nu} u^\mu u^\nu \rangle$$

$$L_A = \frac{F_A}{2\sqrt{2}} \langle A_{\mu\nu} f_-^{\mu\nu} \rangle$$

Antisymmetric tensor formalism

$$V_{\mu\nu}(x) = \begin{pmatrix} \frac{\rho^0}{\sqrt{2}} + \frac{\omega_8}{\sqrt{6}} & \rho^+ & K^{*+} \\ \rho^- & -\frac{\rho^0}{\sqrt{2}} + \frac{\omega_8}{\sqrt{6}} & K^{*0} \\ K^{*-} & \bar{K}^{*0} & -\frac{2\omega_8}{\sqrt{6}} \end{pmatrix}_{\mu\nu}$$

$$L_{VAP} = \sum_{i=1}^5 \lambda_i O^i(V_{\mu\nu}, A^{\mu\nu}, \varphi) = \lambda_1 \langle [V_{\mu\nu}, A^{\mu\nu}] \chi_- \rangle + \dots$$

(Gómez Dumm, Pich, Portolés '04)

$$L^{(P_I=-)}_{R\chi T} = L^{(4)}_{\chi(WZW)} + L_{VJP} + L_{VVP} + L_{VPPP};$$

VMD

$$L_{VJP} = \sum_{i=1}^7 \frac{c_i}{M_V} O^i(V_{\mu\nu}, j^\nu, \partial^\mu \varphi) = \frac{c_5}{M_V} \varepsilon_{\mu\nu\rho\sigma} \langle \{ \nabla_\alpha V^{\mu\nu}, f_+^{\rho\alpha} \} u^\sigma \rangle + \dots$$

(Ruiz-Femenía, Pich, Portolés '03)

$$L_{VVP} = \sum_{i=1}^5 d_i O^i(V_{\mu\nu}, V_{\rho\sigma}, \varphi) = d_1 \varepsilon_{\mu\nu\rho\sigma} \langle \{ V^{\mu\nu}, V^{\rho\alpha} \} \nabla_\alpha u^\sigma \rangle + \dots$$

$$L_{VPPP} = \sum_{i=1}^5 \frac{g_i}{M_V} O^i(V_{\mu\nu}, \varphi) = \frac{g_4}{M_V} \varepsilon_{\mu\nu\alpha\beta} \langle \{ V^{\mu\nu}, u^\alpha u^\beta \} \chi_- \rangle + \dots$$

(Gómez Dumm, Pich, Portolés, R. '07 to appear)

TOOLS : R χ T

$$L_{R\chi T}^{(P_I=+)} = L_{\chi}^{(2)} + L_{V,A}^{kin} + L_V + L_A + L_{VAP};$$

$$L_V = \frac{F_V}{2\sqrt{2}} \langle V_{\mu\nu} f_+^{\mu\nu} \rangle + \frac{iG_V}{\sqrt{2}} \langle V_{\mu\nu} u^\mu u^\nu \rangle$$

$$L_A = \frac{F_A}{2\sqrt{2}} \langle A_{\mu\nu} f_-^{\mu\nu} \rangle$$

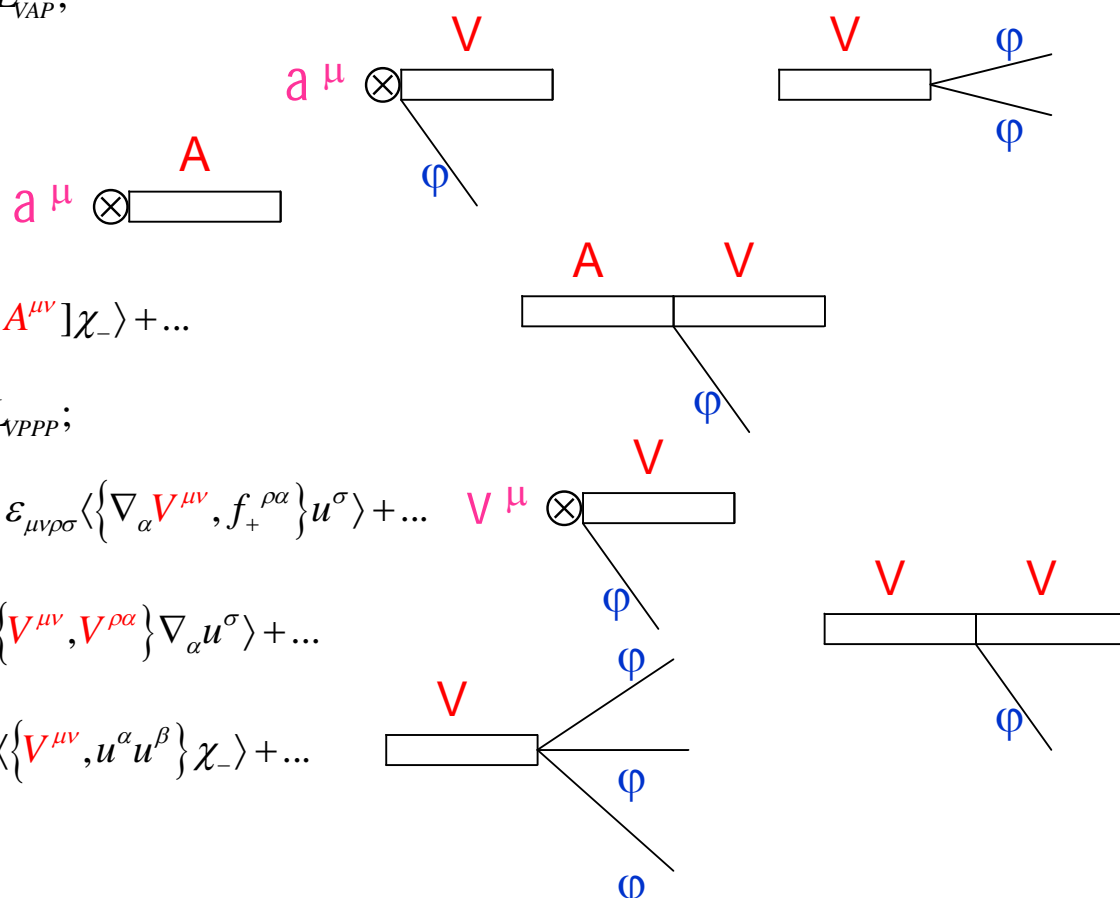
$$L_{VAP} = \sum_{i=1}^5 \lambda_i O^i(V_{\mu\nu}, A^{\mu\nu}, \varphi) = \lambda_1 \langle [V_{\mu\nu}, A^{\mu\nu}] \chi_- \rangle + \dots$$

$$L_{R\chi T}^{(P_I=-)} = L_{\chi(WZW)}^{(4)} + L_{VJP} + L_{VVP} + L_{VPPP};$$

$$L_{VJP} = \sum_{i=1}^7 \frac{c_i}{M_V} O^i(V_{\mu\nu}, j^\nu, \partial^\mu \varphi) = \frac{c_5}{M_V} \varepsilon_{\mu\nu\rho\sigma} \langle \{ \nabla_\alpha V^{\mu\nu}, f_+^{\rho\alpha} \} u^\sigma \rangle + \dots$$

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HADRONIC DECAYS OF THE τ LEPTON

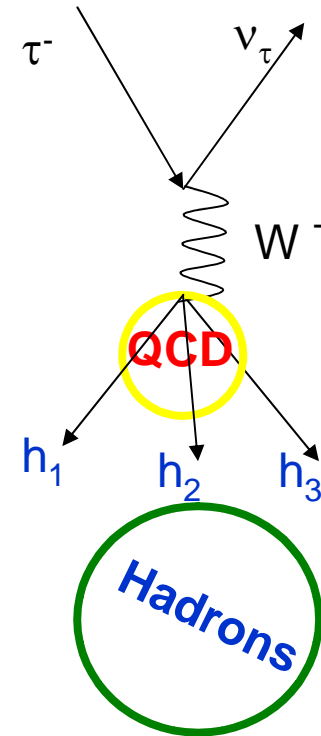
$$M = \frac{G_F}{\sqrt{2}} V_{CKM} \bar{u}(v_\tau) \gamma^\mu (1 - \gamma_5) u(\tau) T_\mu$$

$$\tau^- \rightarrow h_1(p_1) h_2(p_2) h_3(p_3) v_\tau$$

$$(p_1 + p_2 + p_3)^\mu = Q^\mu, \quad V_{1\mu} = \left(g_{\mu\nu} - \frac{Q_\mu Q_\nu}{Q^2} \right) (p_2 - p_1)^\nu, \quad V_{2\mu} = \left(g_{\mu\nu} - \frac{Q_\mu Q_\nu}{Q^2} \right) (p_3 - p_1)^\nu$$

$$T_\mu = V_{1\mu} F_1 + V_{2\mu} F_2 + Q_\mu F_P + i \varepsilon_{\mu\nu\rho\sigma} p_1^\nu p_2^\rho p_3^\sigma F_V$$

$$\frac{d\Gamma}{dQ^2} = \frac{G_F^2 |V_{CKM}|^2}{128(2\pi)^5 M_\tau^3} \int ds dt f(I_{0^-}, I_{1^+}, I_{1^-})$$



(Kühn, Santamaría '90)(KS)

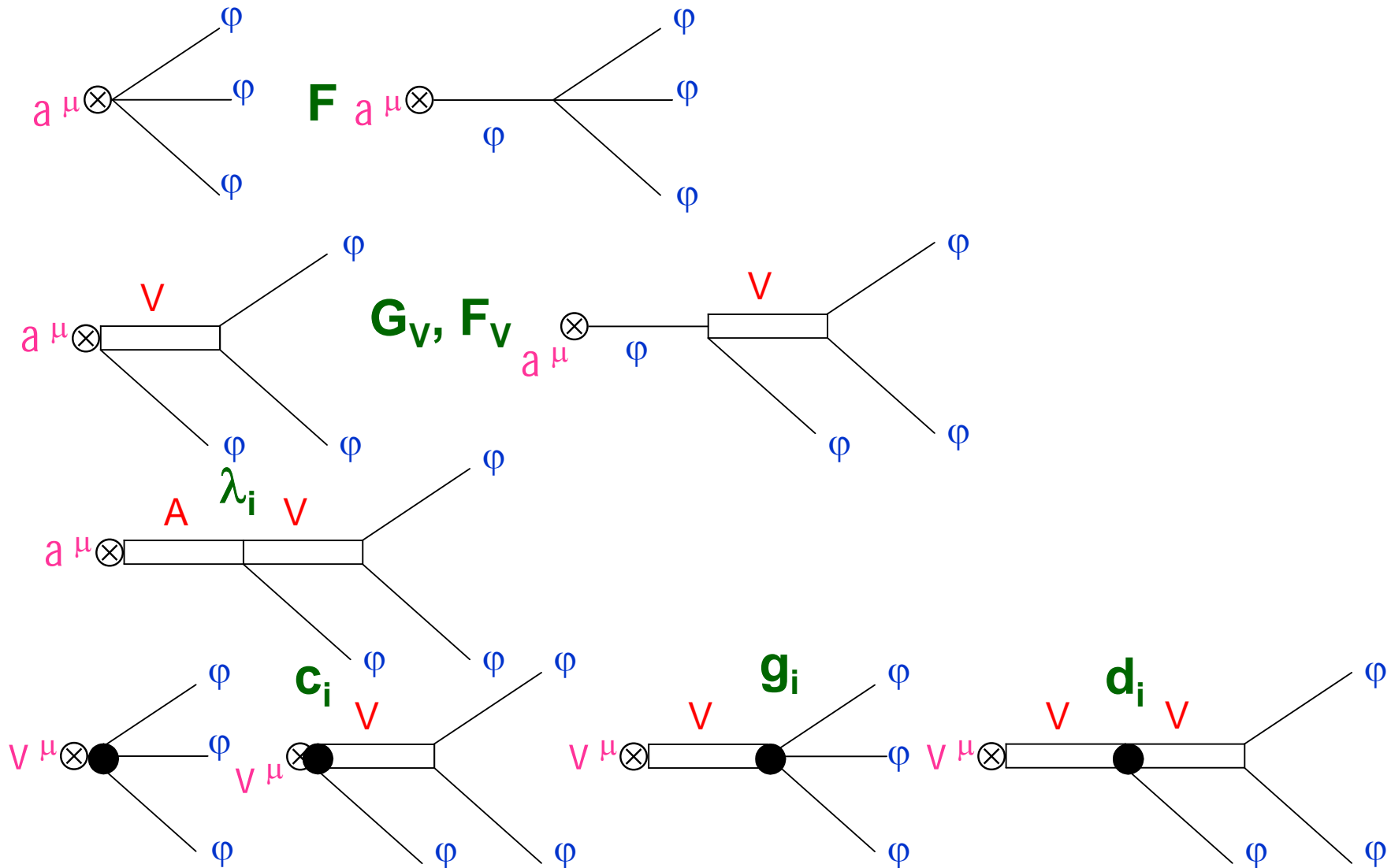
(Gómez-Cadenas, González-García, Pich '90) (GGP)

(Finkemeier, Mirkes '96)(FM)

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$\tau^- \rightarrow (2K \pi)^- \nu_\tau$ in R χ T
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HADRONIC DECAYS OF THE τ LEPTON



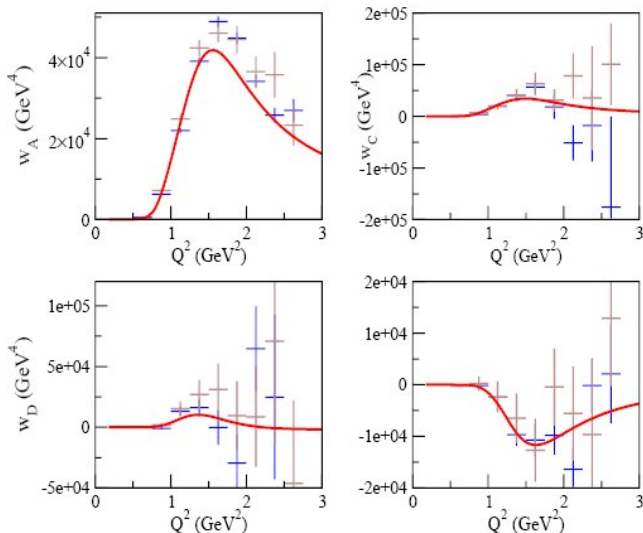
PREVIOUS WORK : $\tau^- \rightarrow (\pi \pi \pi)^- \nu_\tau$

Procedure and results [Gómez Dumm, Pich, Portolés, 2004]

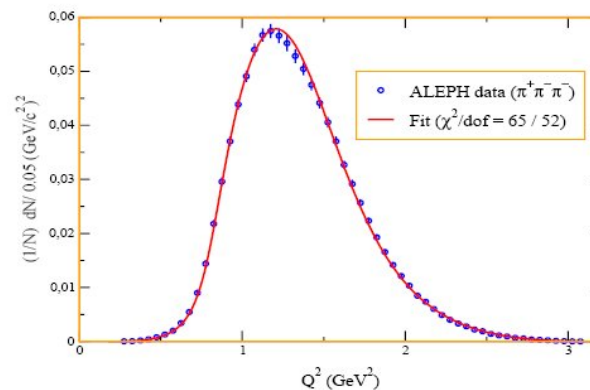
Fit to the spectrum and BR

[ALEPH, 1998] $\tau^- \rightarrow \pi^- \pi^- \pi^+ \nu_\tau$

input



[CLEO-II,2000,(blue)] [OPAL,1997,(grey)]



output

$\tau^- \rightarrow \pi^- \pi^0 \pi^0 \nu_\tau$

Structure Functions

PDG06

$$M_{a_1} = (1.204 \pm 0.007) \text{ GeV}$$

$$\Gamma_{a_1}(M_{a_1}) = (0.48 \pm 0.02) \text{ GeV}$$

Implemented in SHERPA

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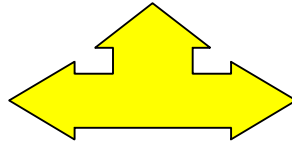
$\tau \rightarrow (2K \pi)^- \nu_\tau$ in R χ T
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$$\tau^- \rightarrow (\pi \pi \pi)^- \nu_{\tau}$$

$\lambda_0??$

(Gómez-Dumm,
Pich, Portolés '04)

12

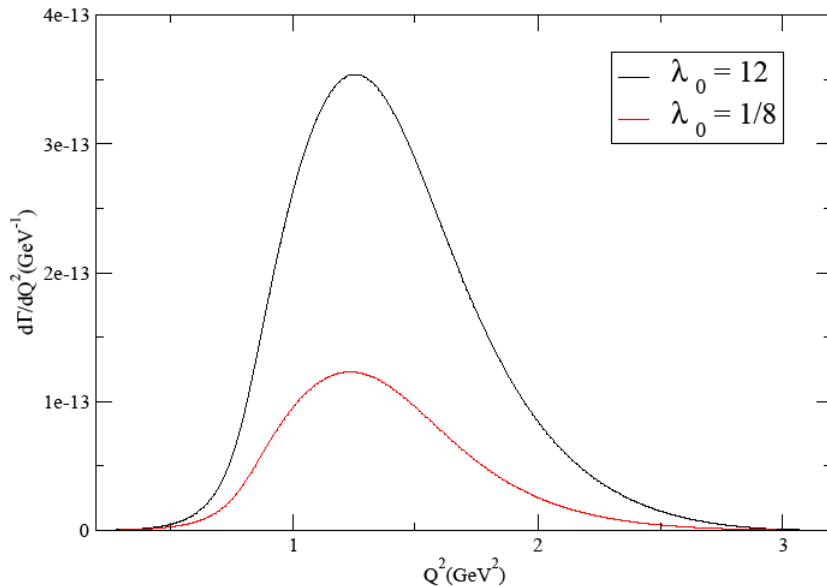


1/8

(Cirigliano,
Ecker, Eidemüller,
Pich, Portolés '04)

<VAP>

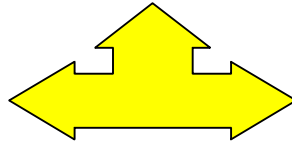
$$\tau^- \rightarrow (\pi \pi \pi)^- \nu_{\tau}$$



$$\underline{\tau^- \rightarrow (\pi \pi \pi)^- \nu_{\tau}}$$

$$\lambda_0 \frac{m_{\pi}^2}{Q^2}$$

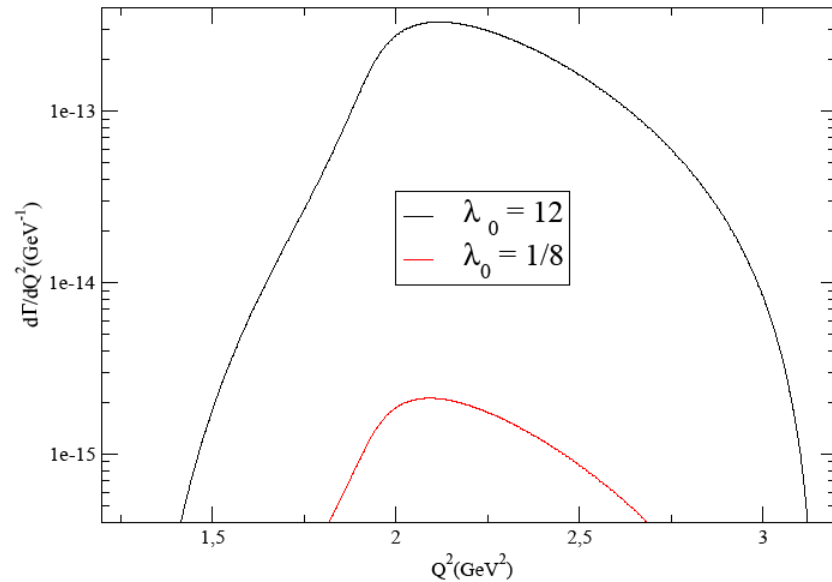
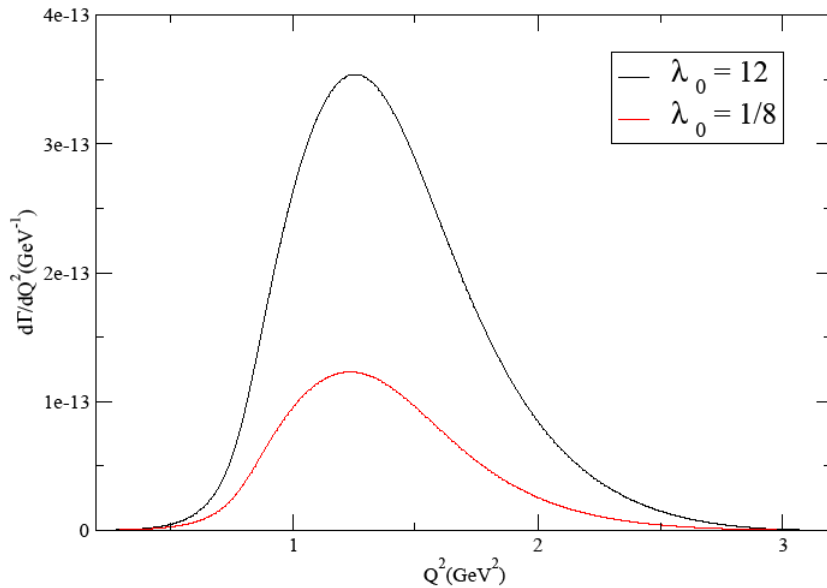
$\lambda_0??$



$$\lambda_0 \frac{m_K^2}{Q^2}$$

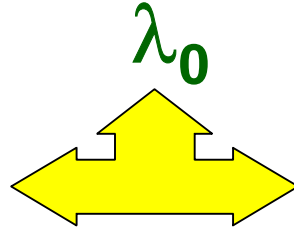
$$\underline{\tau^- \rightarrow (\pi \pi \pi)^- \nu_{\tau}}$$

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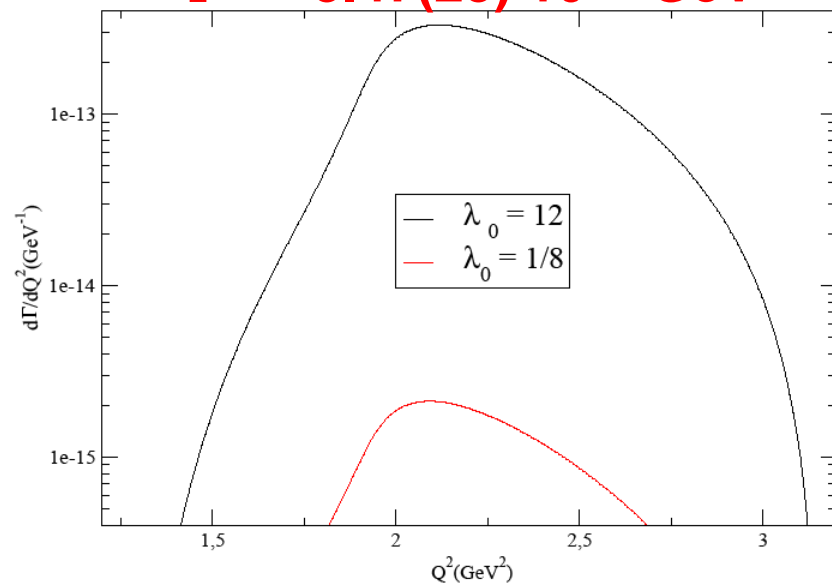
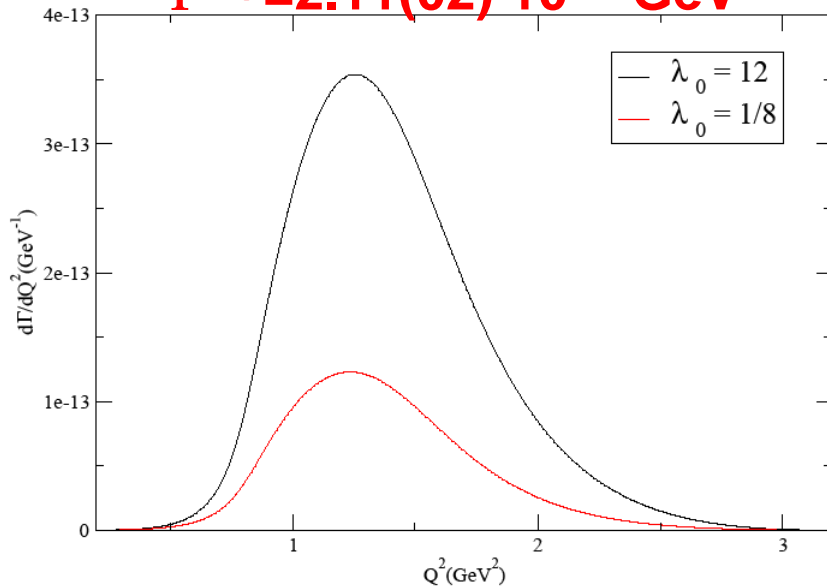
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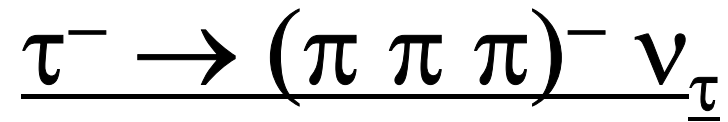
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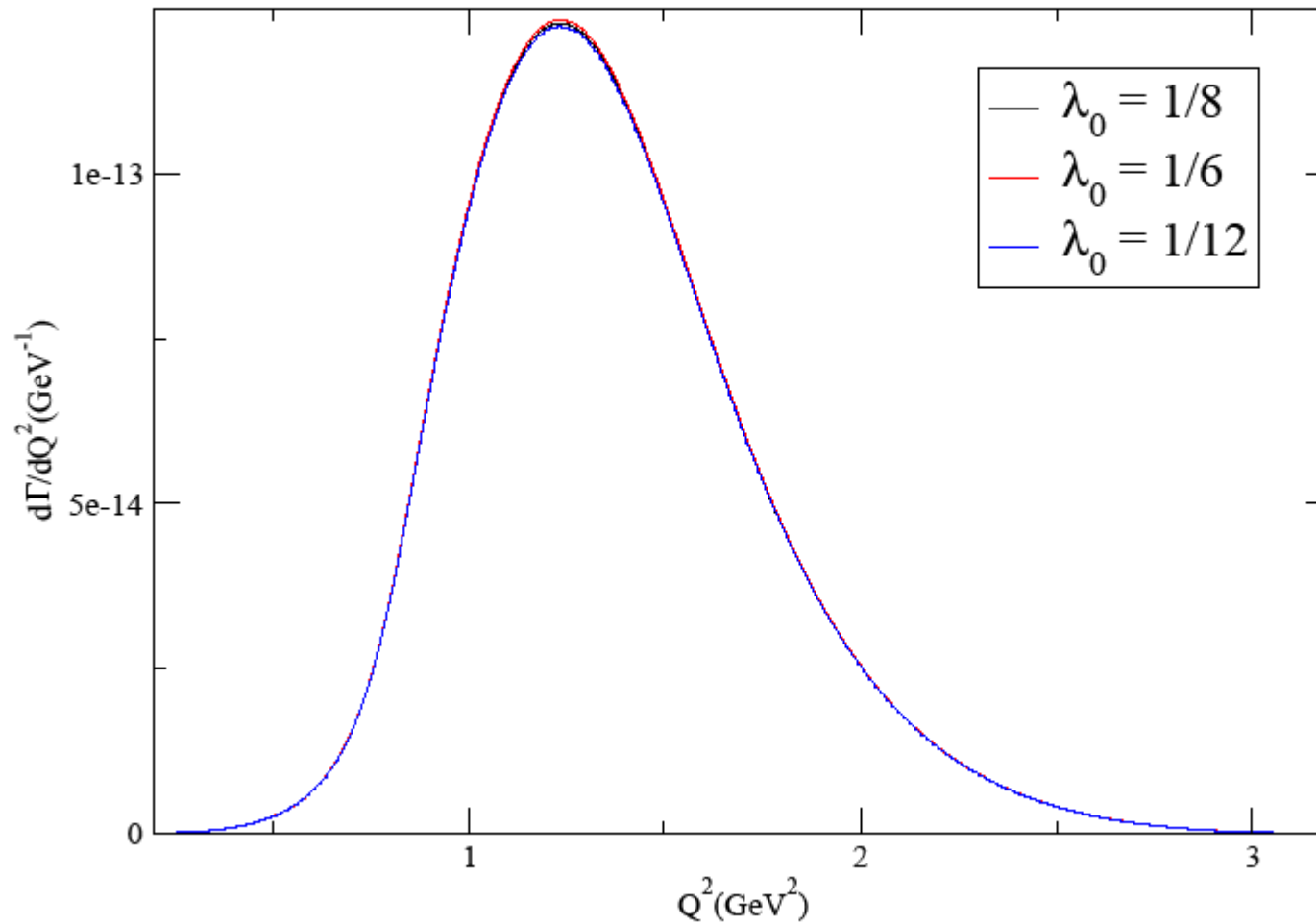
$$\Gamma^{\text{exp}} = 2.11(02) \cdot 10^{-13} \text{ GeV}$$

$$\Gamma^{\text{exp}} = 3.47(23) \cdot 10^{-15} \text{ GeV}$$



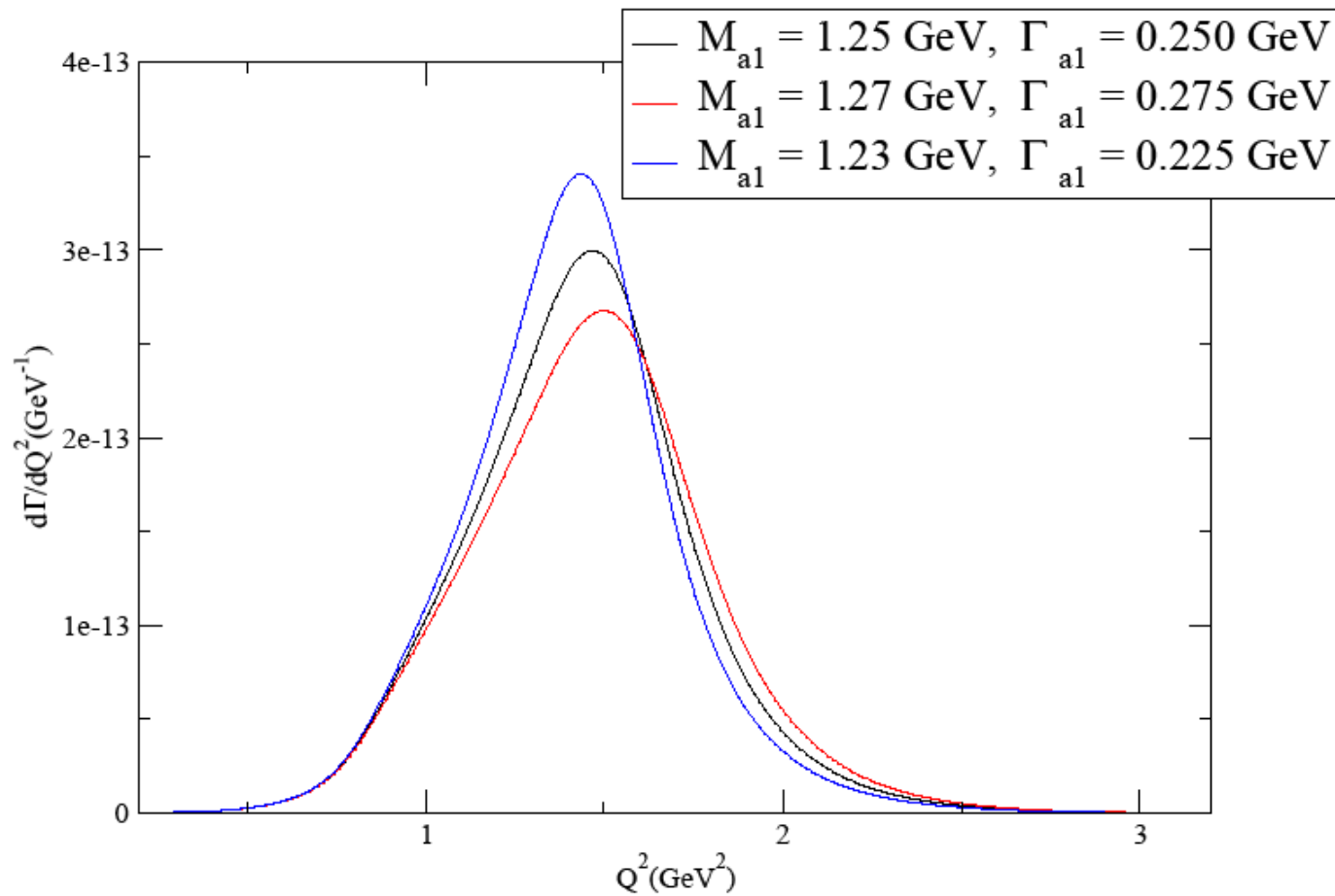


$$\lambda_0 \sim 1/8$$



$$\tau^- \rightarrow (\pi \pi \pi)^- \nu_{\tau}$$

$\lambda_0 = 1/8$ Varying M_{a_1} and Γ_{a_1}



$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

$\tau^- \rightarrow (KK\pi)^- \nu_\tau$ within Resonance Chiral Theory

[Gómez Dumm, Pich, Portolés, Roig, **coming soon**]

- Analysis of both **axial-vector** and **vector** form factors
- Required Brodsky-Lepage behaviour : constraints on unknown couplings of the Lagrangian
- Implemented in **SHERPA**

No experimental contrast

~~[CLEO-III, 2004]~~

Normalization of form factors violating chiral symmetry at leading order

$$c_{1235} = c_1 + c_2 + 8c_3 - c_5$$

$$d_{123} = g_1 + 8d_2 - d_3$$

$$c_{125} = c_1 - c_2 + c_5$$

$$c_{1256} = c_1 - c_2 - c_5 + 2c_6$$

$$g_{123} = g_1 + 2g_2 - g_3$$

Free

$c_{1235}, d_{123}, c_4, g_4, g_5$

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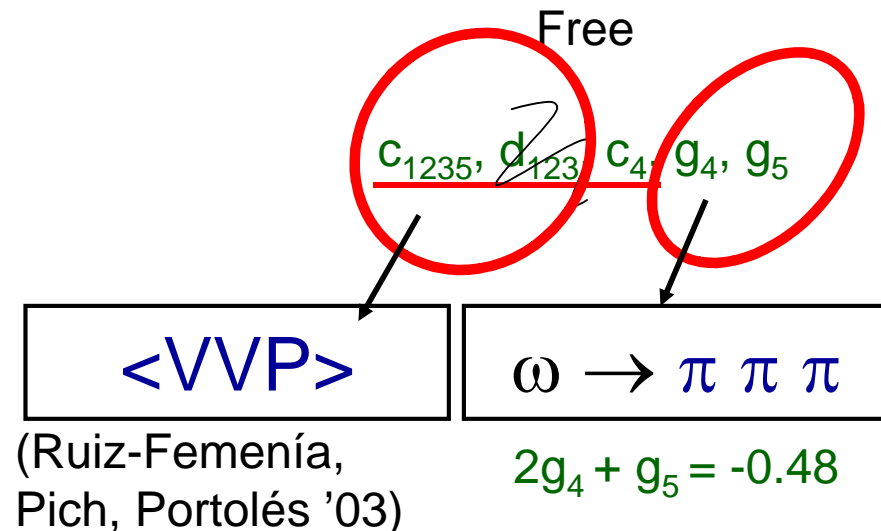
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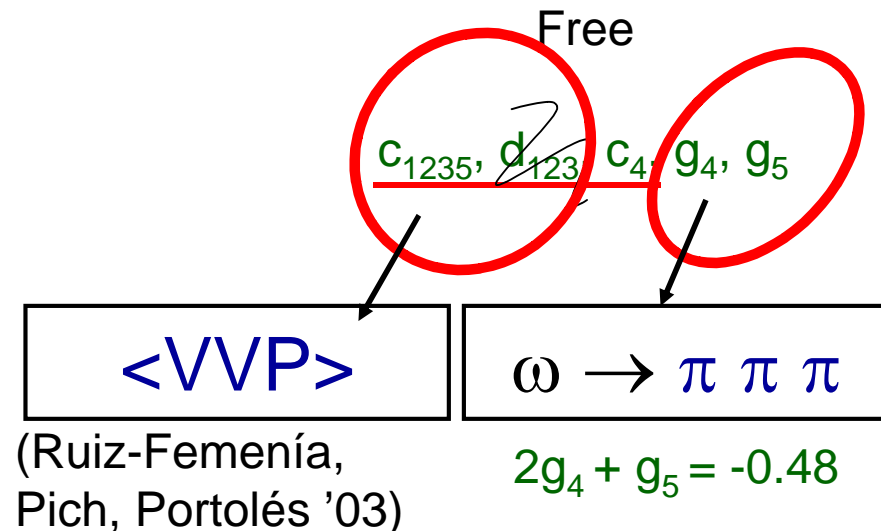
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Normalization of form factors violating chiral symmetry at leading order

Contributions to the decay width

	Axial-Vector	Vector
Our result	~ 37 %	~ 63 %
CLEO-III	~ 50 %	~ 50 %
[GGP,90]	~ 10 %	~ 90 %
[FM,96]	~ 60 %	~ 40 %

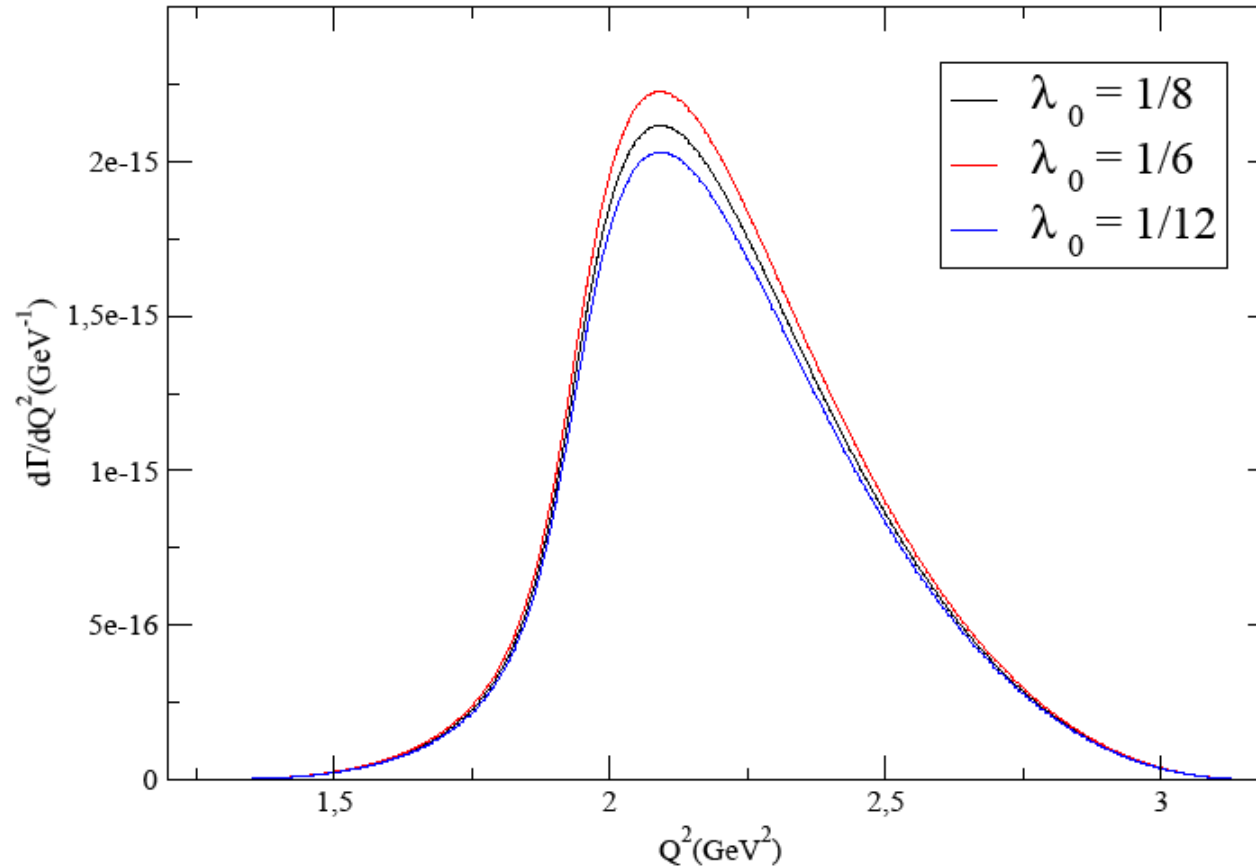


Different predictions for the $K \bar{K} \pi^-$ channels

		χ PT at $O(p^2)$		χ PT at $O(p^4)$		
	N_V/N_A	F^A	F^A	F_4^V	Γ_V	Γ_A
CLEO '04	~ 1.5	Yes				
This work	~ 3					
(Gómez-Cadenas, González-García, Pich '90)	~ 9				-	-
(Finkemeier, Mirkes '96)	~ 0.4					<i>ad-hoc</i>

$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

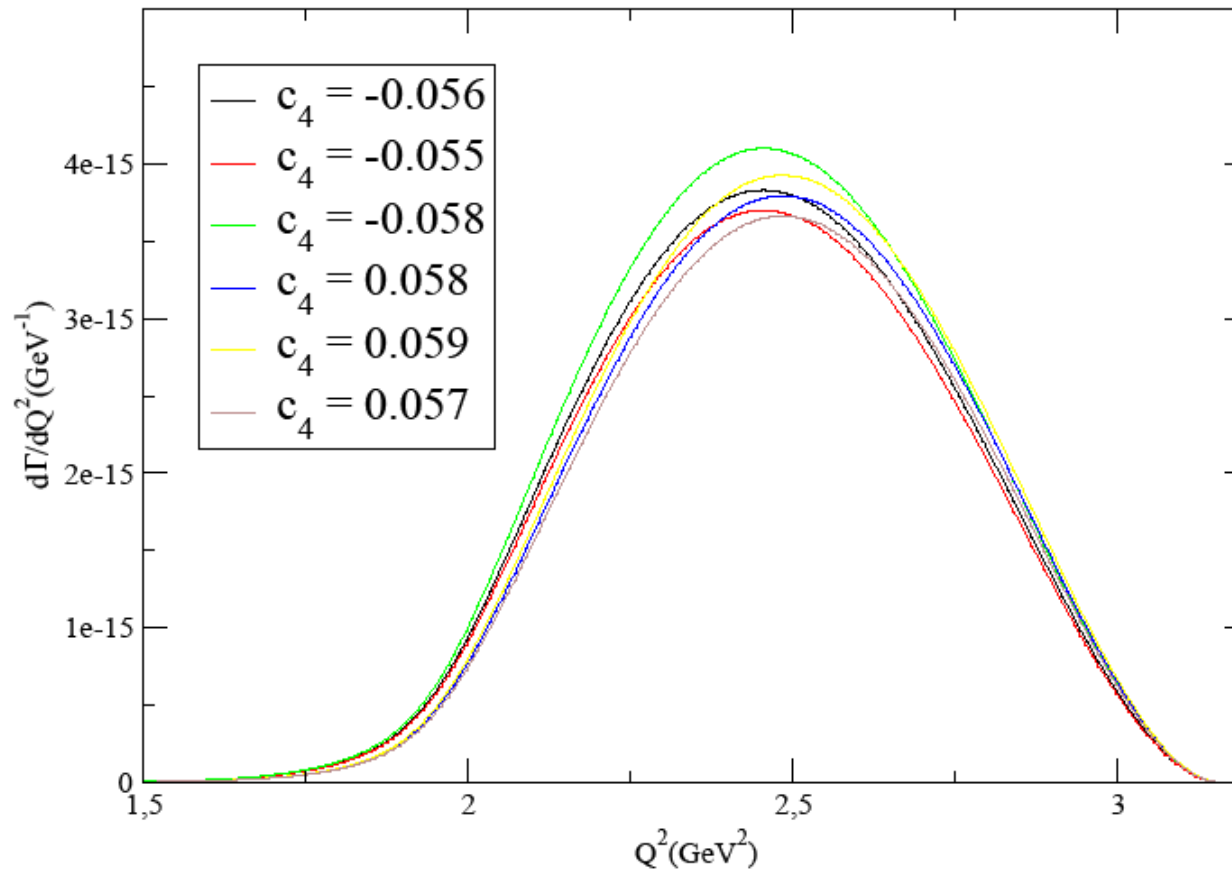
Axial vector current contribution: Dependence on λ_0
 $K^+K^-\pi^-$



$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

Vector current contribution: Dependence on c_4

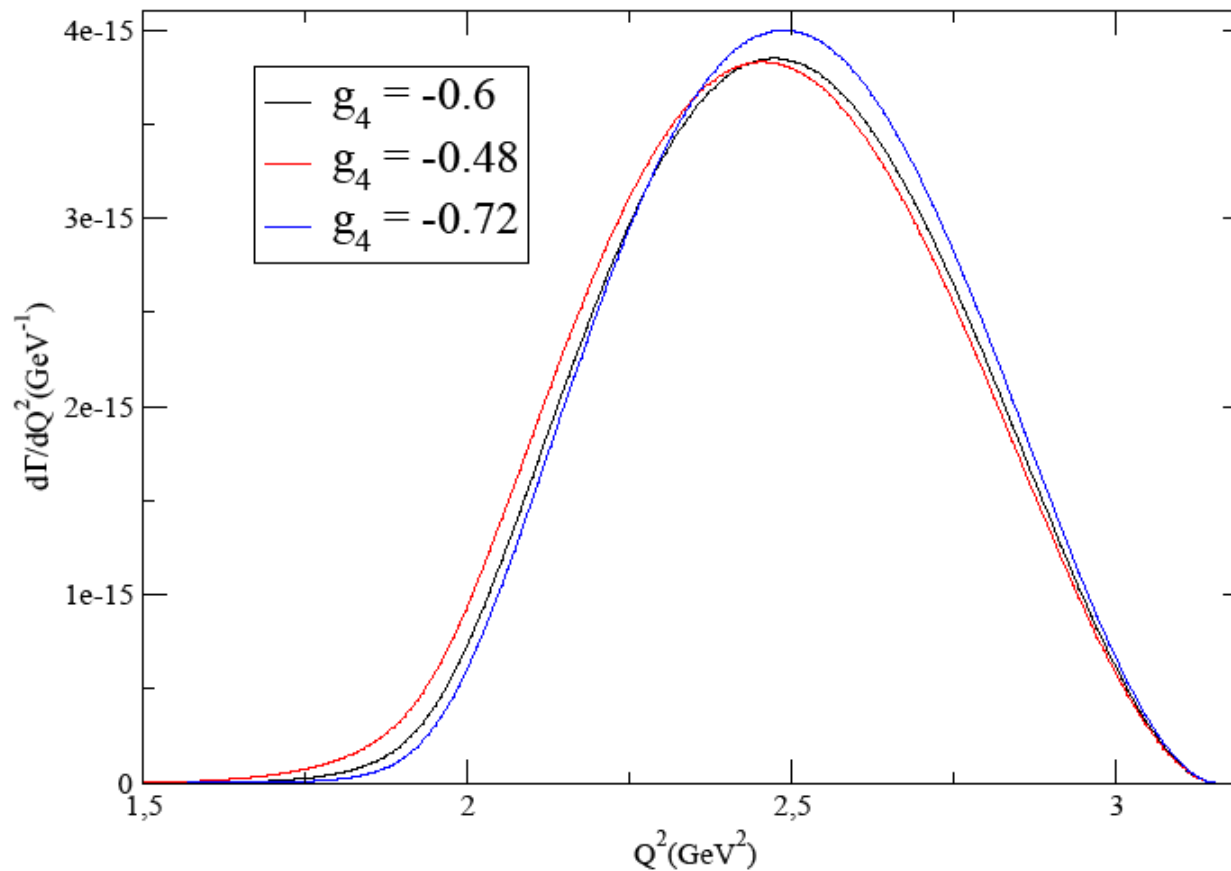
$K^+K^-\pi^-$



$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

Vector current contribution: Dependence on g_4

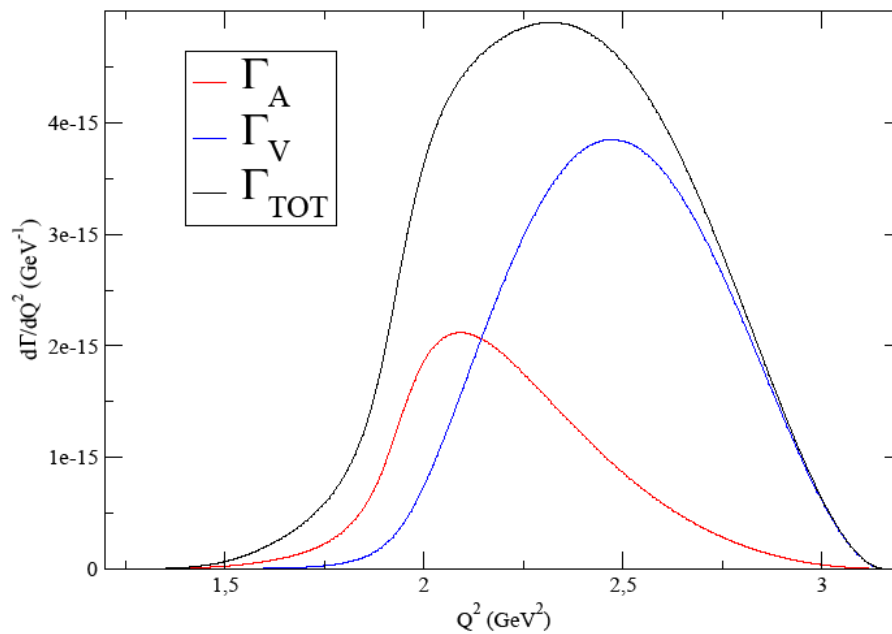
$K^+K^-\pi^-$



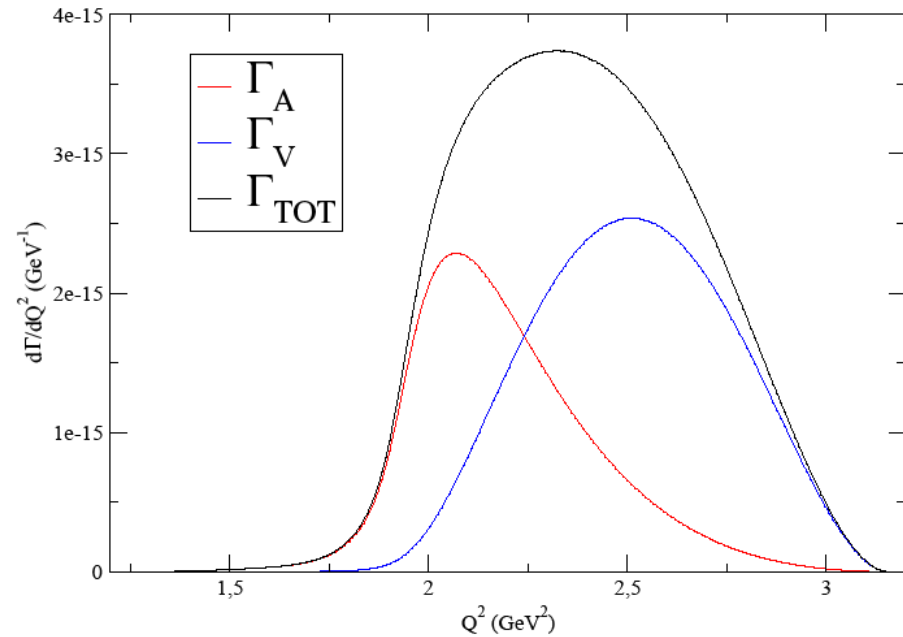
$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

“BEST” FITS

$K^+K^-\pi^-$



$K^-K^0\pi^0$



$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

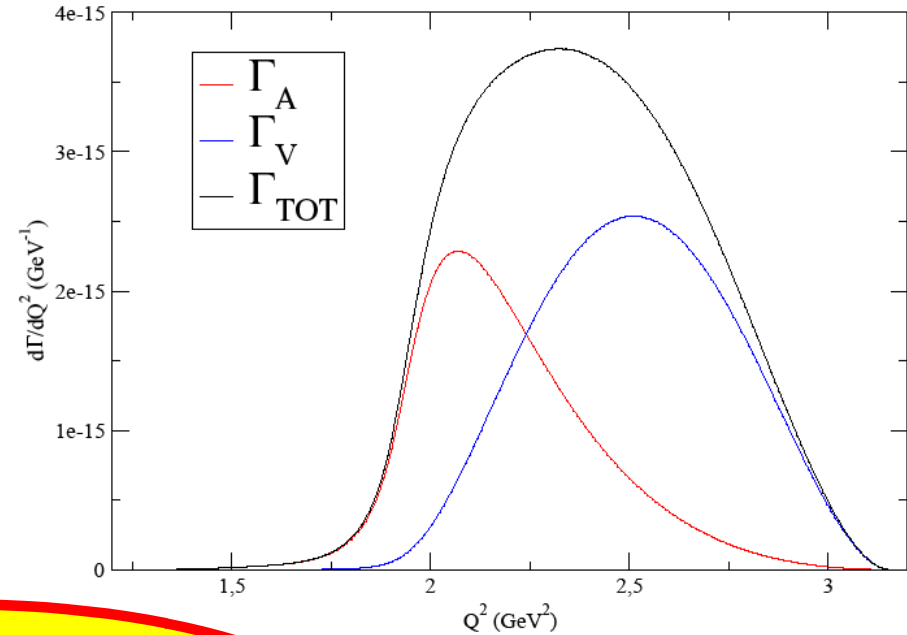
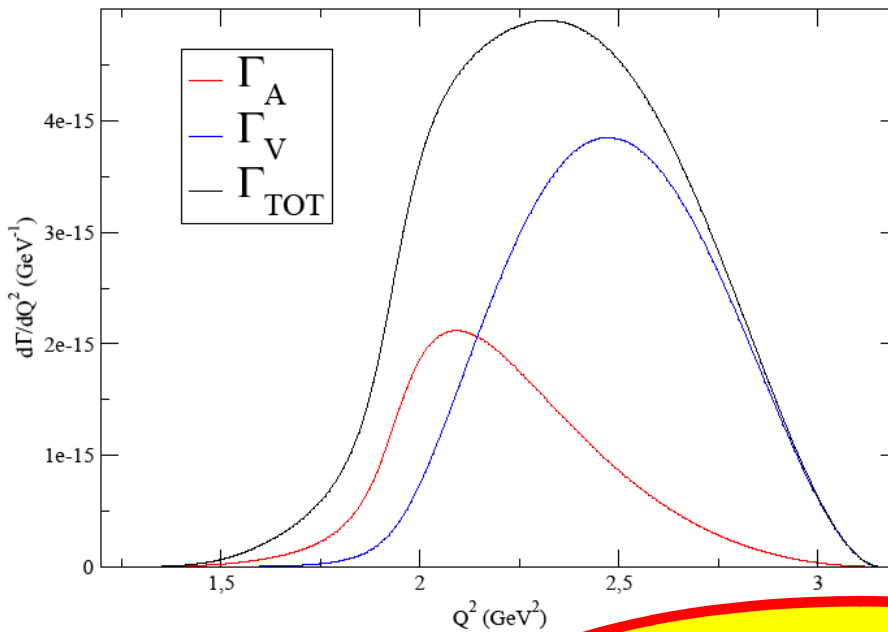


(I. Nugent TAU06)

“BEST” FITS

$K^+K^-\pi^-$

$K^-K^0\pi^0$



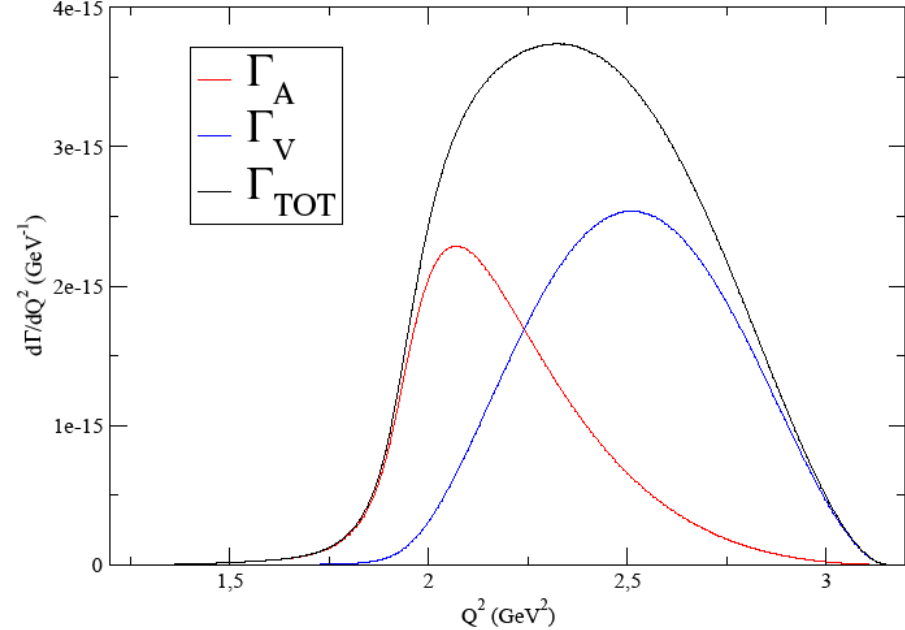
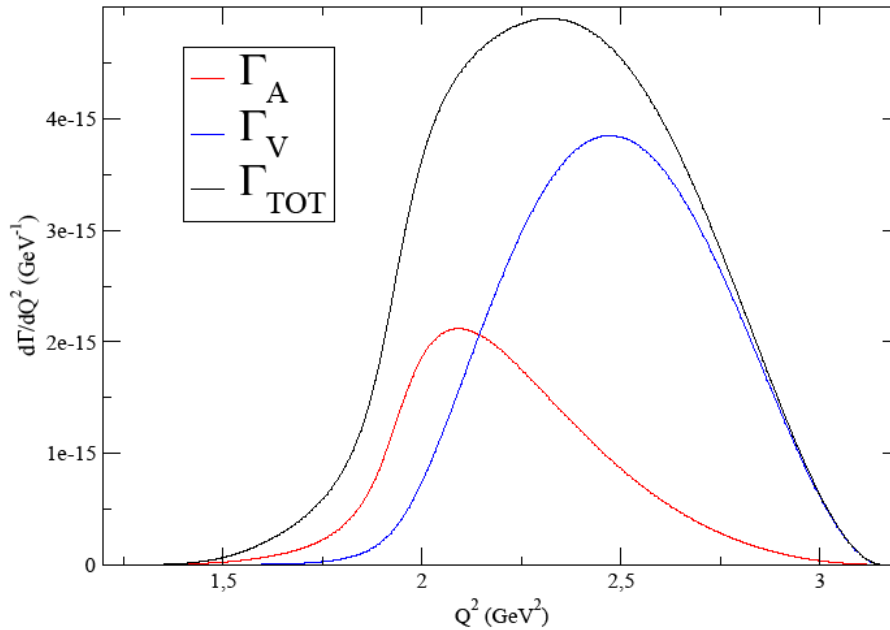
MC's  data
 Our Hadronic matrix
 elements



$K^+K^-\pi^-$

$K^-K^0\pi^0$

“BEST” FITS



$$|c_4| = 0.057 \pm 0.005$$

$$g_4 = -0.48 \pm 0.24$$

$$M_{a1} = 1.25 \pm 0.03 \text{ GeV}$$

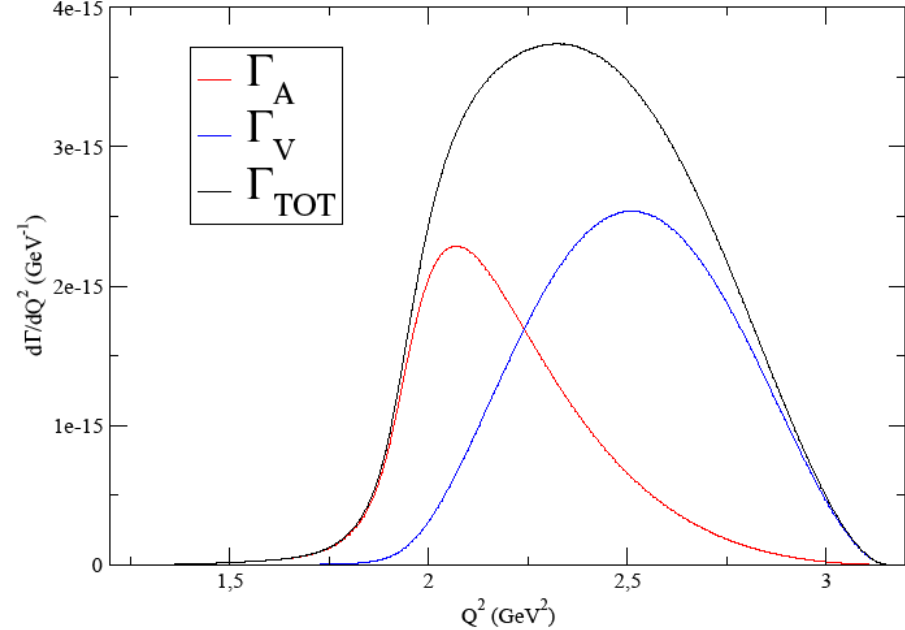
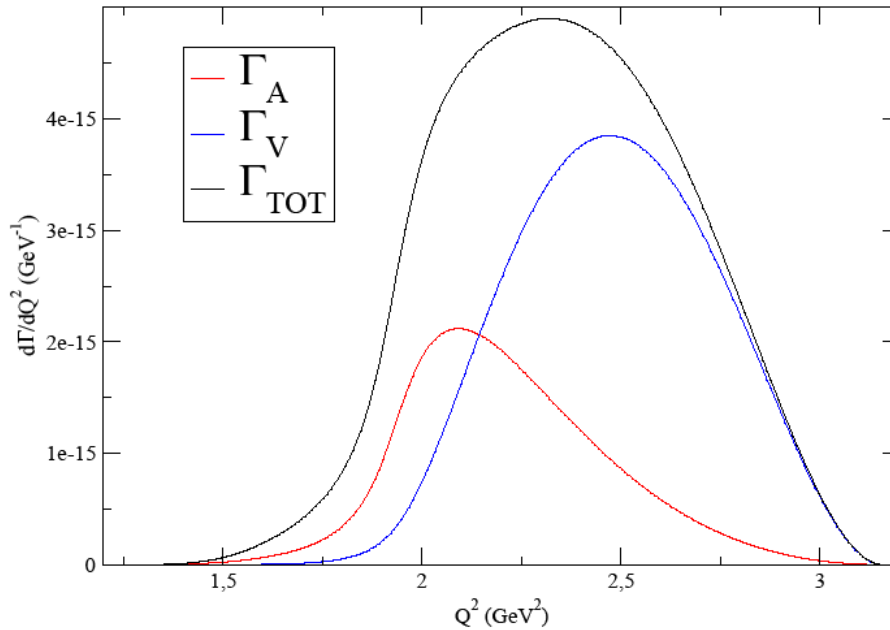
$$\Gamma_{a1}(M_{a1}) = 0.25 \pm 0.03 \text{ GeV}$$

$\tau^- \rightarrow (2K \pi)^- \nu_\tau$

$K^+K^-\pi^-$

$K^-K^0\pi^0$

“BEST” FITS



$$\frac{\Gamma(\tau^- \rightarrow K^+K^-\pi^-\nu_\tau)}{\Gamma(\tau^- \rightarrow K^-K^0\pi^0\nu_\tau)} = 1.7 \pm_{0.3}^{0.1}$$

$$\frac{\Gamma_V(\tau^- \rightarrow (2K\pi)^-\nu_\tau)}{\Gamma_A(\tau^- \rightarrow (2K\pi)^-\nu_\tau)} = 1.7 \pm_{0.6}^{1.0}$$

OUTLOOK: $\tau^- \rightarrow (\underline{h}_1 \underline{h}_2 \underline{h}_3)^- \nu_\tau$

$\underline{h}_1 \underline{h}_2 \underline{h}_3$	F_A	F_V	V_{us}, m_s	m_{ν_τ}	
3π	✓				
$2K \pi$	✓	✓			
$K 2\pi$	✓	✓	✓		
$2\pi \eta$		✓			
$\pi K \eta$	✓	✓	✓		
$3K$	✓	✓	✓	✓	$(\Phi \rightarrow \bar{K}K)$

CONCLUSIONS

- Current analyses of $\tau^- \rightarrow (3\pi^-, K^+K^-\pi^-) \nu_\tau$ data using **TAUOLA** have shown noticeable inconsistencies. We would like to improve the hadronic matrix elements in **TAUOLA**.
- We have studied these decays within **R χ T** with a **Large N_c**-inspired model guided by **QCD**.
- We have revisited $\tau^- \rightarrow 3\pi^- \nu_\tau$ obtaining a better parameterization
- Within the **2K π** channels we have found more br to the channels containing π^-
- We have found Vector Current dominance on the **2K π** channels
- Our results have been implemented in **SHERPA (LHC and TEVATRON)** and will be used by **BELLE**.
- **LHC, BABAR, BELLE, BES-III...** are promising facilities to test our predictions.
- The future looks even more promising and exciting: V_{us} , m_s , m_{V_τ}
and, of course, hadronization of QCD currents.

SKIPPED

SLIDES

$K \bar{K} \pi^-$ Channels within $R_{\chi T}$

Experiment

$$\Gamma = 3.47(23) \cdot 10^{-15} \text{ GeV}$$

vector_1 : (This work)

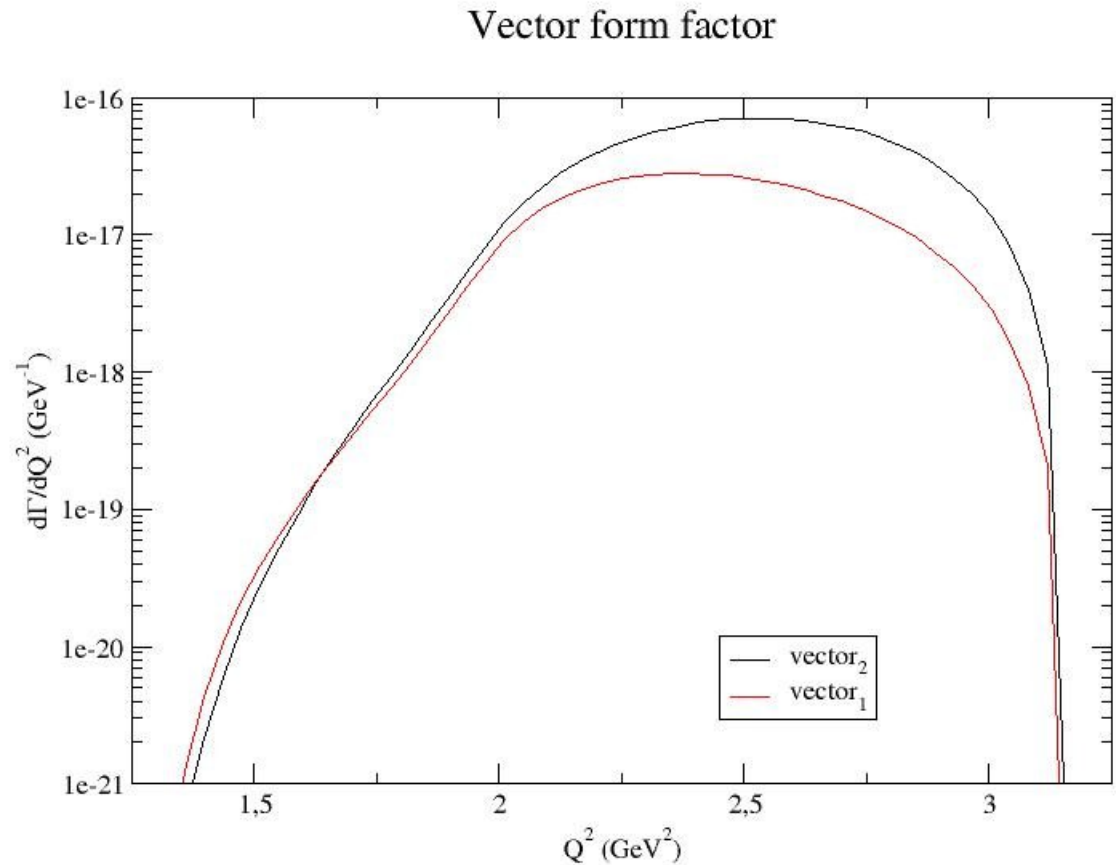
$$\Gamma \sim 1.97 \cdot 10^{-17} \text{ GeV}$$

vector_2 :

(Ruiz-Femenia, Pich, Portolés '03)

$$\Gamma \sim 4.96 \cdot 10^{-17} \text{ GeV}$$

EuroFlavour '07
FlaviAnet



$\tau \rightarrow (2K \pi)^- \nu_\tau$ in $R_{\chi T}$
Pablo Roig Garcés
Instituto de Física Corpuscular

$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

$\tau^- \rightarrow (KK\pi)^- \nu_\tau$ within Resonance Chiral Theory

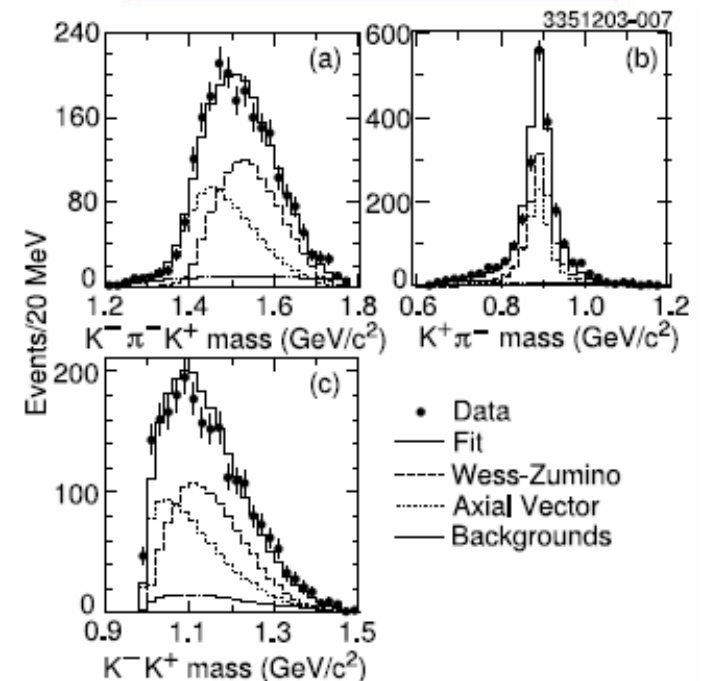
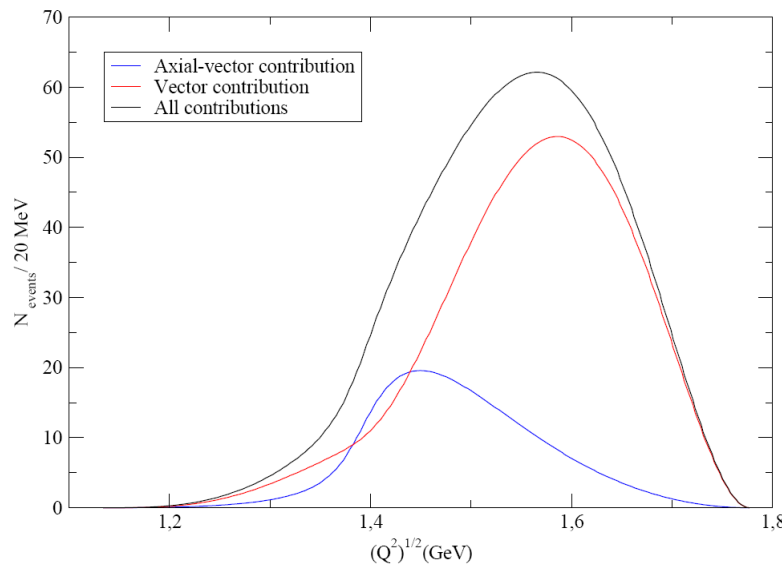
[Gómez Dumm, Pich, Portolés, Roig, **coming soon**]

- Analysis of both **axial-vector** and **vector** form factors
- Required Brodsky-Lepage behaviour : constraints on unknown couplings of the Lagrangian
- Implemented in **SHERPA**

No experimental contrast

~~[CLEO-III, 2004]~~

Normalization of form factors violating chiral symmetry at leading order



$$\underline{\tau^- \rightarrow (2K \pi)^- \nu_\tau}$$

$\tau^- \rightarrow (KK\pi)^- \nu_\tau$ within Resonance Chiral Theory

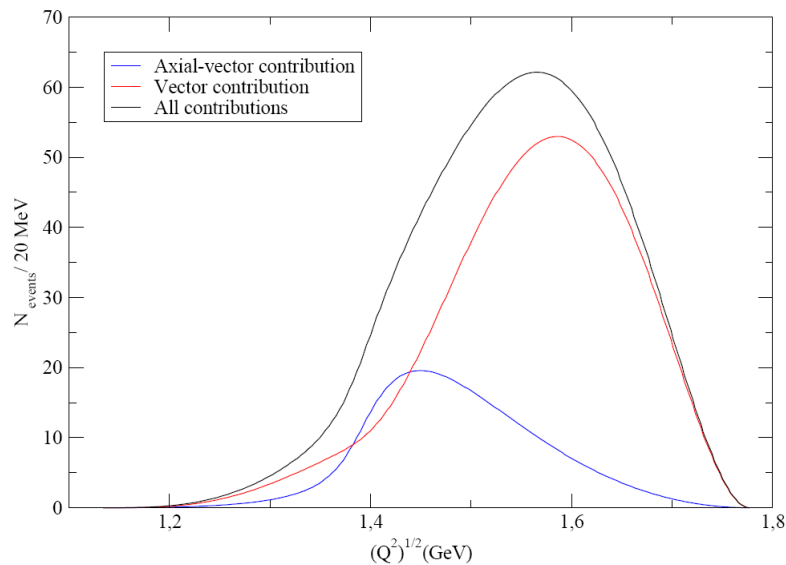
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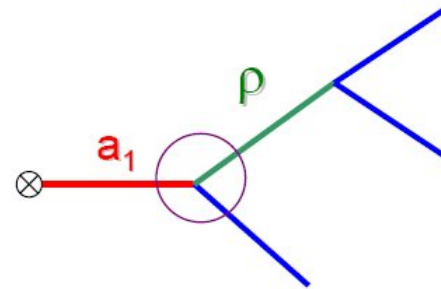
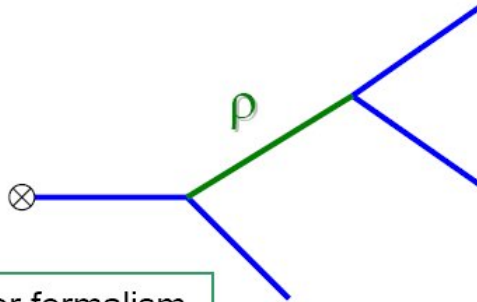
~~[CLEO-III, 2004]~~

Normalization of form factors violating chiral symmetry at leading order



Experiment	# of τ pairs
LEP	$\sim 3 \cdot 10^5$
CLEO	$\sim 1 \cdot 10^7$
BABAR	$\sim 3 \cdot 10^8$
BELLE	$\sim 5 \cdot 10^8$
LHC	$\sim 1 \cdot 10^{12}$

Chiral Resonance Theory + Large- N_C



Tensor formalism
for spin=1 mesons

$$\mathcal{L}_{R\chi T} = \mathcal{L}_{\chi}^{(2)} + \mathcal{L}_V + \mathcal{L}_A + \mathcal{L}_{VAP}$$

$$\mathcal{L}_{\chi}^{(2)} = \frac{\mathbf{F}^2}{4} \langle u_{\mu} u^{\mu} + \chi_{+} \rangle$$

$$\mathcal{L}_V = \frac{F_V}{2\sqrt{2}} \langle V_{\mu\nu} f_{+}^{\mu\nu} \rangle + \frac{iG_V}{\sqrt{2}} \langle V_{\mu\nu} u^{\mu} u^{\nu} \rangle$$

$$\mathcal{L}_A = \frac{F_A}{2\sqrt{2}} \langle A_{\mu\nu} f_{-}^{\mu\nu} \rangle$$

$$\begin{aligned} \mathcal{L}_{VAP} &= \sum_i^5 \lambda_i \mathcal{O}(V_{\mu}, A_{\mu}, \Pi) \\ &= \lambda_1 \langle [V_{\mu\nu}, A_{\mu\nu}] \chi_{-} \rangle + \dots \end{aligned}$$

5 unknown
couplings

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Asymptotic behaviour of $\text{Im}\Pi_{\mu\nu}^A$

[Floratos, Narison, de Rafael, 1979]

$$\text{Im}\Pi_{\mu\nu}^A = \frac{1}{2} \sum_N \int d\rho_N \delta^{(4)}(q - p_N) \langle 0 | A_\mu | N \rangle \langle N | A_\nu^\dagger | 0 \rangle \xrightarrow{(\text{QCD}, q^2 \rightarrow \infty)} \text{Constant}$$

- Asymptotic behaviour of Form Factors (QCD)

$$f_i(\lambda_k) = 0, \quad i = 1, 2$$

3 unknown couplings

- Feynman diagrams : 1 only coupling, λ_0

$$4 \text{ parameters} \left\{ \begin{array}{l} \mathcal{L}_{\text{VAP}} \rightarrow \lambda_0 \\ a_1(1260) \rightarrow M_{a_1}, \Gamma_{a_1}(M_{a_1}^2), \alpha \end{array} \right.$$

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$\tau^- \rightarrow K^+ K^- \pi^- \nu_\tau$ CLEO VERSION

(F. Liu '03, CLEO-III '04)

$$F_V = -\frac{1}{2\sqrt{2}\pi^2 F^3} \sqrt{R_B} \sum_i BW_i$$

$$BW_{V,A}(x = s_i, Q^2) = \frac{M_{V,A}^2}{M_{V,A}^2 - x - i\sqrt{x} \Gamma_{V,A}(x)}$$

$\tau^- \rightarrow K^+ K^- \pi^- \nu_\tau$ CLEO VERSION

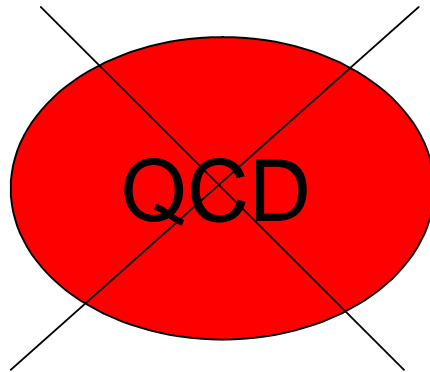
(F. Liu '03, CLEO-III '04)

$$F_V = \underbrace{-\frac{1}{2\sqrt{2}\pi^2 F^3}}_{L^{(4)}_{\chi, WZW}} \sqrt{R_B} \underbrace{\sum_i BW_i}_{\xrightarrow{Q^2 \rightarrow 0} 1}$$

$\tau^- \rightarrow K^+ K^- \pi^- \nu_\tau$ CLEO VERSION

(F. Liu '03, CLEO-III '04)

$$F_V = - \underbrace{\frac{1}{2\sqrt{2}\pi^2 F^3}}_{L^{(4)}_{\chi, WZW}} \sqrt{R_B} \sum_i \underbrace{BW_i}_{\xrightarrow{Q^2 \rightarrow 0} 1}$$



1.80 ± 0.53

Kühn & Santamaría Model

[Kühn, Santamaría, 1990]

✗

QCD

{	KS	$\chi\text{PT } \mathcal{O}(p^2) \checkmark$	$\chi\text{PT } \mathcal{O}(p^4) \boxtimes$
		Vector meson dominance	✓
		Asymptotic behaviour ruled by QCD	✓
		$BW_R = \frac{M_R^2}{M_R^2 - s - i \sqrt{s} \Gamma_R(s)}$?

Example : Vector form factor of the pion

$$F_V(s) = \frac{BW_\rho \left(\frac{1 + \alpha BW_\omega}{1 + \alpha} \right) + \beta BW_{\rho'} + \gamma BW_{\rho''} + \dots}{1 + \beta + \gamma + \dots}$$

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SHERPA

(Simulation of **H**igh **E**nergy **R**eactions of **P**Articles)

- Institute for Theoretical Physics at TU Dresden (Germany) and I3P Durham (UK)



- Group Leader : Frank Krauss



- Tevatron + LHC

- <http://www.physik.tu-dresden.de/~krauss/hep>

- <http://www.sherpa-mc.de/>

Insertion of our matrix elements for hadronic tau decays