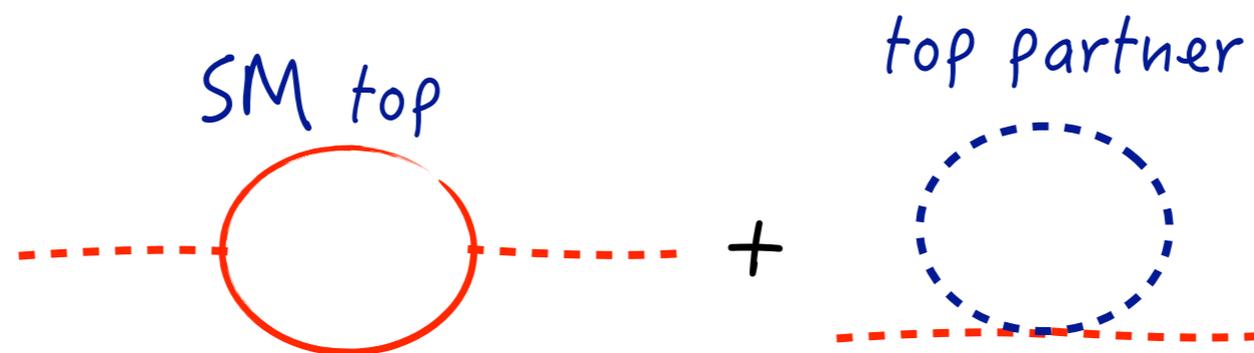
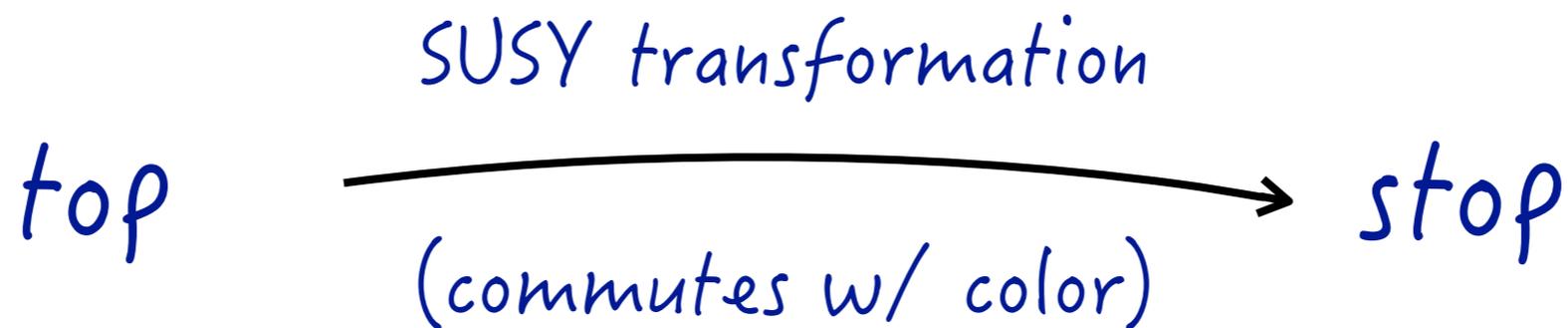


# Neutral Naturalness and the Higgs

Roni Harnik  
Fermilab

# The Tuning Problem

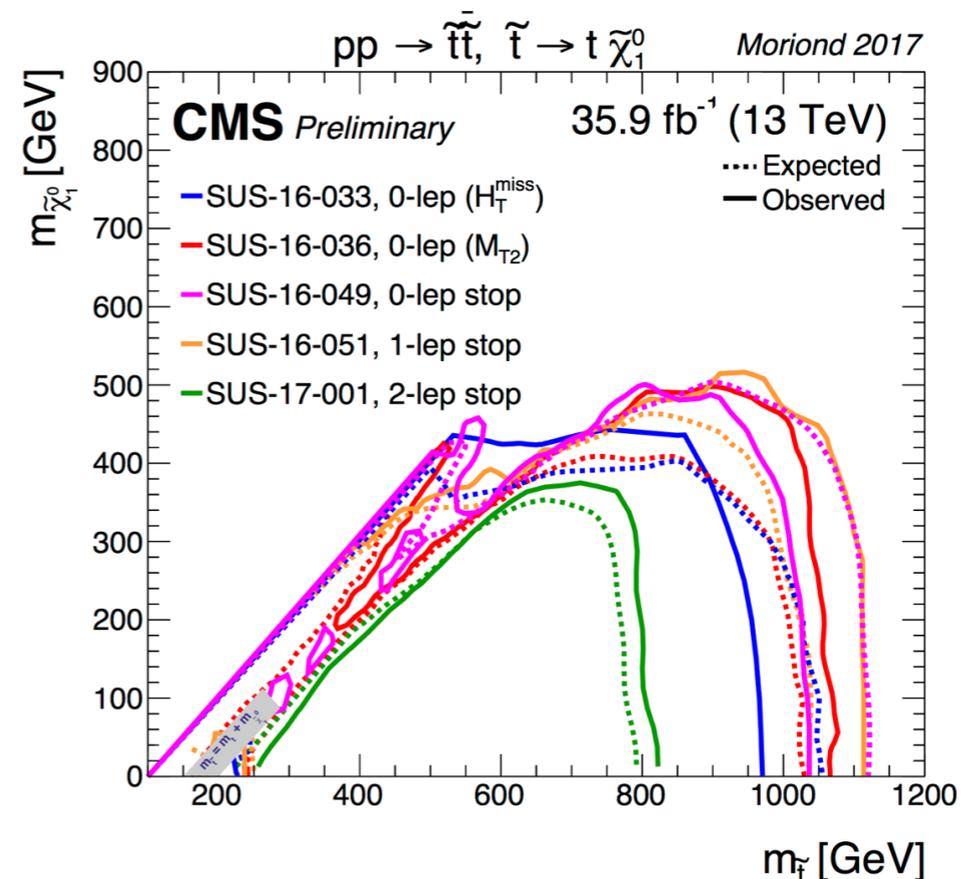
- The weak scale is fine tuned in the SM. a.k.a the hierarchy problem.
- This tuning can be removed by invoking a symmetry. e.g. in SUSY:



# The Tuning Problem

- Naturalness  $\rightarrow$  a colored top partner at  $\approx$ TeV.
- However, so far the LHC discovered just the Higgs and nothing else new..... (may change soon!)

The current manifestation of the tuning problem:

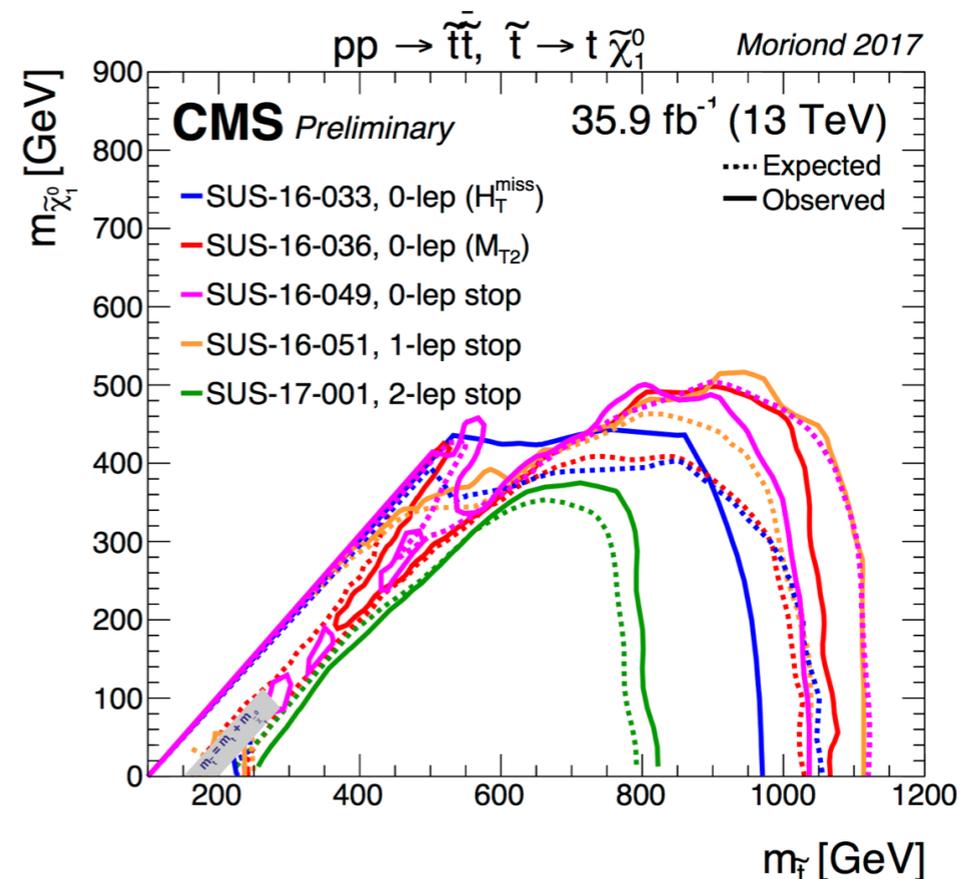


# The Tuning Problem

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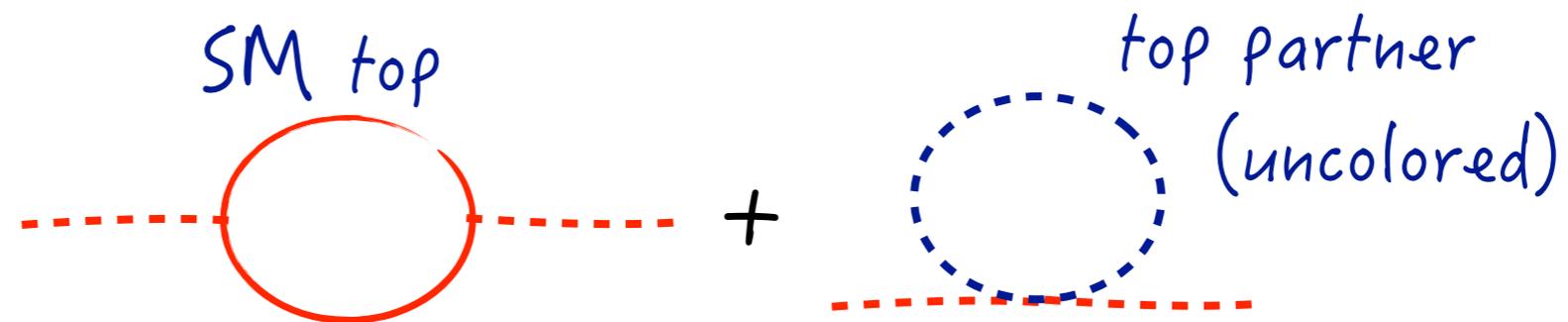
The current manifestation of the tuning problem:

Where the **hell** is everybody?!



# Neutral Naturalness

- Models of Neutral Naturalness address Higgs tuning w/ uncolored partners.

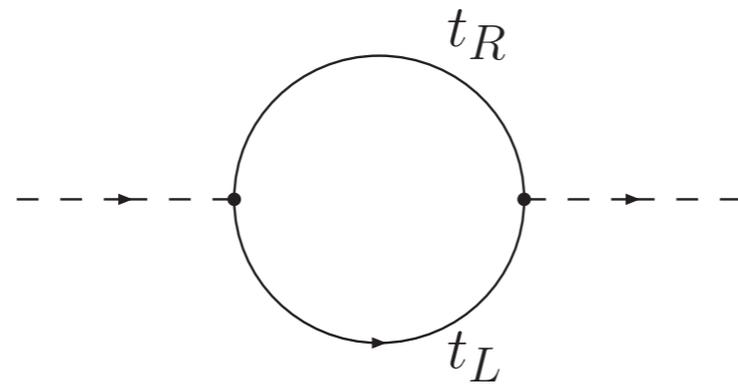


- Higgs Hunters are a good target audience:

Precision and exotic Higgs measurements may be the only LHC discovery channels.

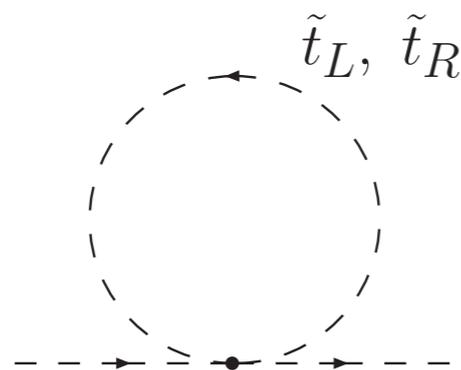
# Neutral Naturalness Models

color factor:



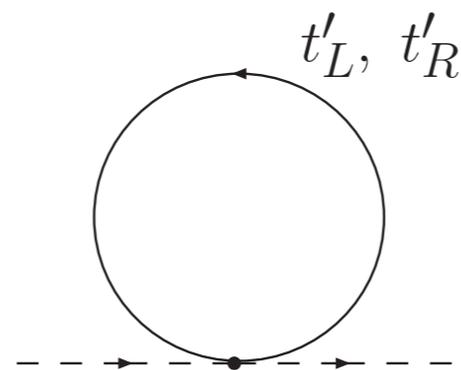
$\times 3$

Standard Model



Supersymmetry

or

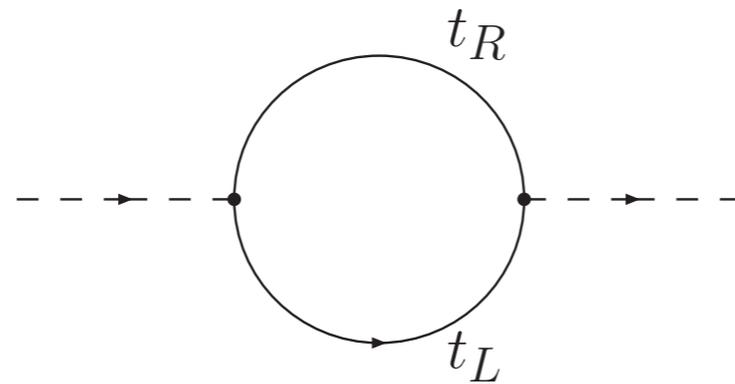


Composite Higgs

$\times 3$

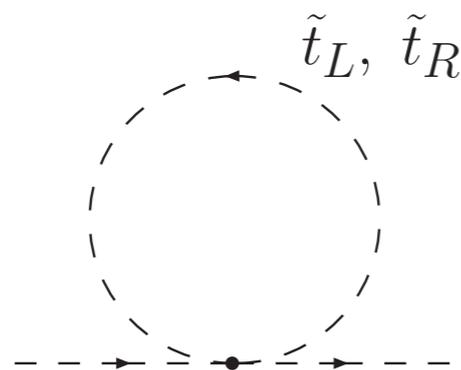
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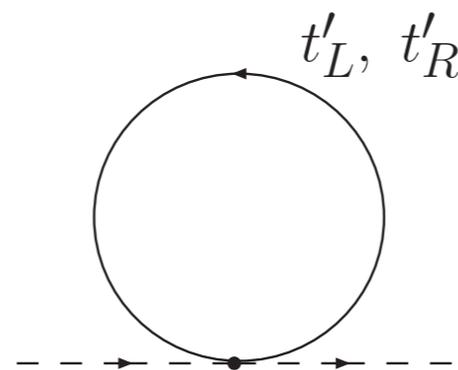
$\times 3$

Standard Model



Supersymmetry

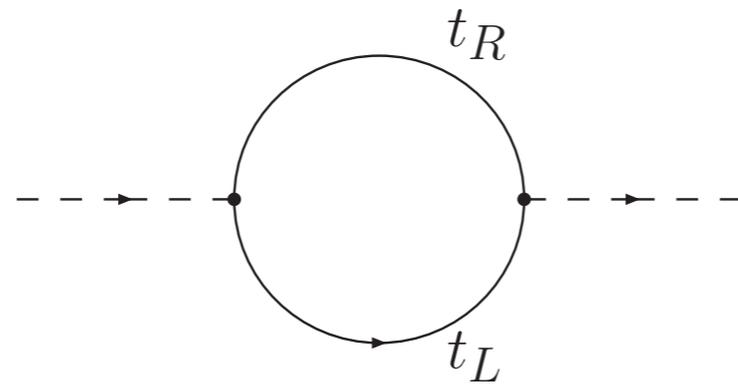
or



Composite Higgs

~~$\times 3$~~   
 $3'$

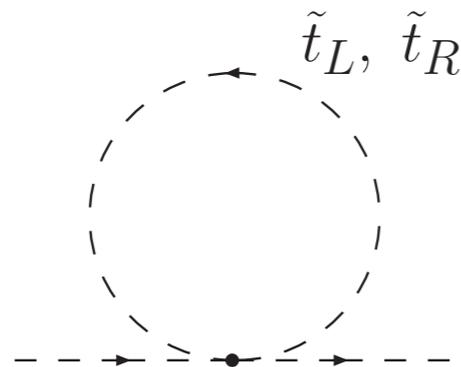
# Neutral Naturalness Models



Standard Model

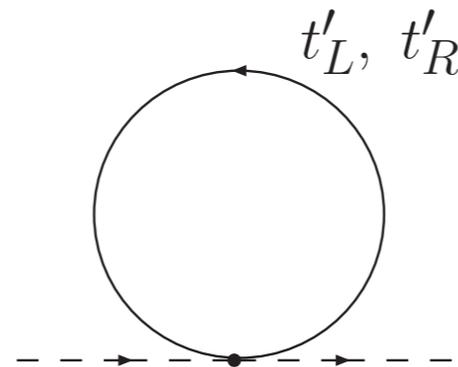
color factor:

$\times 3$



Supersymmetry

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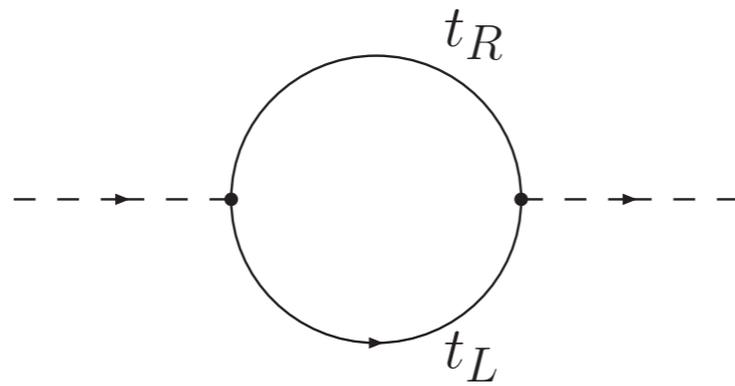


Composite Higgs

~~$\times 3$~~   
 $3'$

symmetry does not commute with color.

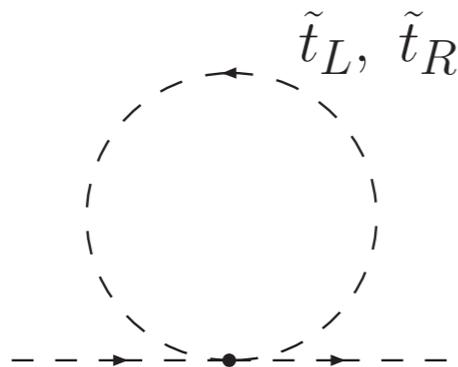
# Neutral Naturalness Models



Standard Model

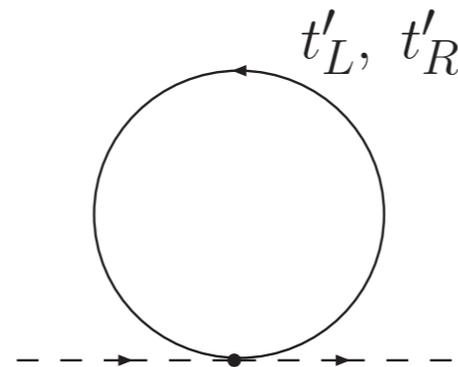
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Supersymmetry

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3'

↓  
**Folded SUSY**

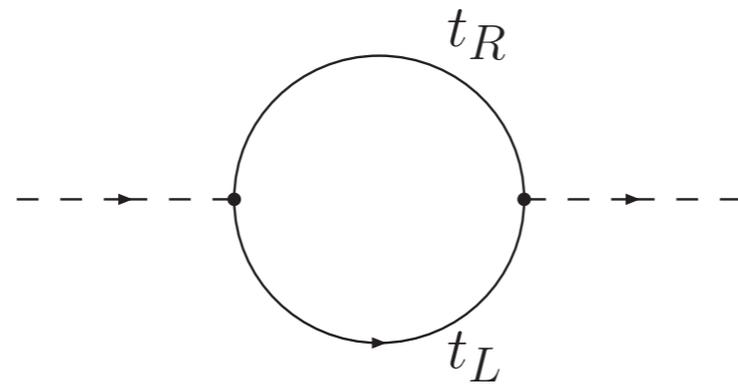
Burdman, Chacko, Goh, RH (06')

↓  
**Twin Higgs**

Chacko, Goh, RH (05')

symmetry does not commute with color.

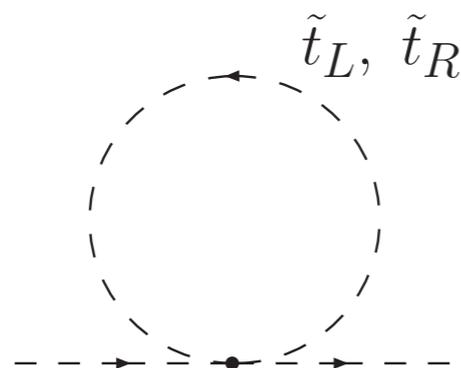
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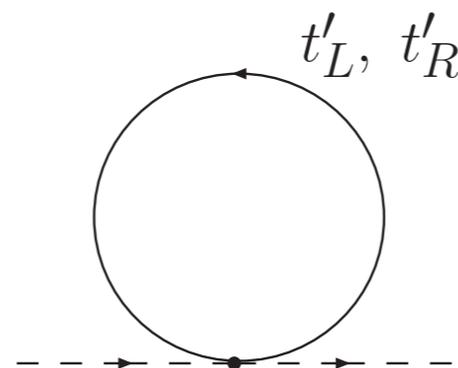
color factor:

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Supersymmetry

or



Composite Higgs

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3'

↓  
**Folded SUSY**

Burdman, Chacko, Goh, RH (06')

↓  
**Twin Higgs**

Chacko, Goh, RH (05')

symmetry does not commute with color.

More Example since: Quirky Little Higgs, Dark Top, Top partners as RH Neutrinos, Orbifold Higgs, Fraternal Twin Higgs,

An example:

Twin Higgs

# Mirror Symmetry

\* Example: Mirror Symmetry

$$(SM_A) \times (SM_B)$$



$Z_2$

An exchange symmetry.  $A \longleftrightarrow B$ .

This  $Z_2$  does not commute with  $SU(3)_{\text{QCD}}$ .  
 Can it play a role in protecting the Higgs? (yes!)

# Twin Higgs (the Disney version)

- \* The hierarchy problem stems from the mass term in the Higgs potential - a quadratic

$$\text{---} \bigcirc \text{---} \propto c \Lambda^2 H^2$$

- \* In a mirror model this becomes -

$$\text{---} \overset{A}{\bigcirc} \text{---} + \text{---} \overset{B}{\bigcirc} \text{---} \propto c \Lambda^2 (H_A^2 + H_B^2)$$

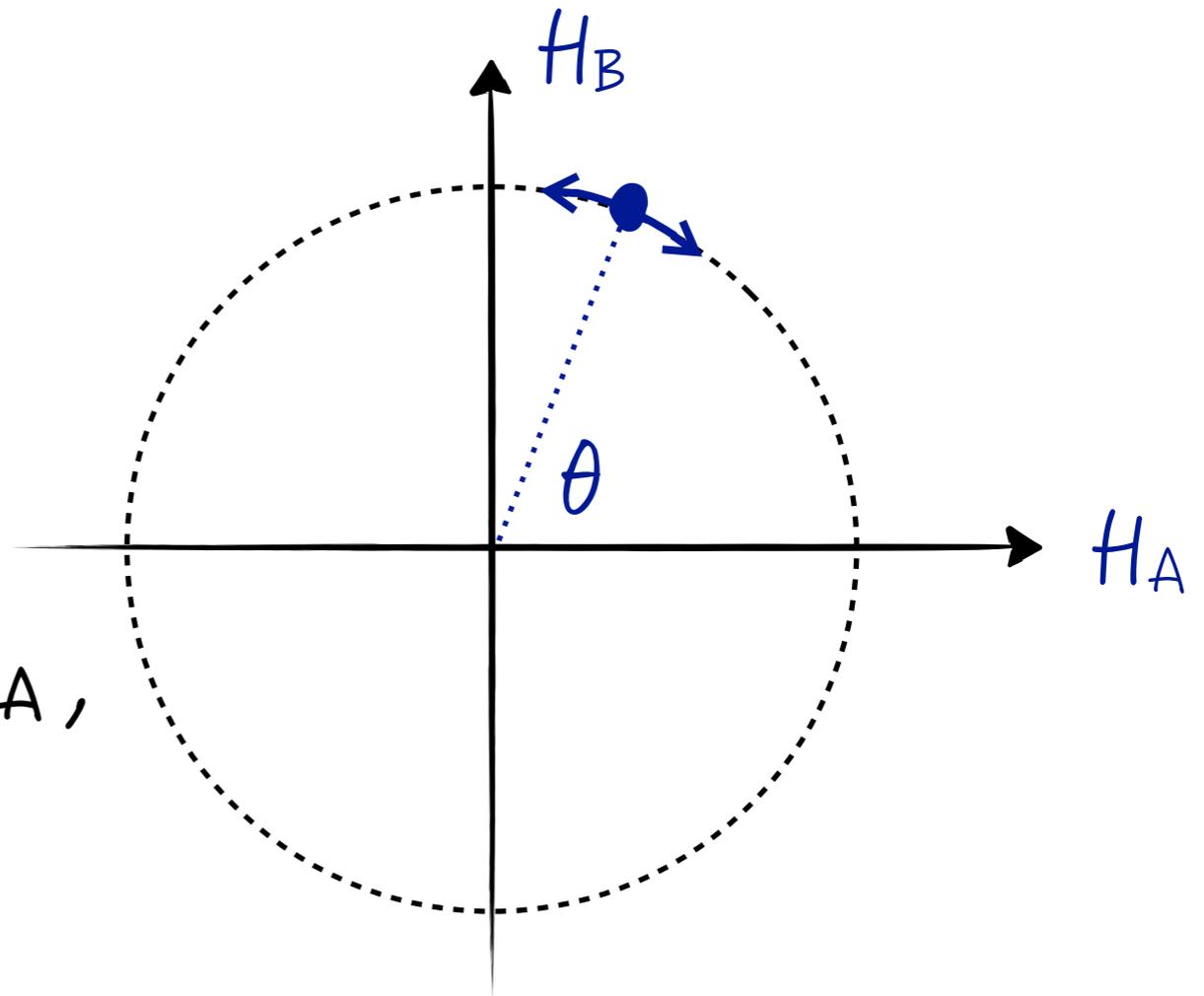
Note:  $(H_A^2 + H_B^2)$  has a higher symmetry of continuously rotating  $H_A$  and  $H_B$  into one another.

# Twin Higgs

- \* The two Higgs doublets form a 4-plet of a global  $SU(4)$  symmetry (or  $SO(8)$  in composite models).
- \*  $SU(4)$  is broken spontaneously at a scale  $f$ .  
(mostly by the other Higgs!)

\* One Goldstone boson is left uneaten.

\* The pNGB is mostly  $H_A$ ,  
(but a bit  $H_B$  too!).



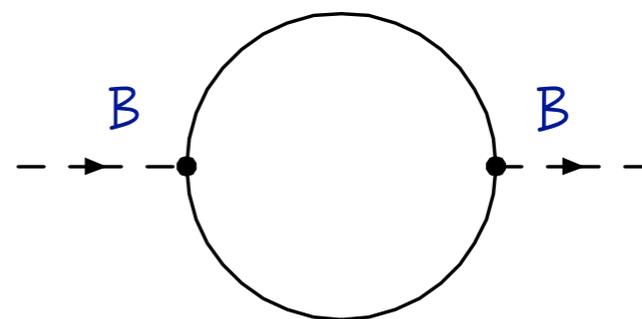
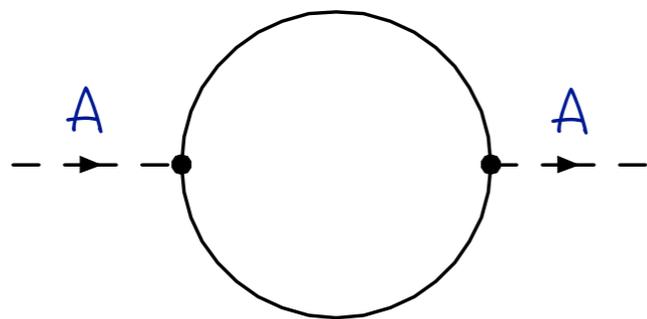
# Cancellation

\*  $Z_2$  guarantees  $(H_A^2 + H_B^2)$  structure at quadratic level.  $SU(4)$  protects the Higgs!

---

\* In diagrams, there is a cancellation between loops.

\* Recall 'Higgs portal' coupling SM's -  
[the Cross term in  $\lambda(H_A^2 + H_B^2)^2$ ]



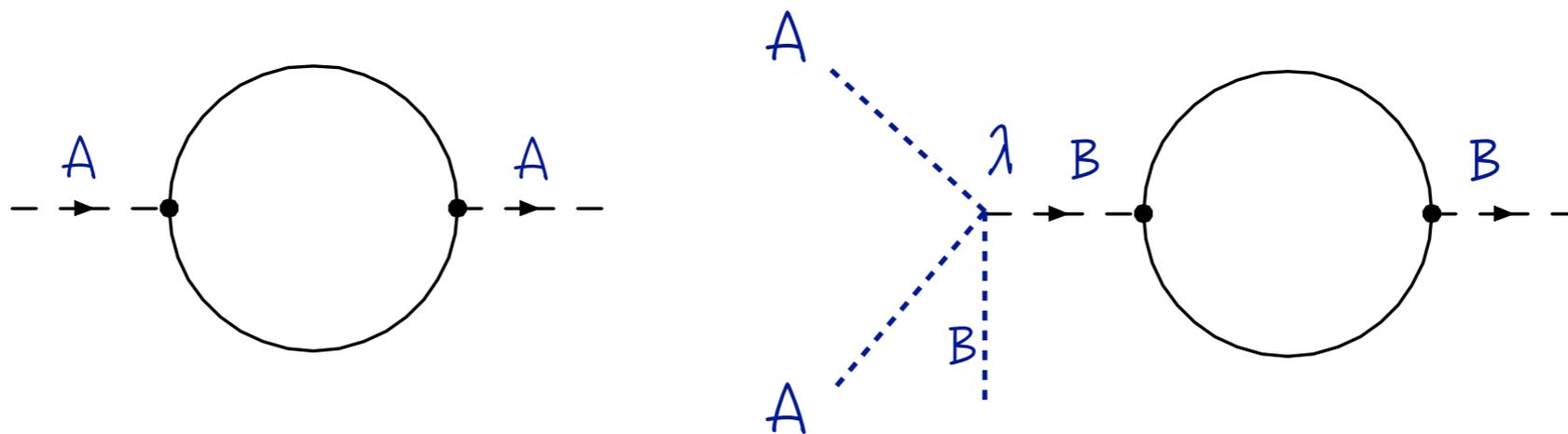
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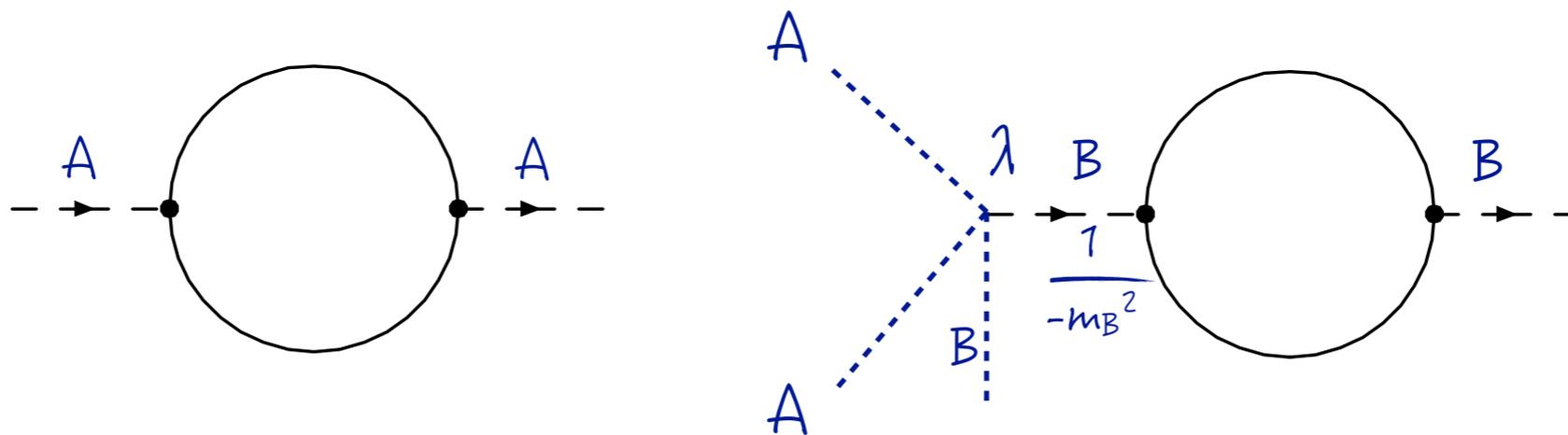
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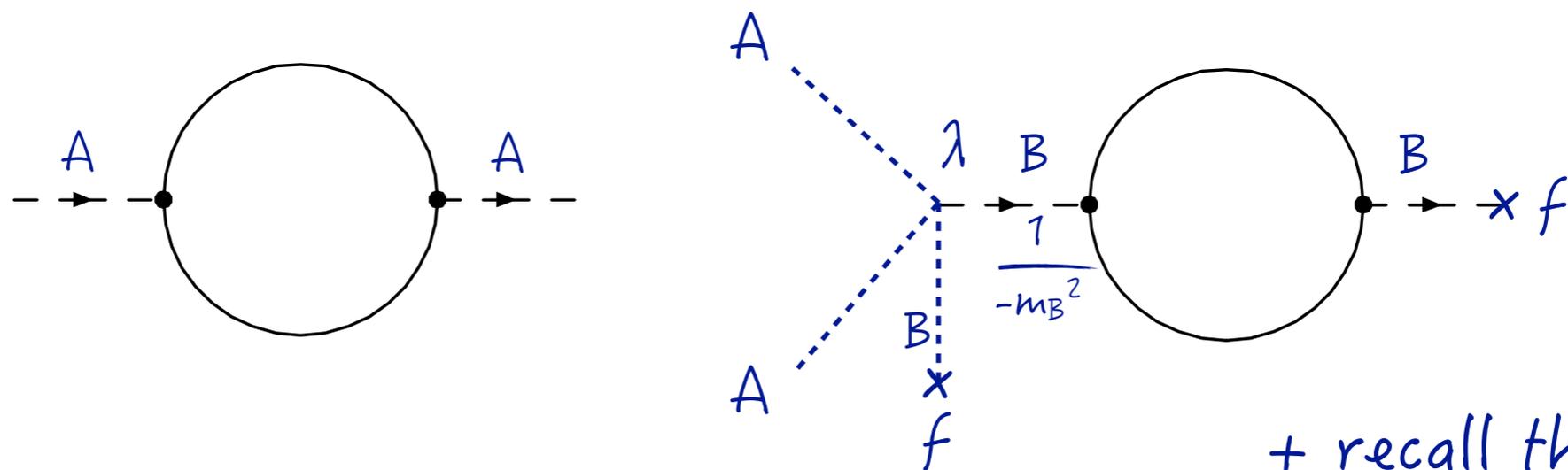


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 [the Cross term in  $\lambda(H_A^2 + H_B^2)^2$ ]



+ recall that  $m_B^2 = \lambda f^2$

# So...

- \* Let's summarize what we have:
  - Higgs is protected by a symmetry.
  - The model is natural up to  $\Lambda$  beyond LHC scale.
  - All new particles below  $\Lambda$  are complete SM singlets!
- \* What's the phenomenology?
  - LHC Run 1 finds the Higgs and nothing else! (check).
  - Now what?...

# Details:

Interest in Neutral Naturalness is re-kindled.

In recent years many issues are studied-

- \* UV Completions

[Craig, Howe][Geller, Telem][Barbieri et al][Low, Wang, Tesi]..

- \* Dark Matter

[March-Russell et al][Craig, Katz][Farina]...

- \* Cosmology

[Barbieri, Hall][Chacko, Craig, Fox, RH][Csacki, Kuflik, Lombardo]

- \* Improved Tuning

[RH, Howe, Kearny][Gregoire et al][Redigolo et al]

- \* LHC Pheno: Higgs is key! (The rest of this talk)

# Twin Higgs Phenomenology

Depends strongly on how  
strictly we enforce the  $Z_2$ :

# Mirror or Fraternal?

\*  $SM_A \times SM_B$ . But what does  $SM_B$  really contain?

Mirror Twin Higgs

$Z_2$  is nearly exact.  
 $SM_B$  is a full copy of our  
 $SM$ , with a slightly higher  
 Higgs vev.

Fraternal Twin Higgs

$Z_2$  only acts on the  
 essentials.  $SM_B$  contains only  
 the 3rd gen and 3-2  
 gauge bosons.

# Mirror or Fraternal?

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Mirror Twin Higgs	Fraternal Twin Higgs
<p><math>Z_2</math> is nearly exact.</p> <p><math>SM_B</math> is a full copy of our SM, with a slightly higher Higgs vev.</p>	<p><math>Z_2</math> only acts on the essentials. <math>SM_B</math> contains only the 3rd gen and 3-2 gauge bosons.</p>
<p>Simple.</p> <p>Cosmology can be fixed</p> <p>LHC elusive (almost)</p>	<p>Minimal but may be theoretically cumbersome.</p> <p>Cosmology's fine.</p> <p>Displaced decays @ LHC</p>

*Identical*

# Higgs Couplings

- \* Recall that the pNGB is only mostly our Higgs boson, but partially theirs.
- \* A mixing angle governing this is  $\theta \sim \frac{v}{f} = \frac{\text{our vev}}{\text{their vev}}$
- \* For both identical and fraternal models, Higgs couplings to SM states are modified by  $\cos(\frac{v}{f})$ .

All SM Higgs  $\sigma$ xBR's are modified by  $\cos^4(\frac{v}{f})$

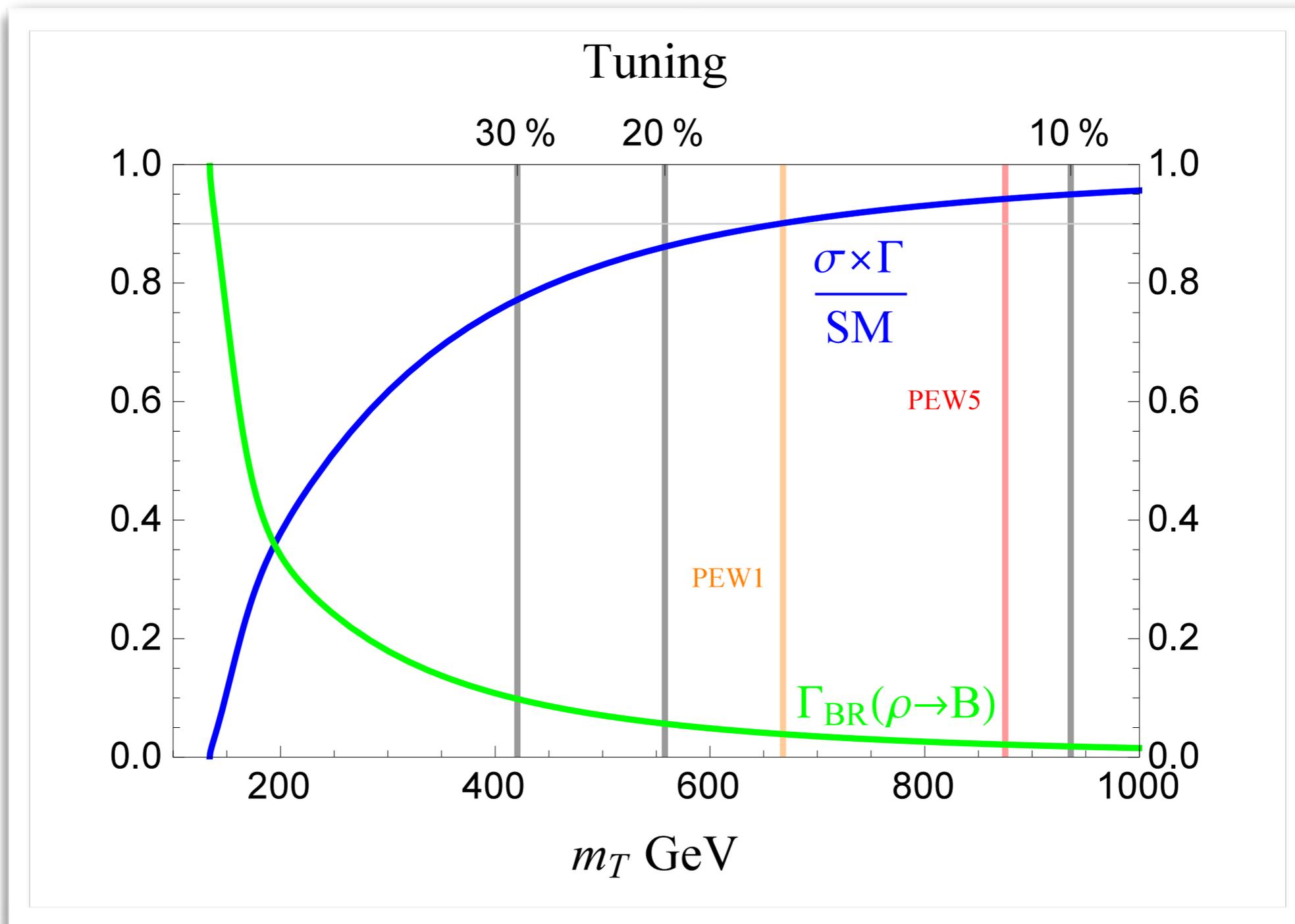
# Identical predictions

- \* In the identical model:  
Roughly  $\sin(\frac{v}{f})$  of the pNGB is the other Higgs.
- \* The 125 GeV Higgs can decay to the other SM.
- \* Twin bottoms cascade down invisible states!

$$BR_{(h \rightarrow inv)} = \sin^2\left(\frac{v}{f}\right)$$

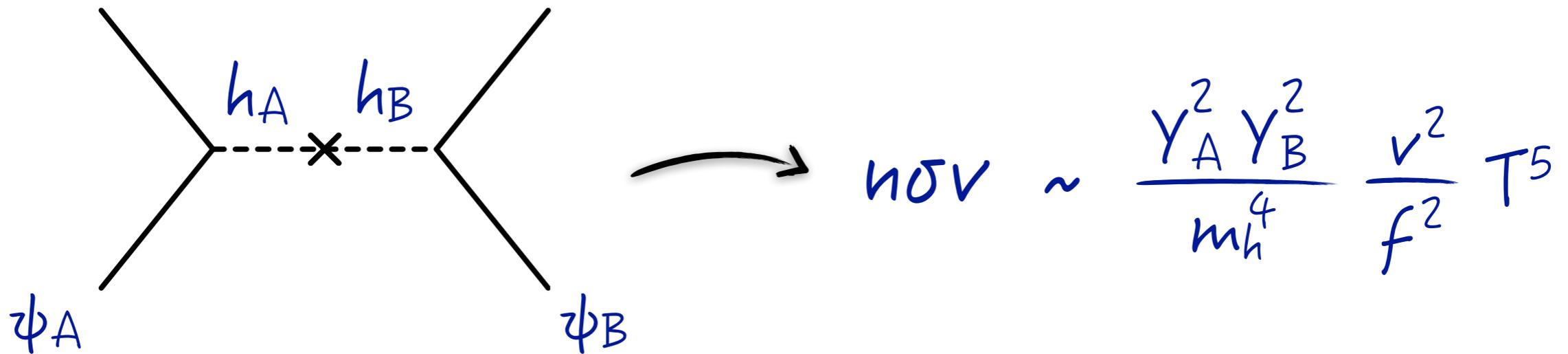
- \* For identical models,  $v/f$ , is setting both Higgs coupling modification and  $BR_{inv}$ . A prediction!

# Identical prediction



# Cosmology

- The Higgs portal keeps the two sectors in thermal eq. down to  $\sim 3$  GeV.



If the Universe was ever above 3 GeV,  
we thermalize the twin sector.

$$\Delta N_{\text{eff}} \sim 6 \quad :-)$$

e.g. [Barbieri, Gregoire, Hall] [Chacko, Craig, Fox, RH]

# Fixing Cosmology

- Easy fix: right handed neutrinos at few GeV.
- The singlet  $N$  is:
  - A-B ambivalent
  - Decouples relativistically
  - Decays preferentially to A sector (like  $f^2/v^2$ )
- $N$  comes to dominate the Universe, decays, and reheats A sector more.

$\Delta N_{\text{eff}}$  is set by  $v/f$  and  $N$  mass & width.  
(predictive!)

# Fraternal

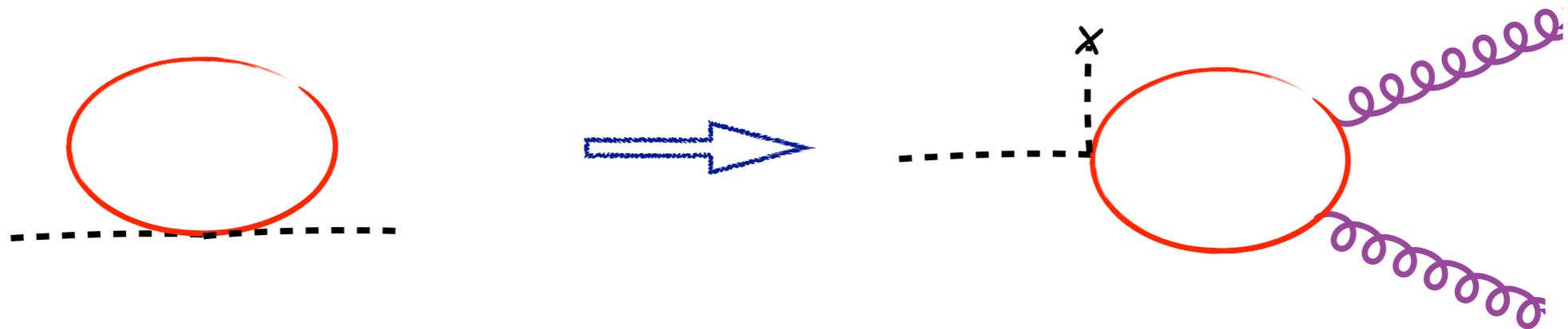
[see also talk by A. Katz]

- \* Invisible decay is reduced.
- \* All SM Higgs  $\sigma \times BR$ 's are modified by  $\cos^2\left(\frac{v}{f}\right)$

and...

# Fraternal: Exotic Decays

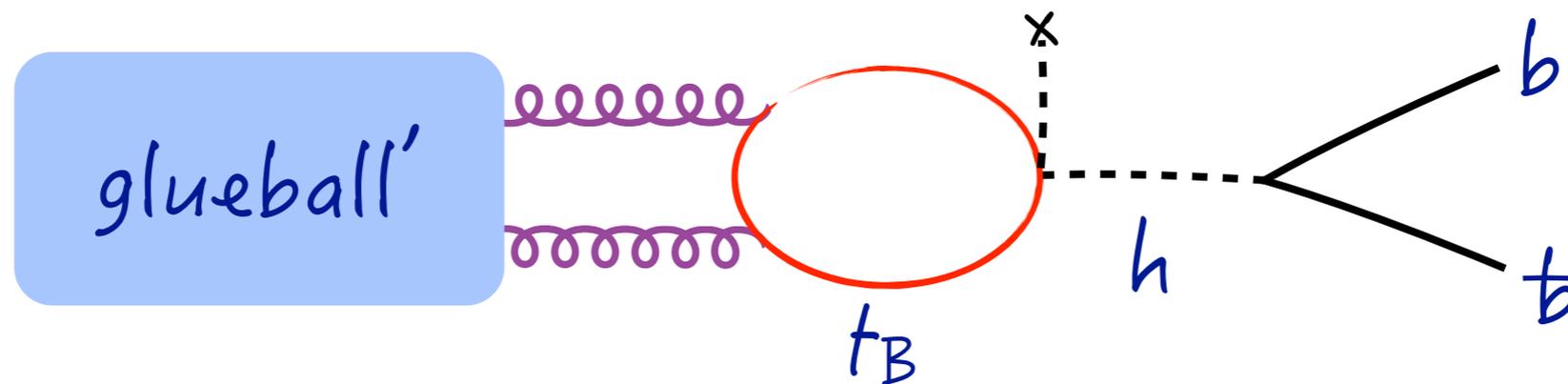
- \* For fraternal model we removed light stuff. But  $gluon_B$  is light. Light gluball!
- \* More generally, if Higgs naturalness done by QCD' colored states:



Higgs decays to glueball<sub>B</sub>'s!  
 Branching fraction of  $\approx 10^{-4}$ .

# Fraternal: Exotic Decays

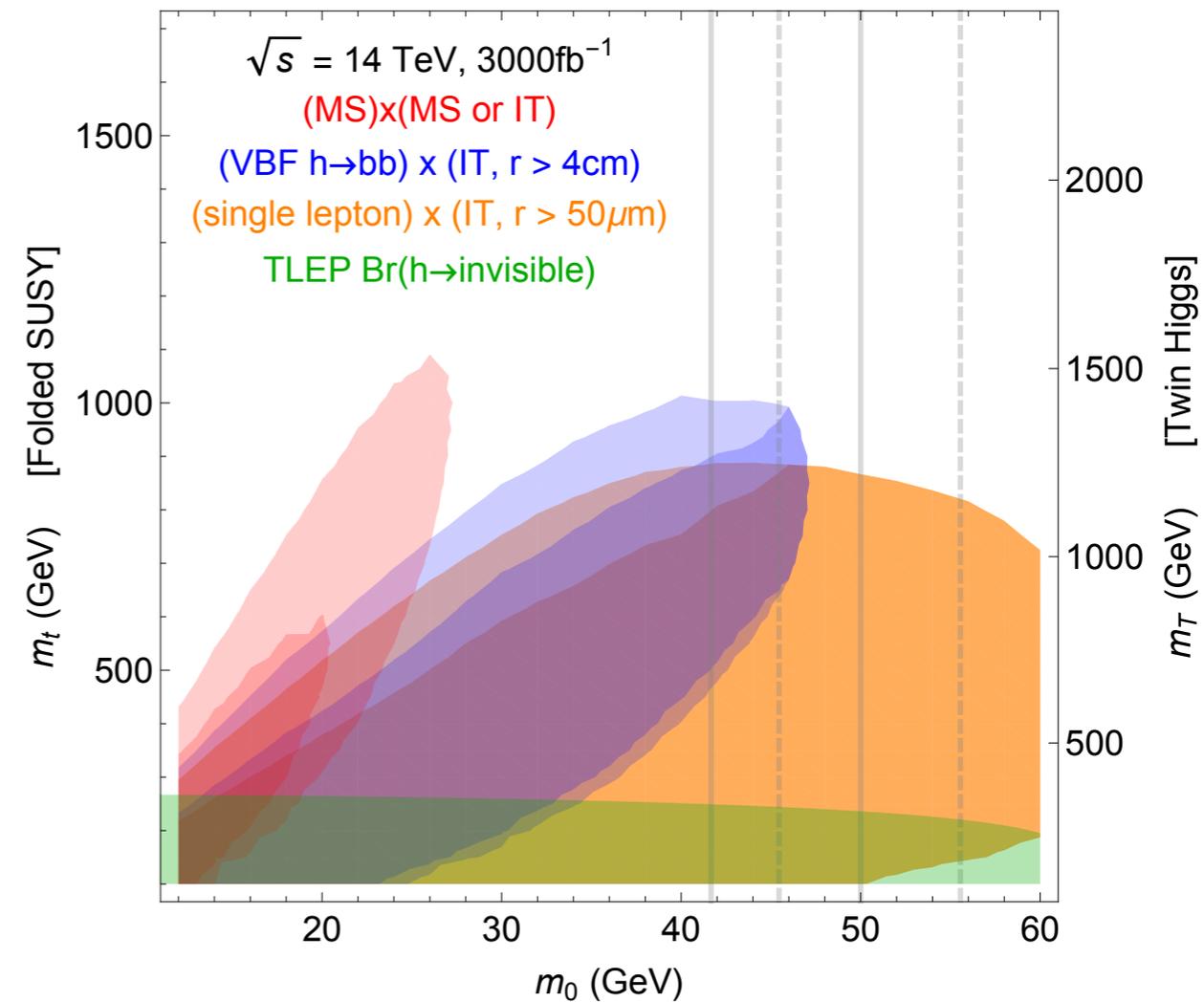
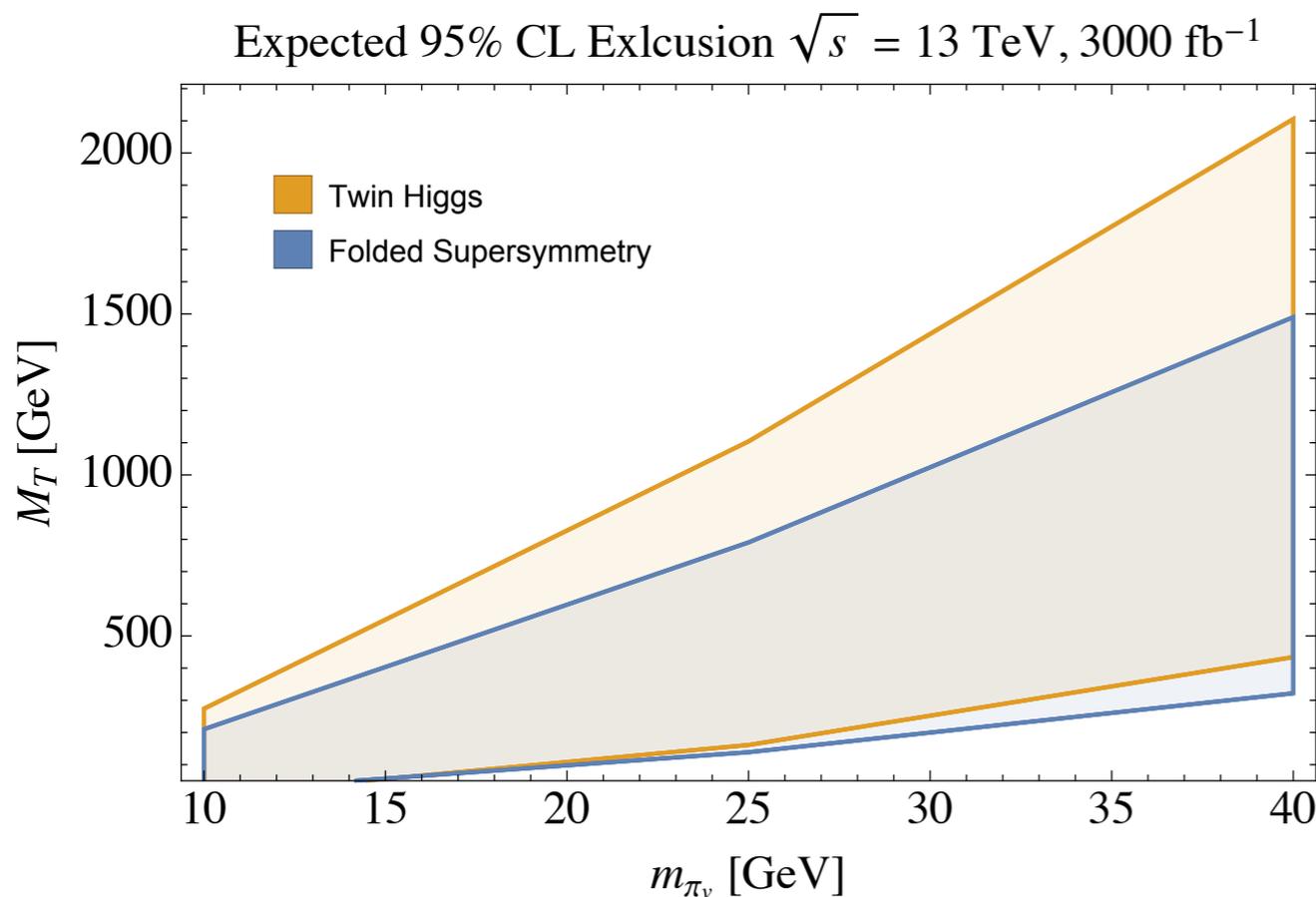
- \* The glueball is not invisible!  
The scalar glueball mixes with the Higgs.
- \* Decays are typically displaced! Low BG!



$c\tau$  depends strongly on parameters.  
ranges for prompt to km.

# Displaced Higgs Decays

- \* Displaced decays are a clean signature!
- \* Recent analyses show impressive reach to top partners at  $\sim$ TeV range!

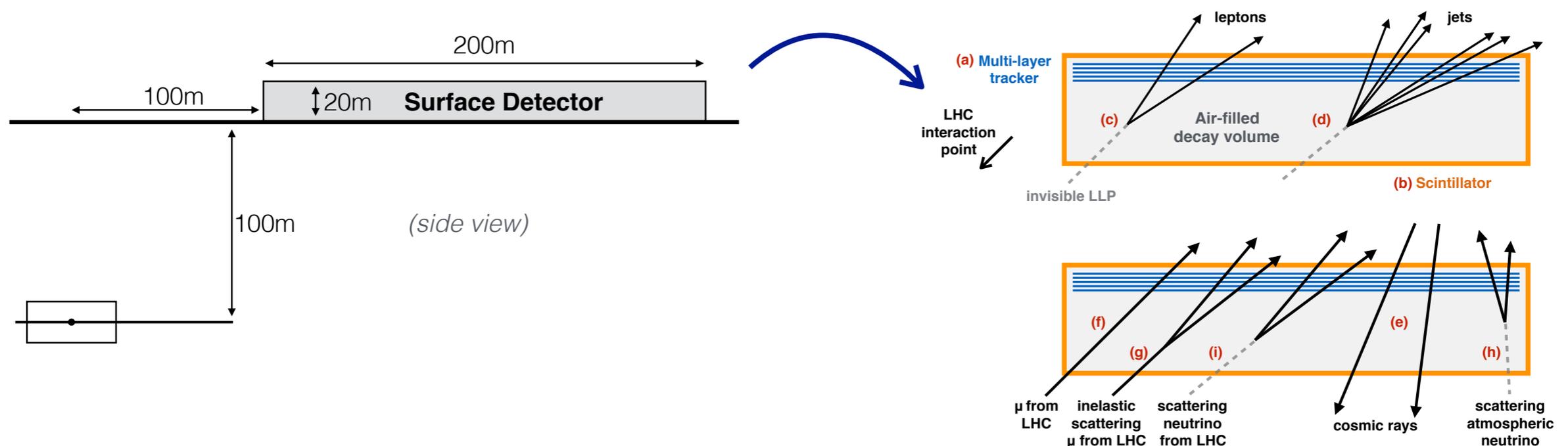


# A New Detector ?

- Pushing the limit on displaced decays in this model lead to a new proposal: **MATHUSLA**

## New Detectors to Explore the Lifetime Frontier

J. P. Chou, D. Curtin, H. J. Lubatti



# Conclusions & Outlook

- \* Models of Neutral Naturalness, e.g. Twin Higgs, elegantly address the “where-is-everybody” problem. Attracting renewed attention.
- \* They motivate a new set of LHC signatures:
  - Higgs Coupling modifications
  - Invisible Higgs decay
  - Exotic and displaced Higgs decays
- \* Higgs measurements can be the key to figuring out the Hierarchy problem!

Deleted Scenes

# Recent Literature

- \* **Twin Dark Matter:** Craig and Katz (15), Knapen and Tsai (16), March Russel et al (15), Farina (15)
- \* **Flavor:** Csaki, Geller, Telem, Weiler (15),
- \* **Dark Photon & Flavor:** RH, Stamou, Zupan (soon)
- \* **Cosmology:** Barbieri, Gregoire, Hall (05), Barbieri, Hall (17), Chacko, Craig, Fox, RH (17), Csacki, Kuflik, Lombardo (17)
- \* **Models:** Cheng Hall Weiner (06), Craig & Howe (14), Geller & Telem (14), Barbieri et al (15), Low, Tesi, Wang (15), Greiore et al (16), RH, Howe Kearny (16), Craig Knapen, Strassler (16), Redgigolo et al (16)

And many more.

# Precision EWK

- \* Precision EW measurements place a constraint on the scale  $f$  but depend on UV completion.
- \* SM Higgs loops contribute to  $S$  &  $T$   
→ modified Higgs couplings are constrained.
- \* Coupling modifications are “made up” by states at cutoff or by heavy Higgs for strong/weak UV completion (respectively).

$$\Delta S \approx \frac{1}{6\pi} \left(\frac{v}{f}\right)^2 \log\left(\frac{m_{h_2}}{m_h}\right) \quad \Delta T \approx -\frac{3}{16\pi \cos^2 \theta_W} \left(\frac{v}{f}\right)^2 \log\left(\frac{m_{h_2}}{m_h}\right)$$

# A Toy Example

- \* A global  $SU(4)$  symmetry w/ one fundamental:

$$V(H) = -m^2 |H|^2 + \lambda |H|^4$$


$$\langle |H|^2 \rangle = \frac{M^2}{2\lambda} \equiv f^2$$


$$SU(4) \longrightarrow SU(3)$$

7 Goldstones

# Cancelation

- \* How does the twin cancelation come about?
- \* Lets think about the theory of Goldstones:  
(a.k.a. broken  $SU(4)$  generators)

$$\Pi = \left( \begin{array}{ccc|c} 0 & 0 & 0 & h_1 \\ 0 & 0 & 0 & h_2 \\ 0 & 0 & 0 & 0 \\ \hline h_1^* & h_2^* & 0 & 0 \end{array} \right)$$

This beast transforms non-linearly under  $SU(4)$ .

For convenience,  
construct a linearly  
transforming combination:

$$H = \begin{pmatrix} H_A \\ H_B \end{pmatrix} = \exp\left(\frac{i}{f}\Pi\right) \begin{pmatrix} 0 \\ 0 \\ 0 \\ f \end{pmatrix}$$

# Cancelation

\* Expanding: 
$$H_A = h \frac{if}{\sqrt{h^\dagger h}} \sin\left(\frac{\sqrt{h^\dagger h}}{f}\right) = ih + \dots,$$
$$H_B = \begin{pmatrix} 0 \\ f \cos\left(\frac{\sqrt{h^\dagger h}}{f}\right) \end{pmatrix} = \begin{pmatrix} 0 \\ f - \frac{1}{2f} h^\dagger h + \dots \end{pmatrix}.$$

\* Back to the top Yukawa:

$$\begin{aligned} \mathcal{L} &\supset y_t H_A \bar{t}_A t_A + y_t H_B \bar{t}_B t_B \\ &= y_t h \bar{t}_A t_A + y_t \left( f - \frac{|h|^2}{2f} \right) \bar{t}_B t_B + \dots \end{aligned}$$

# Cancelation

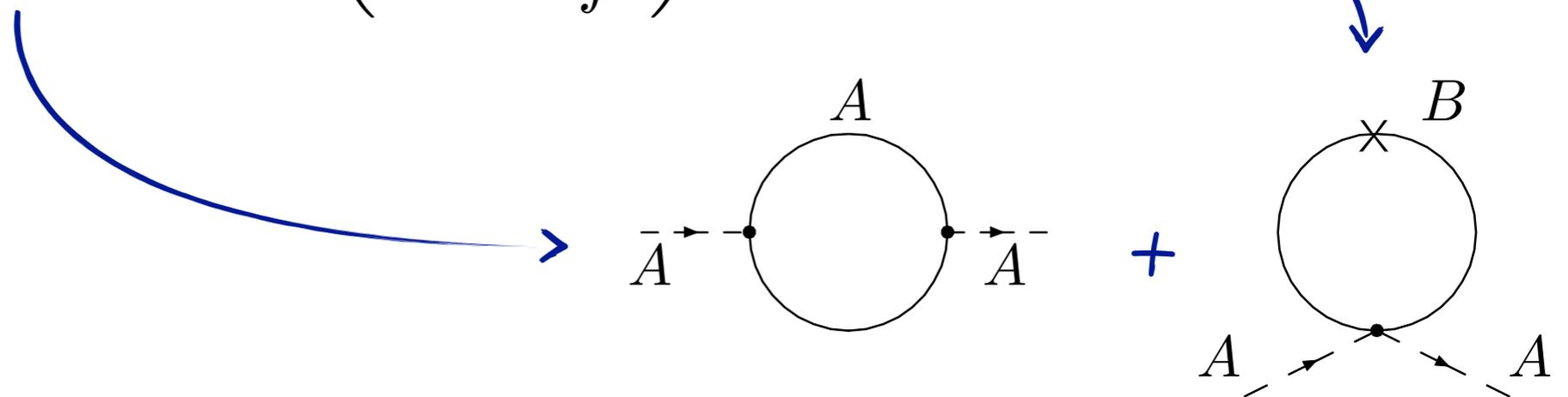
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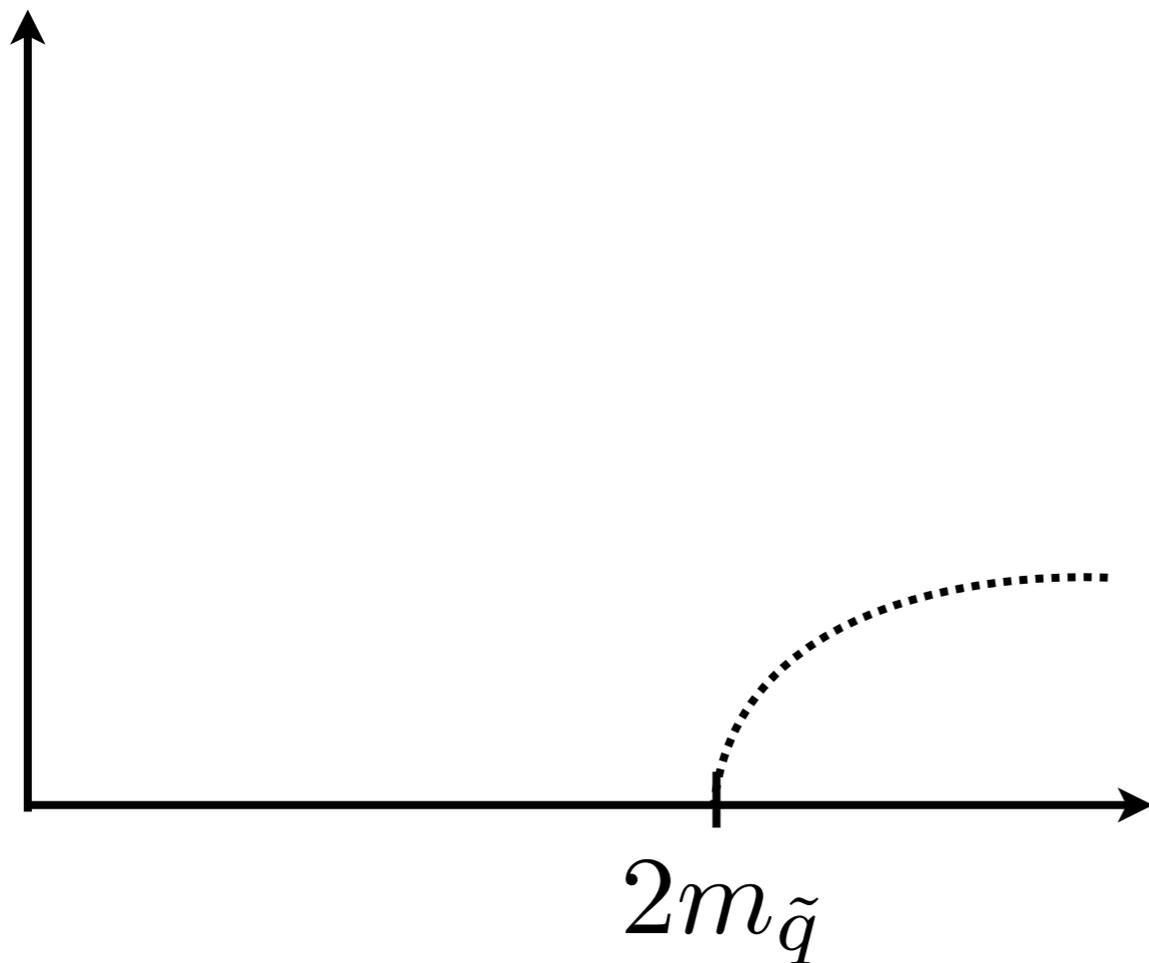
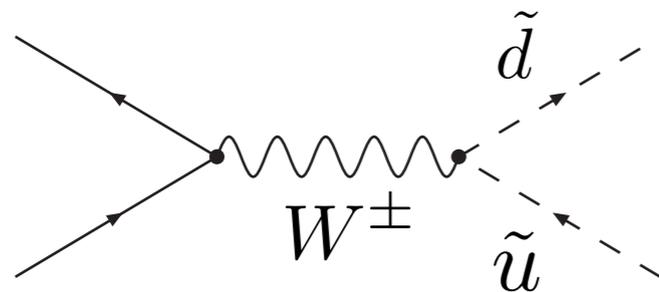
$$\mathcal{L} \supset y_t H_A \bar{t}_A t_A + y_t H_B \bar{t}_B t_B$$

$$= y_t h \bar{t}_A t_A + y_t \left( f - \frac{|h|^2}{2f} \right) \bar{t}_B t_B + \dots$$



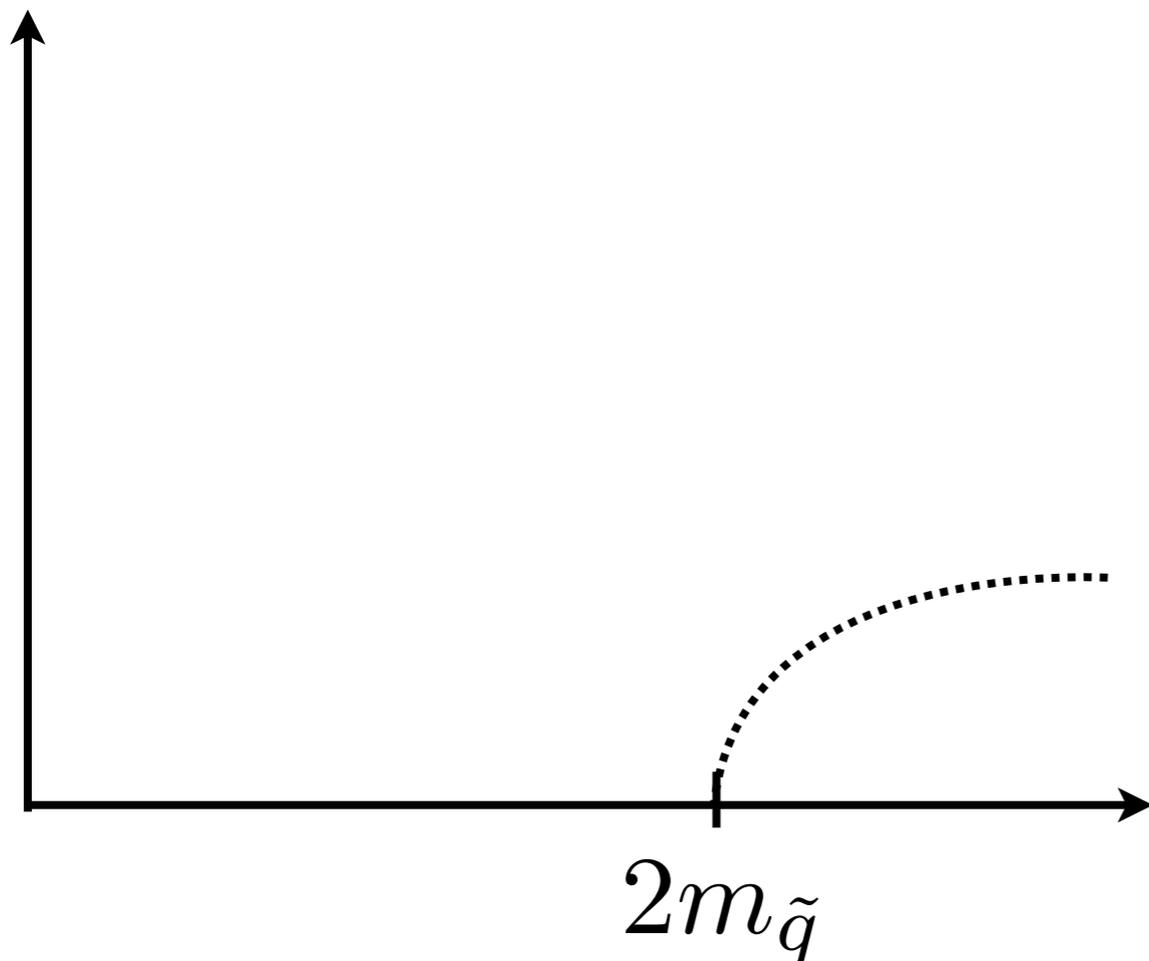
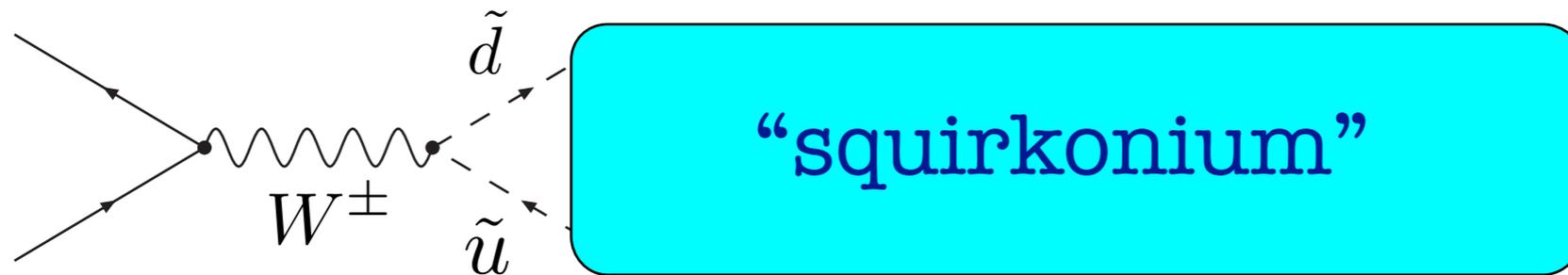
# A Hard Signal

\* Annihilation occurs after radiation:



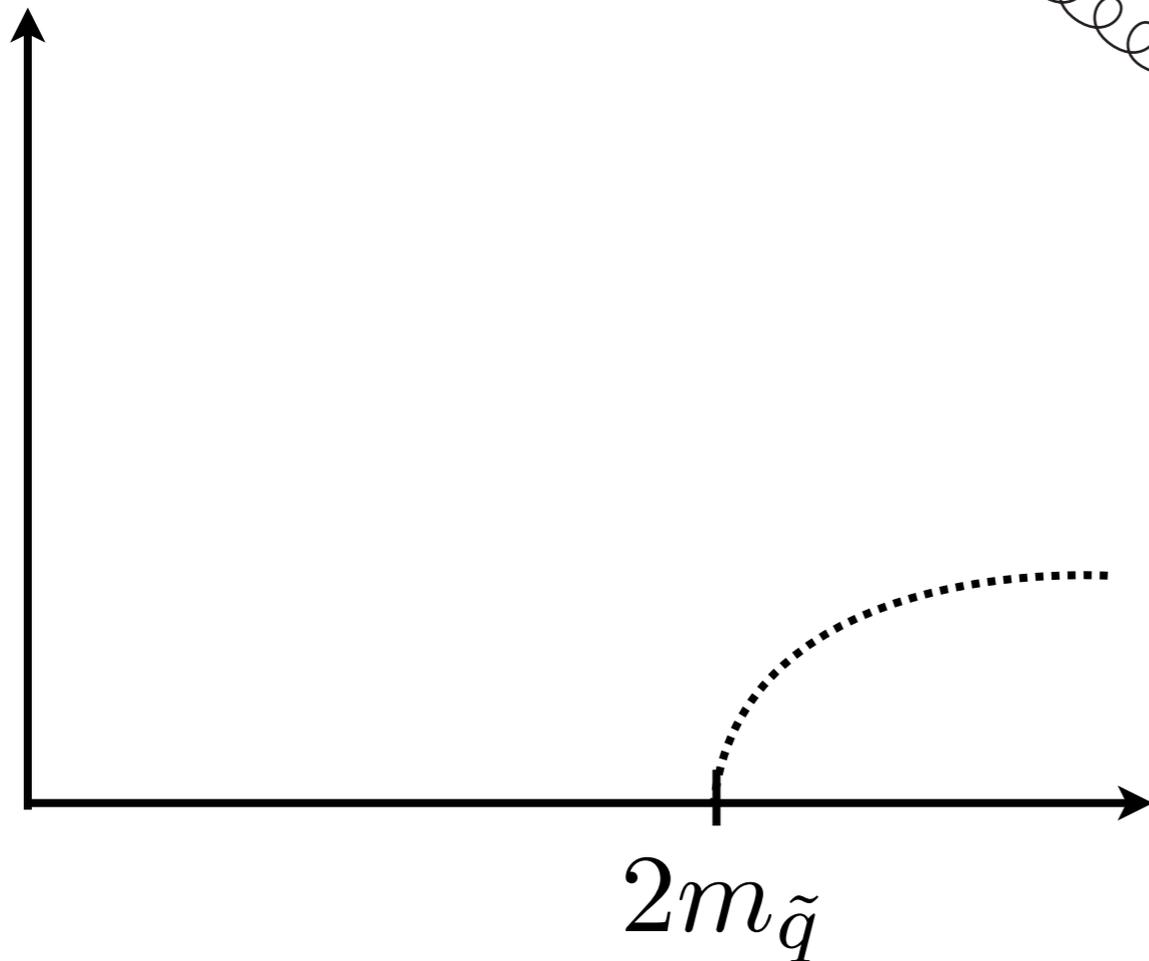
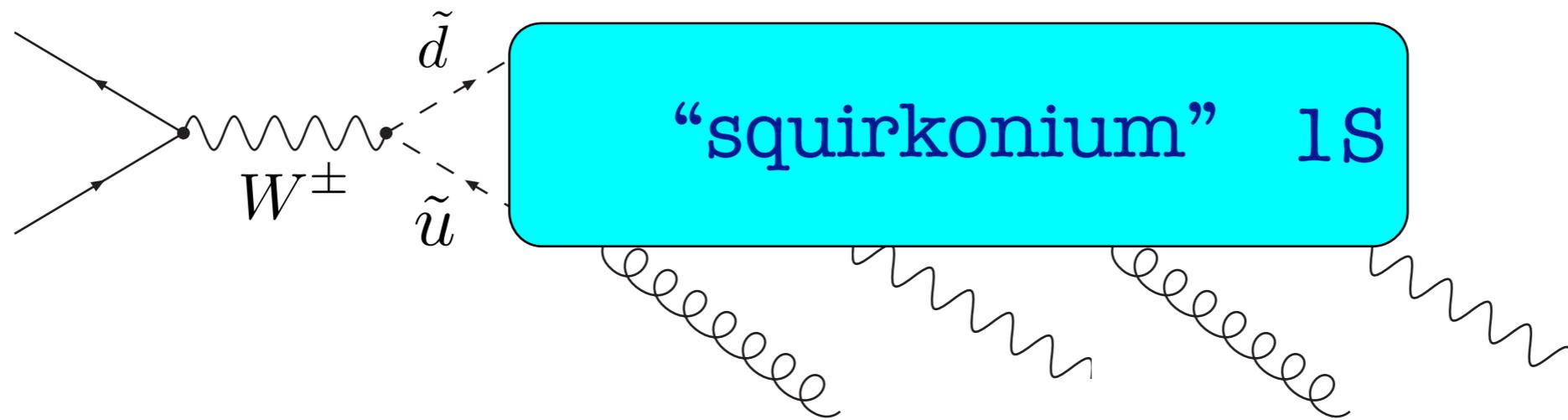
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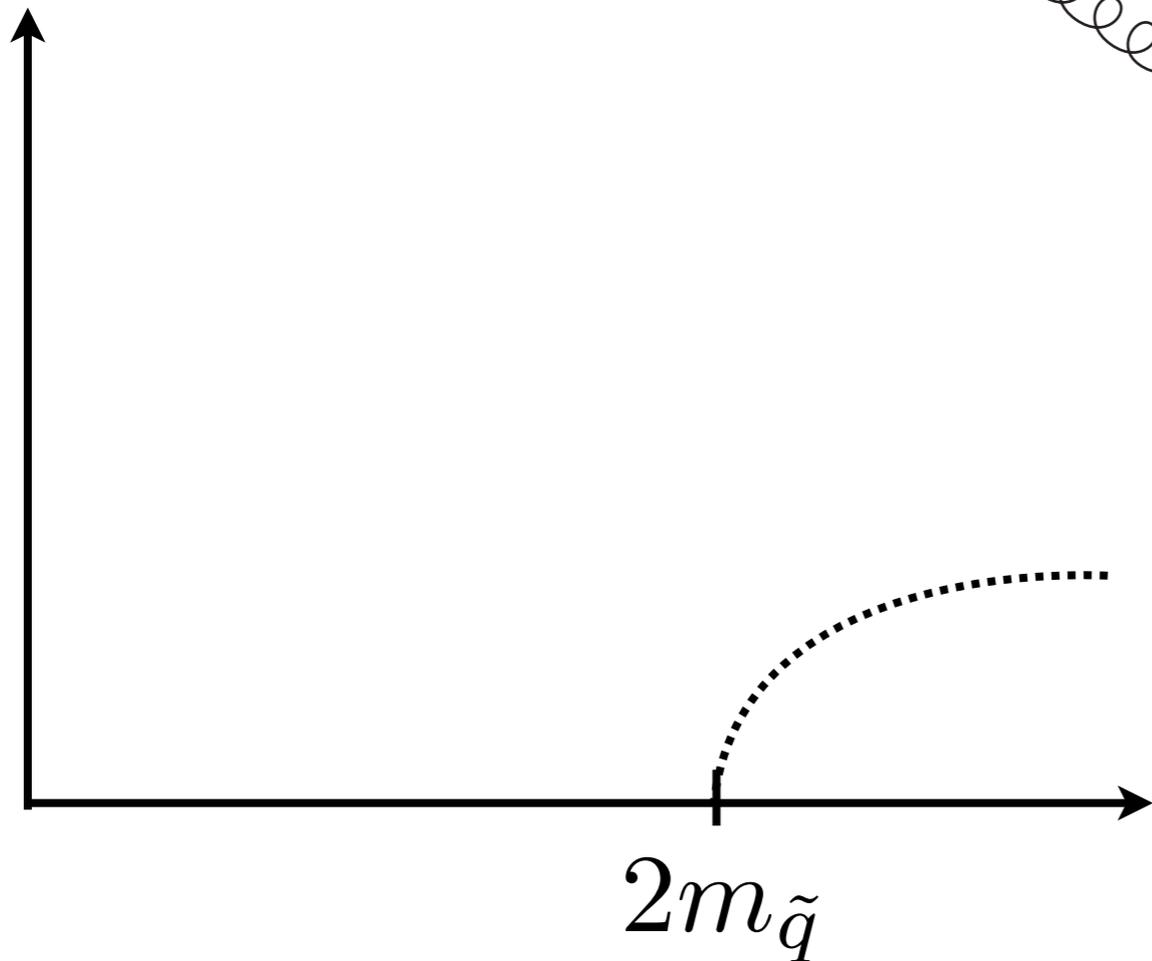
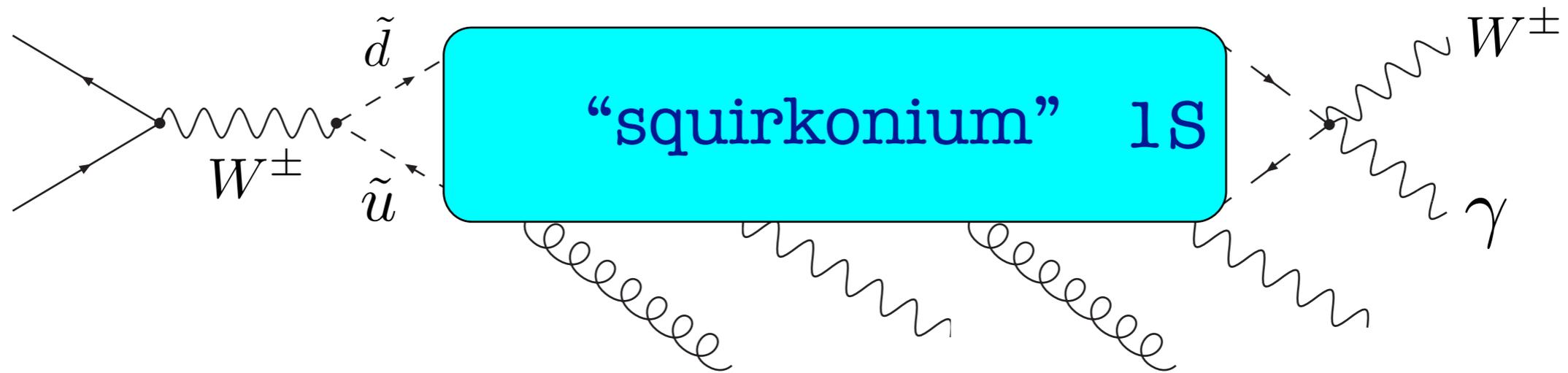
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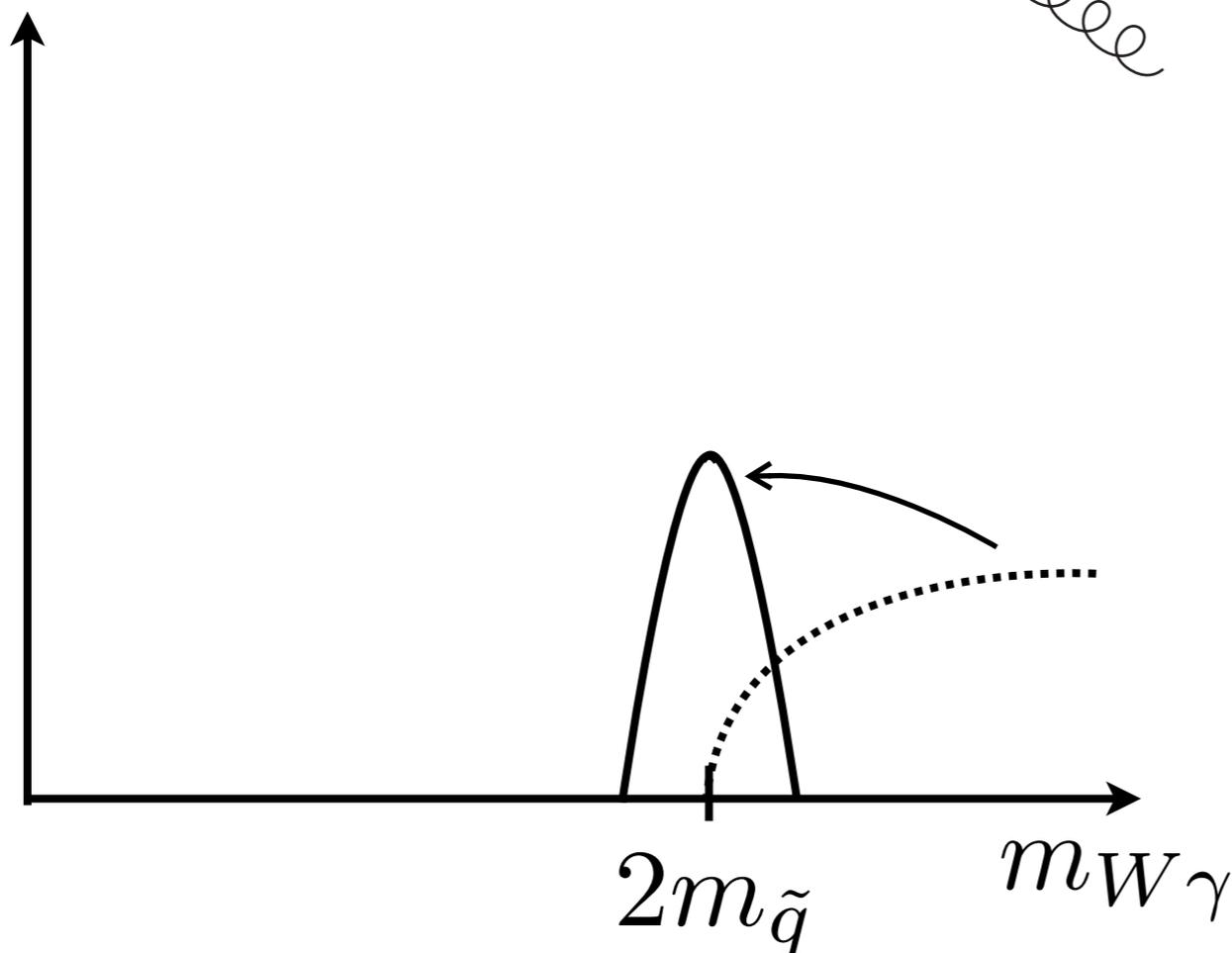
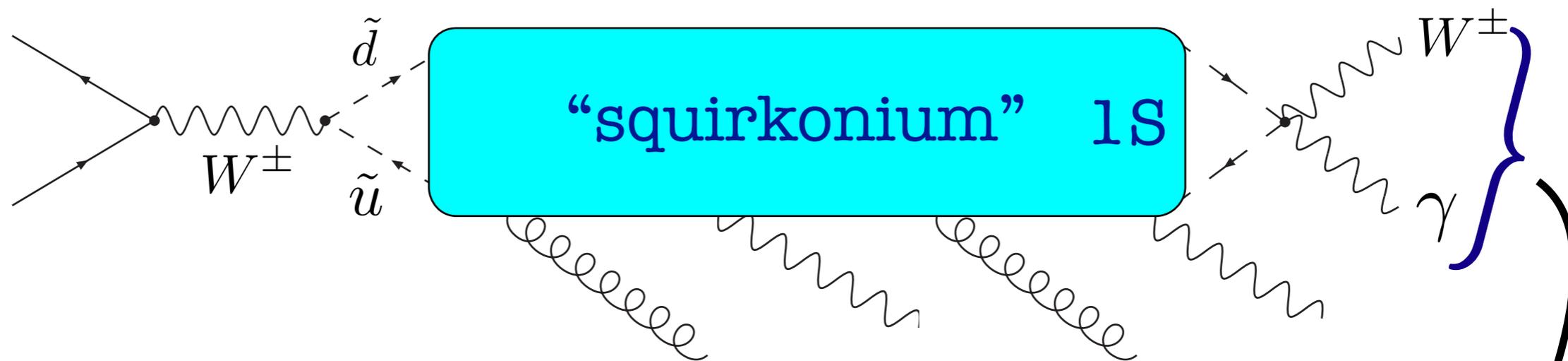
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\* Annihilation occurs after radiation:

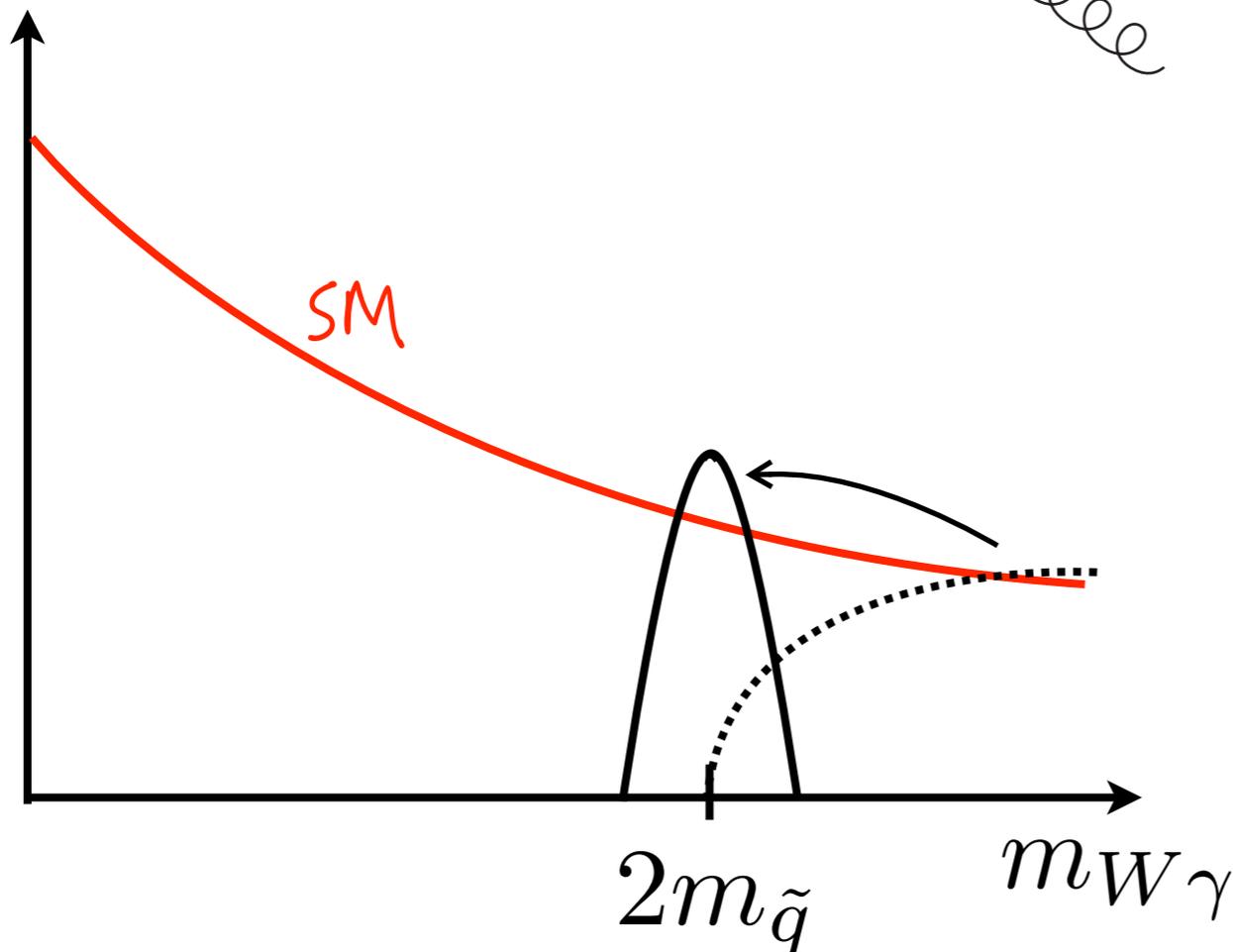
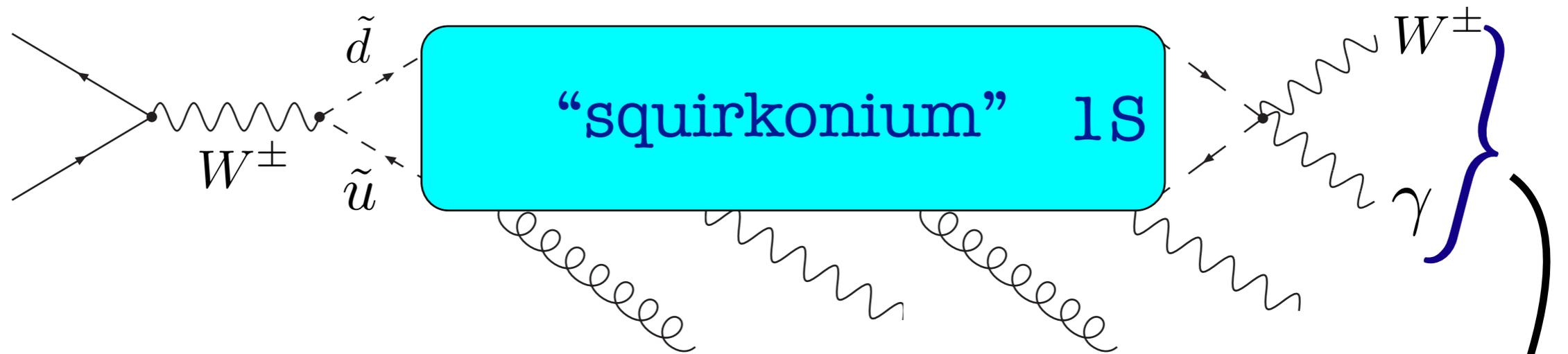


A peak in the invariant mass of  $W+\gamma$

$$m_{W\gamma}^2 = m_{1S}^2 \sim 4m_{\tilde{q}}^2$$

# A Hard Signal

\* Annihilation occurs after radiation:

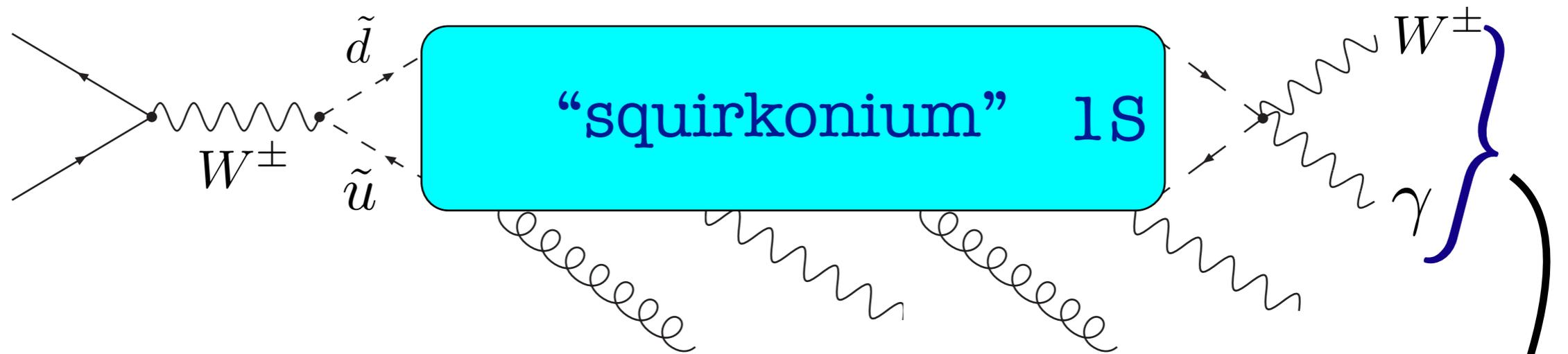


A peak in the invariant mass of  $W+\gamma$

$$m_{W\gamma}^2 = m_{1S}^2 \sim 4m_{\tilde{q}}^2$$

# A Hard Signal

\* Annihilation occurs after radiation:



For ( $\Lambda_{\text{QCD}} \sim$  few GeV) all of this story happens promptly on collider timescales.

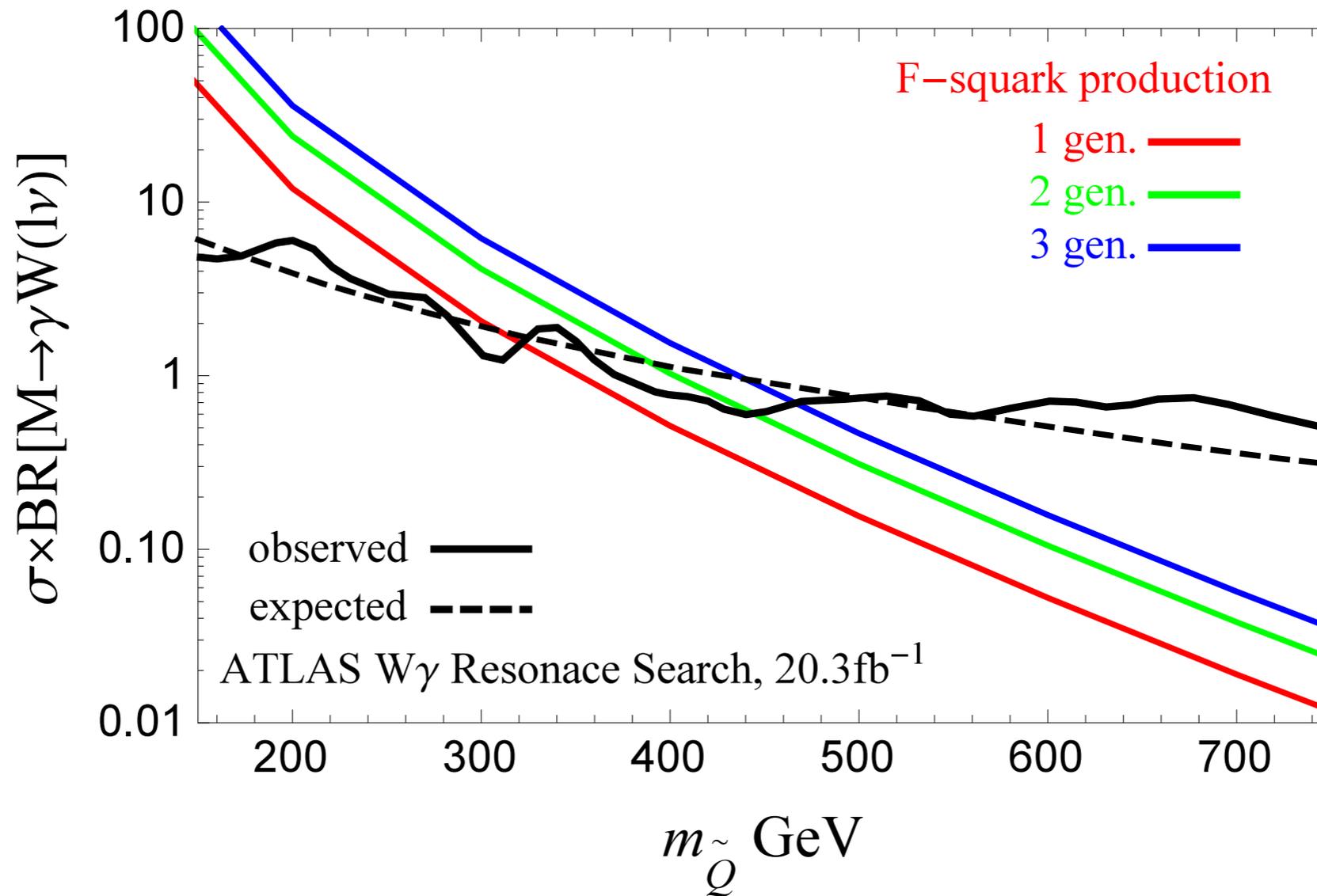
A peak in the invariant mass of  $W+\gamma$

$$m_{W\gamma}^2 = m_{1S}^2 \sim 4m_{\tilde{q}}^2$$

$2m_{\tilde{q}}$   $m_{W\gamma}$

# Current limit

(1411.3310)



\*Assumptions:  
1. Soft stuff does not smear the peak much.  
2. Squirk  $\beta$ -decay is slow.

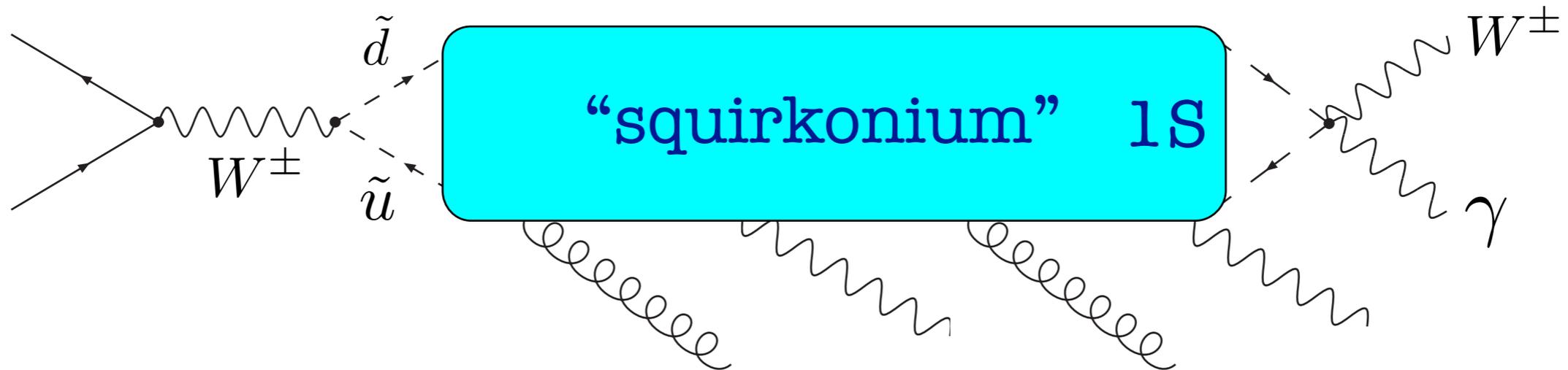
violating these assumptions only weakens the bound.

Using ATLAS  $W\gamma$  analysis.\*

Squirks at 300-450 GeV excluded. Natural models still allowed.

# Soft Stuff

- \* Having found the hard stuff, we can look for the soft stuff.



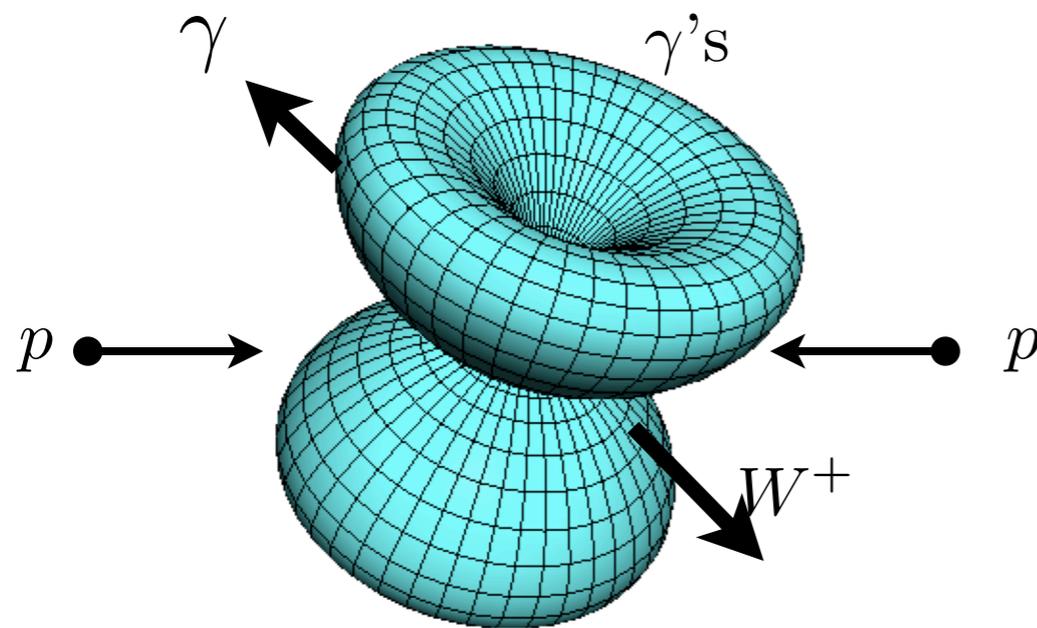
An overly active underlying event...?

(RH and Wizansky 0810.3949)

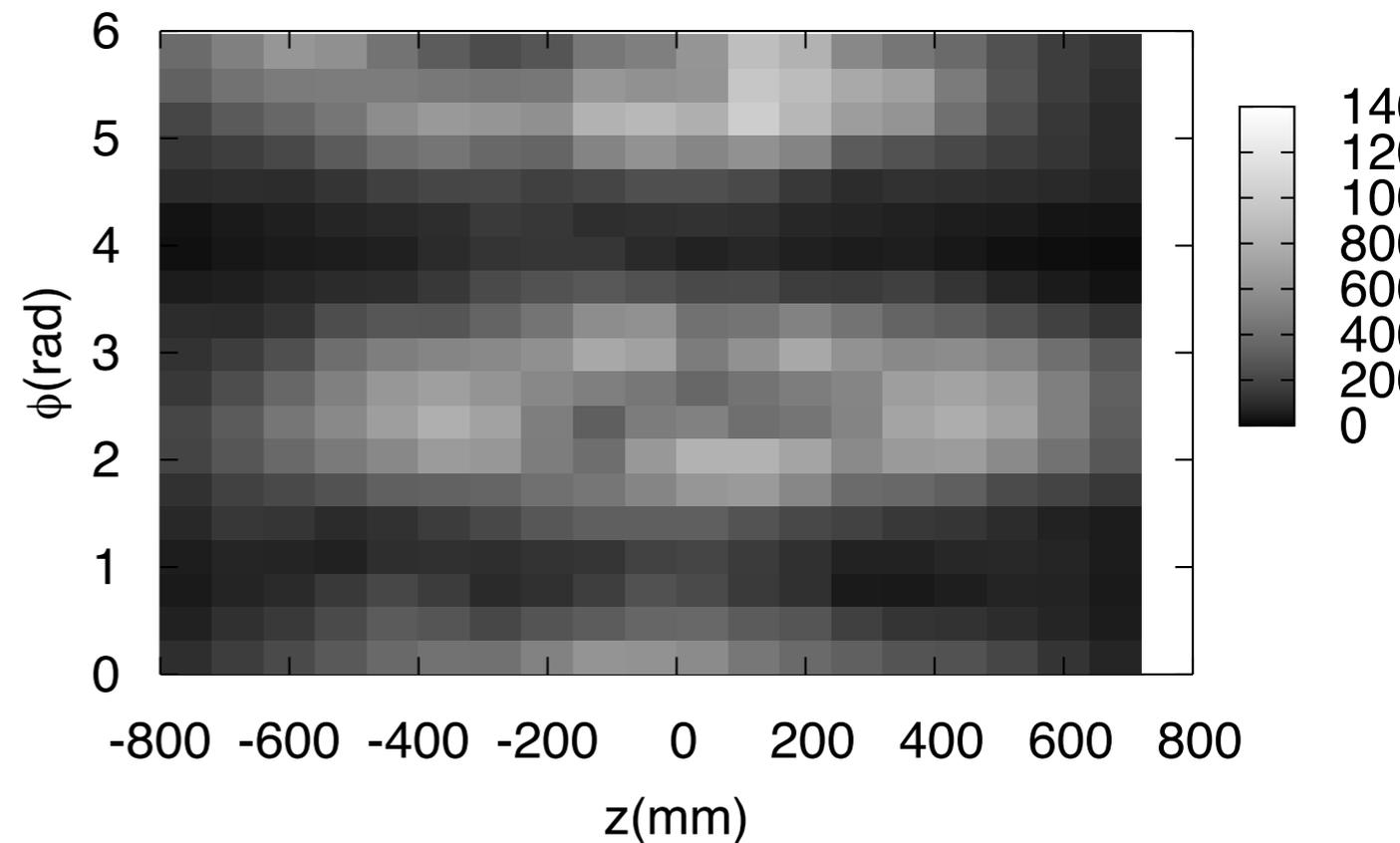
How about Glueball decay?

(see talk by Curtin)

# Photon Radiation



Photons are emitted in a pattern!  
(unless they are dominated by glueballs)



A weird underlying event.