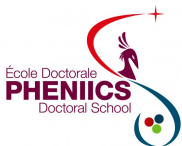


# Generalized Parton Distributions and their covariant extension

Nabil Chouika

Irfu/SPhN, CEA Saclay - Université Paris-Saclay

Pheniics Fest, LAL, Orsay, 30 mai 2017



# Outline

## 1 Introduction to Generalized Parton Distributions

- Definition and properties
- Experimental access

## 2 Modeling Generalized Parton Distributions

- Overlap of Light-cone wave functions
- Double Distributions
- Inversion of Incomplete Radon Transform
- Results

## 3 Conclusion

# Outline

## 1 Introduction to Generalized Parton Distributions

- Definition and properties
- Experimental access

## 2 Modeling Generalized Parton Distributions

- Overlap of Light-cone wave functions
- Double Distributions
- Inversion of Incomplete Radon Transform
- Results

## 3 Conclusion

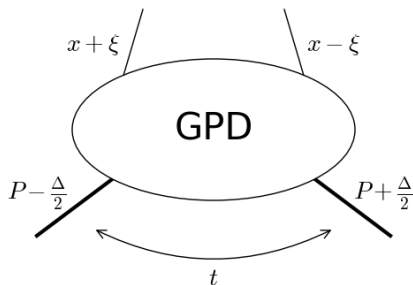
# Definition of GPDs

- Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{i x P^+ z^-} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^+ q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^+=0, z_\perp=0} \quad (1)$$

with:

$$t = \Delta^2, \quad \xi = -\frac{\Delta^+}{2P^+}.$$



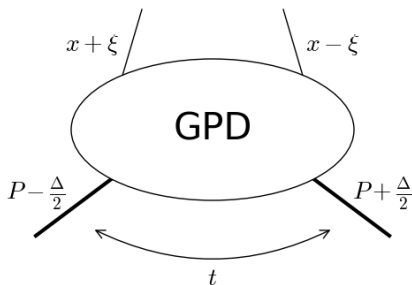
# Definition of GPDs

- Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{i x P^+ z^-} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^+ q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^+=0, z_\perp=0} \quad (1)$$

with:

$$t = \Delta^2, \quad \xi = -\frac{\Delta^+}{2P^+}.$$



- Similar matrix element for gluons.

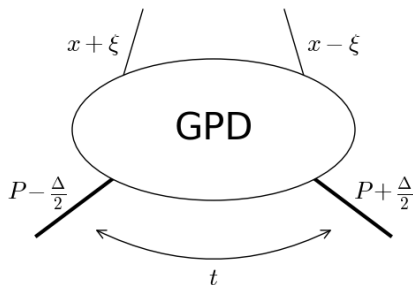
# Definition of GPDs

- Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{i x P^+ z^-} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^+ q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^+=0, z_\perp=0} \quad (1)$$

with:

$$t = \Delta^2, \quad \xi = -\frac{\Delta^+}{2P^+}.$$



- Similar matrix element for gluons.
- More GPDs for spin- $\frac{1}{2}$  hadrons.

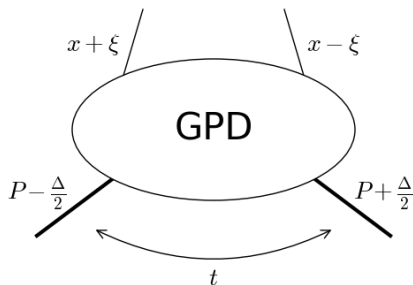
# Definition of GPDs

- Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{i x P^+ z^-} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^+ q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^+=0, z_\perp=0} \quad (1)$$

with:

$$t = \Delta^2, \quad \xi = -\frac{\Delta^+}{2P^+}.$$



- Similar matrix element for gluons.
- More GPDs for spin- $\frac{1}{2}$  hadrons.
- Experimental programs at JLab, COMPASS.

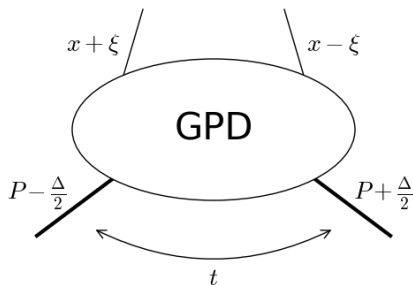
# Definition of GPDs

- Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^+ q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^+=0, z_\perp=0} \quad (1)$$

with:

$$t = \Delta^2, \quad \xi = -\frac{\Delta^+}{2P^+}.$$



- Similar matrix element for gluons.
- More GPDs for spin- $\frac{1}{2}$  hadrons.
- Experimental programs at JLab, COMPASS.

- Impact parameter space GPD (at  $\xi = 0$ ): (Burkardt, 2000)

$$q(x, b_\perp) = \int \frac{d^2 \vec{\Delta}_\perp}{(2\pi)^2} e^{-i \vec{b}_\perp \cdot \vec{\Delta}_\perp} H^q(x, 0, -\vec{\Delta}_\perp^2). \quad (2)$$



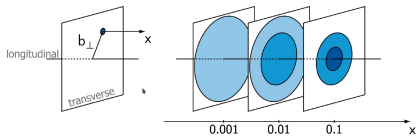
# Definition of GPDs

- Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^+ q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^+=0, z_\perp=0} \quad (1)$$

with:

$$t = \Delta^2, \quad \xi = -\frac{\Delta^+}{2P^+}.$$



- Similar matrix element for gluons.
- More GPDs for spin- $\frac{1}{2}$  hadrons.
- Experimental programs at JLab, COMPASS.

- Impact parameter space GPD (at  $\xi = 0$ ): (Burkardt, 2000)

$$q(x, \vec{b}_\perp) = \int \frac{d^2\vec{\Delta}_\perp}{(2\pi)^2} e^{-i\vec{b}_\perp \cdot \vec{\Delta}_\perp} H^q(x, 0, -\Delta_\perp^2). \quad (2)$$

# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .

# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .
  - ▶ DGLAP:  $|x| > |\xi|$ .

# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .
  - ▶ DGLAP:  $|x| > |\xi|$ .
  - ▶ ERBL:  $|x| < |\xi|$ .

# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .
  - ▶ DGLAP:  $|x| > |\xi|$ .
  - ▶ ERBL:  $|x| < |\xi|$ .
- Link to PDFs and Form Factors:

$$\int dx H^q(x, \xi, t) = F^q(t), \quad (3)$$

$$H^q(x, 0, 0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x). \quad (4)$$

# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .
  - ▶ DGLAP:  $|x| > |\xi|$ .
  - ▶ ERBL:  $|x| < |\xi|$ .
- Link to PDFs and Form Factors:

$$\int dx H^q(x, \xi, t) = F^q(t), \quad (3)$$

$$H^q(x, 0, 0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x). \quad (4)$$

- Polynomiality:

$$\int_{-1}^1 dx x^m H(x, \xi, t) = \text{Polynomial in } \xi. \quad (5)$$

# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .
  - ▶ DGLAP:  $|x| > |\xi|$ .
  - ▶ ERBL:  $|x| < |\xi|$ .
- Link to PDFs and Form Factors:

$$\int dx H^q(x, \xi, t) = F^q(t), \quad (3)$$

$$H^q(x, 0, 0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x). \quad (4)$$

- Polynomiality:

$$\int_{-1}^1 dx x^m H(x, \xi, t) = \text{Polynomial in } \xi. \quad (5)$$

- ▶ From Lorentz invariance.

# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .
  - ▶ DGLAP:  $|x| > |\xi|$ .
  - ▶ ERBL:  $|x| < |\xi|$ .
- Link to PDFs and Form Factors:

$$\int dx H^q(x, \xi, t) = F^q(t), \quad (3)$$

$$H^q(x, 0, 0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x). \quad (4)$$

- Polynomiality:

$$\int_{-1}^1 dx x^m H(x, \xi, t) = \text{Polynomial in } \xi. \quad (5)$$

- ▶ From Lorentz invariance.

- Positivity (in DGLAP): ([Pire et al., 1999](#); [Radyushkin, 1999](#))

$$|H^q(x, \xi, t)| \leq \sqrt{q\left(\frac{x-\xi}{1-\xi}\right) q\left(\frac{x+\xi}{1+\xi}\right)}. \quad (6)$$



# Theoretical constraints on GPDs

Main properties:

- Physical region:  $(x, \xi) \in [-1, 1]^2$ .
  - ▶ DGLAP:  $|x| > |\xi|$ .
  - ▶ ERBL:  $|x| < |\xi|$ .
- Link to PDFs and Form Factors:

$$\int dx H^q(x, \xi, t) = F^q(t), \quad (3)$$

$$H^q(x, 0, 0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x). \quad (4)$$

- Polynomiality:

$$\int_{-1}^1 dx x^m H(x, \xi, t) = \text{Polynomial in } \xi. \quad (5)$$

- ▶ From Lorentz invariance.

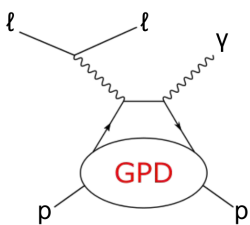
- Positivity (in DGLAP): ([Pire et al., 1999](#); [Radyushkin, 1999](#))

$$|H^q(x, \xi, t)| \leq \sqrt{q\left(\frac{x-\xi}{1-\xi}\right) q\left(\frac{x+\xi}{1+\xi}\right)}. \quad (6)$$

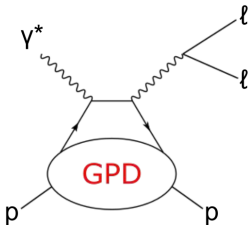
- ▶ Cauchy-Schwarz theorem in Hilbert space.

# Accessing GPDs

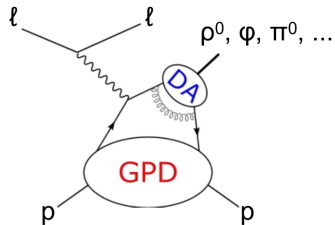
- Exclusive processes:



DVCS



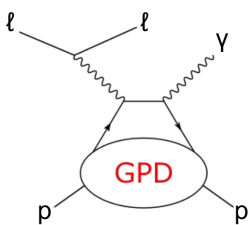
TCS



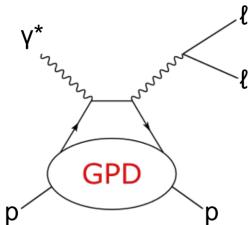
DVMP

# Accessing GPDs

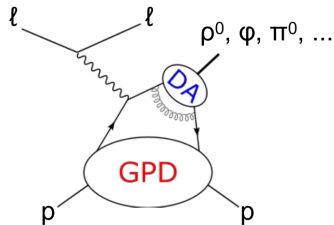
- Exclusive processes:



DVCS



TCS



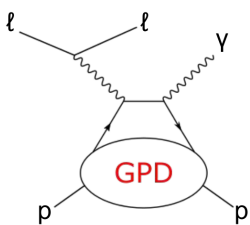
DVMP

- Compton Form Factors: [\(Belitsky et al., 2002\)](#)

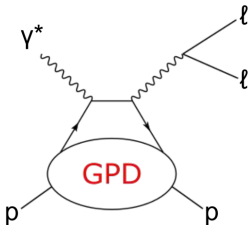
$$\mathcal{F}(\xi, t, Q^2) = \int_{-1}^1 dx C\left(x, \xi, \alpha_S(\mu_F), \frac{Q}{\mu_F}\right) F(x, \xi, t, \mu_F). \quad (7)$$

# Accessing GPDs

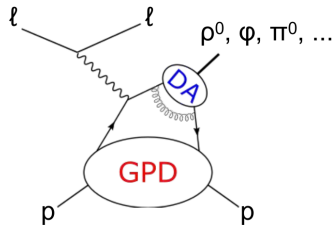
- Exclusive processes:



DVCS



TCS



DVMP

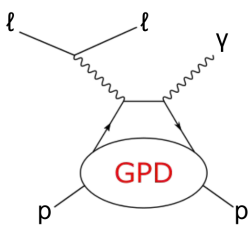
- Compton Form Factors: (Belitsky et al., 2002)

$$\mathcal{F}(\xi, t, Q^2) = \int_{-1}^1 dx C\left(x, \xi, \alpha_S(\mu_F), \frac{Q}{\mu_F}\right) F(x, \xi, t, \mu_F). \quad (7)$$

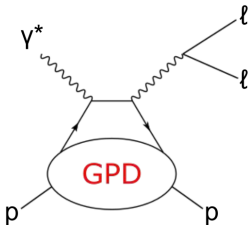
- Observables are convolutions of:

# Accessing GPDs

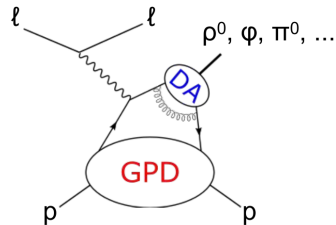
- Exclusive processes:



DVCS



TCS



DVMP

- Compton Form Factors: (Belitsky et al., 2002)

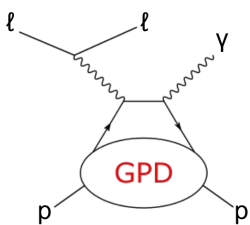
$$\mathcal{F}(\xi, t, Q^2) = \int_{-1}^1 dx C\left(x, \xi, \alpha_S(\mu_F), \frac{Q}{\mu_F}\right) F(x, \xi, t, \mu_F). \quad (8)$$

- Observables are convolutions of:

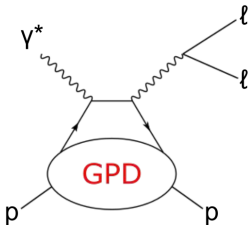
- ▶ a hard-scattering kernel, calculated with perturbative QCD (short distance interactions).

# Accessing GPDs

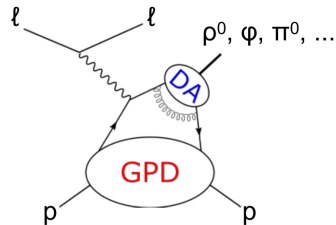
- Exclusive processes:



DVCS



TCS



DVMP

- Compton Form Factors: (Belitsky et al., 2002)

$$\mathcal{F}(\xi, t, Q^2) = \int_{-1}^1 dx C\left(x, \xi, \alpha_S(\mu_F), \frac{Q}{\mu_F}\right) F(x, \xi, t, \mu_F). \quad (9)$$

- Observables are convolutions of:

- ▶ a hard-scattering kernel, calculated with perturbative QCD (short distance interactions).
- ▶ a soft part, i.e. the GPD, with long distance interactions encoded (non-perturbative QCD).

# Outline

## 1 Introduction to Generalized Parton Distributions

- Definition and properties
- Experimental access

## 2 Modeling Generalized Parton Distributions

- Overlap of Light-cone wave functions
- Double Distributions
- Inversion of Incomplete Radon Transform
- Results

## 3 Conclusion

# Overlap of Light-cone wave functions

- A given *hadronic state* is decomposed in a **Fock basis**: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N, \beta} \int [dx]_N [d^2\mathbf{k}_\perp]_N \Psi_{N, \beta}^\lambda(x_1, \mathbf{k}_{\perp 1}, \dots) |N, \beta; k_1, \dots, k_N\rangle, \quad (10)$$

where the  $\Psi_{N, \beta}^\lambda$  are the *Light-cone wave-functions (LCWF)*.



# Overlap of Light-cone wave functions

- A given *hadronic state* is decomposed in a **Fock basis**: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N, \beta} \int [dx]_N [d^2\mathbf{k}_\perp]_N \Psi_{N, \beta}^\lambda(x_1, \mathbf{k}_{\perp 1}, \dots) |N, \beta; k_1, \dots, k_N\rangle, \quad (10)$$

where the  $\Psi_{N, \beta}^\lambda$  are the *Light-cone wave-functions (LCWF)*.

- For example, for the pion:

$$|\pi^+\rangle = \psi_{u\bar{d}}^\pi |u\bar{d}\rangle + \psi_{u\bar{d}g}^\pi |u\bar{d}g\rangle + \dots \quad (11)$$

# Overlap of Light-cone wave functions

- A given *hadronic state* is decomposed in a **Fock basis**: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [dx]_N [d^2\mathbf{k}_\perp]_N \Psi_{N,\beta}^\lambda(x_1, \mathbf{k}_{\perp 1}, \dots) |N, \beta; k_1, \dots, k_N\rangle, \quad (10)$$

where the  $\Psi_{N,\beta}^\lambda$  are the *Light-cone wave-functions (LCWF)*.

- For example, for the pion:

$$|\pi^+\rangle = \psi_{u\bar{d}}^\pi |u\bar{d}\rangle + \psi_{u\bar{d}g}^\pi |u\bar{d}g\rangle + \dots \quad (11)$$

- GPD as an overlap of LCWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^q(x, \xi, t) = \sum_{N,\beta} \sqrt{1-\xi}^{2-N} \sqrt{1+\xi}^{2-N} \sum_a \delta_{a,q} \quad (12)$$

$$\times \int [d\bar{x}]_N [d^2\bar{\mathbf{k}}_\perp]_N \delta(x - \bar{x}_a) \Psi_{N,\beta}^*(\hat{x}'_1, \hat{\mathbf{k}}'_{\perp 1}, \dots) \Psi_{N,\beta}(\tilde{x}_1, \tilde{\mathbf{k}}_{\perp 1}, \dots),$$

in the DGLAP region  $\xi < x < 1$  (pion case).

# Overlap of Light-cone wave functions

- A given *hadronic state* is decomposed in a **Fock basis**: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [dx]_N [d^2\mathbf{k}_\perp]_N \Psi_{N,\beta}^\lambda(x_1, \mathbf{k}_{\perp 1}, \dots) |N, \beta; k_1, \dots, k_N\rangle, \quad (10)$$

where the  $\Psi_{N,\beta}^\lambda$  are the *Light-cone wave-functions (LCWF)*.

- For example, for the pion:

$$|\pi^+\rangle = \psi_{u\bar{d}}^\pi |u\bar{d}\rangle + \psi_{u\bar{d}g}^\pi |u\bar{d}g\rangle + \dots \quad (11)$$

- GPD as an overlap of LCWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^q(x, \xi, t) = \sum_{N,\beta} \sqrt{1-\xi}^{2-N} \sqrt{1+\xi}^{2-N} \sum_a \delta_{a,q} \quad (12)$$

$$\times \int [d\bar{x}]_N [d^2\bar{\mathbf{k}}_\perp]_N \delta(x - \bar{x}_a) \Psi_{N,\beta}^*(\hat{x}'_1, \hat{\mathbf{k}}'_{\perp 1}, \dots) \Psi_{N,\beta}(\tilde{x}_1, \tilde{\mathbf{k}}_{\perp 1}, \dots),$$

in the DGLAP region  $\xi < x < 1$  (pion case).

- Similar result in ERBL ( $-\xi < x < \xi$ ), but with  $N$  and  $N + 2\dots$

# Overlap of Light-cone wave functions

- A given *hadronic state* is decomposed in a **Fock basis**: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [dx]_N [d^2\mathbf{k}_\perp]_N \Psi_{N,\beta}^\lambda(x_1, \mathbf{k}_{\perp 1}, \dots) |N, \beta; k_1, \dots, k_N\rangle, \quad (10)$$

where the  $\Psi_{N,\beta}^\lambda$  are the *Light-cone wave-functions (LCWF)*.

- For example, for the pion:

$$|\pi^+\rangle = \psi_{u\bar{d}}^\pi |u\bar{d}\rangle + \psi_{u\bar{d}g}^\pi |u\bar{d}g\rangle + \dots \quad (11)$$

- GPD as an overlap of LCWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^q(x, \xi, t) = \sum_{N,\beta} \sqrt{1-\xi}^{2-N} \sqrt{1+\xi}^{2-N} \sum_a \delta_{a,q} \quad (12)$$

$$\times \int [d\bar{x}]_N [d^2\bar{\mathbf{k}}_\perp]_N \delta(x - \bar{x}_a) \Psi_{N,\beta}^*(\hat{x}'_1, \hat{\mathbf{k}}'_{\perp 1}, \dots) \Psi_{N,\beta}(\tilde{x}_1, \tilde{\mathbf{k}}_{\perp 1}, \dots),$$

in the DGLAP region  $\xi < x < 1$  (pion case).

- Similar result in ERBL ( $-\xi < x < \xi$ ), but with  $N$  and  $N + 2 \dots$
- GPD is a scalar product of LCWFs:

# Overlap of Light-cone wave functions

- A given *hadronic state* is decomposed in a **Fock basis**: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N, \beta} \int [dx]_N [d^2\mathbf{k}_\perp]_N \Psi_{N, \beta}^\lambda(x_1, \mathbf{k}_{\perp 1}, \dots) |N, \beta; k_1, \dots, k_N\rangle, \quad (10)$$

where the  $\Psi_{N, \beta}^\lambda$  are the *Light-cone wave-functions* (LCWF).

- For example, for the pion:

$$|\pi^+\rangle = \psi_{u\bar{d}}^\pi |u\bar{d}\rangle + \psi_{u\bar{d}g}^\pi |u\bar{d}g\rangle + \dots \quad (11)$$

- GPD as an overlap of LCWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^q(x, \xi, t) = \sum_{N, \beta} \sqrt{1-\xi}^{2-N} \sqrt{1+\xi}^{2-N} \sum_a \delta_{a,q} \times \int [d\bar{x}]_N [d^2\bar{\mathbf{k}}_\perp]_N \delta(x - \bar{x}_a) \Psi_{N, \beta}^*(\hat{x}'_1, \hat{\mathbf{k}}'_{\perp 1}, \dots) \Psi_{N, \beta}(\tilde{x}_1, \tilde{\mathbf{k}}_{\perp 1}, \dots), \quad (12)$$

in the DGLAP region  $\xi < x < 1$  (pion case).

- Similar result in ERBL ( $-\xi < x < \xi$ ), but with  $N$  and  $N + 2 \dots$
- GPD is a scalar product of LCWFs:
  - Cauchy-Schwarz theorem  $\Rightarrow$  **Positivity** fulfilled!

# Overlap of Light-cone wave functions

- A given *hadronic state* is decomposed in a **Fock basis**: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [dx]_N [d^2\mathbf{k}_\perp]_N \Psi_{N,\beta}^\lambda(x_1, \mathbf{k}_{\perp 1}, \dots) |N, \beta; k_1, \dots, k_N\rangle, \quad (10)$$

where the  $\Psi_{N,\beta}^\lambda$  are the *Light-cone wave-functions* (LCWF).

- For example, for the pion:

$$|\pi^+\rangle = \psi_{u\bar{d}}^\pi |u\bar{d}\rangle + \psi_{u\bar{d}g}^\pi |u\bar{d}g\rangle + \dots \quad (11)$$

- GPD as an overlap of LCWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^q(x, \xi, t) = \sum_{N,\beta} \sqrt{1-\xi}^{2-N} \sqrt{1+\xi}^{2-N} \sum_a \delta_{a,q} \quad (12)$$

$$\times \int [d\bar{x}]_N [d^2\bar{\mathbf{k}}_\perp]_N \delta(x - \bar{x}_a) \Psi_{N,\beta}^*(\hat{x}'_1, \hat{\mathbf{k}}'_{\perp 1}, \dots) \Psi_{N,\beta}(\tilde{x}_1, \tilde{\mathbf{k}}_{\perp 1}, \dots),$$

in the DGLAP region  $\xi < x < 1$  (pion case).

- Similar result in ERBL ( $-\xi < x < \xi$ ), but with  $N$  and  $N + 2$ ...
- GPD is a scalar product of LCWFs:
  - ▶ Cauchy-Schwarz theorem  $\Rightarrow$  **Positivity** fulfilled!
  - ▶ Polynomiality not manifest...

# Double Distributions (DDs)

- DD representation of GPDs:

$$H(x, \xi, t) \propto \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi). \quad (13)$$

# Double Distributions (DDs)

- DD representation of GPDs:

$$H(x, \xi, t) \propto \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi). \quad (13)$$

- DD  $f$  is defined on the support  $\Omega = \{|\beta| + |\alpha| \leq 1\}$ .



# Double Distributions (DDs)

- DD representation of GPDs:

$$H(x, \xi, t) \propto \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi). \quad (13)$$

- DD  $f$  is defined on the support  $\Omega = \{|\beta| + |\alpha| \leq 1\}$ .
- **Polynomial** in  $\xi$ :

$$\begin{aligned} \int_{-1}^1 dx x^m H(x, \xi, t) &\propto \int dx x^m \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi) \\ &\propto \int_{\Omega} d\beta d\alpha (\beta + \xi\alpha)^m f(\beta, \alpha, t). \end{aligned} \quad (14)$$

# Double Distributions (DDs)

- DD representation of GPDs:

$$H(x, \xi, t) \propto \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi). \quad (13)$$

- DD  $f$  is defined on the support  $\Omega = \{|\beta| + |\alpha| \leq 1\}$ .
- **Polynomial** in  $\xi$ :

$$\begin{aligned} \int_{-1}^1 dx x^m H(x, \xi, t) &\propto \int dx x^m \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi) \\ &\propto \int_{\Omega} d\beta d\alpha (\beta + \xi\alpha)^m f(\beta, \alpha, t). \end{aligned} \quad (14)$$

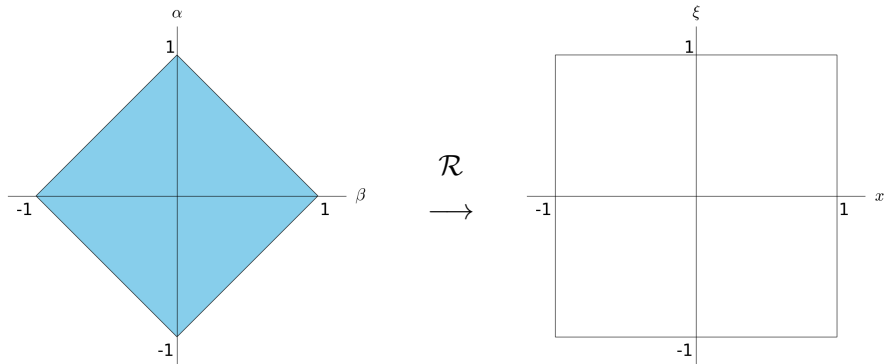
- Positivity not manifest...

# Double Distributions (DDs)

- DD representation of GPDs:

$$H(x, \xi, t) \propto \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi). \quad (13)$$

- Radon Transform: (Radon, 1986; Deans, 1983; Teryaev, 2001)

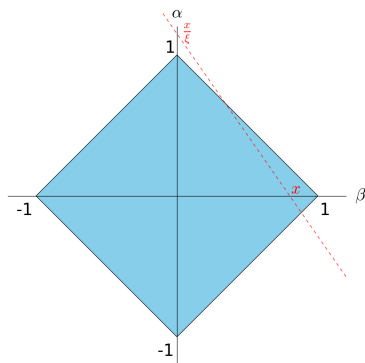
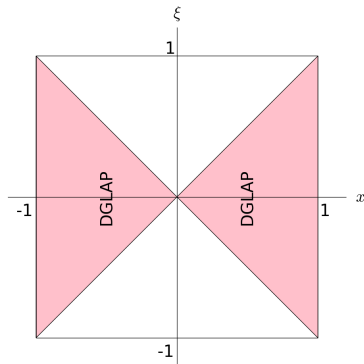


# Double Distributions (DDs)

- DD representation of GPDs:

$$H(x, \xi, t) \propto \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi). \quad (13)$$

- Radon Transform: (Deans, 1983; Teryaev, 2001)

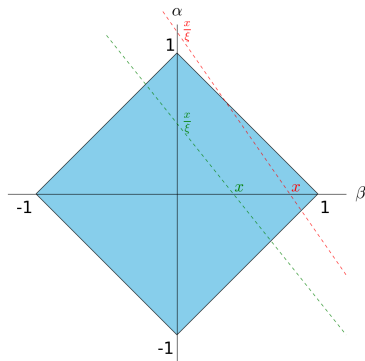
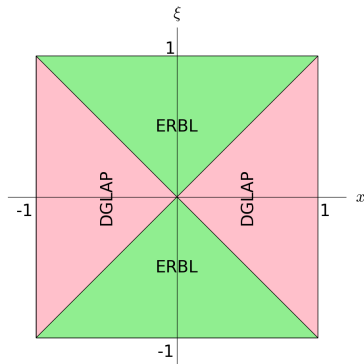

 $\mathcal{R}$   
 $\longrightarrow$ 


# Double Distributions (DDs)

- DD representation of GPDs:

$$H(x, \xi, t) \propto \int_{\Omega} d\beta d\alpha f(\beta, \alpha, t) \delta(x - \beta - \alpha\xi). \quad (13)$$

- Radon Transform: (Deans, 1983; Teryaev, 2001)


 $\mathcal{R}$   
 $\longrightarrow$ 


# From DGLAP GPD to a DD

- In Overlap representation: DGLAP region only (e.g. two-body LCWFs).

# From DGLAP GPD to a DD

- In Overlap representation: DGLAP region only (e.g. two-body LCWFs).
  - ▶ Need ERBL to complete **polynomiality**.

# From DGLAP GPD to a DD

- In Overlap representation: DGLAP region only (e.g. two-body LCWFs).
  - ▶ Need ERBL to complete **polynomiality**.

## Problem

Find  $f(\beta, \alpha)$  on square  $\{|\alpha| + |\beta| \leq 1\}$  such that

$$H(x, \xi)|_{\text{DGLAP}} \propto \int d\beta d\alpha f(\beta, \alpha) \delta(x - \beta - \alpha\xi) .$$



# From DGLAP GPD to a DD

- In Overlap representation: DGLAP region only (e.g. two-body LCWFs).
  - ▶ Need ERBL to complete **polynomiality**.

## Problem

Find  $f(\beta, \alpha)$  on square  $\{|\alpha| + |\beta| \leq 1\}$  such that

$$H(x, \xi)|_{\text{DGLAP}} \propto \int d\beta d\alpha f(\beta, \alpha) \delta(x - \beta - \alpha\xi) .$$

- If model fulfills Lorentz invariance: ([Moutarde, 2015](#))

# From DGLAP GPD to a DD

- In Overlap representation: DGLAP region only (e.g. two-body LCWFs).
  - ▶ Need ERBL to complete **polynomiality**.

## Problem

Find  $f(\beta, \alpha)$  on square  $\{|\alpha| + |\beta| \leq 1\}$  such that

$$H(x, \xi)|_{\text{DGLAP}} \propto \int d\beta d\alpha f(\beta, \alpha) \delta(x - \beta - \alpha\xi) .$$

- If model fulfills Lorentz invariance: ([Moutarde, 2015](#))
  - ▶ DD  $f(\beta, \alpha)$  **exists** (if the GPD behaves well) and is **unique**.

# From DGLAP GPD to a DD

- In Overlap representation: DGLAP region only (e.g. two-body LCWFs).
  - ▶ Need ERBL to complete **polynomiality**.

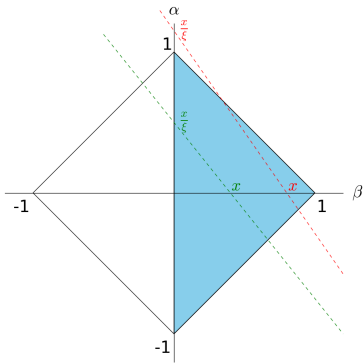
## Problem

Find  $f(\beta, \alpha)$  on square  $\{|\alpha| + |\beta| \leq 1\}$  such that

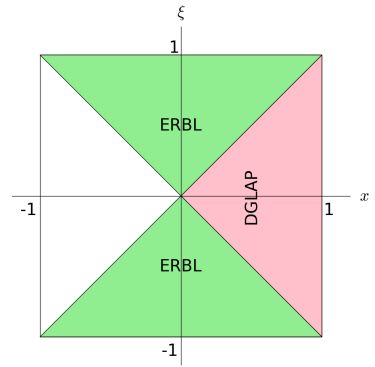
$$H(x, \xi)|_{\text{DGLAP}} \propto \int d\beta d\alpha f(\beta, \alpha) \delta(x - \beta - \alpha\xi).$$

- If model fulfills Lorentz invariance: ([Moutarde, 2015](#))
  - ▶ DD  $f(\beta, \alpha)$  **exists** (if the GPD behaves well) and is **unique**.
  - ▶ We can reconstruct the GPD everywhere.

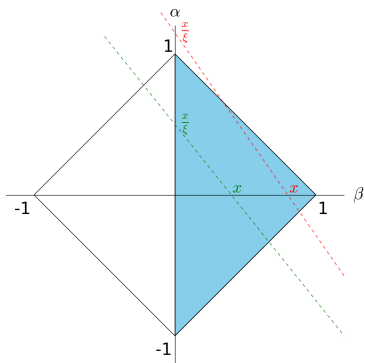
# Numerical Inversion



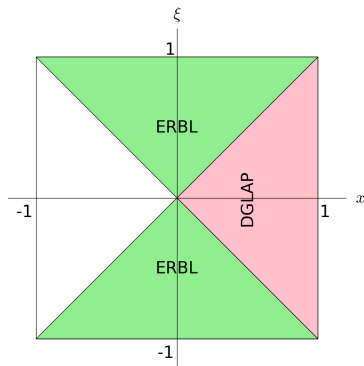
$\mathcal{R}$   
 $\rightarrow$   
 $\leftarrow$   
 ?



# Numerical Inversion

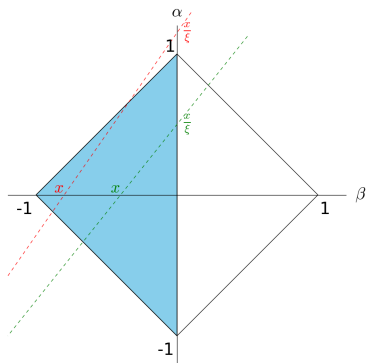


$\mathcal{R}$   
 $\longrightarrow$   
 $\longleftarrow$   
 ?

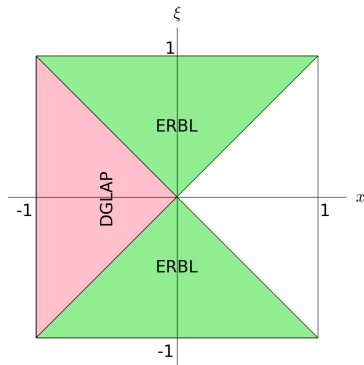


- Quark GPD:  $H(x, \xi) = 0$  for  $-1 < x < -|\xi| \implies f(\beta, \alpha) = 0$  for  $\beta < 0$ .

# Numerical Inversion

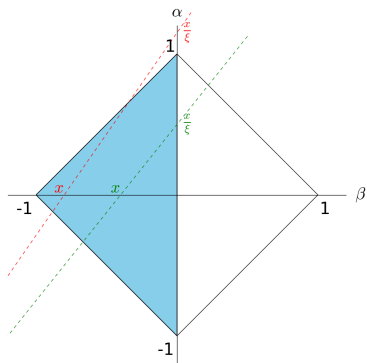


$\mathcal{R}$   
 $\rightarrow$   
 $\leftarrow$   
 ?

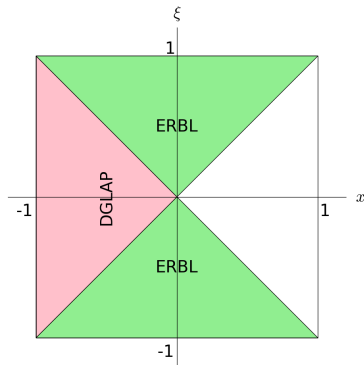


- Quark GPD:  $H(x, \xi) = 0$  for  $-1 < x < -|\xi| \implies f(\beta, \alpha) = 0$  for  $\beta < 0$ .

# Numerical Inversion

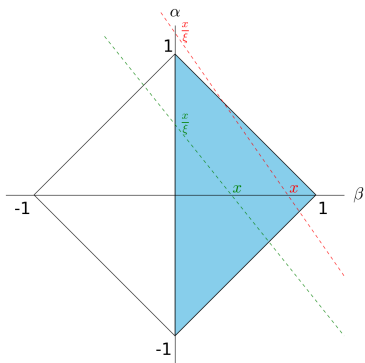


$\mathcal{R}$   
 $\rightarrow$   
 $\leftarrow$   
 ?

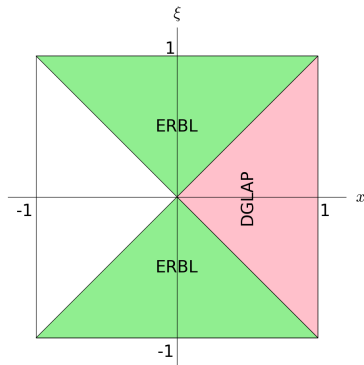


- Quark GPD:  $H(x, \xi) = 0$  for  $-1 < x < -|\xi| \implies f(\beta, \alpha) = 0$  for  $\beta < 0$ .
- Domains  $\beta < 0$  and  $\beta > 0$  are uncorrelated in the DGLAP region.

# Numerical Inversion



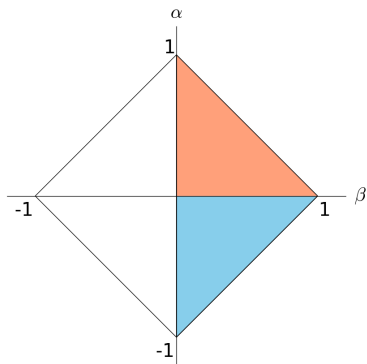
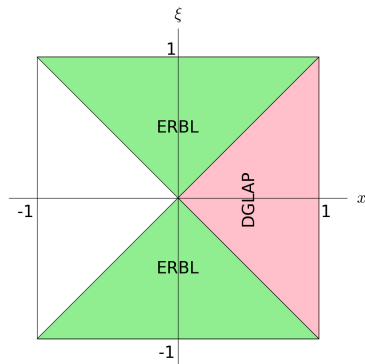
$\mathcal{R}$   
 $\longrightarrow$   
 $\longleftarrow$   
 ?



- Quark GPD:  $H(x, \xi) = 0$  for  $-1 < x < -|\xi| \implies f(\beta, \alpha) = 0$  for  $\beta < 0$ .
- Domains  $\beta < 0$  and  $\beta > 0$  are uncorrelated in the DGLAP region.
- Divide and conquer:
  - ▶ Better numerical stability.
  - ▶ Lesser complexity:  $O(N^p + N^p) \ll O((N + N)^p)$ .



# Numerical Inversion


 $\mathcal{R}$ 
 $\rightarrow$ 
 $\leftarrow$ 
 $?$ 


- Quark GPD:  $H(x, \xi) = 0$  for  $-1 < x < -|\xi| \implies f(\beta, \alpha) = 0$  for  $\beta < 0$ .
- Domains  $\beta < 0$  and  $\beta > 0$  are uncorrelated in the DGLAP region.
- Divide and conquer:
  - ▶ Better numerical stability.
  - ▶ Lesser complexity:  $O(N^p + N^p) \ll O((N + N)^p)$ .
- $\alpha$ -parity of the DD:  $f(\beta, -\alpha) = f(\beta, \alpha)$ .

# Ill-posed problems and Regularization

- The inversion of a Fredholm equation of the first kind

$$\int K(x, y) f(y) dy = g(x) \quad (15)$$

is an ill-posed problem.

# Ill-posed problems and Regularization

- The inversion of a Fredholm equation of the first kind

$$\int K(x, y) f(y) dy = g(x) \quad (15)$$

is an ill-posed problem.

- ▶ The inverse is not continuous: an arbitrarily small variation  $\Delta g$  of the rhs can lead to an arbitrarily large variation  $\Delta f$  of the solution.

# Ill-posed problems and Regularization

- The inversion of a Fredholm equation of the first kind

$$\int K(x, y) f(y) dy = g(x) \quad (15)$$

is an ill-posed problem.

- ▶ The inverse is not continuous: an arbitrarily small variation  $\Delta g$  of the rhs can lead to an arbitrarily large variation  $\Delta f$  of the solution.
- The discrete problem needs to be regularized.

# Ill-posed problems and Regularization

- The inversion of a Fredholm equation of the first kind

$$\int K(x, y) f(y) dy = g(x) \quad (15)$$

is an ill-posed problem.

- ▶ The inverse is not continuous: an arbitrarily small variation  $\Delta g$  of the rhs can lead to an arbitrarily large variation  $\Delta f$  of the solution.
- The discrete problem needs to be regularized.
  - ▶ E.g Tikhonov regularization:  $\min \left\{ \|AX - B\|^2 + \epsilon \|X\|^2 \right\}$ .

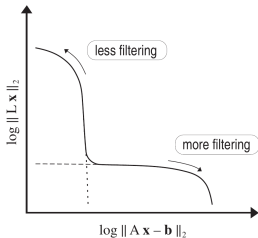
# Ill-posed problems and Regularization

- The inversion of a Fredholm equation of the first kind

$$\int K(x, y) f(y) dy = g(x) \quad (15)$$

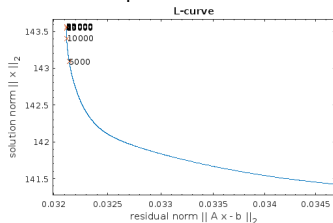
is an ill-posed problem.

- The inverse is not continuous: an arbitrarily small variation  $\Delta g$  of the rhs can lead to an arbitrarily large variation  $\Delta f$  of the solution.
- The discrete problem needs to be regularized.
  - E.g Tikhonov regularization:  $\min \left\{ \|AX - B\|^2 + \epsilon \|X\|^2 \right\}$ .



Theoretical "L-curve": curve parameterized by the regularization factor.

(fig. taken from Ref. [\(Hansen, 2007\)](#))



L-curve with the iteration number as regularization factor.

# Some examples (Dyson-Schwinger model)

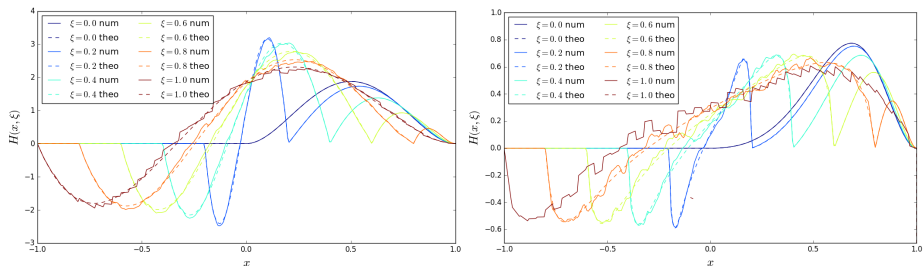


Figure: Extension of GPDs for the pion DSE model of Refs. (Mezrag, 2015; Mezrag et al., 2016). Left:  $t = 0 \text{ GeV}^2$ . Right:  $t = 1 \text{ GeV}^2$ . Comparison to the analytical result.

# Some examples (Spectator model)

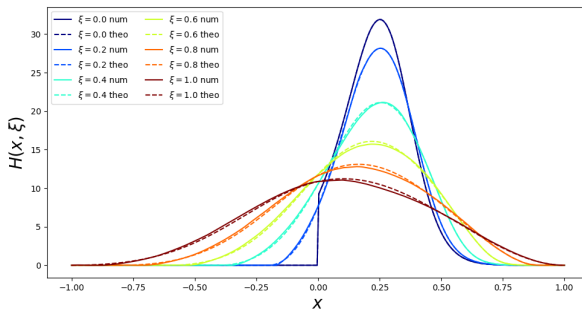


Figure: Extension of GPD E for the nucleon model of Ref. (Hwang and Mueller, 2008). Comparison to the analytical result of the paper.



# Summary

- Generalized Parton Distributions

# Summary

- Generalized Parton Distributions
  - ▶ encode information about the 3D structure of a hadron.

# Summary

- Generalized Parton Distributions
  - ▶ encode information about the 3D structure of a hadron.
  - ▶ are accessible with exclusive processes in experiments: JLab, COMPASS, etc.

# Summary

- Generalized Parton Distributions
  - ▶ encode information about the 3D structure of a hadron.
  - ▶ are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:

# Summary

- Generalized Parton Distributions
  - ▶ encode information about the 3D structure of a hadron.
  - ▶ are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
  - ▶ LCWFs  $\xrightarrow{\text{Overlap}}$  GPD in DGLAP  $\xrightarrow{\text{Inverse Radon Transform}}$  DD  $\xrightarrow{\text{RT}}$  GPD.

# Summary

- Generalized Parton Distributions
  - ▶ encode information about the 3D structure of a hadron.
  - ▶ are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
  - ▶ LCWFs  $\xrightarrow{\text{Overlap}}$  GPD in DGLAP  $\xrightarrow{\text{Inverse Radon Transform}}$  DD  $\xrightarrow{\text{RT}}$  GPD.
  - ▶ Both polynomiality and positivity!

# Summary

- Generalized Parton Distributions

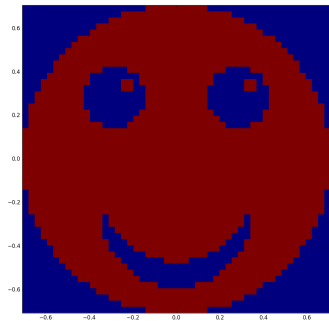
- ▶ encode information about the 3D structure of a hadron.
- ▶ are accessible with exclusive processes in experiments: JLab, COMPASS, etc.

- Systematic procedure for GPD modeling from first principles:

- ▶ LCWFs  $\xrightarrow{\text{Overlap}}$  GPD in DGLAP  $\xrightarrow{\text{Inverse Radon Transform}}$  DD  $\xrightarrow{\text{RT}}$  GPD.
- ▶ Both polynomiality and positivity!
- ▶ Compromise with respect to noise and convergence.

# Summary

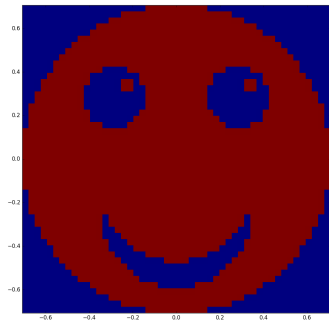
- Generalized Parton Distributions
  - ▶ encode information about the 3D structure of a hadron.
  - ▶ are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
  - ▶ LCWFs  $\xrightarrow{\text{Overlap}}$  GPD in DGLAP  $\xrightarrow{\text{Inverse Radon Transform}}$  DD  $\xrightarrow{\text{RT}}$  GPD.
  - ▶ Both polynomiality and positivity!
  - ▶ Compromise with respect to noise and convergence.
- Thank you!





# Summary

- Generalized Parton Distributions
  - ▶ encode information about the 3D structure of a hadron.
  - ▶ are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
  - ▶ LCWFs  $\xrightarrow{\text{Overlap}}$  GPD in DGLAP  $\xrightarrow{\text{Inverse Radon Transform}}$  DD  $\xrightarrow{\text{RT}}$  GPD.
  - ▶ Both polynomiality and positivity!
  - ▶ Compromise with respect to noise and convergence.
- Thank you!
  - ▶ Any questions?



# Bibliography I

- D. Müller, D. Robaschik, B. Geyer, F. M. Dittes, and J. Hořejši, “Wave functions, evolution equations and evolution kernels from light ray operators of QCD”, *Fortsch. Phys.* **42** (1994) 101–141, arXiv:hep-ph/9812448 [hep-ph].
- A. V. Radyushkin, “Scaling limit of deeply virtual Compton scattering”, *Phys. Lett.* **B380** (1996) 417–425, arXiv:hep-ph/9604317 [hep-ph].
- X.-D. Ji, “Deeply virtual Compton scattering”, *Phys. Rev.* **D55** (1997) 7114–7125, arXiv:hep-ph/9609381 [hep-ph].
- M. Burkardt, “Impact parameter dependent parton distributions and off forward parton distributions for  $\zeta \rightarrow 0$ ”, *Phys. Rev.* **D62** (2000) 071503, arXiv:hep-ph/0005108 [hep-ph], [Erratum: *Phys. Rev.*D66,119903(2002)].
- B. Pire, J. Soffer, and O. Teryaev, “Positivity constraints for off - forward parton distributions”, *Eur. Phys. J.* **C8** (1999) 103–106, arXiv:hep-ph/9804284 [hep-ph].
- A. V. Radyushkin, “Double distributions and evolution equations”, *Phys. Rev.* **D59** (1999) 014030, arXiv:hep-ph/9805342 [hep-ph].
- A. V. Belitsky, D. Mueller, and A. Kirchner, “Theory of deeply virtual Compton scattering on the nucleon”, *Nucl. Phys.* **B629** (2002) 323–392, arXiv:hep-ph/0112108 [hep-ph].
- S. J. Brodsky, T. Huang, and G. P. Lepage, “Hadronic wave functions and high momentum transfer interactions in quantum chromodynamics”, *Conf. Proc.* **C810816** (1981) 143–199.

# Bibliography II

- M. Diehl, T. Feldmann, R. Jakob, and P. Kroll, "The overlap representation of skewed quark and gluon distributions", *Nucl. Phys.* **B596** (2001) 33–65, arXiv:hep-ph/0009255 [hep-ph], [Erratum: Nucl. Phys. B605,647(2001)].
- M. Diehl, "Generalized parton distributions", *Phys. Rept.* **388** (2003) 41–277, arXiv:hep-ph/0307382 [hep-ph].
- J. Radon, "On the determination of functions from their integral values along certain manifolds", *Medical Imaging, IEEE Transactions on* **5** Dec (1986) 170–176.
- S. R. Deans, "The Radon Transform and Some of Its Applications", Wiley-Interscience, 1983.
- O. V. Teryaev, "Crossing and radon tomography for generalized parton distributions", *Phys. Lett.* **B510** (2001) 125–132, arXiv:hep-ph/0102303 [hep-ph].
- H. Moutarde, "Nucleon Reverse Engineering: Structuring hadrons with colored degrees of freedom", 2015.
- P. C. Hansen, "Regularization Tools version 4.0 for Matlab 7.3", *Numerical Algorithms* **46** (2007), no. 2, 189–194.
- C. Mezrag, "Generalised Parton Distributions : from phenomenological approaches to Dyson-Schwinger equations", PhD thesis, IRFU, SPhN, Saclay, 2015.
- C. Mezrag, H. Moutarde, and J. Rodriguez-Quintero, "From Bethe–Salpeter Wave functions to Generalised Parton Distributions", *Few Body Syst.* **57** (2016), no. 9, 729–772, arXiv:1602.07722 [nucl-th].
- D. S. Hwang and D. Mueller, "Implication of the overlap representation for modelling generalized parton distributions", *Phys. Lett.* **B660** (2008) 350–359, arXiv:0710.1567 [hep-ph].