

Bottom Production vs. Theory: Past and Present

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A success story, a few glitches left and, along the way, some good lessons

What shall we mean by “success”?

Take massive Next-to-Leading Order perturbative QCD (+ NLL resummation, where needed) as a reference, and ask for its ability to:

- predict total rates for charm, bottom and top production
- describe differential distributions with the addition of a minimal, self-consistent, and possibly universal set of non-perturbative inputs

A successful comparison will be an agreement between possibly real measurements (i.e. little or no extrapolations/deconvolutions) and QCD predictions, within both experimental and theoretical uncertainties (ren./fact. scales, quark masses, strong coupling, PDFs and FFs,)

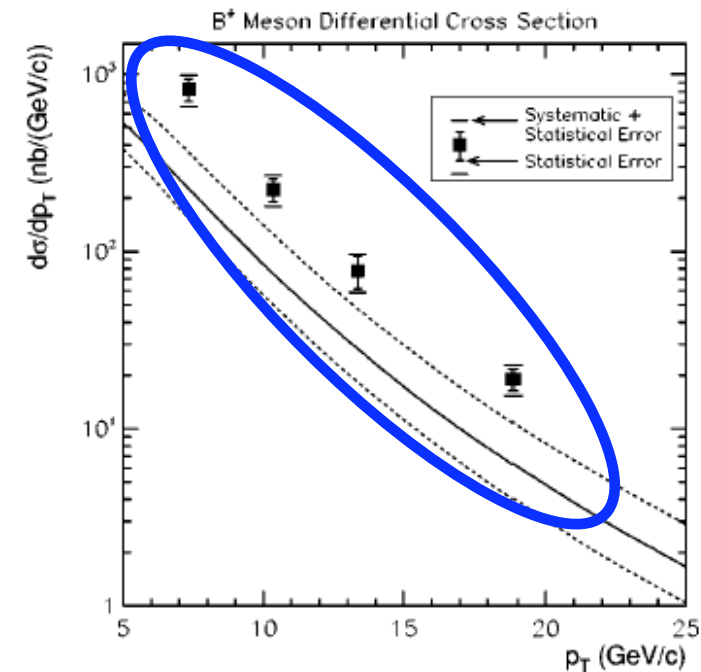
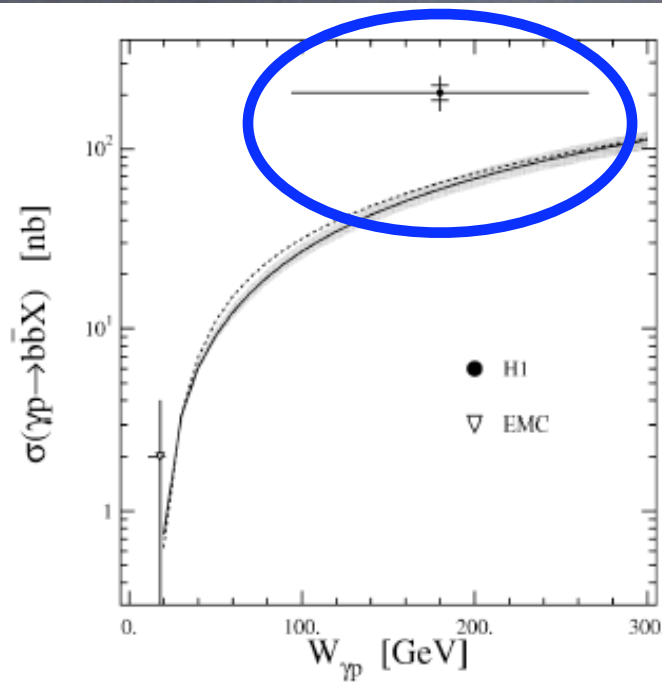
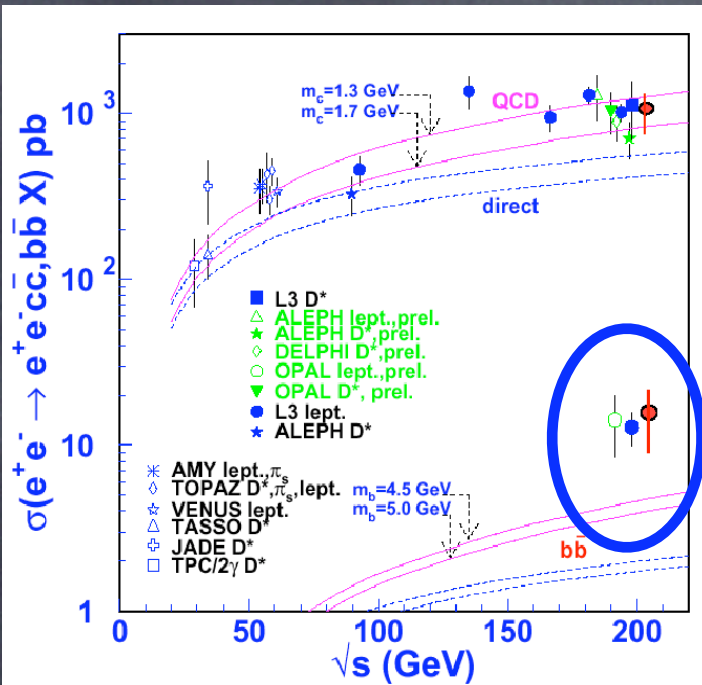
The Past

Bottom: What success?!?

$$\gamma\gamma \rightarrow b\bar{b}$$

$$\gamma p \rightarrow b\bar{b}$$

$$p\bar{p} \rightarrow BX$$



While for charm (large th. unc.) and for top (large expt. unc.) agreement was found, for bottom production discrepancies of 'a factor of three' or so were typically quoted in $\gamma\gamma$, γp and pp

Let's look at these comparisons in detail

NB: the hadroproduction part of this talk draws generously from a seminar that M.L. Mangano gave at Fermilab in January 2004. His full talk, with many more details, can be found at <http://cern.ch/~mlm/talks/Bcrosssection.pdf> and hep-ph/0411020. See also MC hep-ph/0407187

It's worth remembering that most of the perturbative QCD ingredients for heavy quark production have been available for some time now:

Hadroproduction

Nason, Dawson, Ellis, NP B327 (1989) 49, NP B303 (1988) 607

Beenakker, van Neerven, Meng, Schuler, Smith, NP B351 (1991) 507

Photoproduction

Nason, Ellis, NP B312 (1989) 551

Smith, van Neerven, NP B374 (1992) 36

$\Upsilon\Upsilon$

Drees, Kraemer, Zunft, Zerwas, PL B306 (1993) 371

Collinear resummation

Mele, Nason, NP B361 (1991) 626

MC, Greco, NP B421 (1994) 530

Threshold resummation

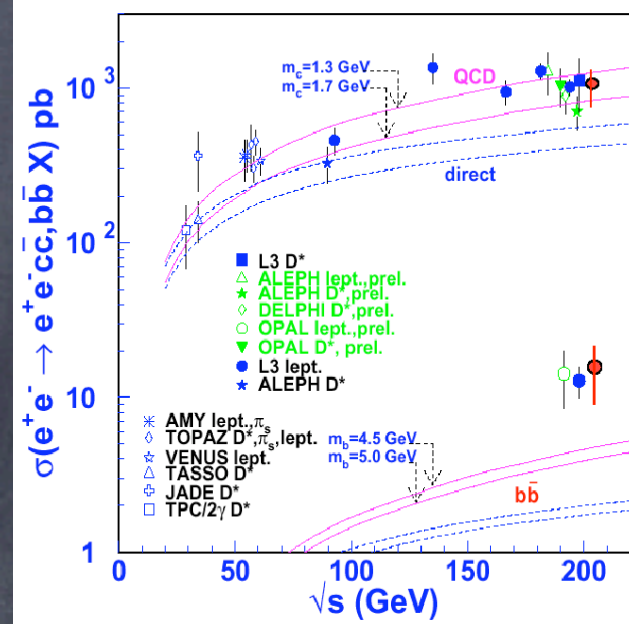
Bonciani, Catani, Mangano, Nason, NP B529 (1998) 424

Laenen, Smith, van Neerven, NP B369 (1992) 543

NLO
(massive)

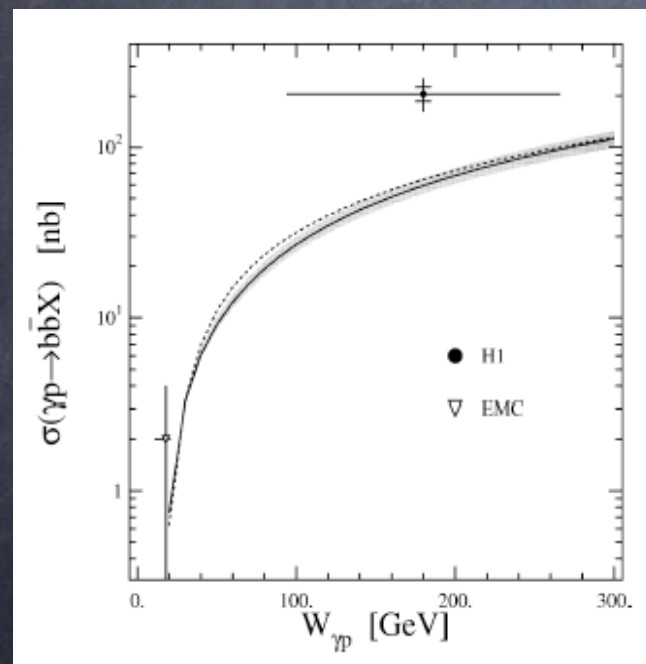
LL/NLL

+ surely many others. Apologies to those I forgot



Will not discuss in detail gamma-gamma and photoproduction results.

Main criticism: most of these results heavily depend on MonteCarlo's, both for extraction of signal and **extrapolation to full phase space** (usually by a large factor)



Hence, they do not fully qualify as real measurements, as they may be biased by theoretical prejudices

News: L3 recently published non-extrapolated measurements

New results from L3 within acceptance region (hep-ex/0507041):

$$\begin{aligned}\sigma(e^+e^- \rightarrow e^+e^-b\bar{b}X)_{\text{electrons}}^{\text{observed}} &= 0.41 \pm 0.08 \pm 0.08 \text{ pb} \\ \sigma(e^+e^- \rightarrow e^+e^-b\bar{b}X)_{\text{muons}}^{\text{observed}} &= 0.56 \pm 0.10 \pm 0.10 \text{ pb}\end{aligned}$$

The first uncertainties are statistical and the second systematic, and arise from the sources discussed below. These cross sections correspond to the phase space of the selected leptons, without any extrapolation: lepton momenta above 2 GeV and polar angles in the ranges $|\cos \theta| < 0.725$ for electrons and $|\cos \theta| < 0.8$ for muons, respectively.

Unfortunately, only compared so far with Monte Carlo results:

$$\begin{aligned}\sigma(e^+e^- \rightarrow e^+e^-b\bar{b}X)_{\text{electrons}}^{\text{CASCADE}} &= 0.11 \pm 0.02 \text{ pb} \\ \sigma(e^+e^- \rightarrow e^+e^-b\bar{b}X)_{\text{muons}}^{\text{CASCADE}} &= 0.14 \pm 0.02 \text{ pb}\end{aligned}$$

Homework: calculate accurate NLO predictions (and estimate th. uncertainties) for OBSERVED cross section

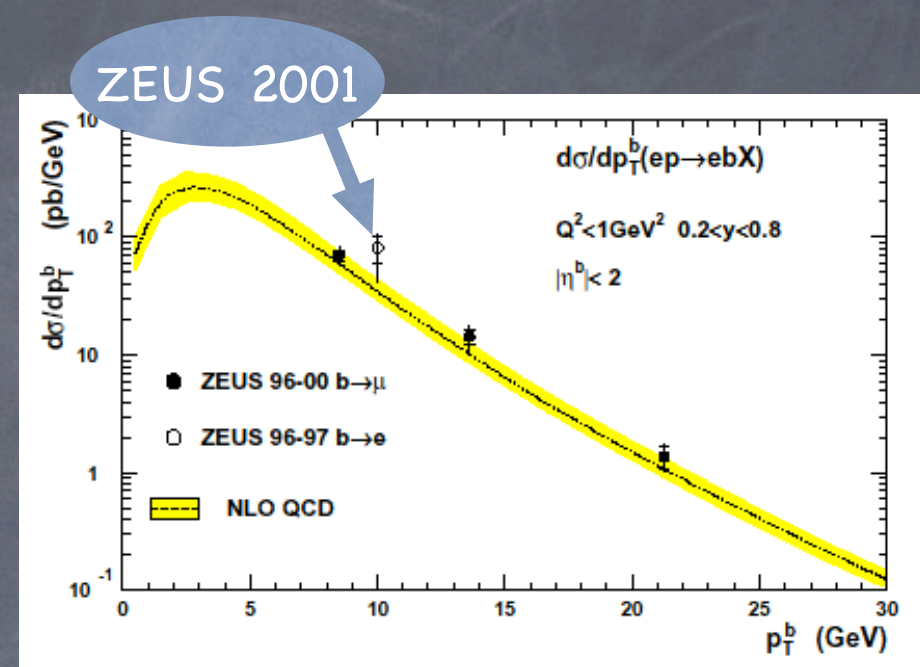
These uncertainties
certainly underestimated

Bottom line: I do not consider these discrepancies as particularly significant.
More recent and more exclusive measurements at HERA are in fair agreement

b Photoproduction

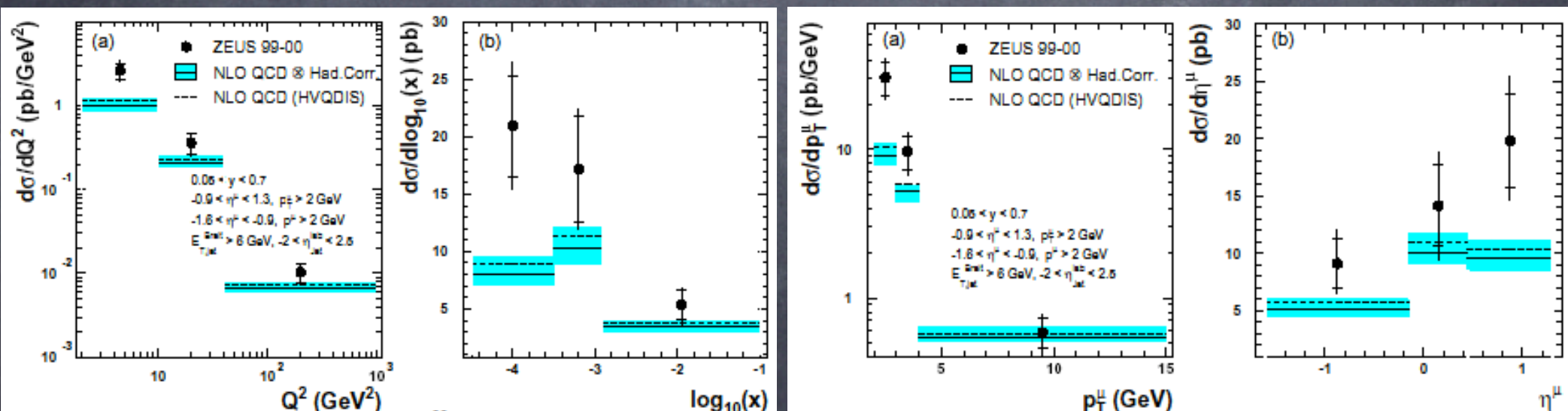
ZEUS 2001 [Eur. Phys. J. C18 (2001) 625]

ZEUS 2003 [hep-ex/0312057]

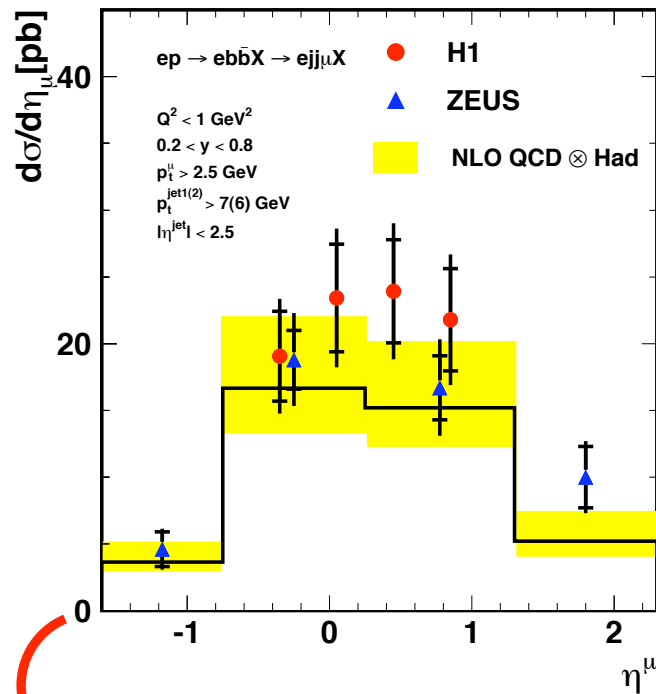


b production in DIS

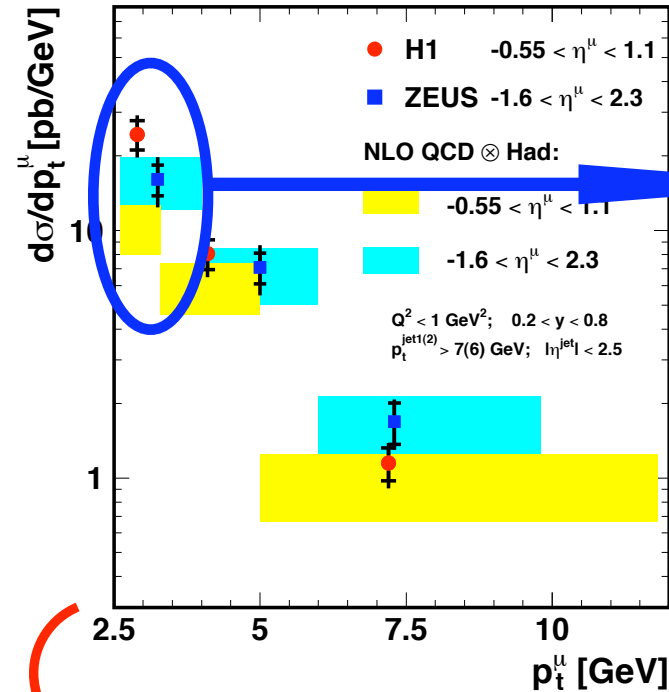
[ZEUS 2004, hep-ex/0405069]



More recent and more exclusive theory/data comparison in photoproduction



Good agreement
H1 vs ZEUS



ZEUS: No excess
at low p_t^μ

Single mildly
troubling point

Things at HERA now look fairly good. But this is '2005 wisdom'.
But what was going on in the past?

In fact, the 'hype of the b excess' around the year 2000 was built on years of measurements apparently supporting such findings. Let's review them.

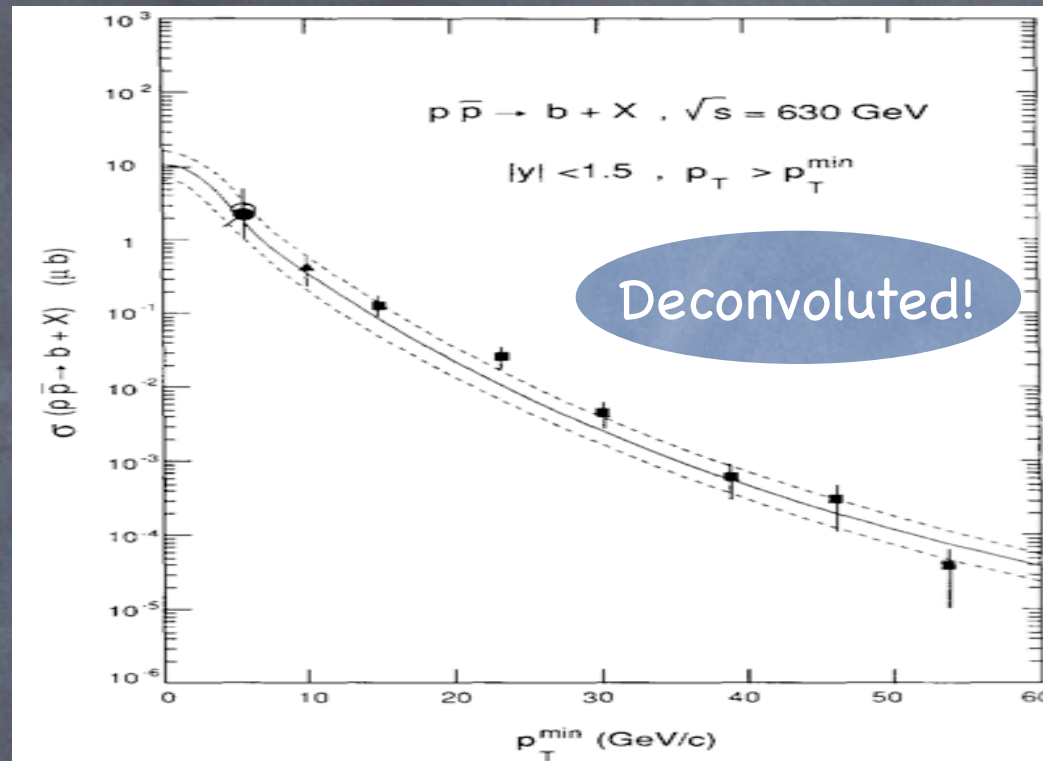
Bottom production in $p\bar{p}$ collisions

UA1 1988–1991

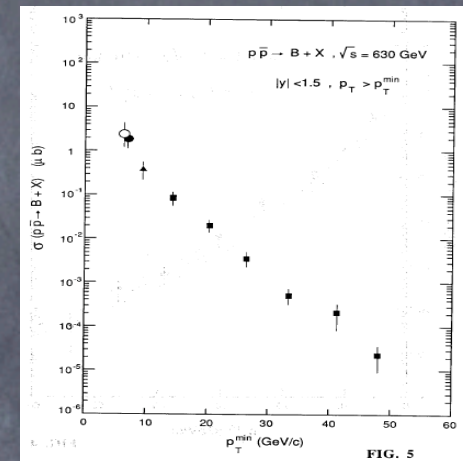
PL B213 (1988) 405

PL B256 (1991) 121

UA1/QCD ~ 1



NB. UA1 also published data for physical particles, B mesons and muons. At that time, they could however not easily be compared to theoretical predictions



CDF 1992

PRL 68 (1992) 3403

$$\sigma(p\bar{p} \rightarrow B^- X; P_T > 9.0 \text{ GeV}/c, |y| < 1.0)$$

$$= 2.8 \pm 0.9(\text{stat}) \pm 1.1(\text{syst}) \mu\text{b}$$

$$\sigma(pp \rightarrow bX; p_T > 11.5 \text{ GeV}, |y| < 1):$$

$$\text{CDF} = 6.1 \pm 1.9 \pm 2.4 \mu\text{b}$$

$$\text{theory} = 1.1 \pm 0.5 \mu\text{b}$$

tion. Our measurement is approximately 1.6 standard deviations above the theoretical calculation.

Deconvoluted!

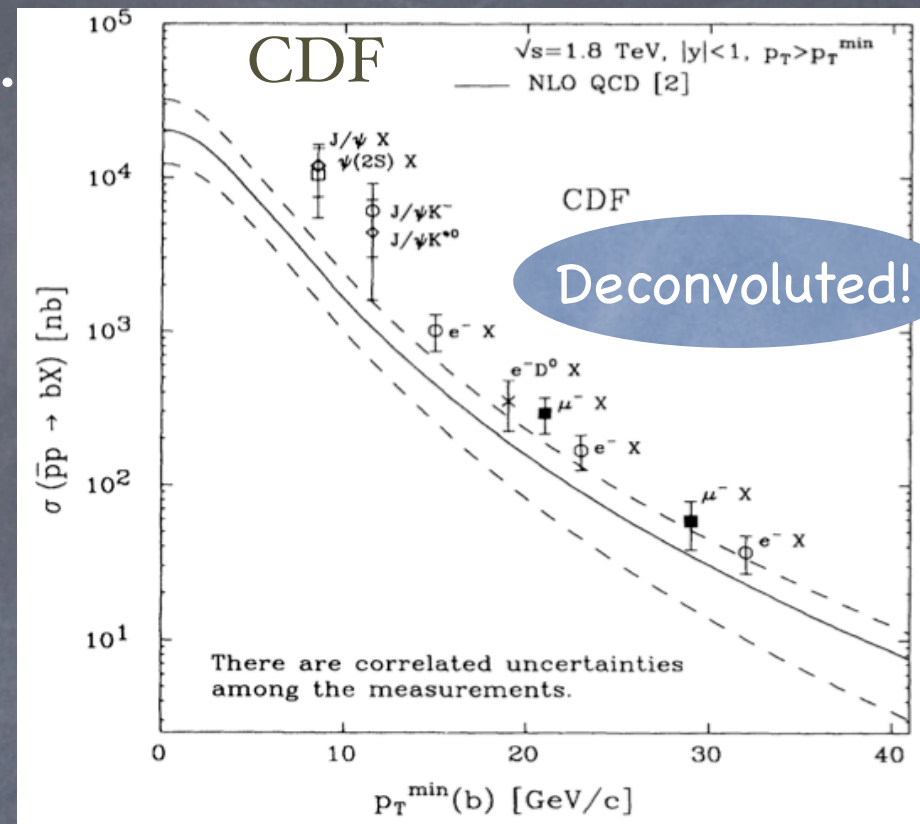
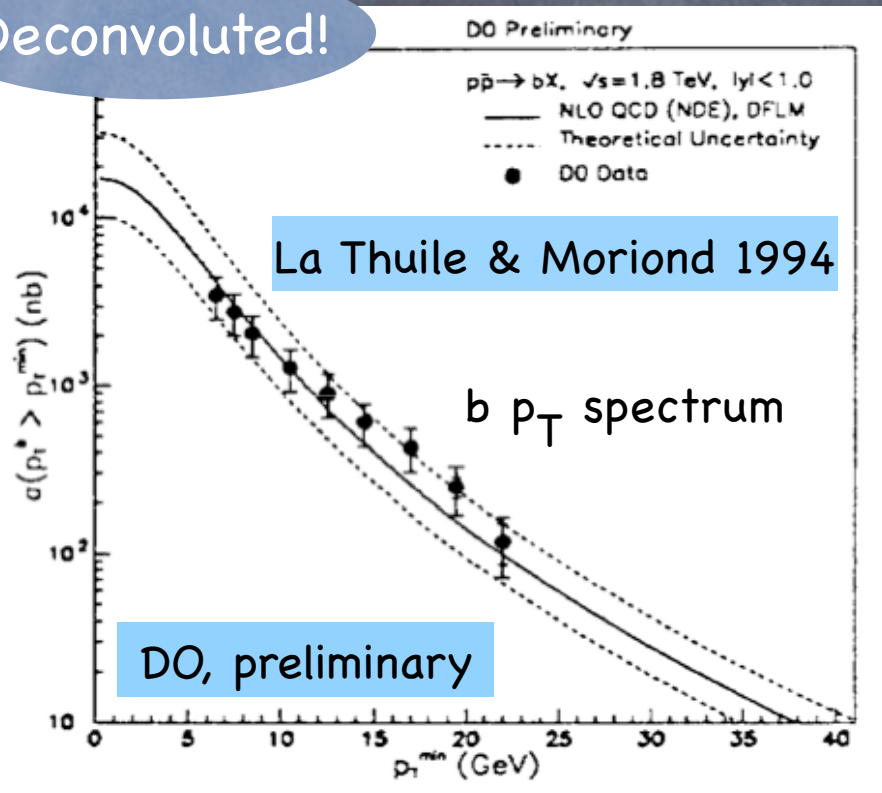
The 'usual' plot enters the stage....

CDF 1993

PRL 71 (1993) 500, PRL 71 (1993) 2396

agreement within the experimental errors. This result supports the conclusion of previous CDF analyses that the next-to-leading order QCD calculation tends to underestimate the inclusive b -quark cross section.

Deconvoluted!



D0 finds however no excess at this stage:
consistent with QCD, barely consistent with CDF

“Real” observables are also measured:

CDF 1995

PRL 75 (1995) 1451

B mesons, NOT deconvoluted
to b quark level

However, how is the theoretical
predictions for B mesons calculated?

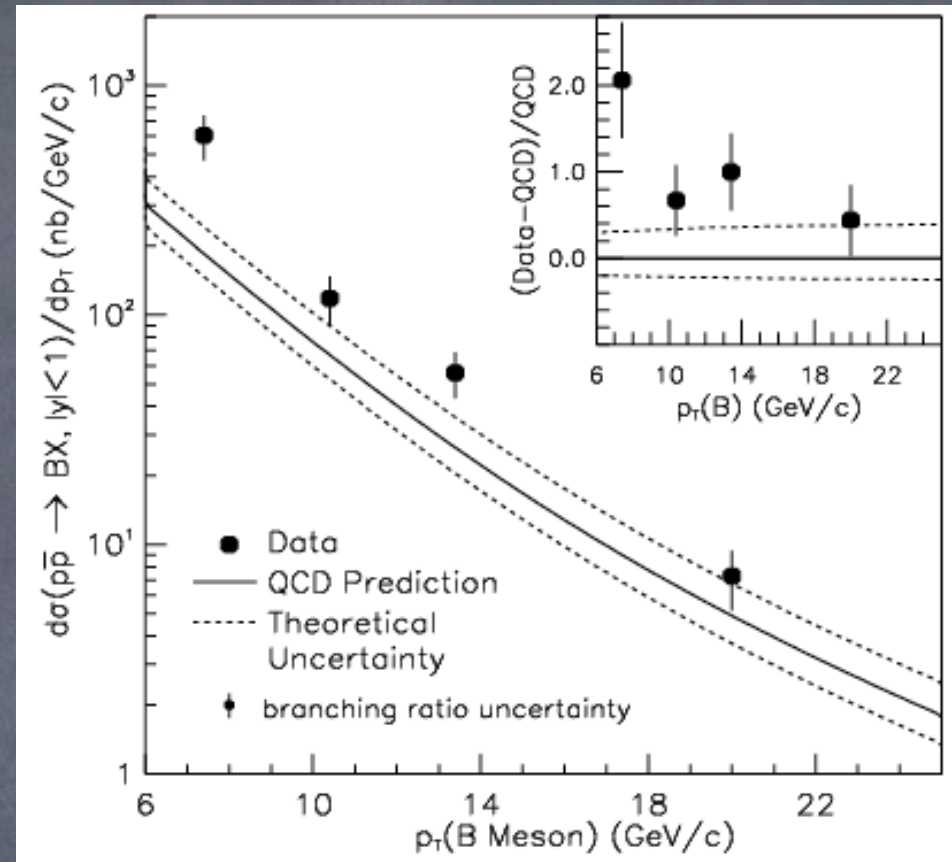


Fig. 2. To determine the level of agreement between the data and the theoretical prediction, the predicted cross section is fitted to the measurements, holding the shape constant and varying the magnitude. The fit yields an overall scale factor of $1.9 \pm 0.2 \pm 0.2$, with a confidence level of 20%. In conclusion, we find that the shape of the B meson differential cross section presented here is adequately described by next-to-leading order QCD, while the absolute rate is at the limits of that predicted by typical variations in the theoretical parameters. It will be interesting

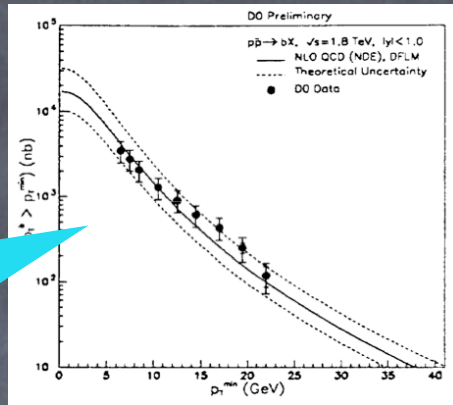
The possible ‘disagreement’ between
data and theory is quantified for
the first time

D0 1995-1996

PRL 74 (1995) 3548

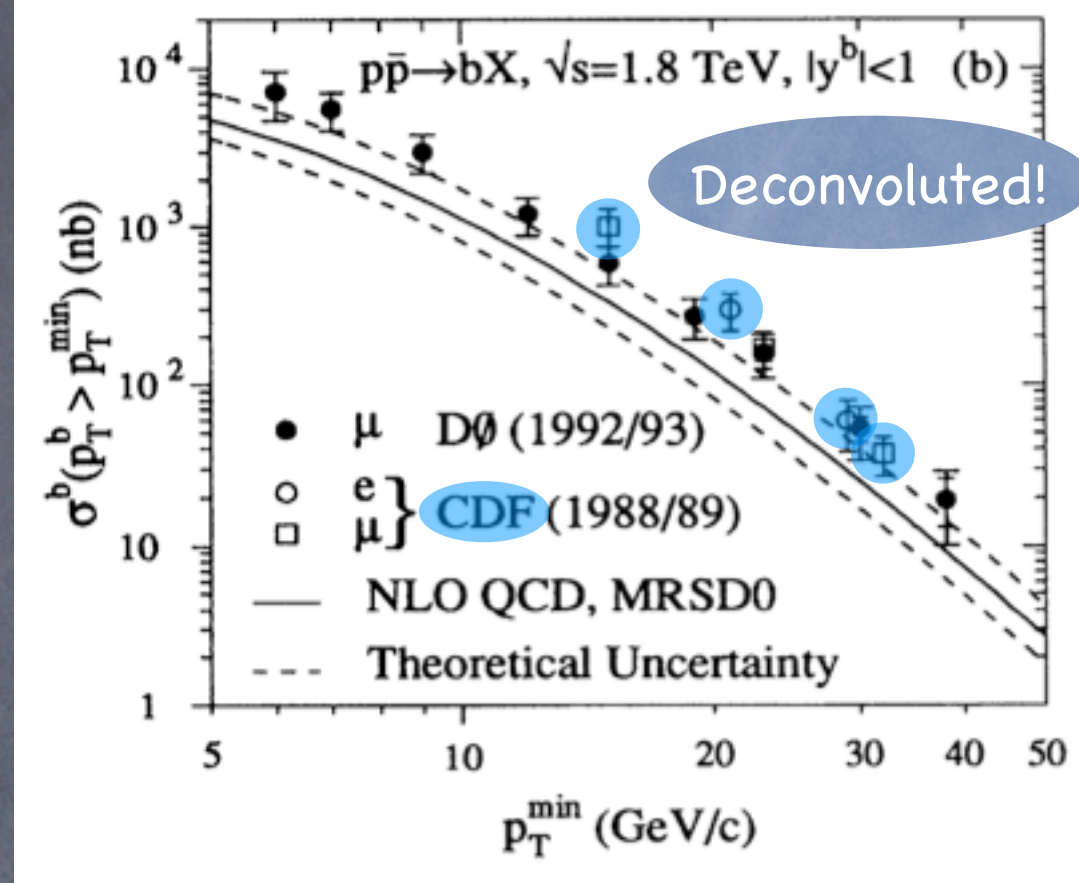
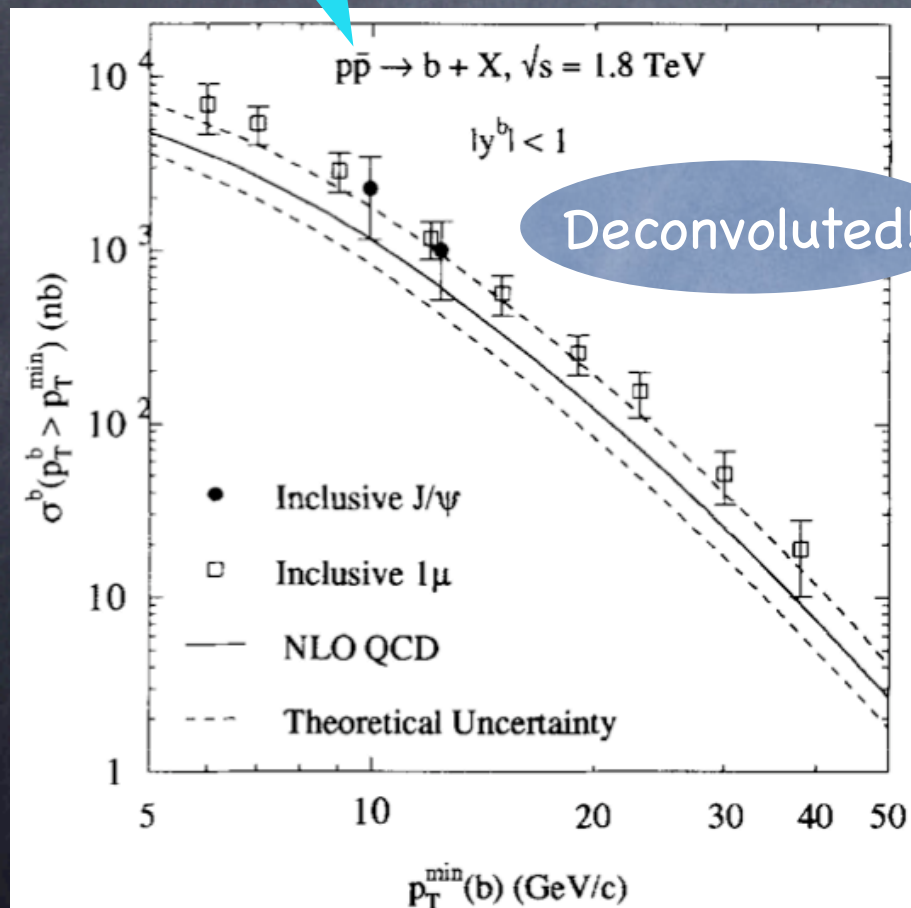
PL B370 (1996) 239

Preliminary



Final

The final D0 data become more CDF-like.



However, they are still compatible with QCD:

Conclusions

tion. Our measurement indicates that, within theoretical uncertainties, the NLO QCD description [1] of heavy flavor production in $p\bar{p}$ at $\sqrt{s} = 1.8$ TeV is adequate for the kinematic range $|y^b| < 1.0$ and $p_T^b > 6$ GeV/c.

A few years later, the data (or the attitude?) change....

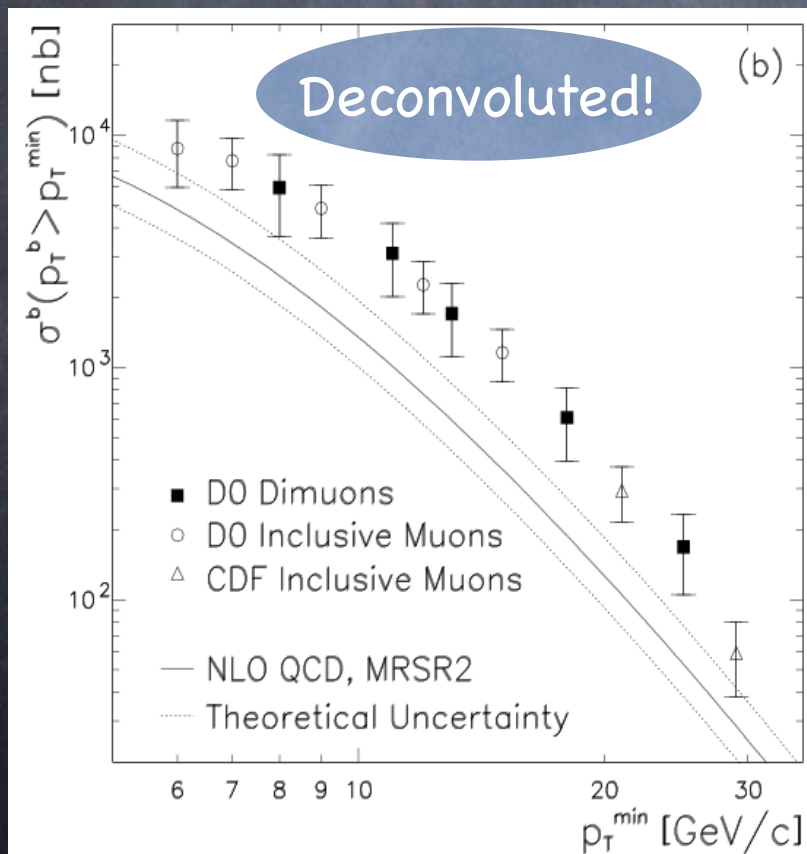
DO 1999–2000

PL B487 (2000) 264

Despite the conclusions of the previous paper ("adequate description"), the previously measured b cross section is now described as being "systematically larger" in the Introduction:

Measurements of the b quark production cross section and $b\bar{b}$ correlations in $p\bar{p}$ collisions provide an important test of perturbative quantum chromodynamics (QCD) at next-to-leading order (NLO). The measured b quark production cross section at $\sqrt{s} = 1.8$ TeV [1–4] is systematically larger than the central values of the NLO QCD predictions [5,6].

tion. Our measurement indicates that, within theoretical uncertainties, the NLO QCD description [1] of heavy flavor production in $p\bar{p}$ at $\sqrt{s} = 1.8$ TeV is adequate for the kinematic range $|y^b| < 1.0$ and $p_T^b > 6$ GeV/c.



This, of course, helps accepting the conclusion that the new data show now a considerable excess:

Conclusions

as described above and is dominated by the variation of the scales. The ratio of the data to the central NLO QCD prediction is approximately three over the entire p_T^{\min} range covered.

approximately three over

Forward muons from b decay

Not quite: TH
systematics not included

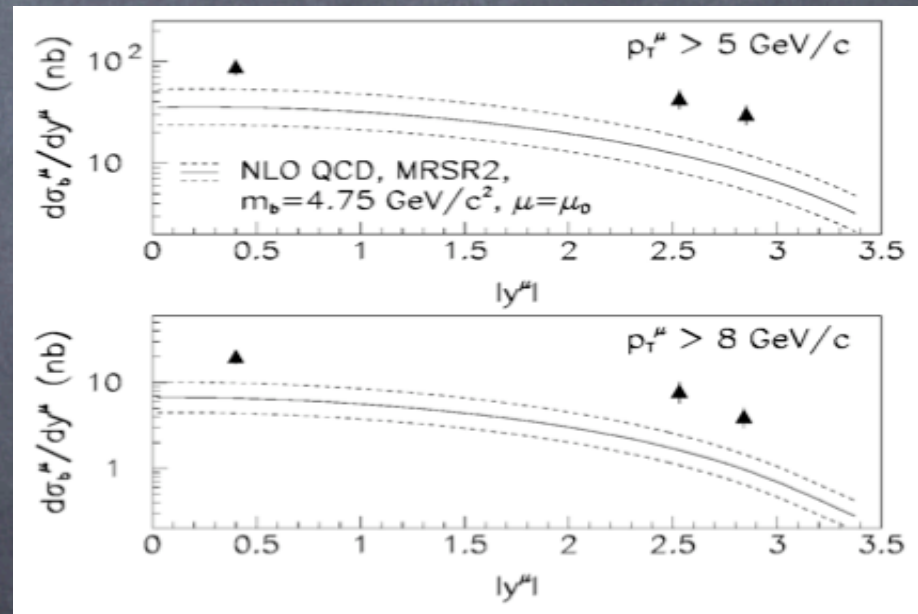
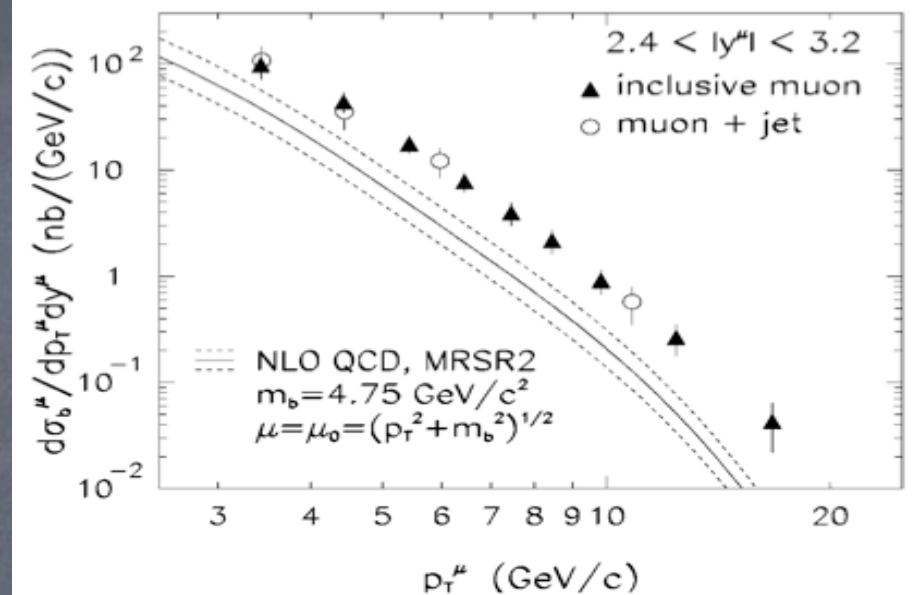
TABLE II. The cross section of muons from b -quark decay compared to NLO QCD. Errors are statistical and systematic added in quadrature.

$p_T^\mu > 5 \text{ GeV}/c$

Rapidity	$\langle y \rangle$	Measured Theory		Ratio
		$\sigma_b^\mu \text{ (nb)}$	$\sigma_b^\mu \text{ (nb)}$	
0.00–0.80	0.40	89 ± 16	36	2.5 ± 0.4
2.40–2.65	2.53	43.5 ± 9.4	12	3.6 ± 0.8
2.65–3.20	2.85	30.5 ± 6.6	8.4	3.6 ± 0.8

$p_T^\mu > 8 \text{ GeV}/c$

Rapidity	$\langle y \rangle$	Measured Theory		Ratio
		$\sigma_b^\mu \text{ (nb)}$	$\sigma_b^\mu \text{ (nb)}$	
0.00–0.80	0.40	20.1 ± 3.7	6.6	3.0 ± 0.6
2.40–2.65	2.53	7.9 ± 2.2	1.6	4.8 ± 1.3
2.65–3.20	2.84	4.1 ± 1.1	0.99	4.0 ± 1.1



Abstract & Conclusions:

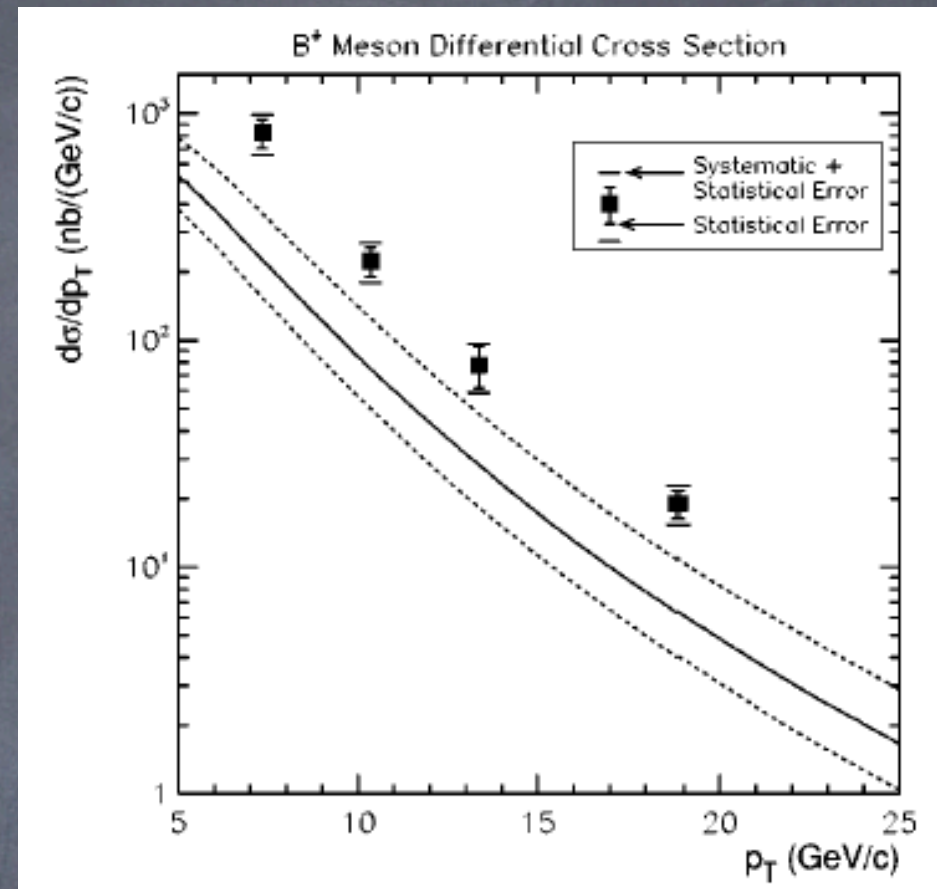
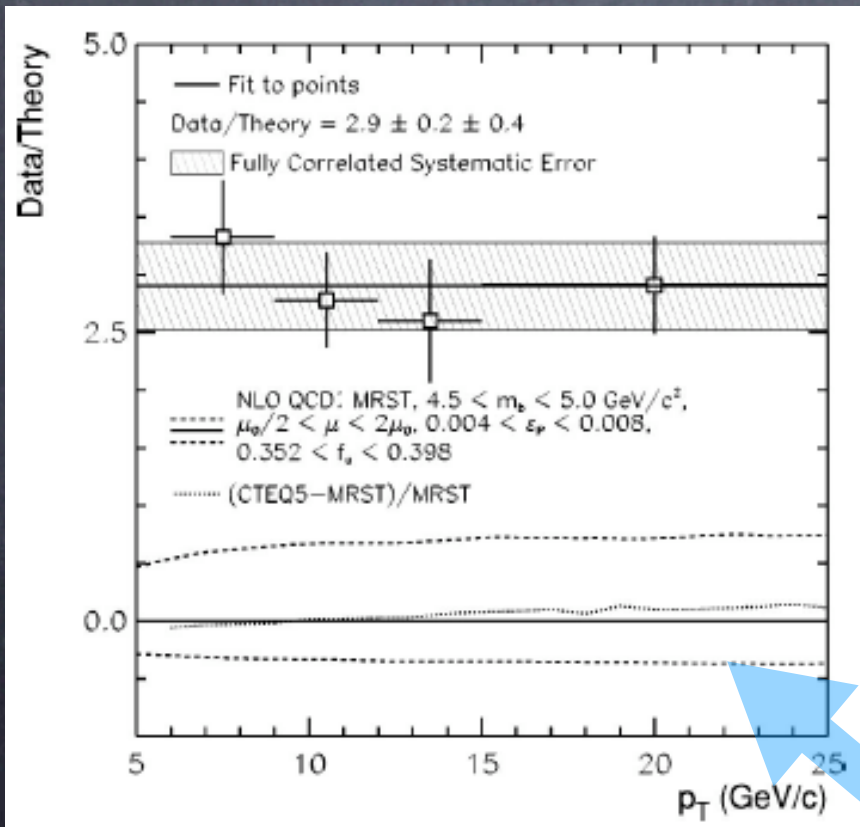
analysis were collected by the D0 experiment at the Fermilab Tevatron. We find that next-to-leading-order QCD calculations underestimate b -quark production by a factor of 4 in the forward rapidity region.

CDF 1998–2002

PRD 65 (2002) 052005

Last CDF Run I result:
B mesons, superseding 1995 result

Data/Theory ratio

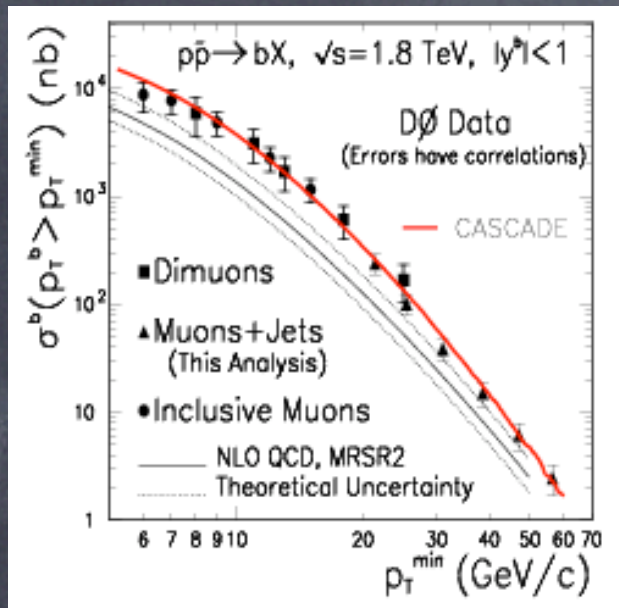


ainties. The differential cross section is measured to be 2.9 ± 0.2 (stat \oplus syst $_{p_T}$) ± 0.4 (syst $_{fc}$) times higher than the NLO QCD predictions with agreement in shape. The first

However, once more, the theoretical uncertainty is not included in the error on the ratio

BTW: being the data points a ratio, shouldn't this band better be around 1 and not 0 !?!

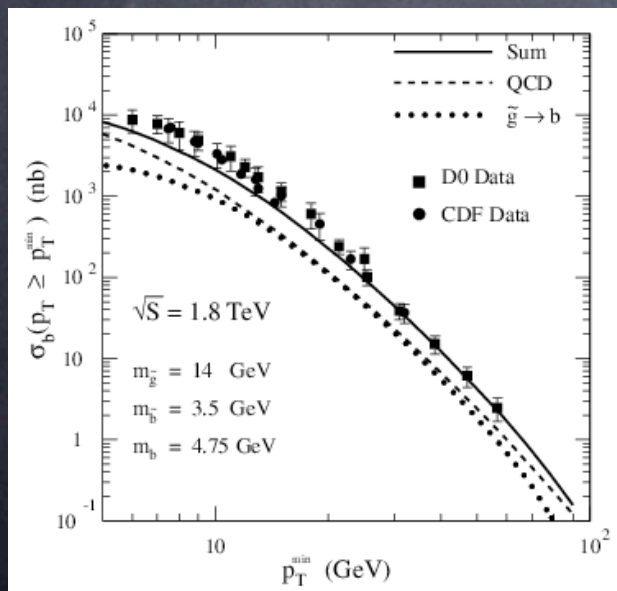
By the years 2001–2002, lots of discrepant data.
Proposed explanations range from the semi-conventional....



H. Jung, CASCADE, [Phys. Rev. D65 (2002) 034015]

MC implementation of small- x dynamics, following CCFM
Main criticism: lack of control of NLO effects

....to the very exotic ones:



Berger, Harris, Kaplan, Sullivan, Tait, Wagner
PRL 86 (2001) 4231

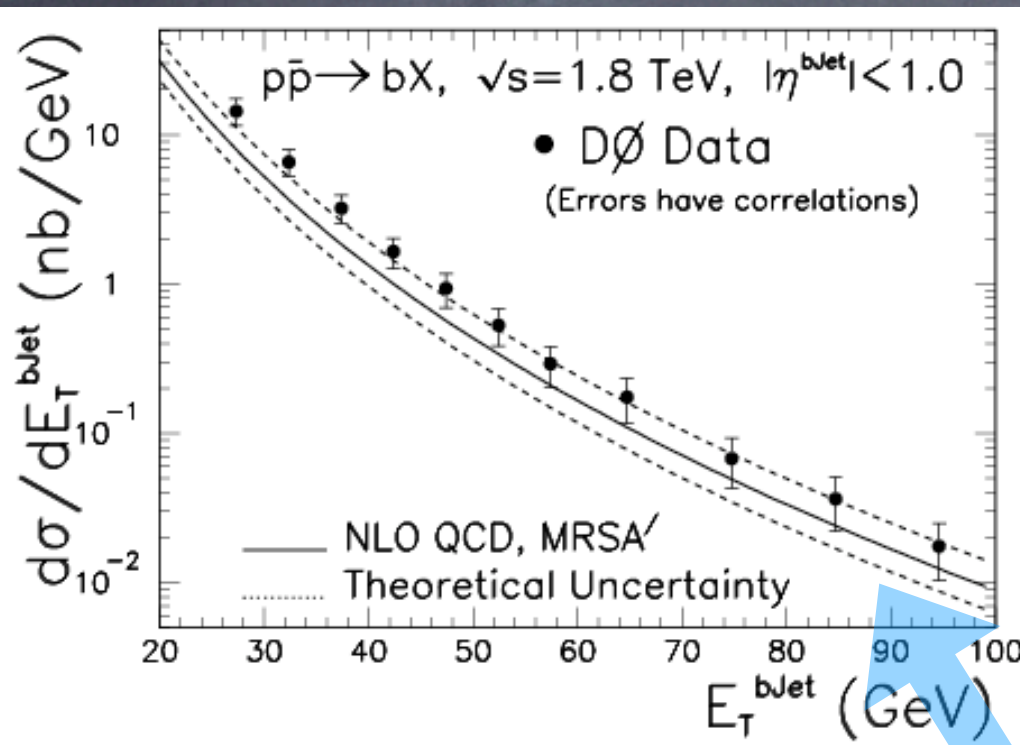
standard model (MSSM) [3]. We postulate the existence of a relatively light gluino \tilde{g} (mass $\approx 12\text{--}16 \text{ GeV}$) that decays into a bottom quark and a light bottom squark \tilde{b} (mass $\approx 2\text{--}5.5 \text{ GeV}$). The \tilde{g} and the \tilde{b} are the spin-1/2 and spin-0 supersymmetric partners of the gluon (g) and bottom quark (b). In our scenario the \tilde{b} is either long-lived or decays hadronically. We obtain good agreement with hadron collider rates of bottom-quark production. Several

NB. Model apparently excluded by
 e^+e^- data, see P. Janot, hep-ph/0403157

D0 2000

PRL 85 (2000) 5068

Most recent D0 results:
b-jets and large p_T b-quarks



b-jets are observable quantities:
no need for a deconvolution

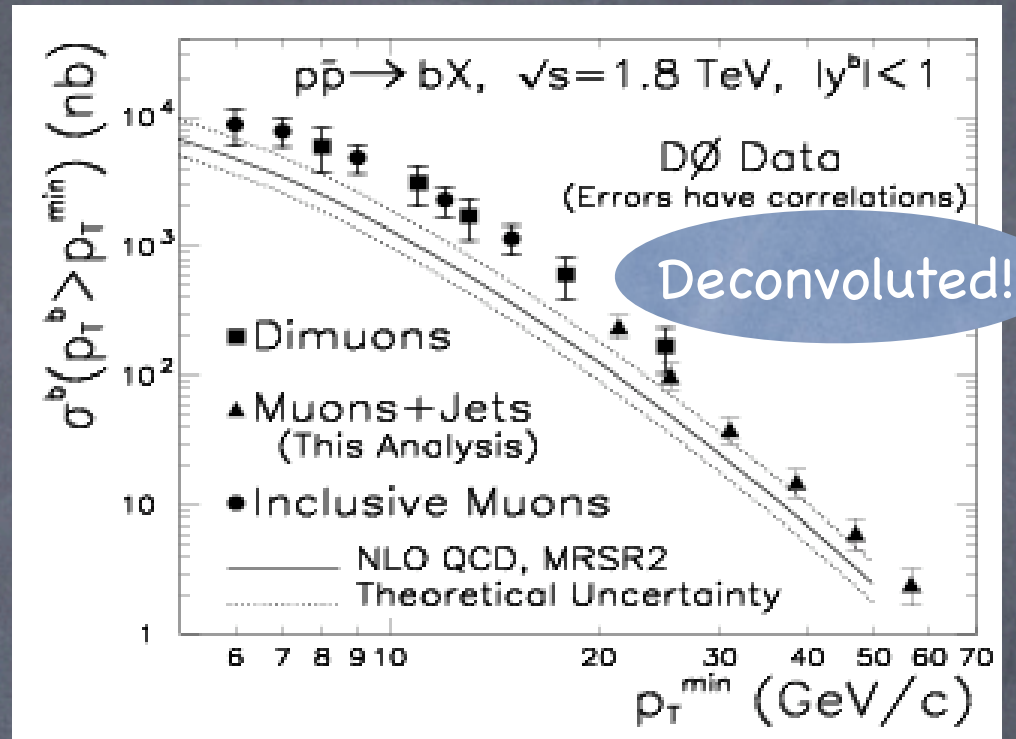


Figure 2 displays the same general pattern of past b production measurements [4–11], with data lying above the central values of the prediction, but comparatively less so in the present case, where general agreement between measurement and the upper band of the theoretical uncertainty is observed.

The Present

Theoretical ingredients of a VCE

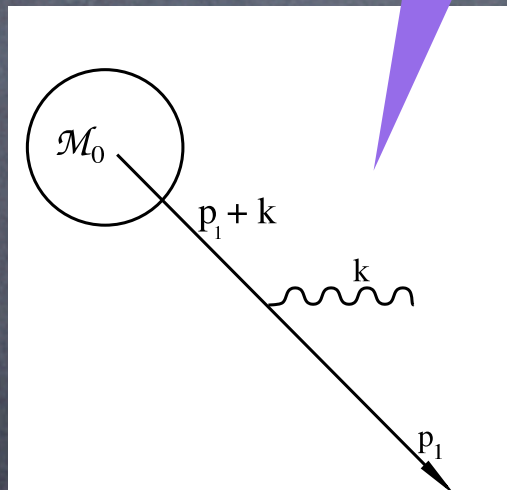
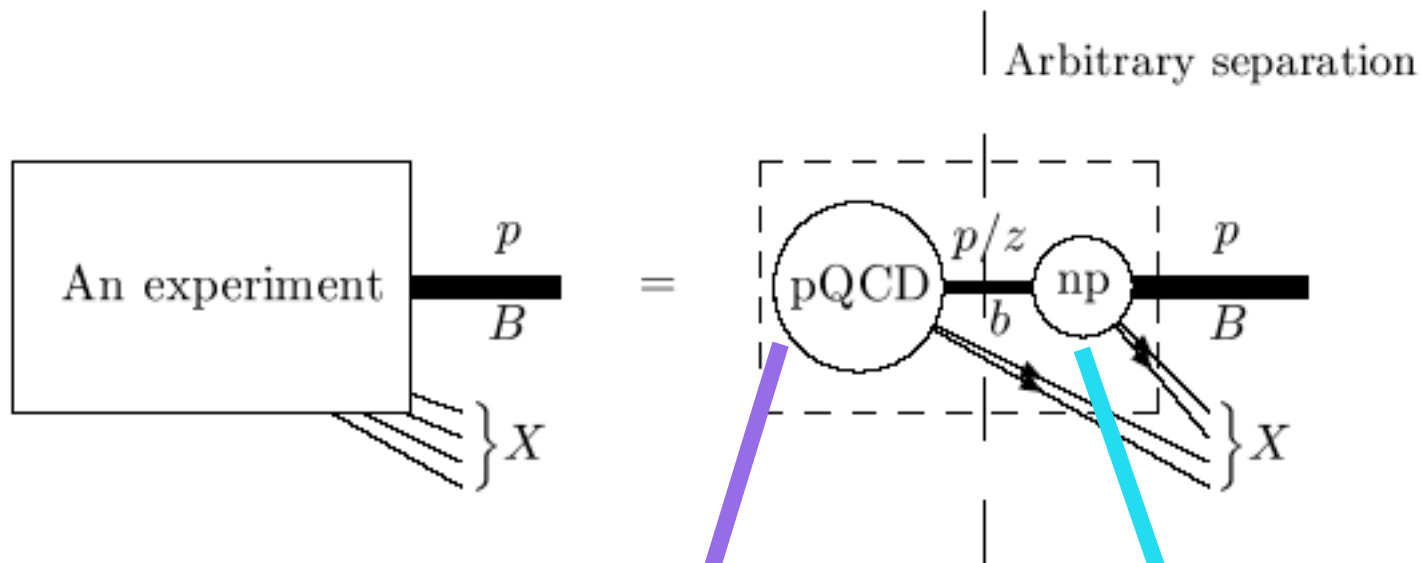
(Very Conventional Explanation)

The prediction for the distribution of a 'real particle' (J/ψ or muon) can be obtained by convoluting:

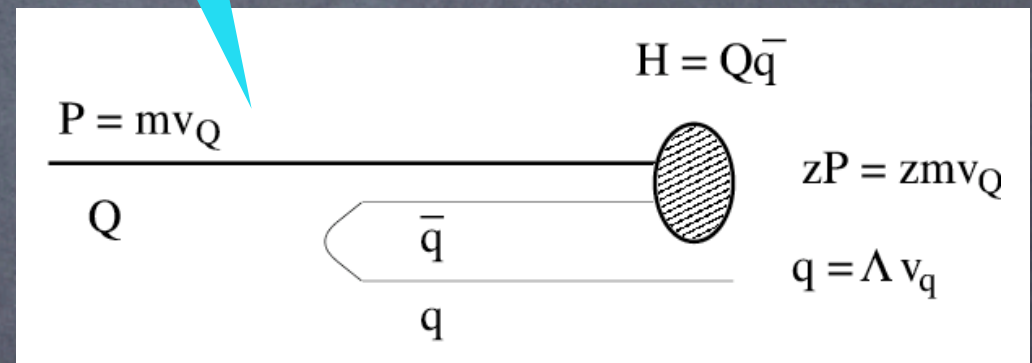
- 1) the NLO (+ NLL = FONLL) calculation for b quarks
- 2) the fragmentation of the b quark into a B meson, $f(b \rightarrow B)$
- 3) the decay of the B meson into the J/ψ or the muon

$$\frac{d\sigma(b \rightarrow B \rightarrow J/\psi)}{dp_T} = \frac{d\sigma(b)}{d\hat{p}_T} \otimes f(b \rightarrow B) \otimes g(B \rightarrow J/\psi)$$

For $f(b \rightarrow B)$ the Peterson, Schlatter, Schmitt, Zerwas form with $\epsilon_b = 0.006$ is used in most experimental papers, following a determination by Chrin made in 1987 (sic) using charm data, $\epsilon_b = m_c^2/m_b^2 \epsilon_c$ rescaling, and LO Montecarlo calculations



Perturbative: gluon radiation



Non-perturbative: hadronization

Not being the b quark a physical particle, $f(b \rightarrow B)$ cannot be a physical observable: its details depend on the perturbative calculation it is interfaced with. A single fragmentation function cannot do for all calculations

Around 1997 [MC, M. Greco, PRD 55 (1997) 7134, M.L. Mangano, lectures on HQ production, hep-ph/9711337] we started arguing that systematics related to fragmentation risked being underestimated, and called for a **stricter consistency between HQ FF determination from e^+e^- data and their use elsewhere**:

For one thing, ϵ_b fitted within a NLO description is smaller than the usual 0.006 value. Hence, a harder Peterson will give a larger cross section in the $p_T > m_b$ region

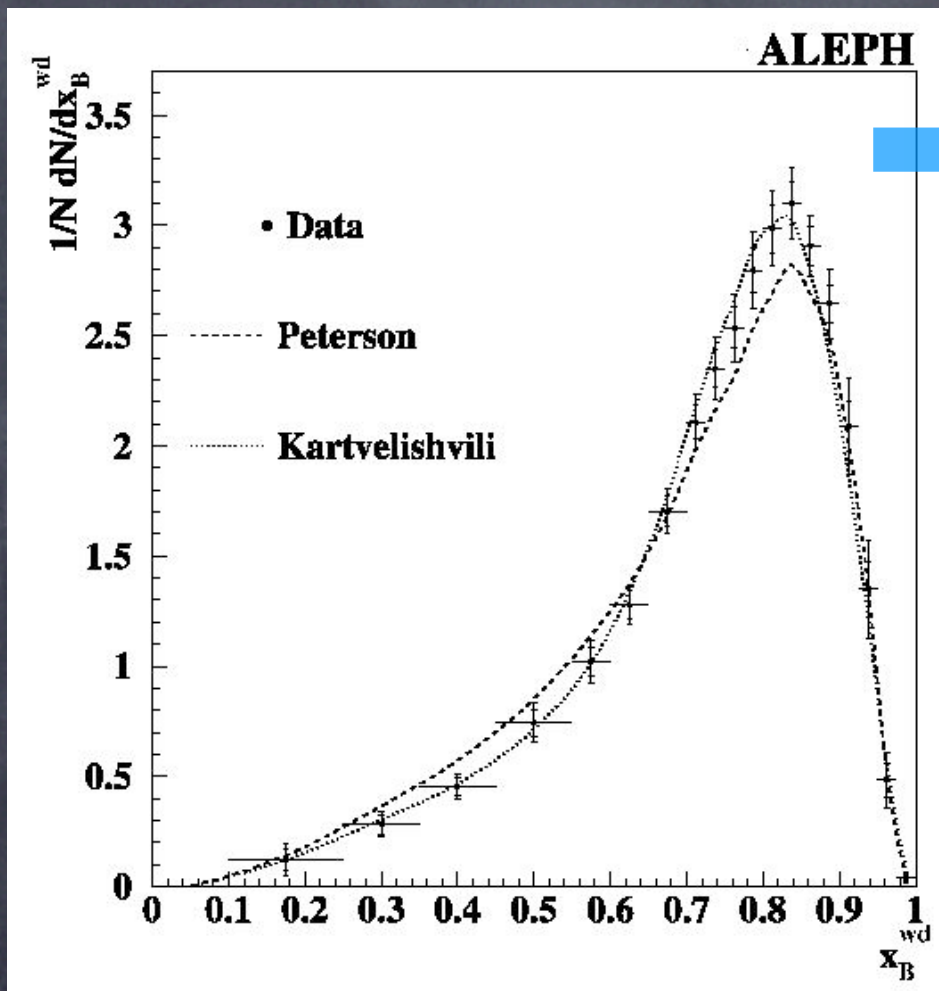
It was also noted that, due to the steeply falling spectrum of the partonic cross section, the transverse momentum distribution in hadronic collisions is sensitive to large moments of the FF, while it is the second moment, $\langle z \rangle$, which is mainly determined from e^+e^- data

$$\text{Assuming } \frac{d\sigma}{d\hat{p}_T} \sim \frac{1}{\hat{p}_T^N} \quad \text{we get} \quad \frac{d\sigma}{dp_T} \sim \int \frac{dz}{z} \left(\frac{z}{\hat{p}_T}\right)^N f(z) = f_N \frac{d\sigma}{d\hat{p}_T}$$

In proton-(anti)proton collisions N is of order 5 for $p_T \sim 10\text{--}20$ GeV. Therefore, **a proper extraction of moments around this one from e^+e^- collisions is more important than a good description of the spectrum**

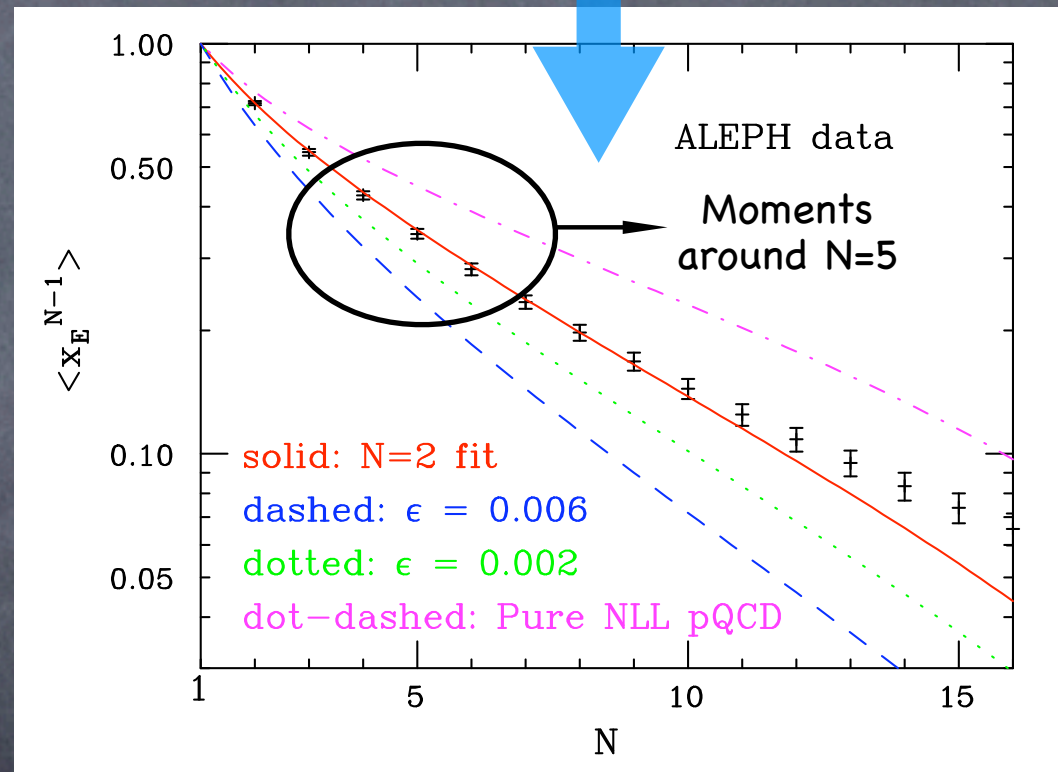
x_E space

Moments space



Distribution of B meson energy fraction

$$\langle x_E^{N-1} \rangle = \int_0^1 x_E^{N-1} f(x_E) dx_E$$



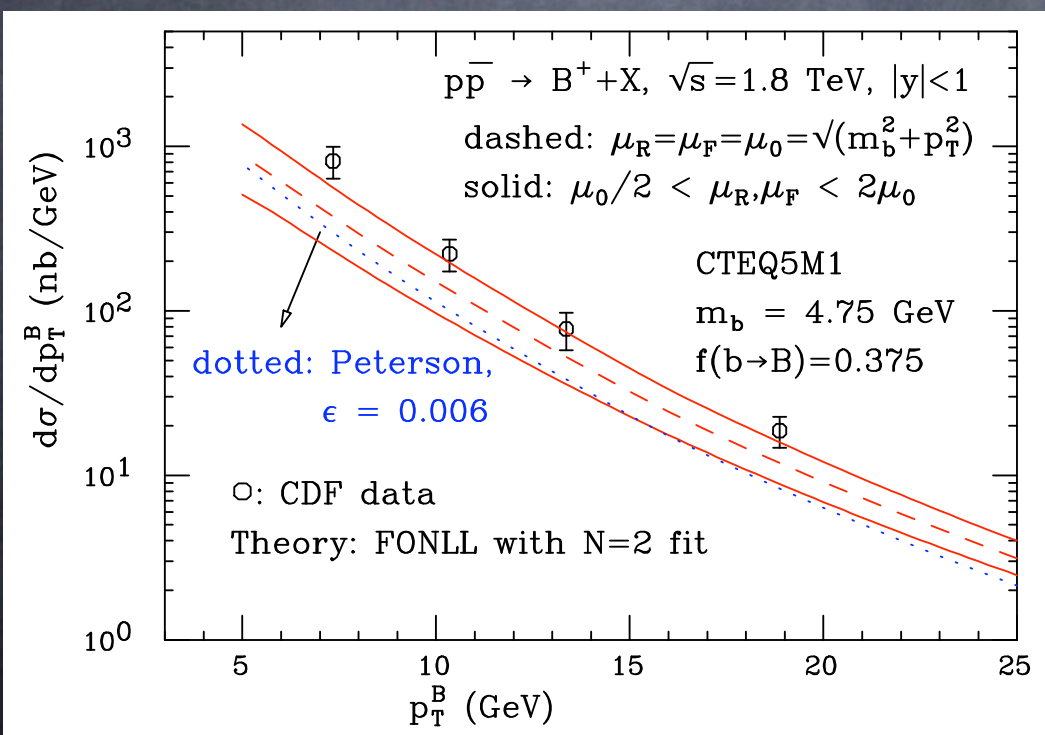
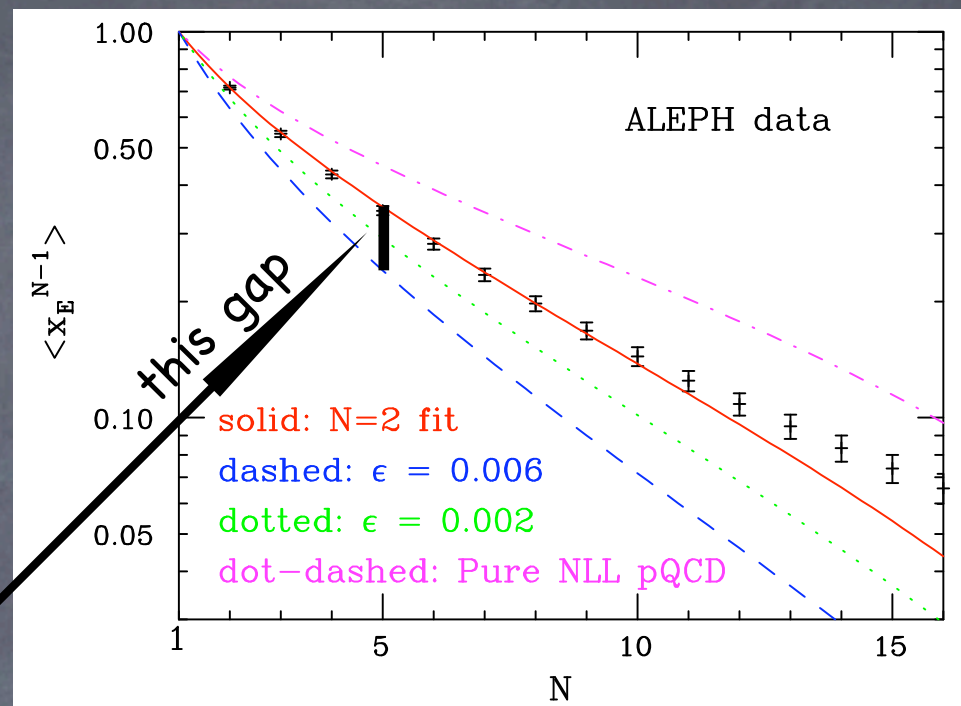
We don't fit this.....

...but rather this.

From the year ~ 2000 accurate enough data on B fragmentation were finally available from LEP, allowing good fits up to $N=10$ or so.

NB. NLL resummed pQCD calculation needed [B. Mele and P. Nason, Nucl. Phys. B361 (1991) 626]

Note that Peterson with $\epsilon_b = 0.006$ underestimates the moments around $N=5$. Its use will consequently underestimate the B cross section



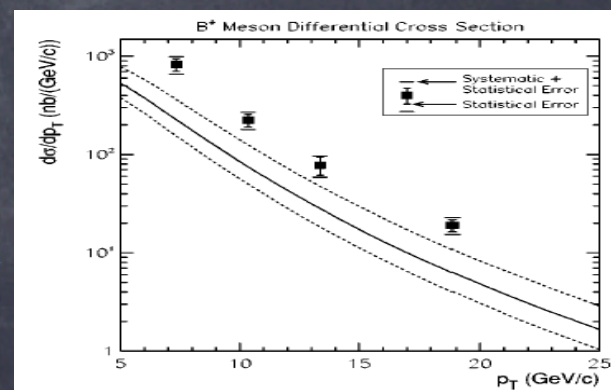
MC and P. Nason, Phys. Rev. Lett. 89 (2002) 122003

With these ingredients, a much better description of the B meson CDF data can be given:

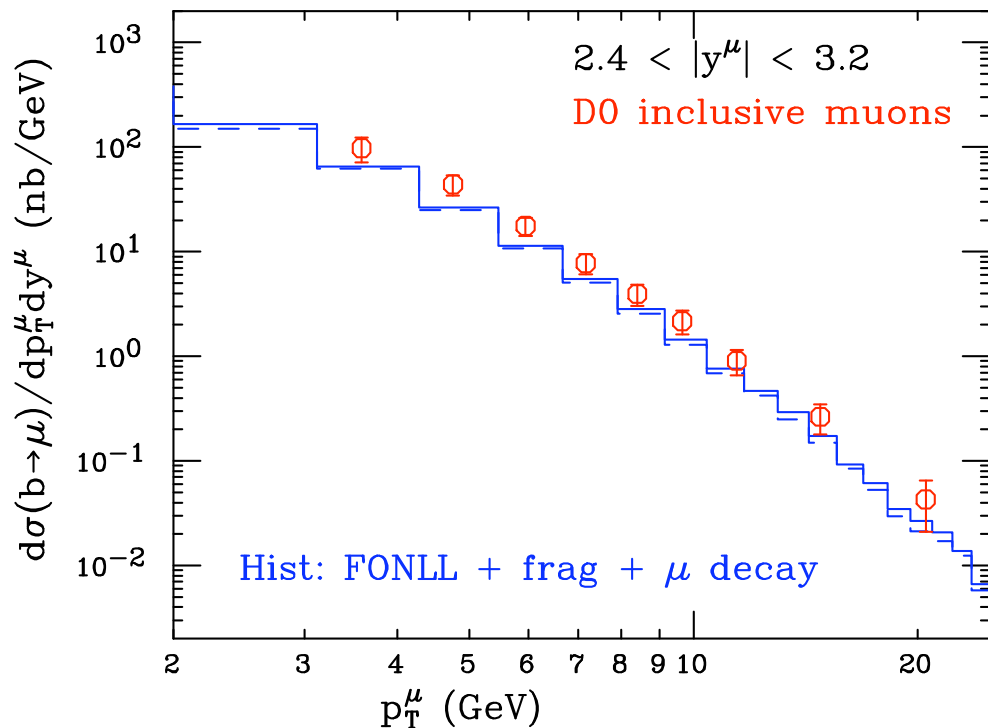
$$\text{Data/Theory} = 1.7 \pm 0.5 (\text{expt.}) \pm 0.5 (\text{th.})$$

i.e. no significant discrepancy

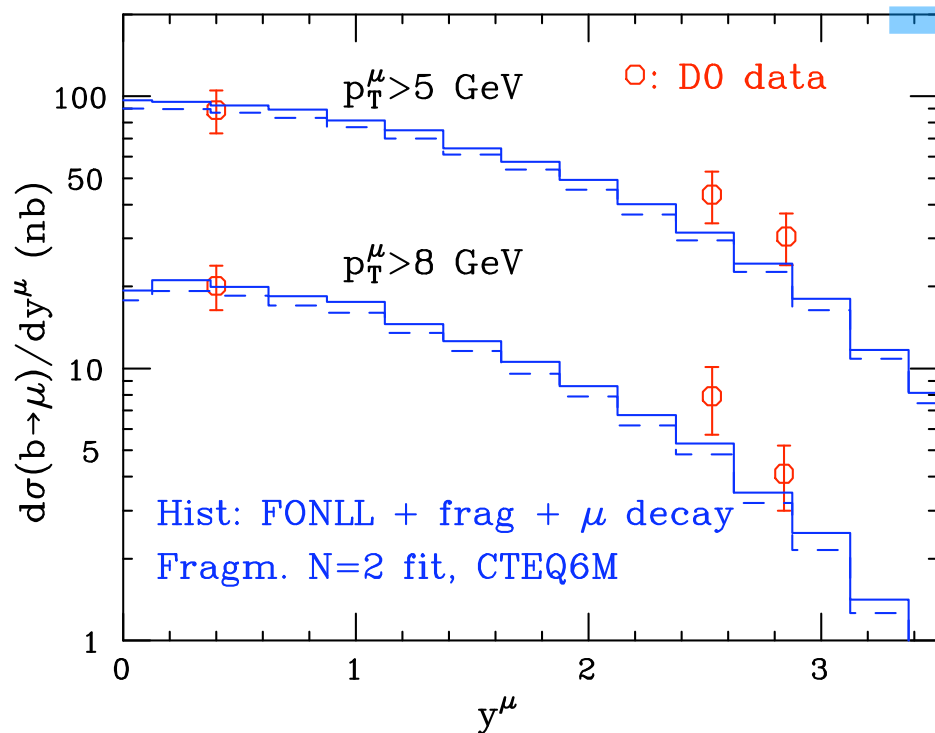
compare



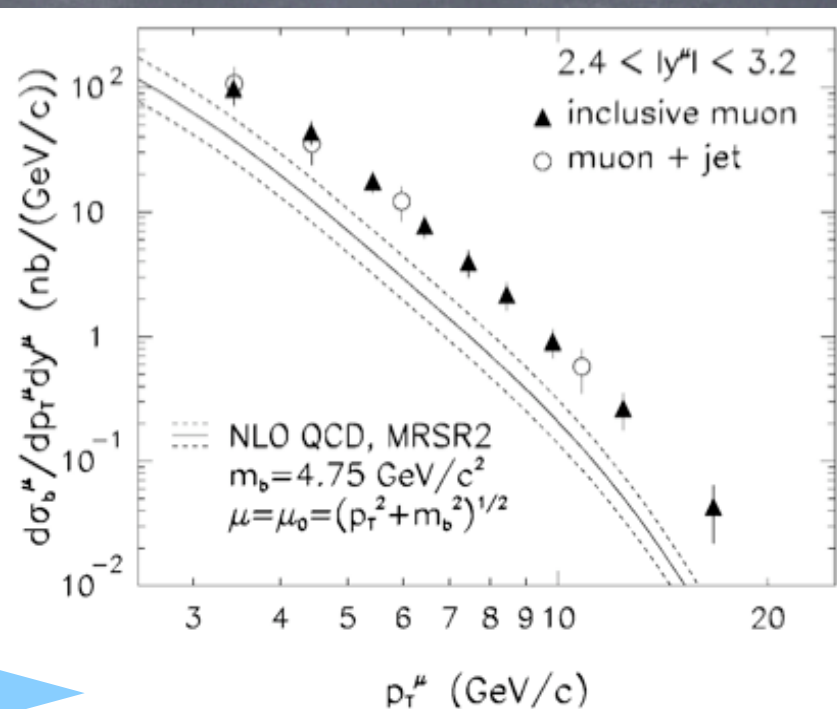
PRELIMINARY



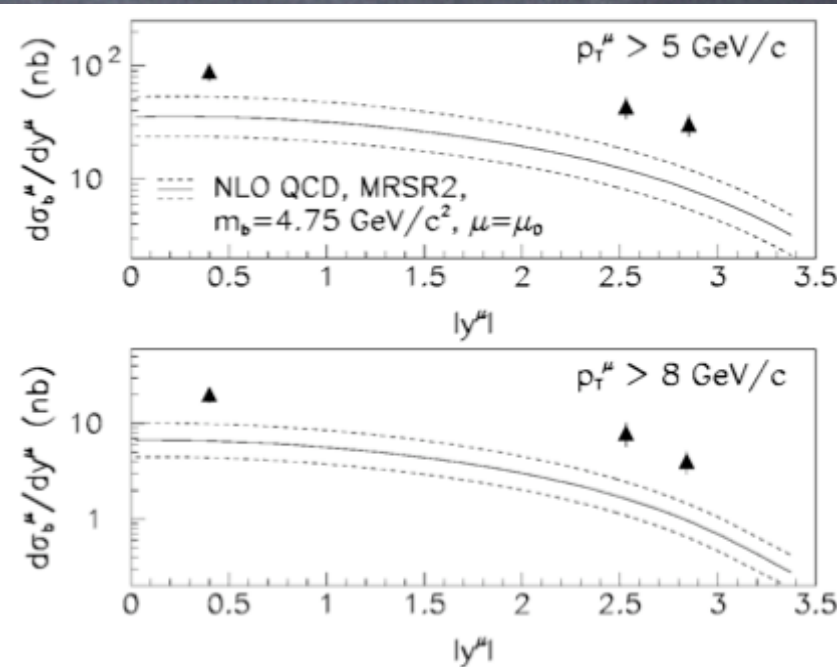
PRELIMINARY



D0 'forward muons from b' data are also now better described

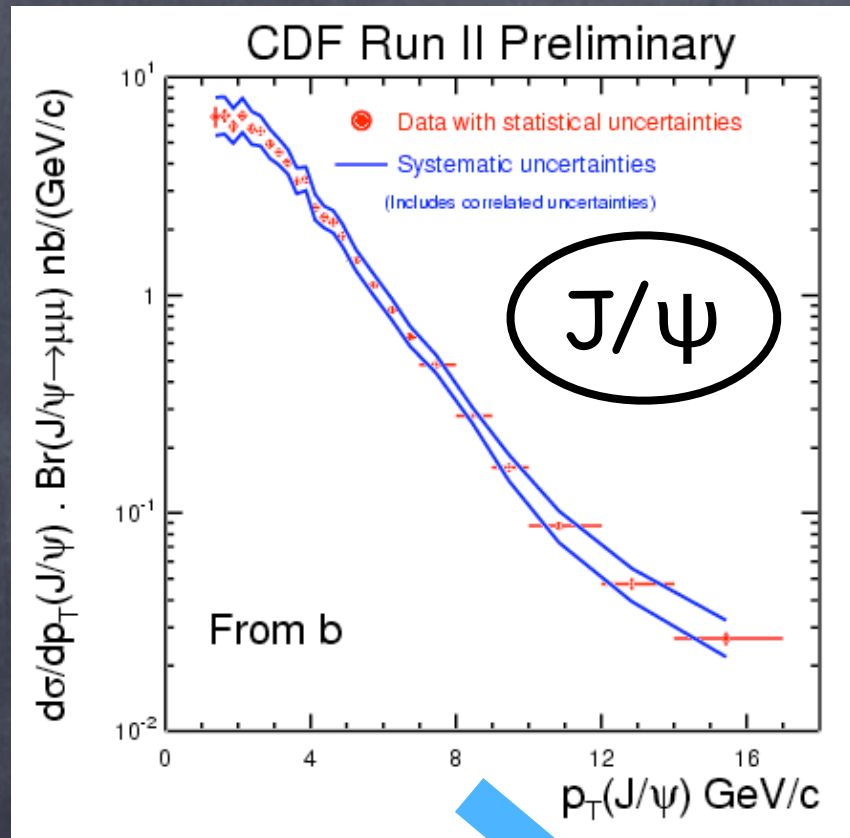


compare

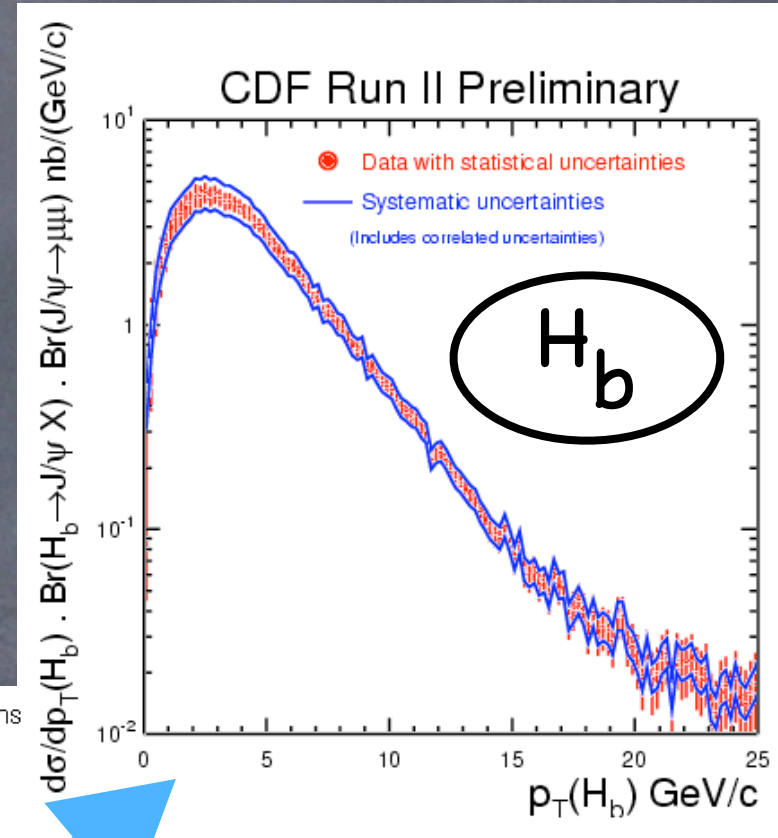
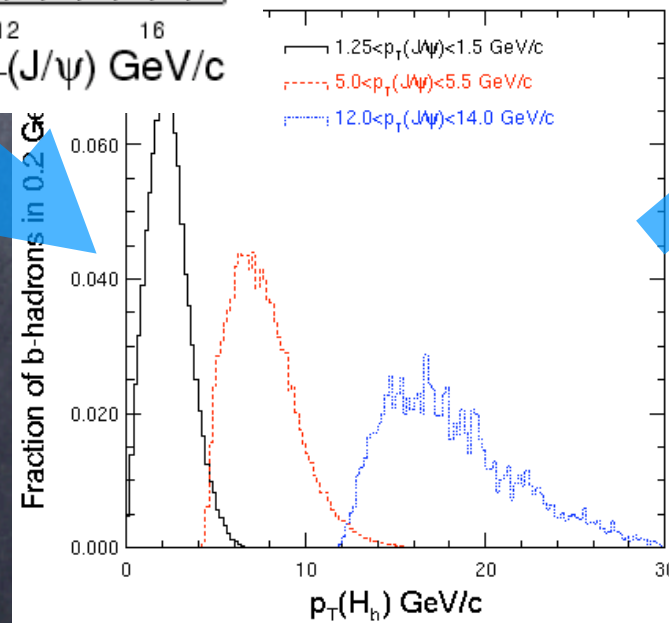


The Tevatron Run II data

CDF published the first preliminary bottom results from Run II data (CDF Note 6285)



Hadron Transverse Momentum Distributions



Simulation of B hadron momentum distribution as a function of the J/ψ momentum

Insofar as QCD effects are concerned, both B hadrons and J/ψ are physical observables

Ingredients of the theoretical prediction

Perturbative items:

- NLO massive calculations
 - NLL resummations
- } FONLL (for LEP + Tevatron)
- Inputs: bottom mass (4.5 – 5 GeV) and α_s ($\Lambda = 0.226$ GeV)
- Uncertainties: ren/fact scale variations

Non-perturbative items:

- gluon and light quarks PDFs
 - b quark to B meson fragmentation
- Input: NLL fit to LEP data (only some moments are important)
- B meson to J/ψ decay spectrum
- Inputs: BR from PDG (1.15 ± 0.06 %)
- Spectrum from CLEO or BABAR
- (detailed knowledge irrelevant due to boost)
- B meson mass (5.3 GeV)

Uncertainties

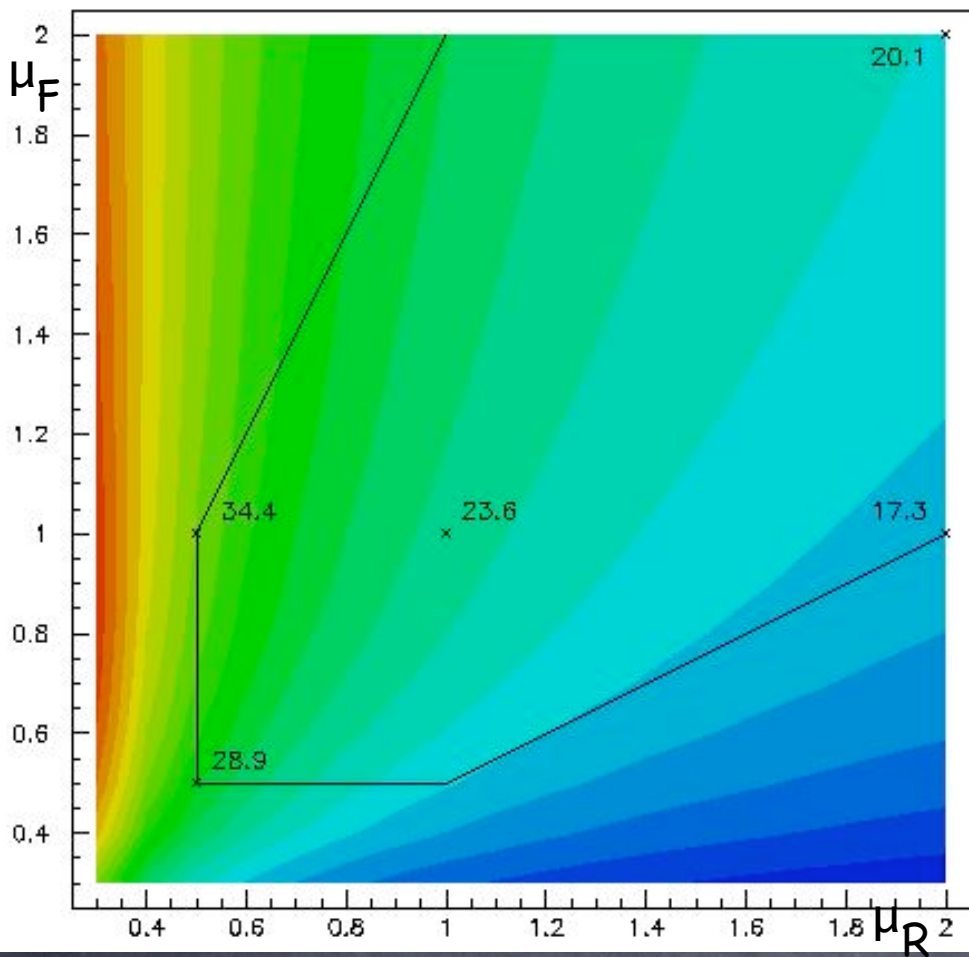
$$\sigma: 28.9 > 23.6 > 20.1 \mu\text{b}$$

$$0.5 < \mu_{R,F}/\mu_0 < 2$$

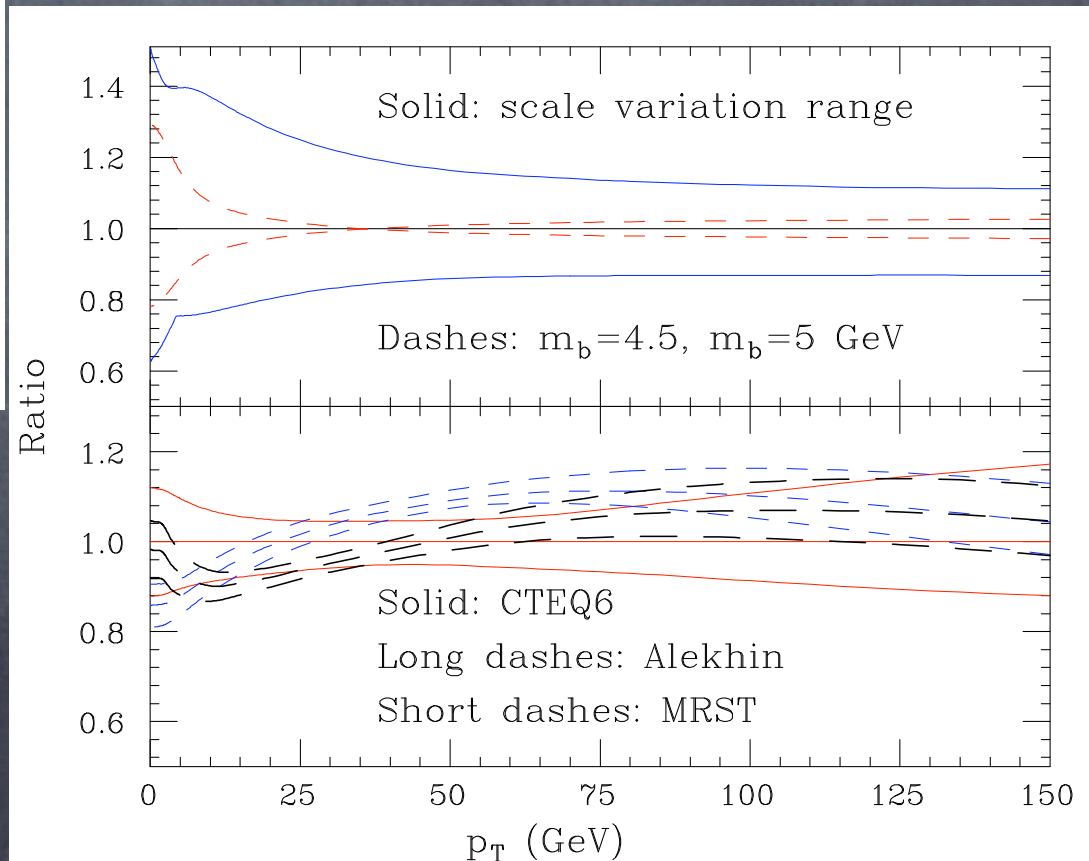
$$\sigma: 34.4 > 23.6 > 17.3 \mu\text{b}$$

$$0.5 < \mu_{R,F}/\mu_0 < 2 \quad \&\& \quad 0.5 < \mu_R/\mu_F < 2$$

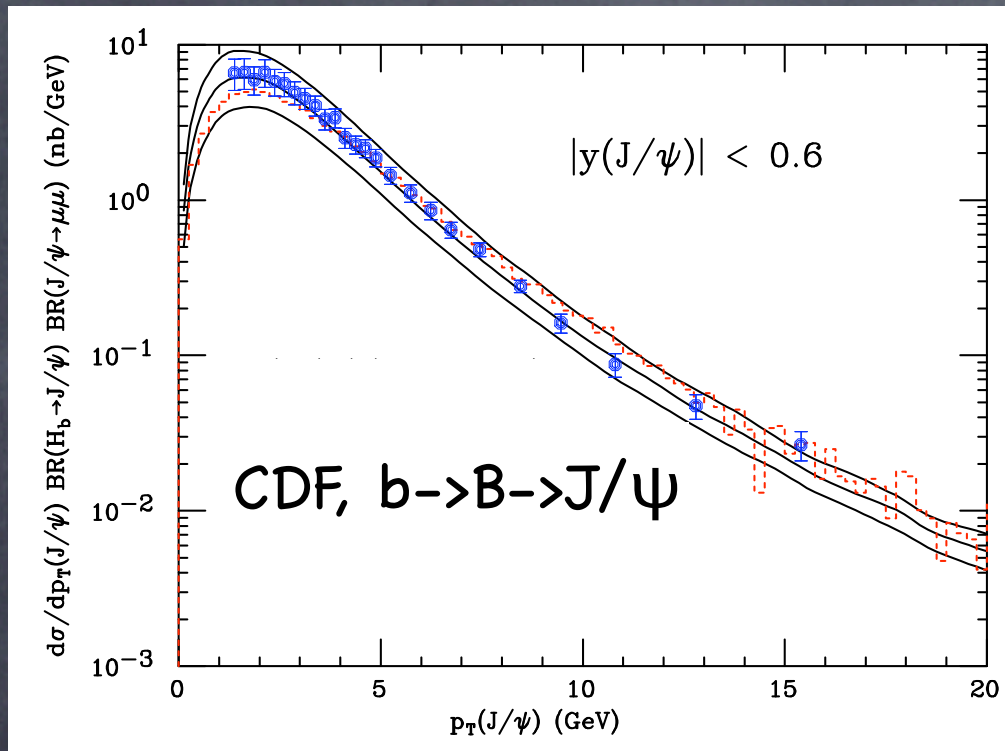
Scale dependence of total
cross section: $\pm 30\text{--}40\%$



Scale dependence and PDFs
uncertainties for transverse
momentum distribution: $\pm 10\text{--}20\%$
at large p_T



2003: CDF Run II preliminary data at 1.96 TeV



$$\sigma(pp \rightarrow H_b \rightarrow \psi; P_{T\psi} > 1.25, |y| < 0.6)$$

$$\sigma_{J/\psi}^{\text{CDF}} = 19.9^{+3.8}_{-3.2 \text{ stat+syst}} \text{ nb}$$

$$\sigma_{J/\psi}^{\text{FONLL}} = 18.3^{+8.1}_{-5.7} \text{ nb}$$

$$\sigma(pp \rightarrow H_b X; P_T > 0, |y| < 0.6) \times B(H_b \rightarrow \psi)$$

$$\sigma_{H_b}^{\text{CDF}} = 24.5^{+4.7}_{-3.9 \text{ stat+syst}} \text{ nb}$$

$$\sigma_{H_b}^{\text{FONLL}} = 22.9^{+9.5}_{-6.8} \text{ nb}$$

$$\sigma(pp \rightarrow b X; P_T > 0, |y| < 1)$$

$$\sigma_b^{\text{CDF}}(|y| < 1) = 29.4^{+6.2}_{-5.4 \text{ stat+syst}} \mu\text{b}$$

$$\sigma_b^{\text{FONLL}}(|y_b| < 1) = 25.0^{+12.6}_{-8.1} \mu\text{b}$$

MC, Frixione, Mangano, Nason, Ridolfi, JHEP 0407 (2004) 033

Theory-Data agreement now almost embarrassing. Fully compatible within errors.

Central values move slightly apart as we go to more 'artificial' cross sections.
Indication of uncertainties and systematics related to deconvolution procedures.

NB. Final data published in PRD 71:032001,2005 (hep-ex/0412071). No significant changes -->

$$\sigma(pp \rightarrow H_b \rightarrow \psi; p_{T\psi} > 1.25, |y| < 0.6)$$

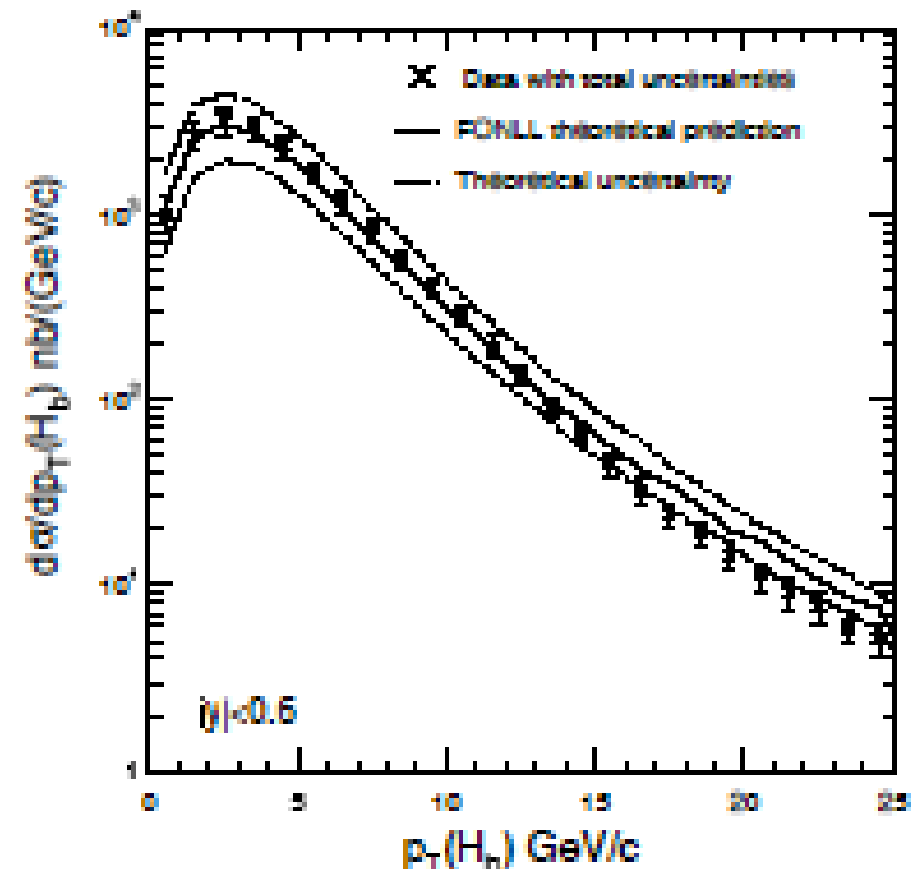
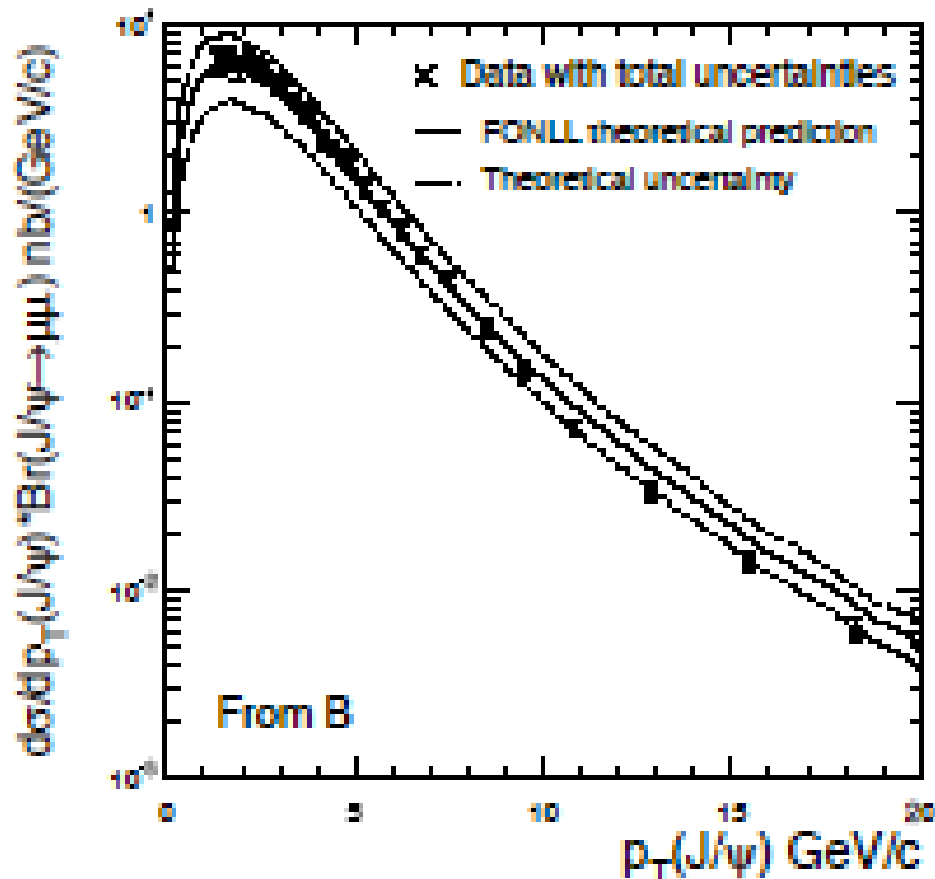
$$\begin{aligned} \sigma(p\bar{p} \rightarrow H_b, H_b \rightarrow J/\psi, p_T(J/\psi) > 1.25 \text{ GeV}/c, |y(J/\psi)| < 0.6) \\ = 0.330 \pm 0.005(stat)^{+0.036}_{-0.033}(syst) \mu b. \end{aligned}$$

$$\sigma_{J/\psi}^{FONLL} = 0.31^{+0.14}_{-0.10} \mu b$$

$$\sigma(pp \rightarrow H_b X; p_T > 0, |y| < 0.6)$$

$$\sigma(p\bar{p} \rightarrow H_b X, |y| < 0.6) = 17.6 \pm 0.4(stat)^{+2.5}_{-2.3}(syst) \mu b.$$

$$\sigma_{(|y| < 0.6)}^{FONLL} = 16.8^{+7.0}_{-5.0} \mu b$$



So, what happened?

How did we go from 'factor of three' excesses to full agreement?

A combination of various factors:

- the real distance between data and theory was actually never this large, once ALL uncertainties were taken into account. Plotting $1\text{-}\sigma$ errors only and discussing central value ratios forgetting errors altogether might have lead to a distorted perception of reality (*'When people quote systematic uncertainties, they usually mean it'. -- M.L. Mangano*)
- both the data and the theory have moved, often legitimately within the uncertainties (which might have been larger than previously thought)
- new measurements without corrections to unphysical particles (ZEUS, CDF) may have minimized the risk of biasing the data. Whatever the reason, they are now in good agreement
- new experimental input (and better use of some of them, e.g. bottom FF) allowed producing more reliable theoretical predictions

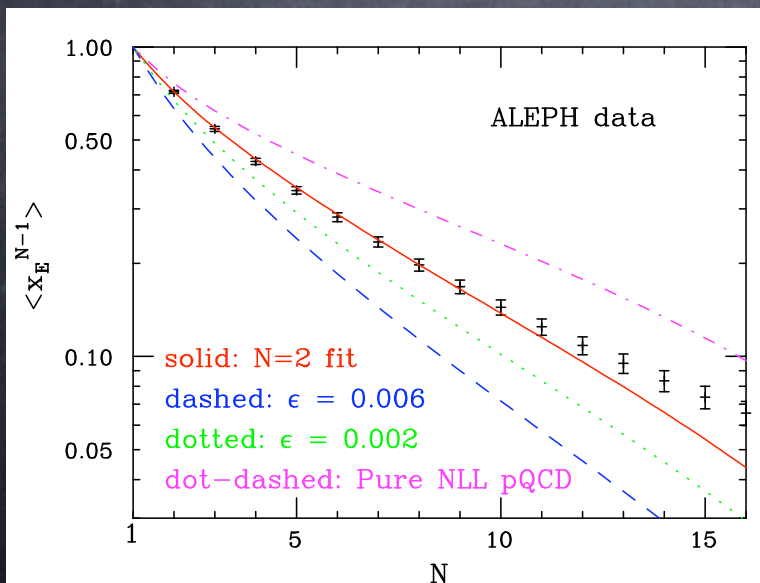
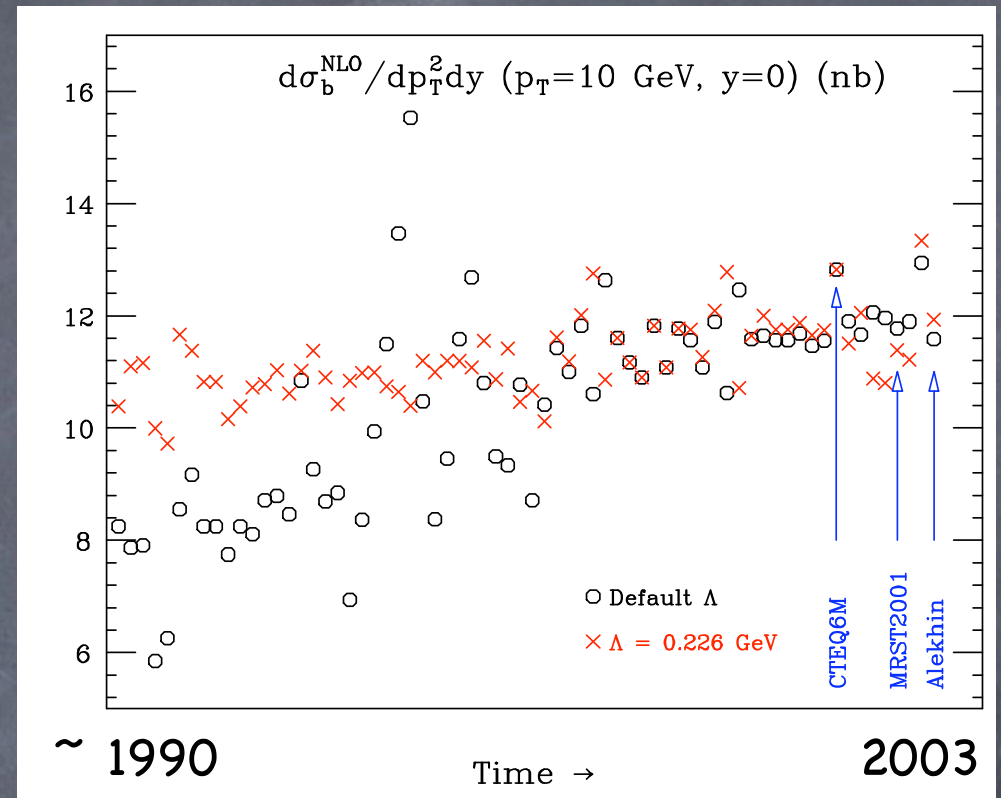
Examples ->

PDFs

Exercise: calculate the b hadroproduction cross section with every PDF set which has ever been published

RESULT: even a factor of two from early sets to modern ones.

NB. a very large part of this discrepancy is due to the evolution of the value for α_s



FFs

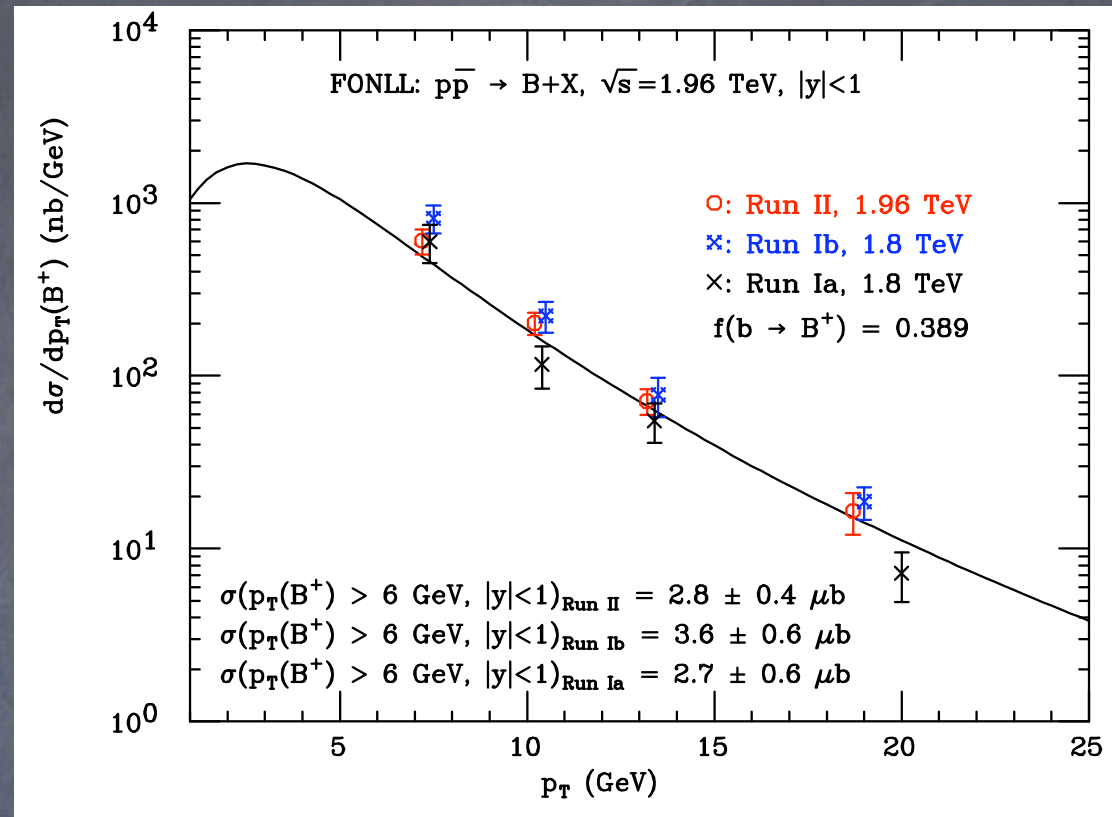
Switching from the usual Peterson with $\epsilon_b = 0.006$ to a FF fitted in moment space increases the large- p_T cross section by 40%.

This is the single most significant increase, and the one not simply due to improved experimental input.

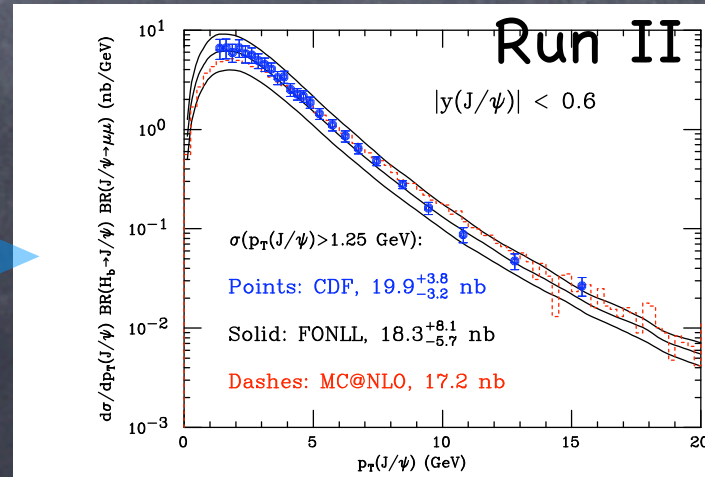
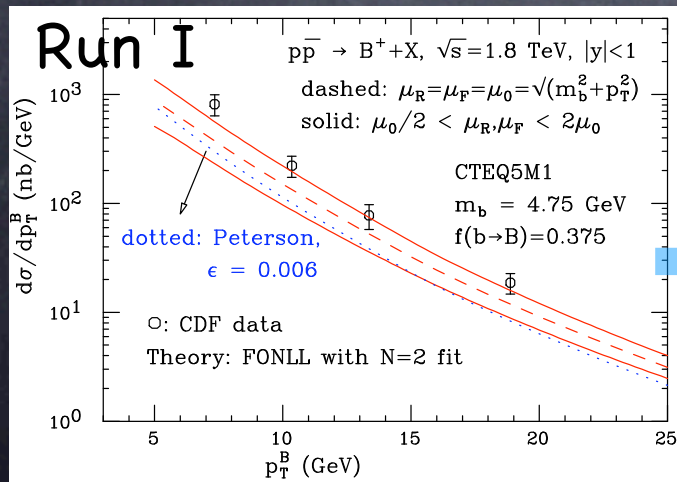
Data

If the input from PDFs and the measurements and extraction of HQ FFs pushed the theory up, the data 'helped' coming down a little:

Compare Run II data to Run I ones:
should be 10% higher, they are
instead about 25% lower

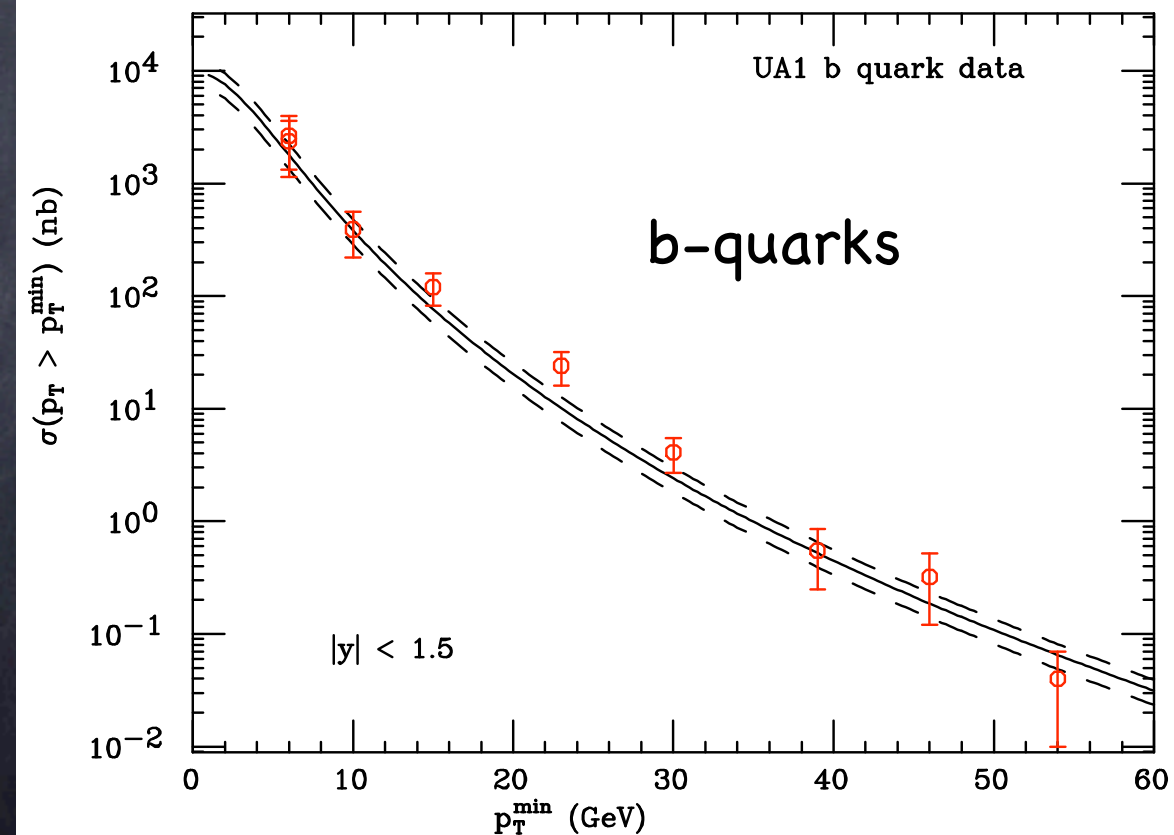
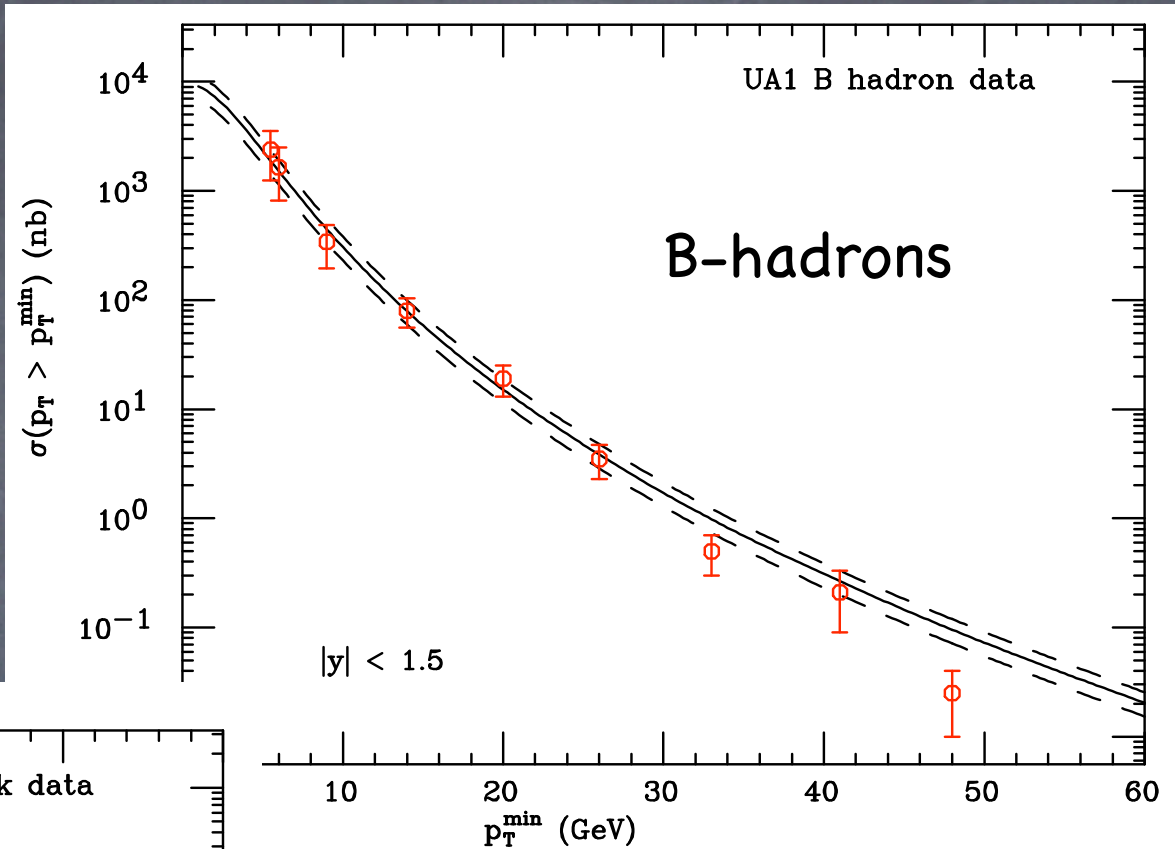


This is the main reason why the same calculation which predicted a CENTRAL VALUE a factor of 1.7 lower than the CENTRAL VALUE of the data, is now in perfect agreement with Run II data



A further 15-20% is given
by updates in the PDFs
(CTEQ5M1 \rightarrow CTEQ6M)

What about old UA1 data?
....everything OK



OLD UA1 data and MODERN theory
appear in good agreement both at the
HADRON and at the QUARK level

NB. Uncertainty bands are only an indication, $\pm 25\%$

Conclusions

- NLO (+NLL) QCD does a good job in predicting real and unbiased bottom hadroproduction data.

Part of the success is due to the possibility of controlling, from the theory side, the whole chain from parton to hadron, carefully matching perturbative and non-perturbative contributions.

Experiments should avoid publishing only deconvoluted/extrapolated quantities, which might include strong biases from MonteCarlo:

“Thou shalt not publish only results for unphysical objects”

- New physics is not needed to explain most of the recent bottom production data, but there is still some room for it within the uncertainties
- Higher order calculations (when?) or further resummations should not change the picture, but may help in reducing the theoretical uncertainties (e.g. small- x effects for total b cross section at the Tevatron)