
Neutralino Annihilation into Quarks Pairs with SUSY-QCD Corrections

Björn Herrmann

Institut für Theoretische Physik und Astrophysik
Universität Würzburg

Work realized in collaboration with Karol Kovařík and Michael Klasen
(to be published)

Orsay, 3 December 2008

Outline

1. Introduction
2. Neutralino annihilation into quarks
3. SUSY-QCD corrections
4. Numerical results
5. Conclusion

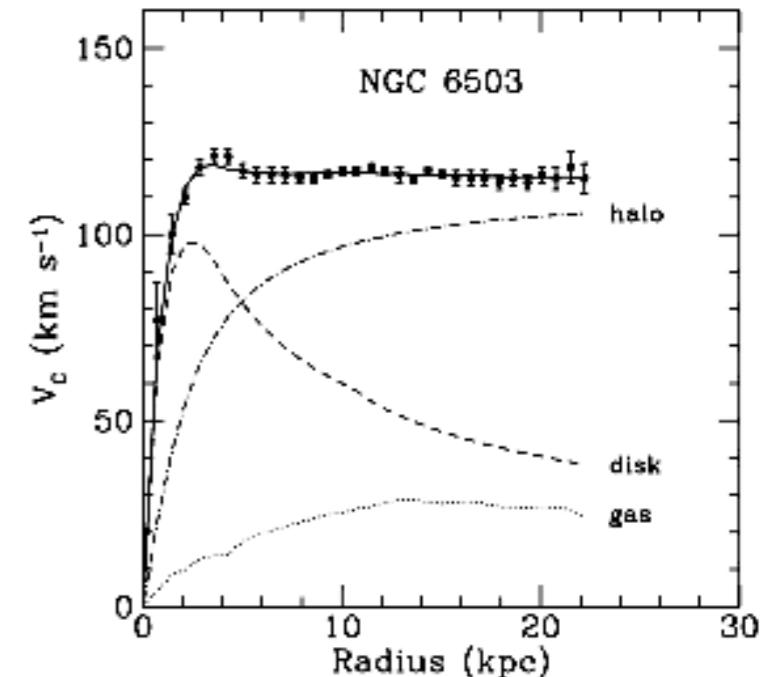
Observational Evidence for Dark Matter

Observational Evidence for Dark Matter

First observational hints

Velocity dispersion and rotation curves

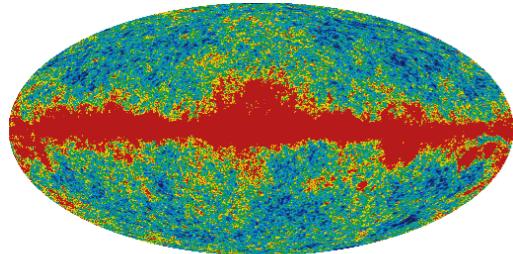
[Zwicky 1933, Rubin et al. 1970]



CMB anisotropies

Cosmological parameters from WMAP mission

[Komatsu et al. (WMAP) 2008]



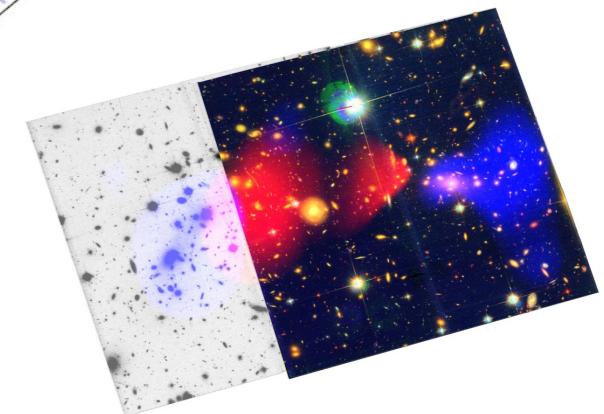
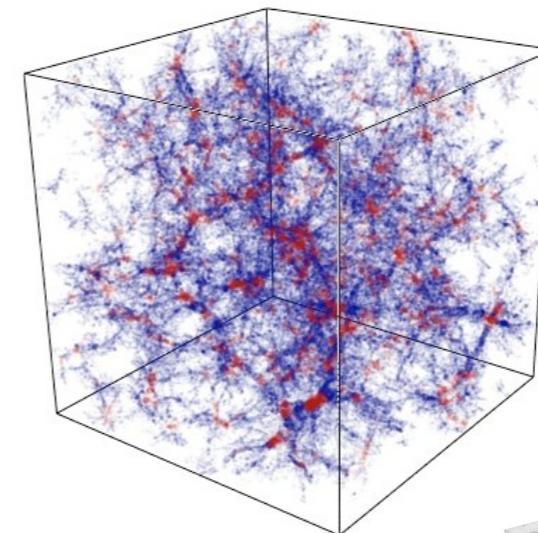
$$\Omega_{\text{tot}} = 1.005 \pm 0.034$$

$$\Omega_{\text{CDM}} = 0.223 \pm 0.013$$

Structure formation

Cold dark matter needed to explain large structures

[Blumenthal et al. 1984]



First direct observation

Observation of “Bullet Cluster” proofs existence of dark matter

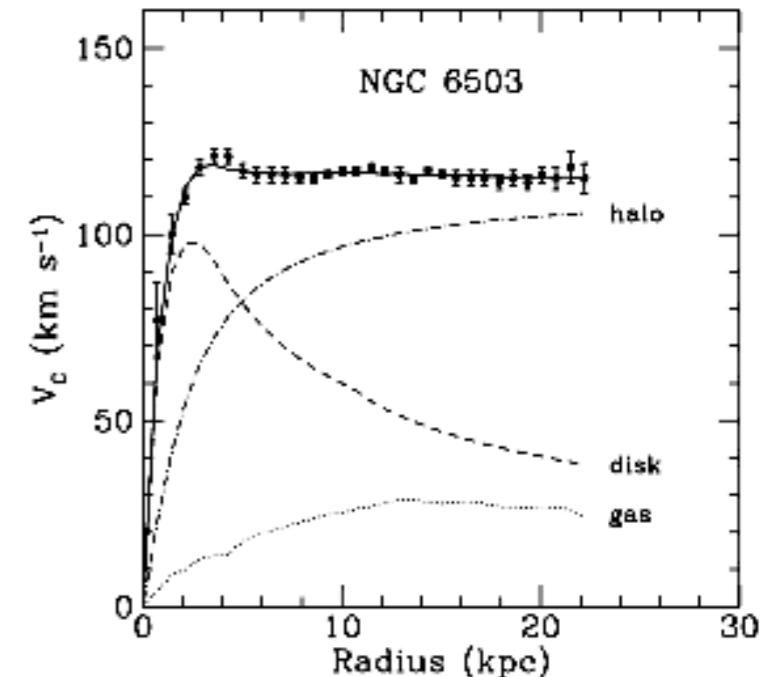
[Clowe et al. 2006]

Observational Evidence for Dark Matter

First observational hints

Velocity dispersion and rotation curves

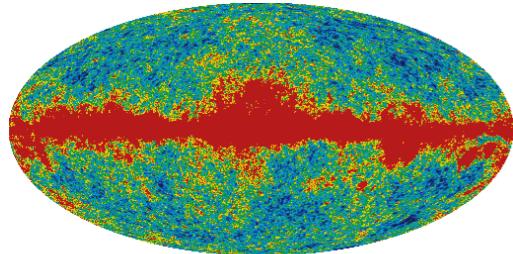
[Zwicky 1933, Rubin et al. 1970]



CMB anisotropies

Cosmological parameters from WMAP mission

[Komatsu et al. (WMAP) 2008]



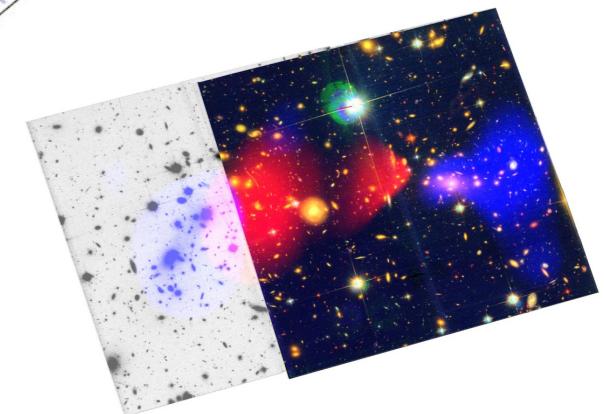
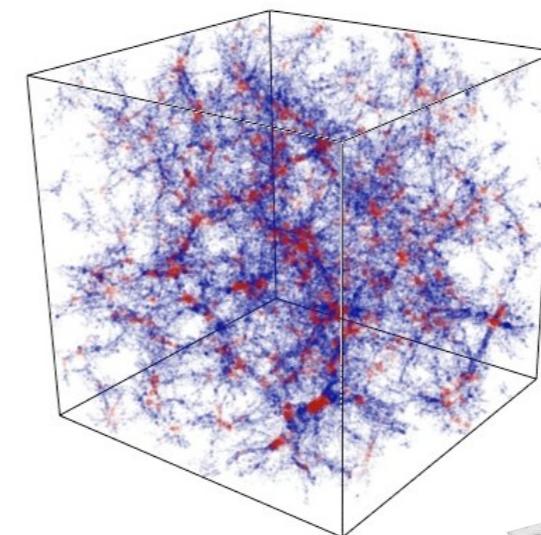
$$\Omega_{\text{tot}} = 1.005 \pm 0.034$$

$$\Omega_{\text{CDM}} = 0.223 \pm 0.013$$

Structure formation

Cold dark matter needed to explain large structures

[Blumenthal et al. 1984]



First direct observation

Observation of “Bullet Cluster” proofs existence of dark matter

[Clowe et al. 2006]

New physics, e.g. Supersymmetry, provides interesting candidates for cold dark matter...

The Minimal Supersymmetric Standard Model (MSSM)

Supersymmetry relates bosons and fermions

$$\mathcal{Q}|b\rangle = |f\rangle \quad \mathcal{Q}|f\rangle = |b\rangle$$

The Minimal Supersymmetric Standard Model (MSSM)

Supersymmetry relates bosons and fermions

$$\mathcal{Q}|b\rangle = |f\rangle \quad \mathcal{Q}|f\rangle = |b\rangle$$

Minimal supersymmetric theory including the Standard Model particles

SM Particles	Spin		Spin		Superpartners
Quarks	$(u_L \ d_L)$	1/2	Q	0	$(\tilde{u}_L \ \tilde{d}_L)$
	u_R^\dagger	1/2	\bar{u}	0	\tilde{u}_R^*
	d_R^\dagger	1/2	\bar{d}	0	\tilde{d}_R^*
Leptons	$(\nu \ e_L)$	1/2	L	0	$(\tilde{\nu} \ \tilde{e}_L)$
	e_R^\dagger	1/2	\bar{e}	0	\tilde{e}_R^*
Higgs	$(H_u^+ \ H_u^0)$	0	H_u	1/2	$(\tilde{H}_u^+ \ \tilde{H}_u^0)$
	$(H_d^0 \ H_d^-)$	0	H_d	1/2	$(\tilde{H}_d^0 \ \tilde{H}_d^-)$
W bosons	W^0, W^\pm	1		1/2	$\tilde{W}^0, \tilde{W}^\pm$
B boson	B^0	1		1/2	\tilde{B}^0
Gluon	g	1		1/2	\tilde{g}
Graviton	G	2		3/2	\tilde{G}
					Gravitino

The Minimal Supersymmetric Standard Model (MSSM)

Supersymmetry relates bosons and fermions

$$\mathcal{Q}|b\rangle = |f\rangle \quad \mathcal{Q}|f\rangle = |b\rangle$$

Minimal supersymmetric theory including the Standard Model particles

SM Particles	Spin		Spin		Superpartners
Quarks	$(u_L \ d_L)$	1/2	Q	0	$(\tilde{u}_L \ \tilde{d}_L)$
	u_R^\dagger	1/2	\bar{u}	0	\tilde{u}_R^*
	d_R^\dagger	1/2	\bar{d}	0	\tilde{d}_R^*
Leptons	$(\nu \ e_L)$	1/2	L	0	$(\tilde{\nu} \ \tilde{e}_L)$
	e_R^\dagger	1/2	\bar{e}	0	\tilde{e}_R^*
Higgs	$(H_u^+ \ H_u^0)$	0	H_u	1/2	$(\tilde{H}_u^+ \ \tilde{H}_u^0)$
	$(H_d^0 \ H_d^-)$	0	H_d	1/2	$(\tilde{H}_d^0 \ \tilde{H}_d^-)$
W bosons	W^0, W^\pm	1		1/2	$\tilde{W}^0, \tilde{W}^\pm$
B boson	B^0	1		1/2	\tilde{B}^0
Gluon	g	1		1/2	\tilde{g}
Graviton	G	2		3/2	\tilde{G}
					Gravitino

Particle mixing after electroweak symmetry breaking gives rise to neutralinos and charginos

Supersymmetric Dark Matter

R-parity conservation implies that lightest supersymmetric particle (LSP) is stable

→ interesting candidate for dark matter

$$P_R = (-1)^{3(B-L)+2s} = \begin{cases} +1 & \text{for Standard Model particles} \\ -1 & \text{for their superpartners} \end{cases}$$

Supersymmetric Dark Matter

R-parity conservation implies that lightest supersymmetric particle (LSP) is stable

→ interesting candidate for dark matter

$$P_R = (-1)^{3(B-L)+2s} = \begin{cases} +1 & \text{for Standard Model particles} \\ -1 & \text{for their superpartners} \end{cases}$$

Here: Focus on minimal supergravity (mSUGRA)

- SUSY breaking mediated between hidden and observable sectors by gravity
- five universal parameters at the grand unification scale

$$m_0, m_{1/2}, A_0, \tan \beta, \text{sgn}(\mu)$$

Supersymmetric Dark Matter

R-parity conservation implies that lightest supersymmetric particle (LSP) is stable

→ interesting candidate for dark matter

$$P_R = (-1)^{3(B-L)+2s} = \begin{cases} +1 & \text{for Standard Model particles} \\ -1 & \text{for their superpartners} \end{cases}$$

Here: Focus on minimal supergravity (mSUGRA)

- SUSY breaking mediated between hidden and observable sectors by gravity
- five universal parameters at the grand unification scale

$$m_0, m_{1/2}, A_0, \tan \beta, \text{sgn}(\mu)$$

Dark matter candidate: Lightest neutralino

- candidate “par excellence” for weakly interacting massive particle (WIMP)

Supersymmetric Dark Matter

R-parity conservation implies that lightest supersymmetric particle (LSP) is stable

→ interesting candidate for dark matter

$$P_R = (-1)^{3(B-L)+2s} = \begin{cases} +1 & \text{for Standard Model particles} \\ -1 & \text{for their superpartners} \end{cases}$$

Here: Focus on minimal supergravity (mSUGRA)

- SUSY breaking mediated between hidden and observable sectors by gravity
- five universal parameters at the grand unification scale

$$m_0, m_{1/2}, A_0, \tan \beta, \text{sgn}(\mu)$$

Dark matter candidate: Lightest neutralino

- candidate “par excellence” for weakly interacting massive particle (WIMP)

Relic density of cold dark matter candidate required to agree with WMAP+SN+BAO data

- identify (dis)favoured regions of the parameter space

$$0.094 \leq \Omega_{\text{CDM}} h^2 \leq 0.136$$

[Hamann et al. 2007]

Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle\sigma_{\text{ann}}v\rangle(n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}}h^2 \propto n \propto \frac{1}{\langle\sigma_{\text{ann}}v\rangle}$$

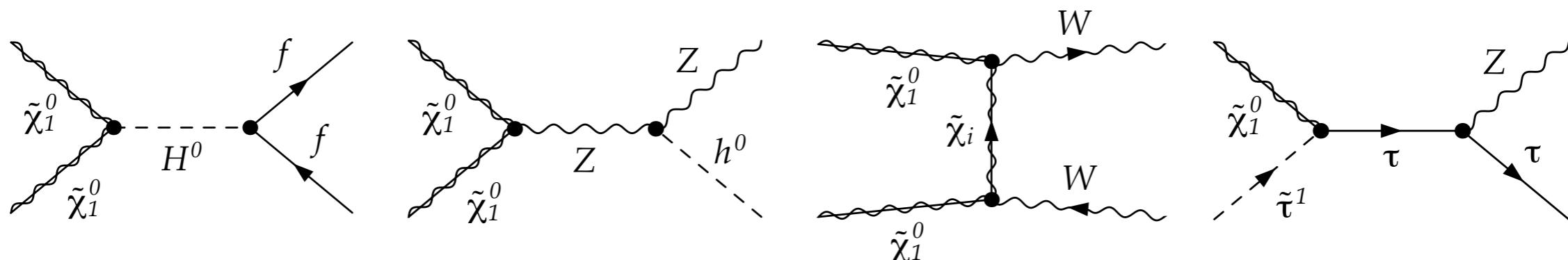
Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\text{ann}} v \rangle (n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}} h^2 \propto n \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle}$$

Cross section σ_{ann} includes all annihilation and coannihilation processes



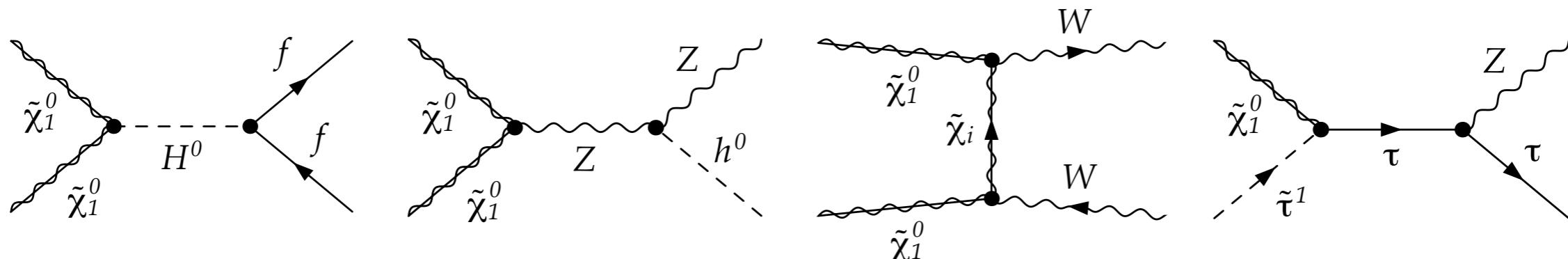
Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\text{ann}} v \rangle (n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}} h^2 \propto n \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle}$$

Cross section σ_{ann} includes all annihilation and coannihilation processes



Thermal average involves velocity distribution of the non-relativistic relic particle

$$\langle \sigma_{\text{ann}} v \rangle = \int dv \tilde{f}(v) \sigma_{\text{ann}} v = \int ds f(s) \sigma_{\text{ann}}(s)$$

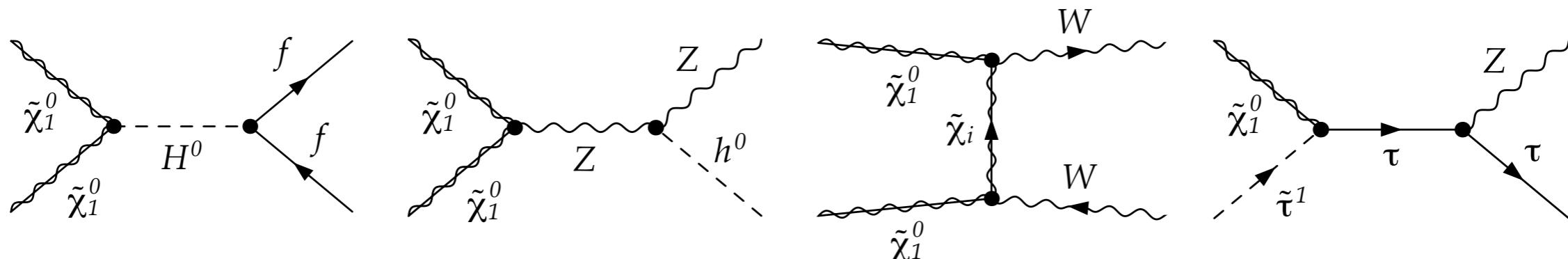
Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\text{ann}} v \rangle (n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}} h^2 \propto n \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle}$$

Cross section σ_{ann} includes all annihilation and coannihilation processes



Thermal average involves velocity distribution of the non-relativistic relic particle

$$\langle \sigma_{\text{ann}} v \rangle = \int dv \tilde{f}(v) \sigma_{\text{ann}} v = \int ds f(s) \sigma_{\text{ann}}(s)$$

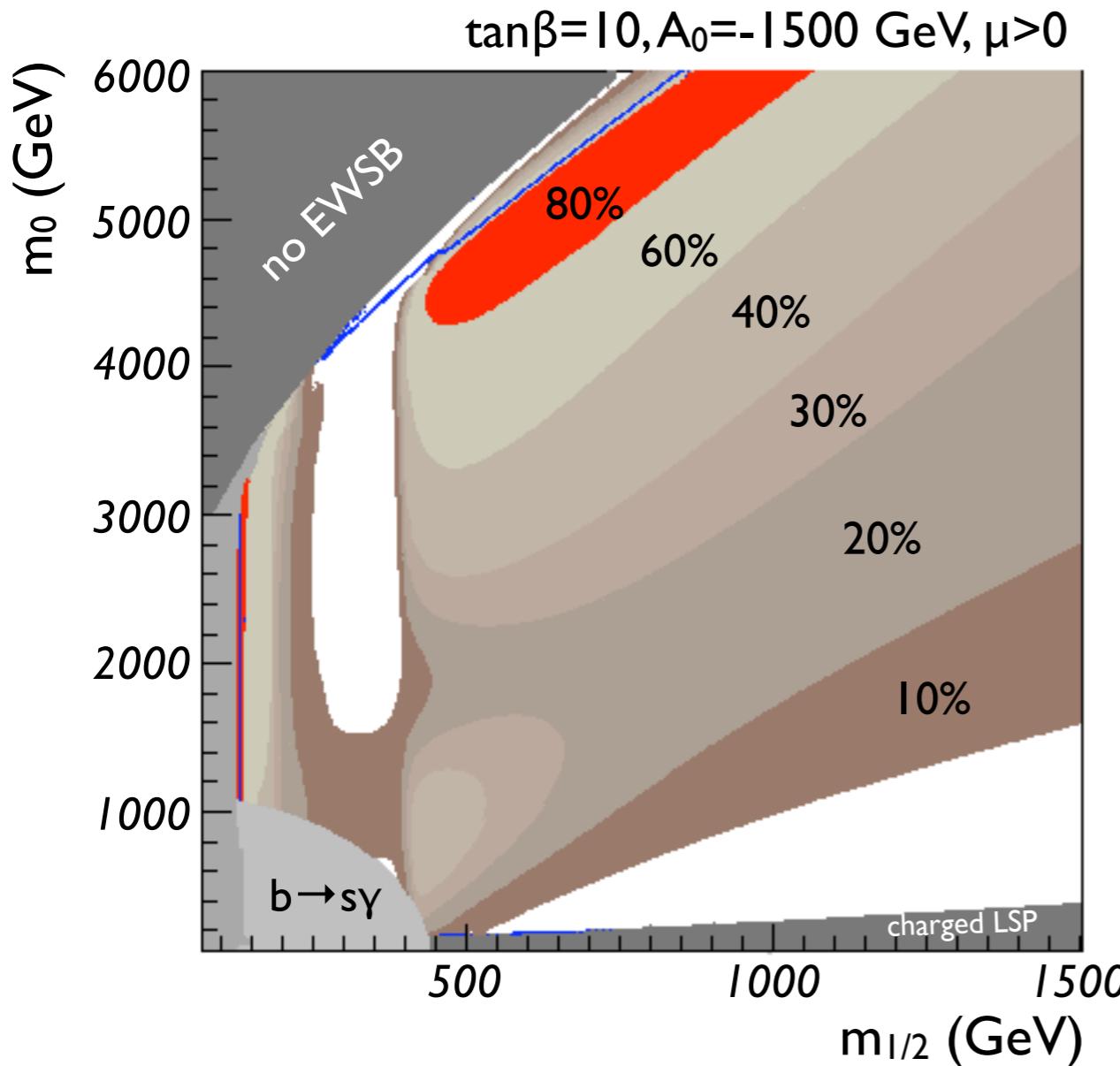
Public codes perform a calculation of the relic density for given scenario

- DarkSUSY (only neutralino)
- micrOMEGAs (all kinds of LSP)

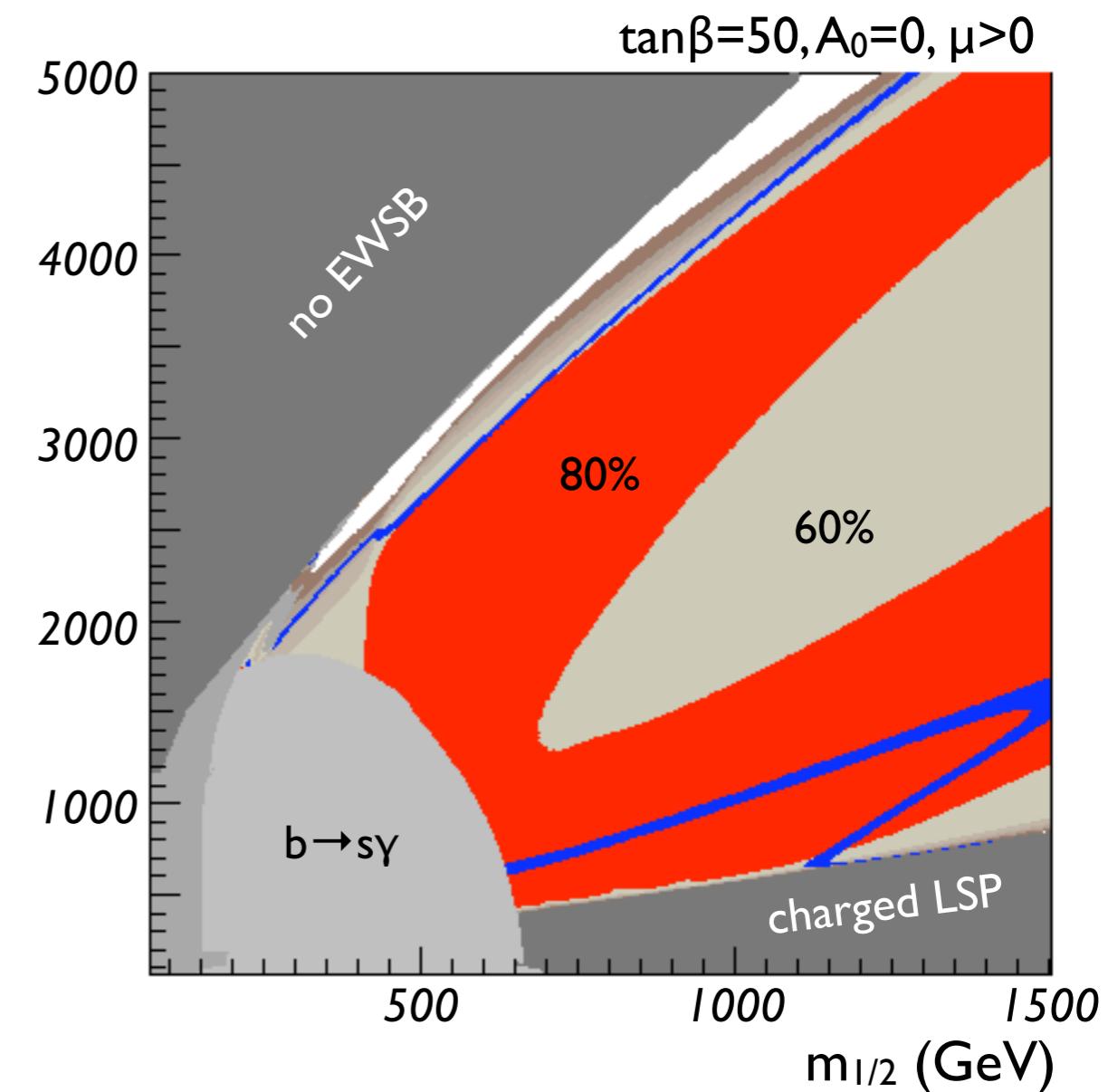
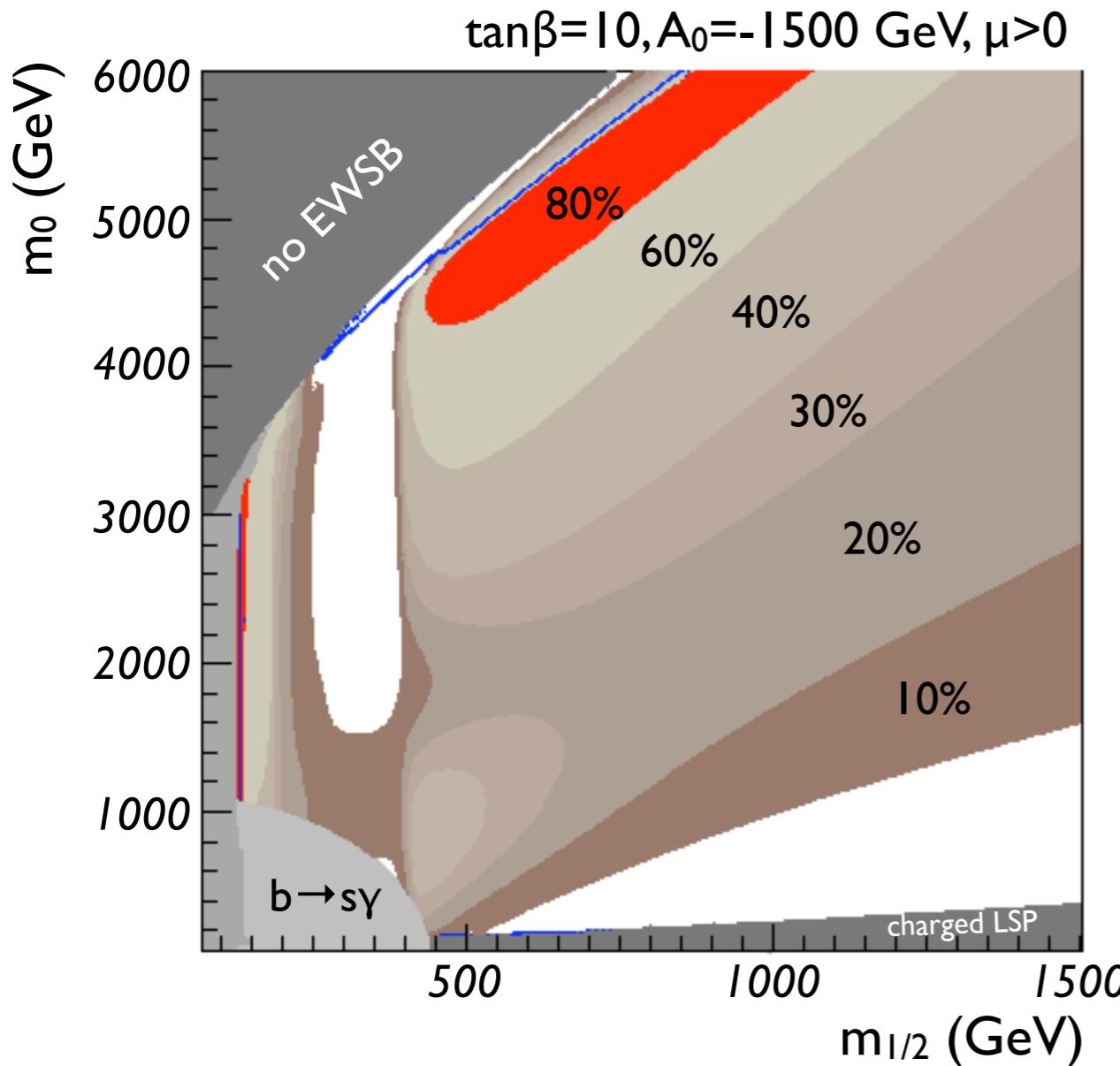
[Gondolo et al. 2004]

[Bélanger et al. 2006]

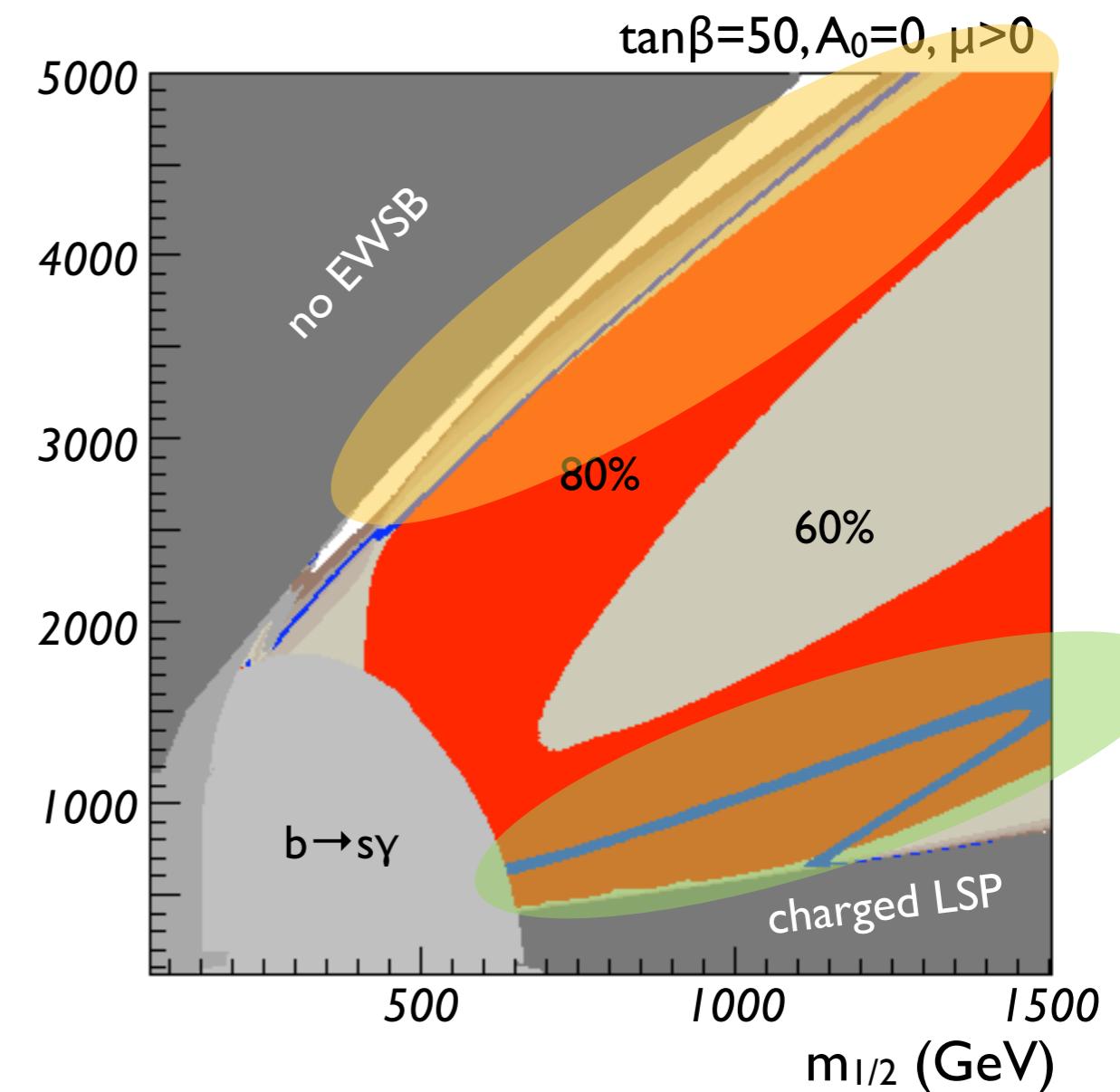
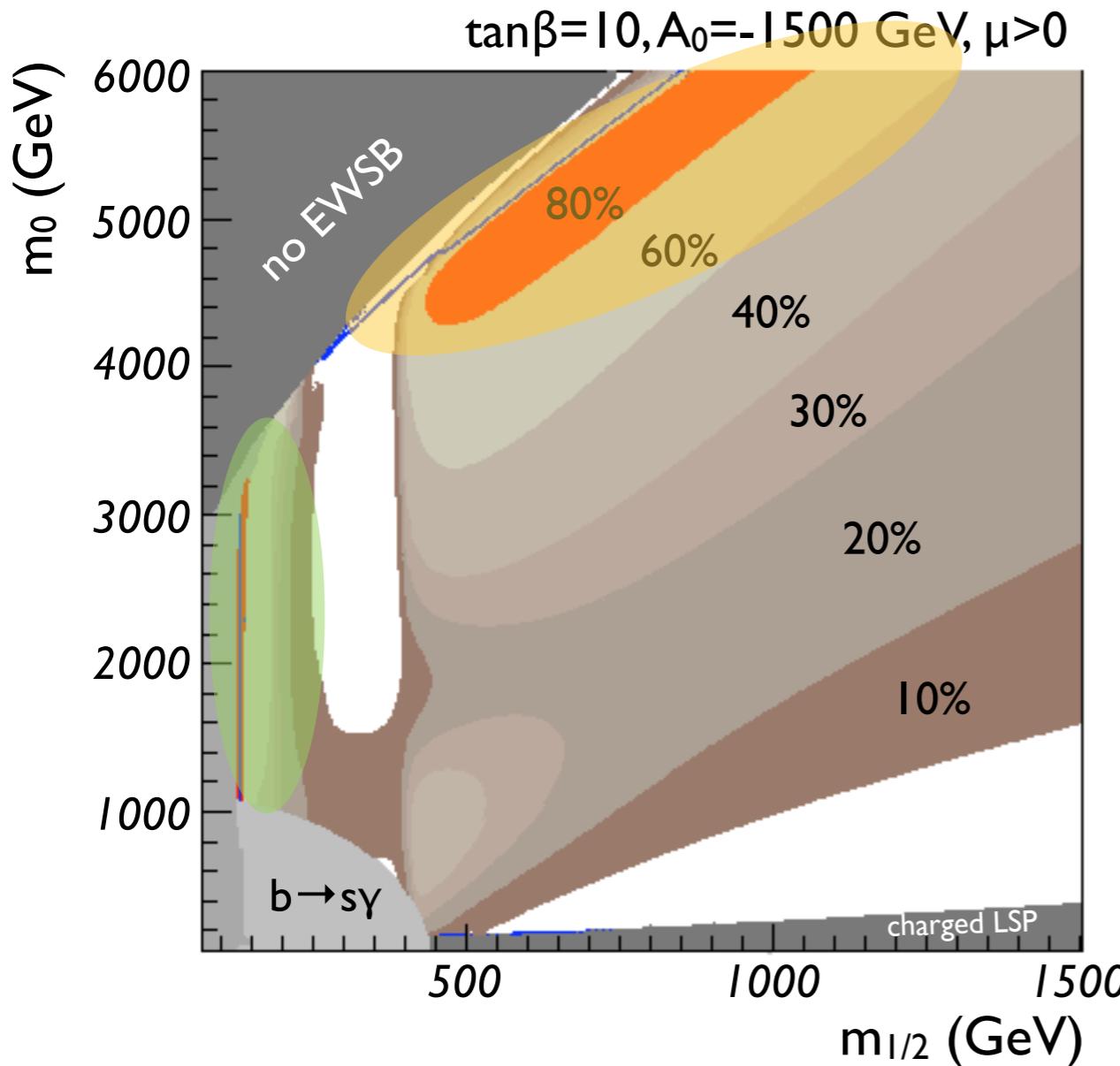
Neutralino Annihilation into Quarks



Neutralino Annihilation into Quarks



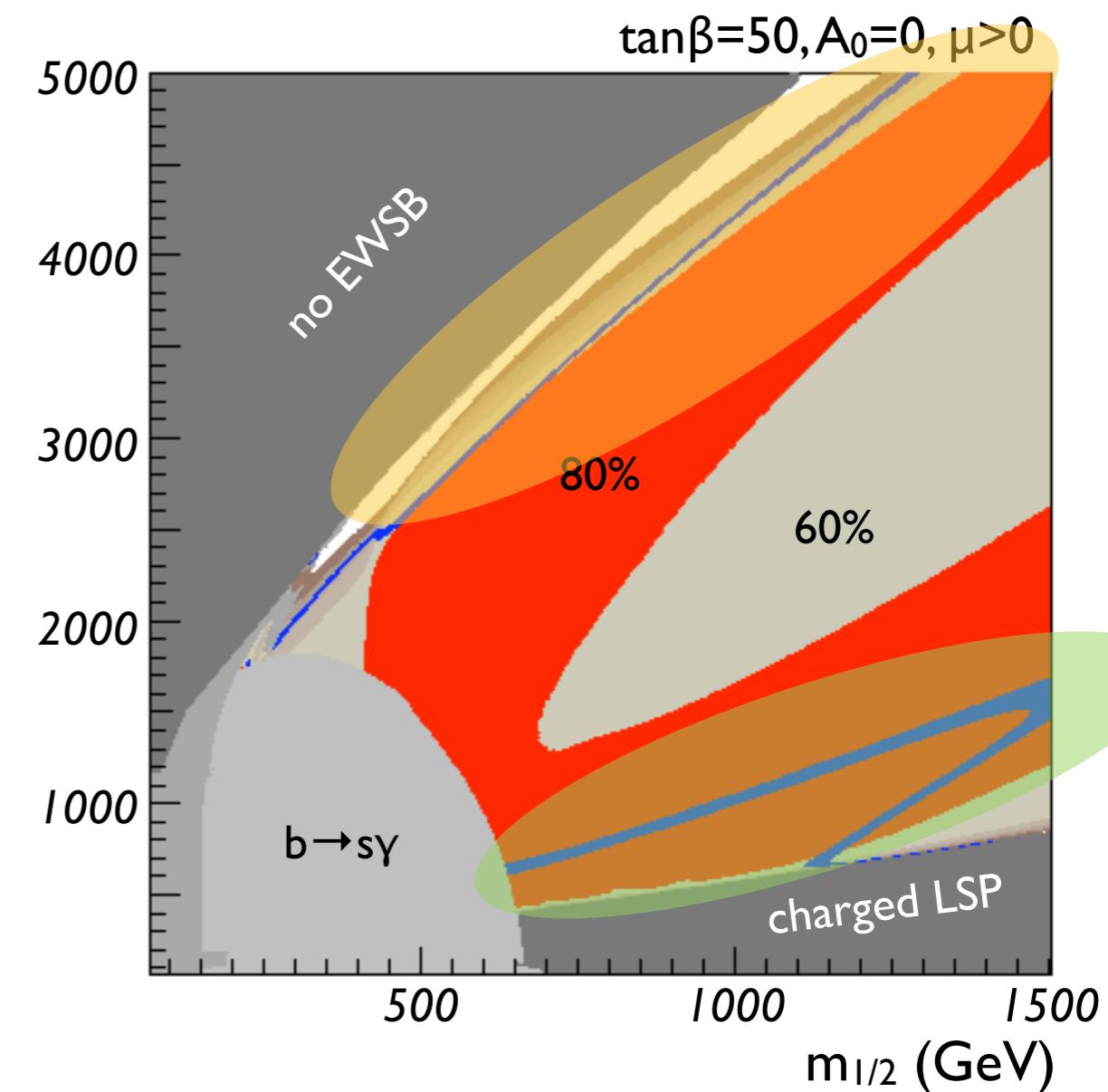
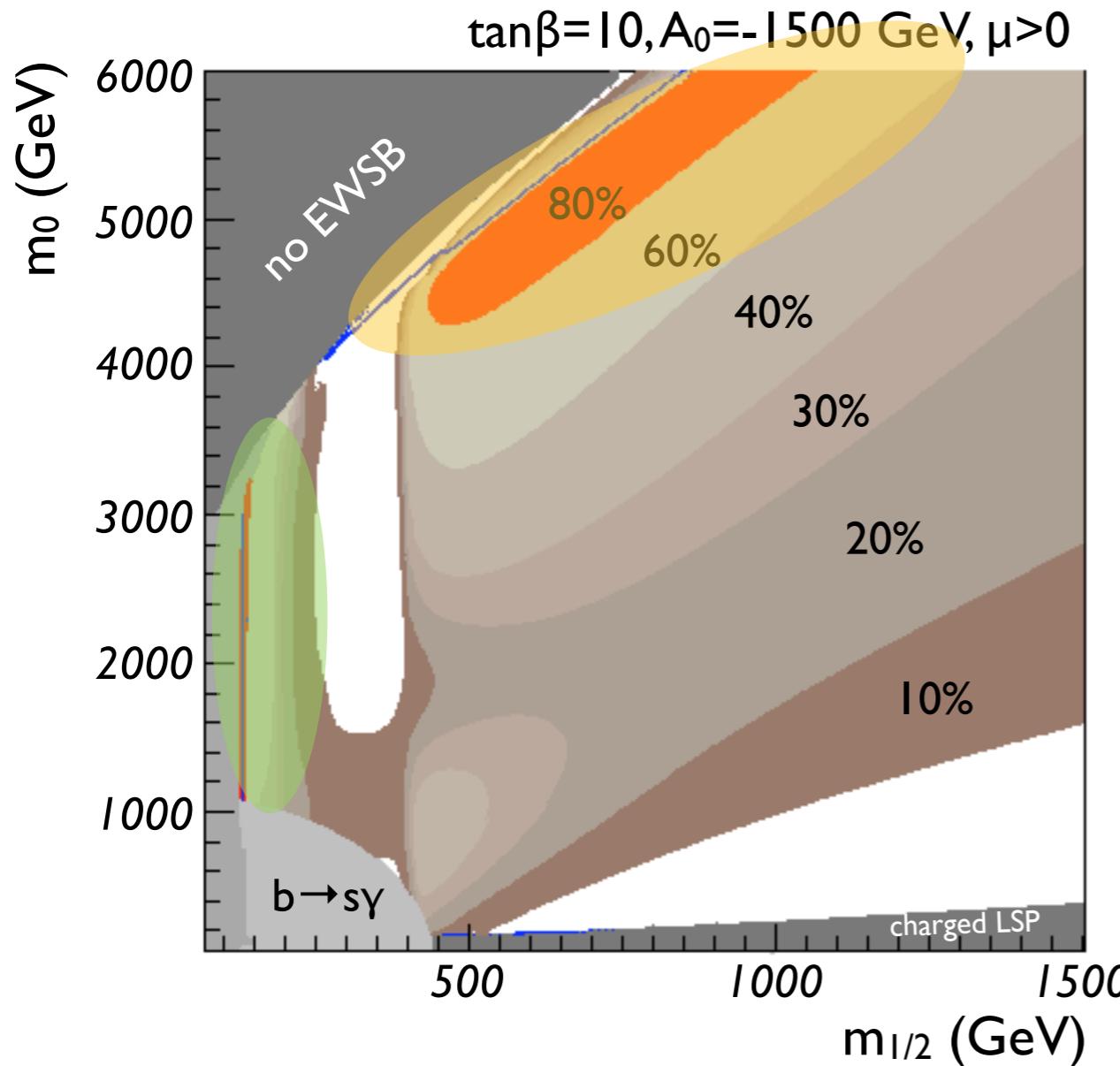
Neutralino Annihilation into Quarks



Interesting regions:

- Focus point region ($t\bar{t}$ dominated)
- Higgs resonances ($b\bar{b}$ dominated)

Neutralino Annihilation into Quarks



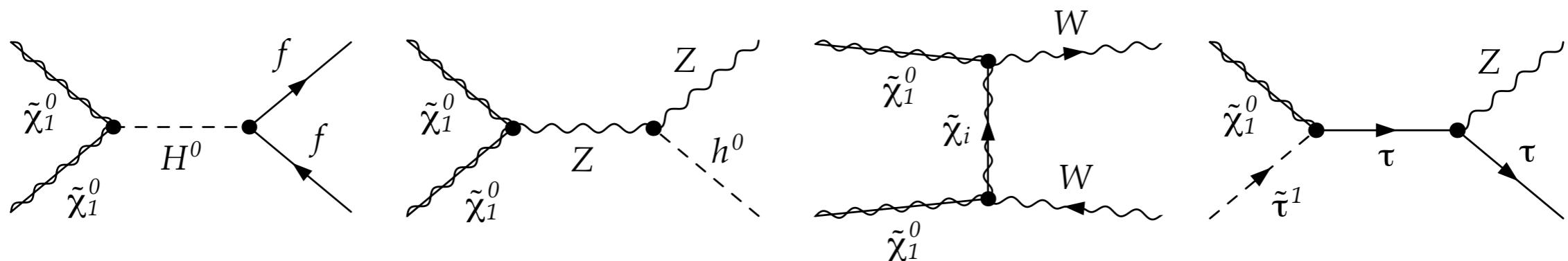
Interesting regions:

- Focus point region ($t\bar{t}$ dominated)
- Higgs resonances ($b\bar{b}$ dominated)

Light quark final states do not play a significant role in mSUGRA scenarios

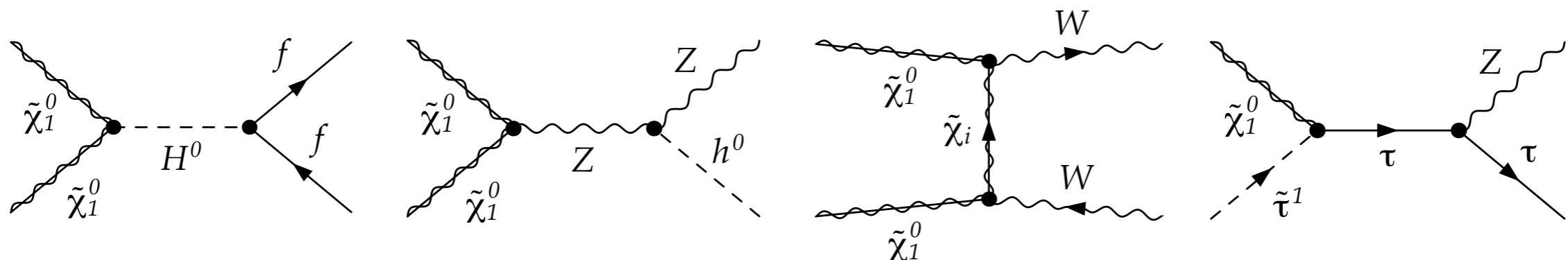
Motivation for SUSY-QCD Corrections

Most (co)annihilation processes only implemented at leading order in public codes



Motivation for SUSY-QCD Corrections

Most (co)annihilation processes only implemented at leading order in public codes

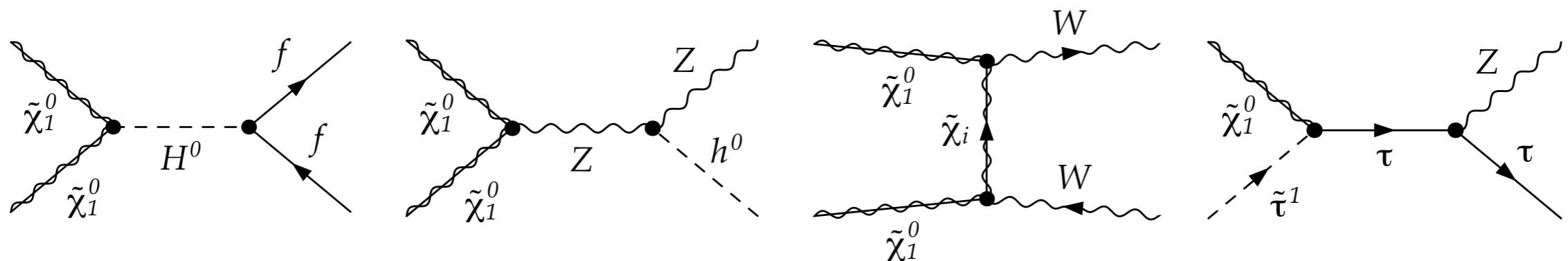


Higher order corrections can have important contributions to cross sections

- QCD corrections significant due to strong coupling constant
- modification of preferred regions in parameter space

Motivation for SUSY-QCD Corrections

Most (co)annihilation processes only implemented at leading order in public codes

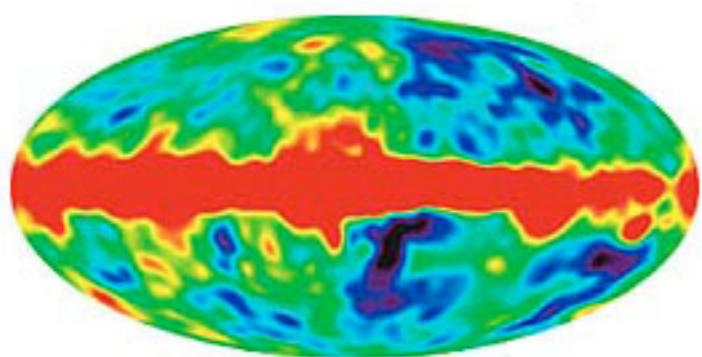


Higher order corrections can have important contributions to cross sections

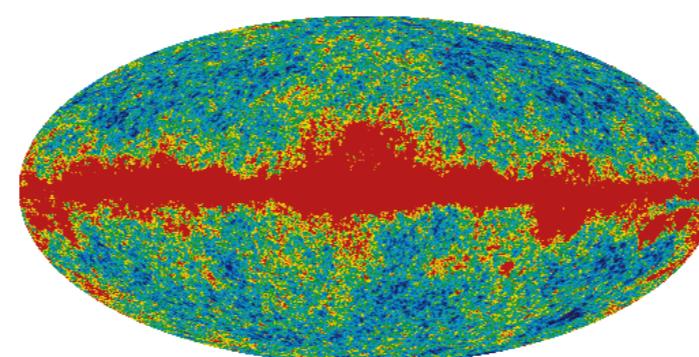
- QCD corrections significant due to strong coupling constant
- modification of preferred regions in parameter space

Planck satellite will deliver new cosmological data in near future

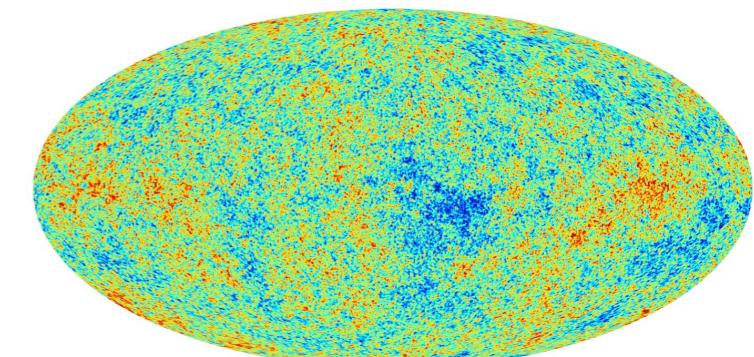
- higher precision in theoretical predictions needed to match experimental improvements



COBE
1989



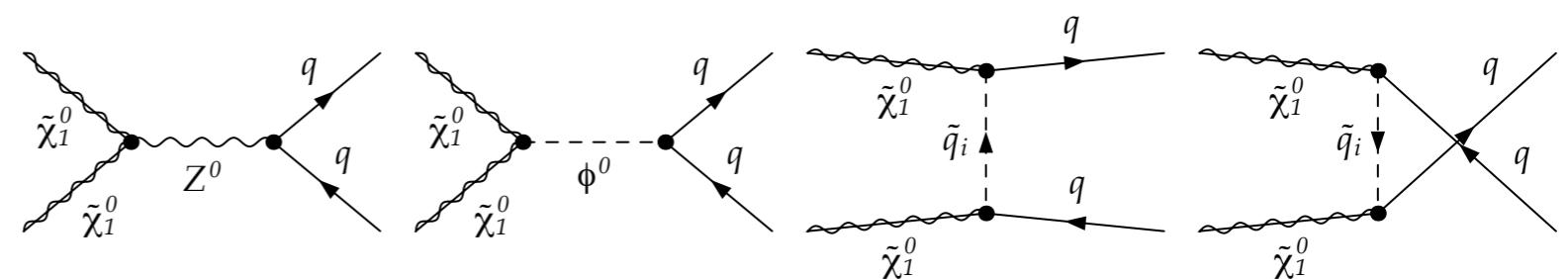
WMAP
2002



Planck
2009 ?

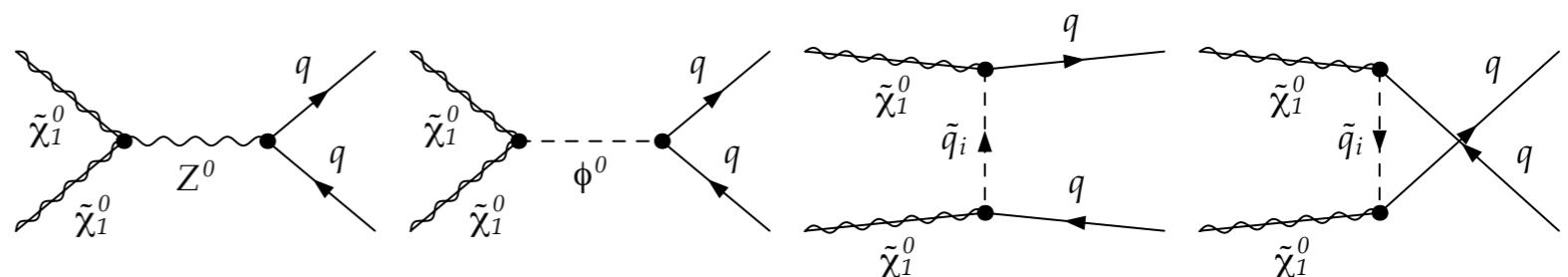
QCD and SUSY-QCD Corrections

Tree-level Feynman diagrams



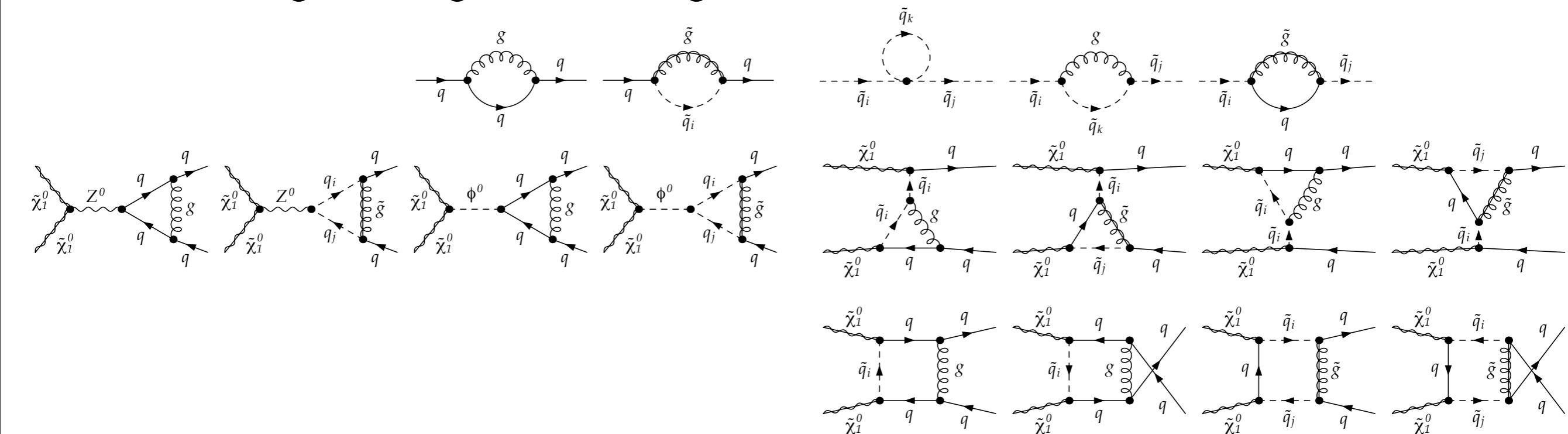
QCD and SUSY-QCD Corrections

Tree-level Feynman diagrams



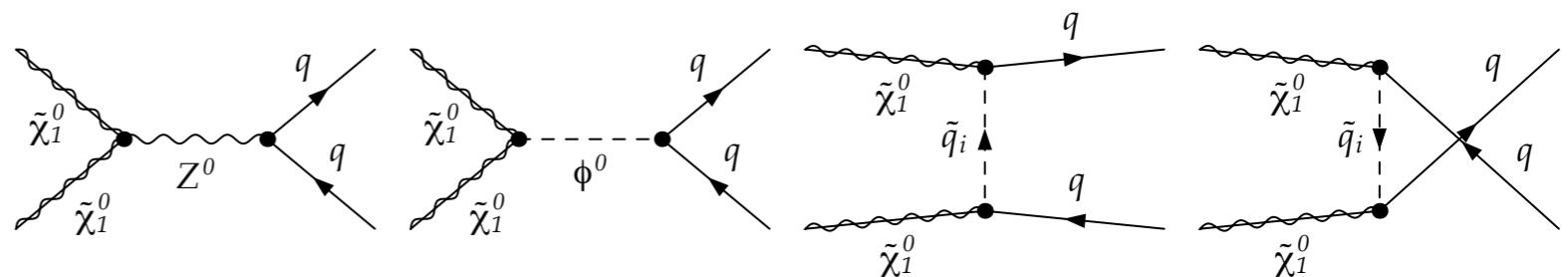
Virtual loop corrections

- Loop integrals calculated in $\overline{\text{DR}}$ regularization scheme (preserves Supersymmetry)
- UV-divergencies regularized through on-shell renormalization



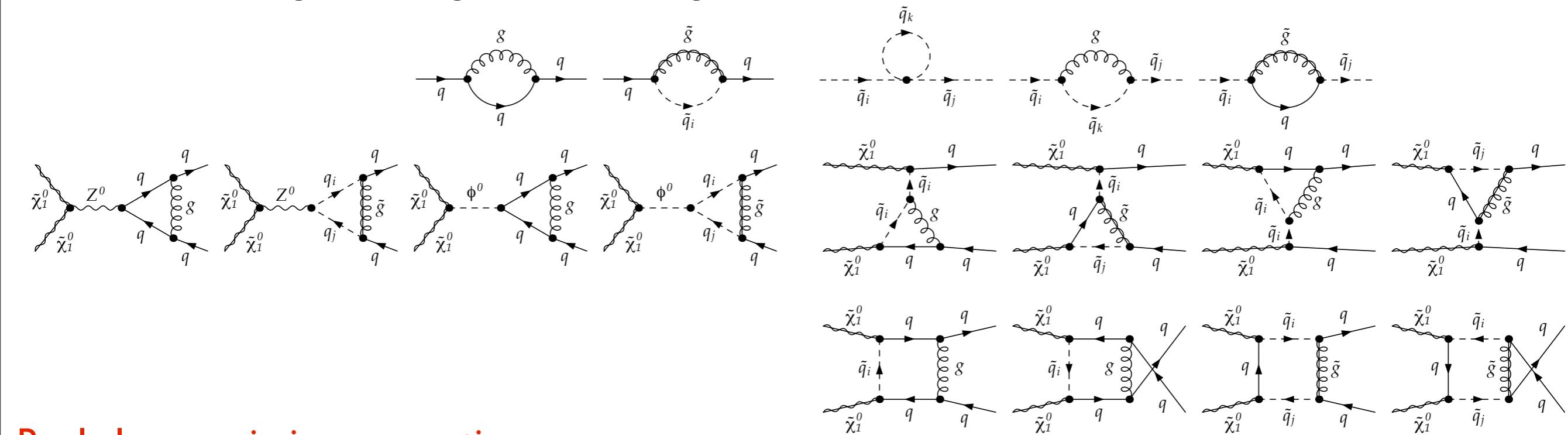
QCD and SUSY-QCD Corrections

Tree-level Feynman diagrams



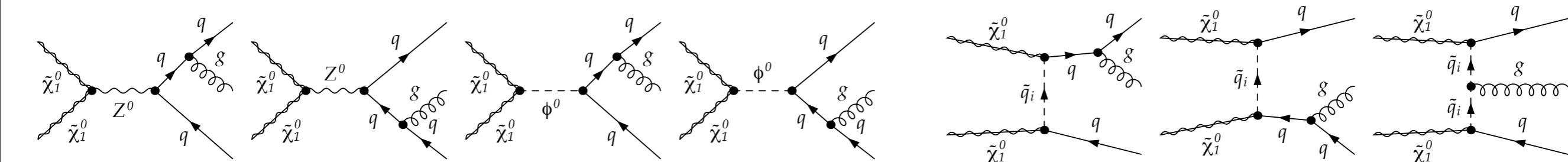
Virtual loop corrections

- Loop integrals calculated in $\overline{\text{DR}}$ regularization scheme (preserves Supersymmetry)
- UV-divergencies regularized through on-shell renormalization



Real gluon emission corrections

- IR-divergencies calculated in $\overline{\text{DR}}$ regularization scheme



QCD and SUSY-QCD Corrections

Dipole subtraction method to combine virtual and real emission contributions

[Catani et al. 2000]

→ auxiliary cross section cancels IR-divergence in both the virtual and real part

$$\sigma_{\text{NLO}} = \left[\sigma_V + \int d\sigma_{\text{aux}} \right]_{\epsilon=0} + \int \left[d\sigma_R - d\sigma_{\text{aux}} \right]_{\epsilon=0}$$

QCD and SUSY-QCD Corrections

Dipole subtraction method to combine virtual and real emission contributions

[Catani et al. 2000]

→ auxiliary cross section cancels IR-divergence in both the virtual and real part

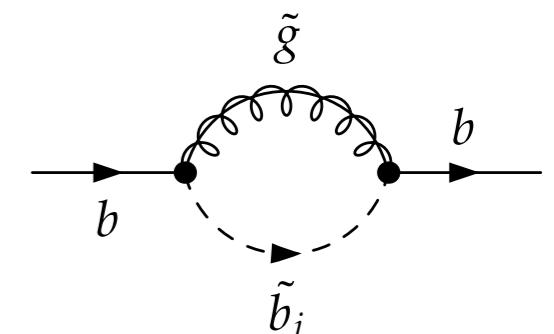
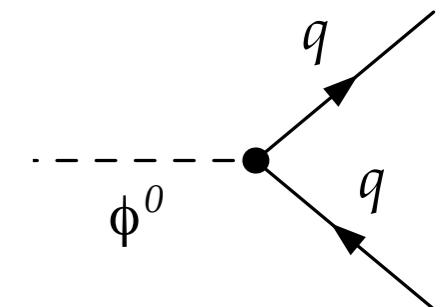
$$\sigma_{\text{NLO}} = \left[\sigma_V + \int d\sigma_{\text{aux}} \right]_{\epsilon=0} + \int \left[d\sigma_R - d\sigma_{\text{aux}} \right]_{\epsilon=0}$$

Yukawa couplings in Higgs exchange diagrams

- Higgs resonances enhance annihilation cross section significantly
- radiative corrections to Higgs boson decay into fermions well known
- include QCD corrections up to $\mathcal{O}(\alpha_s^4)$ [Chetyrkin et al. 1995-2006]
- include SUSY-QCD bottom quark mass resummation (important at large $\tan\beta$ or A_b)

[Carena et al. 2000, Spira et al. 2003]

$$\left(\frac{\Delta m_b}{m_b} \right)_{\tilde{g}\tilde{b}} = \frac{\alpha_s(s)}{\pi} C_F \frac{m_{\tilde{g}}}{2} \left(A_b - \mu \tan \beta \right) I(m_{\tilde{b}_1}^2, m_{\tilde{b}_2}^2, m_{\tilde{g}}^2)$$



QCD and SUSY-QCD Corrections

Dipole subtraction method to combine virtual and real emission contributions

[Catani et al. 2000]

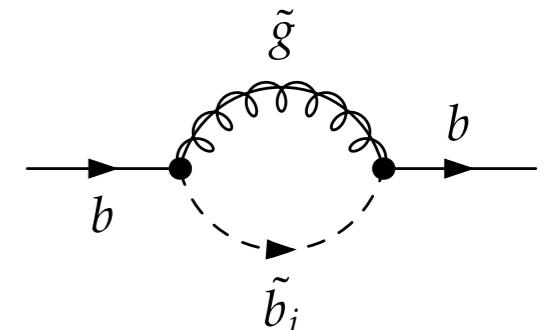
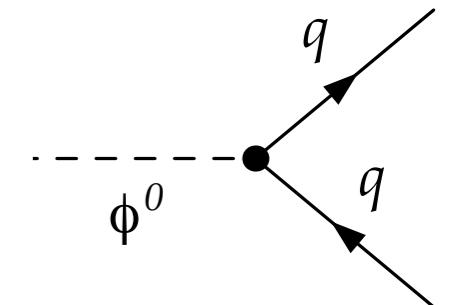
→ auxiliary cross section cancels IR-divergence in both the virtual and real part

$$\sigma_{\text{NLO}} = \left[\sigma_V + \int d\sigma_{\text{aux}} \right]_{\epsilon=0} + \int \left[d\sigma_R - d\sigma_{\text{aux}} \right]_{\epsilon=0}$$

Yukawa couplings in Higgs exchange diagrams

- Higgs resonances enhance annihilation cross section significantly
- radiative corrections to Higgs boson decay into fermions well known
- include QCD corrections up to $\mathcal{O}(\alpha_s^4)$ [Chetyrkin et al. 1995-2006]
- include SUSY-QCD bottom quark mass resummation (important at large $\tan\beta$ or A_b)
[Carena et al. 2000, Spira et al. 2003]

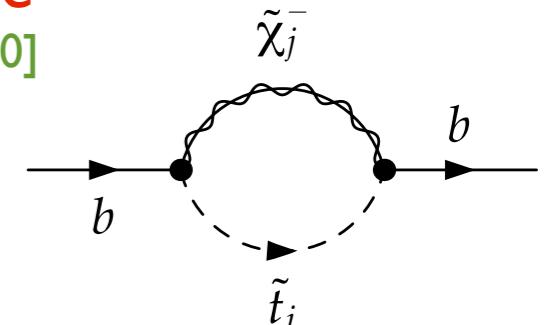
$$\left(\frac{\Delta m_b}{m_b} \right)_{\tilde{g}\tilde{b}} = \frac{\alpha_s(s)}{\pi} C_F \frac{m_{\tilde{g}}}{2} \left(A_b - \mu \tan \beta \right) I(m_{\tilde{b}_1}^2, m_{\tilde{b}_2}^2, m_{\tilde{g}}^2)$$



Stop-chargino loop correction to bottom mass of same order of magnitude

[Carena et al. 2000]

$$\left(\frac{\Delta m_b}{m_b} \right)_{\tilde{t}\tilde{\chi}^\pm} = - \frac{\lambda_t^2}{16\pi^2} \left[A_t \mu \tan \beta - \mu^2 \right] I(\mu^2, m_{\tilde{t}_1}^2, m_{\tilde{t}_2}^2)$$

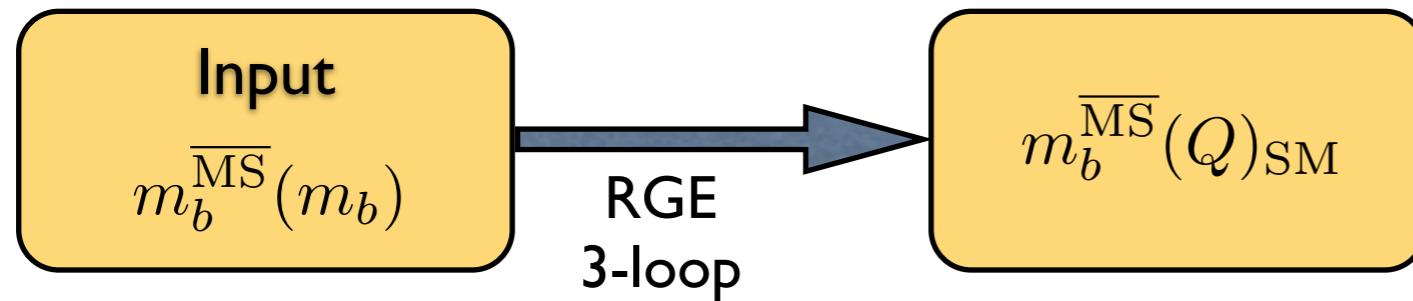


Treatment of Quark Masses

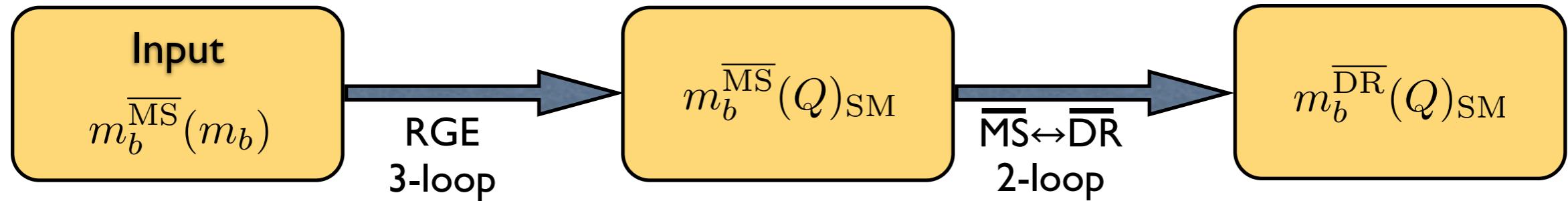
Input

$$m_b^{\overline{\text{MS}}}(m_b)$$

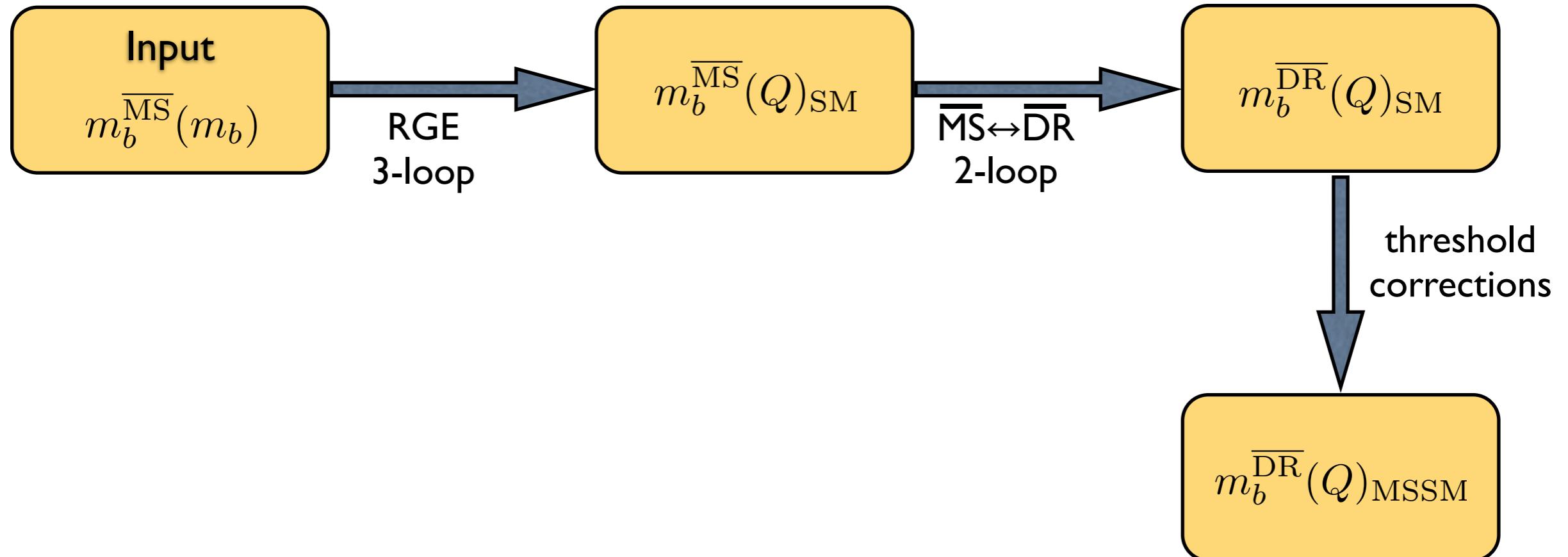
Treatment of Quark Masses



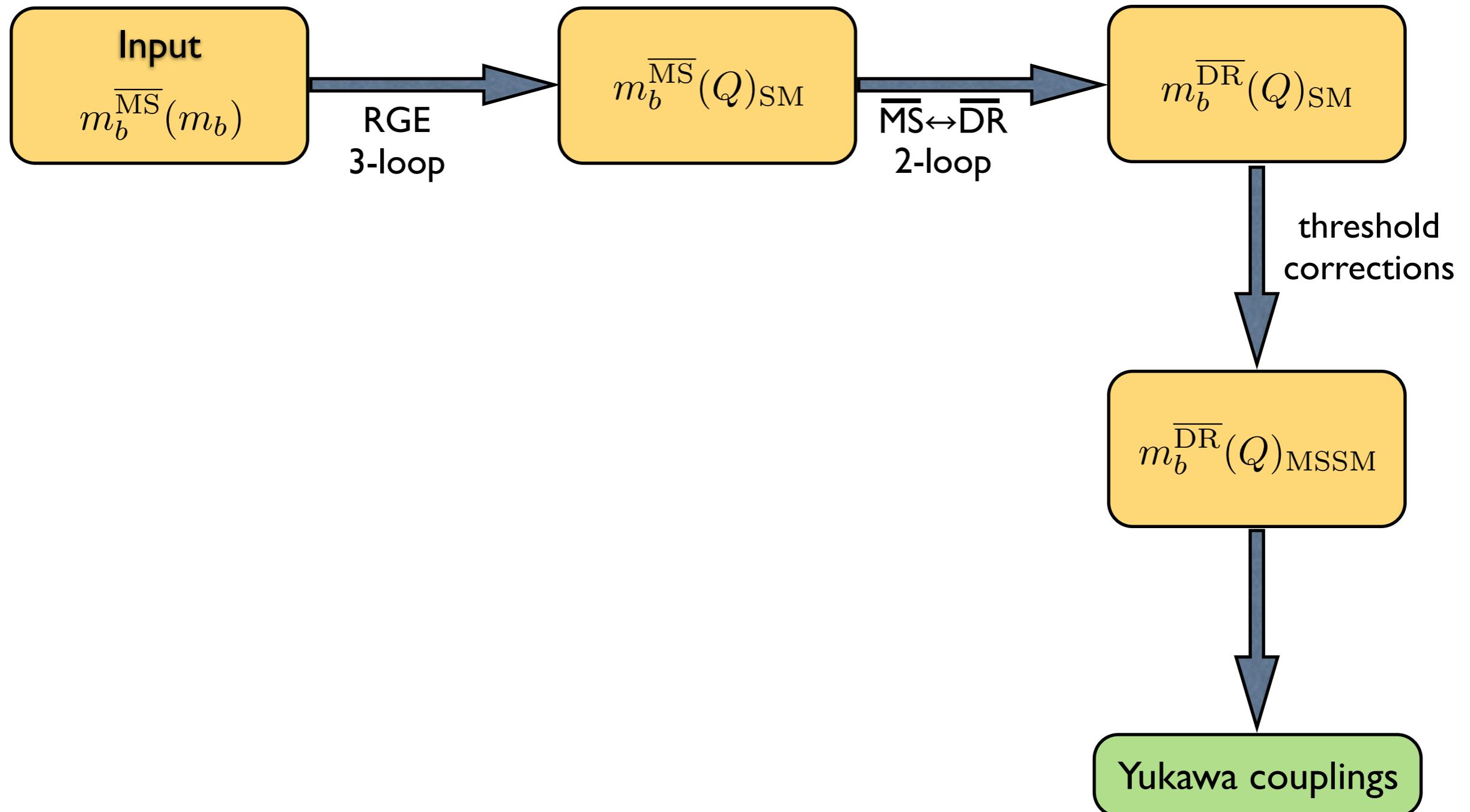
Treatment of Quark Masses



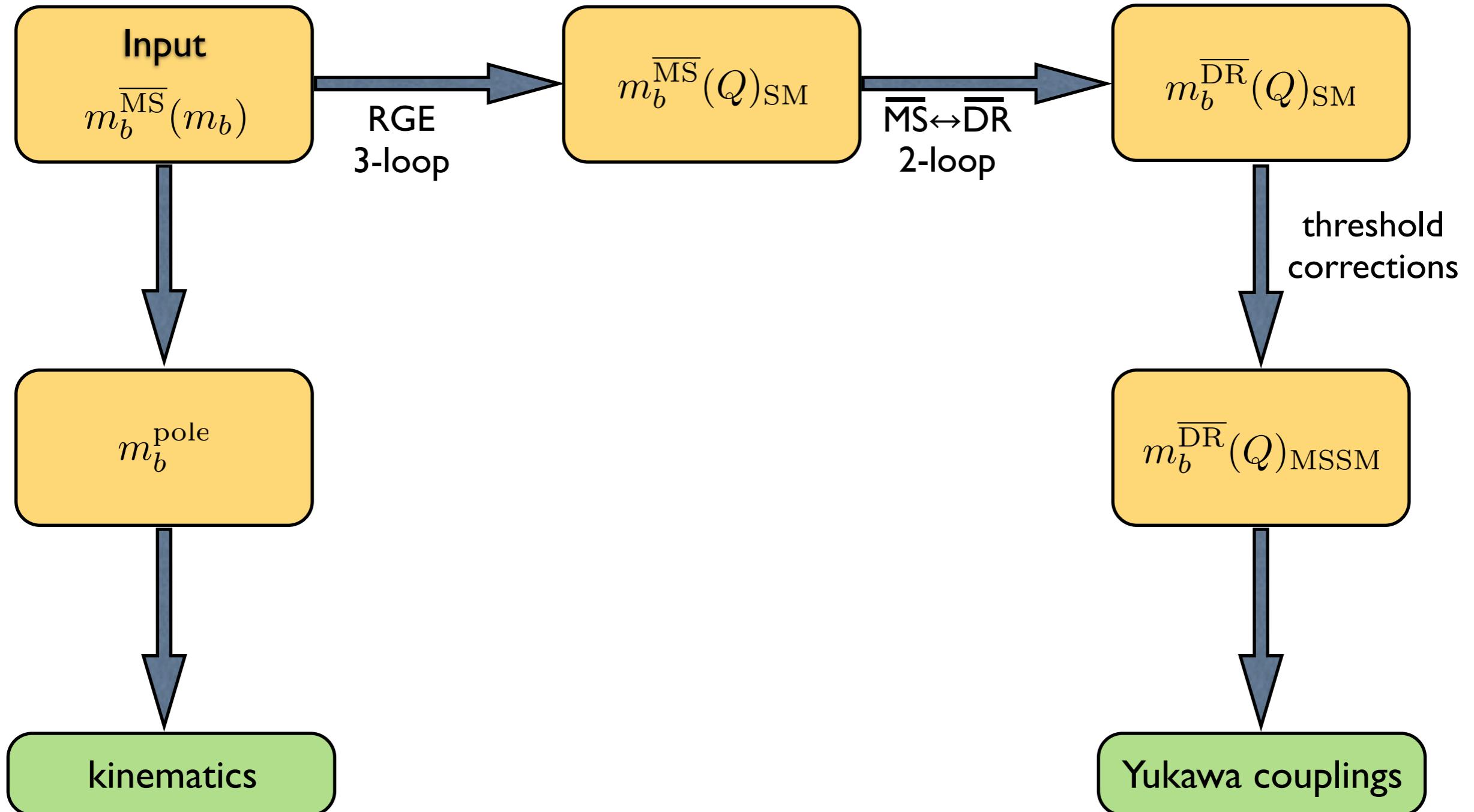
Treatment of Quark Masses



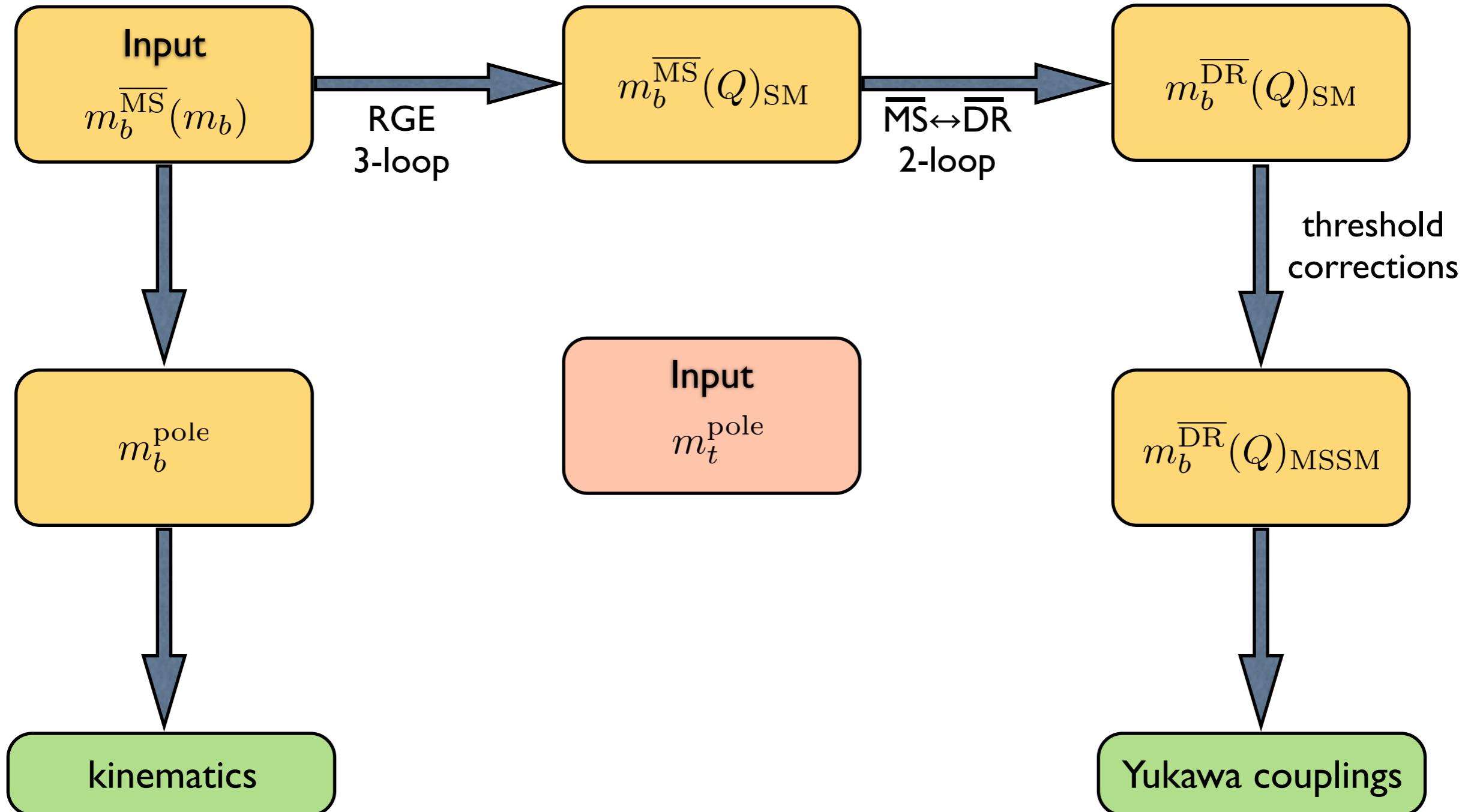
Treatment of Quark Masses



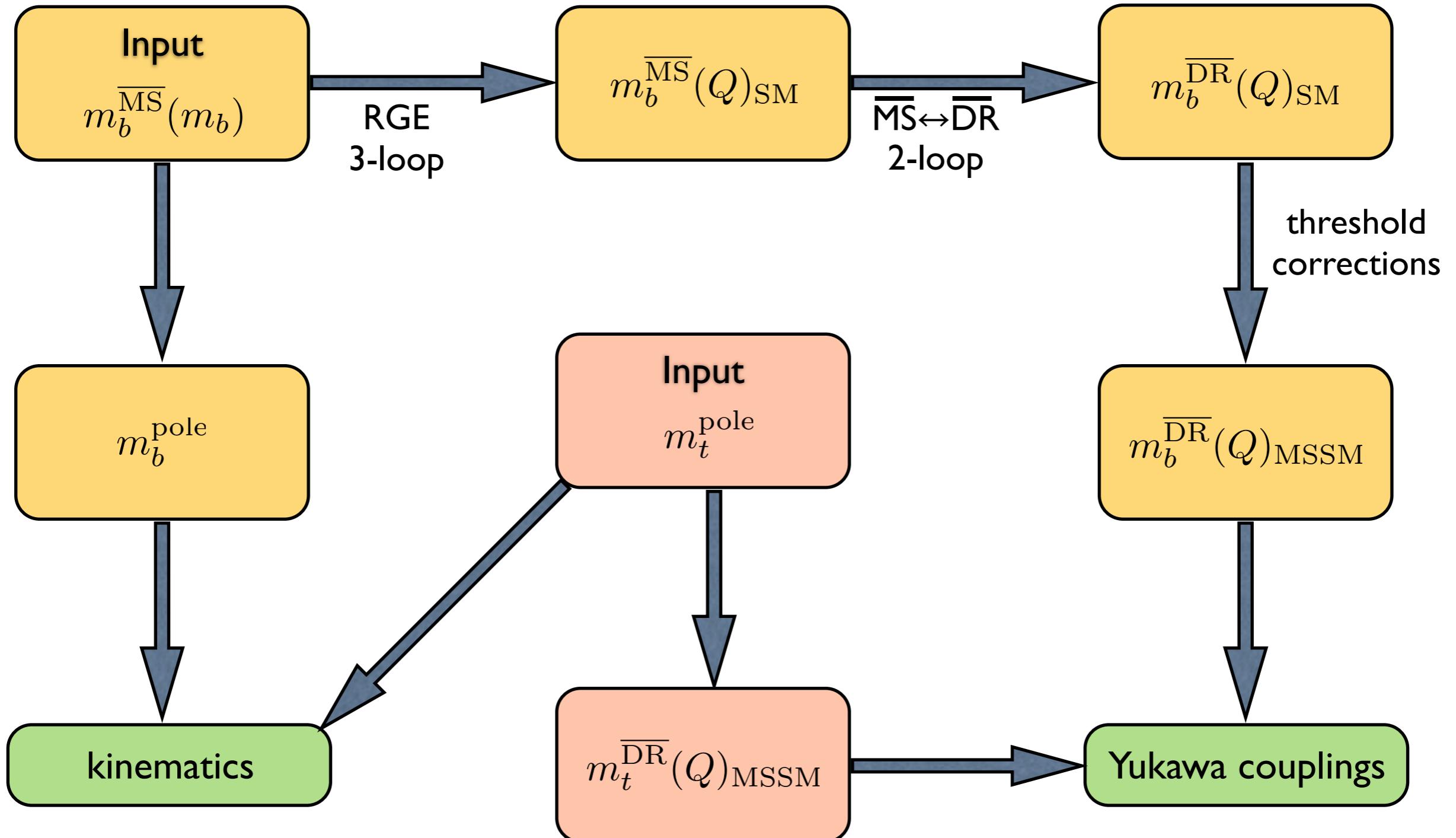
Treatment of Quark Masses



Treatment of Quark Masses

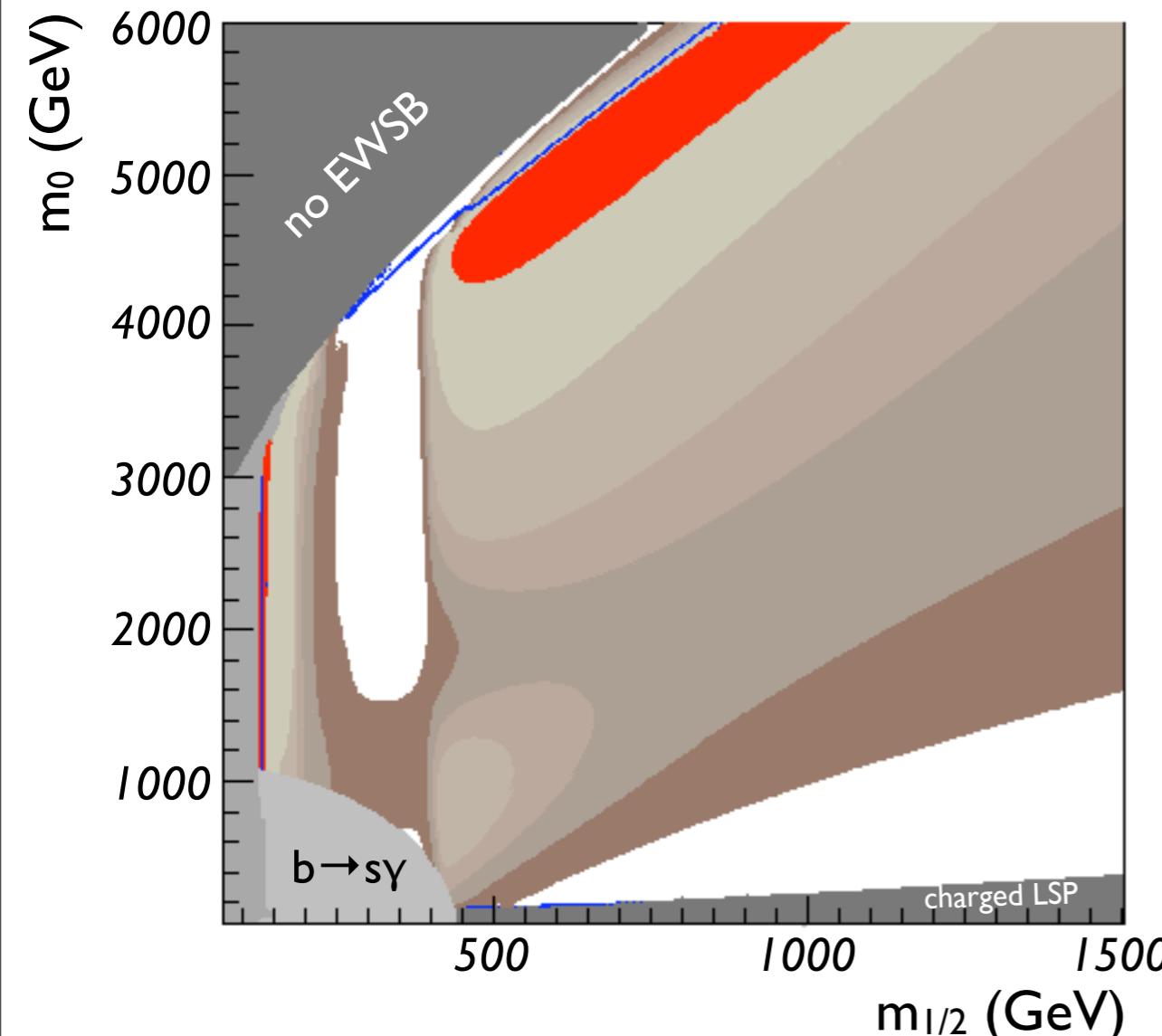


Treatment of Quark Masses



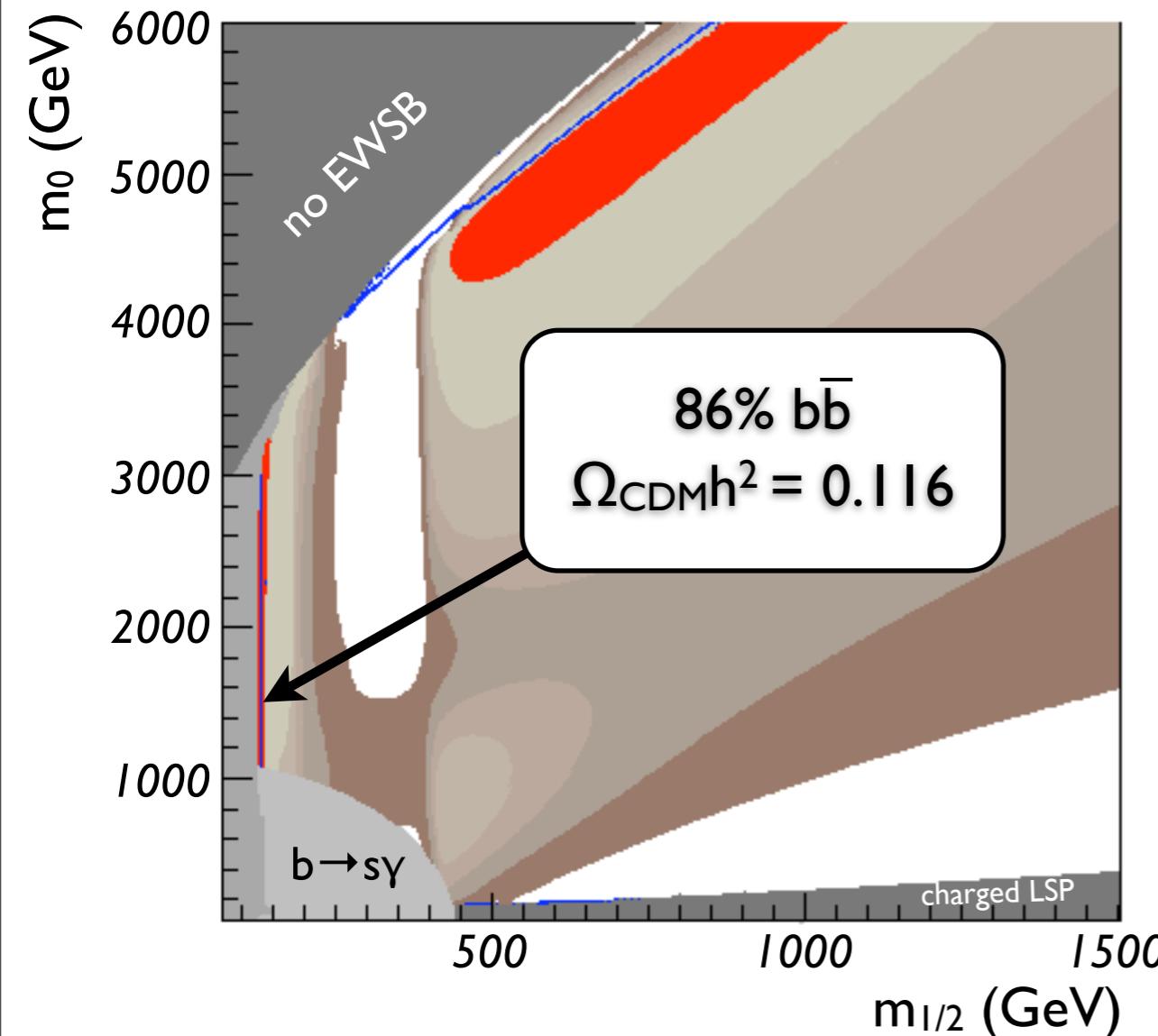
Numerical results: Light Higgs resonance (preliminary)

mSUGRA parameters: $m_0 = 1500$ GeV, $m_{1/2} = 130$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



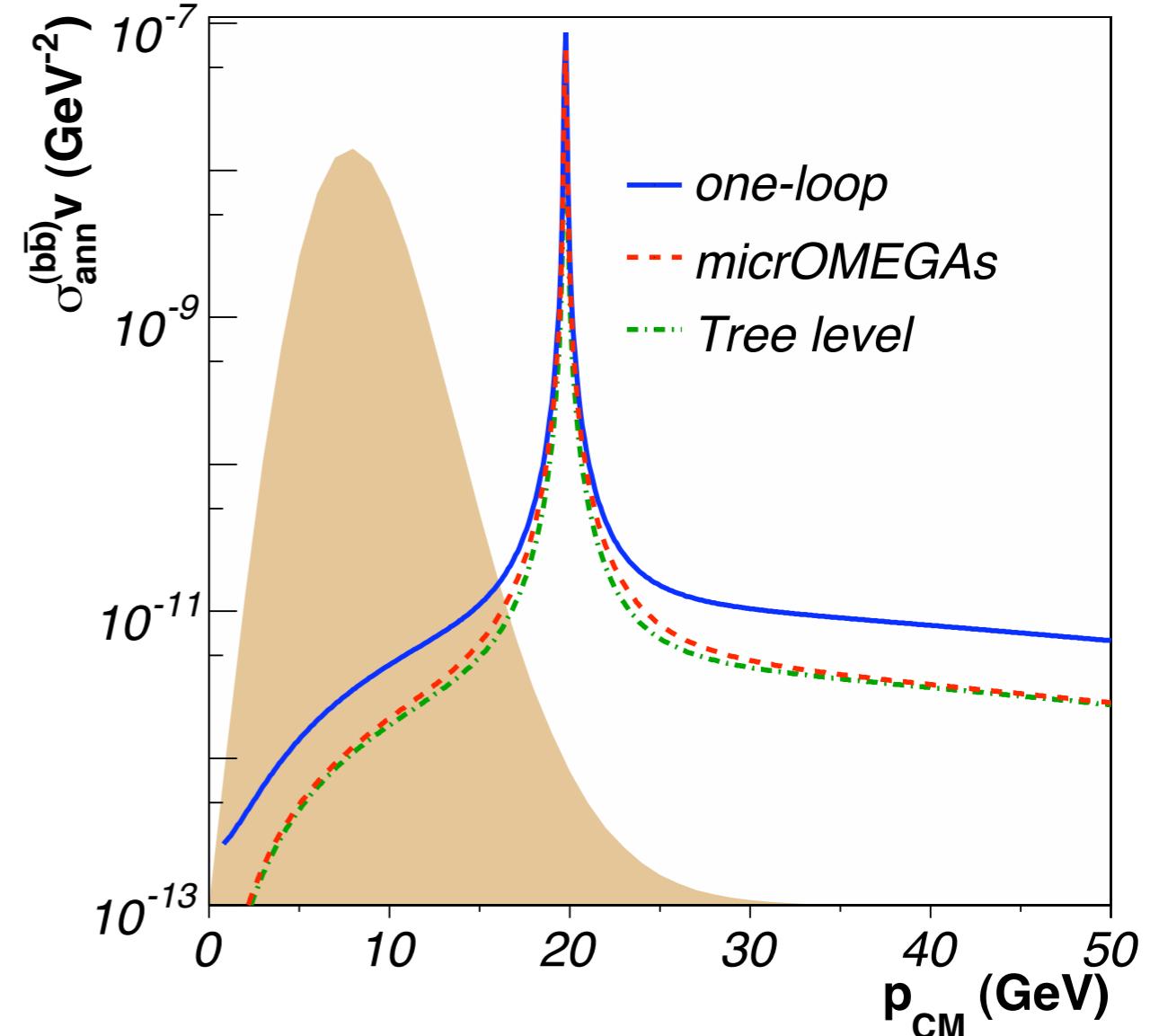
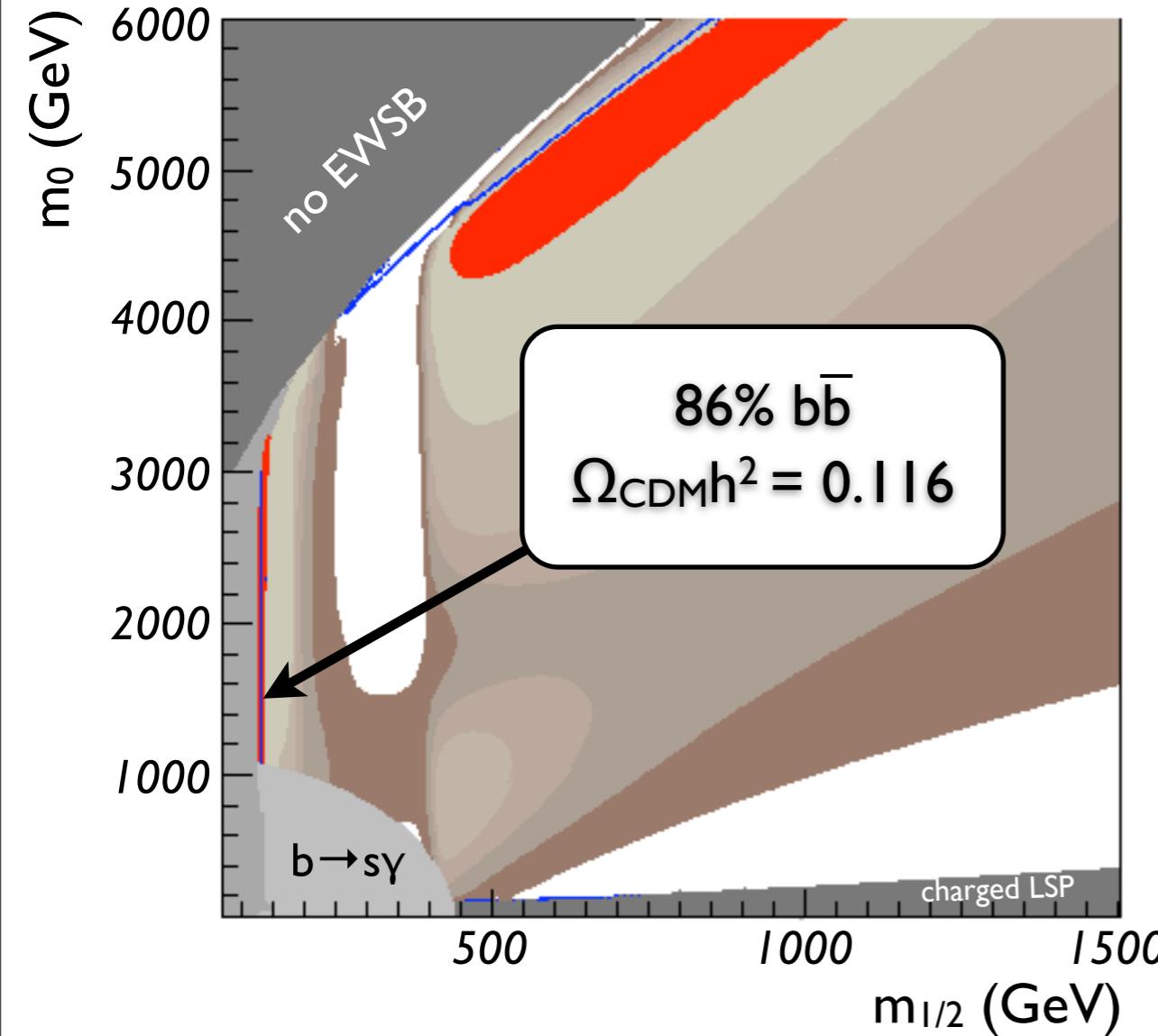
Numerical results: Light Higgs resonance (preliminary)

mSUGRA parameters: $m_0 = 1500$ GeV, $m_{1/2} = 130$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



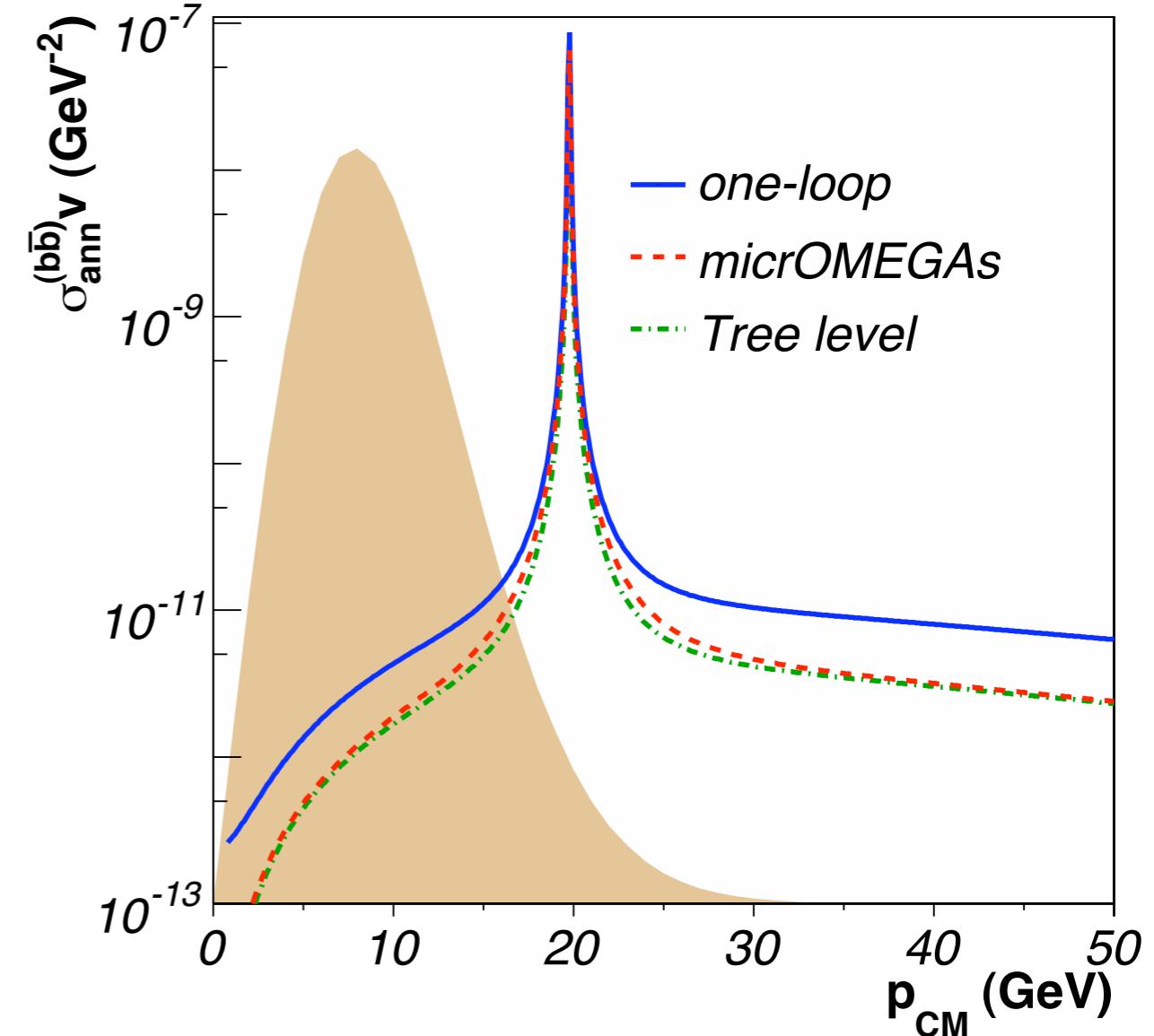
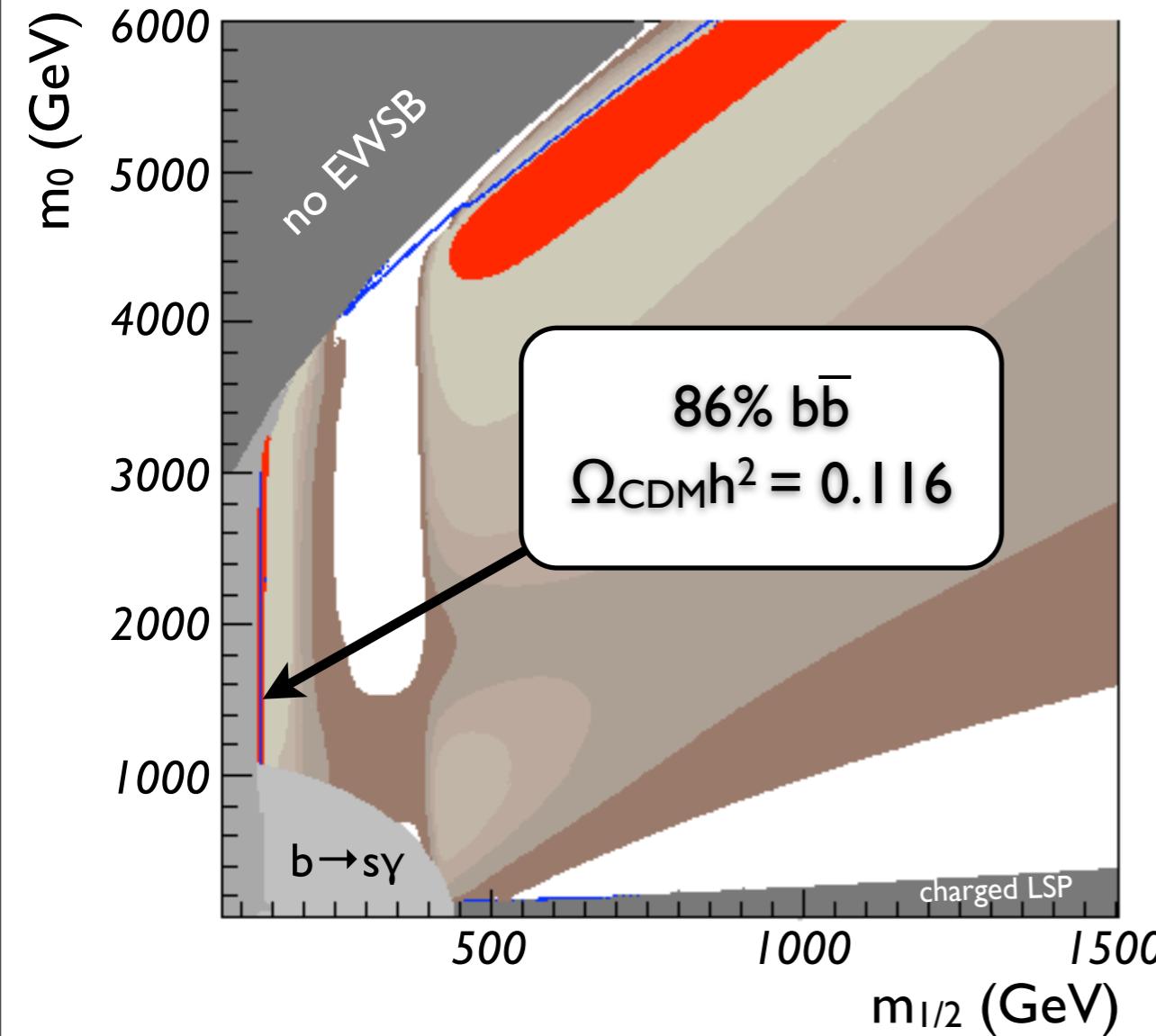
Numerical results: Light Higgs resonance (preliminary)

mSUGRA parameters: $m_0 = 1500$ GeV, $m_{1/2} = 130$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



Numerical results: Light Higgs resonance (preliminary)

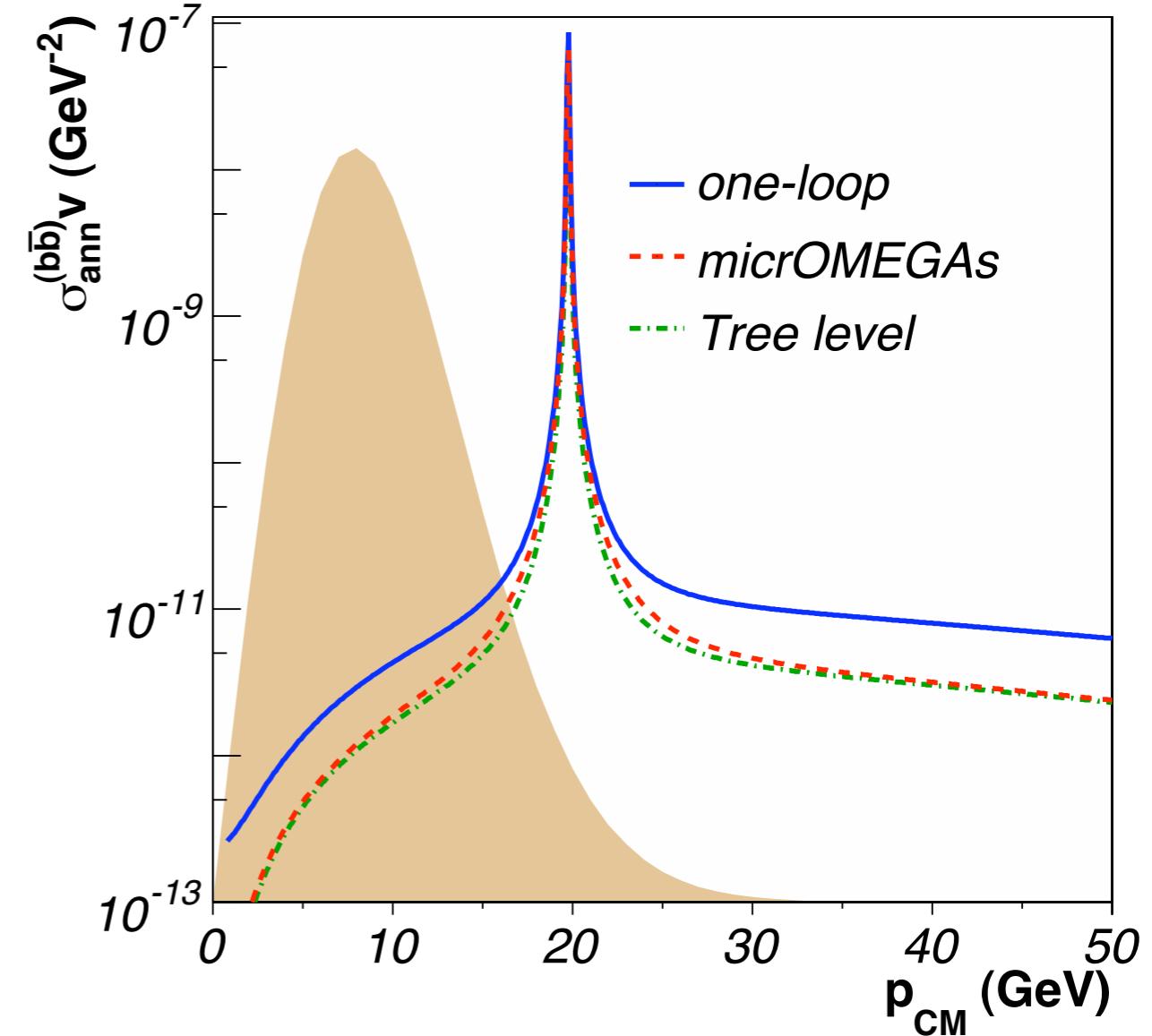
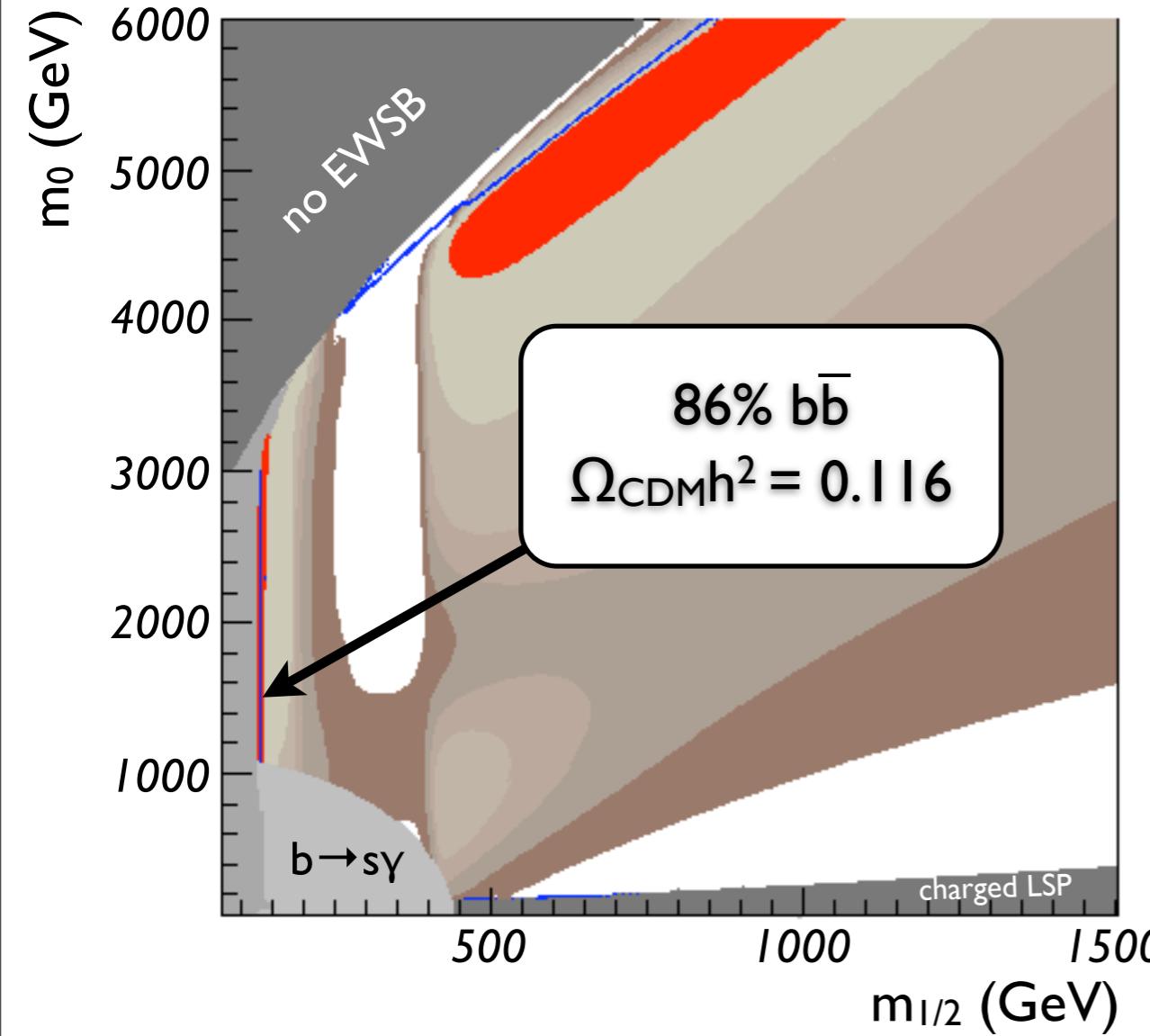
mSUGRA parameters: $m_0 = 1500$ GeV, $m_{1/2} = 130$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



“Effective” Yukawa coupling very good approximation around the Higgs resonance
Position of resonance does not coincide with peak of velocity distribution for thermal average

Numerical results: Light Higgs resonance (preliminary)

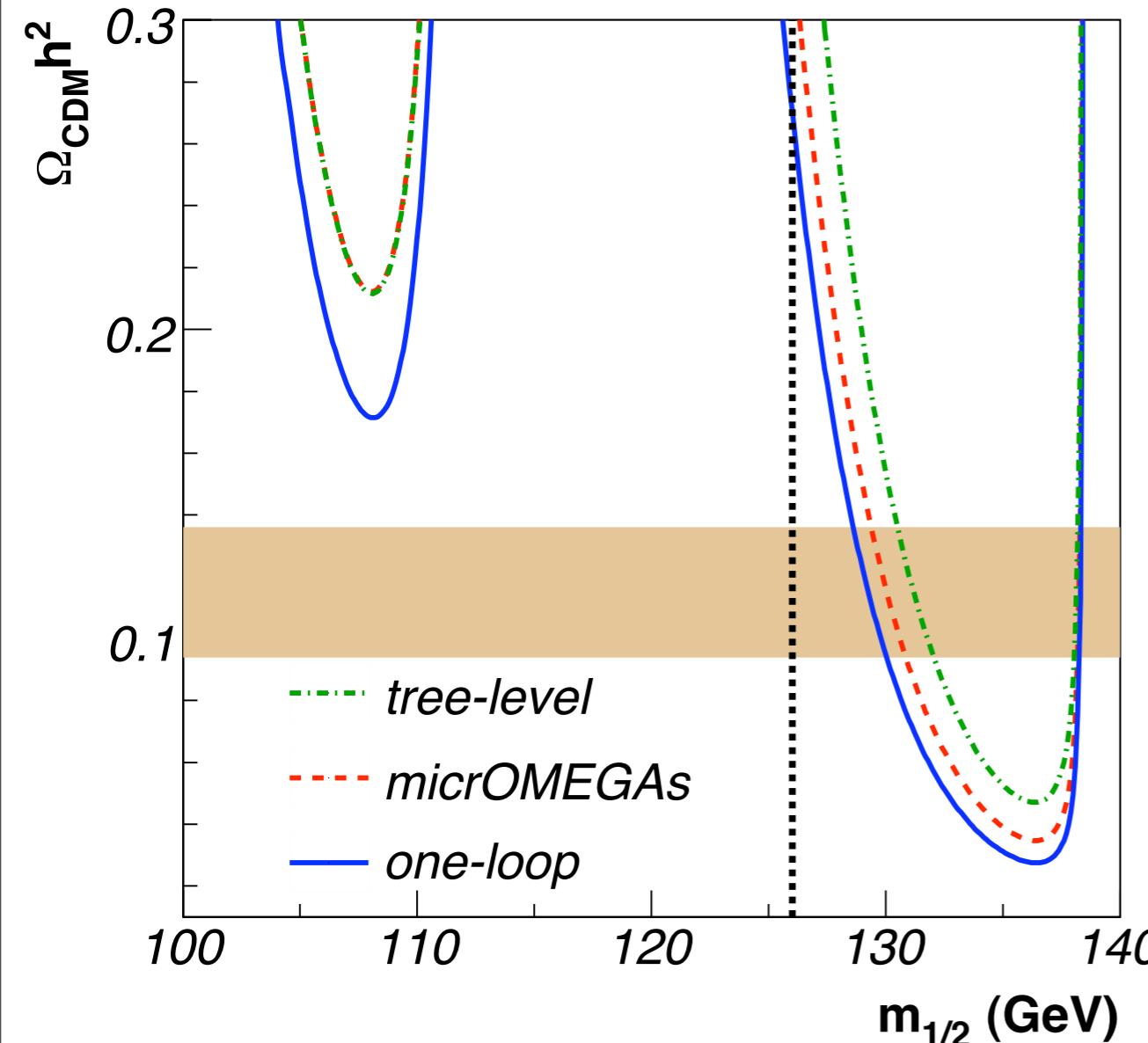
mSUGRA parameters: $m_0 = 1500$ GeV, $m_{1/2} = 130$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



“Effective” Yukawa coupling very good approximation around the Higgs resonance
Position of resonance does not coincide with peak of velocity distribution for thermal average
Cross section enhanced significantly by one-loop QCD and SUSY-QCD corrections

Numerical results: Light Higgs resonance (preliminary)

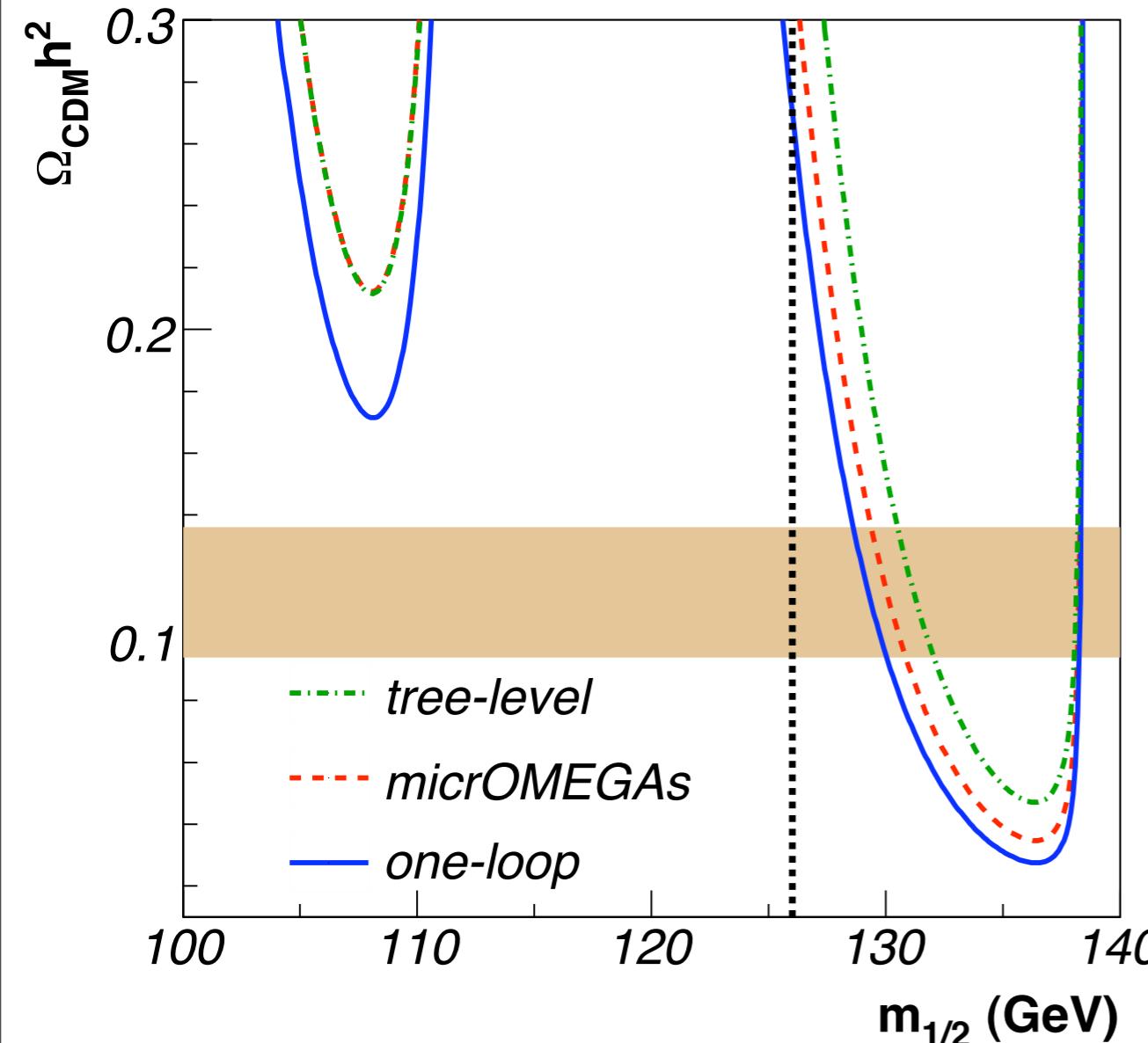
mSUGRA parameters: $m_0 = 1500$ GeV, $m_{1/2} = 130$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



Effect on cross section influences directly the prediction of the relic density

Numerical results: Light Higgs resonance (preliminary)

mSUGRA parameters: $m_0 = 1500$ GeV, $m_{1/2} = 130$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$

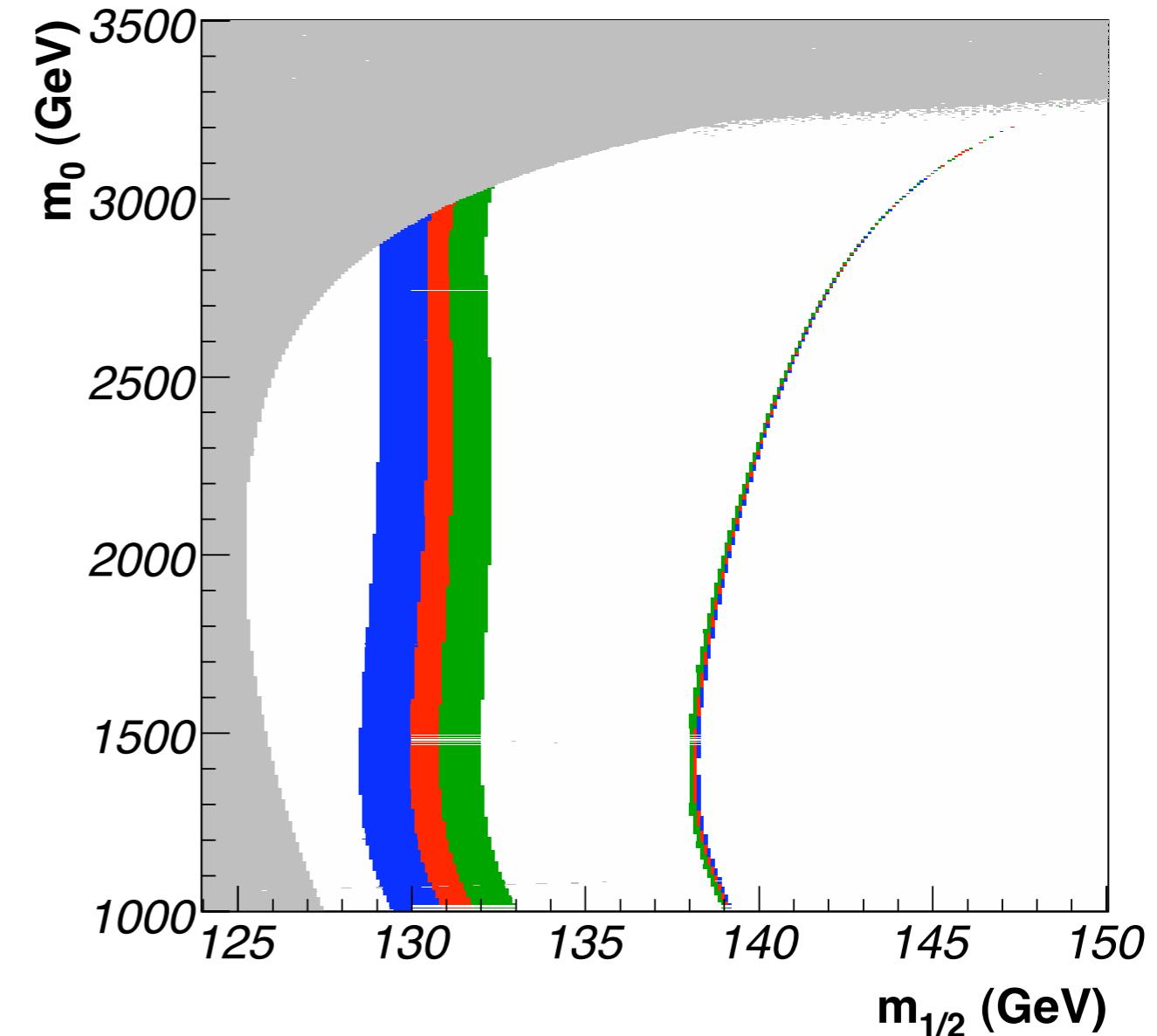
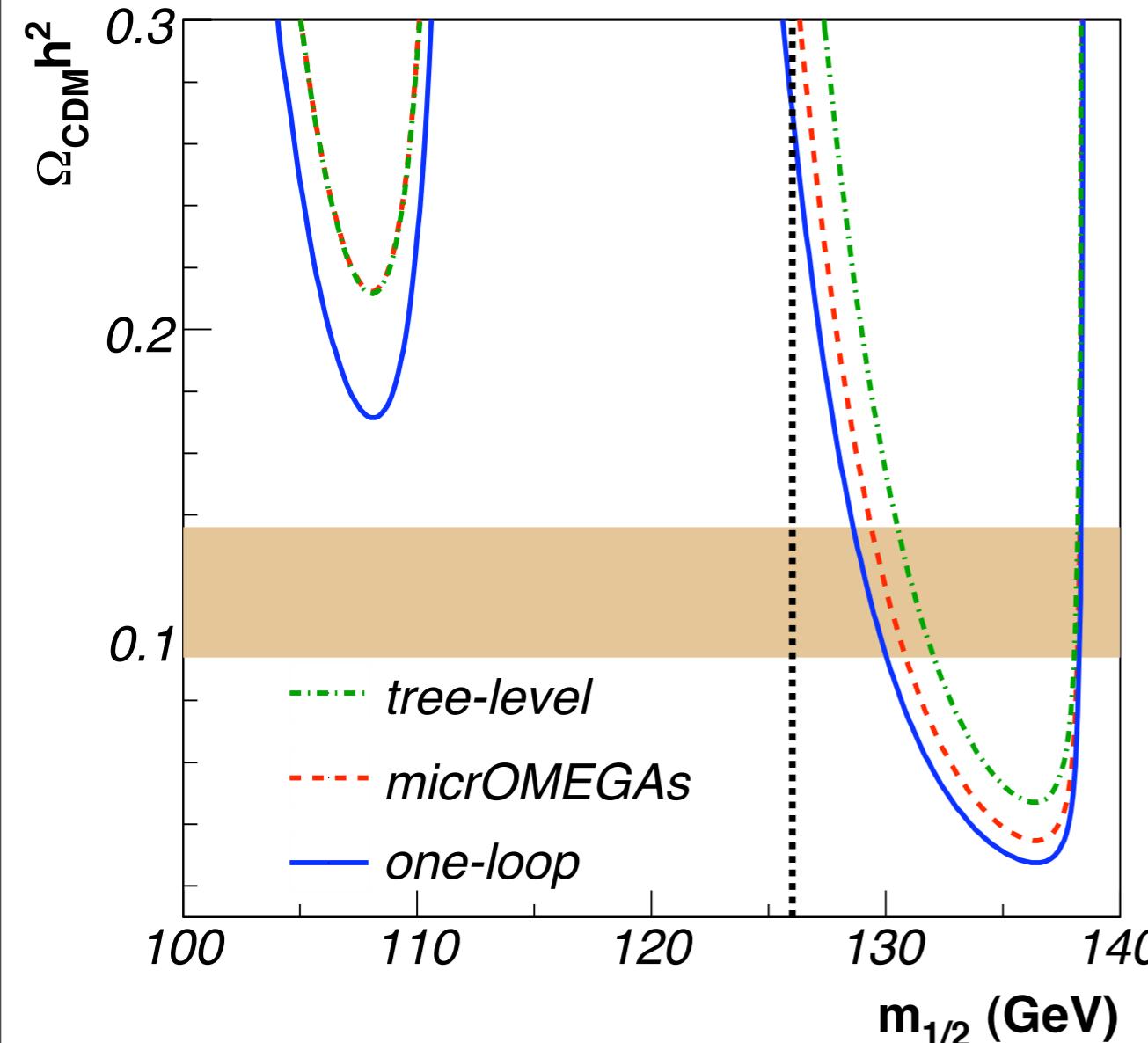


Effect on cross section influences directly the prediction of the relic density

Correction to annihilation through Z-boson exchange significant

Numerical results: Light Higgs resonance (preliminary)

mSUGRA parameters: $m_0=1500$ GeV, $m_{1/2}=130$ GeV, $A_0=-1500$ GeV, $\tan\beta=10$, $\mu>0$

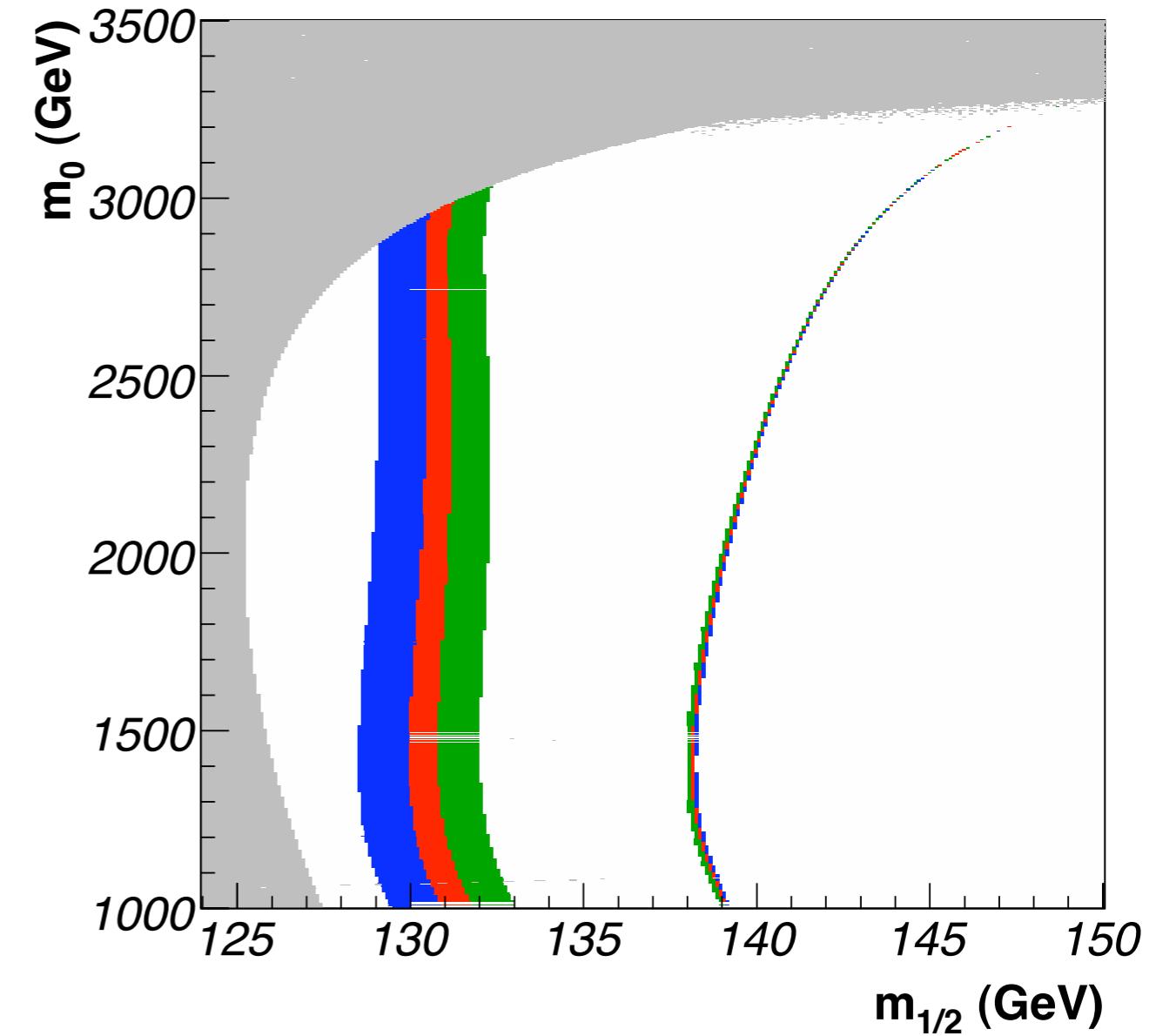
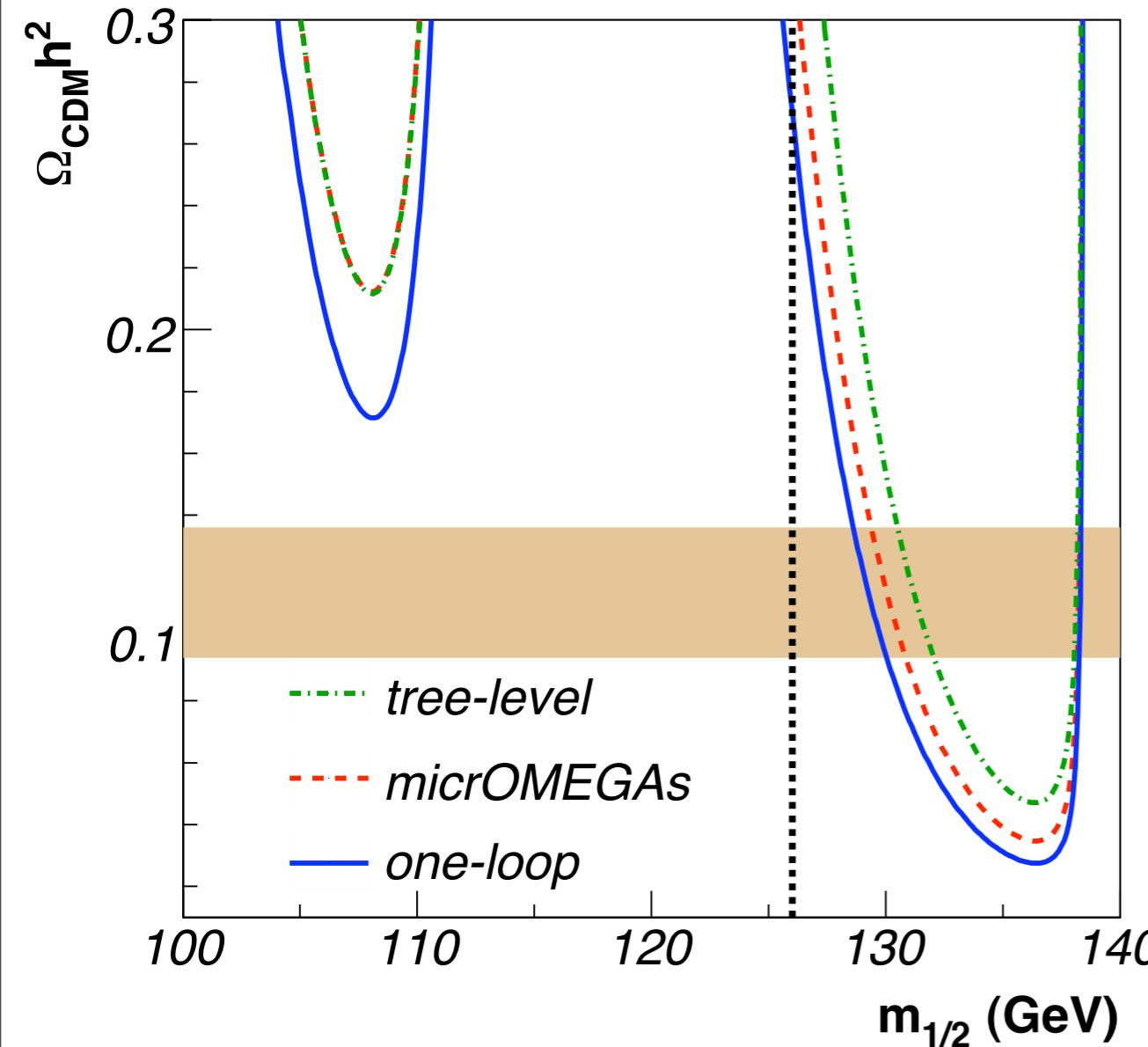


Effect on cross section influences directly the prediction of the relic density

Correction to annihilation through Z-boson exchange significant

Numerical results: Light Higgs resonance (preliminary)

mSUGRA parameters: $m_0=1500$ GeV, $m_{1/2}=130$ GeV, $A_0=-1500$ GeV, $\tan\beta=10$, $\mu>0$



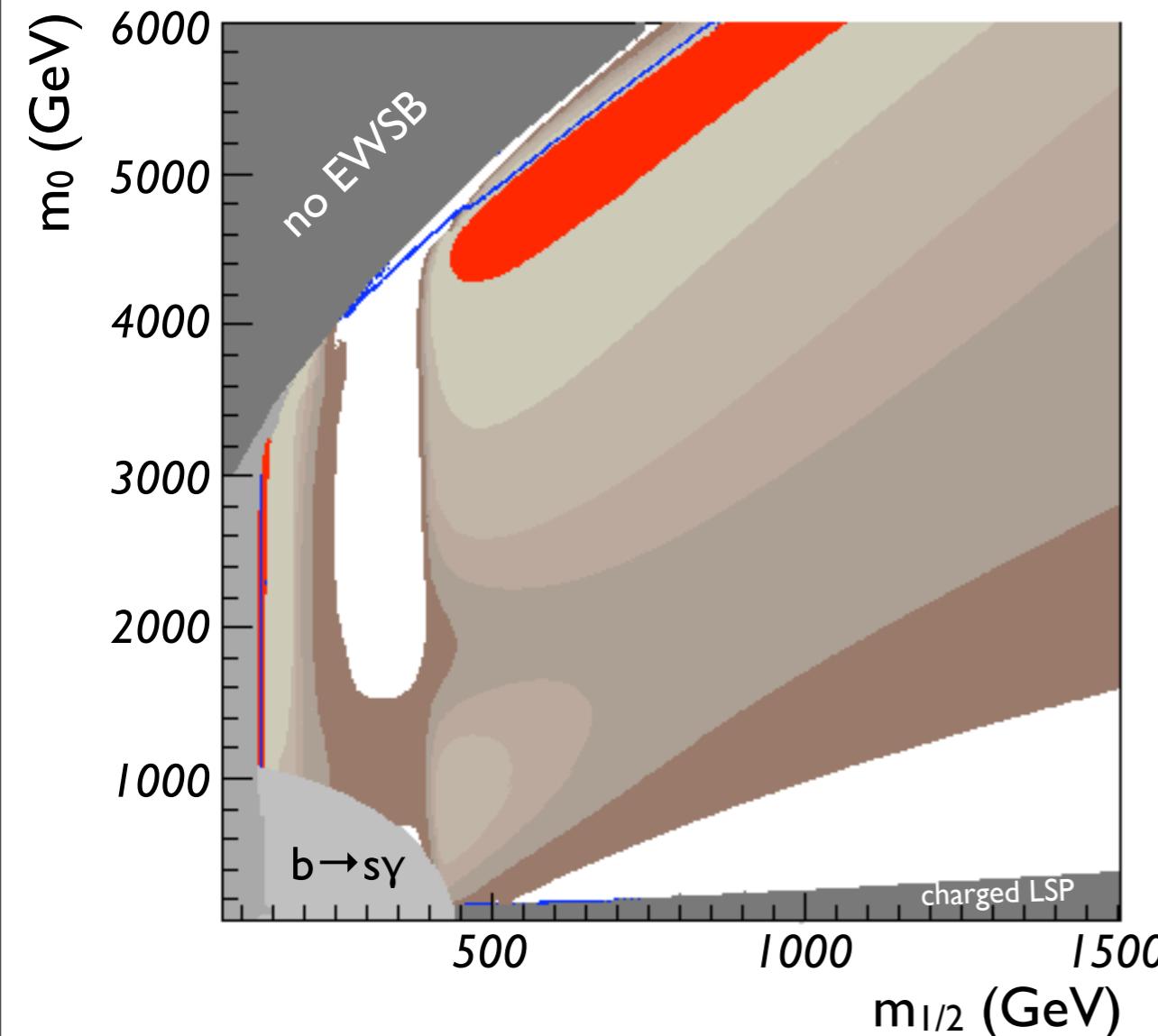
Effect on cross section influences directly the prediction of the relic density

Correction to annihilation through Z-boson exchange significant

Favoured interval shifted towards chargino mass limit $m_{1/2} \geq 126$ GeV

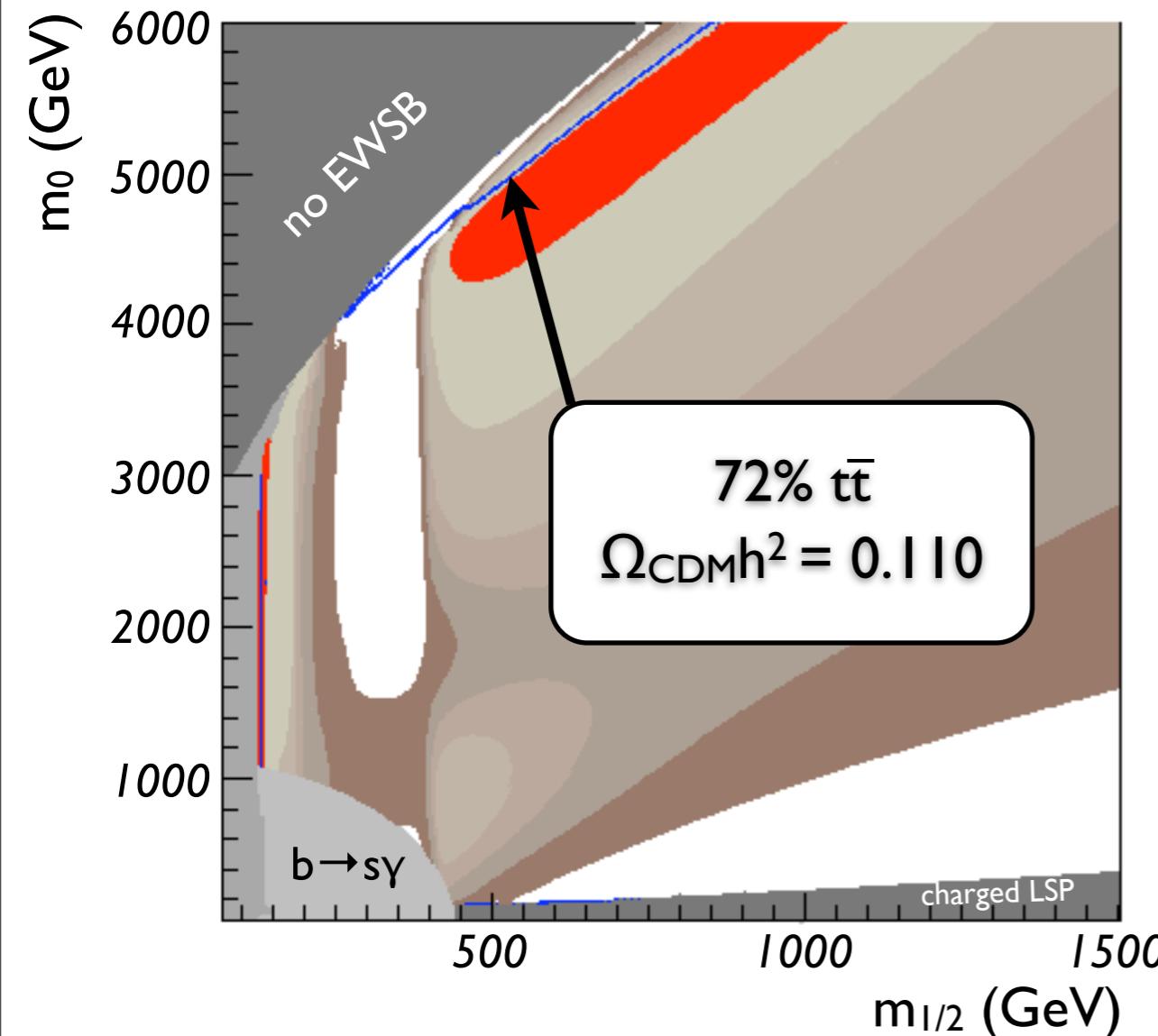
Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=5300$ GeV, $m_{1/2}=625$ GeV, $A_0=-1500$ GeV, $\tan\beta=10$, $\mu>0$



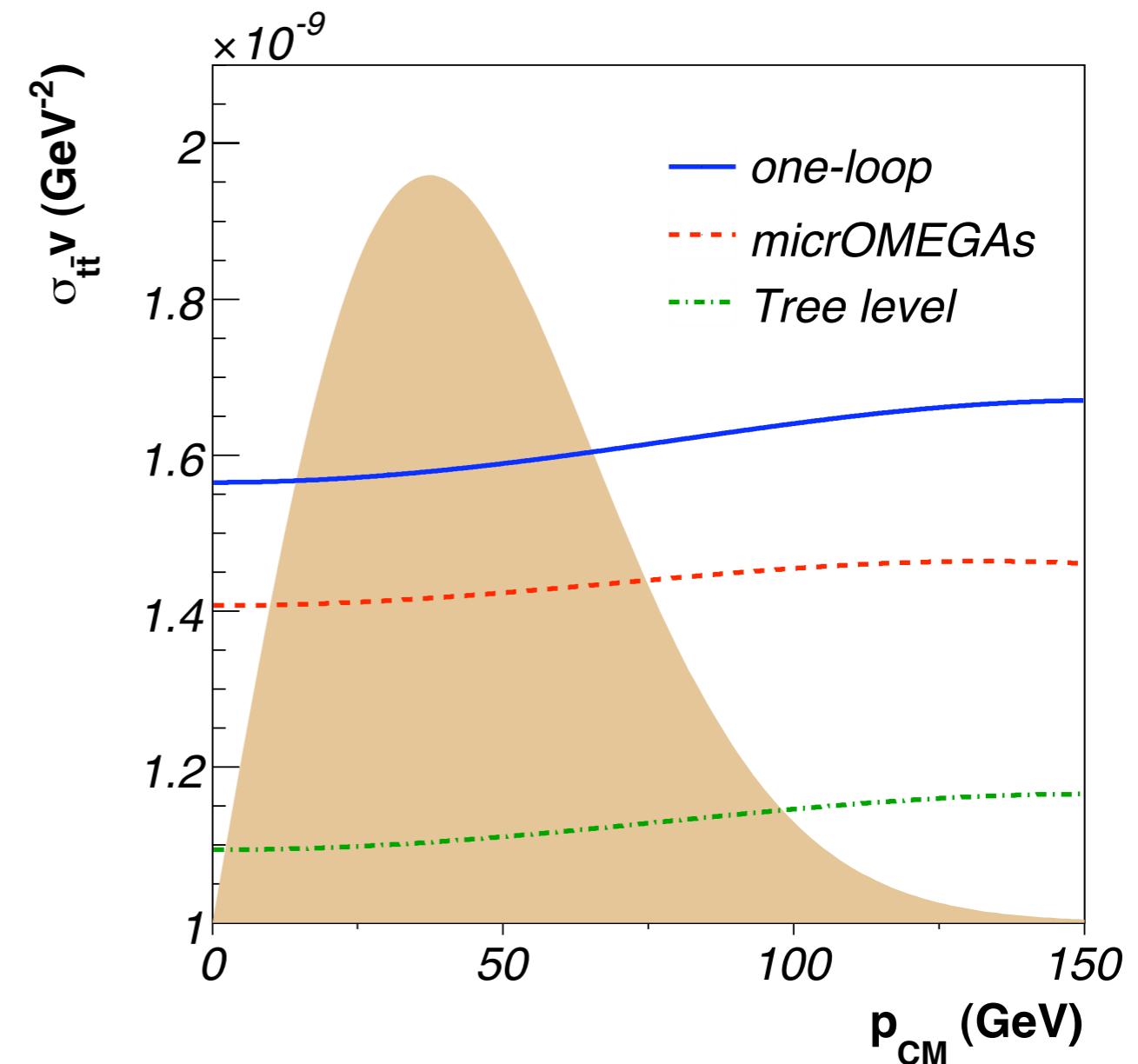
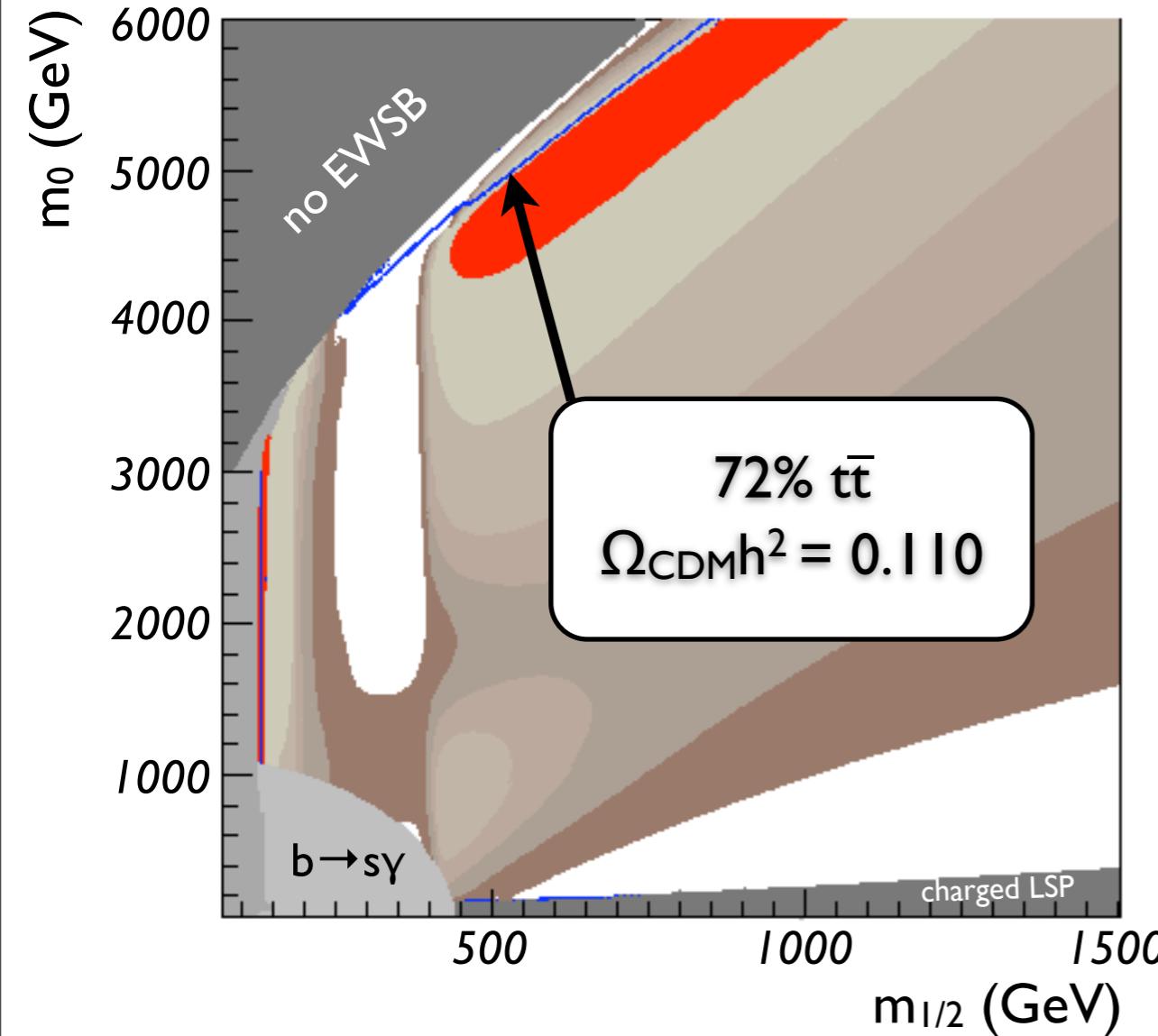
Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0 = 5300$ GeV, $m_{1/2} = 625$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



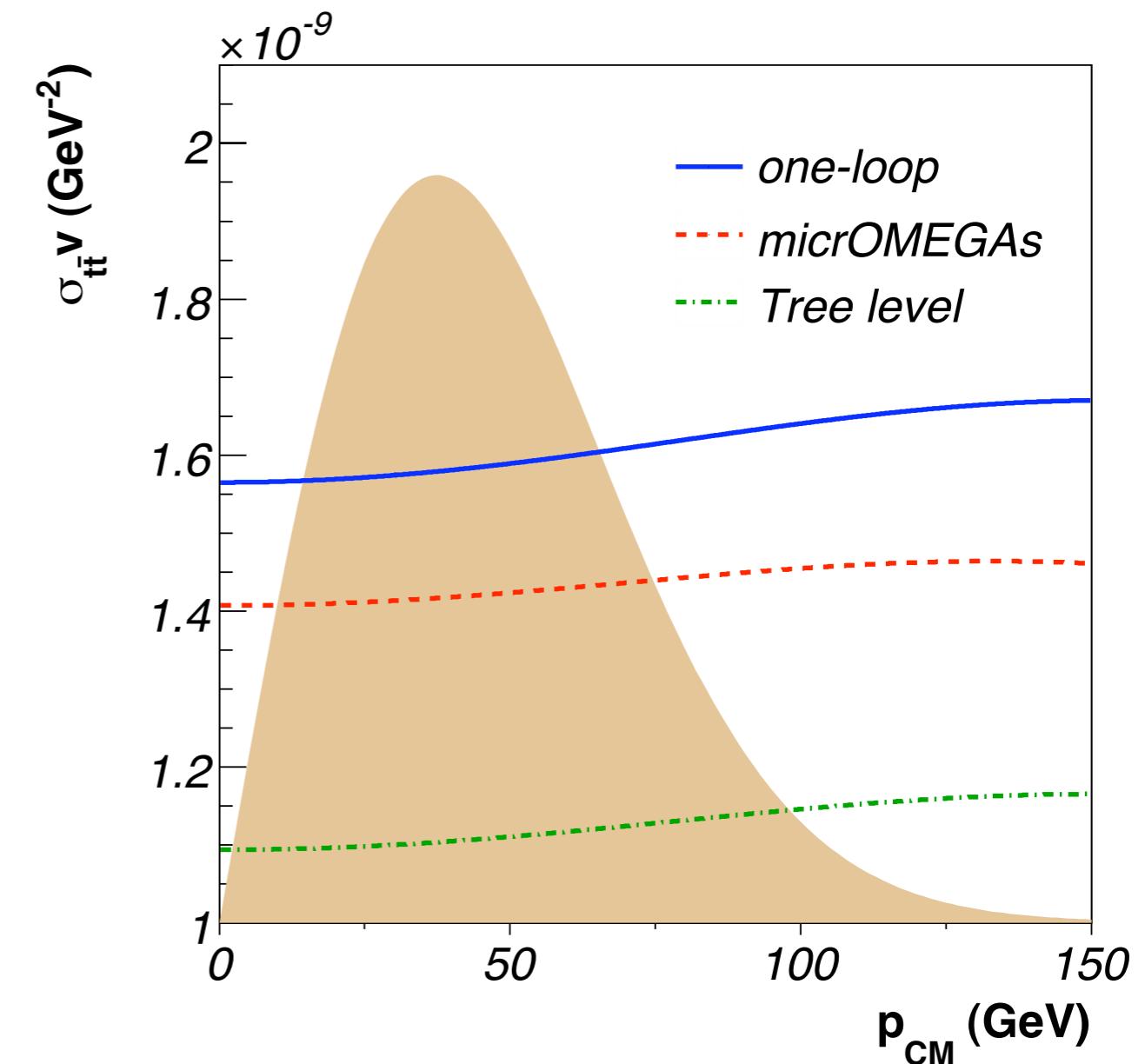
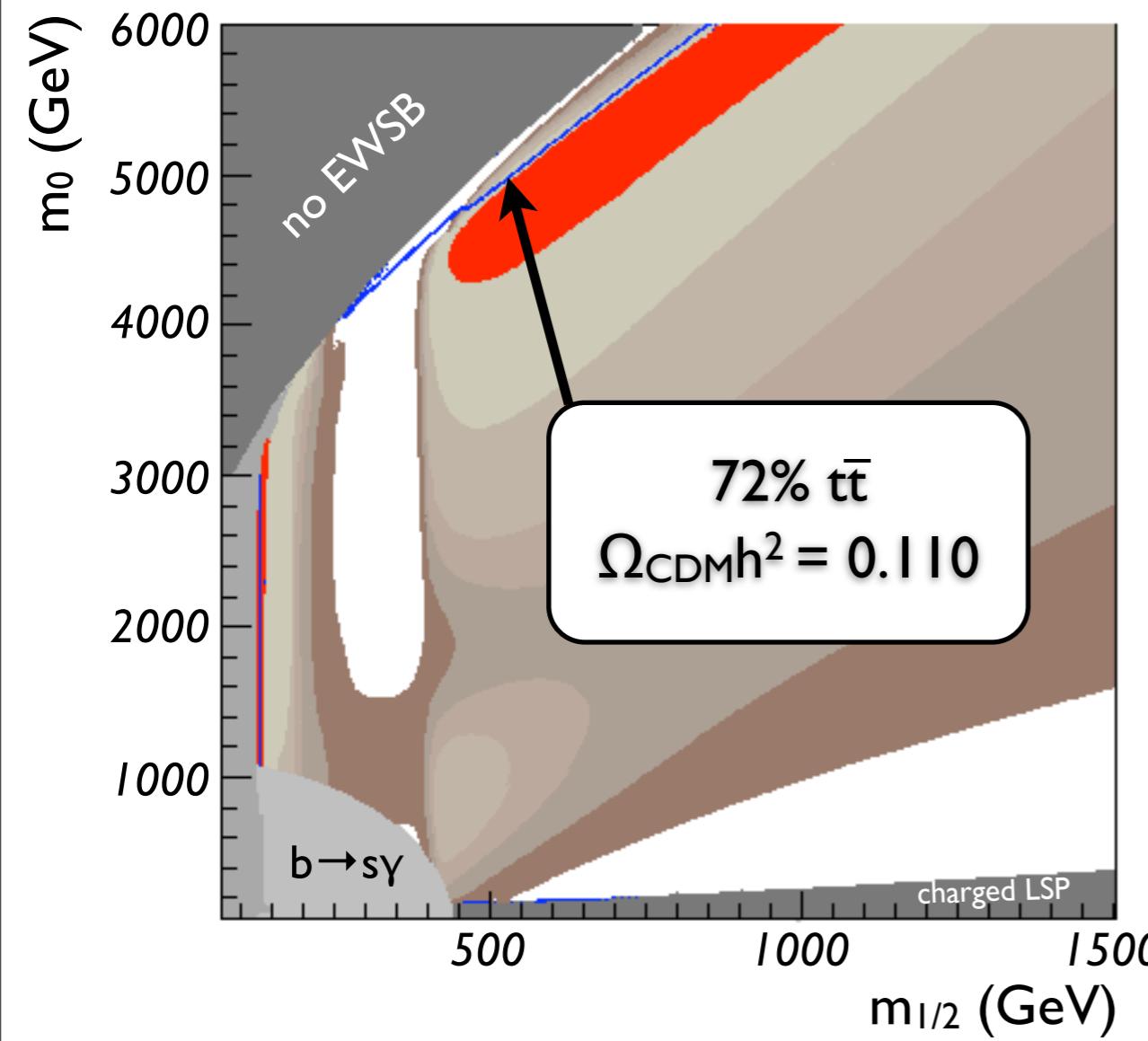
Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0 = 5300$ GeV, $m_{1/2} = 625$ GeV, $A_0 = -1500$ GeV, $\tan\beta = 10$, $\mu > 0$



Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=5300$ GeV, $m_{1/2}=625$ GeV, $A_0=-1500$ GeV, $\tan\beta=10$, $\mu>0$

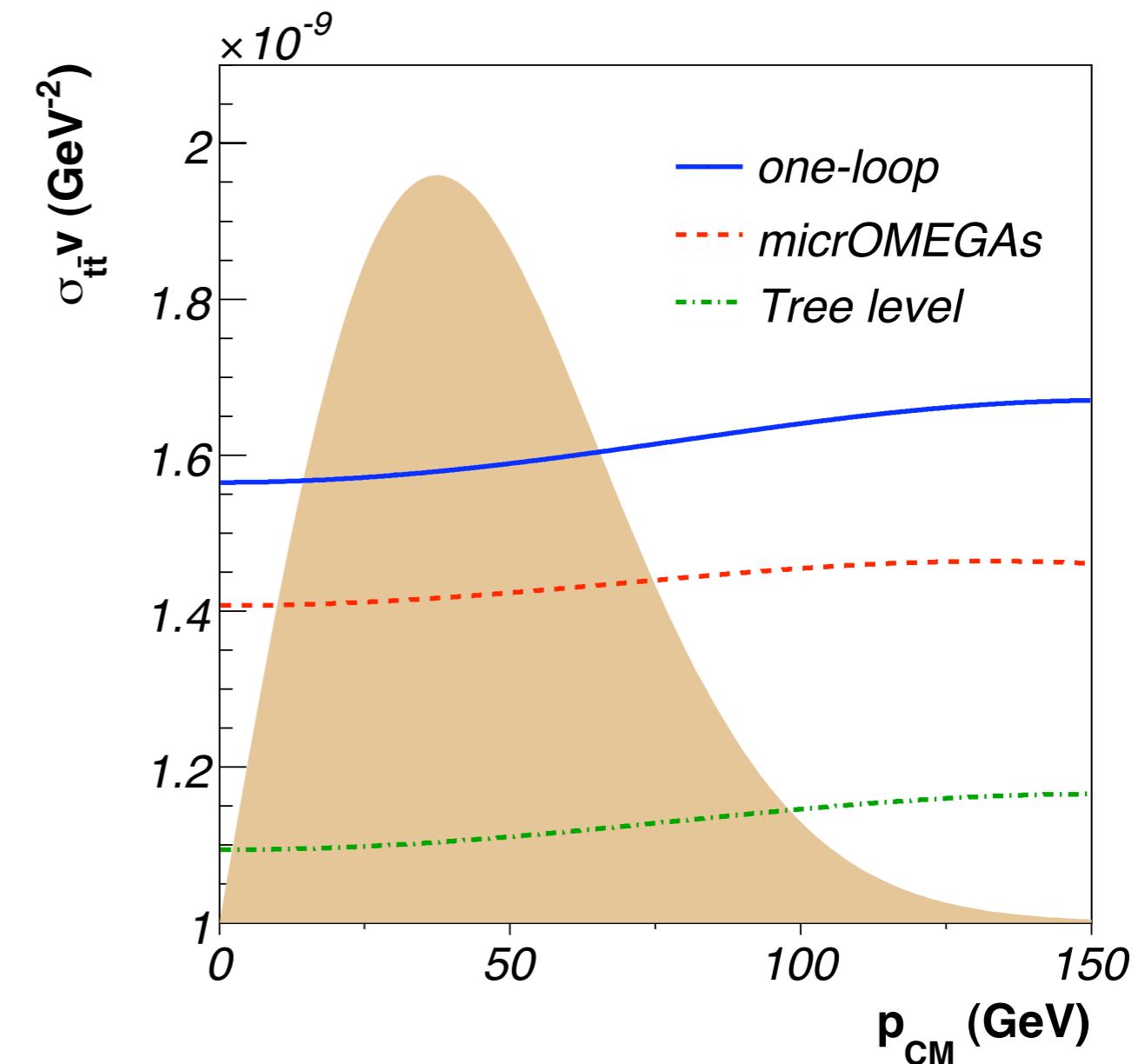
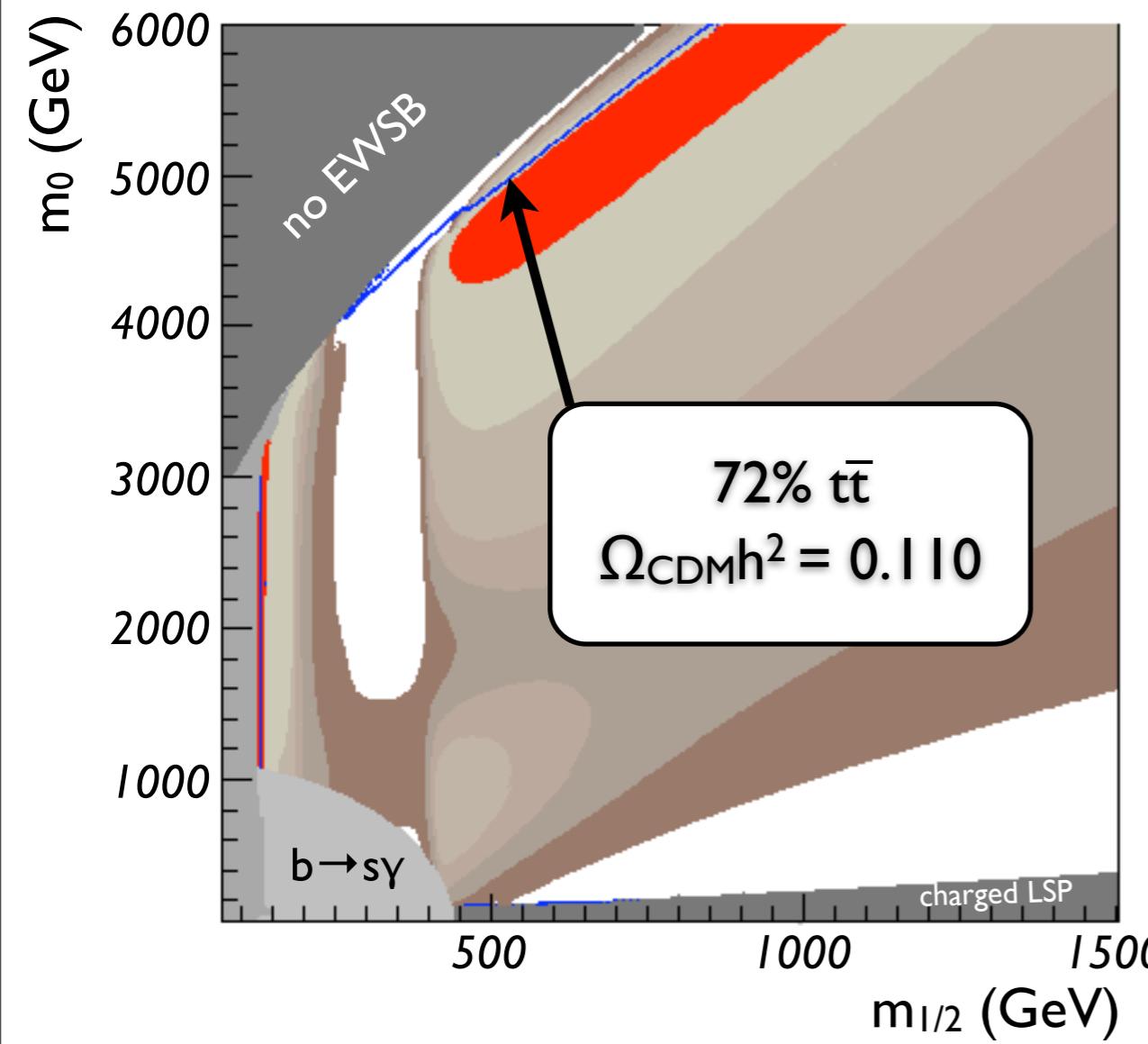


Higgs resonance situated at rather high energies around $p_{cm} \approx 2610$ GeV

Annihilation cross section sizeable due to large top quark mass

Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=5300$ GeV, $m_{1/2}=625$ GeV, $A_0=-1500$ GeV, $\tan\beta=10$, $\mu>0$



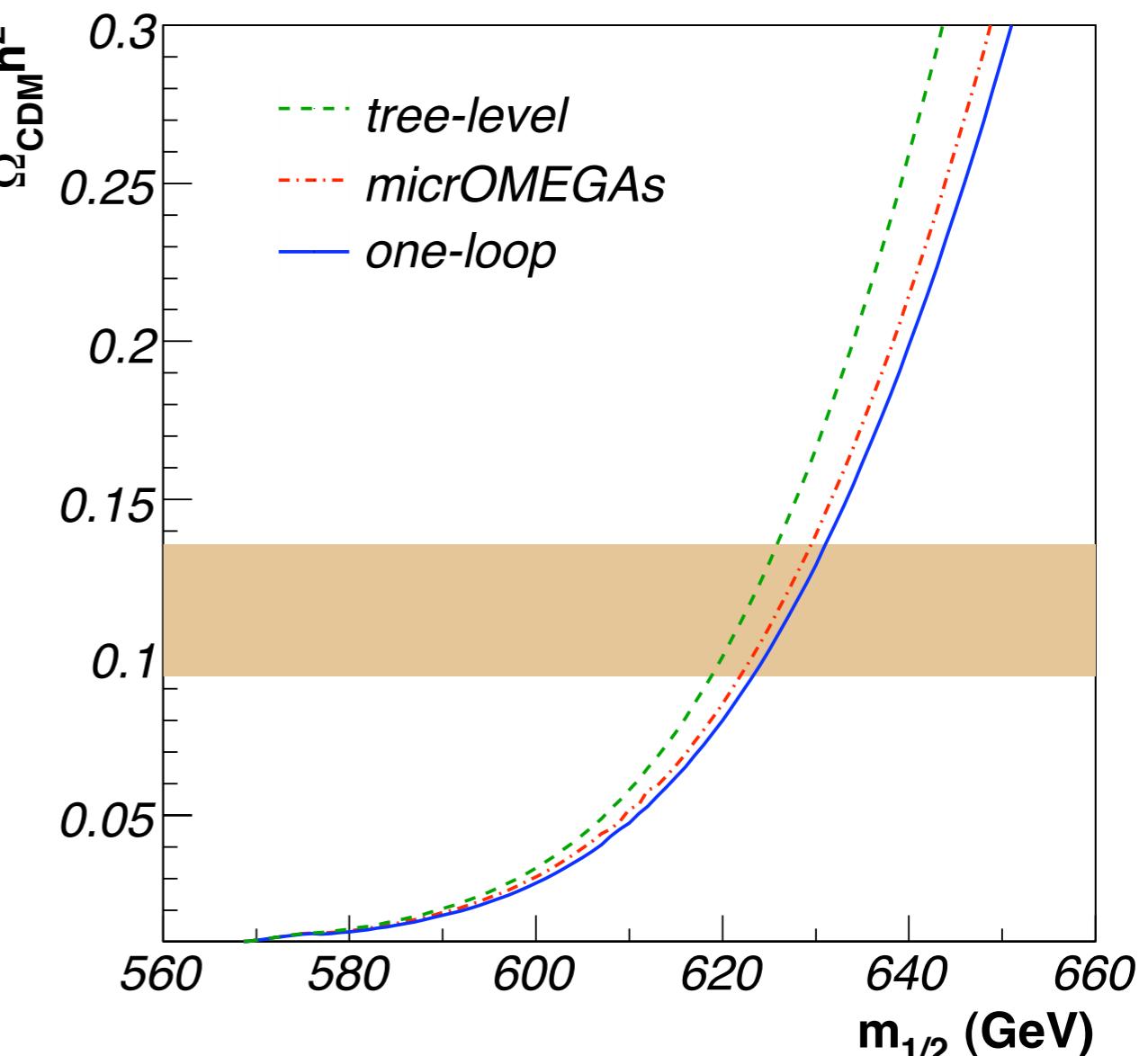
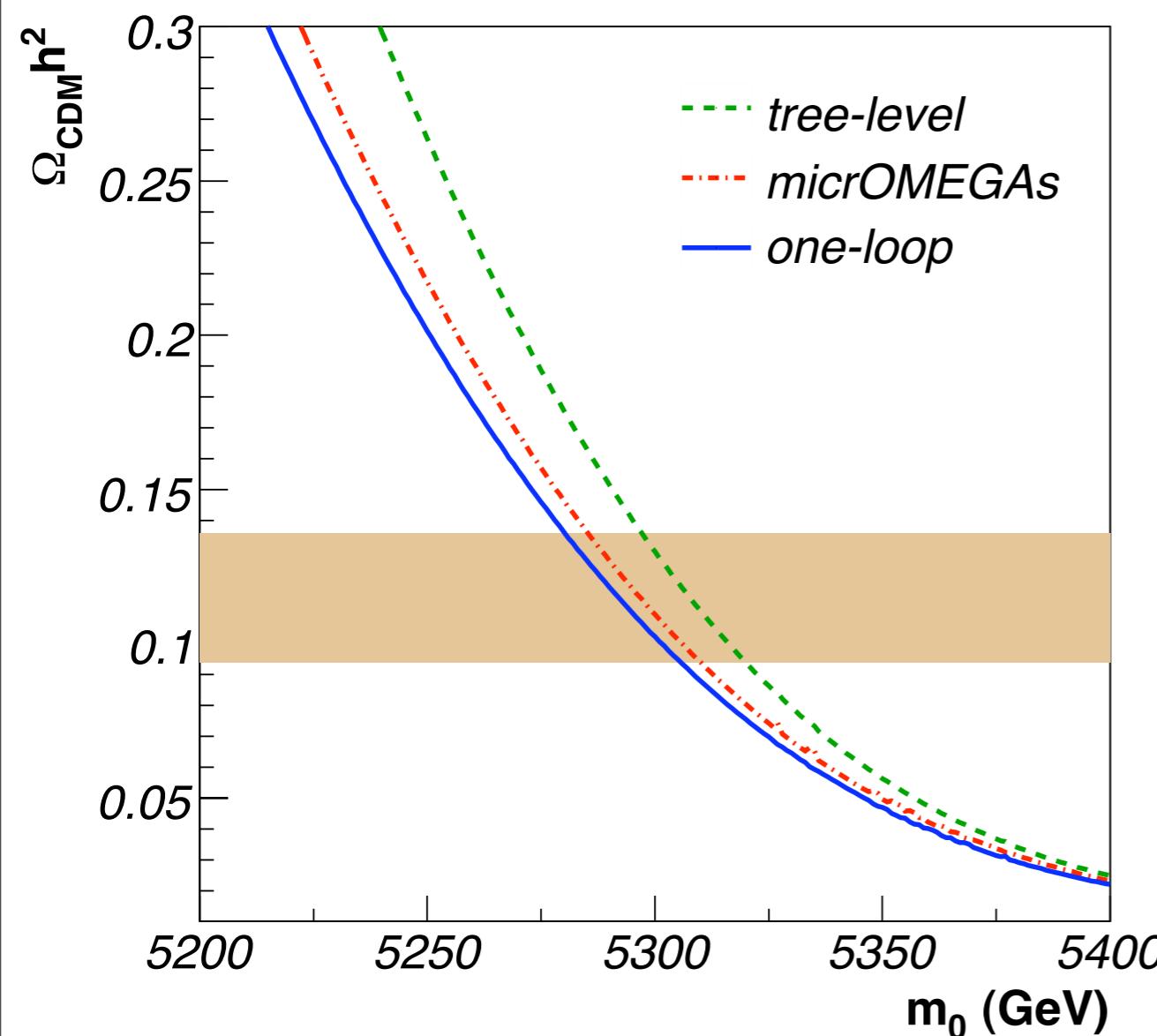
Higgs resonance situated at rather high energies around $p_{cm} \approx 2610$ GeV

Annihilation cross section sizeable due to large top quark mass

Cross section enhanced significantly by one-loop QCD and SUSY-QCD corrections

Numerical results: Focus point region (preliminary)

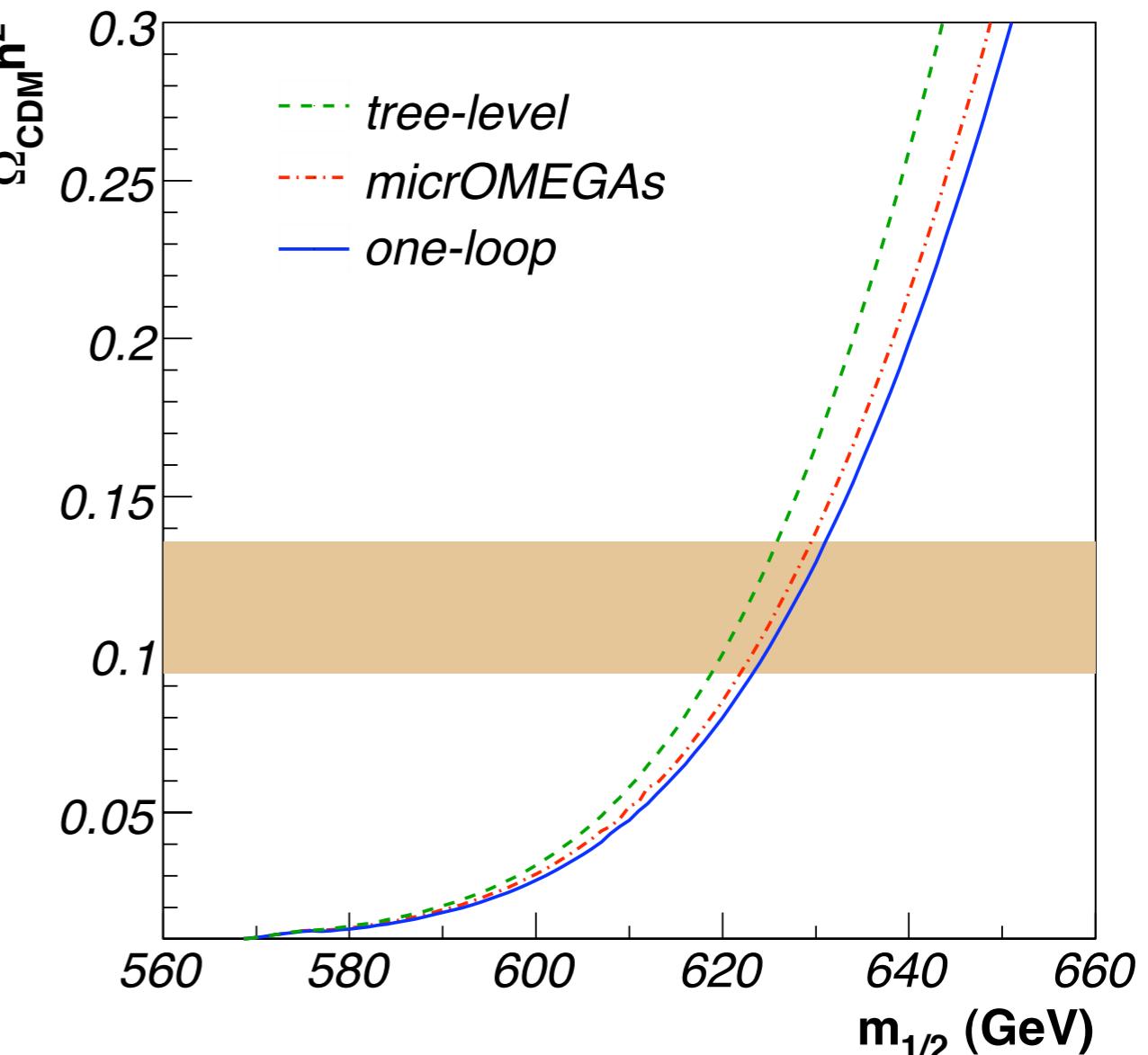
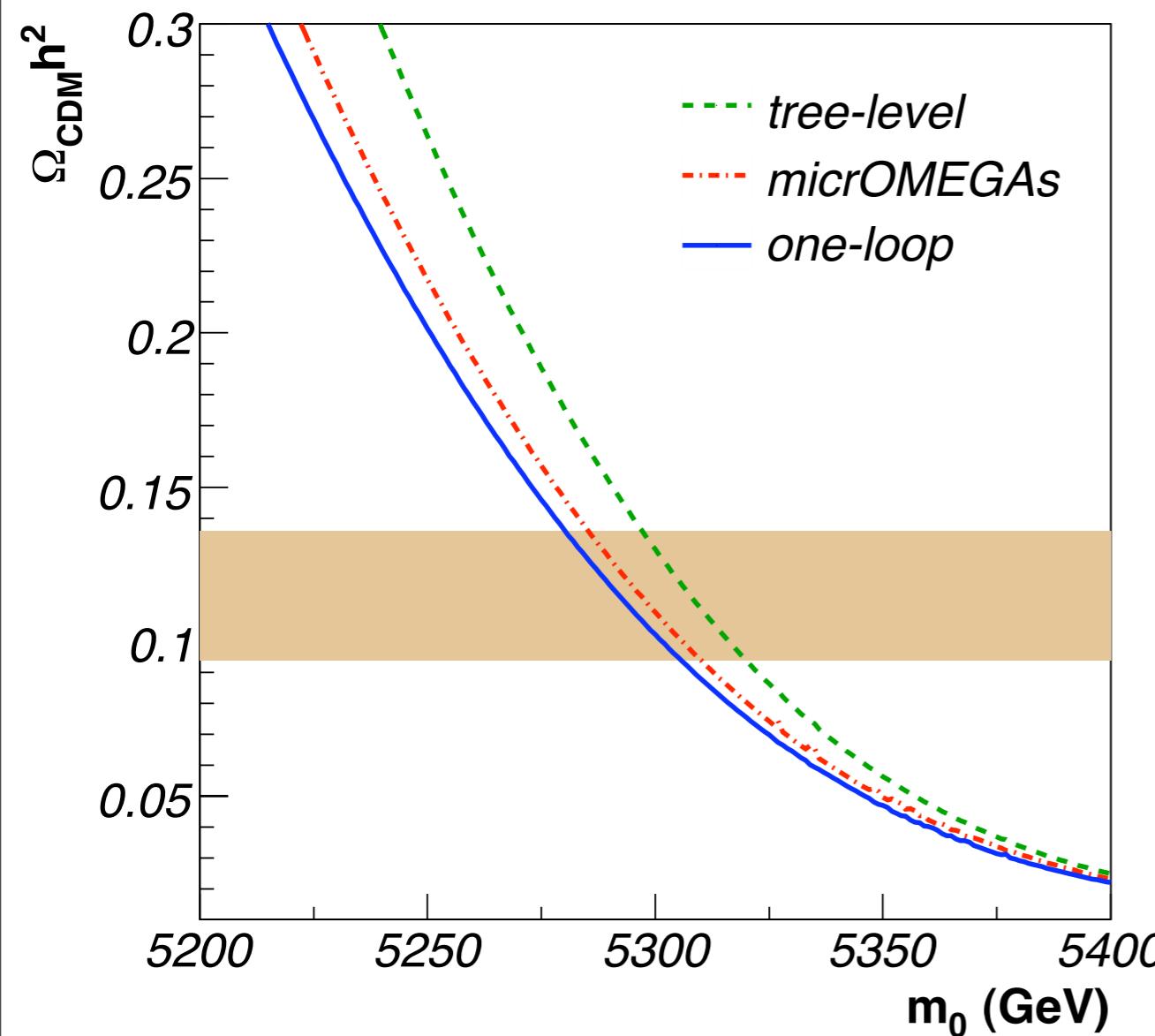
mSUGRA parameters: $m_0=5300$ GeV, $m_{1/2}=625$ GeV, $A_0=-1500$ GeV, $\tan\beta=10$, $\mu>0$



Prediction of neutralino relic density receives sizeable corrections

Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=5300$ GeV, $m_{1/2}=625$ GeV, $A_0=-1500$ GeV, $\tan\beta=10$, $\mu>0$

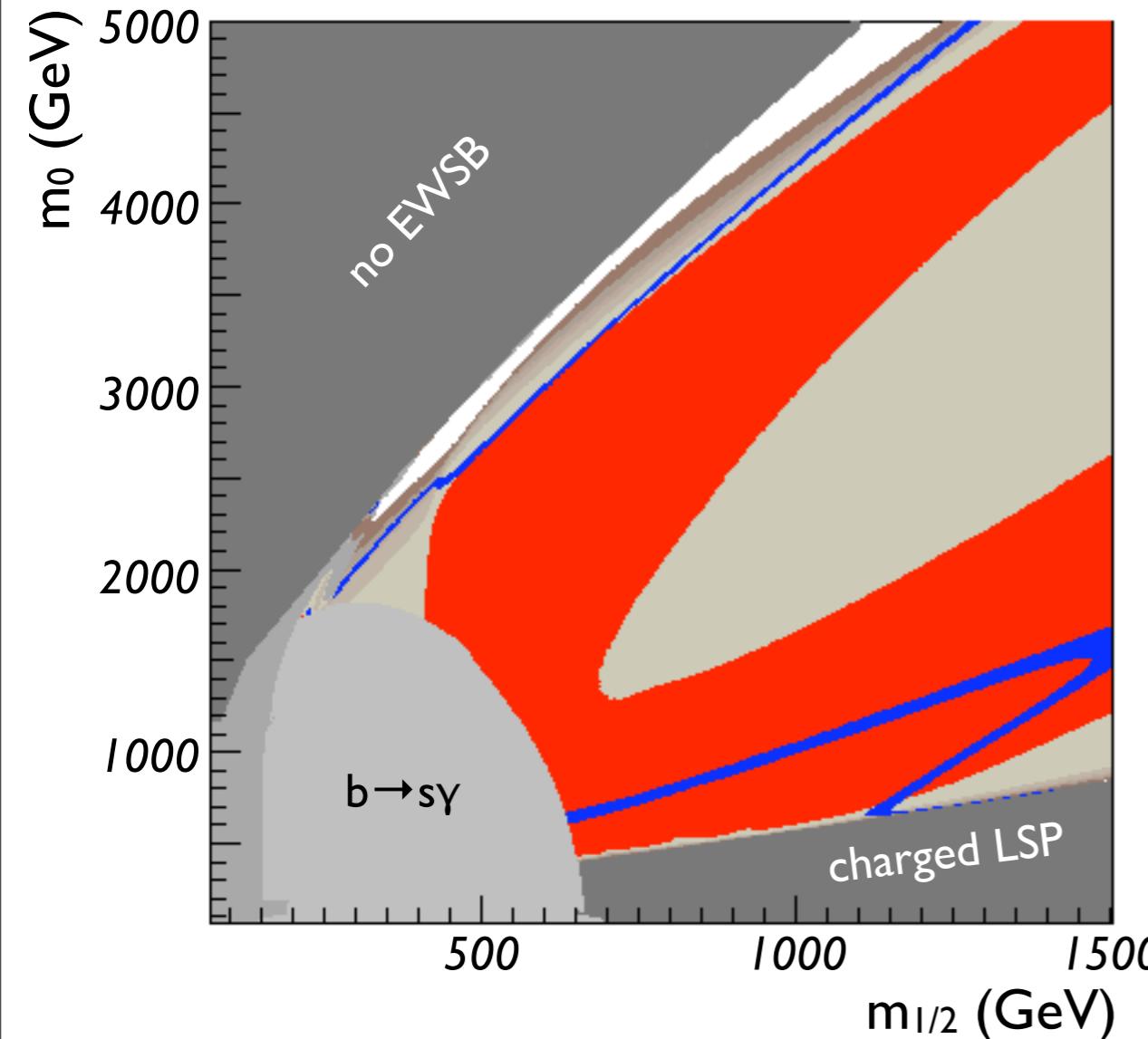


Prediction of neutralino relic density receives sizeable corrections

favoured region shifted to higher scalar masses m_0 and smaller gaugino masses $m_{1/2}$

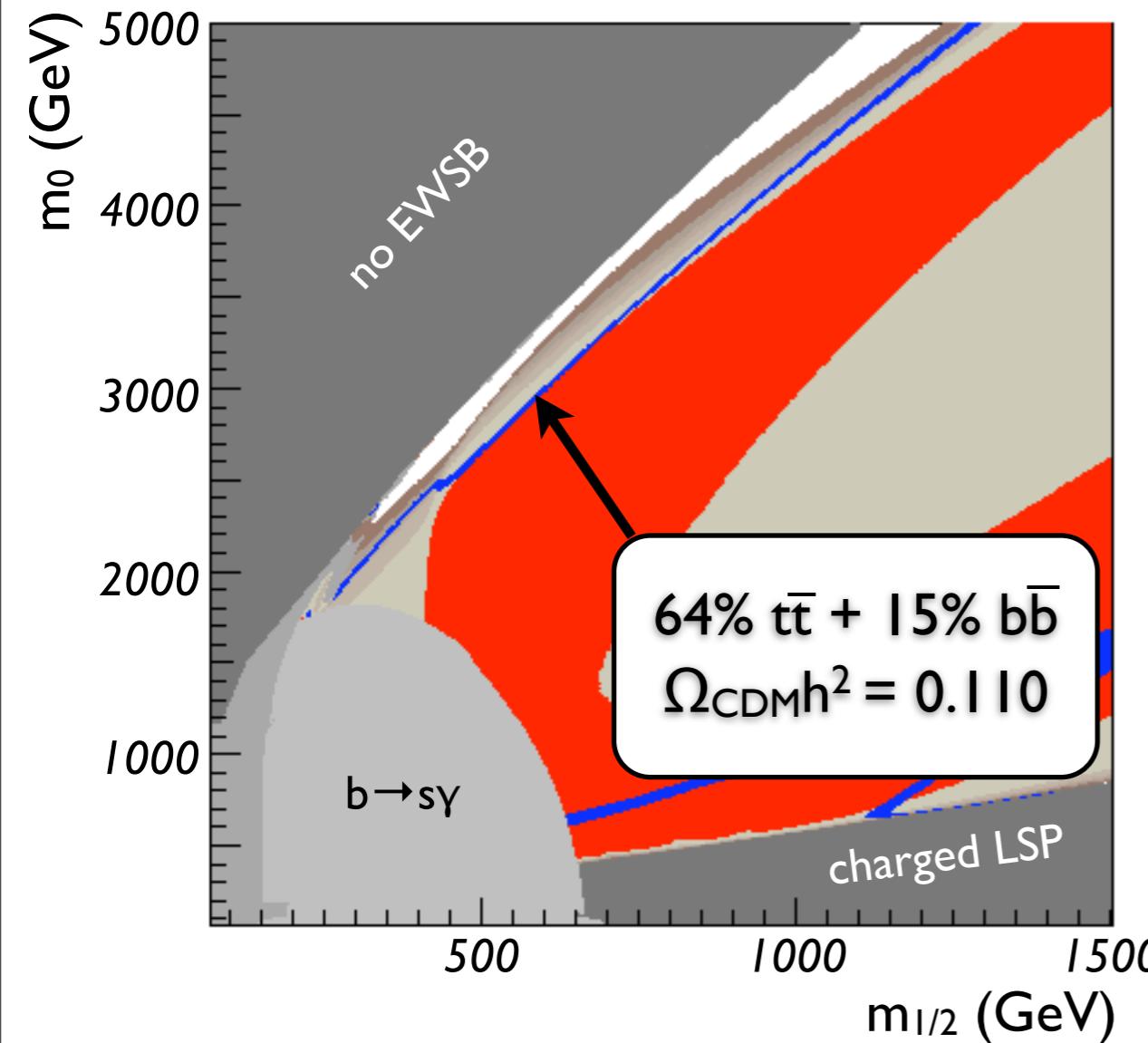
Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0 = 3000$ GeV, $m_{1/2} = 600$ GeV, $A_0 = 0$, $\tan\beta = 50$, $\mu > 0$



Numerical results: Focus point region (preliminary)

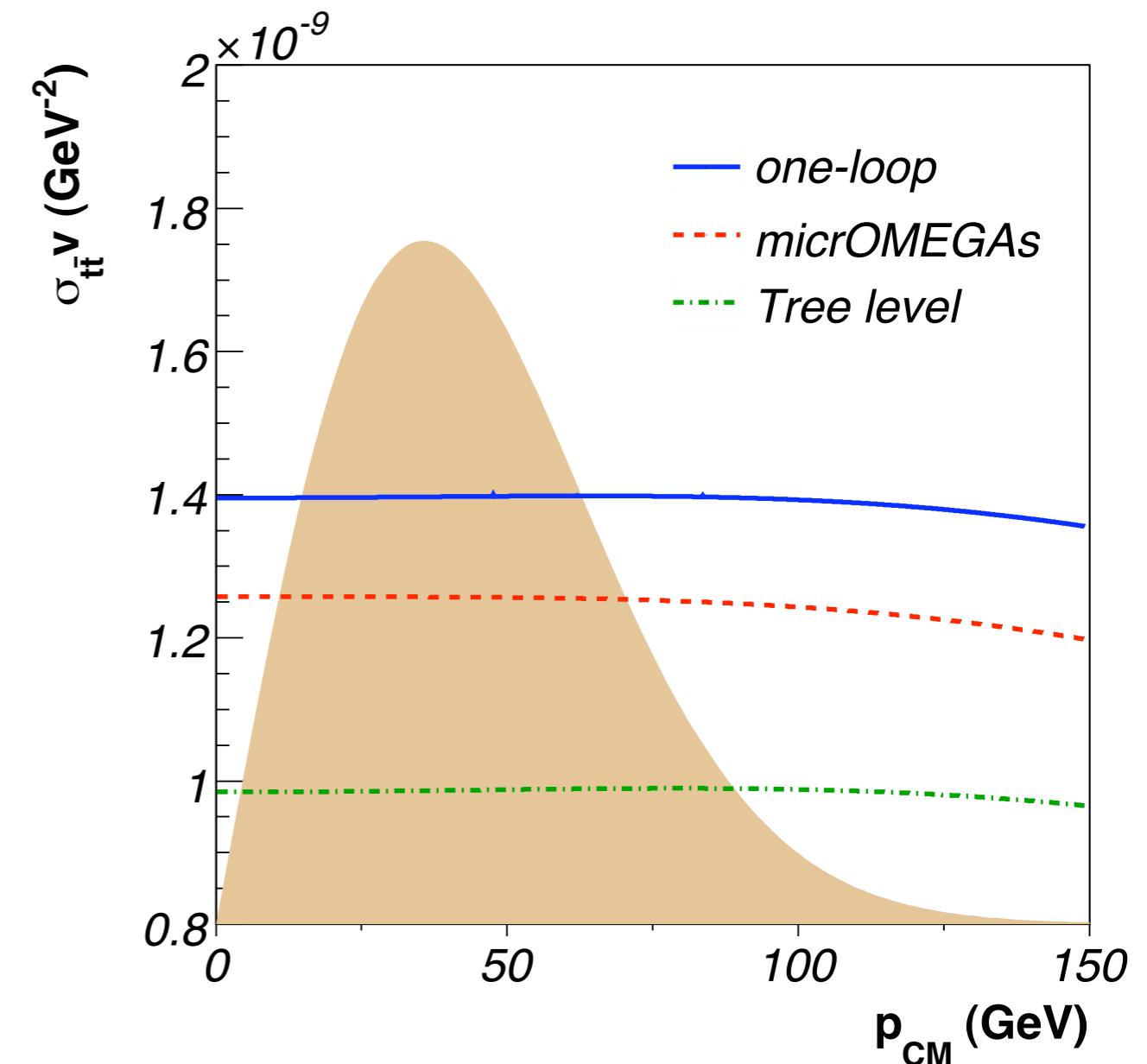
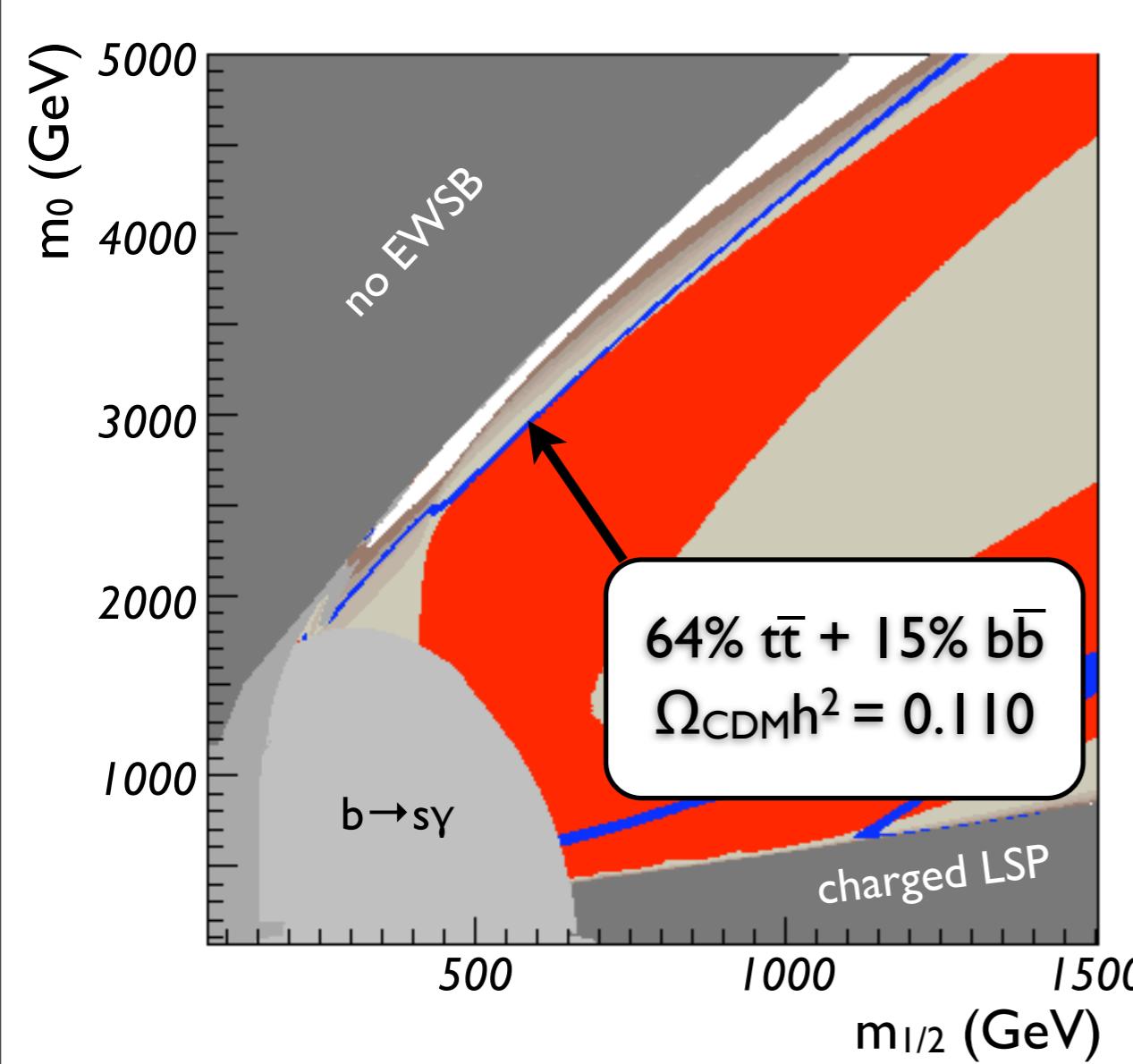
mSUGRA parameters: $m_0 = 3000$ GeV, $m_{1/2} = 600$ GeV, $A_0 = 0$, $\tan\beta = 50$, $\mu > 0$



Important top and bottom quark contributions due to large top mass and large $\tan\beta$

Numerical results: Focus point region (preliminary)

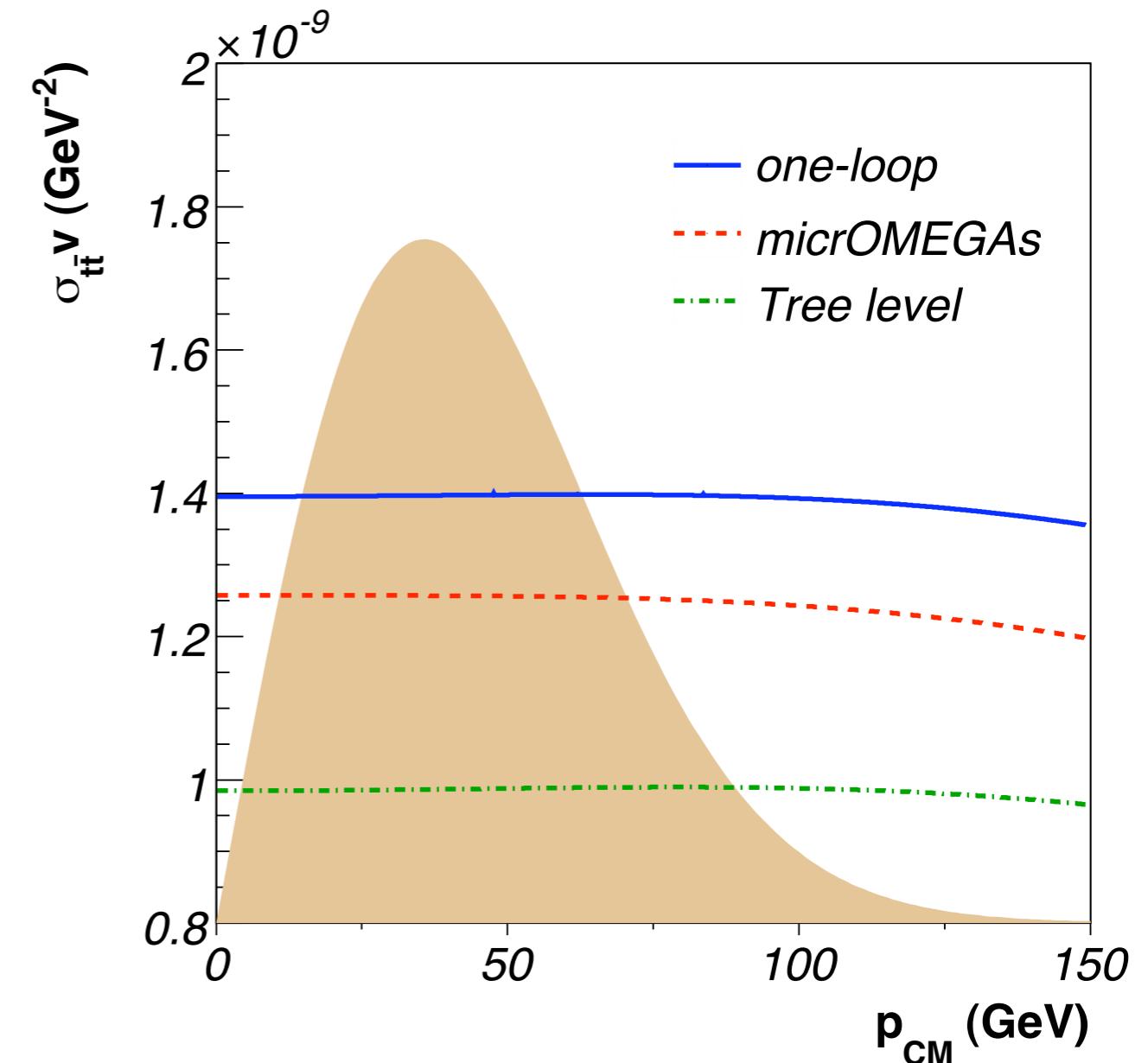
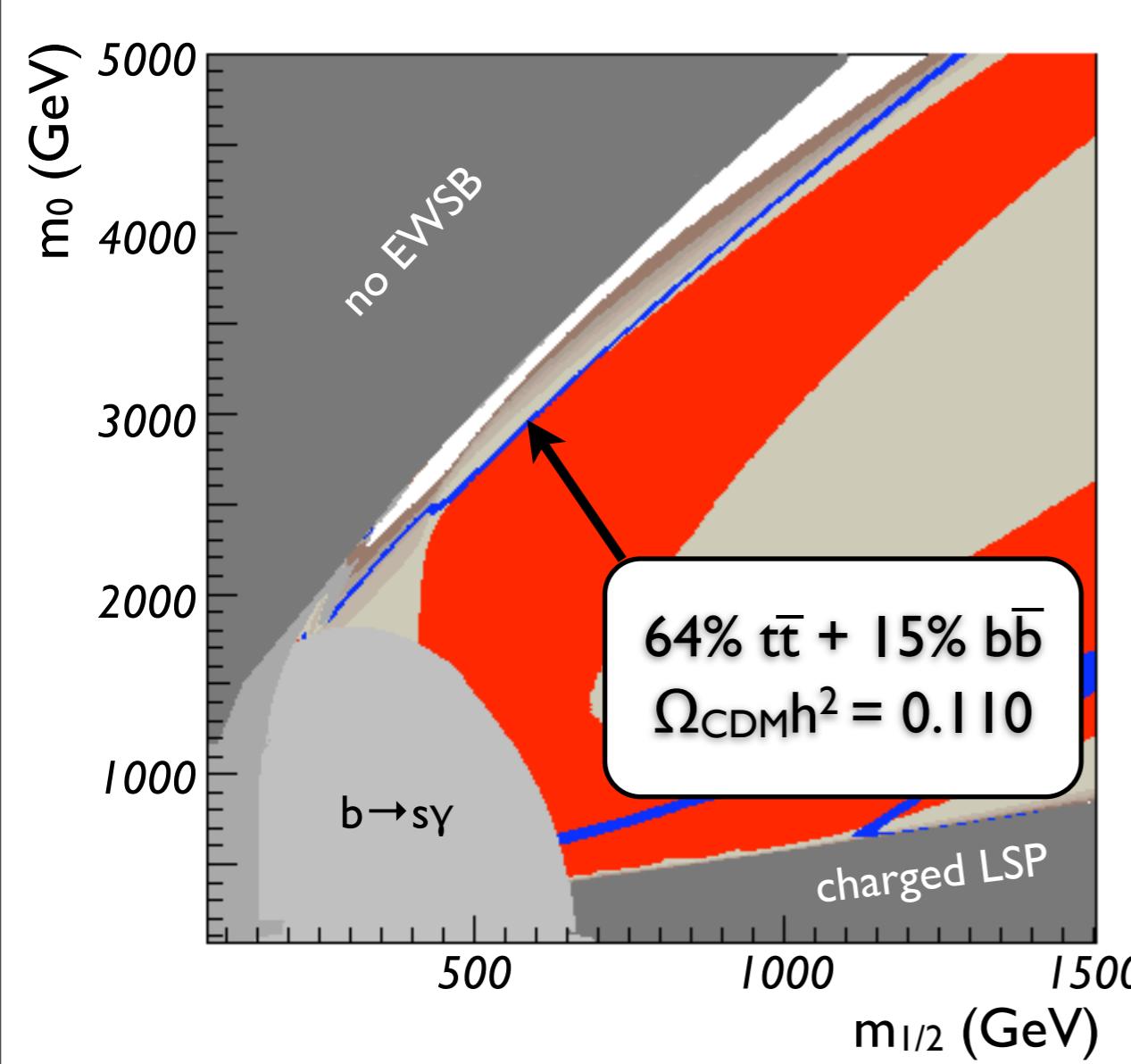
mSUGRA parameters: $m_0=3000$ GeV, $m_{1/2}=600$ GeV, $A_0=0$, $\tan\beta=50$, $\mu>0$



Important top and bottom quark contributions due to large top mass and large $\tan\beta$

Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=3000$ GeV, $m_{1/2}=600$ GeV, $A_0=0$, $\tan\beta=50$, $\mu>0$

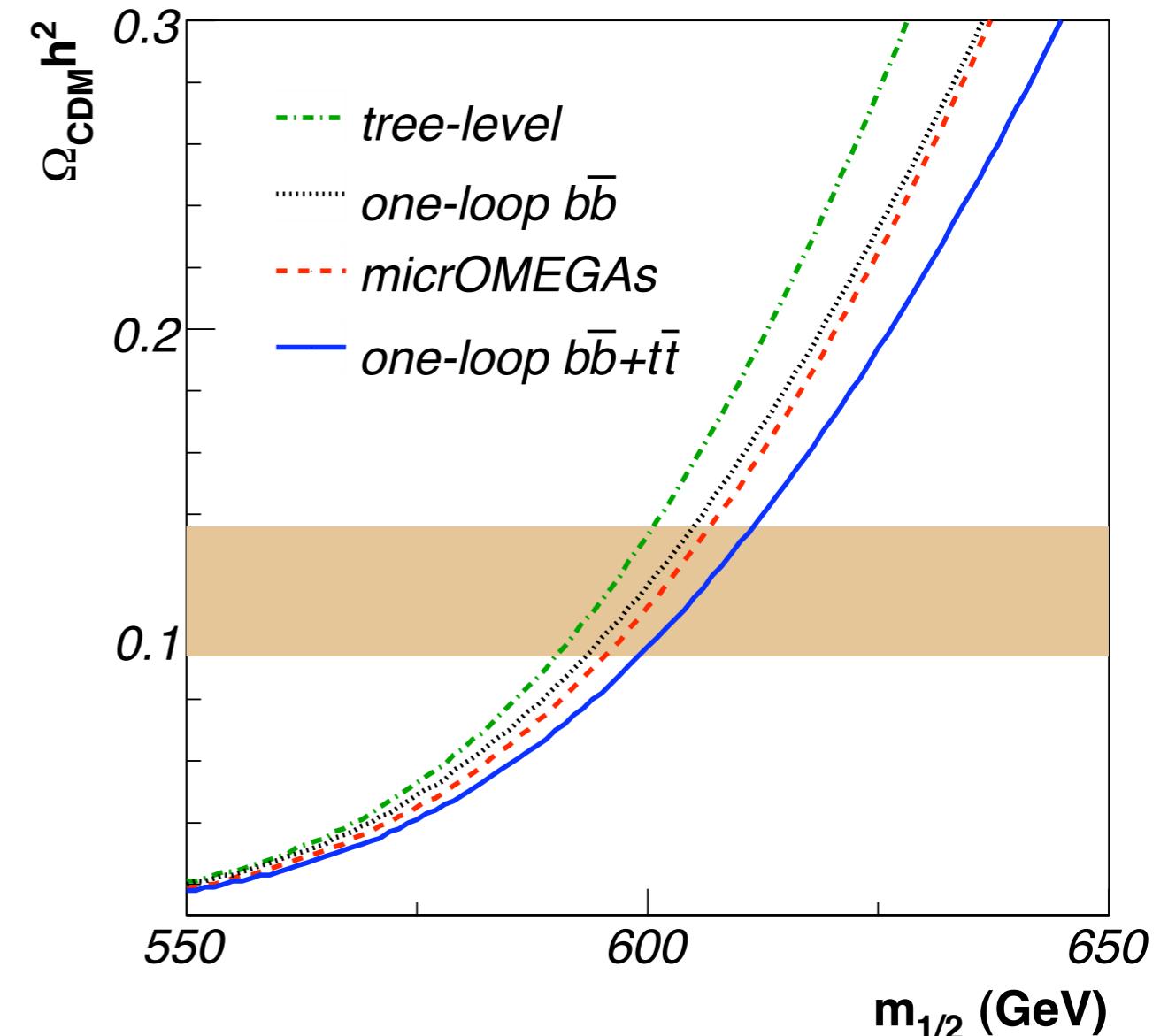
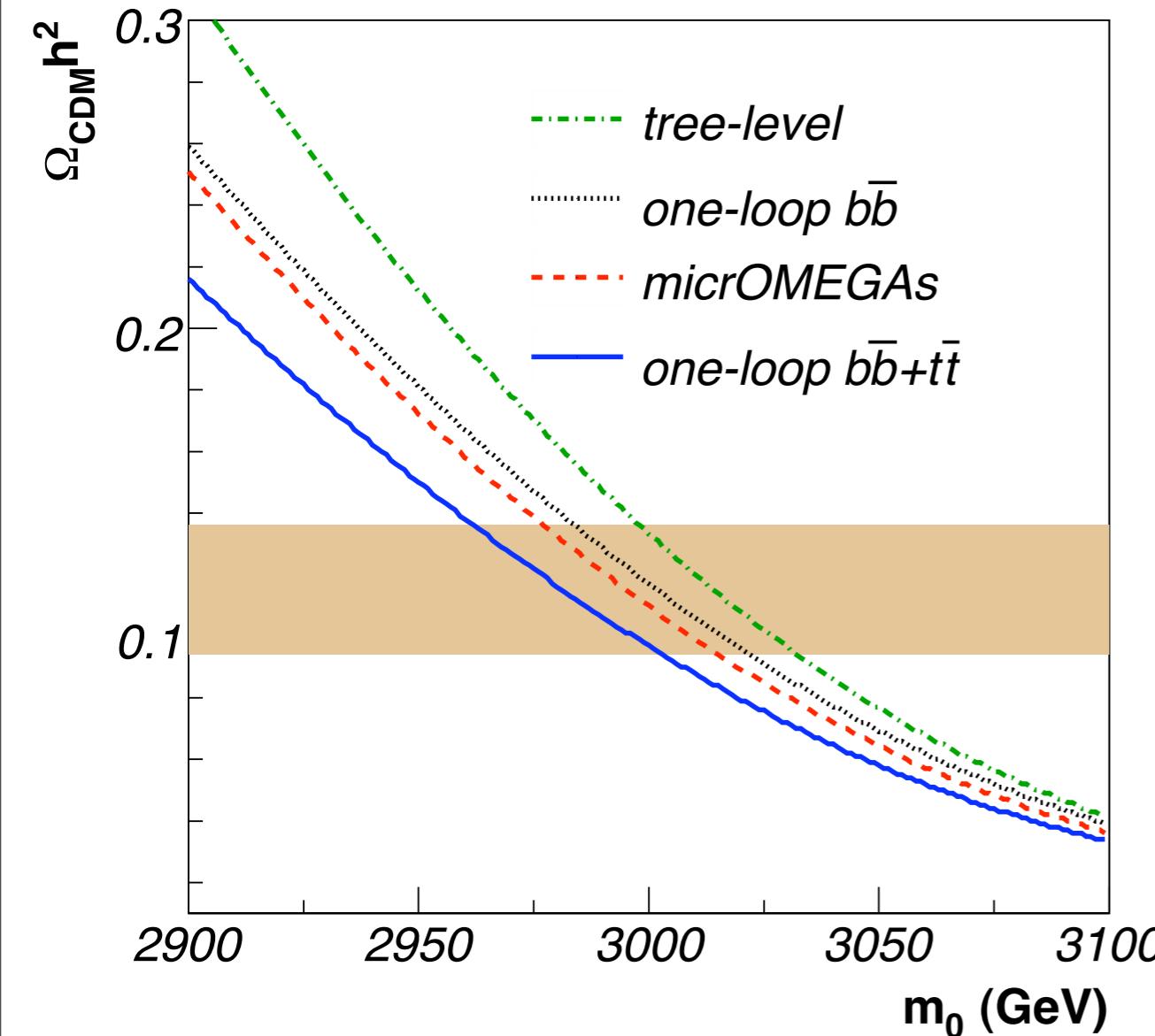


Important top and bottom quark contributions due to large top mass and large $\tan\beta$

Each cross section enhanced significantly by one-loop QCD and SUSY-QCD corrections

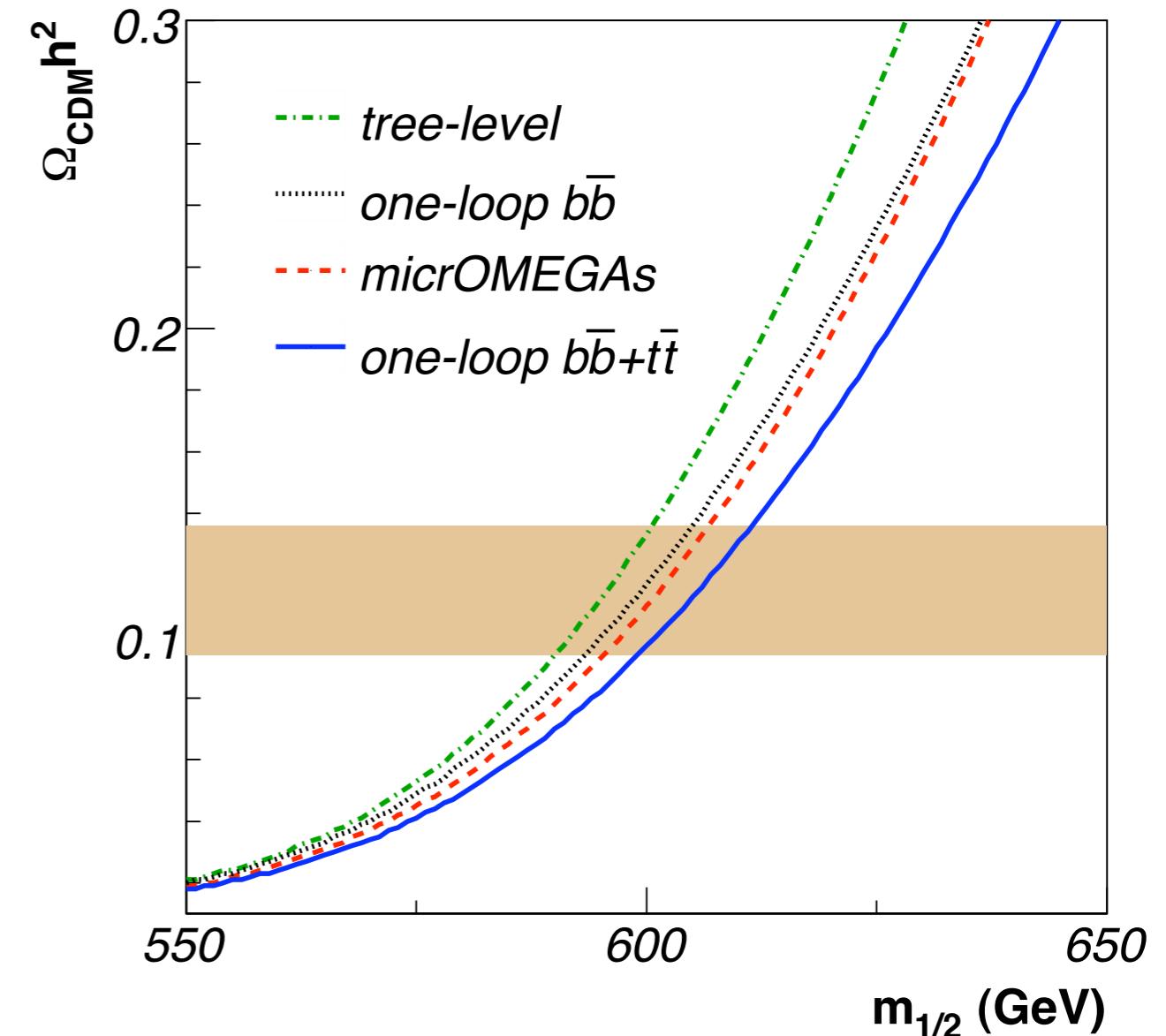
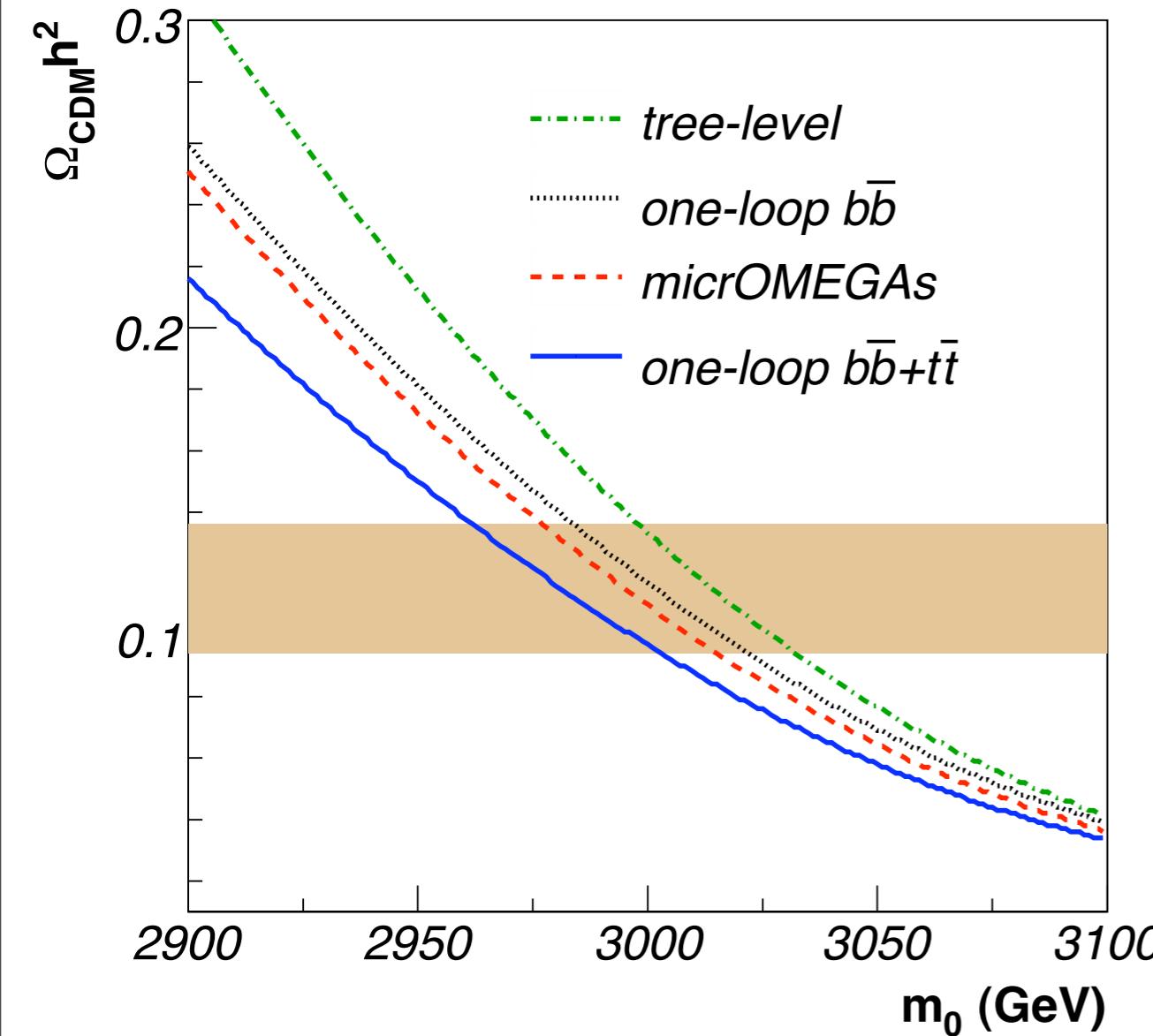
Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=3000$ GeV, $m_{1/2}=600$ GeV, $A_0=0$, $\tan\beta=50$, $\mu>0$



Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=3000$ GeV, $m_{1/2}=600$ GeV, $A_0=0$, $\tan\beta=50$, $\mu>0$

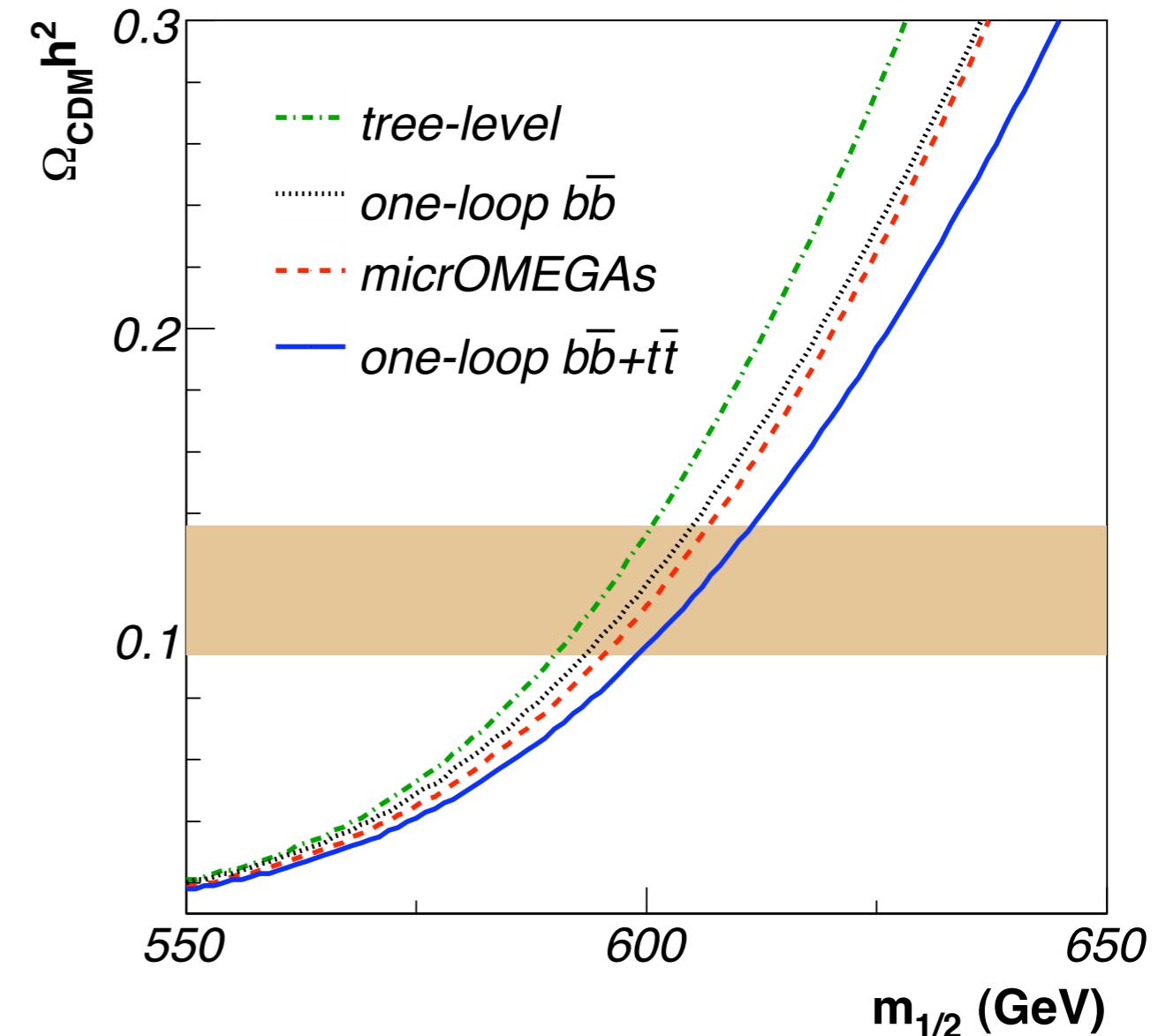
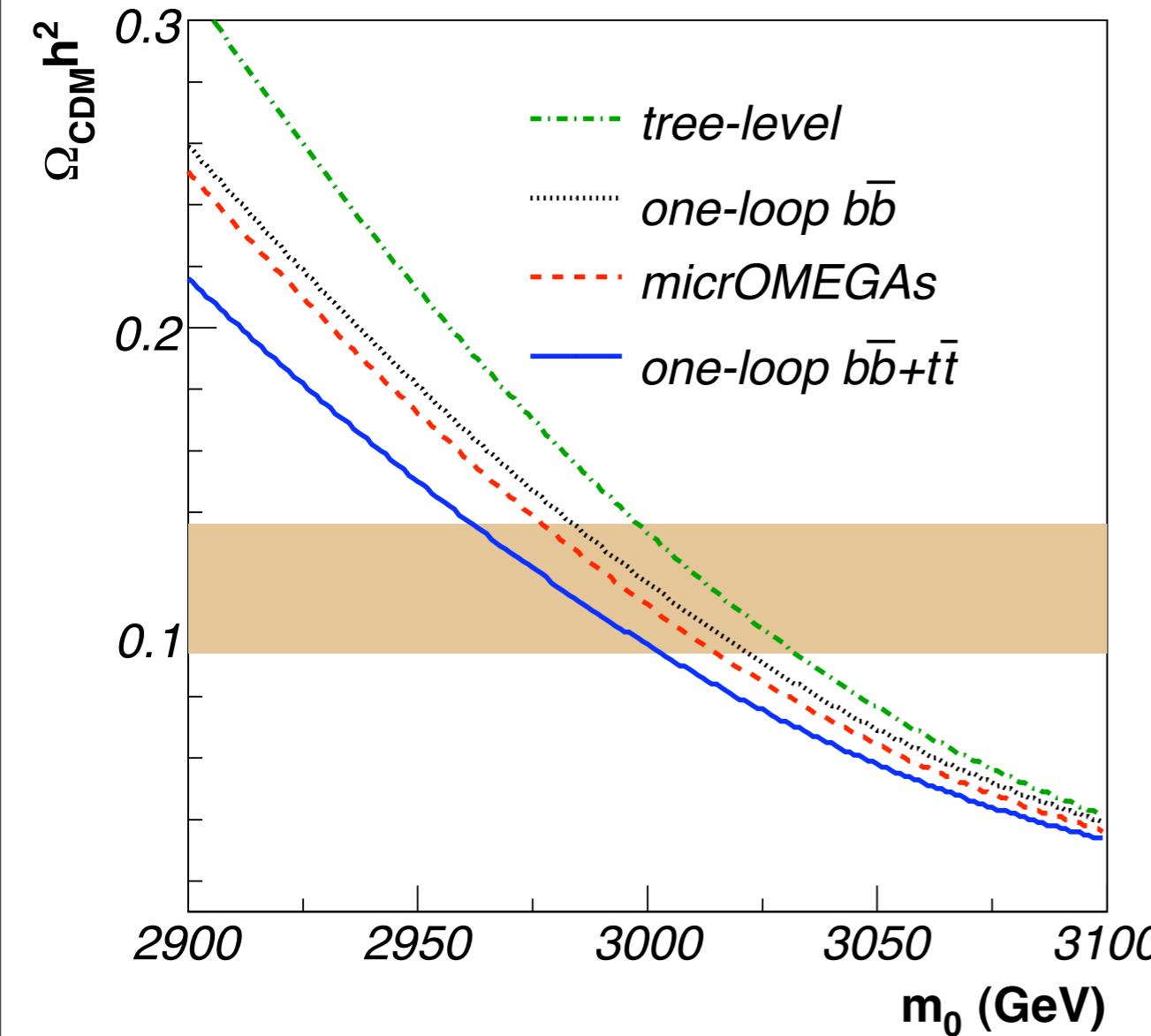


Bottom quark correction very important due to large $\tan\beta$

Top and bottom corrections shift relic density in the same direction

Numerical results: Focus point region (preliminary)

mSUGRA parameters: $m_0=3000$ GeV, $m_{1/2}=600$ GeV, $A_0=0$, $\tan\beta=50$, $\mu>0$



Bottom quark correction very important due to large $\tan\beta$

Top and bottom corrections shift relic density in the same direction

Favoured region shifted to smaller scalar masses m_0 and higher gaugino masses $m_{1/2}$

Conclusion

Conclusion

Dark matter annihilation is an interesting tool to constrain parameter space

- Complementary information with respect to collider and low-energy data
- Neutralino annihilation into “heavy” quark-antiquark pairs
- Radiative corrections important in the light of new cosmological precision measurements

Conclusion

Dark matter annihilation is an interesting tool to constrain parameter space

- Complementary information with respect to collider and low-energy data
- Neutralino annihilation into “heavy” quark-antiquark pairs
- Radiative corrections important in the light of new cosmological precision measurements

SUSY-QCD corrections have significant impact on extraction of SUSY mass parameters

- Neutralino annihilation into “heavy” quark-antiquark pairs
- Calculation of full one-loop QCD and SUSY-QCD corrections
- Significant impact on annihilation cross section and neutralino relic density in mSUGRA

Conclusion

Dark matter annihilation is an interesting tool to constrain parameter space

- Complementary information with respect to collider and low-energy data
- Neutralino annihilation into “heavy” quark-antiquark pairs
- Radiative corrections important in the light of new cosmological precision measurements

SUSY-QCD corrections have significant impact on extraction of SUSY mass parameters

- Neutralino annihilation into “heavy” quark-antiquark pairs
- Calculation of full one-loop QCD and SUSY-QCD corrections
- Significant impact on annihilation cross section and neutralino relic density in mSUGRA

Outlook / Perspectives:

- Study scenarios with important contributions from light quark final states
- Include corrected cross section also in DarkSUSY and compare...
- Application to indirect dark matter detection...
- Implementation of electroweak corrections...