

A new avenue to nuclear fragmentation in a dissipative quantum mean field

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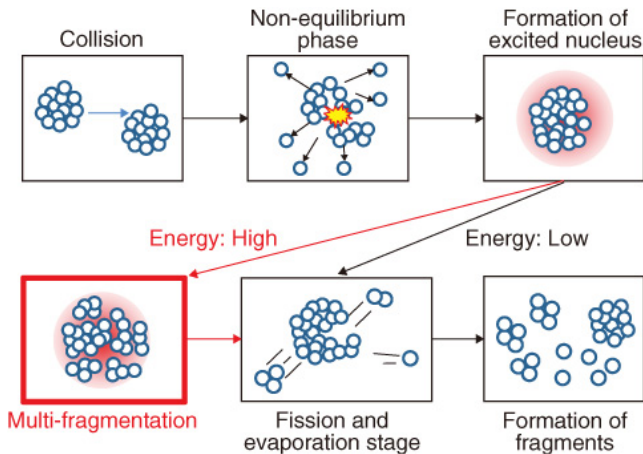
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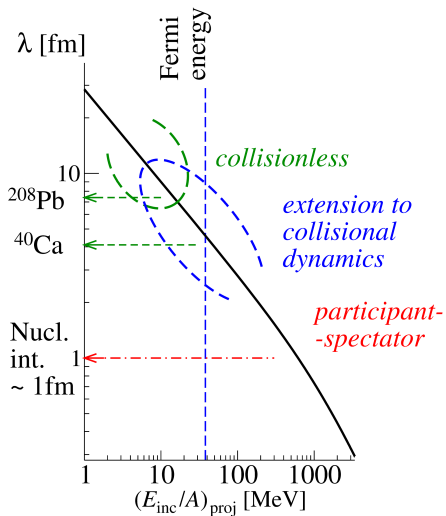
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Introduction



[OGAWA, T. ET AL., JAEA R&D REVIEW 2013,85]

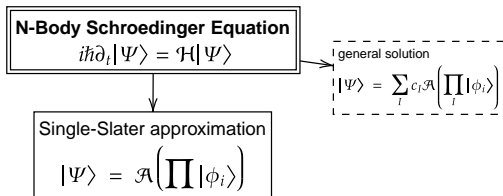
Overview



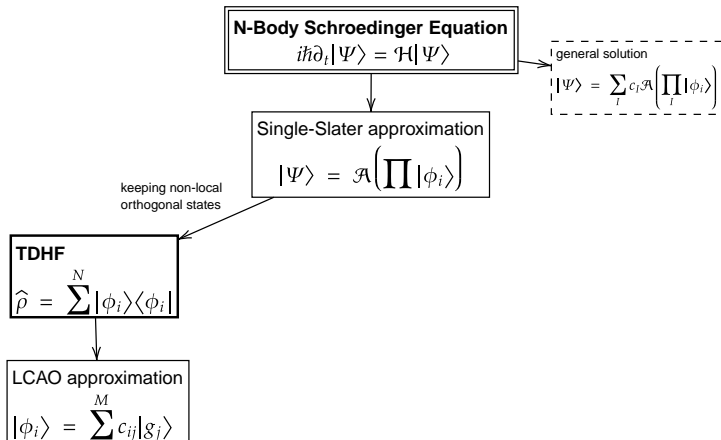
- Collective behaviour at low energy (TDHF)
- Two body nucleonic interaction at intermediate to high energy (MD)
- Interplay between collective and dissipative effects at Fermi energy (Boltzmann)
- Search for theory that can cover Fermi to low energy to study **fragmentation formation!**

Theoretical Background

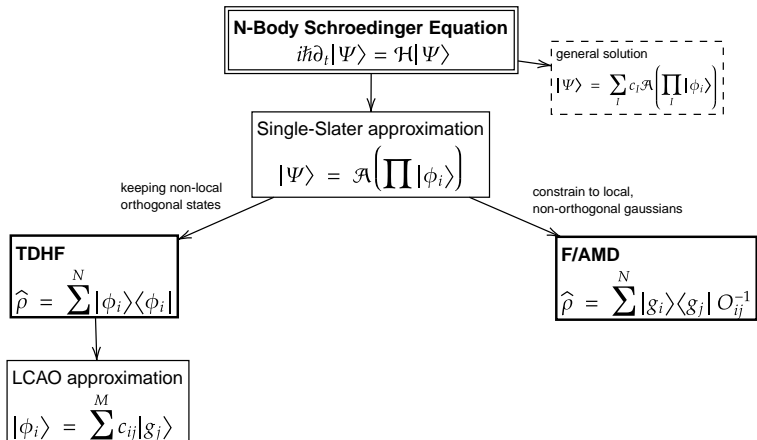
Connections of microscopic models



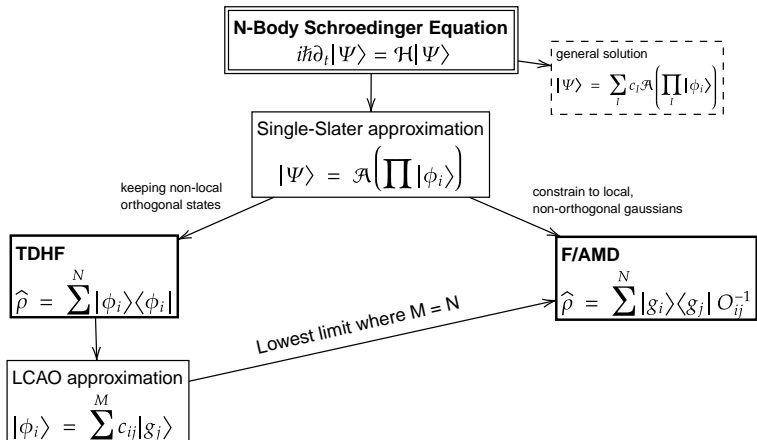
Connections of microscopic models



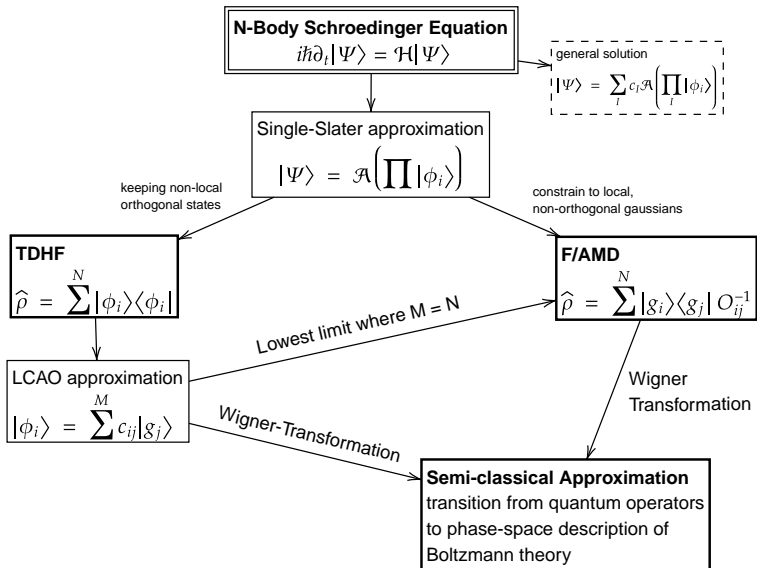
Connections of microscopic models



Connections of microscopic models



Connections of microscopic models



One-Body Dynamics

Stochastic-Collisional-TDHF

$$i\hbar\partial_t\rho_1 \approx [h_1, \rho_1] + \bar{I}_{\text{coll}} + \delta I_{\text{coll}}$$

- Approximation of BBGKY
- \bar{I}_{coll} average collision term
- δI_{coll} fluctuating collision term

Boltzmann-Langevin

$$\partial_t f = \{h, f\} + I_{\text{UU}} + \delta I_{\text{UU}}$$

- Wigner-transform of SC-TDHF
- \bar{I}_{coll} and δI_{coll} contribution replaced by UU analog

Boltzmann-Langevin One-Body (BLOB)

$$\partial_t f - \{h, f\} = I_{\text{UU}} + \delta I_{\text{UU}} = g \int \frac{d\mathbf{p}}{h^3} \int W(AB \leftrightarrow CD) F(AB \rightarrow CD) d\Omega$$

- A,B,C,D: extended equal-isospin phase-space portions

Model

Implementation

Decoherence approximation

[V. DE LA MOTA, IL NUOVO CIMENTO 41 C, 2018, 189]

$$\rho = \sum_i \omega_i |g_i\rangle \langle g_i|$$

Coherent state parametrization

$$g_x(x) = \mathcal{N} \exp\{-a(\chi, \varphi)(x - x_0)^2 + ik_0(x - x_0)\}$$

$$g(\vec{r}) = g_x(x)g_y(y)g_z(z)$$

Mean-field evolution

$$i\hbar\partial_t g_i = \left(\frac{\hbar^2 \vec{k}^2}{2m} + V^{\text{HF}}(\rho) \right) g_i$$

Equation of Motion

$$\dot{x} = \frac{\hbar k_x}{m}$$

$$\dot{k} = -\frac{1}{\hbar} \partial_x \langle V^{\text{HF}} \rangle$$

$$\dot{\chi} = \frac{4\hbar}{m} \gamma \chi$$

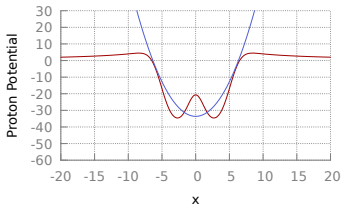
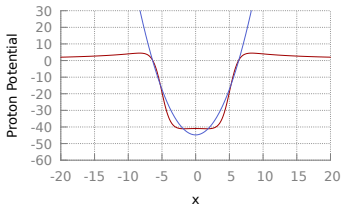
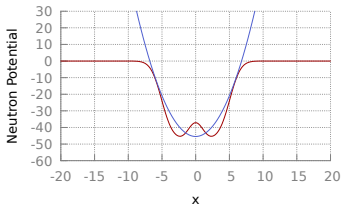
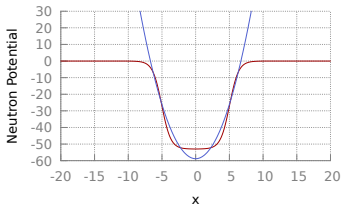
$$\dot{\gamma} = \frac{\hbar}{8m} - \frac{2\hbar\gamma^2}{m} - \frac{1}{\hbar} \partial_x \langle V^{\text{HF}} \rangle$$

V^{HF} : Hartee-Fock potential, Skyrme-type effective force, density and isospin dependent and with surface term (Skt5)

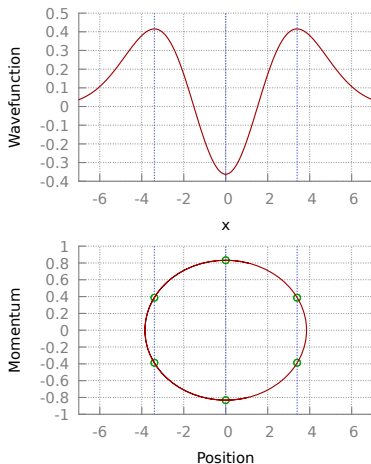
Initialization

Self-consistency iteration routine of $[h_1, \rho_1] = 0$

Demonstration on 58-Ni



Decomposition of the coherent states

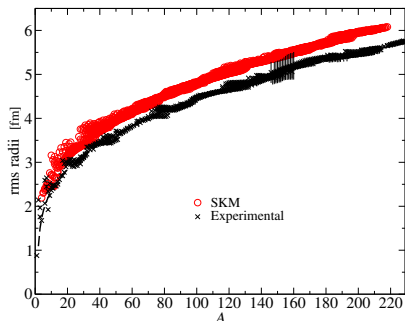
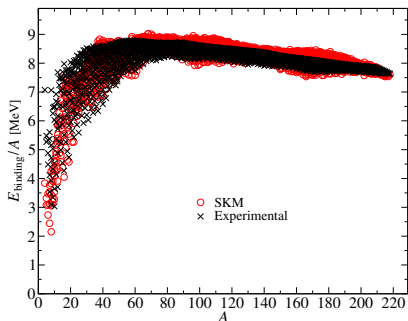


- 1 Find the peaks of reference wavefunction
- 2 Wigner-transform the wavefunction and fit with the Wigner-transform of the coherent states at the corresponding positions

$$\omega_i = \frac{1}{N_{CS,i}}$$

Remark: The Wigner-transform of the coherent states take the form of a rotated gaussian in phasespace

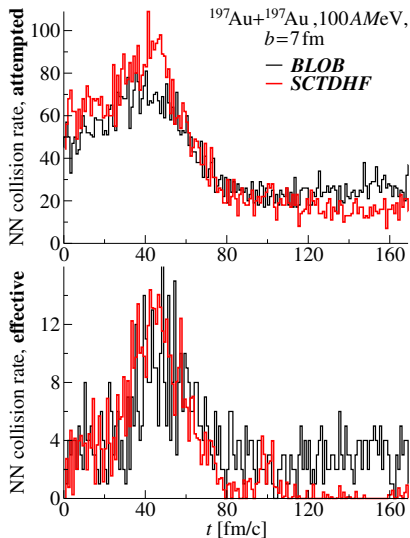
Binding Energy



- slight deviation of the binding energy due to the initialization with the harmonic oscillator states
- Systematic deviation to the experimental size of the nuclei
- Work in progress: cooling procedure

Collision-Fluctuation effects

- **scalarproduct** provides naturally a phasespace metric
- Follow a **BLOB** prescription for the fluctuating collision term
 - 1 Sample a nucleonic representation
 - 2 define nucleonic wavepacket through scalarproduct
 - 3 check for possible collision to attempt
 - 4 perform nucleonic collision considering the overlap in the final state
 - 5 work in progress: cluster recognition



Conclusion

Conclusion

- Range of different mechanism in the final state of heavy-ion collisions are difficult to model
- Current model achieves low energy, due to the stability of TDHF
- Improvements would give TDHF-like model which can give fragments
- Work in progress: Due to the setting we can assign more consistent properties to the residues
- Goal: Nuclear features emerge dynamically without extrapolation in the fragments