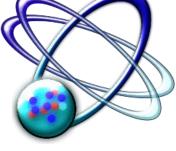


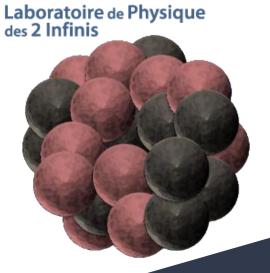
Nuclear structure and α radioactivity



PhyNet

Florian MERCIER





I. The nuclear many body problem

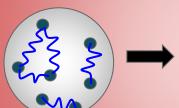
II. Building an interaction

III. Many body interacting problem

IV. Alpha and cluster radioactivity

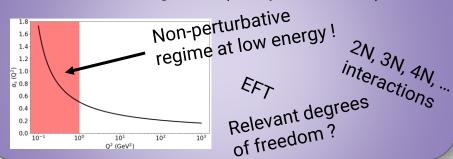
Tackling the nuclear many body problem

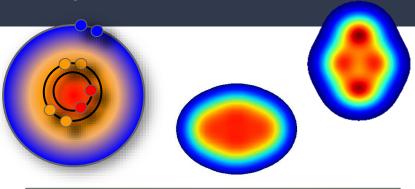
Quantum many body interacting problem!



Many (strongly) interacting particles but not enough for statistical approximation TOOOOO hard to be solved exactly

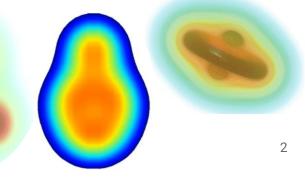
Interaction coming from (non-perturbative) QCD!





Huge phenomenology!

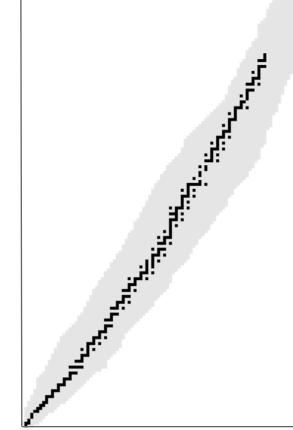
Many emerging properties: deformation, radioactivity, halo, neutron skin, superfluidity, clustering, excitations, ...



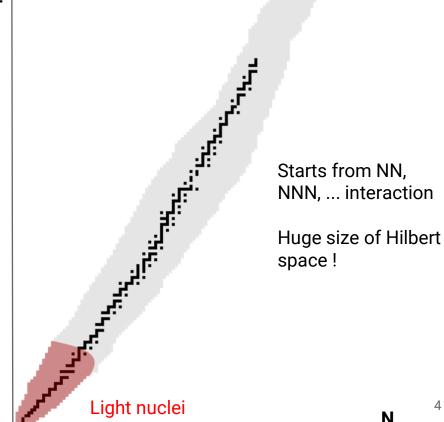
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Nuclear landscape

Z



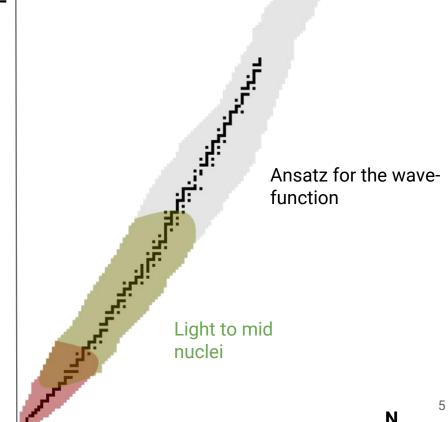
Nuclear landscape



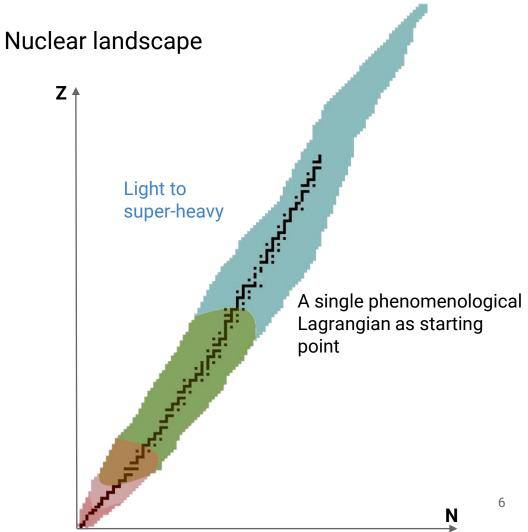
- Ab Initio
- Configuration interaction

Nuclear landscape

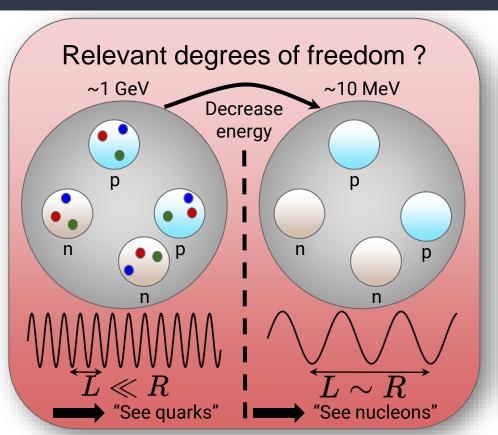


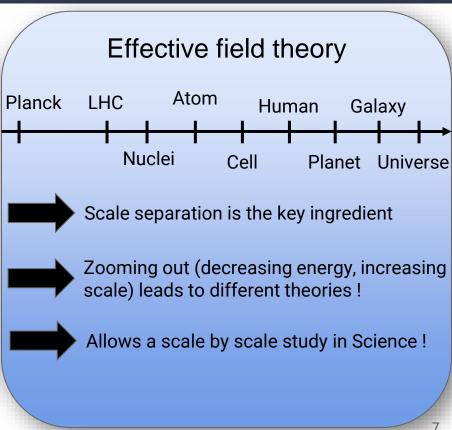


- Ab Initio
- Configuration interaction
- Energy density functional



Building an interaction : what is the idea?





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Many body interacting problem

Strongly interacting many body system

VERY HARD



One-body problem involving independent quasi-particles

Mean-field

Idea:
$$H = \left[\sum_i \frac{p_i^2}{2m}\right] + \left[\sum_{ij} V_{ij} + \sum_{i < j < k} V_{ijk} + \ldots\right]$$

$$= \left[\sum_i \frac{p_i^2}{2m} + U_i\right] + \left[\sum_{ij} V_{ij} + \sum_{i < j < k} V_{ijk} + \ldots - \sum_i U_i\right]$$

$$= H_0 + \mathcal{V}_{res}$$
Such that $|H_0| \gg |\mathcal{V}_{res}|$
As much physics as possible Justify perturbative approach

8

Hartree-Fock theory

Symmetries

$$H = H_0 + \mathcal{V}_{res}$$
 with a certain \c{U}_i

We know that the Hamiltonian *H* possesses some symmetries

$$[H,N]=0$$
 Particle number conservation

$$[H,J]=0$$
 Invariance under rotation

Let us choose the same symmetries for H_0

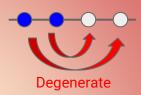
$$[H_0, N] = 0$$
 $[H_0, J] = 0$

Problem

Perturbation theory



Improvement by going at higher order



degenerate state

$$oldsymbol{E} = \mathcal{E} + \sum_{n>0}^{\bullet} rac{|\langle \phi_n | \mathcal{V}_{res} | \phi_0
angle|^2}{\mathcal{E} - \mathcal{E}_n} + \dots$$



We are missing an essential ingredient in \mathcal{V}_{res} !

Change U_i to change the physical content

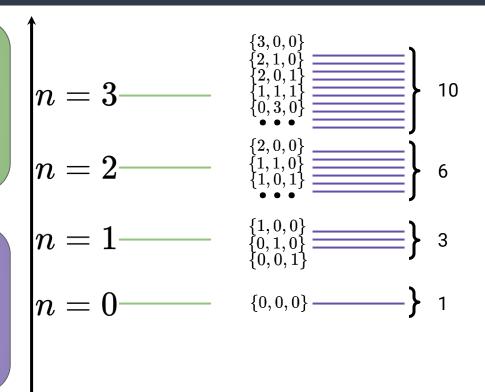
How to raise degeneracies?

The isotropic (symmetric) harmonic oscillator

$$egin{aligned} E_n^{sym} &= \hbar \omega (n_x + n_y + n_z + 3/2) \ &= \hbar \omega (n + 3/2) \end{aligned}$$

The anisotropic (asymmetric) harmonic oscillator

$$egin{aligned} E_{n_xn_yn_z}^{asym} &= \hbar(\omega_x n_x + 1/2) + \hbar(\omega_y n_y + 1/2) \ &+ \hbar(\omega_z n_z + 1/2) \end{aligned}$$





Hartree-Fock-Bogoliubov theory

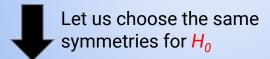
Symmetries

$$H = H_0 + \mathcal{V}_{res}$$
 with a certain \c{U}_i

We know that the Hamiltonian *H* possesses some symmetries

$$[H,N]=0$$
 Particle number conservation

$$[H,J]=0$$
 Invariance under rotation



$$[\mathbf{H_0}, N] = 0 \quad [\mathbf{H_0}, J] = 0$$

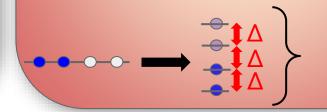
Problem Solution

Assume $igl[H_0', N igr]
eq 0$ Particle number not conserved anymore!

Particles can create pairs which behave as single degree of freedom

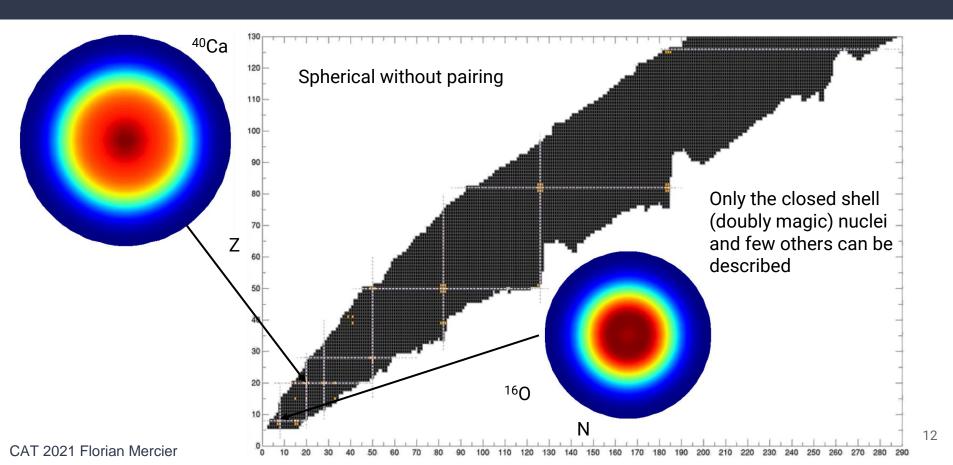
As for spin pairing, it leads to a splitting of the energy and a new minimum energy configuration!

$$E = \epsilon \pm \Delta$$



Non degenerate anymore!

Impact of symmetry breaking: no breaking



Impact of symmetry breaking: U(1) breaking

What are we missing then?

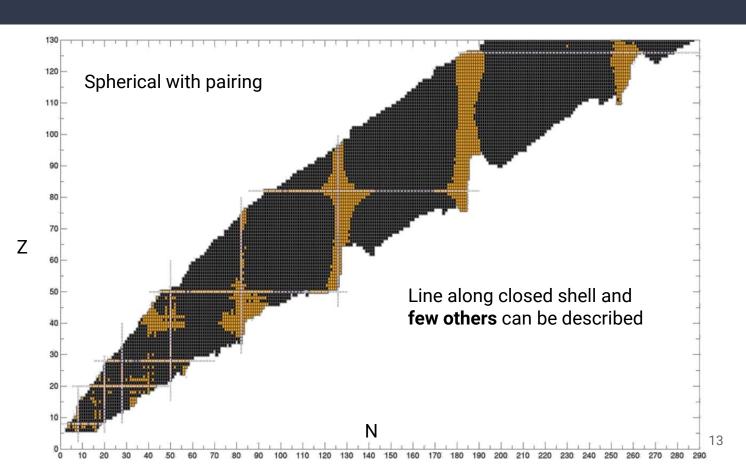


Let us break other symmetries!

$$[extbf{\emph{H}}_0,J]
eq 0$$

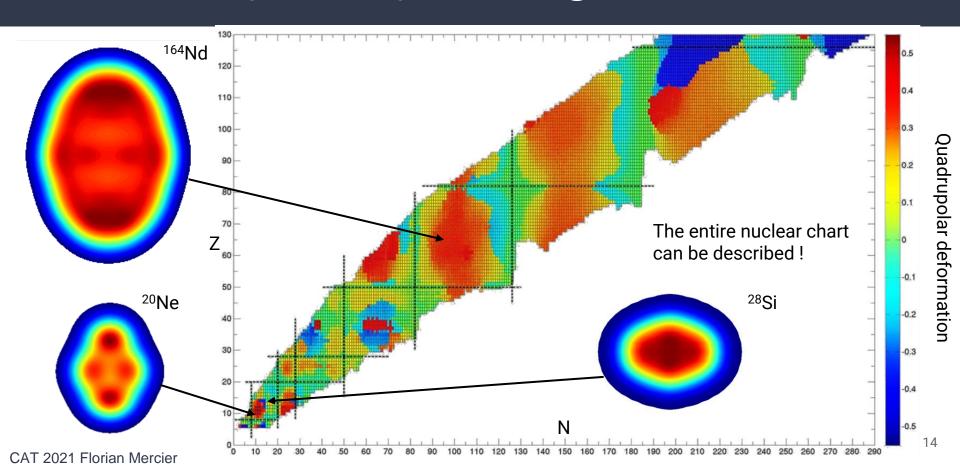


Nuclei can now be deformed

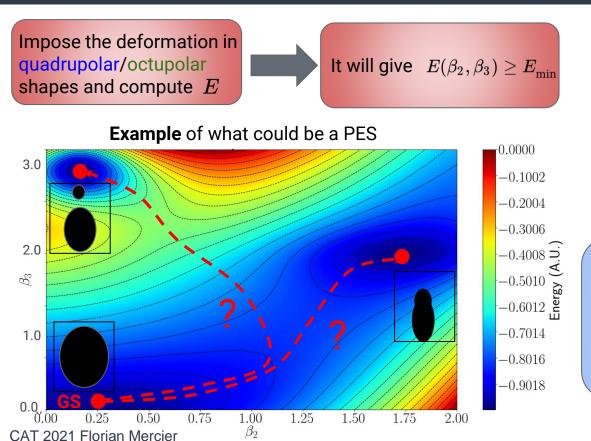


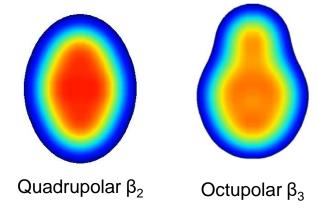
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Impact of symmetry breaking: U(1) and SO(3)



Deformations and Potential Energy Surface

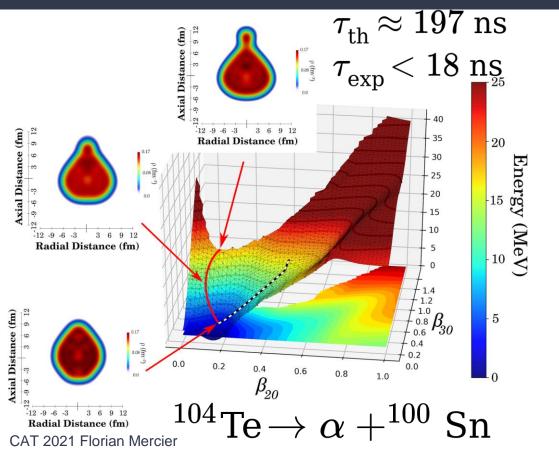


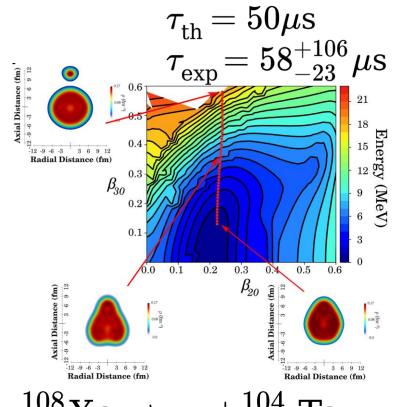


- Can shapes evolve from GS to α emission ?
- Is it possible to compute the associated probability?

Alpha decay chain $108 \, \mathrm{Xe} \rightarrow 104 \, \mathrm{Te} \rightarrow 100 \, \mathrm{Sn}$

F. Mercier and al., Phys. Rev. C 102, 011301(R) (2020)





Other results for heavier nuclei

F. Mercier and al., under review