

# UHE neutrino-nucleon cross section from IceCube-Gen2

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UNIVERSITY OF  
COPENHAGEN



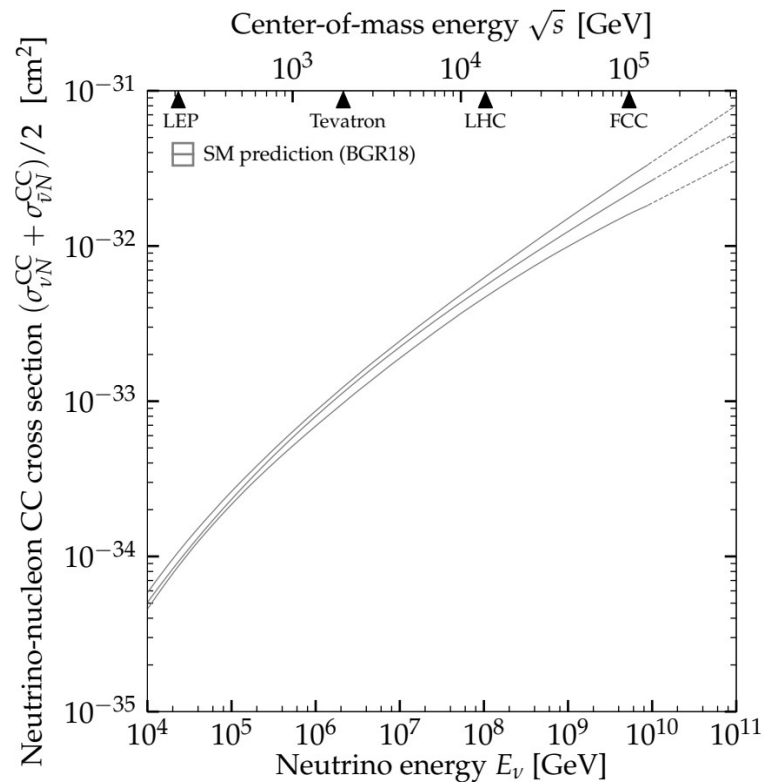
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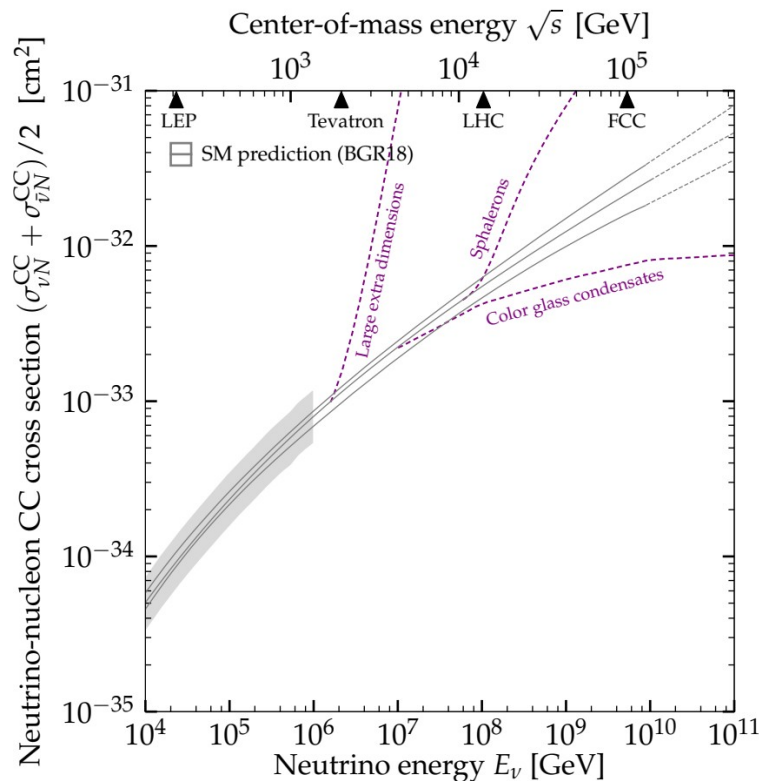
# MAIN OBJECTIVES

- **Our main goal** is to prepare the most **detailed prediction** of the measurement capabilities of the **neutrino-nucleon cross section at the  $\sim EeV$**  scale with the next generation of neutrino telescopes

# Neutrino-nucleon cross section measurements today

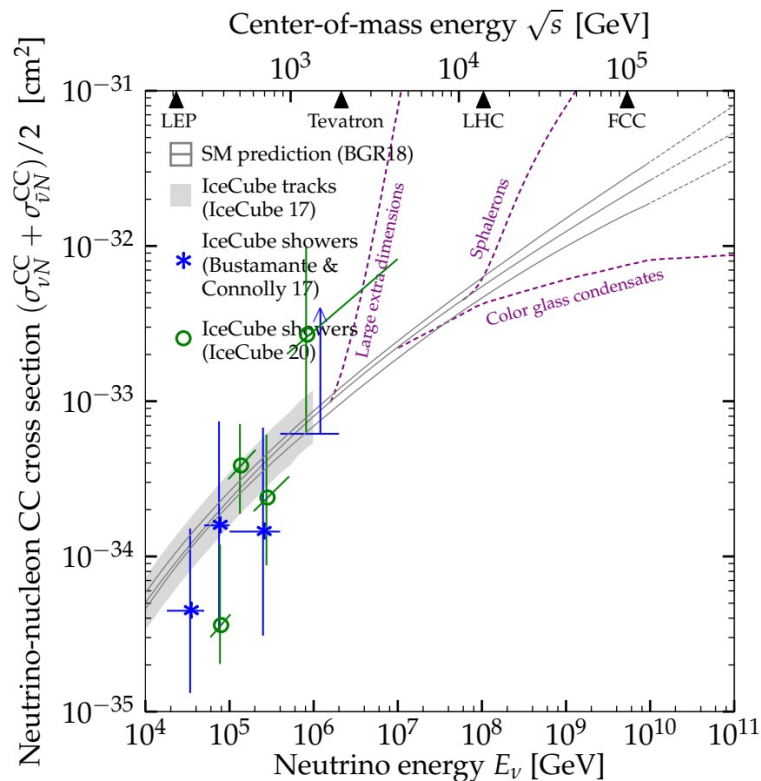


# Neutrino-nucleon cross section measurements today



Potential exciting new Physics at ultra-high energies

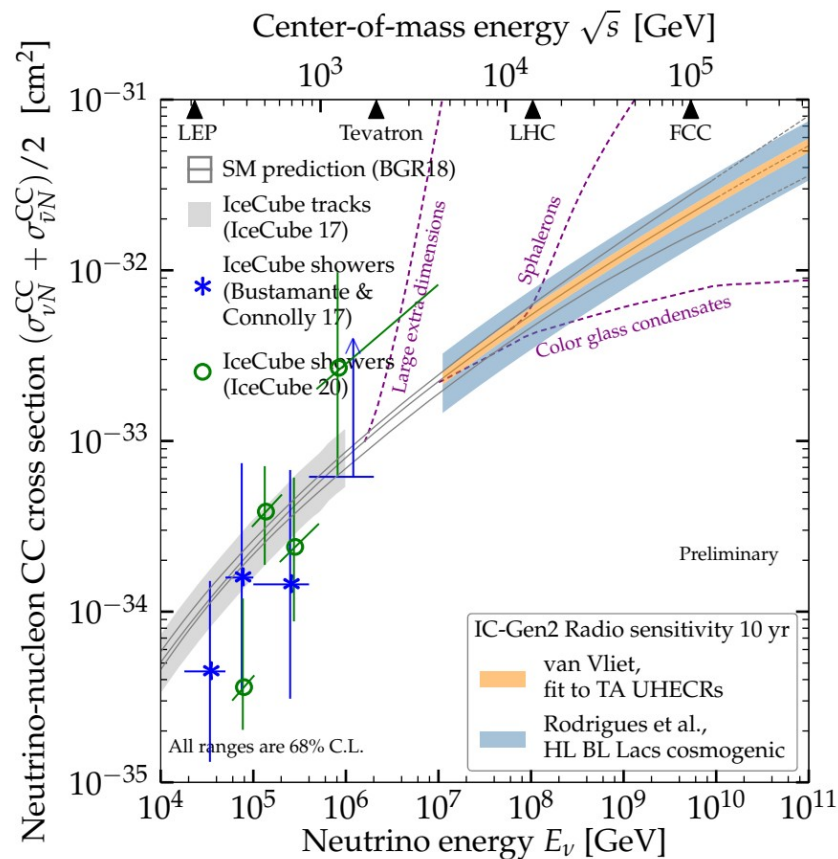
# Neutrino-nucleon cross section measurements today



# THE METHOD: Model-dependent $\rightarrow$ independent

Relax model dependency in stages:

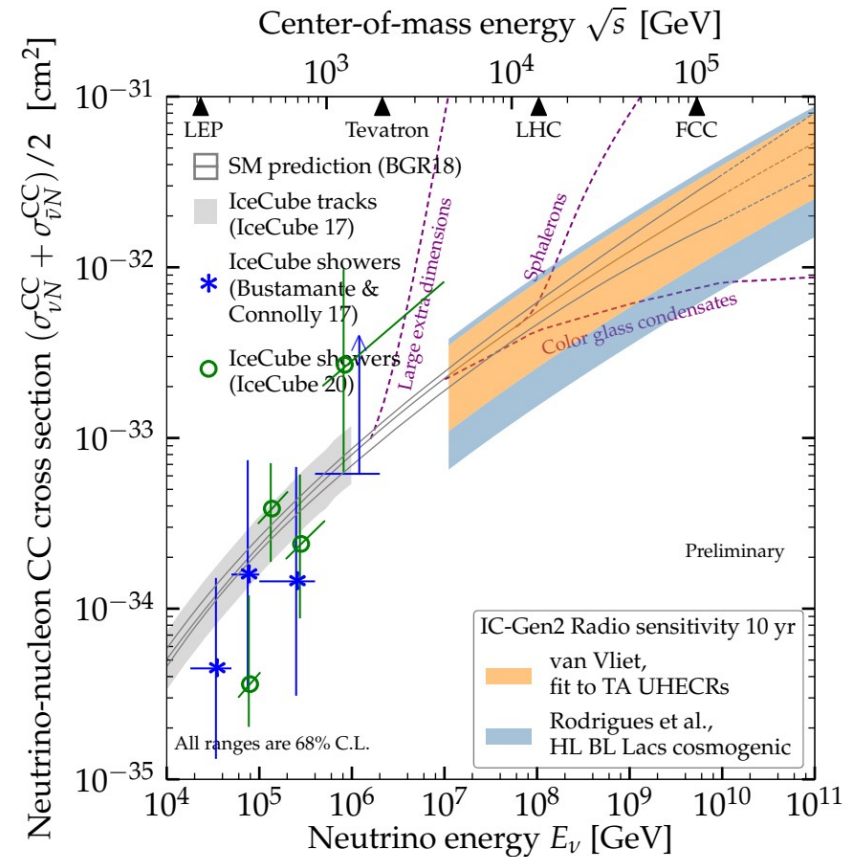
- 1) **Perfectly known UHE $\nu$  flux.**
- 2) Add a normalization uncertainty to the assumed UHE $\nu$ .
- 3) Measure the flux and the cross section simultaneously.



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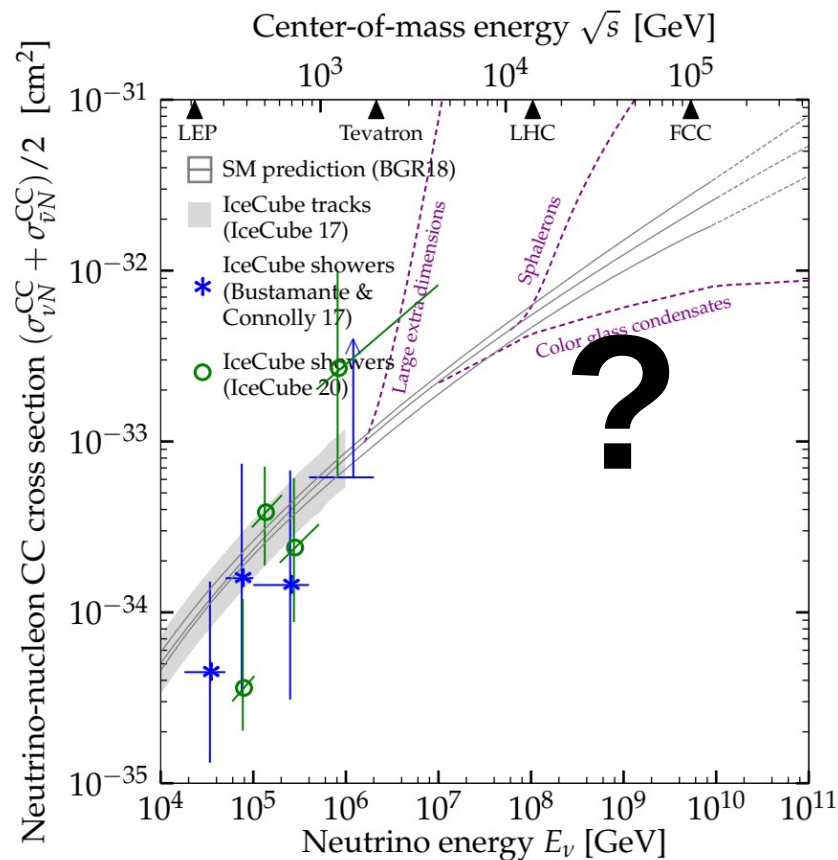
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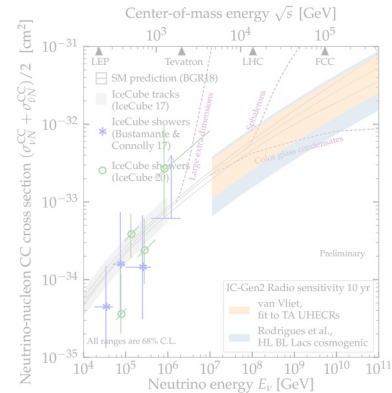
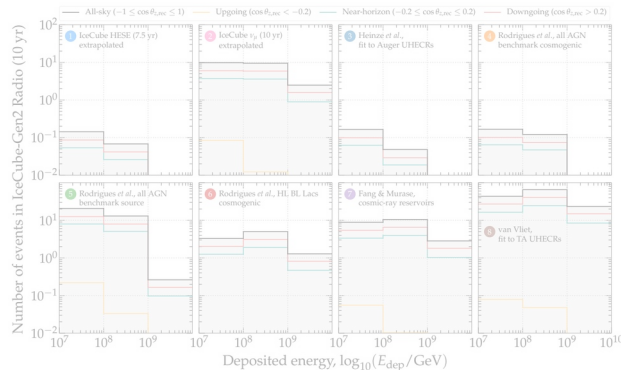
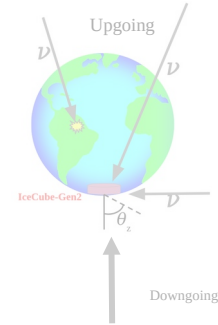
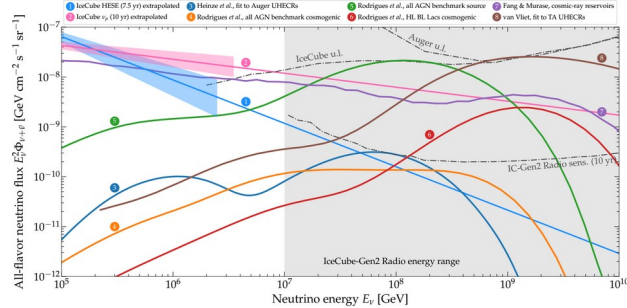
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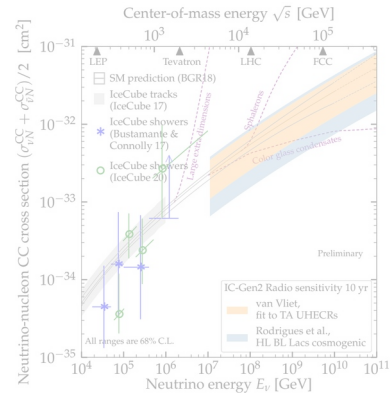
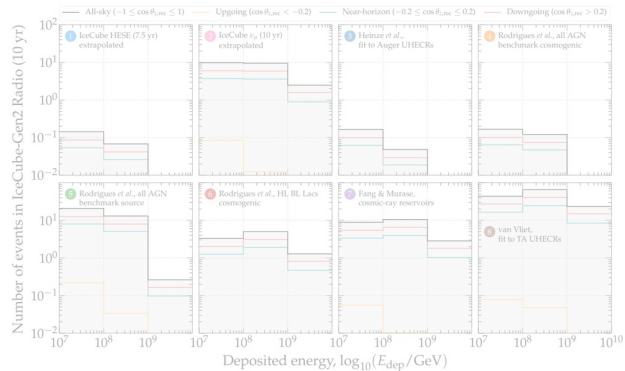
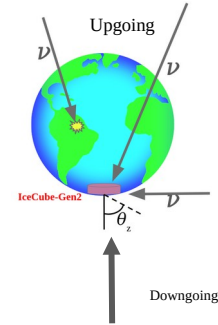
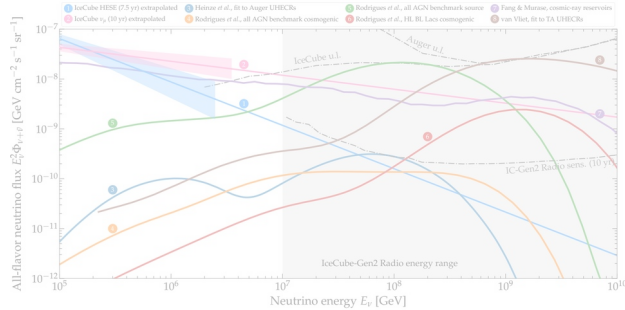




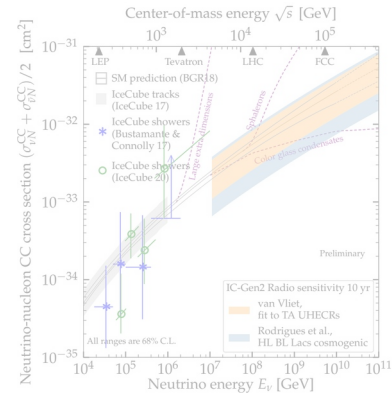
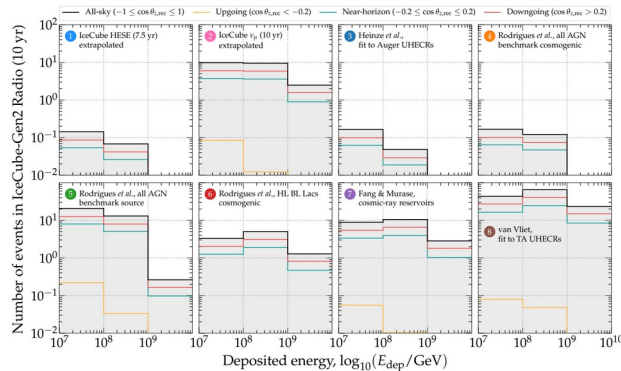
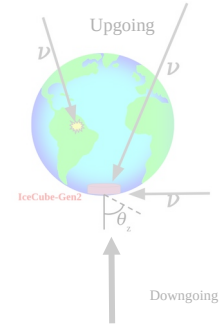
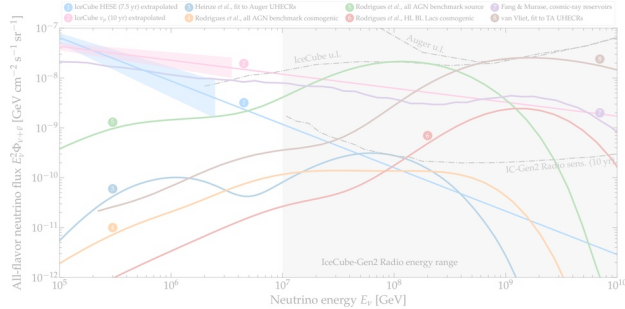
# WORKFLOW: Each step of this work



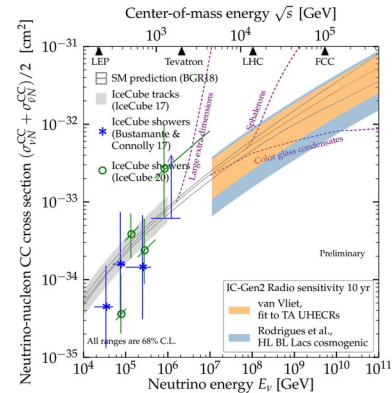
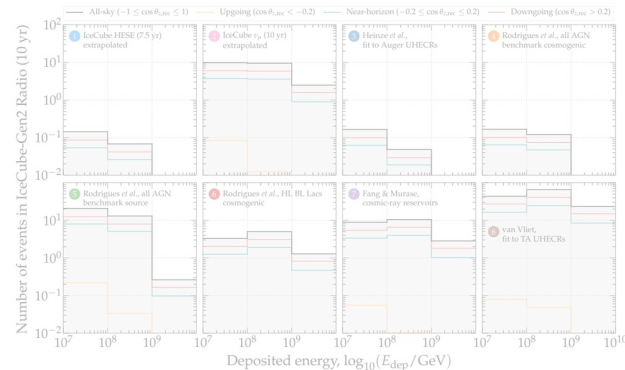
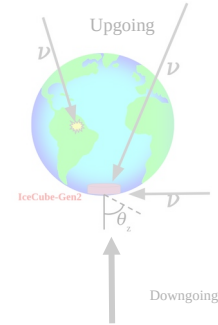
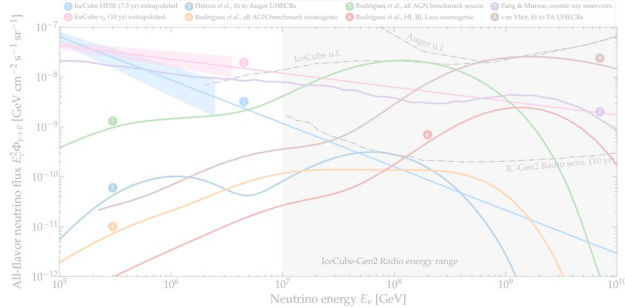
# WORKFLOW: Each step of this work



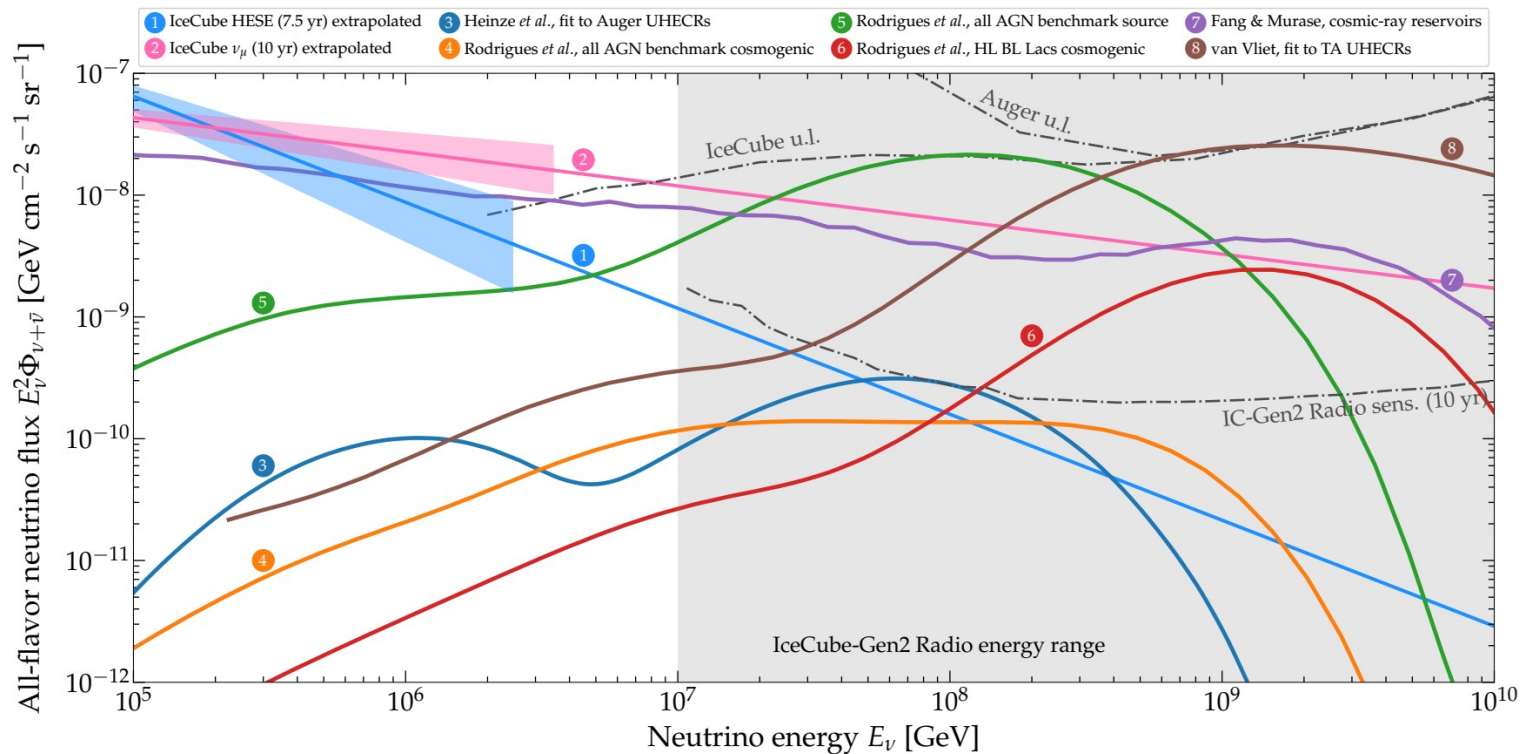
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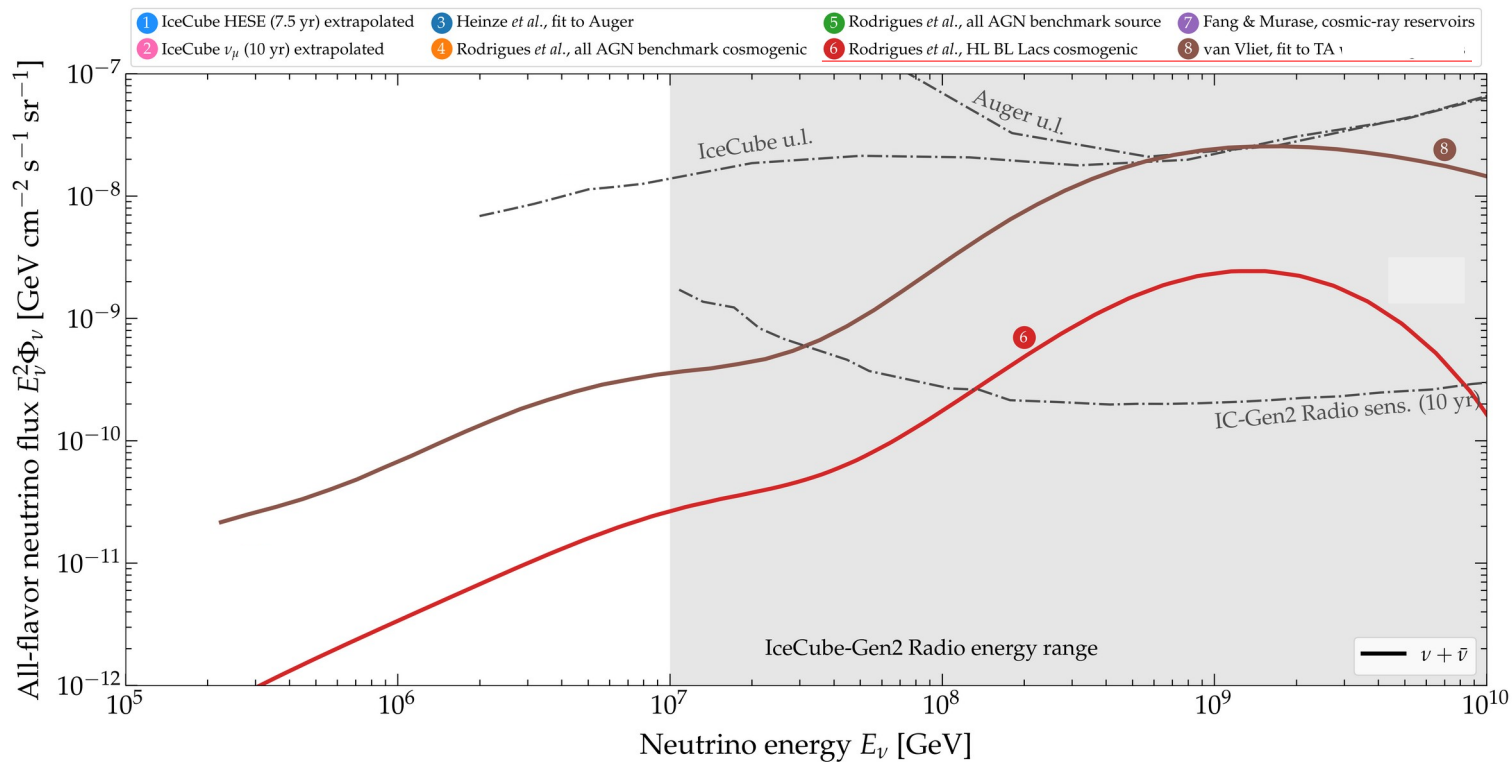
# WORKFLOW: Each step of this work



# The UHE neutrino flux

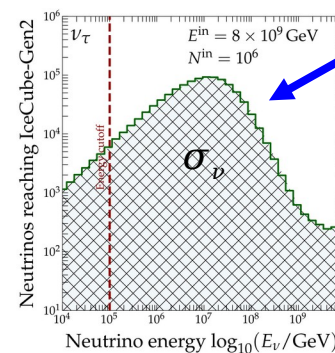
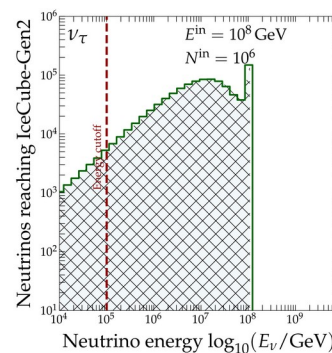
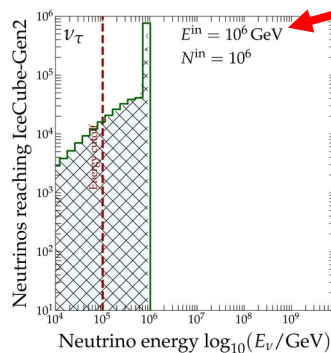
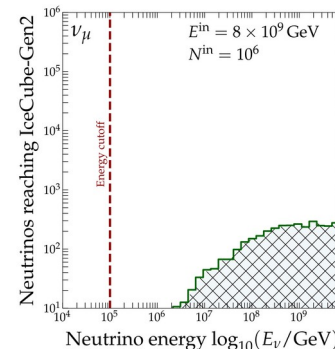
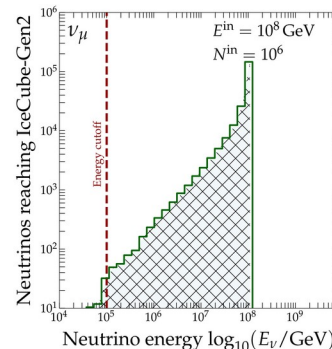
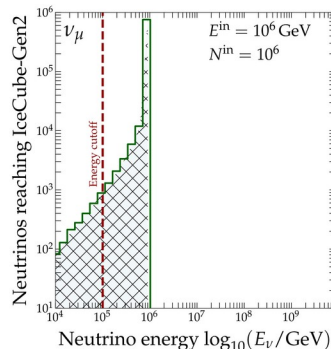


# The UHE neutrino flux



# NuPropEarth: an in-Earth neutrino propagation tool

- Monte Carlo neutrino propagation
- Leading interaction DIS
- + sub-leading contributions
- Most updated  $\sigma_{\nu N}$  predictions
- Earth density: PREM model

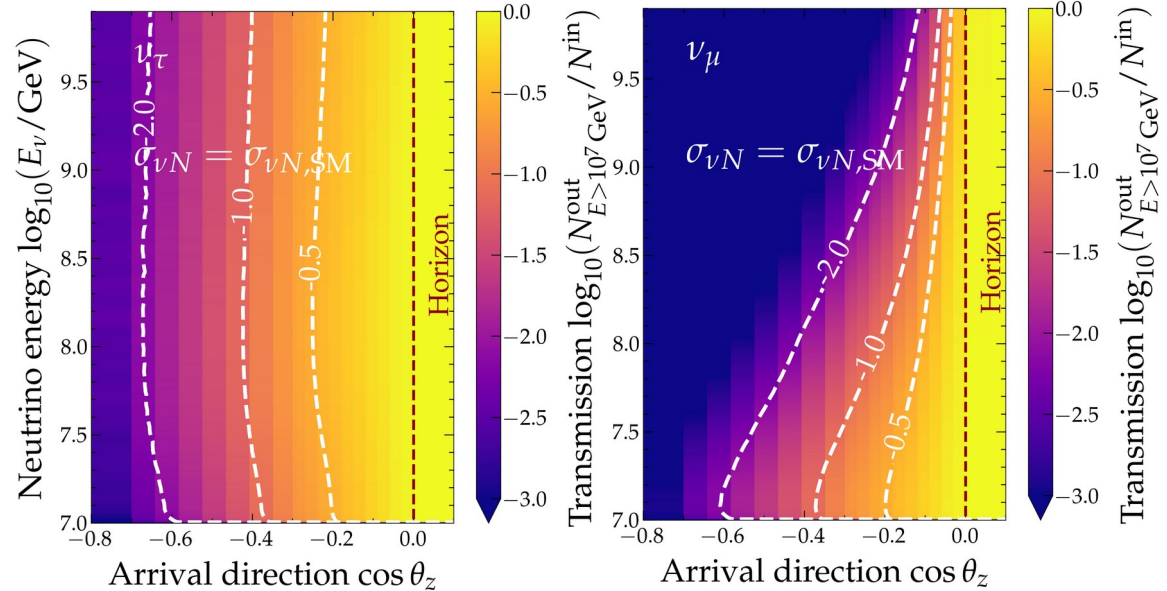
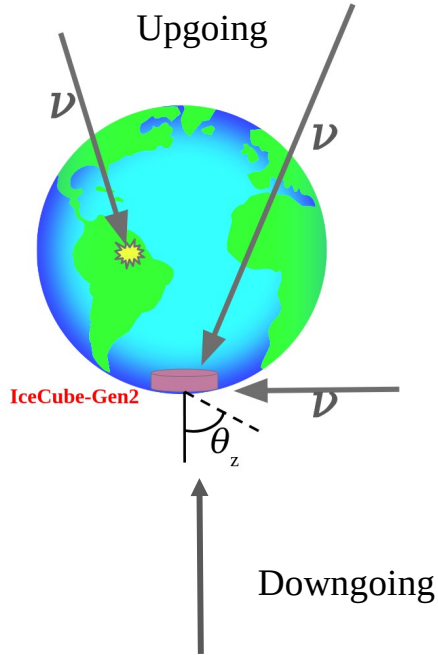


Garcia et. al., JCAP 2020

Higher injected neutrino energy

# Monochromatic neutrino beam propagation

- Monochromatic neutrino beams
- Various directions at different energies
- Store how they cascade down at the detector.







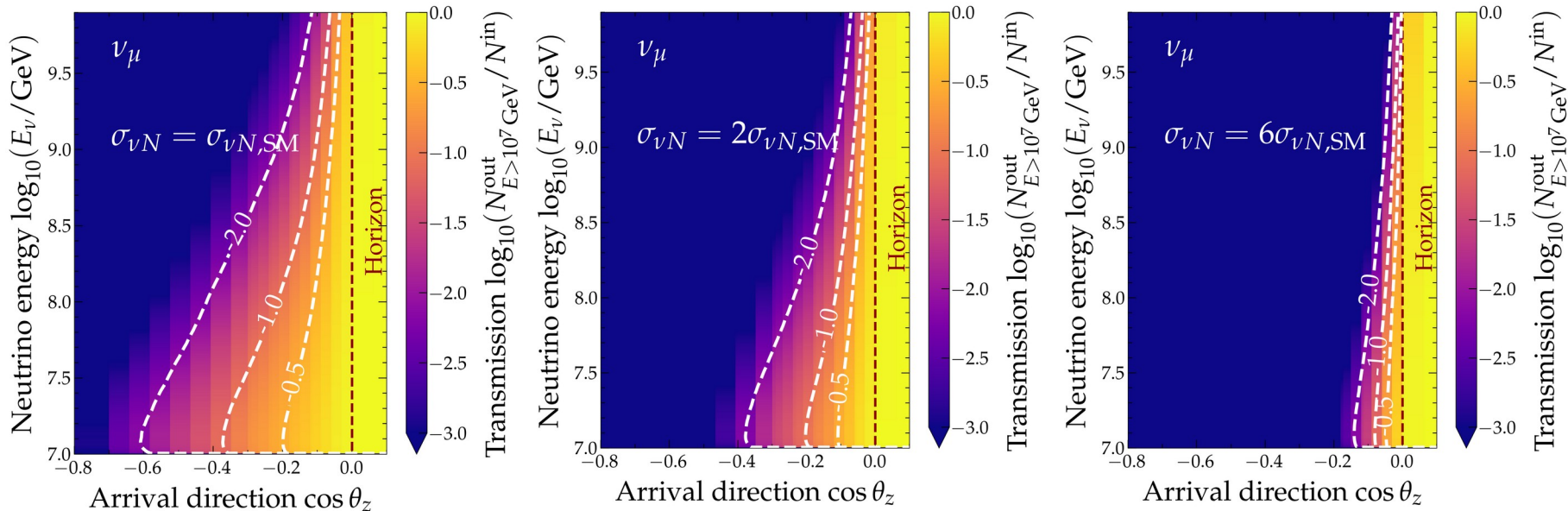
# Propagate neutrino for $\sigma \neq \sigma_{\text{SM}}$ , $\mathbf{f} = [0.1, 10]$

Cross section  $\sigma_{\text{SM}}$

$\sigma = \sigma_{\text{SM}}$

$\sigma = 2\sigma_{\text{SM}}$

$\sigma = 6\sigma_{\text{SM}}$



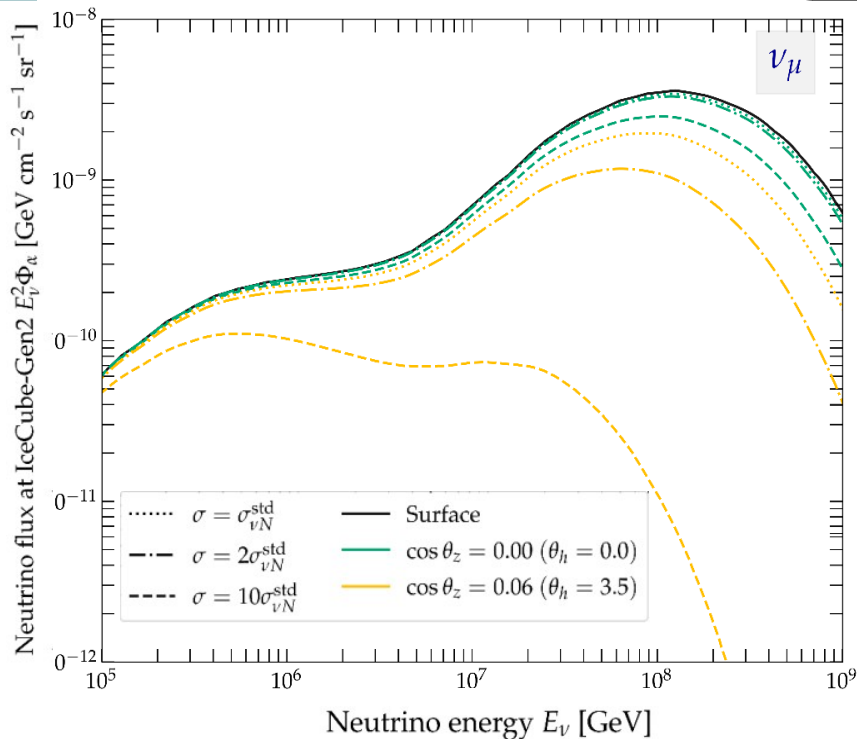
# Flux propagation from precomputed tables

## Flux propagation:

We sum the transmission histograms for every energy weighted by the value of the flux at each energy

## The result:

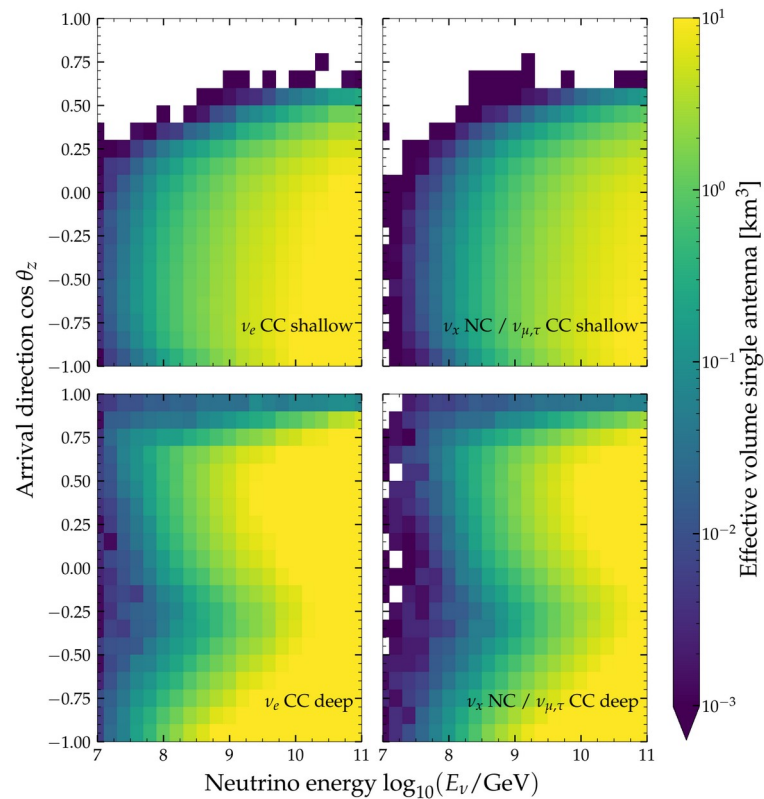
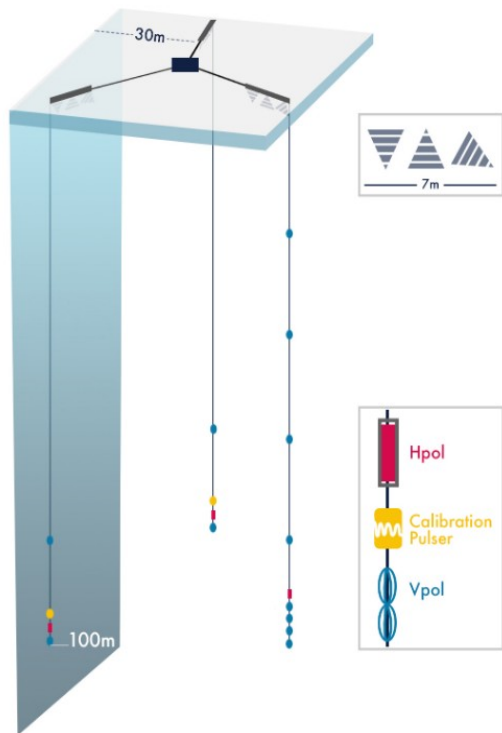
Propagated neutrino flux for each flavor at different directions



# Detector response: Effective volume

## IceCube-Gen2 Radio

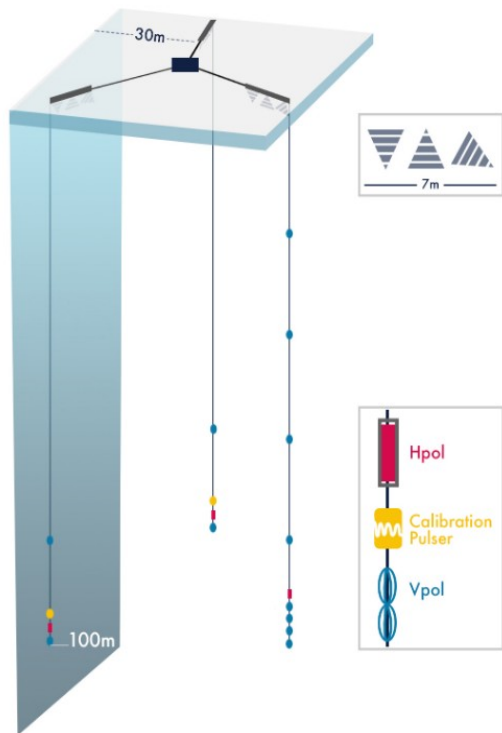
M.G. Aartsen et. al., *J. Phys. G.* 2021



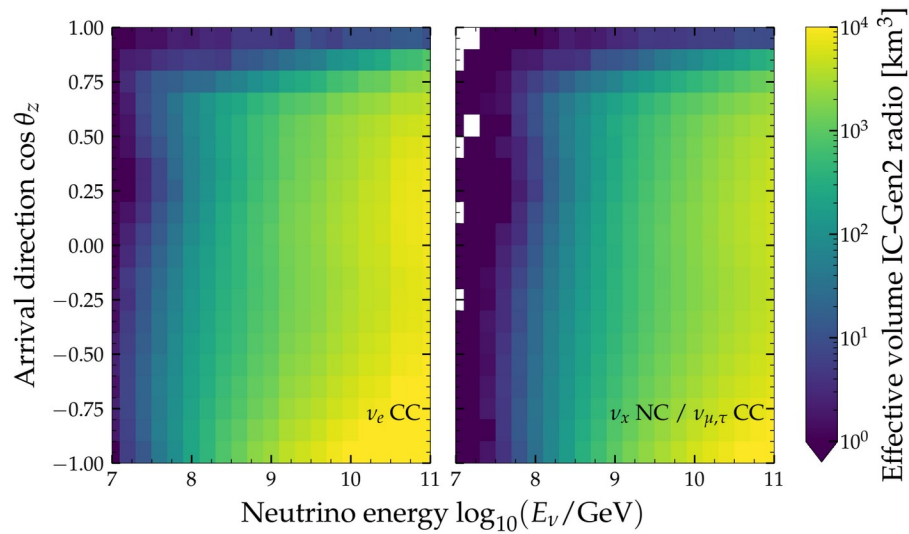
# Predicted event rate at IceCube-Gen2 radio

## IceCube-Gen2 Radio

M.G. Aartsen et. al., *J. Phys. G.* 2021



- The full detector is made of a total of 144 hybrid (shallow + deep) and 169 shallow-only stations



# Predicted event rate at IceCube-Gen2 radio

$$\text{Event rate: } N \sim T \Phi_{\nu} \sigma_{\nu N} e^{-\tau_{\nu N}(E_{\nu}, \theta_z)}$$

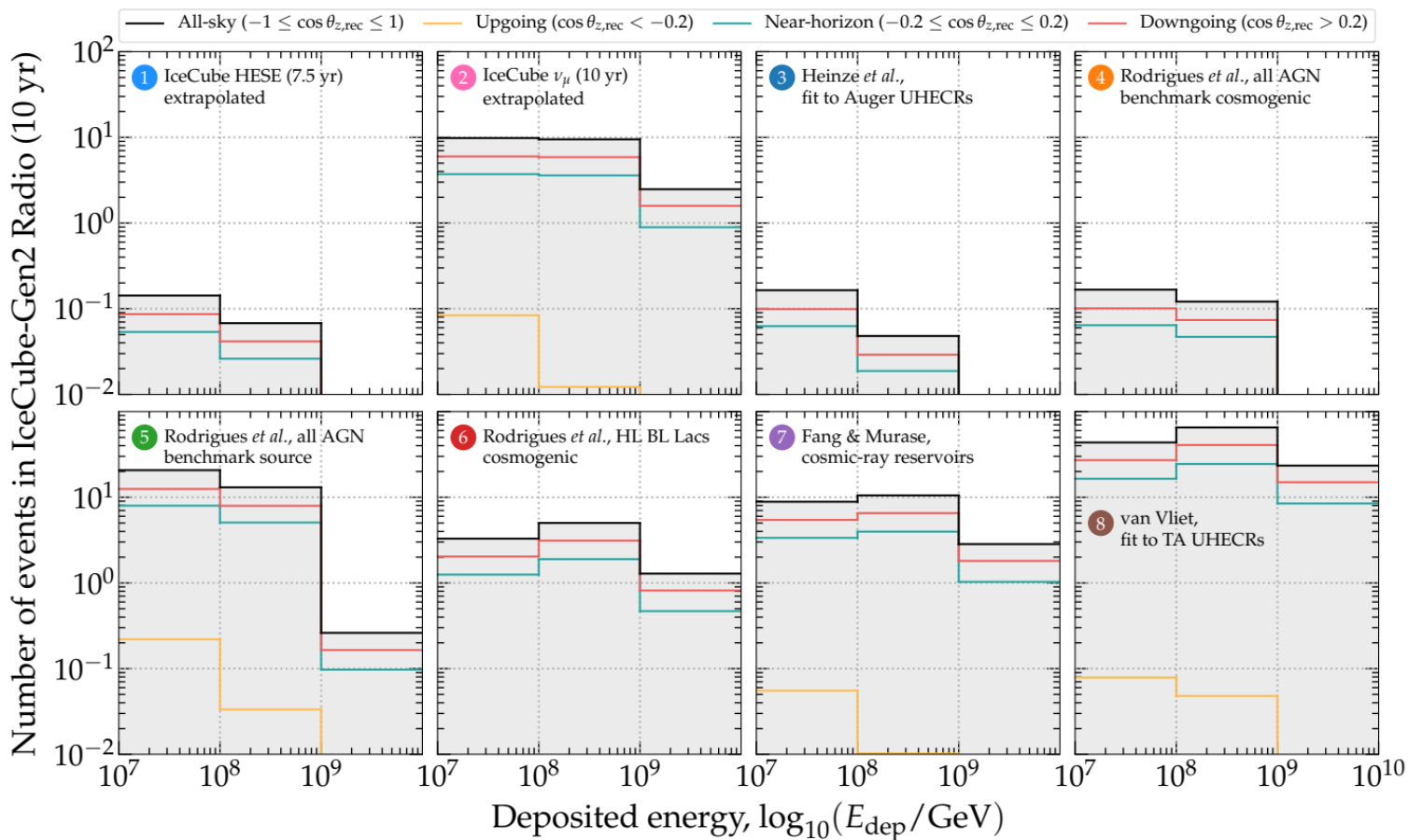
Increases the probability of a neutrino detection

Decreases probability of a neutrino transmission

- We account for energy and direction smearing through the resolution functions\*
- We write the event rate in terms of experimentally measurable quantities ( $E_{\text{dep}}$ ,  $\cos \theta_{z, \text{rec}}$ )
- IC-Gen2 effective volume simulated for this study without Earth attenuation (we include it ourselves)

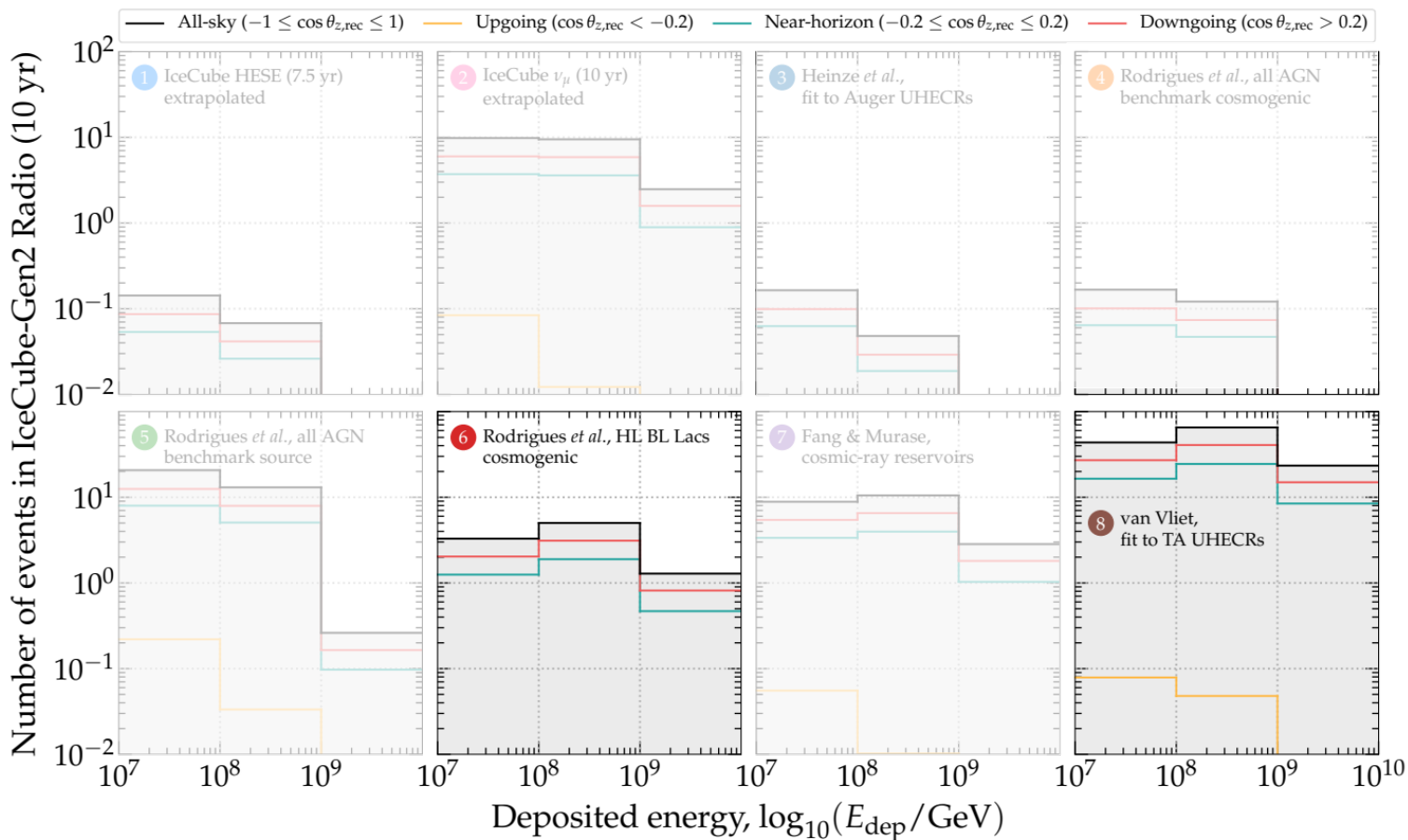
\* See backup for full expression

# The predicted event rates after 10 years of IC-Gen2



\* See backup for full expression

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# How well can we measure $\sigma_{\nu N}$ ? - Perfectly known flux

- Bayesian analysis
- Unbinned Poissonian likelihood:

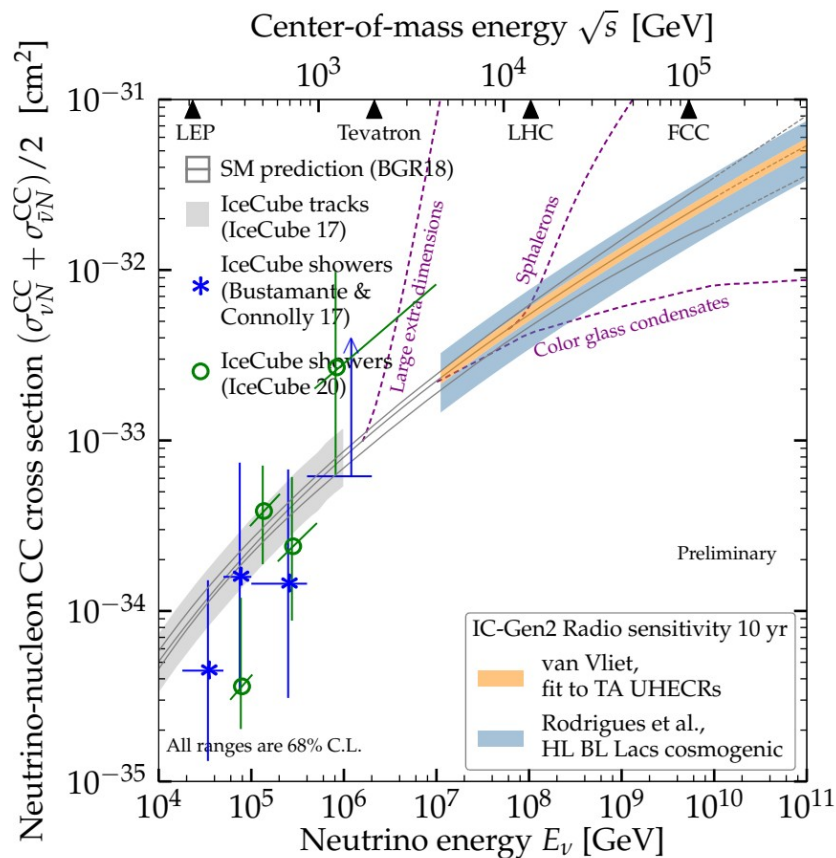
$$\mathcal{L}(f, N_{\text{obs}}) = \frac{e^{-[N(f)+N_{\text{bkg}}]} [N(f) + N_{\text{bkg}}]^{N_{\text{obs}}}}{N_{\text{obs}}!}$$

- Fix  $N_{\text{obs}}$  and maximize for  $f = \sigma/\sigma_{SM}$
- Credible regions from the posterior
- Flat  $f$  prior

$N_{\text{obs}}$  = Observed events assuming  $f=1$

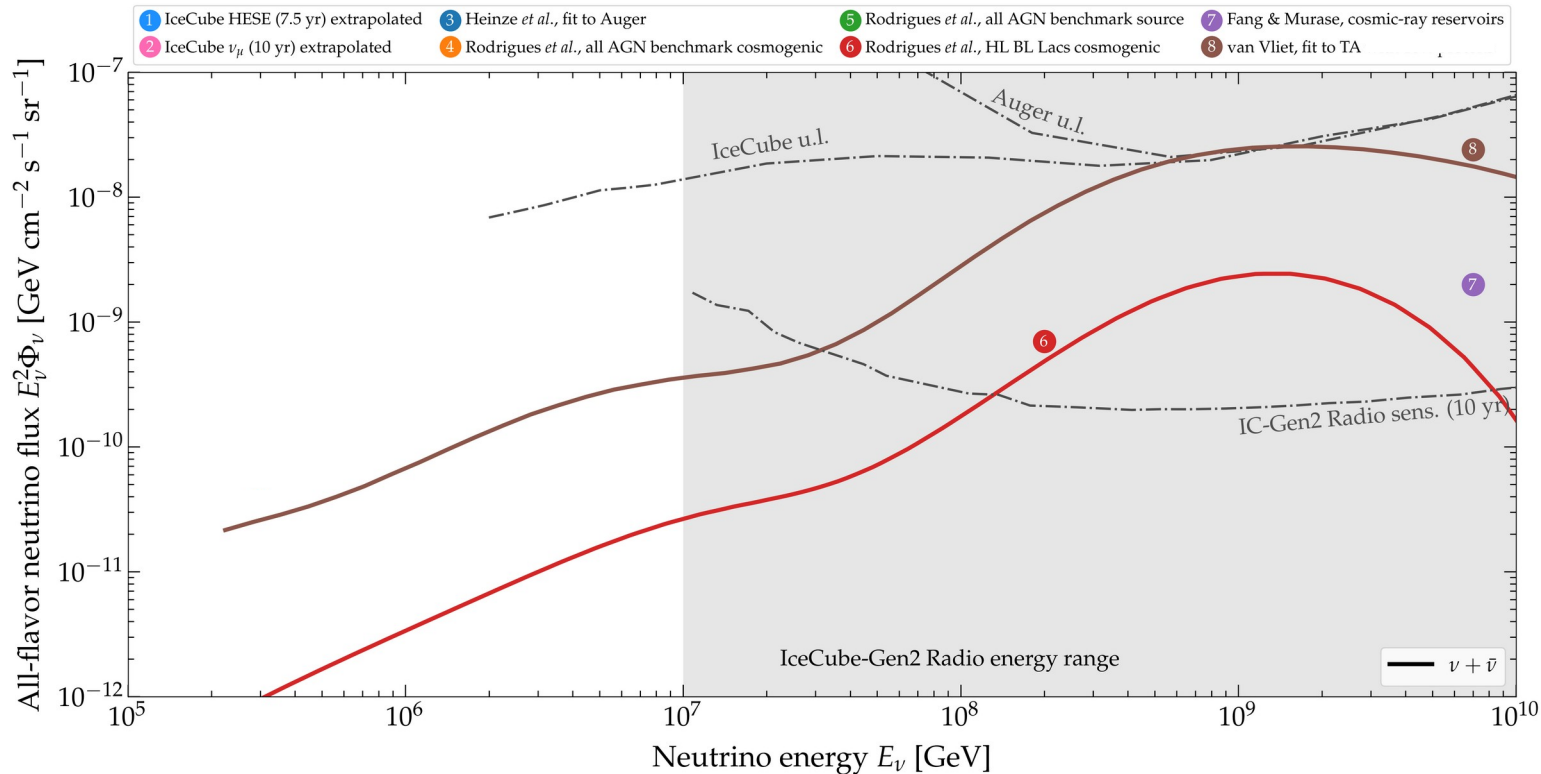
$N_{\text{bkg}}$  = Background events (**atm muons**)

$N(f)$  = Predicted number of events for different  $f$



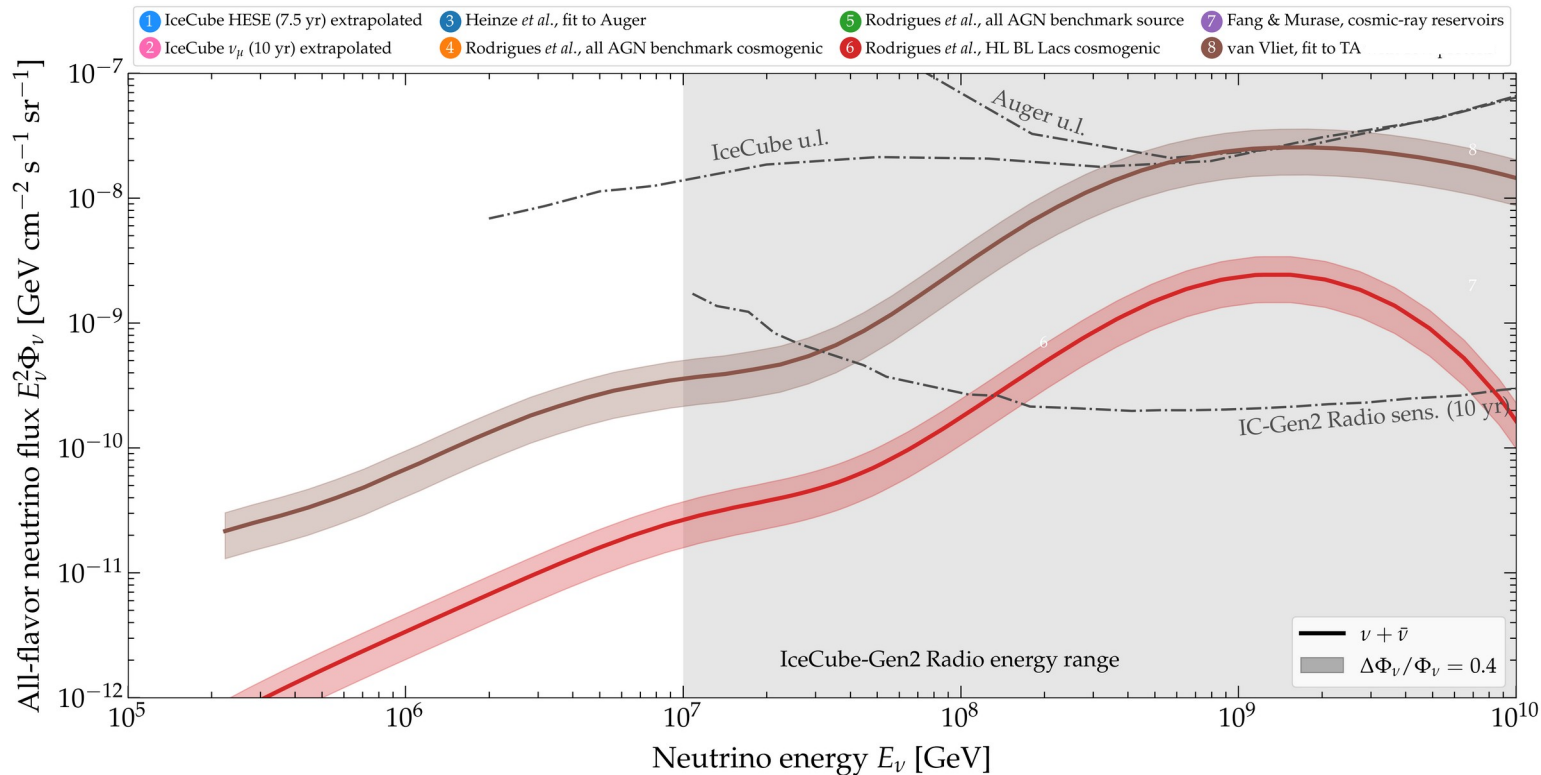
# But assuming we know the flux perfectly is aggressive

At this point we have obtained our results under the assumption of a known flux



# But assuming we know the flux perfectly is aggressive

We add an error in the flux normalization to account for  
and unknown flux on a first instance



# An analysis binned on reconstructed arrival direction

- The data set is divided in 3 different bins for  $\cos \theta_{z, \text{rec}}$ .
- The normalization uncertainty is taken into account as a Gaussian prior.
- As the exposure time increases, the measurement improves but eventually saturates

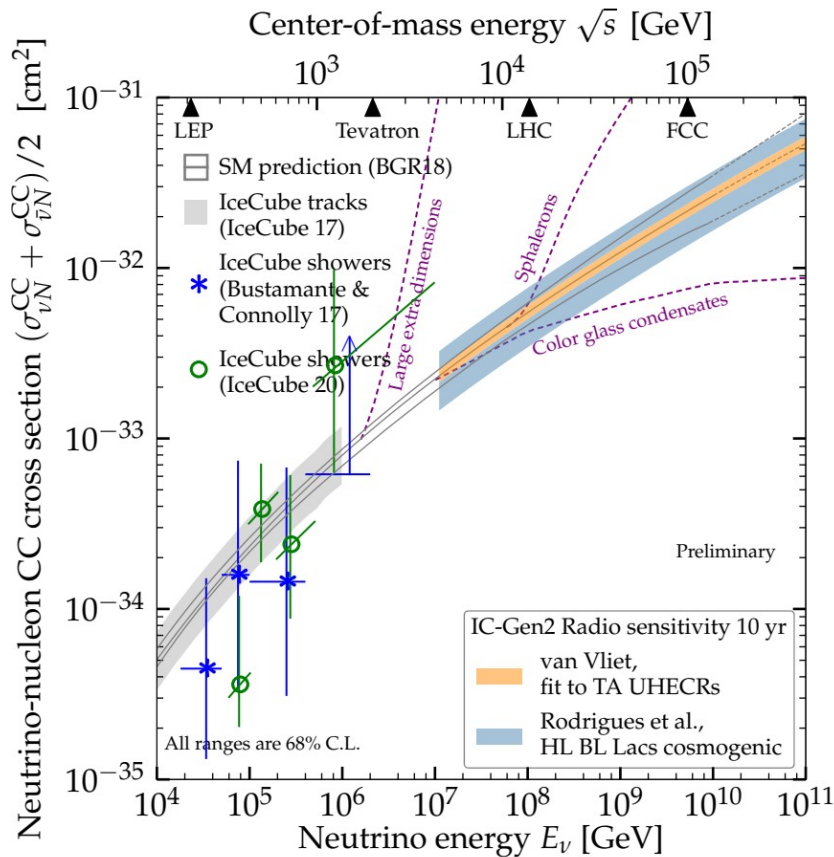
$$\Delta\Phi/\Phi = 0.4$$

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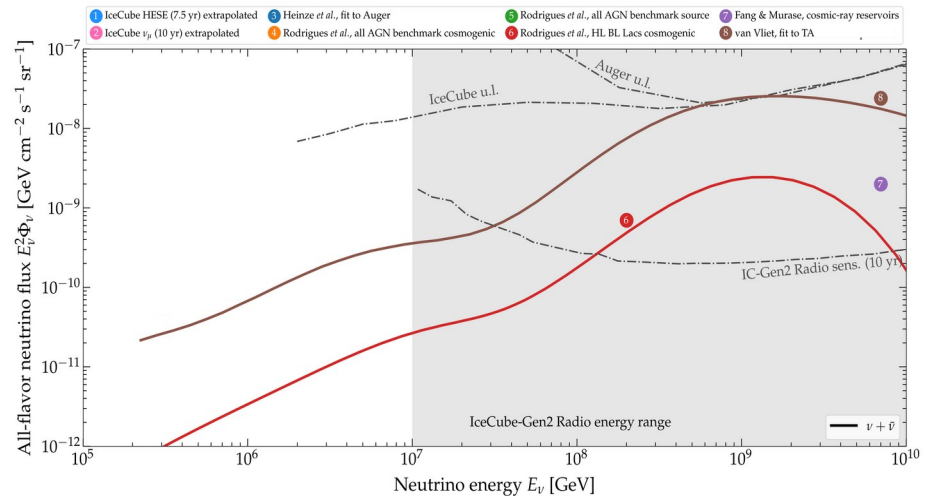
$$\Delta\Phi/\Phi = 0.1$$

# How well can we measure $\sigma_{\nu N}$ in the next decade?

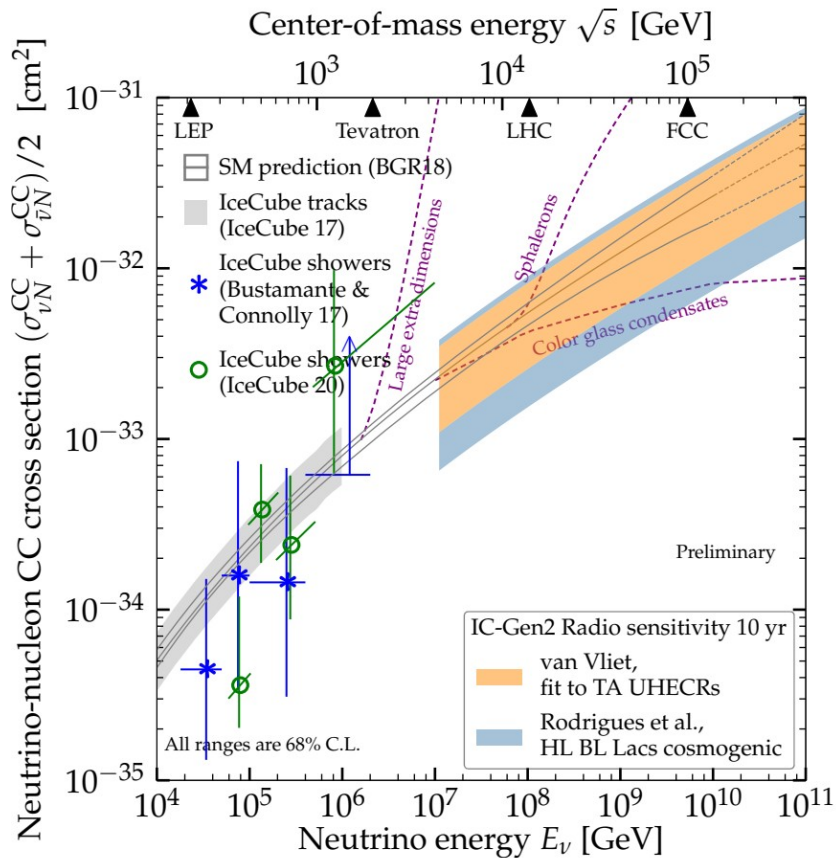


$$\sigma/\sigma_{\text{SM}} = 1 + 0.395 - 0.375$$

$$\sigma/\sigma_{\text{SM}} = 1 + 0.101 - 0.085$$

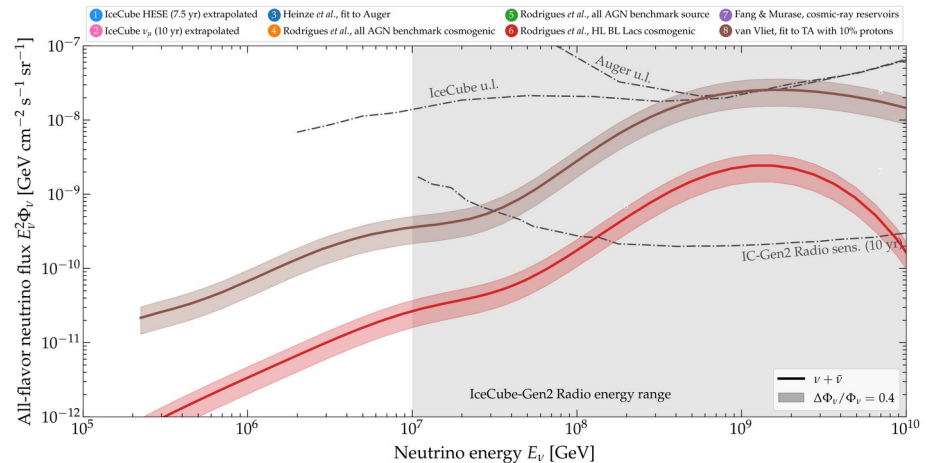


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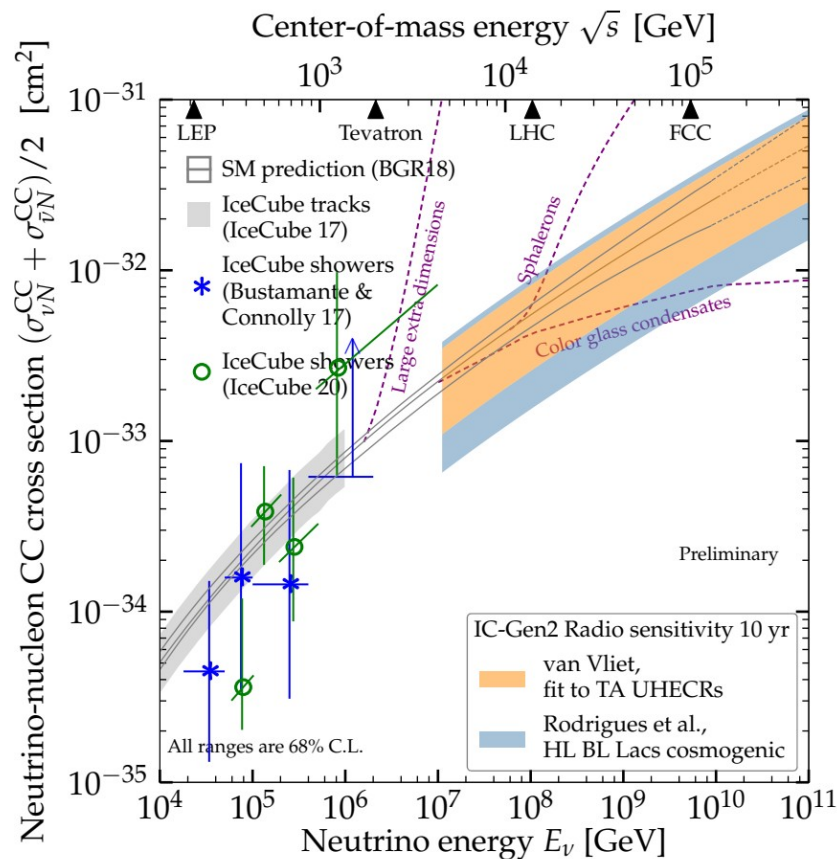
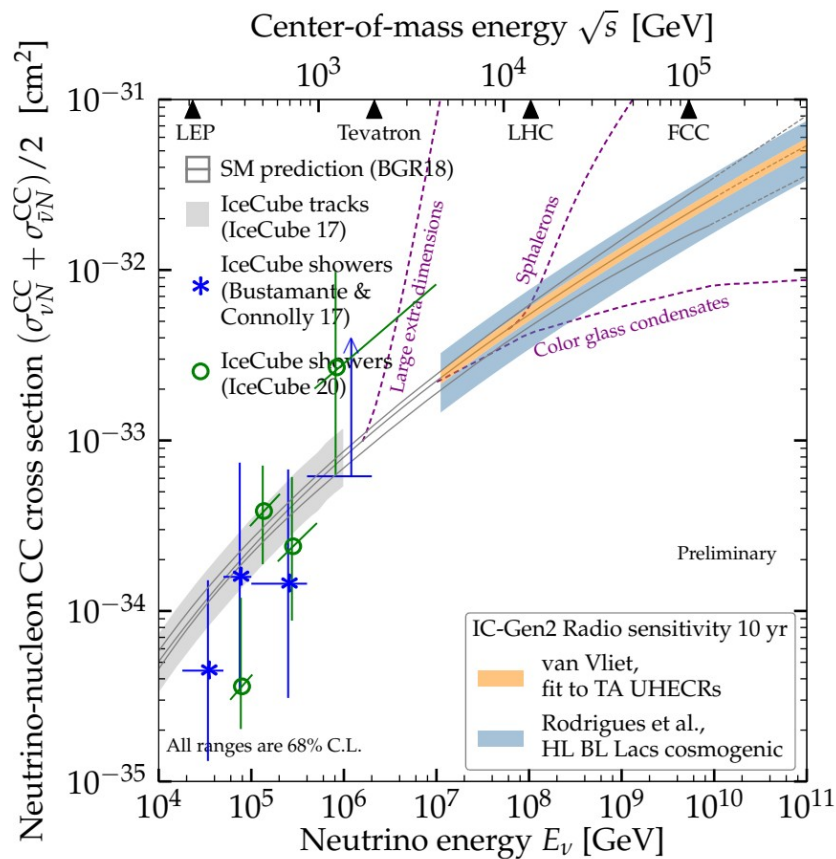


$$\sigma/\sigma_{\text{SM}} = 1 + 0.631 - 0.720$$

$$\sigma/\sigma_{\text{SM}} = 1 + 0.492 - 0.531$$



# How well can we measure $\sigma_{\nu N}$ in the next decade?

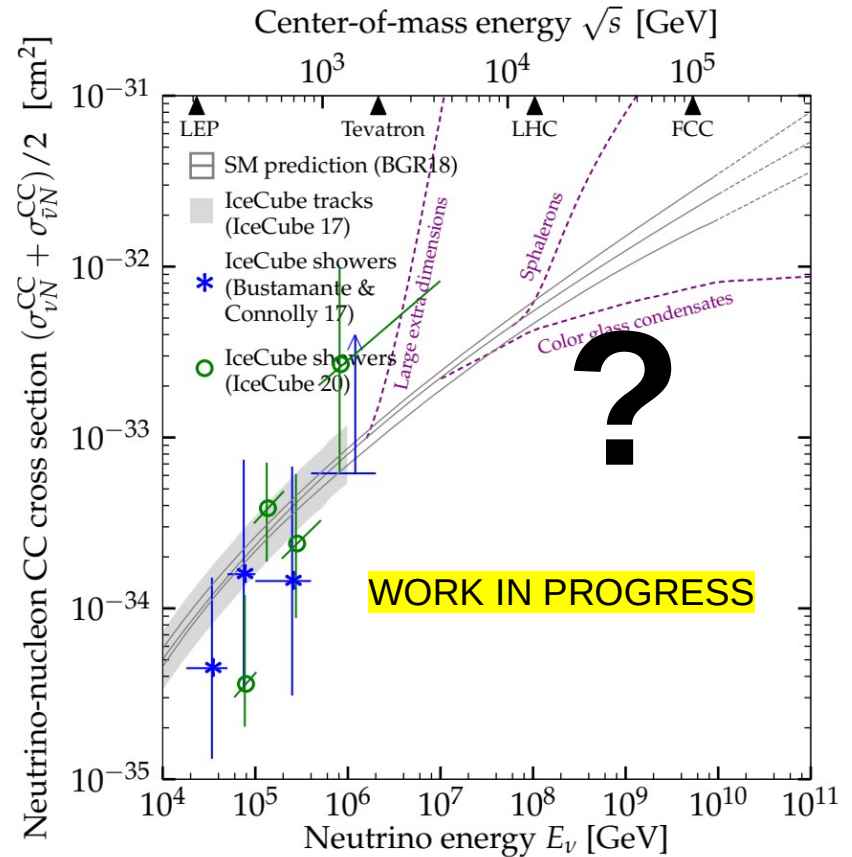




# Back to the most general situation

Relax model dependency in stages:

- 1) Perfectly known UHE $\nu$  flux.
- 2) Add a normalization uncertainty to the assumed UHE $\nu$ .
- 3) Measure the flux and the cross section **simultaneously**.



## Conclusions & take-home message

- The radio component of IceCube-Gen2 has good sensitivity to UHE  $\nu$ .
- Good enough for physics studies if the flux is high enough.
- We have developed a detailed end-to-end machinery for event rates a physics studies.
- Cross section measurement measurement depend on the uncertainty of the flux
- For high fluxes a more clever strategy is being follow to measure cross section and flux simultaneously.

# Backup slides

# Event rate full computation

Differential event rate per energy and arrival direction

$$\frac{d^3 N_{\nu\alpha}^{\text{CC}}}{dE_\nu dy d\cos\theta_z} = 2\pi T N_{\text{Av}} \frac{\rho_{\text{ice}}}{M_{\text{ice}}} V_{\text{eff},\nu\alpha}^{\text{CC}}(E_\nu, \cos\theta_z) \times \frac{d\sigma_{\nu N}^{\text{CC}}(E_\nu, y)}{dy} \Phi_{\nu\alpha}^{\text{det}}(E_\nu, \cos\theta_z)$$

Diff. event rate in terms of deposited energy and reconstructed arrival direction

$$\frac{d^2 N_{\nu\alpha}^{\text{CC}}}{dE_{\text{dep}} d\theta_{z,\text{rec}}} = \int_{-1}^{+1} d\cos\theta_z \int_{E_{\text{dep}}}^{\infty} dE_\nu \int_0^1 dy \frac{d^3 N_{\nu\alpha}^{\text{CC}}}{dE_\nu dy d\cos\theta_z} R_E[E_{\text{dep}}, E_{\text{sh},\nu\alpha}^{\text{CC}}(E_\nu, y)] R_{\theta_z}(\theta_{z,\text{rec}}, \theta_z)$$

The total number of events is obtained from integration over deposited energy and reconstructed arrival direction

