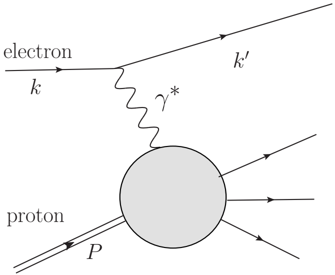
Generalised Partons Distributions: recent developments and perspectives

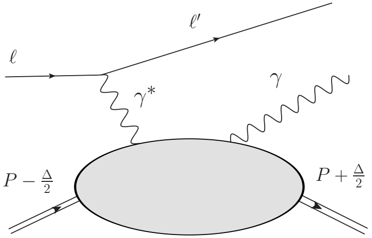
Cédric Mezrag

CEA Saclay, Irfu DPhN

May 31^{st}, 2021

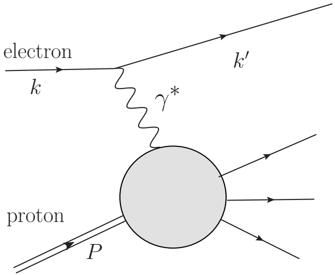
Why deep exclusive processes ?

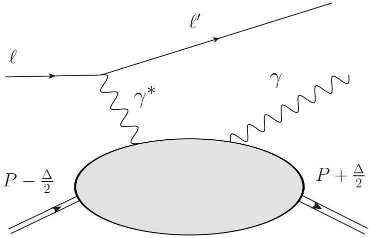




- Deep exclusive processes are generally more difficult to measure than inclusive ones

Why deep exclusive processes ?





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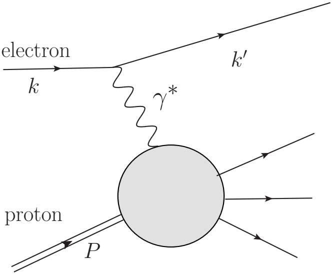
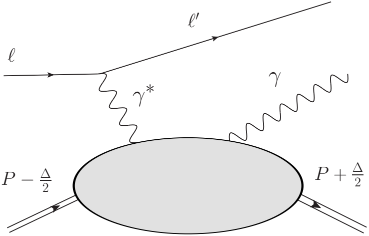
- Reason : we require *not* to break the proton $\rightarrow$ small cross sections

Why deep exclusive processes ?





- Deep exclusive processes are generally more difficult to measure than inclusive ones
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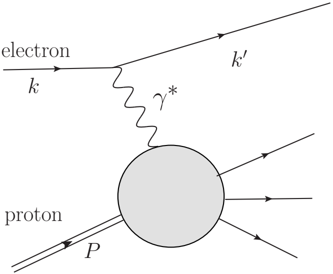
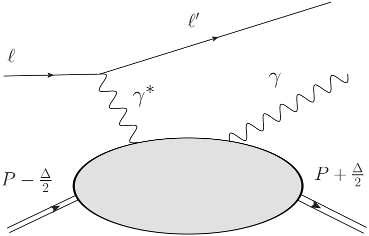
A curse and a blessing

Not breaking the proton allows one to study the distribution of quarks and gluons in coordinate space

Why deep exclusive processes ?





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New correlators

In exclusive processes, usual PDF are traded for Generalised Parton Distributions (GPDs)

Nucleon tomography through Generalised Partons Distributions (GPDs)

Generalised Parton Distributions

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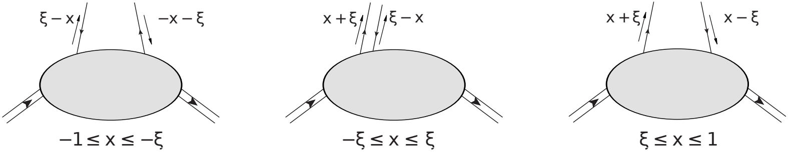
  - "hadron-parton" amplitudes which depend on three variables (x, ξ, t) and a scale μ,



  - ★ x: average momentum fraction carried by the active parton
  - ★ ξ: skewness parameter $\xi \simeq \frac{x_B}{2-x_B}$
  - ★ t: the Mandelstam variable

- Generalised Parton Distributions (GPDs):
  - "hadron-parton" amplitudes which depend on three variables (x, ξ, t) and a scale μ,
  - are defined in terms of a non-local matrix element,

$$\frac{1}{2}\int \frac{e^{ixP^+z^-}}{2\pi}\langle P+\frac{\Delta}{2}|\bar{\psi}^q(-\frac{z}{2})\gamma^+\psi^q(\frac{z}{2})|P-\frac{\Delta}{2}\rangle \mathrm{d}z^-|_{z^+=0,\mathbf{z}=0}$$
$$= \frac{1}{2P^+}\left[H^q(x,\xi,t)\bar{u}\gamma^+u + E^q(x,\xi,t)\bar{u}\frac{i\sigma^{+\alpha}\Delta_\alpha}{2M}u\right].$$

$$\frac{1}{2}\int \frac{e^{ixP^+z^-}}{2\pi}\langle P+\frac{\Delta}{2}|\bar{\psi}^q(-\frac{z}{2})\gamma^+\gamma_5\psi^q(\frac{z}{2})|P-\frac{\Delta}{2}\rangle \mathrm{d}z^-|_{z^+=0,\mathbf{z}=0}$$
$$= \frac{1}{2P^+}\left[\tilde{H}^q(x,\xi,t)\bar{u}\gamma^+\gamma_5 u + \tilde{E}^q(x,\xi,t)\bar{u}\frac{\gamma_5\Delta^+}{2M}u\right].$$

D. Müller *et al.*, Fortsch. Phy. 42 101 (1994)
X. Ji, Phys. Rev. Lett. **78**, 610 (1997)
A. Radyushkin, Phys. Lett. **B380**, 417 (1996)

4 GPDs without helicity transfer + 4 helicity flip GPDs

Generalised Parton Distributions

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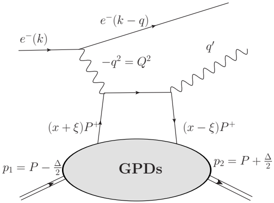
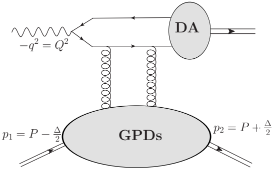
  - can be split into quark flavour and gluon contributions,
  - are related to PDF in the forward limit $H(x, \xi = 0, t = 0; \mu) = q(x; \mu)$
  - are universal, *i.e.* are related to the Compton Form Factors (CFFs) of various exclusive processes through convolutions

$$\mathcal{H}(\xi, t) = \int \mathrm{d}x \; C(x, \xi) H(x, \xi, t)$$





Properties

- Polynomiality Property:

$$\int_{-1}^{1} \mathrm{d}x\, x^m H^q(x,\xi,t;\mu) = \sum_{j=0}^{\left[\frac{m}{2}\right]} \xi^{2j} C^q_{2j}(t;\mu) + mod(m,2)\xi^{m+1} C^q_{m+1}(t;\mu)$$

X. Ji, J.Phys.G 24 (1998) 1181-1205
A. Radyushkin, Phys.Lett.B 449 (1999) 81-88

Special case :

$$\int_{-1}^{1} \mathrm{d}x\; H^q(x,\xi,t;\mu) = F_1^q(t)$$

Lorentz Covariance

Properties

- Polynomiality Property:

Lorentz Covariance

- Positivity property:

$$\left| H^q(x,\xi,t) - \frac{\xi^2}{1-\xi^2} E^q(x,\xi,t) \right| \leq \sqrt{\frac{q\left(\frac{x+\xi}{1+\xi}\right) q\left(\frac{x-\xi}{1-\xi}\right)}{1-\xi^2}}$$

A. Radysuhkin, Phys. Rev. **D59**, 014030 (1999)
B. Pire *et al.*, Eur. Phys. J. **C8**, 103 (1999)
M. Diehl *et al.*, Nucl. Phys. **B596**, 33 (2001)
P.V. Pobilitsa, Phys. Rev. **D65**, 114015 (2002)

Positivity of Hilbert space norm

- Polynomiality Property:

Lorentz Covariance

- Positivity property:

Positivity of Hilbert space norm

- Support property:

$$x \in [-1; 1]$$

M. Diehl and T. Gousset, Phys. Lett. **B428**, 359 (1998)

Relativistic quantum mechanics

Properties

- Polynomiality Property:

Lorentz Covariance

- Positivity property:

Positivity of Hilbert space norm

- Support property:

Relativistic quantum mechanics

- Scale evolution property
  $\rightarrow$ generalization of DGLAP and ERBL evolution equations

D. Müller *et al.*, Fortschr. Phys. 42, 101 (1994)

Renormalization

Properties

- Polynomiality Property:

Lorentz Covariance

- Positivity property:

Positivity of Hilbert space norm

- Support property:

Relativistic quantum mechanics

- Scale evolution property

Renormalization

Problem

- There is hardly any model fulfilling *a priori* all these constraints.

see J.-M. Morgado Chavez talk this afternoon

- Lattice QCD computations remain very challenging.

- In the limit $\xi \to 0$, one recovers a density interpretation:
  - 1D in momentum space (x)
  - 2D in coordinate space $\vec{b}_\perp$ (related to t)

M. Burkardt, Phys. Rev. **D62**, 071503 (2000)

Interpretation of GPDs I

2+1D structure of the nucleon

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- Possibility to extract density from experimental data

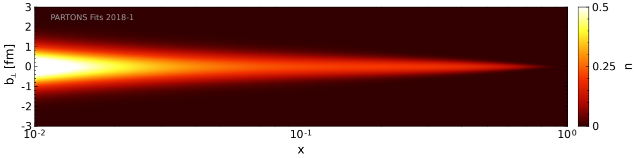

figure from H. Moutarde *et al.*, EPJC 78 (2018) 890

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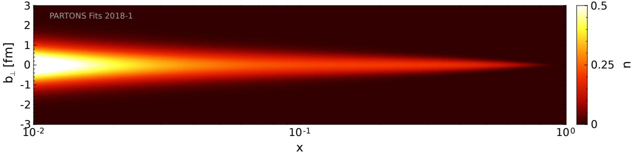

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- Correlation between x and $b_\perp$ $\to$ going beyond PDF and FF.

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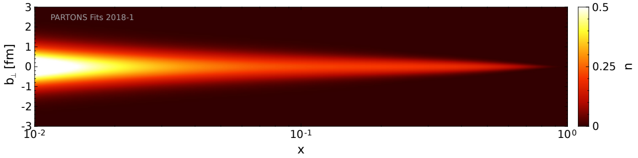

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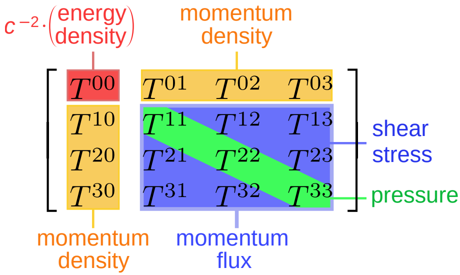
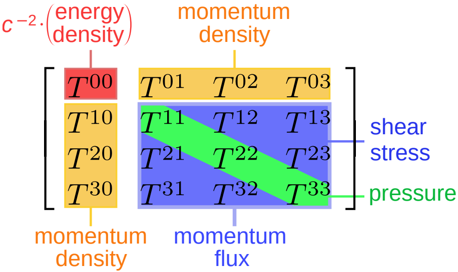
- Correlation between x and $b_\perp \to$ going beyond PDF and FF.
- Caveat: no experimental data at $\xi = 0$
  $\to$ extrapolations (and thus model-dependence) are necessary



How energy, momentum, pressure are shared between quarks and gluons

Caveat: renormalization scheme and scale dependence

C. Lorcé *et al.*, PLB 776 (2018) 38-47,
M. Polyakov and P. Schweitzer,
IJMPA 33 (2018) 26, 1830025
C. Lorcé *et al.*, Eur.Phys.J.C 79 (2019) 1, 89



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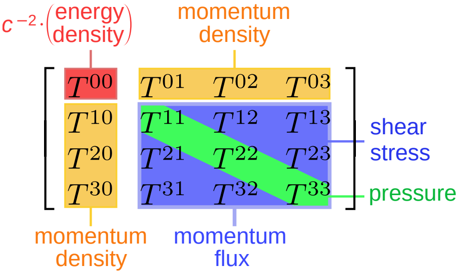
$$\langle p', s' | T^{\mu\nu}_{q,g} | p, s \rangle = \bar{u} \left[P^{\{\mu}\gamma^{\nu\}} A_{q,g}(t;\mu) + \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu}\Delta^2}{M} C_{q,g}(t;\mu) + M g^{\mu\nu} \bar{C}_{q,g}(t;\mu) + \frac{P^{\{\mu} i\sigma^{\nu\}\Delta}}{2M} B_{q,g}(t;\mu) + \frac{P^{[\mu} i\sigma^{\nu]\Delta}}{2M} D_{q,g}(t;\mu) \right] u$$



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$$\int_{-1}^{1} \mathrm{d}x\, x\, H_q(x,\xi,t;\mu) = A_q(t;\mu) + (2\xi)^2 C_q(t;\mu)$$

$$\int_{-1}^{1} \mathrm{d}x\, x\, E_q(x,\xi,t;\mu) = B_q(t;\mu) - (2\xi)^2 C_q(t;\mu)$$

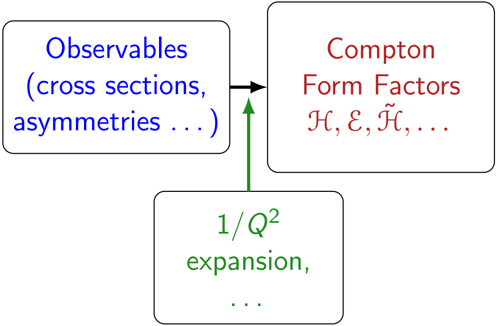
- Ji sum rule
- Fluid mechanics analogy

X. Ji, PRL 78, 610-613 (1997)
M.V. Polyakov PLB 555, 57-62 (2003)

Experimental connection to GPDs

Observables
(cross sections,
asymmetries . . .)

Experimental connection to GPDs



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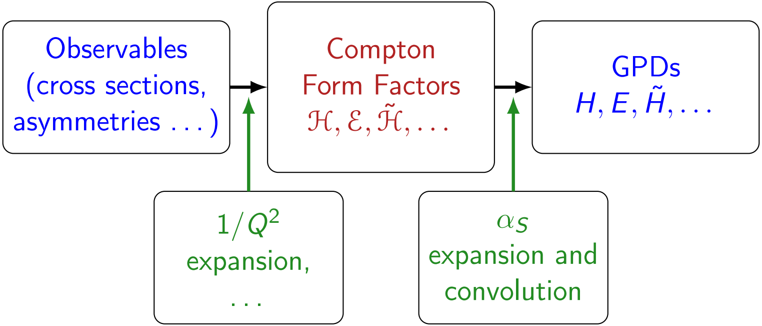

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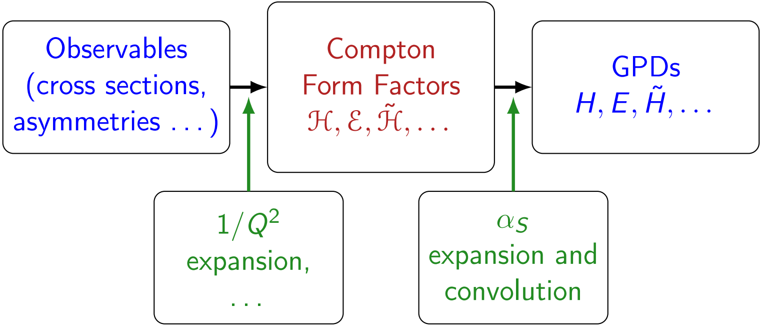
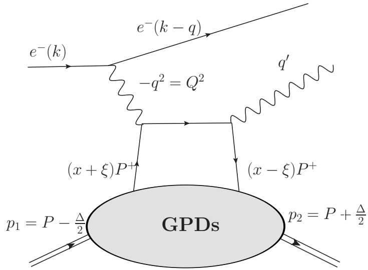


- CFFs play today a central role in our understanding of GPDs
- Extraction generally focused on CFFs

Deep Virtual Compton Scattering



- Best studied experimental process connected to GPDs

  $\rightarrow$ Data taken at Hermes, Compass, JLab 6, JLab 12

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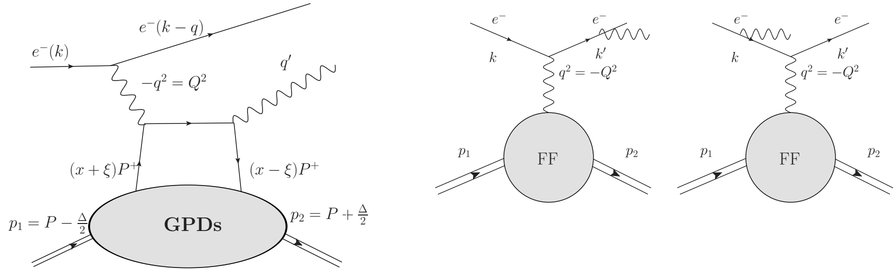

- Best studied experimental process connected to GPDs
  → Data taken at Hermes, Compass, JLab 6, JLab 12
- Interferes with the Bethe-Heitler (BH) process
  - Blessing: Interference term boosted w.r.t. pure DVCS one
  - Curse: access to the angular modulation of the pure DVCS part difficult

M. Defurne *et al.*, Nature Commun. 8 (2017) 1, 1408

QCD corrections to DVCS

- At LO, the DVCS coefficient function is a QED one

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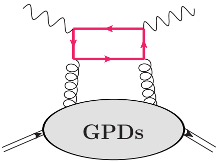
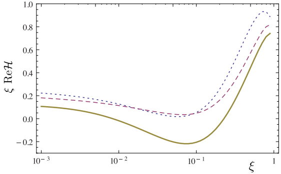
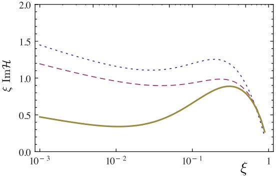
- At NLO, gluon GPDs play a significant role in DVCS







H. Moutarde *et al.*, PRD 87 (2013) 5, 054029

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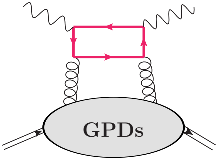

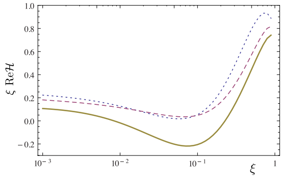

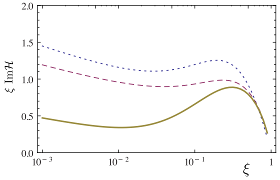

H. Moutarde *et al.*, PRD 87 (2013) 5, 054029

- Recent N2LO studies, impact needs to be assessed

V. Braun *et al.*, JHEP 09 (2020) 117

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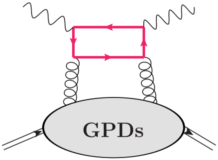

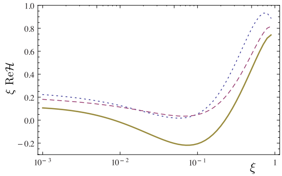

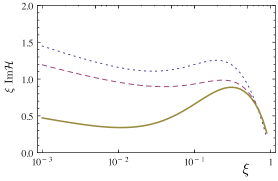

H. Moutarde *et al.*, PRD 87 (2013) 5, 054029

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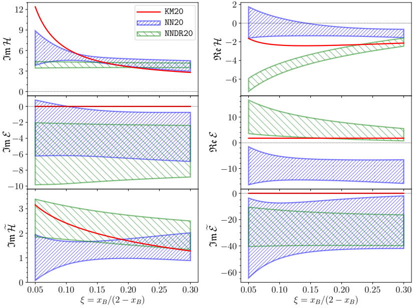
V. Braun *et al.*, JHEP 09 (2020) 117

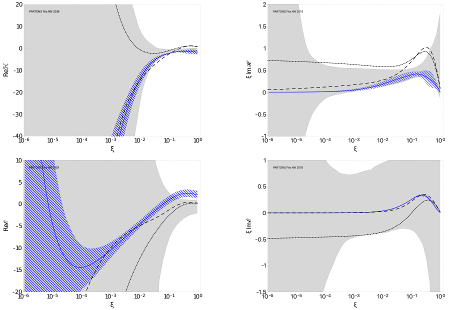
- Evolution equations: needs for open evolution codes
  $\rightarrow$ dedicated round table this afternoon

A. Vinnikov, hep-ph/0604248
V. Bertone *et al.*, in preparation

Recent CFF extractions



M. Cuič *et al.*, PRL 125, (2020), 232005



H. Moutarde *et al.*, EPJC 79, (2019), 614

- Recent effort on bias reduction in CFF extraction (ANN)

  additional ongoing studies, J. Grigsby *et al.*, arXiv:2012.04801

- Studies of ANN architecture to fulfil GPDs properties (dispersion relation,polynomiality,. . .)
- Recent efforts on propagation of uncertainties (allowing impact studies for JLAB12, EIC and EICC)

  see *e.g.* H. Dutrieux *et al.*, arXiv:2105.09245

Dispersion relation and the D-term

- At all orders in α_S, dispersion relations relate the real and imaginary parts of the CFF.

I. Anikin and O. Teryaev, PRD 76 056007
M. Diehl and D. Ivanov, EPJC 52 (2007) 919-932

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- For instance at LO:

$$Re(\mathcal{H}(\xi,t)) = \frac{1}{\pi}\int_{-1}^{1}\mathrm{d}x\ Im(\mathcal{H}(x,t))\left[\frac{1}{\xi-x}-\frac{1}{\xi+x}\right] + \underbrace{2\int_{-1}^{1}\mathrm{d}\alpha\frac{D(\alpha,t)}{1-\alpha}}_{\text{Independent of } \xi}$$

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- For instance at LO:

$$\underbrace{Re(\mathcal{H}(\xi,t))}_{\text{Extracted from data}} = \frac{1}{\pi}\int_{-1}^{1} \mathrm{d}x \underbrace{Im(\mathcal{H}(x,t))}_{\text{Extracted from data}} \left[\frac{1}{\xi - x} - \frac{1}{\xi + x}\right] + 2\int_{-1}^{1} \mathrm{d}\alpha \frac{D(\alpha,t)}{1-\alpha}$$

- $D(\alpha,t)$ is related to the EMT (pressure and shear forces)

M.V. Polyakov PLB 555, 57-62 (2003)

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M.V. Polyakov PLB 555, 57-62 (2003)

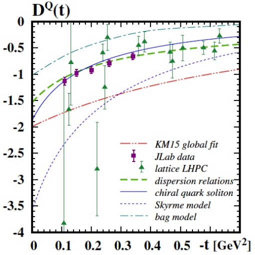


figure from M. Polyakov and P Schweitzer, IJMPA 33 (2018) 26, 1830025
data from Burkert *et al.*, Nature 557 (2018) 7705, 396-399

- First attempt from JLab 6 *GeV* data
  Warning: Systematic uncertainties are not reported on figure

Burkert *et al.*, Nature 557 (2018) 7705, 396-399

- Tensions with other studies

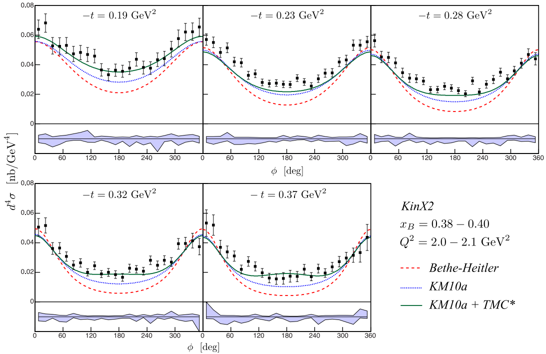
K. Kumericki, Nature 570 (2019) 7759, E1-E2
H. Moutarde *et al.*, Eur.Phys.J.C 79 (2019) 7, 614
H. Dutrieux *et al.*, Eur.Phys.J.C 81 (2021) 4

- Model dependence, scheme/scale dependence

Finite t corrections

Kinematical corrections in t/Q^2 and M^2/Q^2

V. Braun *et al.*, PRL 109 (2012), 242001



M. Defurne *et al.* PRC 92 (2015) 55202

- Sizeable even for $t/Q^2 \sim 0.1$
- Not currently included in global fits.

DVCS on Nuclei

- DVCS off the deuteron

F. Cano *et al.*, EPJA 19 (2004) 423
M. Benali *et al.*, Nature Phys. 16 (2020) 2, 191-198

  - Incoherent scattering : DVCS off the quasi-free neutron
    → significant step toward flavour separation

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  - Coherent scattering : probing partons inside a deuteron
    → Spin 1 target: richer spin structure → more GPDs
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- DVCS off He^4

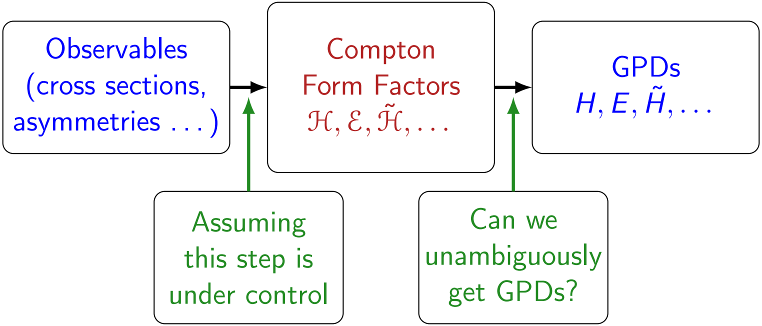
M. Hattawy *et al.*, PRL 119 (2017) 20, 202004

  - Coherent scattering on a scalar target
    $\rightarrow$ Less spin structure $\rightarrow$ less GPDs
  - Incoherent scattering: information on the structure of a bound nucleon

S. Fucini *et al.*, Phys.Rev.C 102 (2020) 065205

The DVCS deconvolution problem I

From CFF to GPDs



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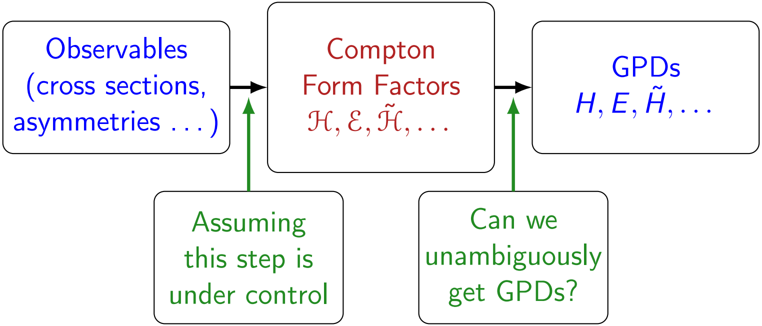

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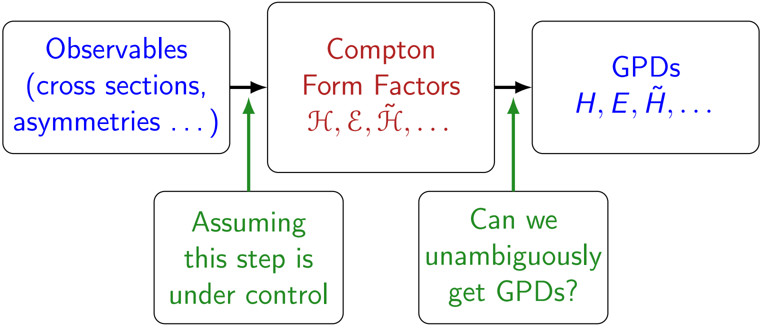

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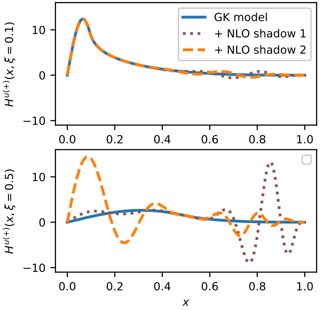
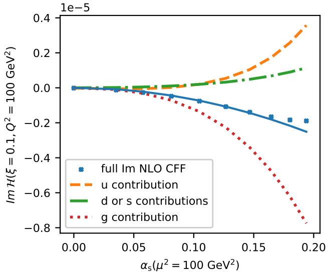
- Are QCD corrections improving the situation?

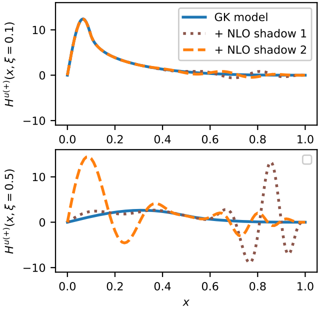


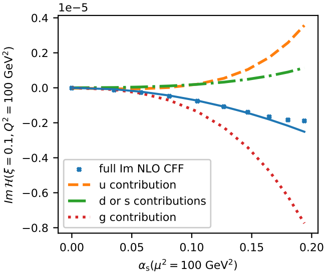


- NLO analysis of shadow GPDs:
  - Cancelling the line $x = \xi$ is necessary but **no longer** sufficient
  - Additional conditions brought by NLO corrections reduce the size of the "shadow space"...
  - ... but do not reduce it to 0 $\rightarrow$ NLO shadow GPDs

H. Dutrieux *et al.*, arXiv:2104.03836

The DVCS deconvolution problem II





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  - it was argued that evolution would solve this issue

    A. Freund PLB 472, 412 (2000)
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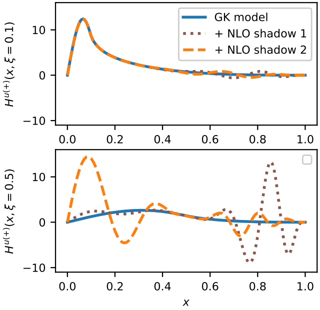

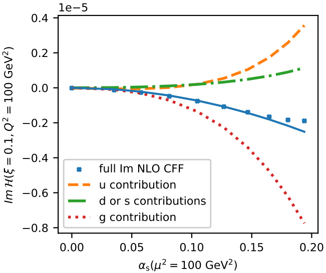

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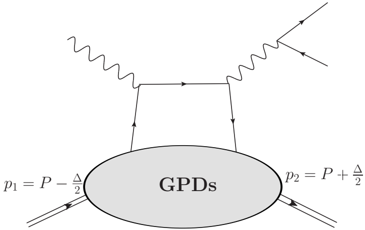
Multichannel Analysis required to fully determine GPDs

Timelike Compton Scattering



- Amplitude related to the DVCS one ($Q^2 \to -Q^2, \dots$)
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  E. Berger *et al.*, EPJC 23 (2002) 675
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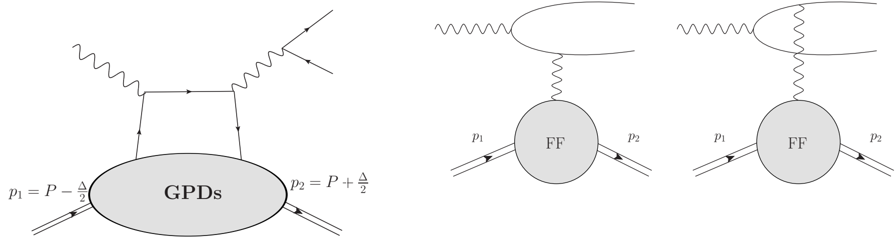

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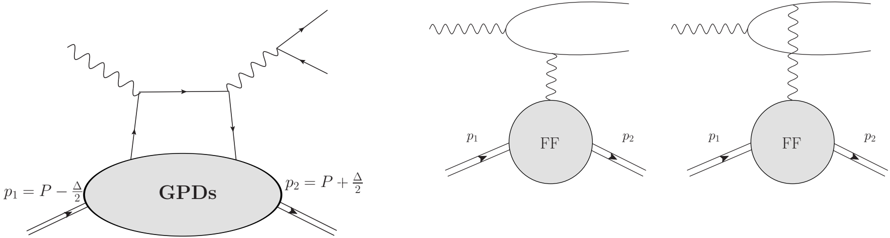

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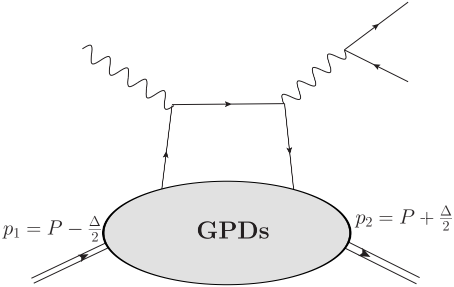
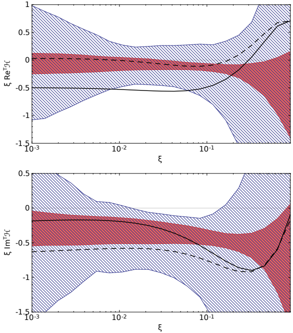
- Interferes with the Bethe-Heitler (BH) process
- Same type of final states as exclusive quarkonium production

TCS: Recent results



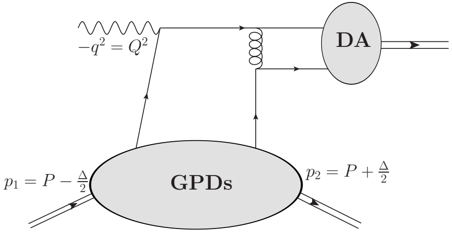


O. Grocholski *et al.*, EPJC 80, (2020) 61

- DVCS Data-driven prediction for TCS at LO and NLO
- First experimental measurement at JLab through forward-backward asymmetry (interference term)

P. Chatagnon *et al.*, in preparation

Deep Virtual Meson Production



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- Same GPDs than previously
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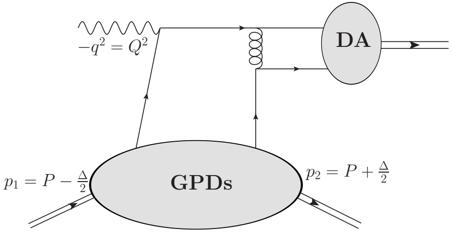

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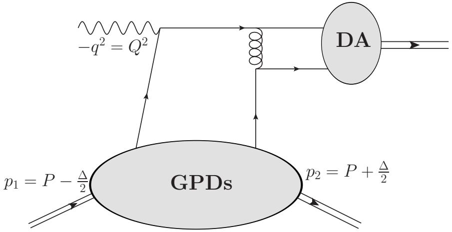

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- Factorisation proven $\neq$ factorisation visible at achievable Q^2
  - Leading-twist dominance at a given Q^2 is process-dependent $\rightarrow$ for DVMP it can change between mesons.
  - At JLab kinematics, higher-twist contributions are very strong $\rightarrow$ hide factorisation of σ_L

Status of DVMP

- π^0 electroproduction
  - $\sigma_T > \sigma_L$ at JLab 6 and likely at JLab 12 kinematics ($Q^2 = 8.3 GeV^2$)
    M. Dlamini *et al.*, arXiv:2011.11125
  - No extraction of σ_L at JLab 12 yet
  - Model-dependent treatment of σ_T using higher-twist contributions
    S. V. Goloskokov and P. Kroll, EPJC 65, 137 (2010)
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    see *e.g.* L. Favart, EPJA 52 (2016) 6, 158
  - $\sigma_T \neq 0$ though $\rho_{0;T}$ production vanishes at leading twist
    $\rightarrow$ No LT access to chiral-odd GPDs.

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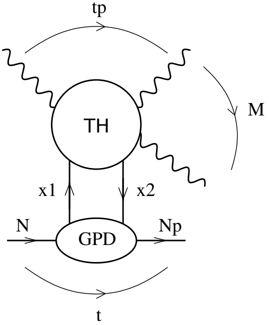
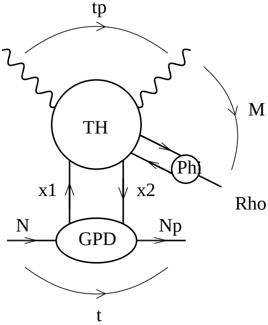
DVMP is as interesting as challenging
Additional data would be more than welcome

New channels: Multiparticle production



A. Pedrak *et al.*, PRD 96 (2017) 7, 074008



R. Boussarie *et al.*, JHEP 02 (2017) 054

- New combination of CFFs $\rightarrow$ welcome in global fits.
- LT access to chiral-odd GPDs in the (γ, ρ) case.
- Electroproduction done for $\gamma\gamma$.
- Additional particle in the final state
  - more difficult experimentally $\rightarrow$ need higher luminosity
  - more degrees of freedom $\rightarrow$ solution to the deconvolution problem?

PARTONS and Gepard

PARTONS

partons.cea.fr

B. Berthou *et al.*, EPJC 78 (2018) 478

Gepard

calculon.phy.hr/gpd/server/index.html

K. Kumericki, EPJ Web Conf. 112 (2016) 01012

see talk this afternoon by K. Kumericki

- Similarities : NLO computations, BM formalism, ANN, ...
- Differences : models, evolution, ...

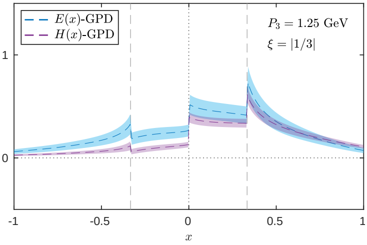
Physics impact

These integrated softwares are the mandatory path toward reliable multichannel analyses.

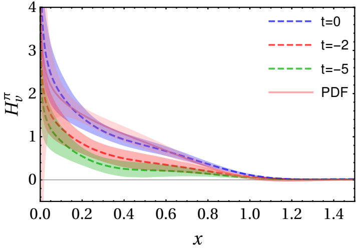
Lattice-QCD computations

- Lattice practicionner used to compute matrix element of local operators
  $\rightarrow$ Mellin moments of GPDs
- new techniques allow to extrapolate euclidean matrix element on the lightcone

X. Ji Phys. Rev. Lett., 2013, 110, 262002
Y.-Q. Ma and J.-W. Qiu Phys. Rev. D, 2018, 98, 074021
A. Radyushkin, Phys. Rev. D, 2017, 96, 034025



C. Alexandrou *et al.*, PRL 125 (2020) 26, 262001



J.-W. Chen Nucl.Phys.B 952 (2020) 114940

GPDs addendum

- Phenomenological parametrisations (KM, GK, VGG, MSW, . . .)
- Extension of the relation between CFF and observables

B. Kriesten *et al.* PRD 101 (2020) 054021

- New modelling efforts

N. Chouika *et al.*, EPJC 77, (2017) 906
S. Rodini, B. Pasquini *et al.*, in preparation

- Meson GPDs and virtual particle measurement

D. Amrath *et al.*,
(see talk by José Manuel Morgado Chavez this afternoon)

- Continuum QCD computations

see e.g. Jin-Lin Zhang *et al.*, arXiv:2009.11384
A. Freese *et al.*, Phys.Rev.C 101 (2020) 3, 035203

- Exclusive charmonium production
- Transition GPDs
- . . .

Conclusions

Summary

- After 25 years, GPDs formalism is well established ...
- ... but the GPDs themselves remain poorly known
- due to the high luminosity requirements
- the situation may change with JLAB 12, EIC and EICC

Perspectives

- Significant efforts in phenomenology remain be done (CFF and GPD)
- Multichannel analysis could help solving the deconvolution problem
- Ab-initio computations may provide insights in the next decade

In the perspective of EIC and EICC, a lot of work remains to be done to exploit the forthcoming data.

Thank you for your attention