

The $2N+1$ body problem: an impurity immersed in a spin $\frac{1}{2}$ fermionic superfluid

Frédéric Chevy

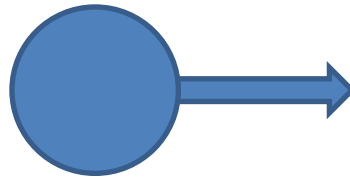
Laboratoire Kastler Brossel

C. Salomon, I. Ferrier-Barbur, M. Delehaye, S. Laurent, M. Pierce (Exp)
X. Leyronas, R. Alhyder, A. Bigué (Theory)



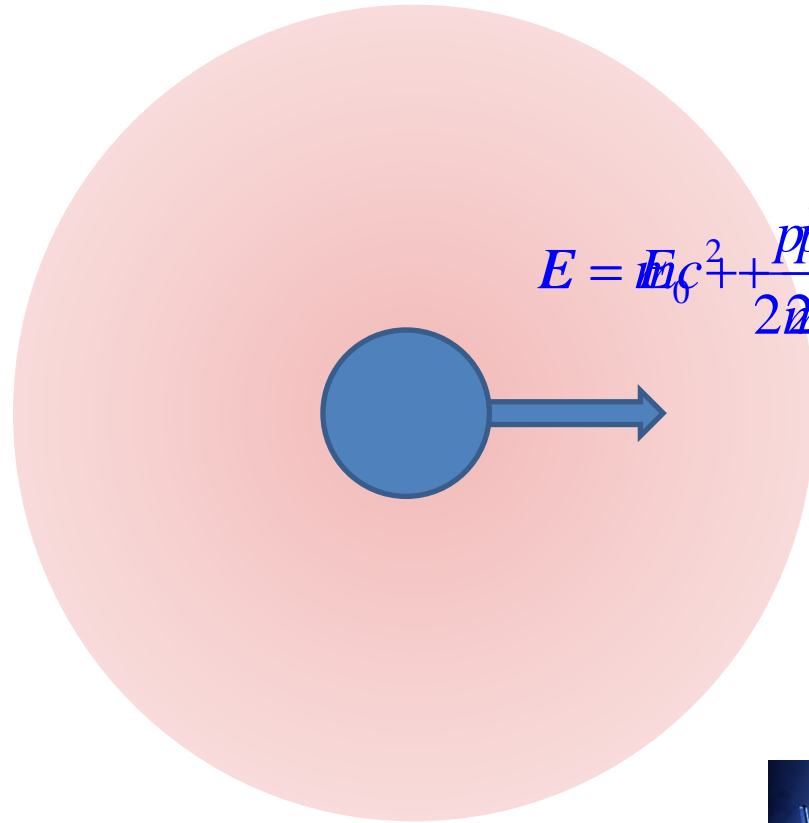
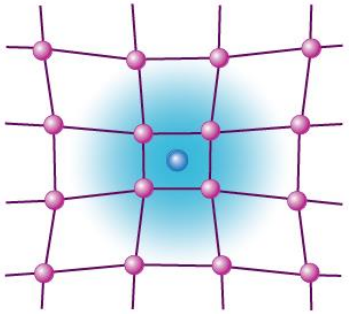
THE POLARON PROBLEM

$$E = mc^2 + \frac{p^2}{2m} + \dots$$

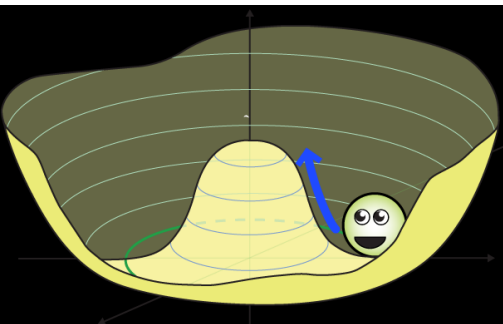


THE POLARON PROBLEM

Landau-Pekar-Frölich



Higgs mechanism



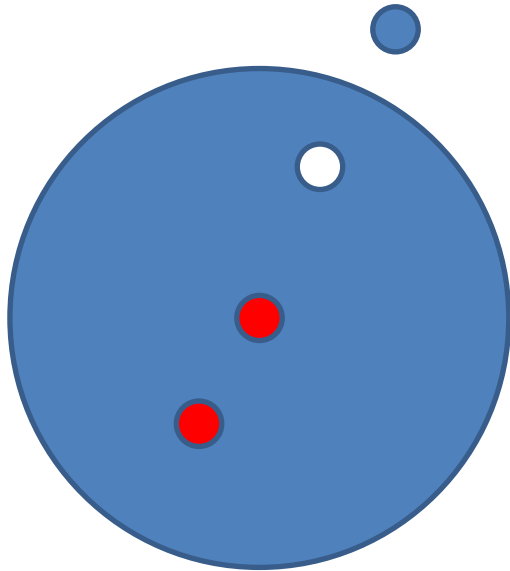
Dilute solutions



Ultracold atoms: Fermi and Bose polarons

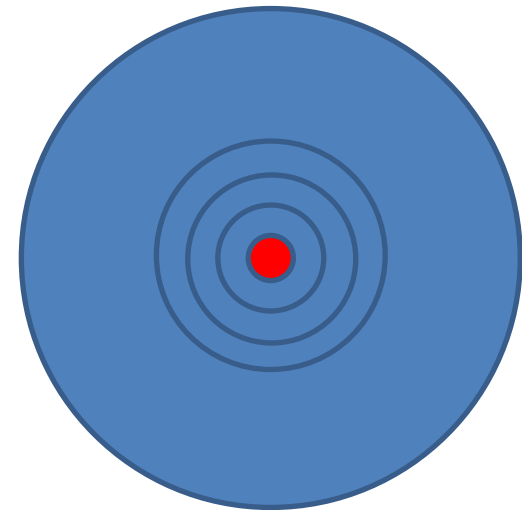
Ultracold regime: de Broglie wavelength \gg range of the interatomic potential (**zero-range limit**)

Fermi polaron (ENS, MIT, Innsbruck, Technion...): impurity immersed in an ideal gas of spin polarized fermions



Stable for any $k_F a$ (Moser & Seringer, 2016)

Bose polaron (Aarhus, JILA): impurity immersed in an weakly interacting Bose-Einstein condensate



Efimov effect: polaron unstable for large na^3



POLARON PHASE DIAGRAM

Impurity-medium coupling

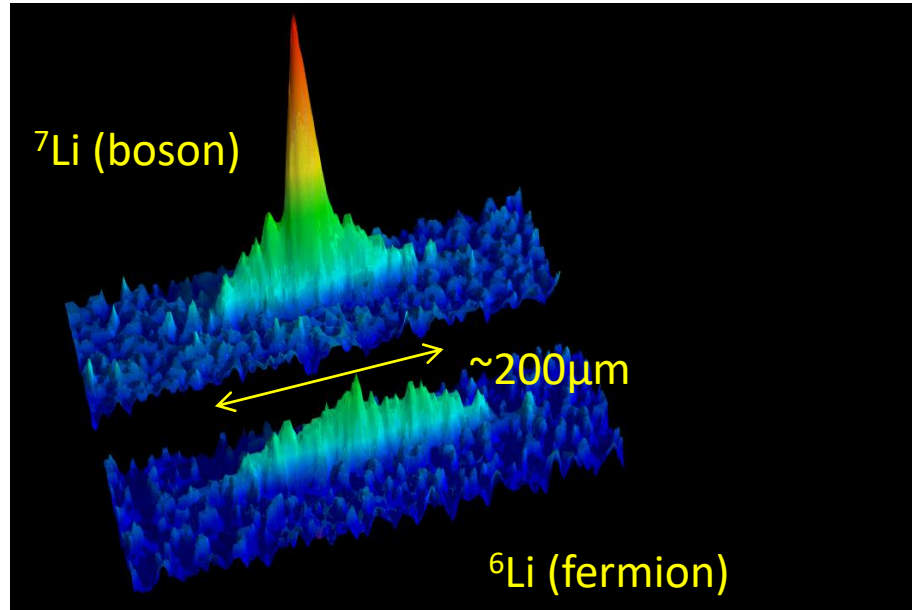
Fermi polaron (ENS, MIT, Innsbruck...); Bose polaron (Aarhus, JILA); Kondo problem...

Intra medium coupling

POLARON PHASE DIAGRAM

Impurity-medium coupling

Fermi polaron (ENS, MIT, Innsbruck...); Bose polaron (Aarhus, JILA); Kondo problem...



Impurity weakly coupled to a strongly correlated medium (ENS, Seattle, Tokyo, Shanghai...)

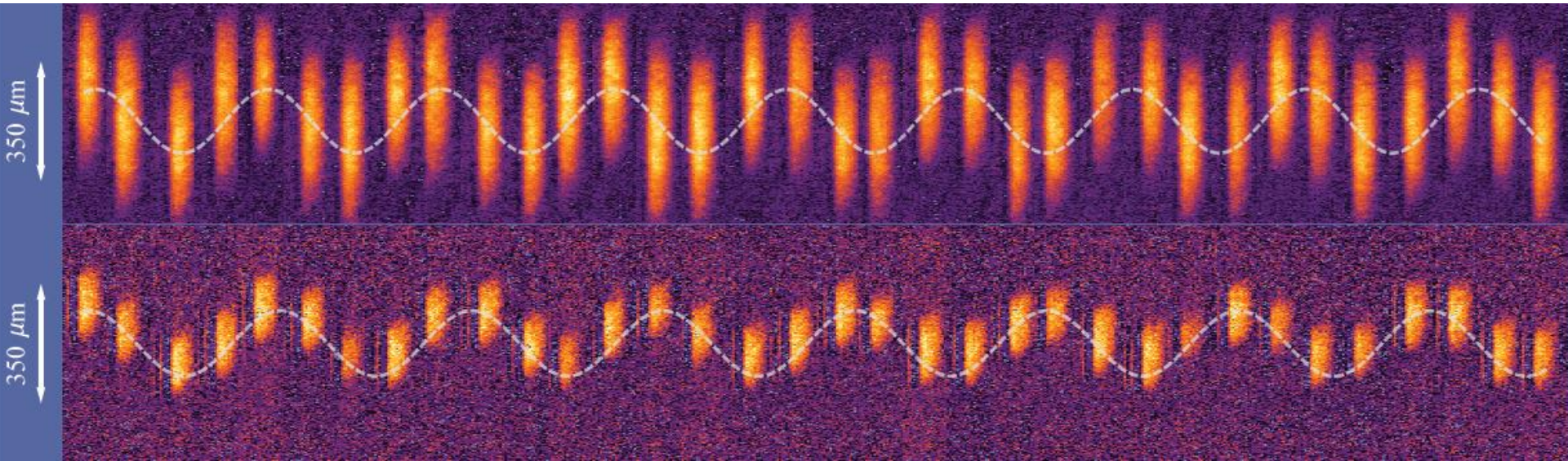
Intra medium coupling

Dynamics of the impurity

I. Ferrier-Barbut *et al.*, Science **345**, 1035 (2014)

See also C. Salomon's talk

DYNAMICS OF THE MIXTURES



$$\omega_6 = 2\pi \times 16.80(2)\text{Hz}$$

$$\tilde{\omega}_6 = 2\pi \times 16.80(1)\text{Hz}$$

$$\omega_7 = 2\pi \times 15.27(2)\text{Hz}$$

$$\tilde{\omega}_7 = 2\pi \times 15.00(1)\text{Hz}$$

Single Superfluid

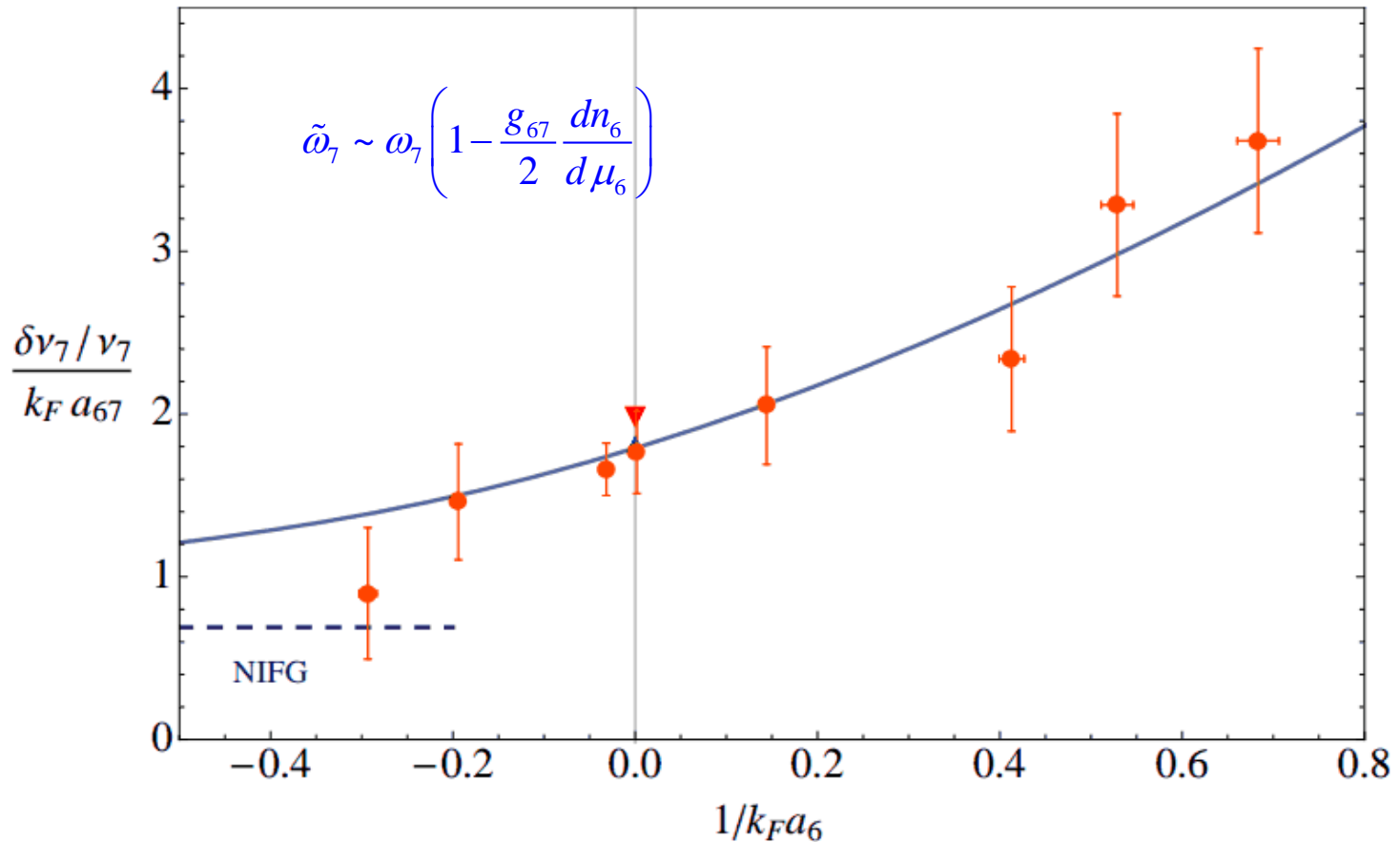
Coupled Superfluids

$$\text{Ratio} = (7/6)^{1/2} = (m_7/m_6)^{1/2}$$

FREQUENCY SHIFT

$$V_{\text{eff},7} = V(\mathbf{r}) + g_{67} n_6(\mu_6(\mathbf{r})) \quad \mu_6(\mathbf{r}) = \mu_6^0 - V(\mathbf{r}) \quad (\text{Local Density Approximation})$$

$$\approx g_{67} n_6(\mu_6^0) + V(\mathbf{r}) \left(1 - g_{67} \frac{\partial n_6}{\partial \mu_6} \right)$$



Beyond weak coupling

M. Pierce *et al.* PRL **123**, 080403 (2019)

A. Bigué *et al.*, *in preparation*

Energy of the polaron: beyond mean-field contributions

Contact interaction $V(r) = g'_{IF} \delta(\mathbf{r})$

$$g'_{IF}(\Lambda) = g_{IF} + \frac{g_{IF}^2}{\Omega} \sum_{q < \Lambda} \frac{1}{2\varepsilon_q} + \dots$$

Second-order perturbation theory

$$\Delta E = g_{IF} n_F + \frac{g_{IF}^2 n_F}{\Omega} \sum_{q < \Lambda} \left[\frac{1}{2\varepsilon_q} - \chi(q, \varepsilon_q) \right]$$

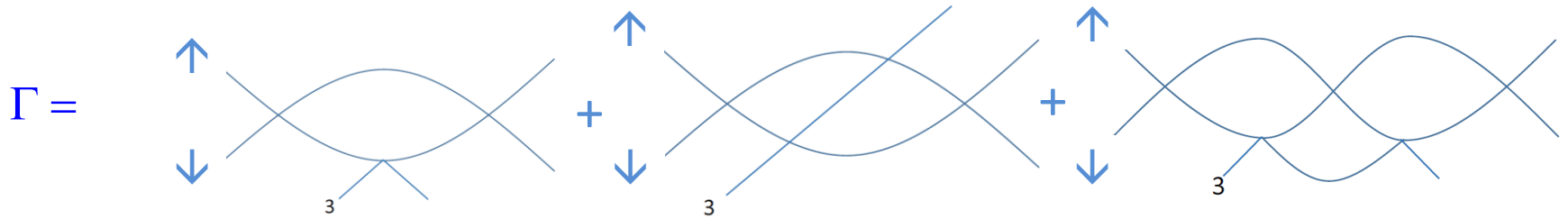
UV (log) divergent!

(see for instance BCS)

« Dynamical compressibility »
of the superfluid

Back to the three-body problem

Same divergence in the 3-body scattering (\approx Wu beyond mean-field correction for BECs)



$= g'_{BF}$ (Born) \Rightarrow log-divergent

$= \frac{g_{BF}}{1 + ika_{BF}}$ (Faddeev) \Rightarrow regular

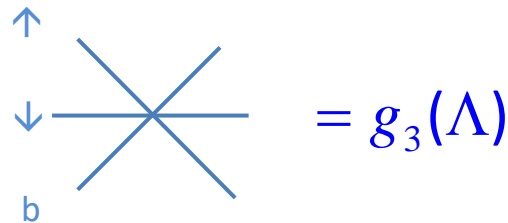
$$\Gamma_{\text{Born}} - \Gamma_{\text{Faddeev}} = g_{BF}^2 \kappa \left(\frac{m_B}{m_F} \right) \log(\Lambda R_{3b}) + o(1)$$

$$\kappa(7/6) = \frac{4}{9\sqrt{3}} + \frac{1}{3\pi} - \frac{\sqrt{3}}{2\pi^2}$$

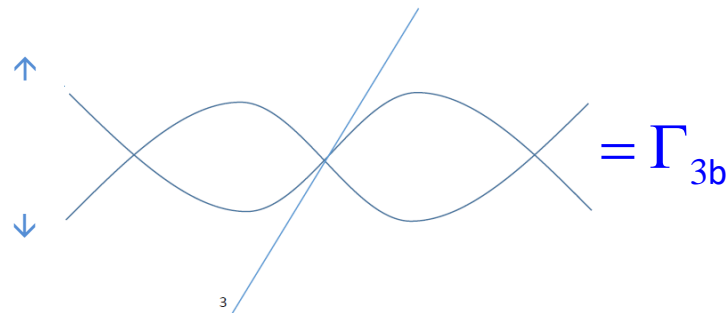
$$R_{3b} \simeq 1.5a_{BF} \text{ } ({}^7\text{Li}/{}^6\text{Li})$$

Curing the three-body problem

Add an explicit 3-body contact interaction to regularize the 3-body scattering amplitude (Braaten & Nieto '99)



Contribution to the 3-body amplitude



$$\Gamma_{\text{Faddeev}} = \Gamma_{\text{Born}} + \Gamma_{3b} \Rightarrow g_3(\Lambda) \left(\frac{1}{\Omega} \sum_{k < \Lambda} \frac{1}{2\varepsilon_k} \right)^2 = -g_{BF}^2 \kappa \left(\frac{m_B}{m_F} \right) \log(\Lambda R_{3b}) + o(1)$$

Curing the polaron energy

Treat the **three-body interaction perturbatively** for the polaron

$$E = g_{\text{IF}} n_{\text{F}} \left[1 + k_{\text{F}} a_{\text{IF}} F \left(\frac{1}{k_{\text{F}} a_{\text{FF}}} \right) + \frac{a_{\text{IF}} C_{\text{F}}}{N} \kappa \left(\frac{m_{\text{B}}}{m_{\text{F}}} \right) \ln(k_{\text{F}} R_{3\text{b}}) + \dots \right]$$

$$F \left(\frac{1}{k_{\text{F}} a_{\text{FF}}} \right) = \frac{4\pi}{k_{\text{F}}} \left[\frac{\hbar^2}{\mu} \int_{q < \Lambda} \frac{d^3 \mathbf{q}}{(2\pi)^3} \left(\frac{1}{2\varepsilon_q} - \chi(q, \varepsilon_q) \right) + \frac{C_{\text{f}}}{N} \kappa \left(\frac{m_{\text{B}}}{m_{\text{F}}} \right) \ln(\Lambda / k_{\text{F}}) \right]$$

Expression **does not make any assumption on the properties of the many-body background**

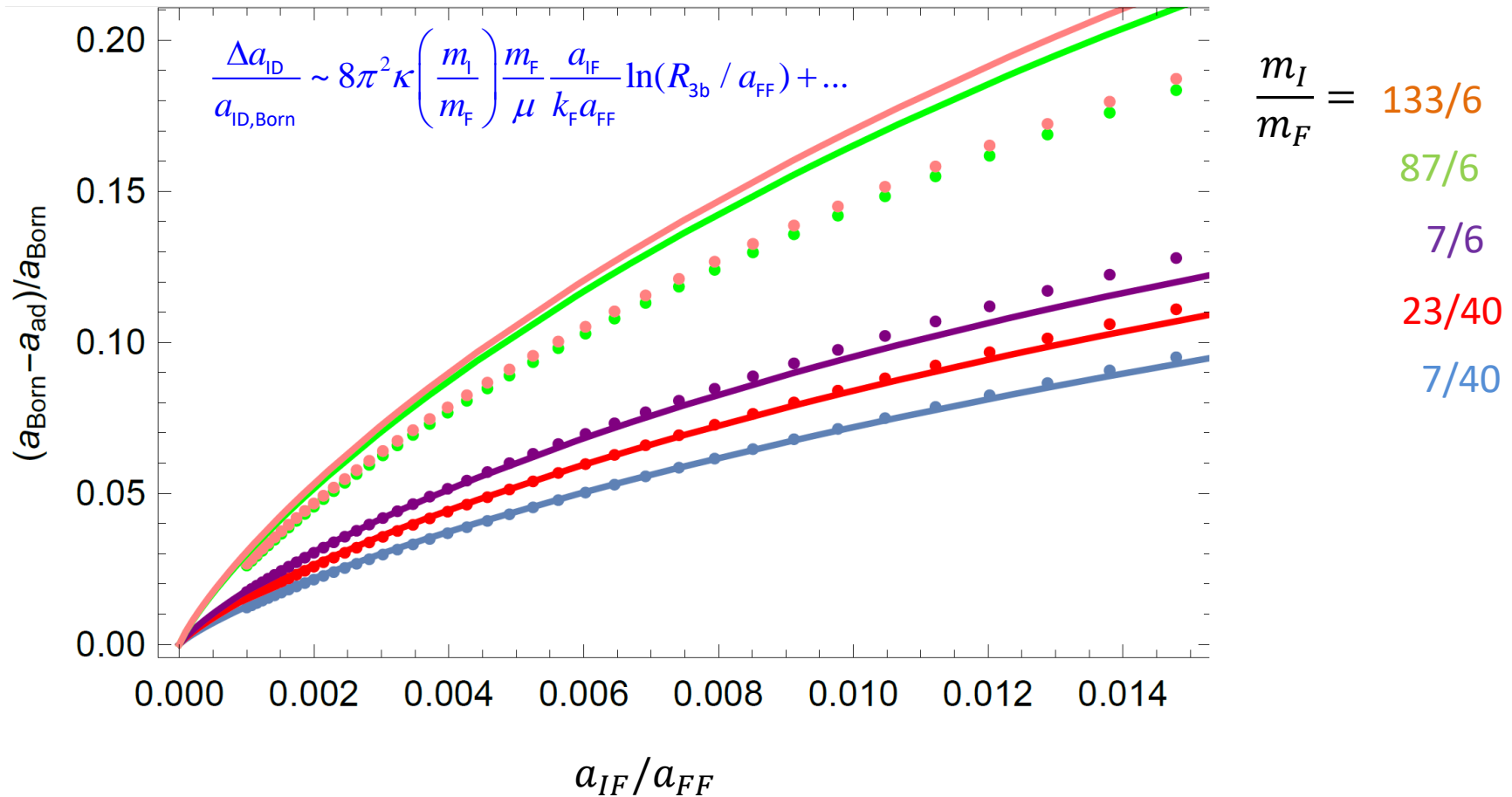
Asymptotic expressions

$$a_{\text{FF}} \rightarrow 0^- \Rightarrow \text{Fermi polaron problem} \Rightarrow F(-\infty) = \frac{3}{4\pi}$$

$$a_{\text{FF}} \rightarrow 0^+ \Rightarrow \text{Bose polaron problem} \Rightarrow E = g_{\text{ID}} n_{\text{D}}$$

$$\Rightarrow F(+\infty) \sim 8\pi^2 \kappa \left(\frac{m_{\text{B}}}{m_{\text{F}}} \right) \frac{m_{\text{F}}}{\mu} \frac{\ln(k_{\text{F}} a_{\text{FF}})}{k_{\text{F}} a_{\text{FF}}}$$

Cross-check: impurity-dimer scattering length



Calculation of F

$$F\left(\frac{1}{k_F a_{FF}}\right) = \frac{4\pi}{k_F} \left[\frac{\hbar^2}{\mu} \int_{q < \Lambda} \frac{d^3 \mathbf{q}}{(2\pi)^3} \left(\frac{1}{2\varepsilon_q} - \chi(q, \varepsilon_q) \right) + \frac{C_f}{N} \kappa \left(\frac{m_B}{m_F} \right) \ln(\Lambda / k_F) \right]$$

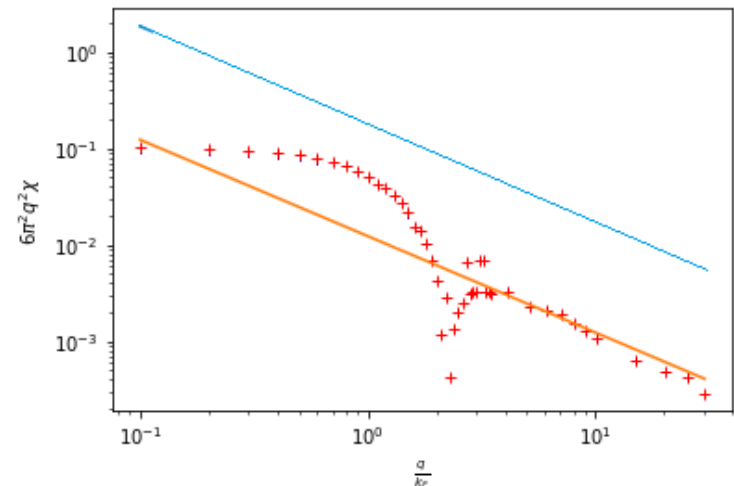
F depends only on the **properties of the medium**.

Use **BCS/Mean-field** theory to calculate χ .

- describes correctly the pair-breaking sector but misses collective modes (Bogoliubov Anderson modes) -> eg breakdown of the F-sum rule.

Regularization using 3-body interactions **fails within** BCS approximation

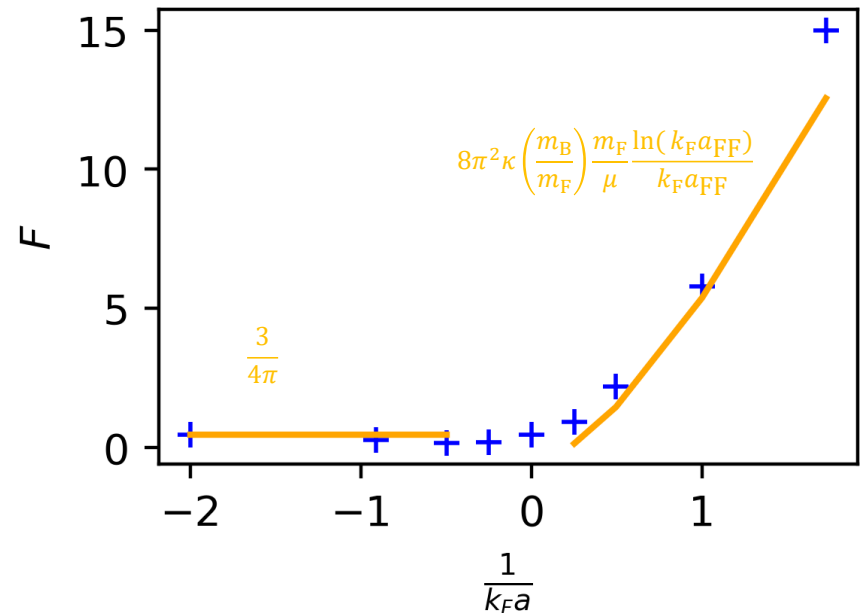
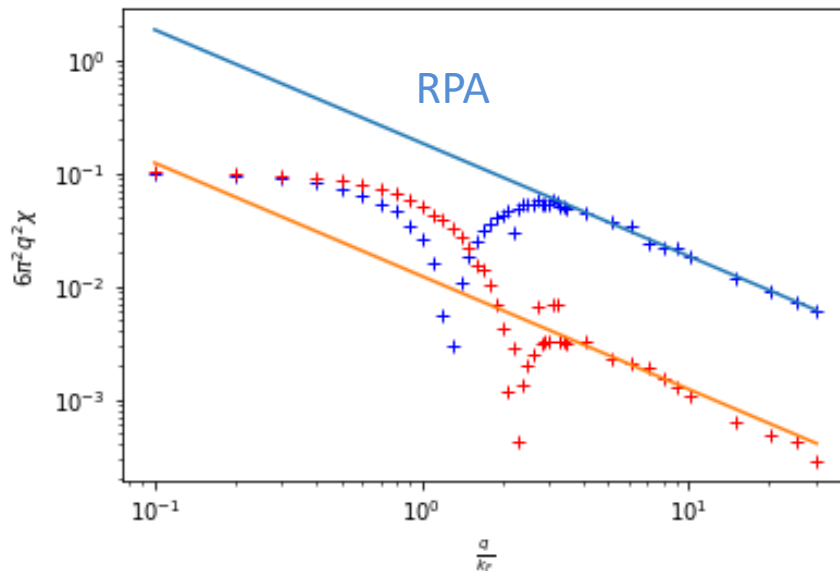
$$\chi(q, \varepsilon_q^{(r)}) = \frac{1}{\varepsilon_q^{(r)}} \left[1 - \pi^2 \kappa(\eta) \frac{m_f C_f}{m_r N q} + \dots \right]$$



Random Phase Approximation

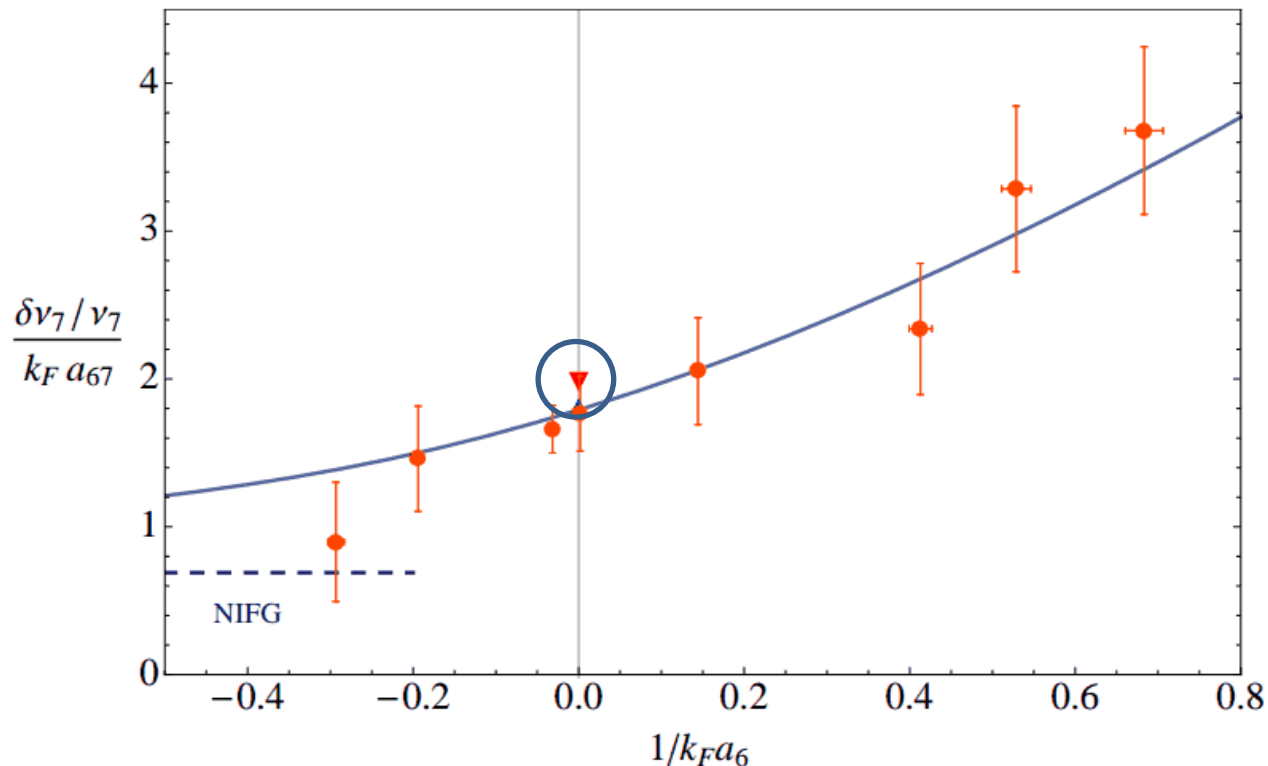
RPA: solve dynamic equations self-consistently by allowing the gap to vary (Minguzzi *et al. 01*, Combescot *et al. 06*,...).

- Static properties still given by BCS theory but recovers the collective modes
- Correct scaling for the compressibility at large q .



THE NEXT STEPS

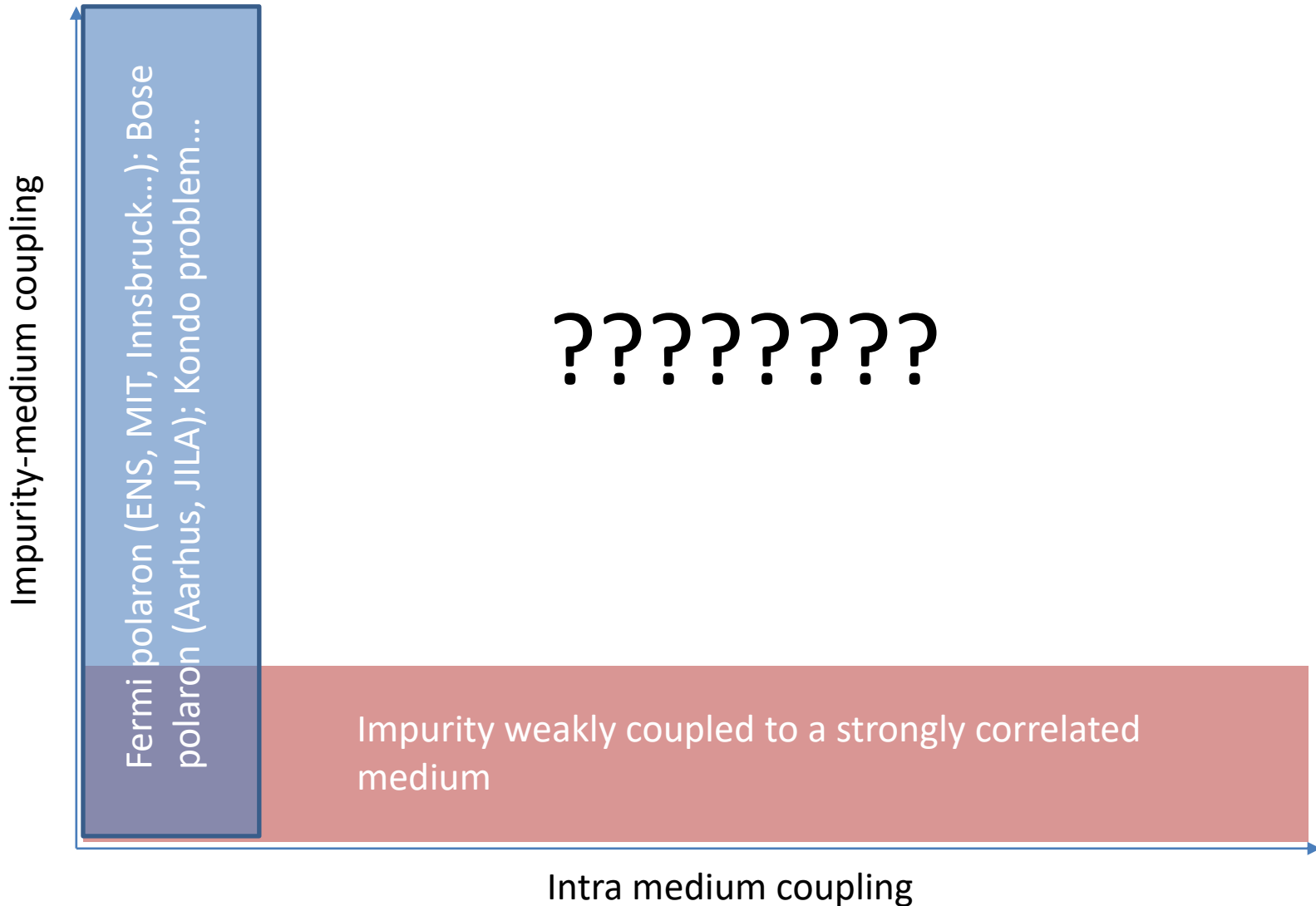
- **Challenge:** is it possible to observe the beyond mean-field corrections using high precision oscillation frequency measurements?



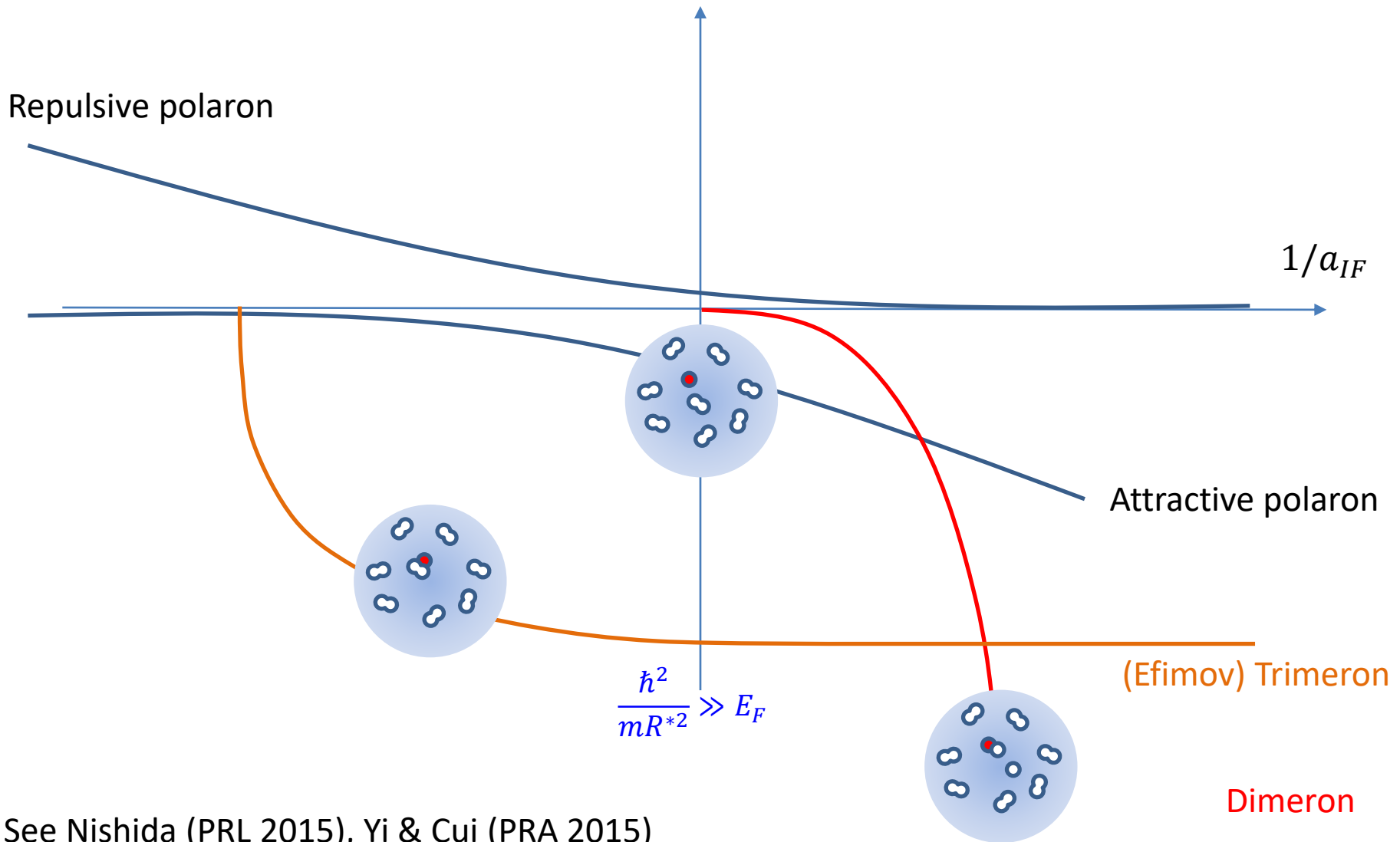
Conclusion and outlook

Beyond the weakly interacting regime, R. Alhyder *et al.* PRA **102**, 033322 (2020)

POLARON PHASE DIAGRAM



Polaron, dimeron, trimeron



See Nishida (PRL 2015), Yi & Cui (PRA 2015)

