



Laboratoire de Physique des 2 Infinis



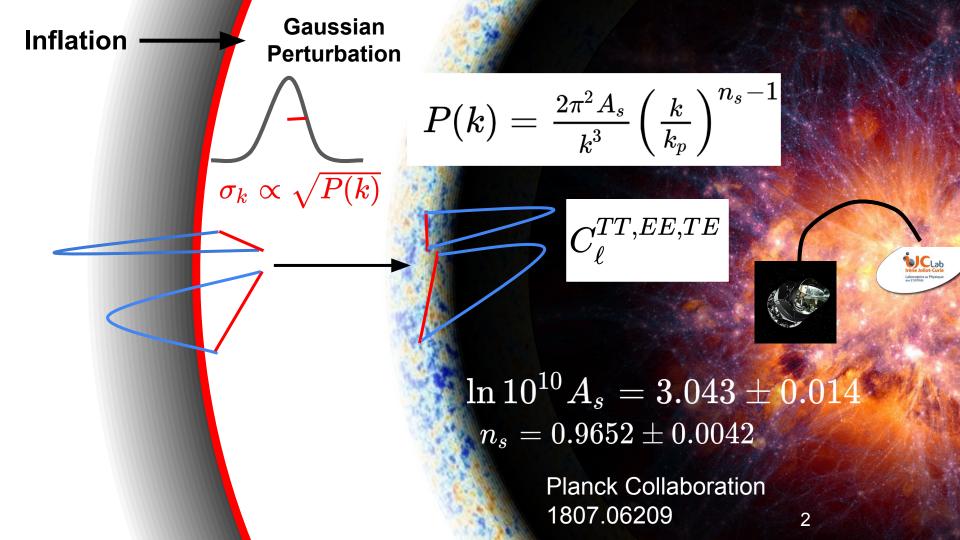


Second-order relativistic initial conditions for N-body simulations

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**Nearly-Gaussian** Inflation **Perturbation** Creminelli and Zaldarriaga 0407059  $B(k_1,k_2,k_3 
ightarrow 0) = (n_s-1)P(k_1)P(k_2)$  $B(k_1,k_2,k_3)$  $f_{NL} \sim 10^{-2}$ JCLab  $B^{IJK}_{\ell_1,\ell_2,\ell_3}$  $f_{NL}=-0.9\pm5.1$ **Planck Collaboration** 1807.06209 2

#### Inflation

### Future:

P(k)

 $B(k_1,k_2,k_3)$ 

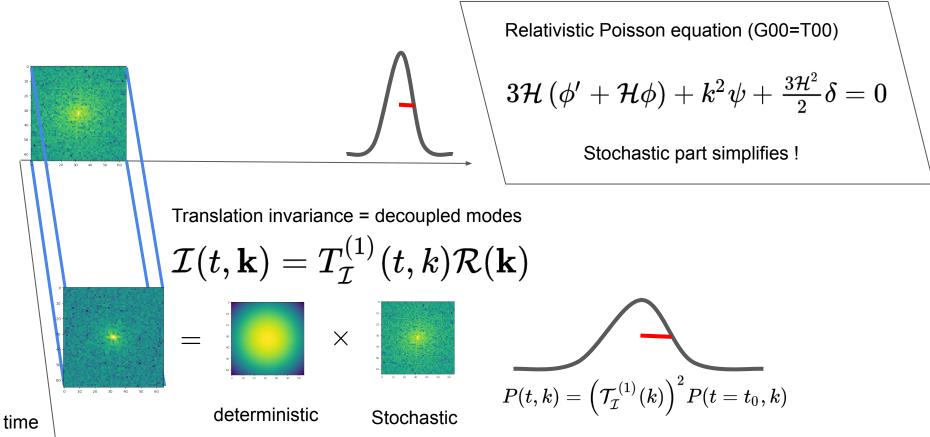
 $f_{NL} = -12 \pm 21$ SDSS Collaboration 2106.13725

Euclid  $\sigma_{f_{NL}} \sim 1$ Desjacques et al 1611.09787 SPHEREX  $\sigma_{f_{NL}} \sim 0.1$ SKA  $\sigma_{f_{NL}} \sim 0.1$ Karagiannis et al 1801.09280 and 1911.03964 Doré et al. 1412.4872

2

JCLab

## Linear evolution



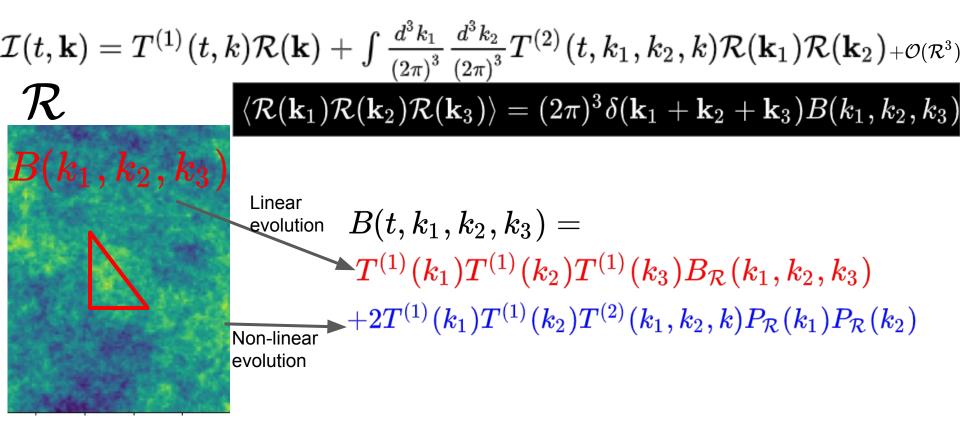
# **Quadratic Evolution**

First order equation with second order fields

Quadratic source term

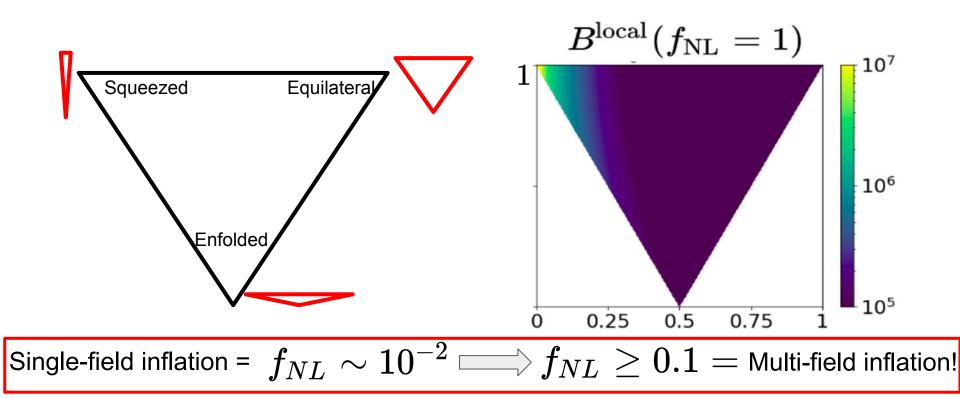
$$egin{aligned} & \mathbf{3}\mathcal{H}\left(\phi'+\mathcal{H}\phi
ight)+k^{2}\psi+rac{3\mathcal{H}^{2}}{2}\delta=3\mathcal{H}^{2}\phi^{2}-rac{1}{2}
abla_{i}\phi
abla^{i}\phi+2\phi\Delta\phi\ & 
abla_{i}\phi
abla^{i}\phi
ightarrow\intrac{d^{3}k_{1}}{(2\pi)^{3}}rac{d^{3}k_{2}}{(2\pi)^{3}}igg[-\mathbf{k}_{1}\cdot\mathbf{k}_{2}igg]\phi(\mathbf{k}_{1})\phi(\mathbf{k}_{2})\delta(\mathbf{k}-\mathbf{k}_{1}-\mathbf{k}_{2})\ & \phi\Delta\phi
ightarrow\intrac{d^{3}k_{1}}{(2\pi)^{3}}rac{d^{3}k_{2}}{(2\pi)^{3}}igg[-k_{2}^{2}igg]\phi(\mathbf{k}_{1})\phi(\mathbf{k}_{2})\delta(\mathbf{k}-\mathbf{k}_{1}-\mathbf{k}_{2}) \end{aligned}$$

Quadratic level: the bispectrum evolution



#### Squeezed limit

 $B(t,k_1,k_2,k_3) = T^{(1)}(k_1)T^{(1)}(k_2)T^{(1)}(k_3)B_{\mathcal{R}}(k_1,k_2,k_3)$ 

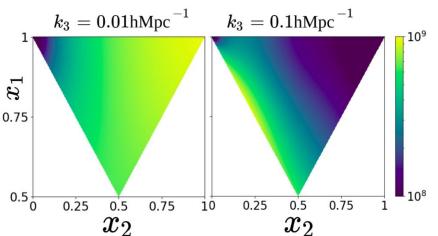


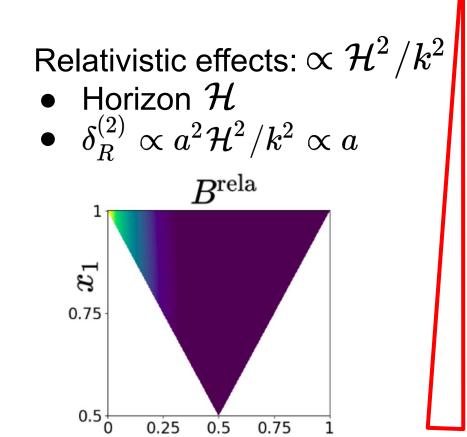
#### Relativistic effects

 $B(t,k_1,k_2,k_3) = +2T^{(1)}(k_1)T^{(1)}(k_2)T^{(2)}(k_1,k_2,k)P_{\mathcal{R}}(k_1)P_{\mathcal{R}}(k_2)$ 

Newtonian limit:

- No horizon
- ullet  $\delta^{(1)} \propto a$
- ullet  $\delta^{(2)} \propto a^2$

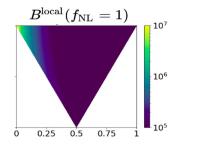




# $B(t, k_1, k_2, k_3) = T^{(1)}(k_1)T^{(1)}(k_2)T^{(1)}(k_3)B_{\mathcal{R}}(k_1, k_2, k_3) \\ +2T^{(1)}(k_1)T^{(1)}(k_2)T^{(2)}(k_1, k_2, k)P_{\mathcal{R}}(k_1)P_{\mathcal{R}}(k_2)$

• Degenerated in momentum space

$$T^{(1)}(t) \propto a$$

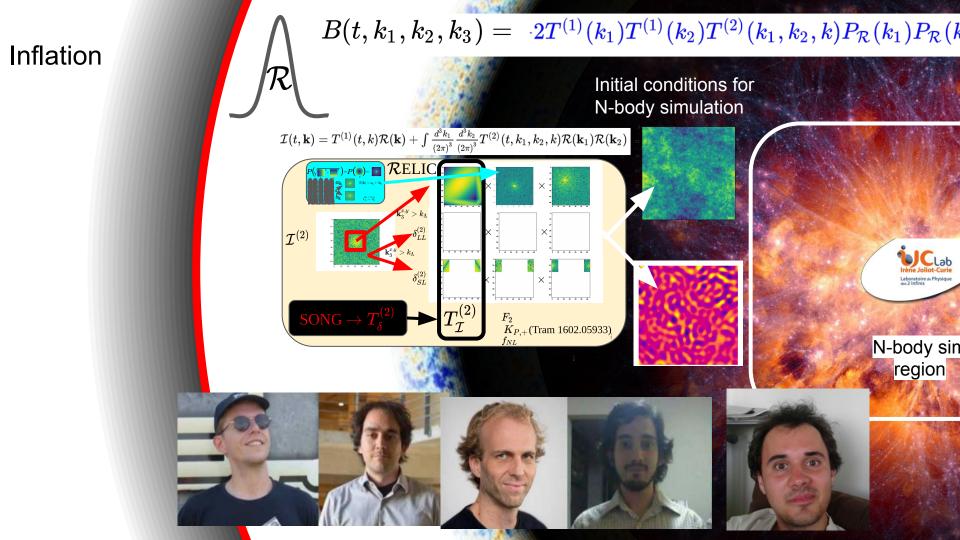


 $B^{
m rela}$ 

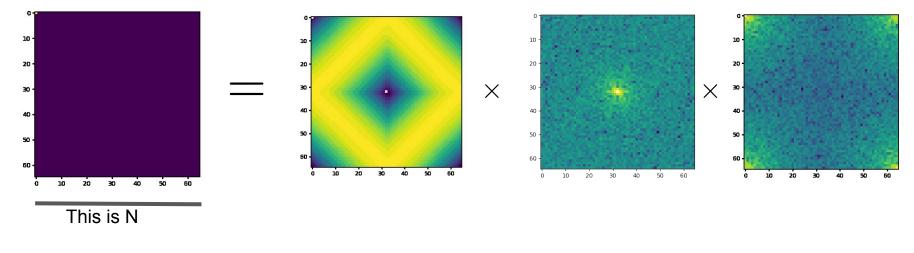
0.5

0.75

- Degenerated in time  ${T_R^{(2)}(t)\propto a\over T_R^{(2)}(t)\propto a}$
- Important for large scales, hence in the squeezed limit
- Small scales highly non-linear N-body-simulation

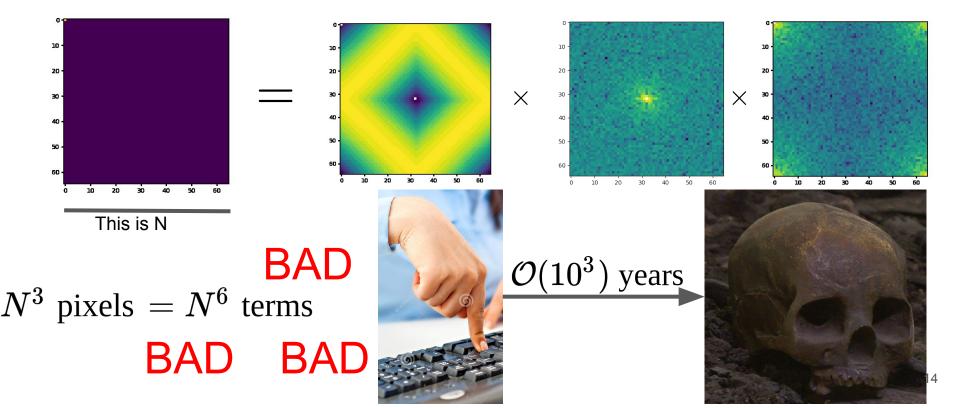


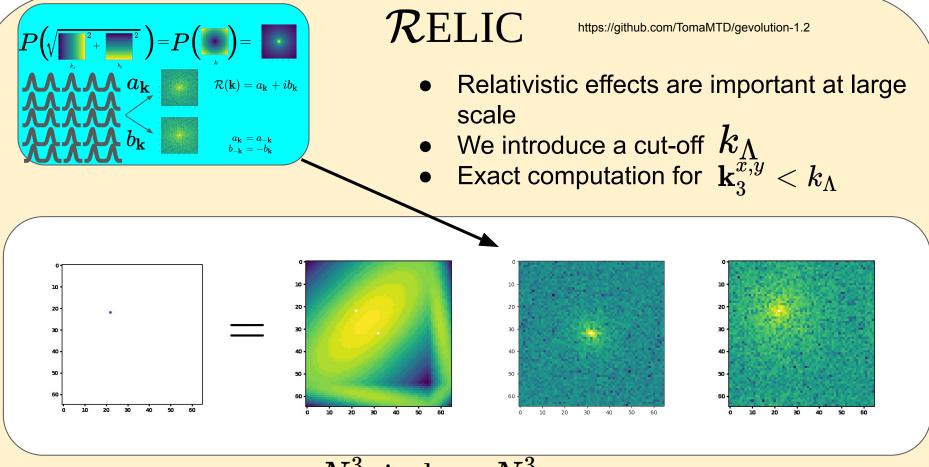
$$\mathcal{I}^{(2)}(\mathbf{k}_3) = \int rac{d^3k_1}{\left(2\pi
ight)^3} \, T_{\mathcal{I}}^{(2)}(k_1, |\mathbf{k}_3 - \mathbf{k}_1|, k_3) \mathcal{R}^{(1)}(\mathbf{k}_1) \mathcal{R}^{(1)}(\mathbf{k}_3 - \mathbf{k}_1)$$



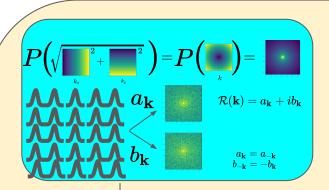
1 pixel =  $N^3$  terms  $\longrightarrow N^3$  pixels =  $N^6$  terms

$$\mathcal{I}^{(2)}(\mathbf{k}_3) = \int rac{d^3k_1}{\left(2\pi
ight)^3} \, T_{\mathcal{I}}^{(2)}(k_1, |\mathbf{k}_3 - \mathbf{k}_1|, k_3) \mathcal{R}^{(1)}(\mathbf{k}_1) \mathcal{R}^{(1)}(\mathbf{k}_3 - \mathbf{k}_1)$$





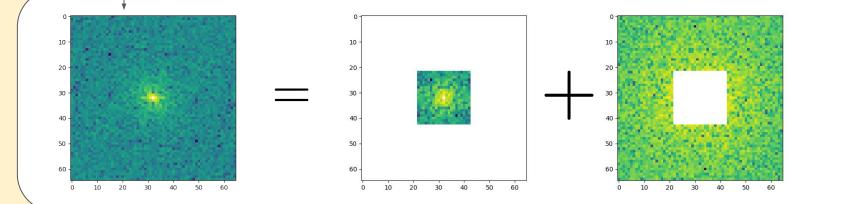
 $N_\Lambda^3$ pixels  $imes N^3$ terms

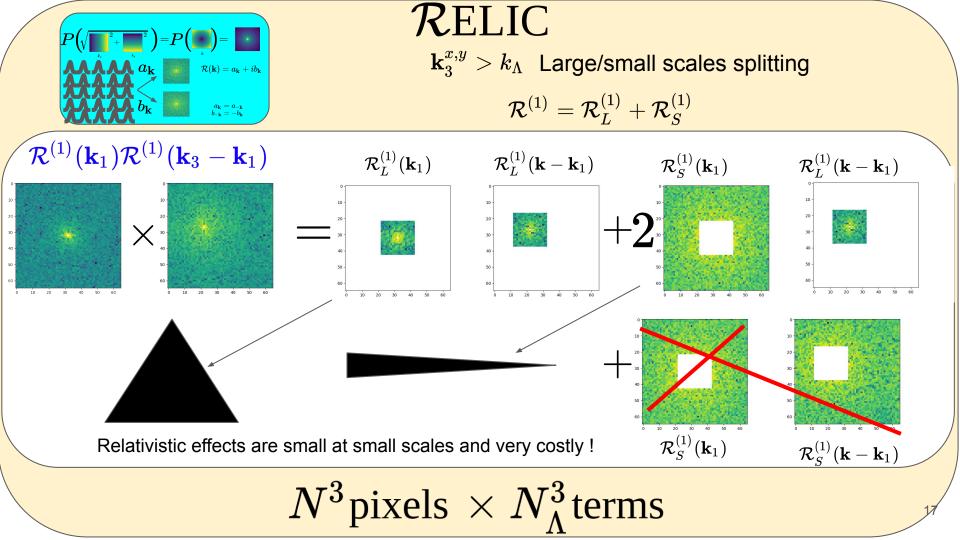


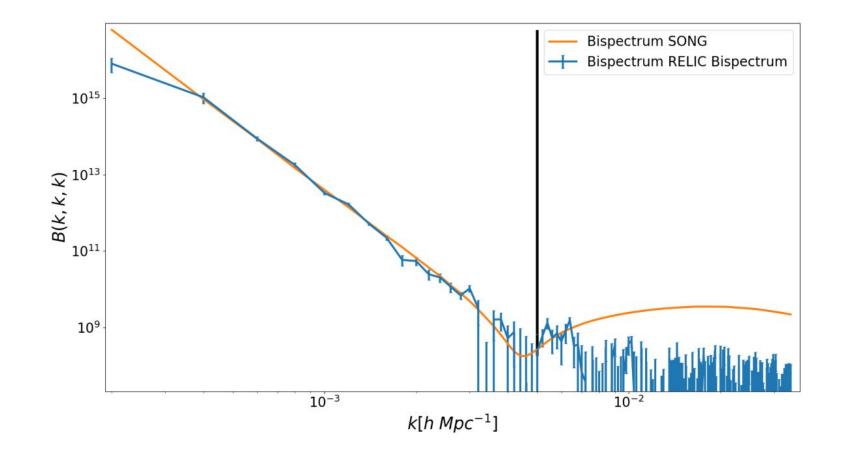
# $\mathcal{R}$ ELIC

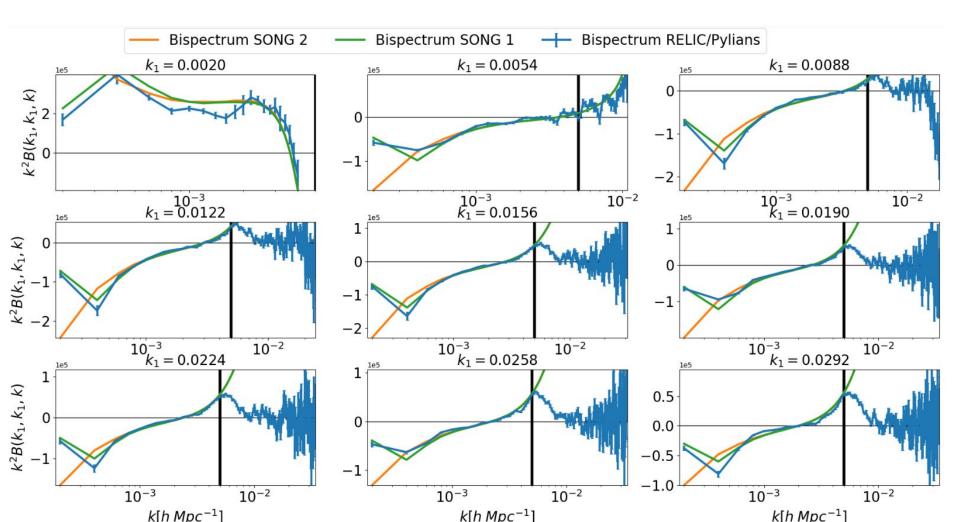
 $\mathbf{k}_3^{x,y} > k_\Lambda$  Large/small scales splitting

$$\mathcal{R}^{(1)}(k) = W(k) \mathcal{R}^{(1)}(k) + (1-W(k)) R^{(1)}(k)$$









## Merci !