

Coalescence approach in phase space for HQ

Statistical factor colour-spin-isospin

Parton Distribution function

Hadron Wigner function

$$\frac{dN_{Hadron}}{d^2 p_T} = g_H \int \prod_{i=1}^n p_i \cdot d\sigma_i \frac{d^3 p_i}{(2\pi)^3} f_q(x_i, p_i) f_W(x_1, \dots, x_n; p_1, \dots, p_n) \delta(p_T - \sum_i p_{iT})$$

Thermal+flow for u,d,s ($p_T < 3$ GeV)

$$\frac{dN_{q,\bar{q}}}{d^2 p_T} \sim \exp\left(-\frac{\gamma_T(m_T - p_T \cdot \beta_T + \mu_q)}{T}\right)$$

$$V = \pi R^2 \tau \cosh(y_z) , R(\tau_f) = R_0(1 + 0.5 \beta_{max} \tau_f)$$

$$\beta(r) = \frac{r}{R} \beta_{max}$$

Total V now set to SHM, A. Andronic JHEP(2021)035

+ quenched minijets for u,d,s ($p_T > 3$ GeV)

For Charm from the studies of R_{AA} and v_2 of D-meson to determine the Space Diffusion coeff.: from parton simulations solving relativistic Boltzmann transport equation
In pp it is FONNL distribution

Coalescence evaluated in a fireball

Space-momentum-time correlation over the freeze-out hypersurface of a transport simulation are **not fully transferred**

→ dimension set by experimental constraints (if any)

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Wigner function <-> Wave function

$$\Phi_M^W(\mathbf{r}, \mathbf{q}) = \int d^3 r' e^{-i\mathbf{q}\cdot\mathbf{r}'} \varphi_M\left(\mathbf{r} + \frac{\mathbf{r}'}{2}\right) \varphi_M^*\left(\mathbf{r} - \frac{\mathbf{r}'}{2}\right)$$

$\varphi_M(\mathbf{r})$ meson wave function

Assuming gaussian wave function

$$f_M(x_1, x_2; p_1, p_2) = A_W \exp\left(-\frac{x_{r1}^2}{\sigma_r^2} - p_{r1}^2 \sigma_r^2\right)$$

For baryon $N_q=3$

$$f_H(\dots) = \prod_{i=1}^{N_q-1} A_W \exp\left(-\frac{x_{ri}^2}{\sigma_{ri}^2} - p_{ri}^2 \sigma_{ri}^2\right)$$

Note: only σ_r coming from $\varphi_M(\mathbf{r})$ or $\sigma_r^* \sigma_p = 1$

valid for harmonic oscillator with $V(r) \rightarrow \sigma_r^* \sigma_p > 1$

Wigner function width fixed by root-mean-square charge radius from quark model

| Meson | $\langle r^2 \rangle_{ch}$ | σ_{p1} | σ_{p2} |
|-----------------------|----------------------------|---------------|---------------|
| $D^+ = [c\bar{d}]$ | 0.184 | 0.282 | — |
| $D_s^+ = [\bar{s}c]$ | 0.083 | 0.404 | — |
| Baryon | $\langle r^2 \rangle_{ch}$ | σ_{p1} | σ_{p2} |
| $\Lambda_c^+ = [udc]$ | 0.15 | 0.251 | 0.424 |
| $\Xi_c^+ = [usc]$ | 0.2 | 0.242 | 0.406 |
| $\Omega_c^0 = [ssc]$ | -0.12 | 0.337 | 0.53 |

C.-W. Hwang, EPJ C23, 585 (2002);
C. Albertus et al., NPA 740, 333 (2004)

$$\begin{aligned} \langle r^2 \rangle_{ch} &= \frac{3}{2} \frac{m_2^2 Q_1 + m_1^2 Q_2}{(m_1 + m_2)^2} \sigma_{r1}^2 \\ &\quad + \frac{3}{2} \frac{m_3^2 (Q_1 + Q_2) + (m_1 + m_2)^2 Q_3}{(m_1 + m_2 + m_3)^2} \sigma_{r2}^2 \end{aligned} \quad (8)$$

$\sigma_{ri} = 1/\sqrt{\mu_i \omega}$ Harmonic oscillator relation

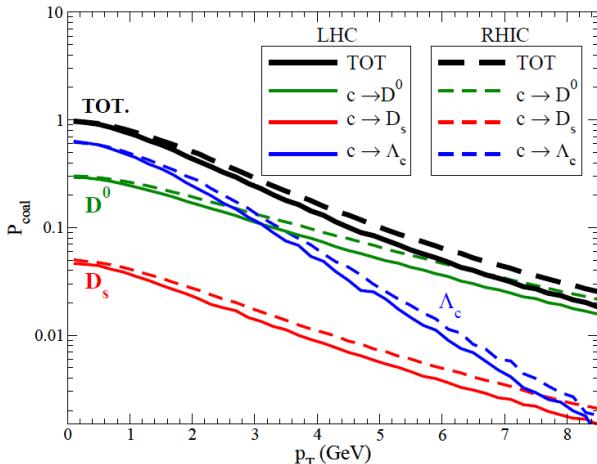
$$\mu_1 = \frac{m_1 m_2}{m_1 + m_2}, \quad \mu_2 = \frac{(m_1 + m_2) m_3}{m_1 + m_2 + m_3}.$$

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$$f_H(\dots) = \prod_{i=1}^{N_q-1} A_W \exp\left(-\frac{x_{ri}^2}{\sigma_{ri}^2} - p_{ri}^2 \sigma_{ri}^2\right)$$



- Normalization in $f_W(\dots)$ fixed by requiring $P_{coal}(p \rightarrow 0) = 1$ Not in standard coalescenceothers modify by hand σ_r to enforce confinement for a charm at rest in the medium.

- The charm not “coalescing” undergo fragmentation:

$$\frac{dN_{had}}{d^2 p_T dy} = \sum \int dz \frac{dN_{fragm}}{d^2 p_T dy} \frac{D_{had/c}(z, Q^2)}{z^2}$$

charm number conserved at each p_T