DIFFERENTIAL HIGGS BOSON PRODUCTION

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MOTIVATION

- Theoretical predictions need to keep up with the ever-increasing precision of experimental measurements
- ▶ Need to understand the SM background in order to resolve **new physics**



Example: Higgs inclusive: $8\% \to 3\%$ expected experimental uncertainty at 3000 fb^{-1} . The PDF uncertainty on the theoretical prediction cannot be neglected anymore.



HIGGS AT N3LO

▶ 2015 : First Prediction for the inclusive Higgs done with threshold expansion at N³LO [C. Anastasiou, C. Duhr, F. Dulat, F. Herzog, B. Mistlberger]

▶ 2016: Vector-Boson Fusion Higgs production in QCD at three loops

[F. A. Dreyer, A. Karlberg]

2018

Exact Higgs inclusive

[B. Mistlberger]

- Study of threshold expansion for Differential Higgs
- [F. Dulat, B. Mistlberger, AP]

Vector-Boson Fusion Higgs Pair Production

[F. A. Dreyer, A. Karlberg]

2019

- ullet Higgs differential using q_T subtraction [L. Cieri, X. Chen, T. Gehrmann, E. Glover, A. Huss]
- Higgs rapidity distribution using threshold expansion [F. Dulat, B. Mistlberger, AP]

▶ 2020

- Higgs Boson Production in Bottom-Quark Fusion [c. Duhr, F. Dulat, B. Mistlberger]
- ...matching the 4- and 5-flavour schemes [c. Duhr, F. Dulat, V. Hirschi, B. Mistlberger]

2021

Fully Differential Higgs Boson Production

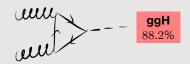
[X. Chen, T. Gehrmann, E. Glover, A. Huss, B. Mistlberger, AP]

Higgs pT Spectrum [G. Billis, B. Dehnadi, M. A. Ebert, J. K. L. Michel, F. J. Tackmann]



GLUON FUSION

ggH is a **loop-induced** process and represents the largest correction to the cross-section at **13 TeV**



The computation is performed in the **infinite top mass** approximation **Effective theory:**



- ► Remove one loop!
- ▶ Good approximation: $\delta_t^{NNLO} \sim 0.7\%$
- To be combined with mass corrections, EW corrections, etc...



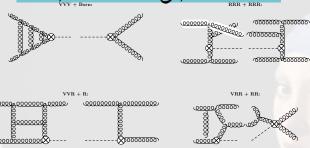
HIGGS RAPIDITY DISTRIBUTION



The **real radiation** is treated **inclusively** and the Higgs boson is resolved only in its rapidity

$$\frac{\mathrm{d}\hat{\sigma}_{ij\to H+X}}{\mathrm{d}Y} = \int \mathrm{d}p_h \frac{\mathrm{d}\phi_n}{\mathrm{d}\phi_n} \left| \mathcal{M}_{ij\to H+X} \right|^2 \delta\left(Y - \frac{1}{2}\log\left(\frac{E+p_z}{E-p_z}\right)\right)$$

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$$\frac{\mathrm{d}\hat{\sigma}_{ij\to H+X}}{\mathrm{d}Y} = \int \mathrm{d}p_h \mathrm{d}\phi_n \left| \mathcal{M}_{ij\to H++X} \right|^2 \delta\left(Y - \frac{1}{2}\log\left(\frac{E+p_z}{E-p_z}\right)\right)$$

RAPIDITY DISTRIBUTION

The general form of the rapidity distribution can be written as:

$$\frac{\mathsf{d}\sigma_{PP\to H+X}}{\mathsf{d}Y} = \hat{\sigma}_0 \sum_{ij} \int_0^1 \mathsf{d}x_1 \mathsf{d}x_2 \mathsf{d}y_1 \mathsf{d}y_2 f_i(y_1) f_j(y_2) \delta(\tau - x_1 x_2 y_1 y_2)$$

$$\times \delta \left(\mathsf{Y} - \frac{1}{2} \log \left(\frac{\mathsf{x}_1 \mathsf{y}_1}{\mathsf{x}_2 \mathsf{y}_2} \right) \right)$$
 Partonic $\eta_{ij}(\mathsf{x}_1, \mathsf{x}_2)$,

Where we define the partonic cross-section in terms of a power series in the **strong coupling** constant:

$$\eta_{ij}(\mathbf{x}_1,\mathbf{x}_2) = \sum_{k=0}^{3} \left(\frac{\alpha_{S}}{\pi}\right)^k \eta_{ij}^{(k)}(\mathbf{x}_1,\mathbf{x}_2).$$



ROAD TO COMPUTATION

Computing the rapidity distribution analytically is a hard challenge!

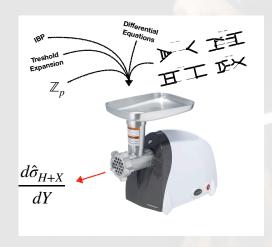
$$\frac{\mathrm{d}\hat{\sigma}_{ij\to H+X}}{\mathrm{d}\mathbf{Y}}$$

Divide and Conquer:

 Perform expansion around the production threshold.
 Already a success for the inclusive N3LO

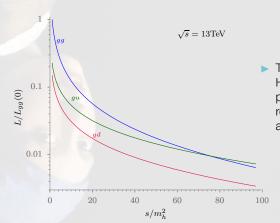
$$\bar{z} = 1 - z = 1 - \frac{m_H^2}{s} \sim 0$$

Validate the truncation of the threshold expansion



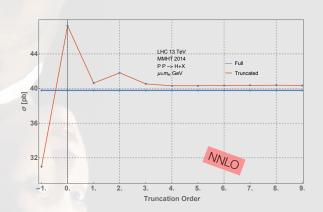
PARTON LUMINOSITY

$$L(z) = \int_{\frac{\tau}{z}}^{1} \frac{dx}{x} f_{i}(x) f_{j}\left(\frac{\tau}{zx}\right).$$



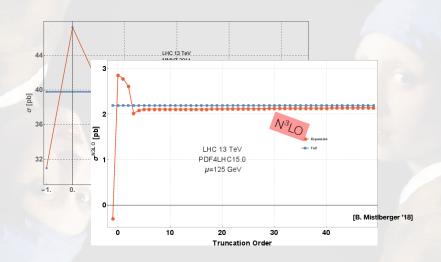
The probability of producing the Higgs boson as a function of the partonic center of mass is reduced as the energy moves away from the threshold

INCLUSIVE THRESHOLD EXPANSION





INCLUSIVE THRESHOLD EXPANSION





THRESHOLD EXPANSION

In dimensional regularization the expression for the partonic cross-section takes the form,

$$\begin{split} \eta_{ij}^{(3)}(\mathbf{X}_1,\mathbf{X}_2) &= \eta_{ij}^{(3)}\delta(1-\mathbf{X}_1)\delta(1-\mathbf{X}_2) \\ &+ \sum_{n,m=1}^{3} \underbrace{ \begin{array}{c} \mathbf{Distributions} \\ (1-\mathbf{X}_1)^{-1-m\epsilon} \ (1-\mathbf{X}_2)^{-1-n\epsilon} \end{array} }_{\eta_{ij}^{(3,m,n)}(\mathbf{X}_1,\mathbf{X}_2)}, \end{split}} \end{split}}$$

Different sectors of the loop momentum give rise to different m, n exponent

$$\int_0^1 dx (1-x)^{-1+a\epsilon} f(x) = \int_0^1 dx \left[\frac{\delta(1-x)}{a\epsilon} + \sum_{n=0}^\infty \frac{(a\epsilon)^n}{n!} \left[\frac{\log^n(1-x)}{1-x} \right]_+ \right] f(x)$$

REACHING BEYOND THRESHOLD EXPANSION

Obtain finite expressions with a suitable mass factorization and ultraviolet renormalization counter term $CT_n^{(3)}$:

$$\eta_{\mathit{ij}}^{(3)}(\textit{x}_{1}, \textit{x}_{2}) = \lim_{\epsilon \to 0} \left[\eta_{\mathit{ij}, \mathsf{bare}}^{(3)}(\textit{x}_{1}, \textit{x}_{2}) + \textit{CT}_{\mathit{ij}}^{(3)}(\textit{x}_{1}, \textit{x}_{2}) \right]$$

- Use the fact that poles in the dimensional regulator ϵ cancel to impose further constraints on the partonic functions.
- Fix most of the logarithmically enhanced terms
- Smaller set of expressions that need threshold expansion

Integrate over the rapidity to match the inclusive x-section,

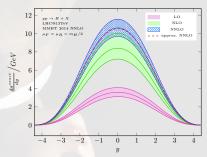
$$\eta_{ij}^{(3),incl.}(z) = \int dY \eta_{ij}^{(3)}(x_1,x_2).$$

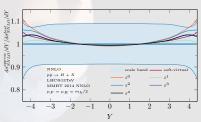
- Strong check on the differential partonic cross-section
- Agreement between the two threshold expansions for all computed orders!



RESULTS

Threshold at NNLO:



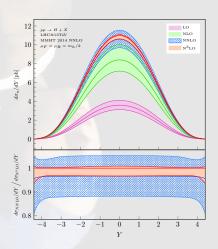


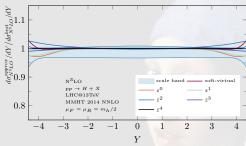
- Threshold expansion using Differential equations and expansion by regions
- Obtain 6 terms in the expansion
- Rescale to the inclusive result
- The approximation performs well for central rapidities |Y| < 3</p>
- Consistent improvement by including more terms
- To access the missing information from high energy contributions and fill the gap to the exact NNLO we need other tools.



RESULTS

Threshold at N3LO:





- Consistent behaviour between NNLO and N3LO regarding threshold expansion!
- uniform throughout the entire rapidity range
- Scale variation uncertainty reduced to [-3.4%,+0.9%]
- **Agreement** with a independent computation relying on q_T -subtraction

[Cieri,Chen,Gehrmann,Glover,Huss]



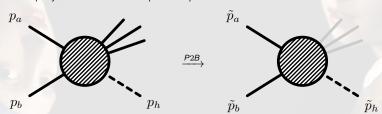
PROJECTION TO BORN

X. Chen, T.Gehrmann, E.W.N. Glover, A. Huss, B. Mistlberger, AP [2102.07607]

The master formula for the **Projection to Born** (P2B):

$$\frac{\mathrm{d}\sigma_{H}^{N3LO}}{\mathrm{d}\mathcal{O}} = \left(\frac{\mathrm{d}\sigma_{H+jet}^{N2LO}}{\mathrm{d}\mathcal{O}} - \frac{\mathrm{d}\sigma_{H+jet}^{N2LO}}{\mathrm{d}\tilde{\mathcal{O}}}\right) + \frac{\mathrm{d}\sigma_{H}^{N3LO}}{\mathrm{d}\tilde{\mathcal{O}}}$$

- Phase space singularities associated with fully unresolved configurations are cancelled identically.
- The projection to the born phase space $\mathcal{O} \xrightarrow{P2B} \tilde{\mathcal{O}}$ is defied as:

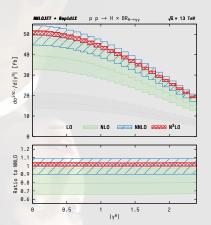


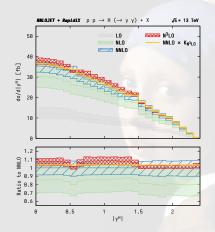
$$\tilde{\rho}_a = \xi_a \rho_a, \quad \tilde{\rho}_b = \xi_b \rho_b, \quad \tilde{\rho}_h = \tilde{\rho}_a + \tilde{\rho}_b$$

$$\tilde{\rho}_a^2 = \tilde{\rho}_b^2 = 0, \quad \tilde{\rho}_b^2 = \rho_b^2, \quad \tilde{\gamma}_b^2 = \gamma_b^2,$$



RESULTS: DI-PHOTON DECAY





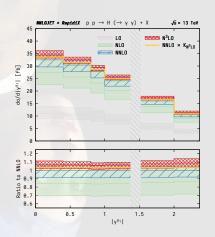
Fiducial cuts:

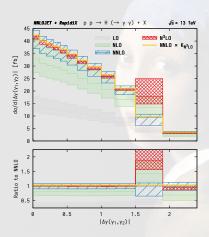
$$\rho_{\tau}^{\gamma_1} > 0.35 \times m_{\gamma\gamma}, \quad \rho_{\tau}^{\gamma_2} > 0.25 \times m_{\gamma\gamma},$$

 $|\eta^{\gamma}| < 2.37$ excluding $1.37 < |\eta^{\gamma}| < 1.52$



RESULTS: DI-PHOTON DECAY







SUMMARY

- We computed the Higgs boson rapidity distribution at N3LO
- ▶ We observe **stabilisation** of the perturbative expansion in the strong coupling and a significant reduction in the scale variation of the cross-section to [-3.4%,+0.9%]
- N3LO corrections are uniform throughout the entire rapidity range, well estimated by a K-factor
- Combined the rapidity distribution at N3LO with the Higgs + jet at NNLO distribution using projection to born to compute the fully differential Higgs.
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